Signals of the QCD Phase Transition in Compact Stars

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Outline

- Introduction: Nuclear Equation of State and \bullet **Supernovae**
- Constraints on Quark Matter from Pulsar Masses
- QCD Phase Transition and Supernovae \bullet
- **Summary** \bullet

Introduction: Nuclear Equation of State and Supernovae

Nuclear Equation of State as Input in Astrophysics

- supernovae simulations: $T=1\text{--}50$ MeV, $n=10^{-10}$ $^{\rm o}\!\!-\!\!2n_0$ D
- proto-neutron star: $T=1\text{--}50$ MeV, $n=10^{-3}$ £ $^{\rm 5-10}n_0$
- global properties of neutron stars: $T=0,\,n=10^{-3}$ $^{\rm 5-10}n_0$ D

neutron star mergers: $T=0\text{--}100$ MeV, $n=10^{-10}$ ∙ $^{\rm o}\!\!-\!\!10n_0$

New Equations of State for Supernovae

(Hempel, Fischer, JSB, Liebendörfer 2012, Steiner, Hempel, Fischer 2012)

- new equation of state for supernova simulations availableD
- check with new pulsar mass limit £
- based on models for describing nuclear properties (nucleons only)
- improvement in NSE (Hempel and JSB 2010)

Neutrino Luminosities for a 40 M_{\odot} Progenitor

(Hempel, Fischer, JSB, Liebendörfer 2012)

- neutrino luminosities for ^a40 M_{\odot} progenitor star
	- collapse to ^a black hole within1s
- collapse time depends onequation of state
- differences in particular in ν_{μ} \bullet and ν_τ spectra

Correlation between maximum mass and collapse time

(Hempel, Fischer, JSB, Liebendörfer 2012)

- relation between collapse timescale to ^a black hole and compact star masses£
- cases for cold case (crosses), for constant entropy per baryon $S/A = 4$ £ (circles), and from simulation (boxes) for different EoS

constant $S/A = 4$ describes maximum mass achieved in simulations quite well

Profiles just before collapse(Hempel, Fischer, JSB, Liebendörfer 2012)

- radial profiles of density, abundances, entropy, temperature
- **Situation just before collapse** to ^a black hole
- **P** rather constant entropy per baryon of $S/A = 4$
- extremely high densities and∙ temperatures (about 100MeV!)

Phase Diagram of Quantum Chromodynamics QCD

Early universe at zero density and high temperature£

- neutron star matter at small temperature and high density£
- first order phase transition at high density (not deconfinement!)£
- probed by heavy-ion collisions at GSI, Darmstadt (FAIR)

Structure of a Neutron Star (Fridolin Weber)

A Quark Star? (NASA press release 2002)

NASA news release 02-082: "Cosmic X-rays reveal evidence for new form of matter"^a quark star?

Neutron Star versus Strange Quark Star

(Chandra X-Ray Center, 2002)

Selfbound Star versus Ordinary Neutron Star

(Hartle, Sawyer, Scalapino (1975!))

selfbound stars:

- vanishing pressure at ^a finiteenergy density
- mass-radius relation starts at the£ origin (ignoring ^a possible crust)
- arbitrarily small masses and radii D possible

neutron stars:

- bound by gravity, finite pressure forD all energy density
- mass-radius relation starts at large \bullet radii
- minimum neutron star mass: **D** $M\sim 0.1M_\odot$ with $R\sim 200$ km

Hybrid Stars in the effective mass bag model

(Schertler et al. (2000))

- hybrid star: consists of hadronic and quark matter£
- three phases possible: hadronic, mixed phase and pure quark phase£
- composition depends crucially on the parameters as the bag constant B £ (and on the mass)

Matching to low density EoS

Two possibilities for ^a first-order chiral phase transition:

- A weakly first-order chiral transition (or no true phase transition), \bullet ⇒ one type of compact star:
bristal atare masaguerade as n hybrid stars masquerade as neutron stars
- A strongly first-order chiral transition⇒ two types of compact stars:
e naw atable aslution with emall ^a new stable solution with smaller masses and radii

Third Family of Compact Stars (Gerlach 1968)

(Glendenning, Kettner 2000; Schertler, Greiner, JSB, Thoma 2000)

third solution to the TOV equations besides white dwarfs and neutron stars, ∙ solution is stable!

- generates stars more compact than neutron stars∙
- possible for any first order phase transition!

Signals for Quark Matter/Phase Transition?

- delayed collapse of ^a proto-neutron star to ^a black hole \bullet (Thorsson, Prakash, Lattimer, 1994)
- spontaneous spin-up of pulsars (Glendenning, Pei, Weber, 1997) \bullet
- mass-radius relation: rising twins (Schertler et al., 2000)∙
- rapidly rotating pulsars due to r-mode *stability* window ∙
- enhanced cooling of neutron stars \bullet
- collapse of ^a neutron star to the third family? (gravitational \bullet waves, γ -rays, neutrinos)
- gravitational wave signals of phase transitions from neutron star \bullet mergers?
- **Secondary shock wave in supernova explosions (Sagert, Fischer** et al. 2009)

Signals for Strange Stars?

. . .

similar masses and radii, cooling, surface (crust), . . . but look for

- extremely small mass, small radius stars (includes strangelets)
- strange dwarfs: small and light white dwarfs with ^a strange star∙ core (Glendenning, Kettner, Weber, 1995)
- super-Eddington luminosity from bare, hot strange stars (Pageand Usov, 2002)
- quark novae with modified r-process nucleosynthesis (Jaikumar, \bullet Meyer, Otsuki, Ouyed 2007)
- gamma-ray bursts by conversion to strange quark matter (GRBs \bullet without ^a supernova, late x-ray emission, long quiescent times)

Constraints on quark matter from pulsar masses

Constraints on the Mass–Radius Relation(Lattimer and Prakash 2004)

spin rate from PSR B1937+21 of 641 Hz: $R < 15.5$ km for $M = 1.4 M_{\odot}$

- Schwarzschild limit (GR): $R>2GM=R_s$ L
- causality limit for EoS: $R > 3 G M$ £
- mass limit from PSR J1614-2230 (red band): $M=(1.97\pm0.04)M_\odot$

Quark Star Masses: Unpaired Case

Use free gas of quarks with ^a term from interactions and from ^avacuum energy:

$$
\Omega_{QM} = \sum_{i=u,d,s,e} \Omega_i + \frac{3\mu^4}{4\pi^2} (1 - a_4) + B_{eff}
$$

- Effective model with an expansion in the chemical potential μ \bullet
- Two parameters: effective bag constant B_{eff} and interaction parameter a_4
- 2-flavour constraint: nuclei do not collapse to (u,d) quark matter!
- 3-flavour constraint: strange (u,d,s) quark matter shall be more \bullet stable than nuclear matter, so that selfbound quark stars dubbedstrange stars can exist

Quark Star Masses: Unpaired Case

- Kepler line: mass shedding limit for 716 Hz \bullet (highest observed pulsar frequency)
- green region: allowed parameter space from maximum pulsar massL
- corrections from interactions are needed $(a_4< 1)$ to be compatible with D observations!

Quark Star Masses: effects of quark pairing

Add to ^a free gas of quarks terms from interaction, from pairing andfrom an vacuum energy:

$$
\Omega_{CFL} = \frac{6}{\pi^2} \int_0^{\nu} dp \, p^2 (p - \mu) + \frac{3}{\pi^2} \int_0^{\nu} dp \, p^2 (\sqrt{p^2 + m_s^2} - \mu)
$$

$$
+ (1 - a_4) \frac{3\mu^4}{4\pi^2} - \frac{3\Delta^2 \mu^2}{\pi^2} + B_{eff}
$$

where $\nu = 2\mu$ $\sqrt{\mu^2}$ $- m_z^2$ $_{s}^{2}/3.$

- Δ : gap energy of the color-superconducting phase \bullet (normally $\Delta\leq100$ MeV)
- fix strange quark mass to m_s $_{s} = 100$ MeV \bullet
- set for simplicity a_4 $_4 = 0$

Quark Star Masses: effects of quark pairing

- two constraints on quark matter: 2-flavour and 3-flavour line
- green region: allowed parameter space from maximum pulsar massD
- a gap of at least $\Delta = 20$ MeV is needed to be compatible with observations £
- pulsar masses above $1.9M_\odot$ start to constrain QCD parameters! D
- additional interactions needed for pulsar masses well above $2.3M_{\odot}$

Hybrid Stars with ^a stiff nuclear EoS

- nuclear phase: relativistic mean field model with parameter set NL3 (fitted to∙ properties of nuclei)
- match with Gibbs (lines) or Maxwell construction (shaded area)D
- solid lines: pure quark matter cores, dashed lines: mixed phase cores£

Hybrid Stars with ^a soft nuclear EoS

- \bullet nuclear phase: relativistic mean field model with parameter set TM1 (fitted toproperties of nuclei)
- match with Gibbs (lines) or Maxwell construction (shaded area)D
- solid lines: pure quark matter cores, dashed lines: mixed phase cores£
- no pure quark cores compatible with data for ^a soft nuclear EoS

Hybrid Stars with ^a NJL model

(Bonanno and Sedrakian 2011)

- uses Nambu-Jona-Lasinio model for quark matter£
- matches to nuclear EoS with hyperons (RMF with set NL3)L
- 2SC quark matter below green line£
- $\delta=R_{\rm CFL}/R$: amount of CFL quark matter ∙

Color-superconducting quark matter in the NJL model

$$
p = \frac{1}{2\pi^2} \sum_{i=1}^{18} \int_0^{\Lambda} dk \, k^2 |\epsilon_i| + 4K \sigma_u \sigma_d \sigma_s - \frac{1}{4G_D} \sum_{c=1}^3 |\Delta_c|^2
$$

$$
-2G_S \sum_{\alpha=1}^3 \sigma_\alpha^2 + \frac{1}{4G_V} \omega_0^2 + p_e
$$

use Nambu–Jona-Lasinio model for describing quark matter \bullet

- describes both dynamical quark masses (quark condensates $\sigma)$ \bullet and the color-superconducting gaps Δ (Rüster et al. (2005))
- parameters: cutoff, scalar and vector coupling constants $G_S, \, G_V,$ £ diquark coupling G_D , 't Hooft term coupling K
- fixed to hadron masses, pion decay constant, free: G_{D} $_D$ and G_V \bullet

Phases in Quark Matter (Rüster et al. (2005))

- first order phase transition based on symmetry arguments! £
- phases of color superconducting quark matter in β equilibrium: £
- normal (unpaired) quark matter (NQ)L
- two-flavor color superconducting phase (2SC), gapless 2SC phase £
- color-flavor locked phase (CFL), gapless CFL phase, metallic CFL phaseD (Alford, Rajagopal, Wilczek, Reddy, Buballa, Blaschke, Shovkovy, Drago, Rüster, Rischke, Aguilera, Banik, Bandyopadhyay, Pagliara, . . .)) and the contract of the con

QCD Phase Transition and Supernovae

Historical Notes:

- De Rujula 1987: May ^a supernova bang twice? (two neutrino peaks \bullet from SN1987A delayed by ⁵ hours)
- Hatsuda 1987: formation of ^a strange star within 1s! \bullet
- Gentile et al. 1993: hydro simulation with ^a phase transition (second \bullet shock wave, but no neutrinos included)
- Drago and Tambini 1999: prompt bounce by strange quark matter∙ formation
- Nakazato, Sumiyoshi, Yamada 2008: SN simulation for $100M_\odot$ with £ phase transition (no second shock wave)

Proto-neutron star evolution with quarks

(J. Pons, A. Steiner, M. Prakash, J. Lattimer (2001))

- standard lore for the onset \bullet of the quark phase in core-collapse supernovae: during evolution of theproto-neutron star
- timescale for quark matter∙ to appear(see volume fraction χ): typically $(5\,$ -20)s (due to ^a large bag $\frac{1}{\sqrt{2}}$ constant, $B^{1/4}>1$ 4 4 $>$ 180 MeV!)
	- supernova collapsetimescale: milliseconds(with SASI 600 ms?)
- quark matter appears well ◢ after bounce?

Phase Transition to Quark Matter for Astros

(Irina Sagert and Giuseppe Pagliara)

- quark matter appears at low density due to β -equilibrium for a bag constant of £ $B^{1/4} = 165 \text{ MeV}$
- low critical density for low Y_p due to nuclear asymmetry energy ∙
- quark matter favoured at finite temperature£

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- low critical density for low Y_p due to nuclear asymmetry energy ∙
- quark matter favoured at finite temperature \bullet
- production of quark matter in supernovae at bounce possible! £

Check: Mass-Radius Diagram of Cold Neutron Stars

(Sagert, Fischer, Hempel, Pagliara, JSB, Thielemann, Liebendörfer 2011)

- presence of quark matter can change drastically the mass-radius diagram£
- $\frac{1}{\sqrt{2}}$ $\frac{1}{\sqrt{2}}$ maximum mass: $1.56 M_\odot$ $(B^1$ $^4 = 162$ MeV), $1.5 M_{\odot}$ (B^1) $^4=165$ MeV) € \rightarrow too low! need α_s corrections!

Check: Phase Transition for Heavy-Ion Collisions

(Irina Sagert and Giuseppe Pagliara)

- no β -equilibrium (just up-/down-quark matter) \bullet
- large critical densities in particular for isospin-symmetric matter£ (proton fraction $Y_p=0.5)$
- production of ud-quark matter unfavoured for HICs at small T and high density ∙
- no contradiction with heavy-ion data! p_{max} is a set of the set

Implications for Supernovae – Explosion!

(Sagert, Hempel, Pagliara, JSB, Fischer, Mezzacappa, Thielemann, Liebendörfer, 2009)

- velocity profile of ^a supernova for different times (around 250ms) \bullet
- formation of ^a core of pure quark matter produces ^a second shock waveL

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velocity profile of ^a supernova for different times (around 250ms) \bullet

formation of ^a core of pure quark matter produces ^a second shock wave

leads to an delayed explosion \bullet

Implications for Supernova – Neutrino-Signal!

(Sagert, Hempel, Pagliara, JSB, Fischer, Mezzacappa, Thielemann, Liebendörfer, 2009)

- temporal profile of the emittedneutrinos out of the supernova
- thick lines: without, thin lines: ∙ with ^a phase transition
- pronounced second peak of anti-neutrinos due to theformation of quark matter
- peak location and height determined by the critical density andstrength of the QCD phase transition

Supernova Explosion – Parameter dependence

^a moment of black hole formation

^bblack hole formation before positive explosion energy is achieved

(Sagert, Hempel, Pagliara, JSB, Fischer, Mezzacappa, Thielemann, Liebendörfer, 2009)

- supernova simulation runs for different parameters
- appearance of the quark core at $t_\mathrm{pb}=200$ to 500 ms
- results (t_pb , baryonic mass and explosion energy) are significantly ∙ sensitive to the location of the QCD phase transition (bag constant)
- problem: so far explosion only for cases with too low ^a maximum mass!

Detection with Neutrino Detectors

(Dasgupta, Fischer, Horiuchi, Liebendörfer, Mirizzi, Sagert, JSB, arXiv:0912.2568)

- detection of neutrinos from ^a SN with SuperK (left) and IceCube (right)£
- mostly sensitive to antineutrinos by inverse β decay reactions $(\bar{\nu}_e P \to n e^+)$
- take spectrum from supernova simulation, SN at distance of 10 kpc£
- highly sensitive to second burst from QCD phase transition!

The Future: CBM@FAIR and NICA

(Klähn, Blaschke, Weber 2011)

- left: equation of state and flow constraints, £ right: compatible mass-radius relations and astrophysical constraints
- higher baryon densities achieved at higher bombarding energy∙
- probing densities beyond $2-3n_0$
- QCD phase transition can occur in the core of neutron \bullet stars
	- ═⇒ new family of compact stars possible, explosive
phenomena phenomena
- **•** transition can be present during a supernova, shortly after the first bounce⇒ second shock forms, visible in a a second peak in
the (anti-)neutrino signal_gravitational waves (?) the (anti-)neutrino signal, gravitational waves (?), r-process nucleosynthesis (?) . . .
- **o** to stimulate your fancies: color superconducting phase change transport properties – implications forSN neutrino spectra?

X-Ray burster EXO 0748–676 and Quark Matter

- analysis of Özel (Nature 2006): $M \geq 2.10 \pm 0.28 M_{\odot}$ and $R \geq 13.8 \pm 1.8$ km, \bullet claims: 'unconfined quarks do not exist at the center of neutron stars'!
- reply by Alford, Blaschke, Drago, Klähn, Pagliara, JSB (Nature 445, E7 \bullet (2007)): limits rule out soft equations of state, not quark stars or hybrid stars!
- multiwavelength analysis of Pearson et al. (2006): data more consistent with \bullet $M=1.35M_\odot$ than with $M=2.1M_\odot$

Fits to X-Ray Burster Spectra

(Suleimanov, Poutanen, Revnivtsev, Werner 2011)

- x-ray burster with photospheric radius expansion
- assume (color-corrected) black-body emission and Eddington flux at 'touch-down' (Ozel 2006): simple model fit fails above ^a certain distance!
- large correction from model atmosphere composition

Mass-Radius Constraints from X-Ray Burster and Binaries

(Steiner, Lattimer, Brown 2011)

- fit to three x-ray burster data with photospheric radius expansion and three€ quiescent x-ray binaries (from previous analysis)
- relax constraint at 'touch-down' to be on the surface $(r_{ph}\gg R)$
- strong constraint on radius relation (left: combined fit, right: separate fits)_{- p.44}

Mass-Radius Constraints from Isolated Neutron Stars

(Hambaryan, Suleimanov, Schwope, Neuhäuser, Werner, Potekhin 2011)

- isolated neutron star, pulses in x-rays£
- phase space resolved x-ray spectroscopyD
- fit to geometry of hot spot etc. including redshift z
- resulting compactness: $(M/M_{\odot})/(R/{\rm km}) = 0.087 \pm 0.004$ £
- indication for ^a stiff equation of state

RXJ 1856: Neutron Star or Quark Star? (Trümper et al. (2003), Ho et al. (2007))

- two-component blackbody: small soft temperature, so as not to spoil the x-ray£
- this implies ^a rather LARGE radius so that the optical flux is right! £
- lower limit for radiation radius: $R_\infty = R/\sqrt{1-\frac{4}{3}}$ $2GM/R = 17$ km (d/140 pc)
- from parallax measurement: distance $d = 123(+11, -15)pc$ ∙ (Walter, Eisenbeiss, Lattimer, Kim, Hambaryan, Neuhaeuser 2011)

Neutron Star Radii from Neutron Star Mergers

(Bauswein and Janka, 2012)

- neutron star merger simulation with 3D smoothed particle hydro code using \bullet conformal flatness approximation
- strong correlation with peak frequency in gravitational waves and neutron starD radius rather insensitive to masses of neutron stars
- measurable with advanced LIGO in ^a few years

Strangeness in Supernova Matter: Hyperons

C. Ishizuka, A. Ohnishi, K. Tsubakihara, K. Sumiyoshi, S. Yamada (2008)

- supernova matter for $Y_c = 0.4$ with constant entropy/baryon ratio S/B £
- hyperon fraction at bounce $T\sim20$ MeV: about 0.1% D
- thermally produced strangeness, hyperons are in β -equilibrium! £

Nucleation Timescales for strange quark matter

(B. W. Mintz, E. Fraga, G. Pagliara, JSB, arXiv:0910.3927)

- nucleation of strange quark matter via fluctuations in strangeness∙
- timescales for different surface tensions and densitiesD 4 $\frac{4}{i}/(4\pi^2$ (quark EoS used: $p=(1\,$ $- \,c) \mu$ $^{2}))$

bubble nucleation within 1 km 3 within 100 ms for $\sigma < 20$ MeV fm $^{-2}$ ∙

Gravitational Wave Amplitudes

(Abdikamalov, Dimmelmeier, Rezzolla, Miller 2008)

- amplitude of gravitational wave signal from collapsing neutron stars at 10kpc
- above sensitivity of present (LIGO, VIRGO) and well above future (Advanced£ LIGO) detector
- events in Virgo cluster (20 Mpc) needs probably third generation detectors

Gravitational Wave Background from Phase Transitions

(Sigl 2007)

- Phase transitions in neutron stars generate gravitional waves (blue band, for∙ anisotropic neutrino emission: solid blue line)
- background of such gravitational waves detectable with future space€ detectors!
- signal larger than the one for conventional type II Supernovae (dashed line)and from inflation (dash-dotted lines)) and the contract of the con

Gravitational wave signals from hybrid stars

(Oechslin, Uryū, Pogosyan, Thielemann 2004)

- **Fourier spectra of gravitational** waves
- **•** increasing initial mass from top to bottom
	- solid line: neutron star, dashedline: hybrid star
- different spectra for hybrid \bullet stars!