



New Developments in
Collective Neutrino Oscillations
in Supernovae

Georg G. Raffelt

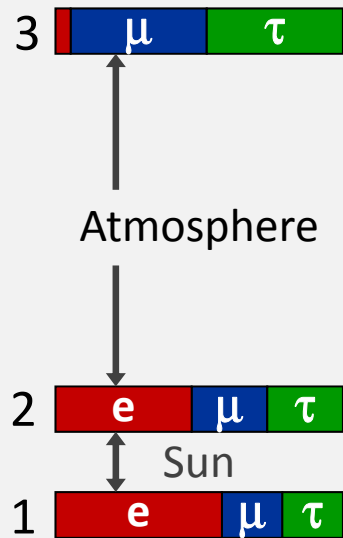
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Three-Flavor Neutrino Parameters

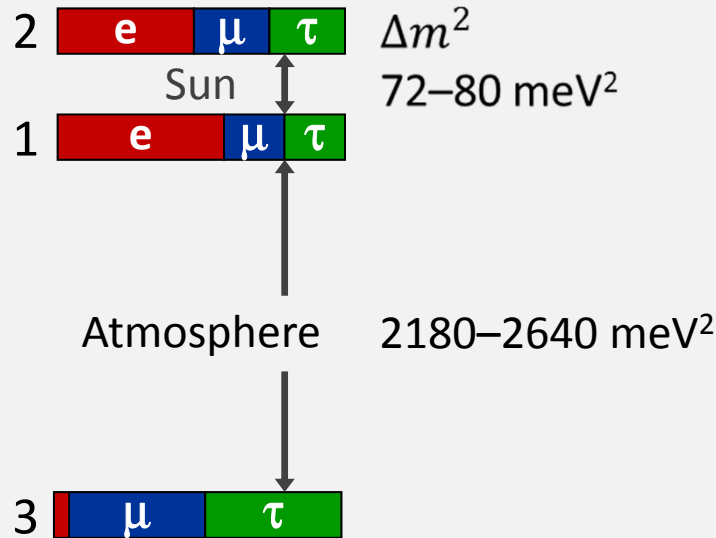
Three mixing angles $\theta_{12}, \theta_{13}, \theta_{23}$ (Euler angles for 3D rotation), $c_{ij} = \cos \theta_{ij}$, a CP-violating “Dirac phase” δ , and two “Majorana phases” α_2 and α_3

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{39^\circ < \theta_{23} < 53^\circ \text{ Atmospheric/LBL-Beams}} \underbrace{\begin{pmatrix} c_{13} & 0 & e^{-i\delta} s_{13} \\ 0 & 1 & 0 \\ -e^{i\delta} s_{13} & 0 & c_{13} \end{pmatrix}}_{7^\circ < \theta_{13} < 11^\circ \text{ Reactor}} \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{33^\circ < \theta_{12} < 37^\circ \text{ Solar/KamLAND}} \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\frac{\alpha_2}{2}} & 0 \\ 0 & 0 & e^{i\frac{\alpha_3}{2}} \end{pmatrix}}_{\text{Relevant for } 0\nu 2\beta \text{ decay}} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Normal



Inverted



Tasks and Open Questions

- Precision for all angles
- CP-violating phase δ ?
- Mass ordering?
(normal vs inverted)
- Absolute masses?
(hierarchical vs degenerate)
- Dirac or Majorana?

3400 citations

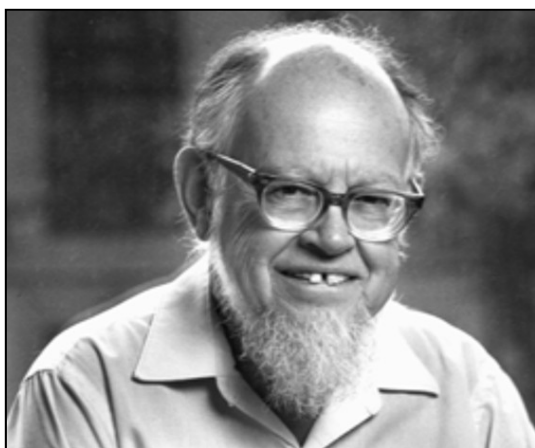
Neutrino oscillations in matter

L. Wolfenstein

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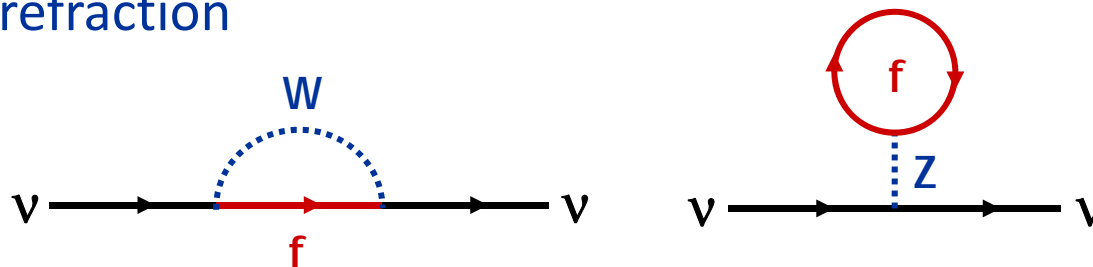
(Received 6 October 1977; revised manuscript received 5 December 1977)

The effect of coherent forward scattering must be taken into account when considering the oscillations of neutrinos traveling through matter. In particular, for the case of massless neutrinos for which vacuum oscillations cannot occur, oscillations can occur in matter if the neutral current has an off-diagonal piece connecting different neutrino types. Applications discussed are solar neutrinos and a proposed experiment involving transmission of neutrinos through 1000 km of rock.



Lincoln Wolfenstein

Neutrinos in a medium suffer flavor-dependent refraction

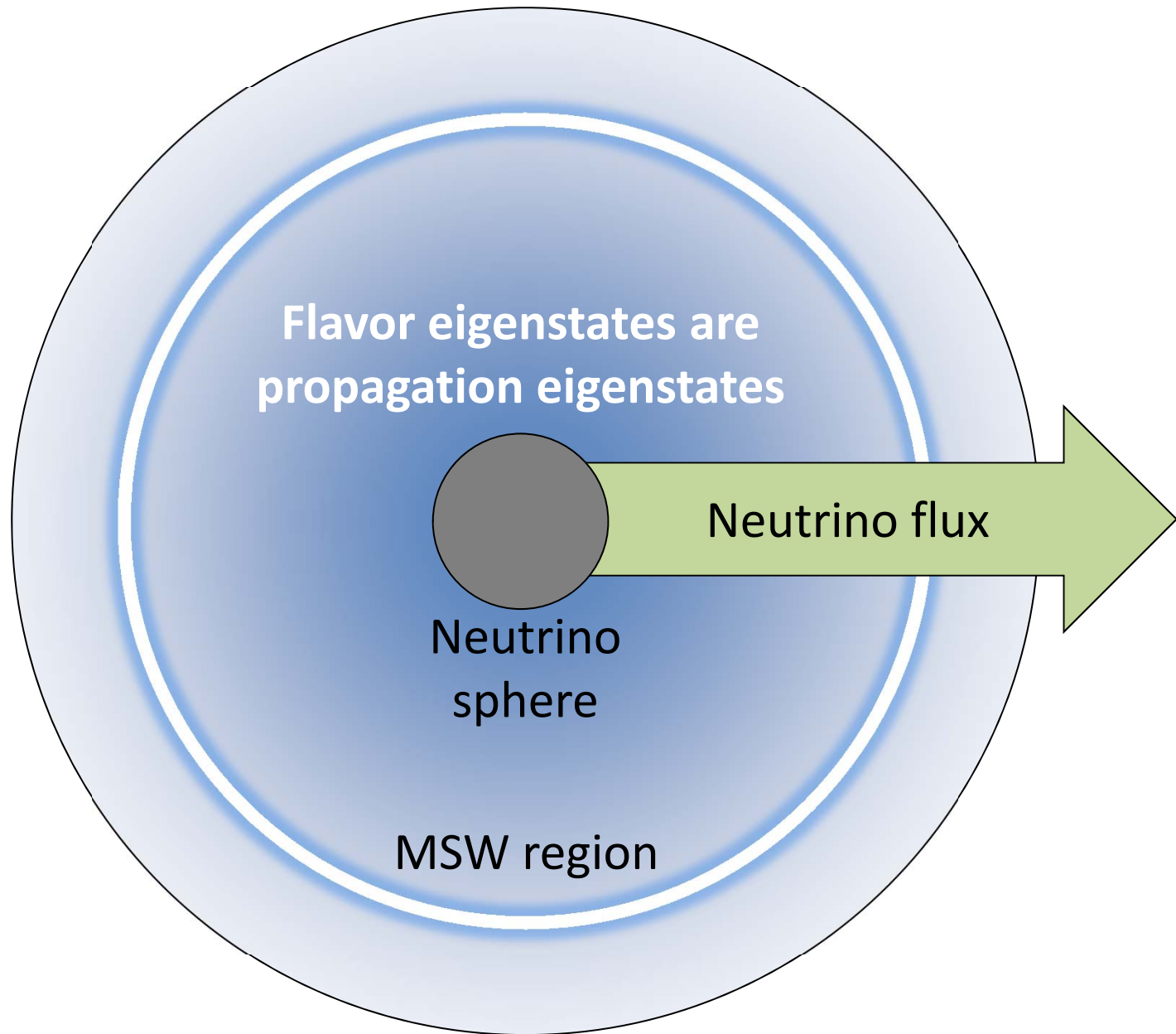


$$V_{\text{weak}} = \sqrt{2}G_F \times \begin{cases} N_e - N_n/2 & \text{for } \nu_e \\ -N_n/2 & \text{for } \nu_\mu \end{cases}$$

Typical density of Earth: 5 g/cm³

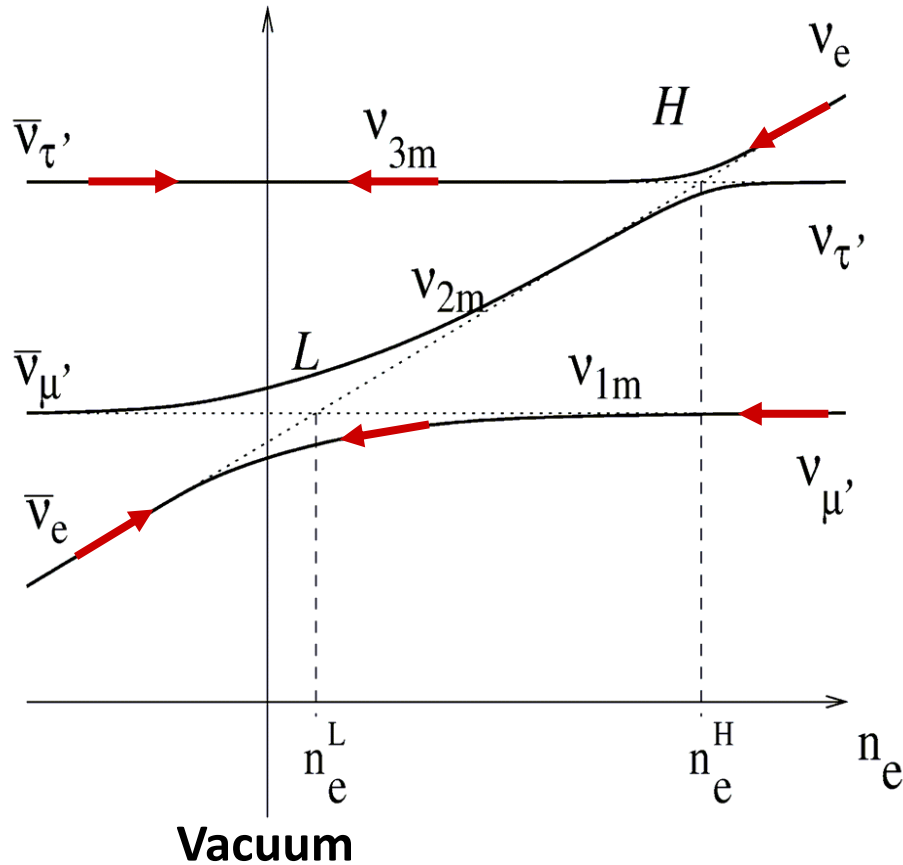
$$\Delta V_{\text{weak}} \approx 2 \times 10^{-13} \text{ eV} = 0.2 \text{ peV}$$

Flavor Oscillations in Core-Collapse Supernovae

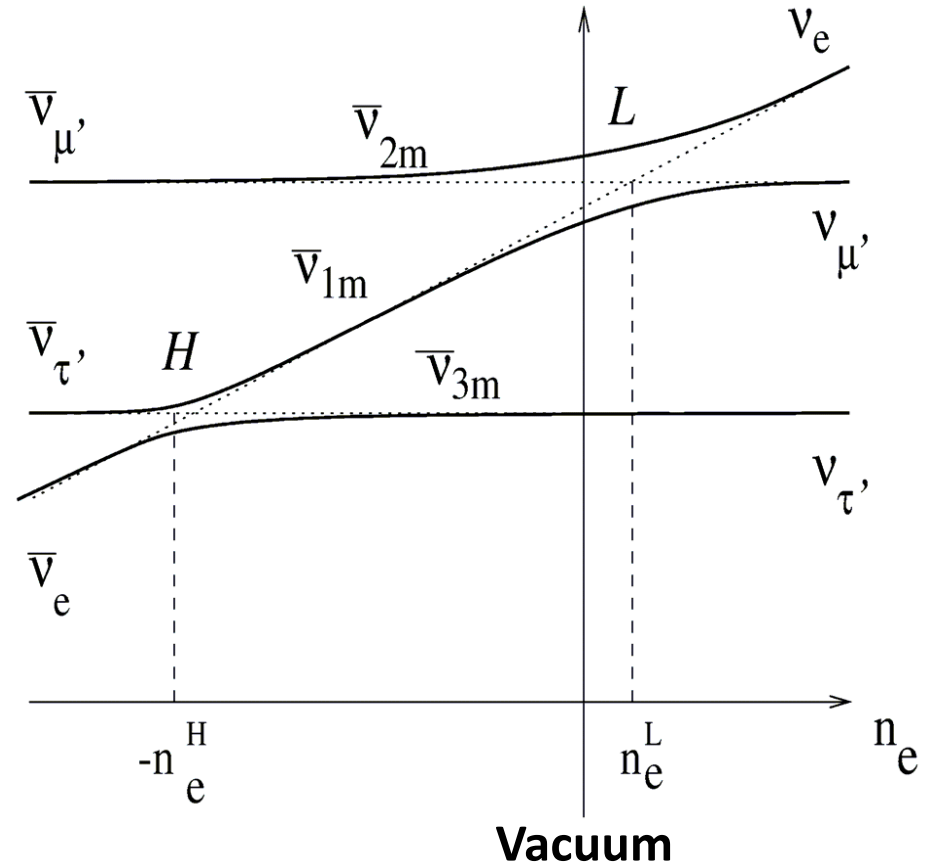


Three-Flavor Eigenvalue Diagram

Normal mass hierarchy



Inverted mass hierarchy



Dighe & Smirnov, Identifying the neutrino mass spectrum from a supernova neutrino burst, astro-ph/9907423

Signature of Flavor Oscillations

	1-3-mixing scenarios		
	A	B	C
Mass ordering	Normal (NH)	Inverted (IH)	Any (NH/IH)
$\sin^2 \theta_{13}$	$\gtrsim 10^{-3}$		$\lesssim 10^{-5}$
MSW conversion	adiabatic		non-adiabatic
ν_e survival prob.	0	$\sin^2 \theta_{12} \approx 0.3$	$\sin^2 \theta_{12} \approx 0.3$
$\bar{\nu}_e$ survival prob.	$\cos^2 \theta_{12} \approx 0.7$	0	$\cos^2 \theta_{12} \approx 0.7$
$\bar{\nu}_e$ Earth effects	Yes	No	Yes
May distinguish mass ordering			

Assuming collective effects are not important during accretion phase
(Chakraborty et al., arXiv:1105.1130, Sarikas et al. arXiv:1109.3601)

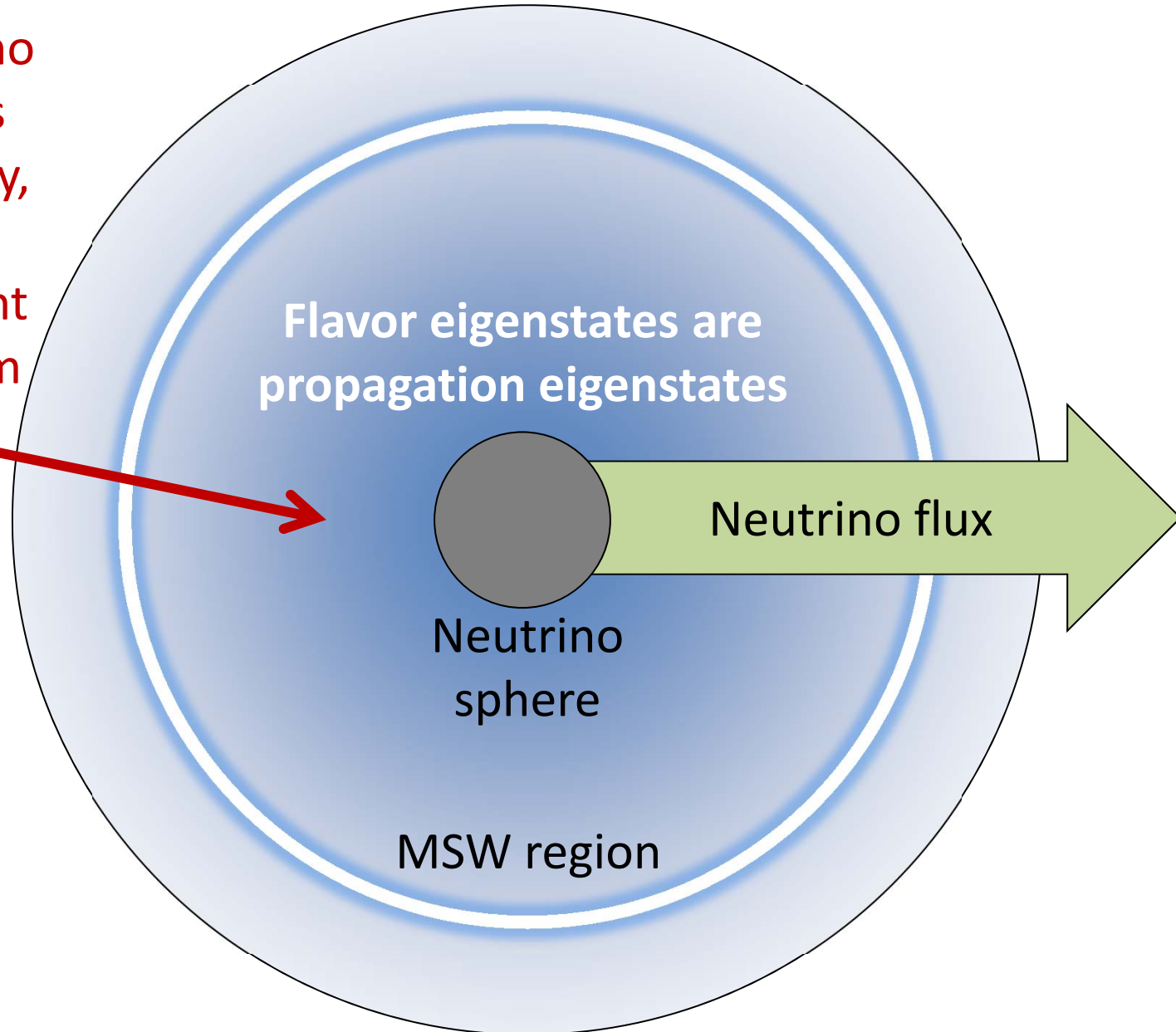
Signature of Flavor Oscillations

	1-3-mixing scenarios		
	A	B	C
Mass ordering	Normal (NH)	Inverted (IH)	Any (NH/IH)
$\sin^2 \theta_{13}$	$\theta_{13} \approx 9^\circ$		$\lesssim 10^{-5}$
MSW conversion	adiabatic		non-adiabatic
ν_e survival prob.	0	$\sin^2 \theta_{12} \approx 0.3$	$\sin^2 \theta_{12} \approx 0.3$
$\bar{\nu}_e$ survival prob.	$\cos^2 \theta_{12} \approx 0.7$	0	$\cos^2 \theta_{12} \approx 0.7$
$\bar{\nu}_e$ Earth effects	Yes	No	Yes
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Assuming collective effects are not important during accretion phase
(Chakraborty et al., arXiv:1105.1130, Sarikas et al. arXiv:1109.3601)

Flavor Oscillations in Core-Collapse Supernovae

Neutrino-neutrino refraction causes a flavor instability, flavor exchange between different parts of spectrum

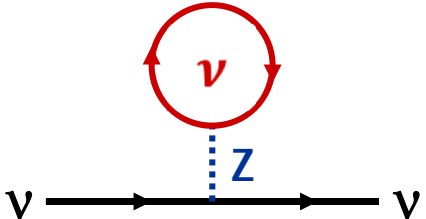


Flavor-Off-Diagonal Refractive Index

2-flavor neutrino evolution as an effective 2-level problem

$$i \frac{\partial}{\partial t} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = H \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}$$

Effective mixing Hamiltonian

$$H = \frac{M^2}{2E} + \sqrt{2}G_F \begin{pmatrix} N_e - \frac{N_n}{2} & 0 \\ 0 & -\frac{N_n}{2} \end{pmatrix} + \sqrt{2}G_F \begin{pmatrix} N_{\nu_e} & N_{\langle \nu_e | \nu_\mu \rangle} \\ N_{\langle \nu_\mu | \nu_e \rangle} & N_{\nu_\mu} \end{pmatrix}$$


Mass term in flavor basis: causes vacuum oscillations

Wolfenstein's weak potential, causes MSW "resonant" conversion together with vacuum term

Flavor-off-diagonal potential, caused by flavor oscillations. (J.Pantaleone, PLB 287:128,1992)

Flavor oscillations feed back on the Hamiltonian: Nonlinear effects!

Collective Supernova Nu Oscillations since 2006

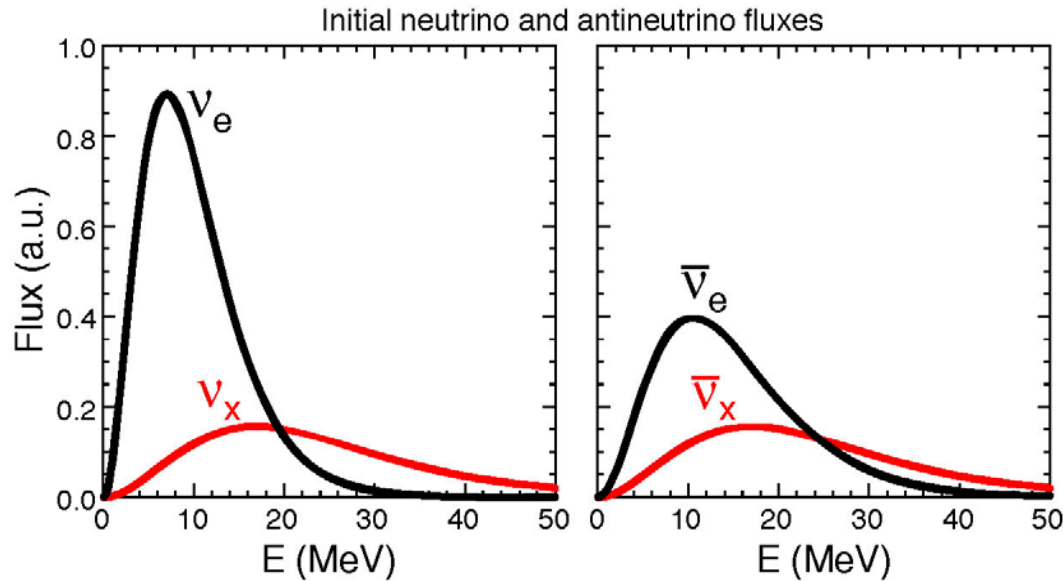
Two seminal papers in 2006 triggered a torrent of activities

Duan, Fuller, Qian, astro-ph/0511275, Duan et al. astro-ph/0606616

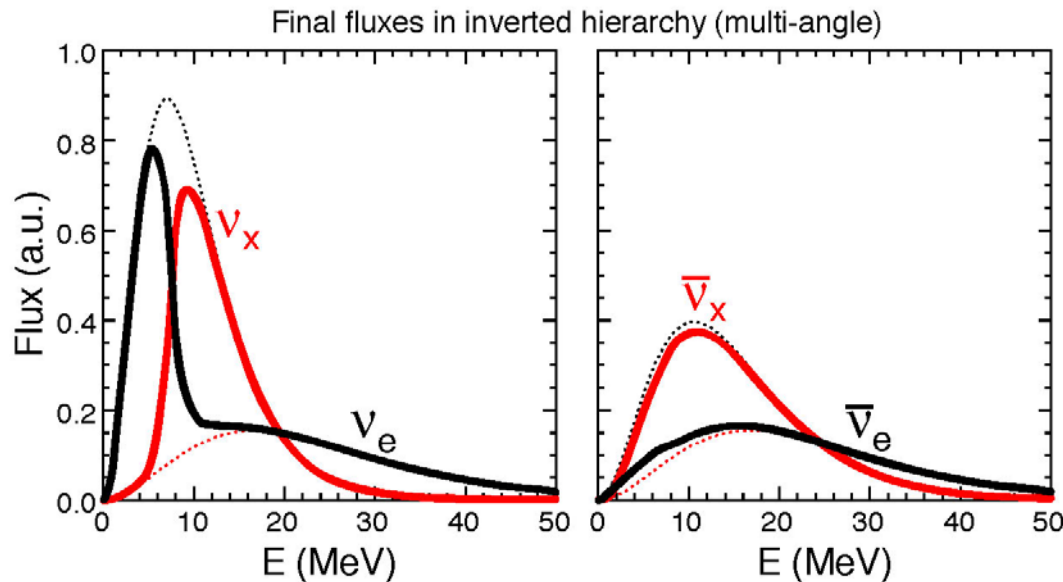
Balantekin, Gava & Volpe [0710.3112]. Balantekin & Pehlivan [astro-ph/0607527]. Blennow, Mirizzi & Serpico [0810.2297]. Cherry, Fuller, Carlson, Duan & Qian [1006.2175, 1108.4064]. Cherry, Wu, Fuller, Carlson, Duan & Qian [1109.5195]. Cherry, Carlson, Friedland, Fuller & Vlasenko [1203.1607]. Chakraborty, Choubey, Dasgupta & Kar [0805.3131]. Chakraborty, Fischer, Mirizzi, Saviano, Tomàs [1104.4031, 1105.1130]. Choubey, Dasgupta, Dighe & Mirizzi [1008.0308]. Dasgupta & Dighe [0712.3798]. Dasgupta, Dighe & Mirizzi [0802.1481]. Dasgupta, Dighe, Raffelt & Smirnov [0904.3542]. Dasgupta, Dighe, Mirizzi & Raffelt [0801.1660, 0805.3300]. Dasgupta, Mirizzi, Tamborra & Tomàs [1002.2943]. Dasgupta, Raffelt & Tamborra [1001.5396]. Dasgupta, O'Connor & Ott [1106.1167]. Duan, Fuller, Carlson & Qian [astro-ph/0608050, 0703776, 0707.0290, 0710.1271]. Duan, Fuller & Qian [0706.4293, 0801.1363, 0808.2046, 1001.2799]. Duan, Fuller & Carlson [0803.3650]. Duan & Kneller [0904.0974]. Duan & Friedland [1006.2359]. Duan, Friedland, McLaughlin & Surman [1012.0532]. Esteban-Pretel, Mirizzi, Pastor, Tomàs, Raffelt, Serpico & Sigl [0807.0659]. Esteban-Pretel, Pastor, Tomàs, Raffelt & Sigl [0706.2498, 0712.1137]. Fogli, Lisi, Marrone & Mirizzi [0707.1998]. Fogli, Lisi, Marrone & Tamborra [0812.3031]. Friedland [1001.0996]. Gava & Jean-Louis [0907.3947]. Gava & Volpe [0807.3418]. Galais, Kneller & Volpe [1102.1471]. Galais & Volpe [1103.5302]. Gava, Kneller, Volpe & McLaughlin [0902.0317]. Hannestad, Raffelt, Sigl & Wong [astro-ph/0608695]. Wei Liao [0904.0075, 0904.2855]. Lunardini, Müller & Janka [0712.3000]. Mirizzi, Pozzorini, Raffelt & Serpico [0907.3674]. Mirizzi & Serpico [1111.4483]. Mirizzi & Tomàs [1012.1339]. Pehlivan, Balantekin, Kajino & Yoshida [1105.1182]. Pejcha, Dasgupta & Thompson [1106.5718]. Raffelt [0810.1407, 1103.2891]. Raffelt & Sigl [hep-ph/0701182]. Raffelt & Smirnov [0705.1830, 0709.4641]. Raffelt & Tamborra [1006.0002]. Sawyer [hep-ph/0408265, 0503013, 0803.4319, 1011.4585]. Sarikas, Raffelt, Hüdepohl & Janka [1109.3601]. Sarikas, Tamborra, Raffelt, Hüdepohl & Janka [1204.0971]. Saviano, Chakraborty, Fischer, Mirizzi [1203.1484]. Wu & Qian [1105.2068].

Spectral Split

Initial
fluxes at
neutrino
sphere



After
collective
trans-
formation



Figures from
Fogli, Lisi,
Marrone & Mirizzi,
arXiv:0707.1998

Explanations in
Raffelt & Smirnov
arXiv:0705.1830
and 0709.4641
Duan, Fuller,
Carlson & Qian
arXiv:0706.4293
and 0707.0290

Three Ways to Describe Flavor Oscillations

Schrödinger equation in terms of “flavor spinor”

$$i\partial_t \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = H \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \frac{\Delta m^2}{2E} \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{pmatrix} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}$$

Neutrino flavor density matrix

$$\rho = \begin{pmatrix} \langle \nu_e | \nu_e \rangle & \langle \nu_e | \nu_\mu \rangle \\ \langle \nu_\mu | \nu_e \rangle & \langle \nu_\mu | \nu_\mu \rangle \end{pmatrix}$$

Equivalent commutator form of Schrödinger equation

$$i\partial_t \rho = [H, \rho]$$

Expand 2×2 Hermitean matrices in terms of Pauli matrices

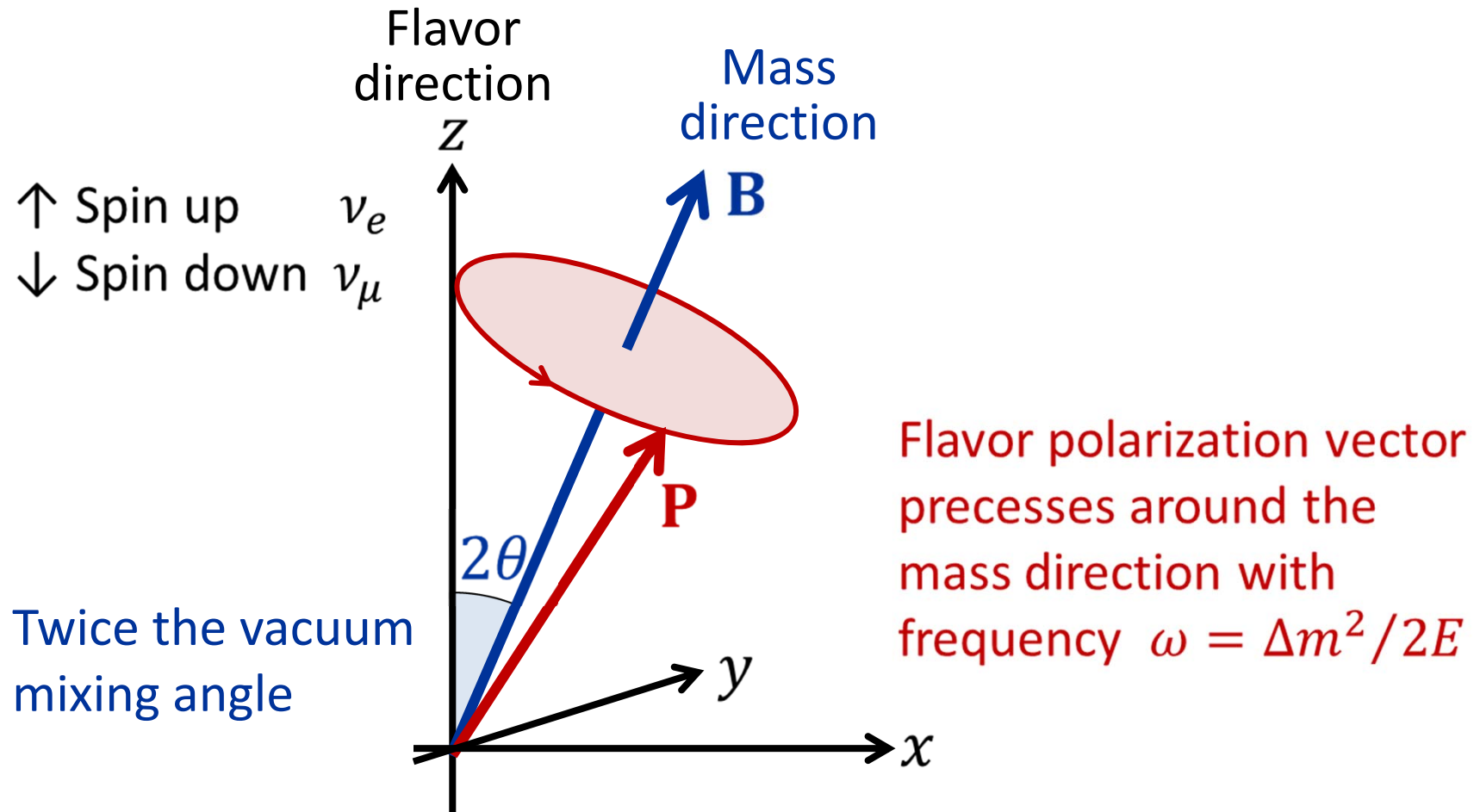
$$\rho = \frac{1}{2} [\text{Tr}(\rho) + \mathbf{P} \cdot \boldsymbol{\sigma}] \quad \text{and} \quad H = \frac{\Delta m^2}{2E} \mathbf{B} \cdot \boldsymbol{\sigma} \quad \text{with} \quad \mathbf{B} = (\sin 2\theta, 0, \cos 2\theta)$$

Equivalent spin-precession form of equation of motion

$$\dot{\mathbf{P}} = \omega \mathbf{B} \times \mathbf{P} \quad \text{with} \quad \omega = \frac{\Delta m^2}{2E}$$

\mathbf{P} is “polarization vector” or “Bloch vector” or “flavor isospin vector”

Flavor Oscillation as Spin Precession



Collective Nu Oscillations as a Many-Body Problem

Hamiltonian for interacting “flavor spins” (*classical* in mean-field approach)

$$H = \sum_{i=1}^N \omega_i \mathbf{B} \cdot \mathbf{P}_i + \sqrt{2} G_F N_e \mathbf{L} \cdot \sum_{i=1}^N \mathbf{P}_i + \mu \sum_{i,j=1}^N (1 - \cos \theta_{ij}) \mathbf{P}_i \cdot \mathbf{P}_j$$

Unit vector
in mass direction

Unit vector
in flavor direction

Multi-angle effects from
current-current structure

“Spin-pairing H” for isotropic system (or single angle), ignoring matter effect

$$H = \sum_{i=1}^N \omega_i \mathbf{B} \cdot \mathbf{P}_i + \mu \mathbf{P}_{\text{tot}}^2$$

BCS theory (using Anderson’s pseudo-spin), nuclear physics, ...

Integrable system (as many “Gaudin invariants” as spins)

→ Pehlivan, Balantekin, Kajino & Yoshida [arxiv:1105.1182] for introduction

N-mode coherent solutions (“Normal and anomalous solitons”)

- Emil Yuzbashian, Phys. Rev. **B** 78, 184507 (2008) Super-conductivity (BCS)
- Georg Raffelt, Phys. Rev. **D** 83, 105022 (2011) Collective Nus

Synchronized Flavor Oscillations

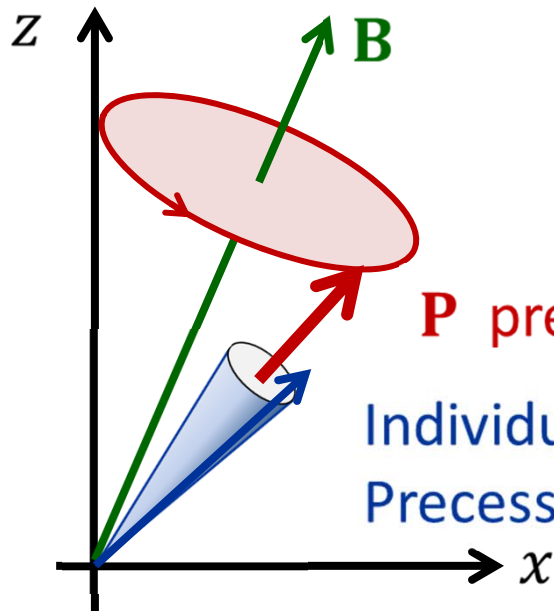
Precession equation for each ν mode with energy E , i.e. $\omega = \Delta m^2/2E$

$$\dot{\mathbf{P}}_\omega = \underbrace{(\omega\mathbf{B} + \lambda\mathbf{L} + \mu\mathbf{P})}_{\mathbf{H}_{\text{eff}}} \times \mathbf{P}_\omega \quad \text{with} \quad \lambda = \sqrt{2}G_{\text{F}}N_e \quad \text{and} \quad \mu = \sqrt{2}G_{\text{F}}N_\nu$$

Total flavor spin of entire ensemble

$$\mathbf{P} = \sum_\omega \mathbf{P}_\omega \quad \text{normalize} \quad |\mathbf{P}_{t=0}| = 1$$

Individual spins do not remain aligned – feel “internal” field $\mathbf{H}_{\nu\nu} = \mu\mathbf{P}$



Synchronized oscillations for large neutrino density $\mu \gg \delta\omega$

\mathbf{P} precesses with ω_{sync} for large ν density

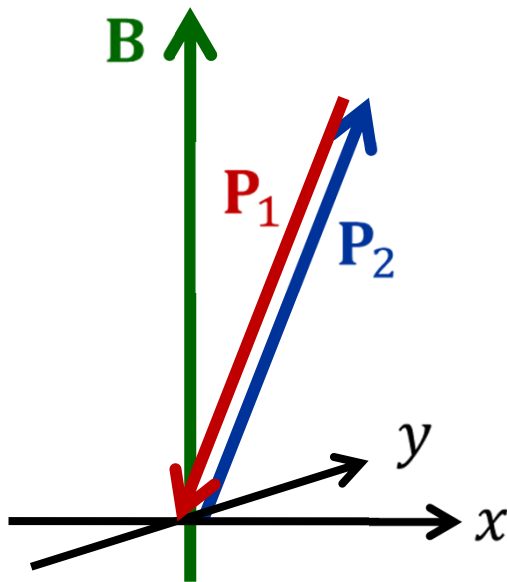
Individual \mathbf{P}_ω “trapped” on precession cones
Precess around \mathbf{P} with frequency $\sim \mu$

Instability in Flavor Space

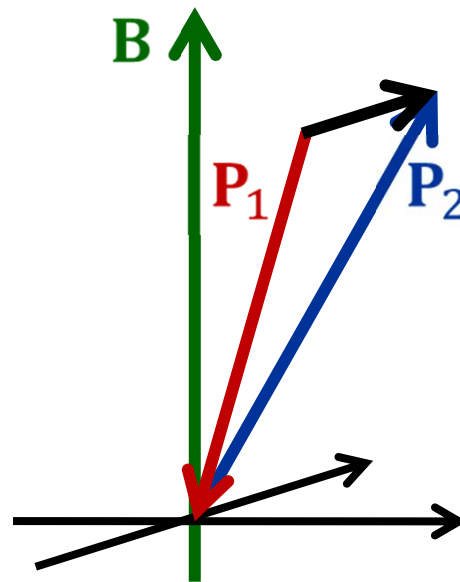
Two-mode example in co-rotating frame, initially $\mathbf{P}_1 = \downarrow$, $\mathbf{P}_2 = \uparrow$ (flavor basis)

$$\dot{\mathbf{P}}_1 = [-\omega \mathbf{B} + \mu (\mathbf{P}_1 + \mathbf{P}_2)] \times \mathbf{P}_1$$

$$\dot{\mathbf{P}}_2 = [+ \omega \mathbf{B} + \underbrace{\mu (\mathbf{P}_1 + \mathbf{P}_2)}_{0 \text{ initially}}] \times \mathbf{P}_2$$



- Initially aligned in flavor direction and $\mathbf{P} = 0$
- Free precession $\pm \omega$



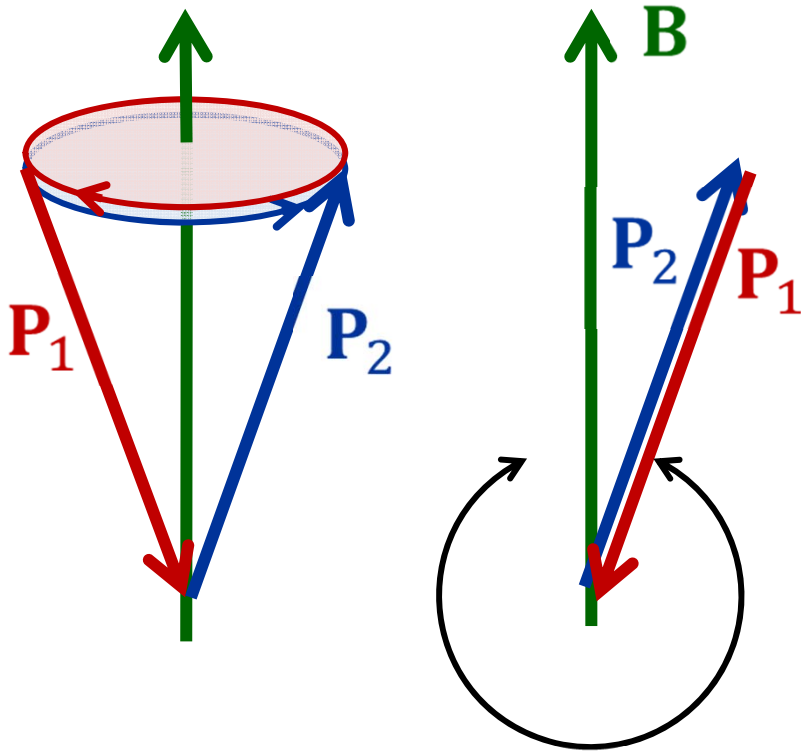
- After a short time, transverse \mathbf{P} develops by free precession

$$\mathbf{P} = \mathbf{P}_1 + \mathbf{P}_2$$

Matter effect transverse to mass and flavor directions
Both \mathbf{P}_1 and \mathbf{P}_2 tilt around \mathbf{P} if μ is large

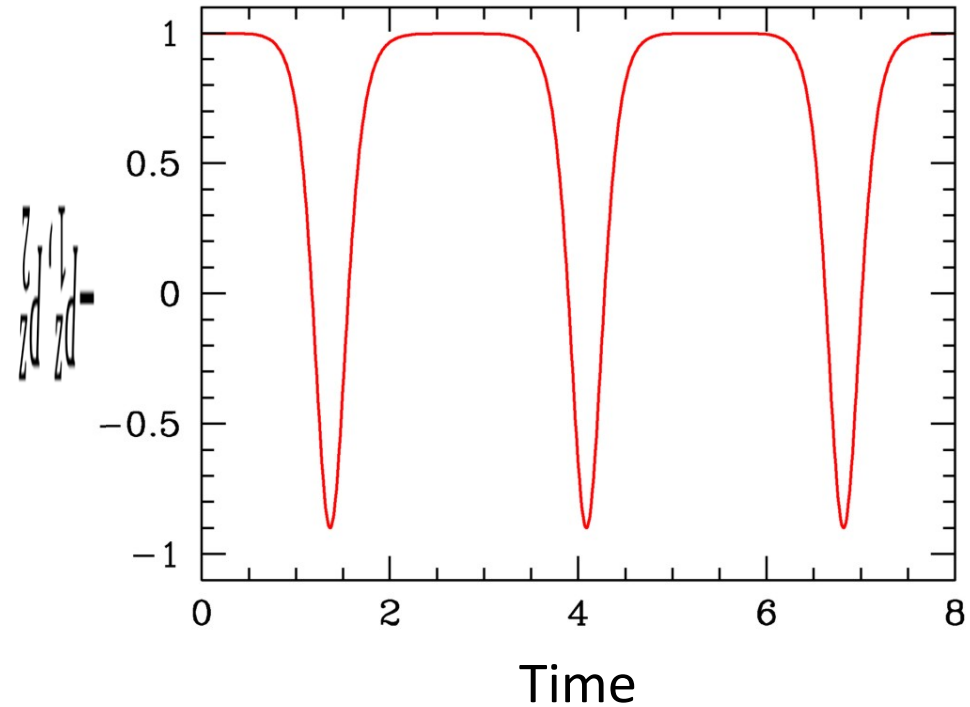
Two Spins with Opposite Initial Orientation

No interaction ($\mu = 0$)
Free precession in
opposite directions



Strong interaction
($\mu \rightarrow \infty$)
Pendular motion

Even for very small mixing angle,
large-amplitude flavor oscillations

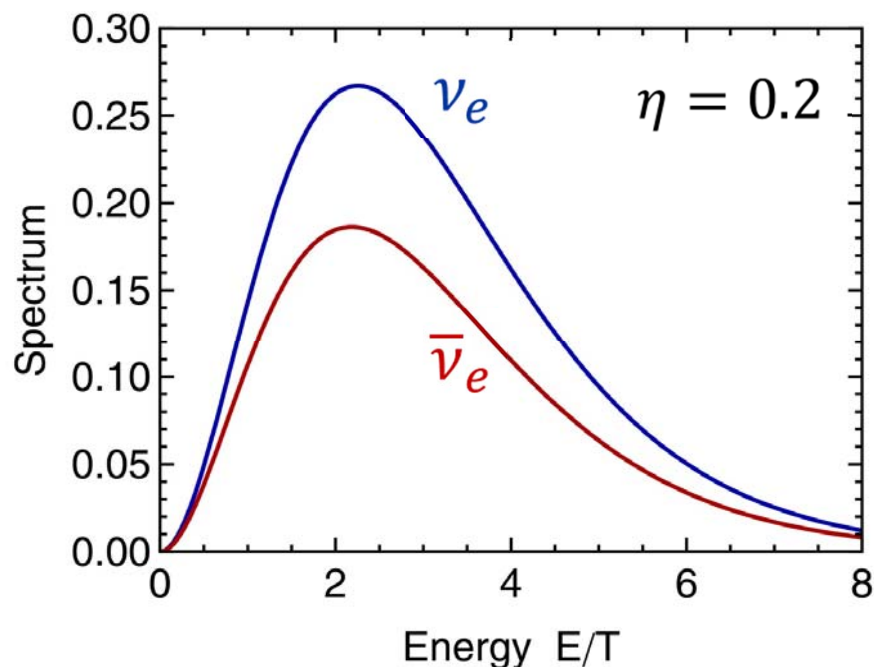


Inverse-Energy Spectrum

Fermi-Dirac energy spectrum

$$\frac{dN}{dE} \propto \frac{E^2}{e^{E/T - \eta} + 1}$$

η degeneracy parameter, $-\eta$ for $\bar{\nu}$



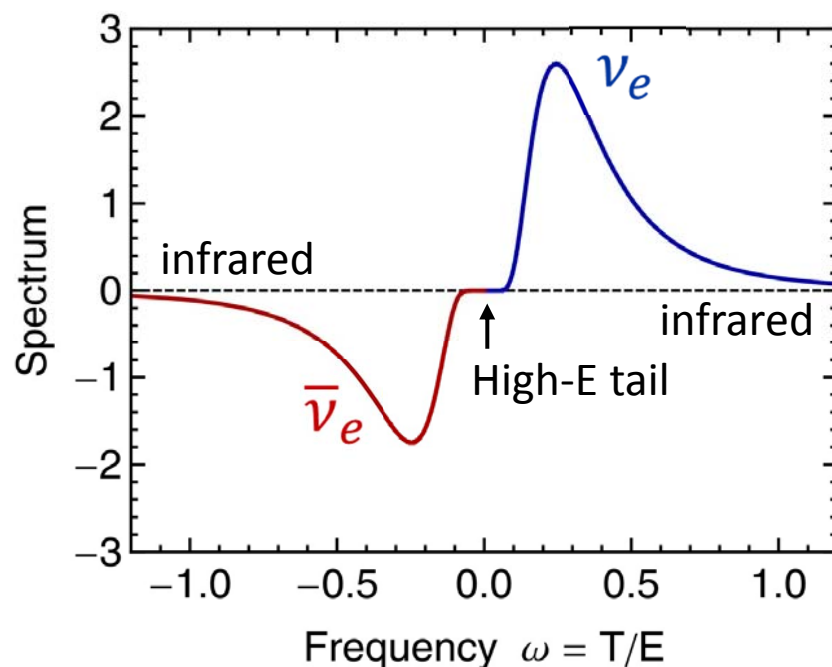
Spectrum in terms of $\omega = T/E$

- Antineutrinos $E \rightarrow -E$
- and dN/dE negative

(flavor isospin convention)

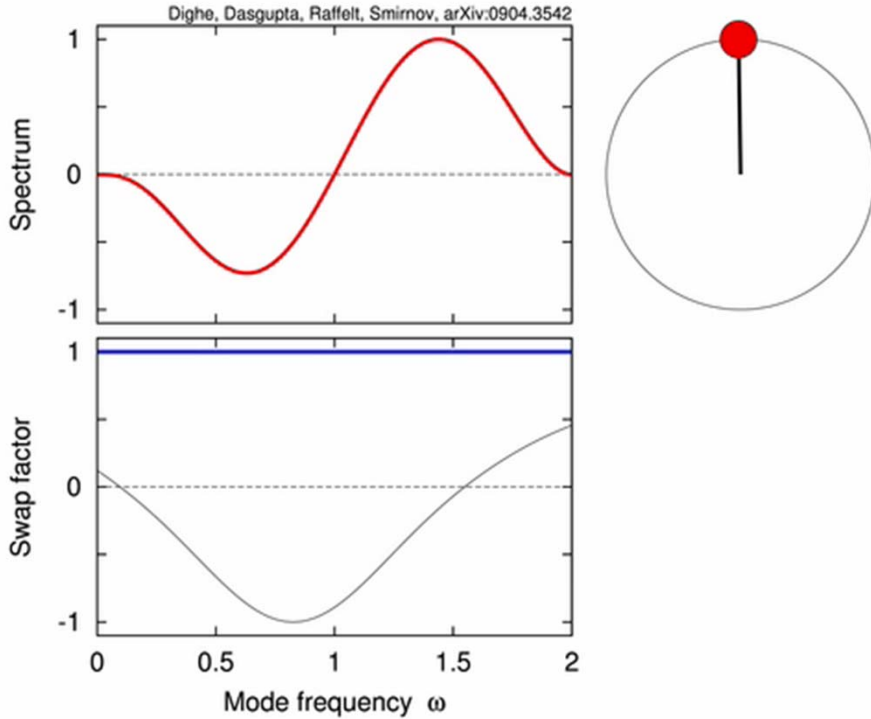
$$\omega > 0: \nu_e = \uparrow \quad \text{and} \quad \nu_\mu = \downarrow$$

$$\omega < 0: \bar{\nu}_e = \downarrow \quad \text{and} \quad \bar{\nu}_\mu = \uparrow$$

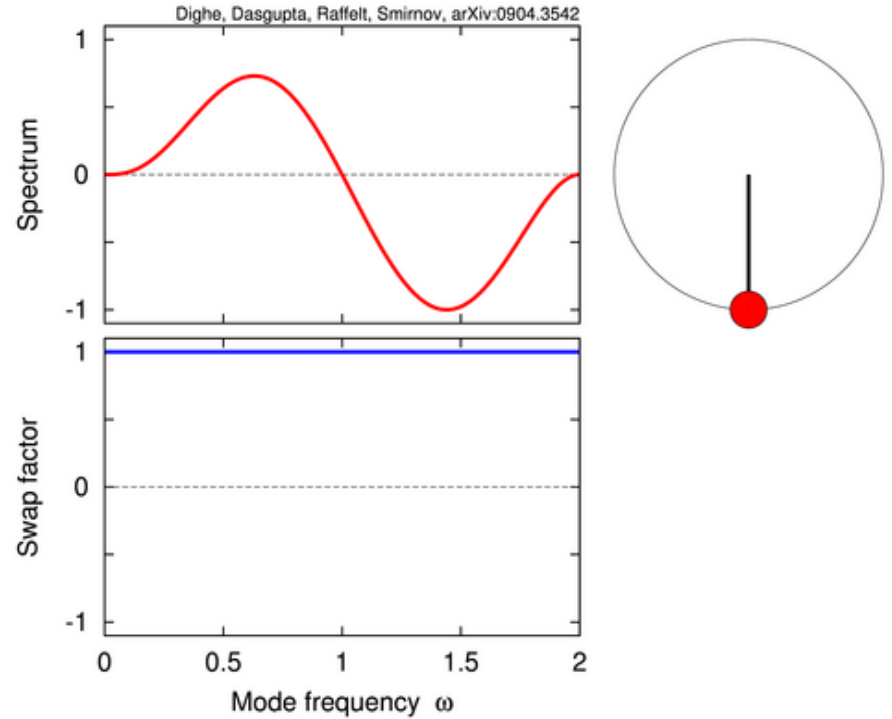


Flavor Pendulum

Single “positive” crossing (IH)
(potential energy at a maximum)



Single “negative” crossing (NH)
(potential energy at a minimum)



Dasgupta, Dighe, Raffelt & Smirnov, arXiv:0904.3542

For movies see <http://www.mppmu.mpg.de/supernova/multisplits>

Self-Induced Flavor Conversion

- Flavor content exchanged between different momentum modes (or ν s and anti- ν s changing together)
 - No net flavor conversion of ensemble
 - Instability required to get started:
Exponential growth of the off-diagonal density matrix parts
- **Linearized Stability Analysis**
(first stressed by Ray Sawyer)

Sawyer, arXiv:0803.4319 – Banerjee, Dighe & Raffelt, arXiv:1107.2308

Linearized Stability Analysis

Schrödinger equation for flavor matrices of neutrino fluxes $\Phi_{\omega,u}$

$\omega = \pm \Delta m^2 / 2E$ $u = \sin^2(\text{emission angle})$ $v_u = \text{radial velocity at } r$

$$i\partial_r \Phi_{\omega,u} = \left[\frac{\omega + \sqrt{2}G_F N_\ell}{v_u} + \frac{\sqrt{2}G_F}{4\pi r^2} \int d\omega' du' \Phi_{\omega',u'} \frac{1 - v_u v_{u'}}{v_u v_{u'}} , \Phi_{\omega,u} \right]$$

Linearize in small off-diagonal flux terms and Fourier transform

$$\Phi_{\omega,u} = \frac{g_{\omega,u}}{2} \begin{pmatrix} 1 & Q_{\omega,u} e^{-i\Omega r} \\ Q_{\omega,u}^* e^{i\Omega r} & -1 \end{pmatrix}$$

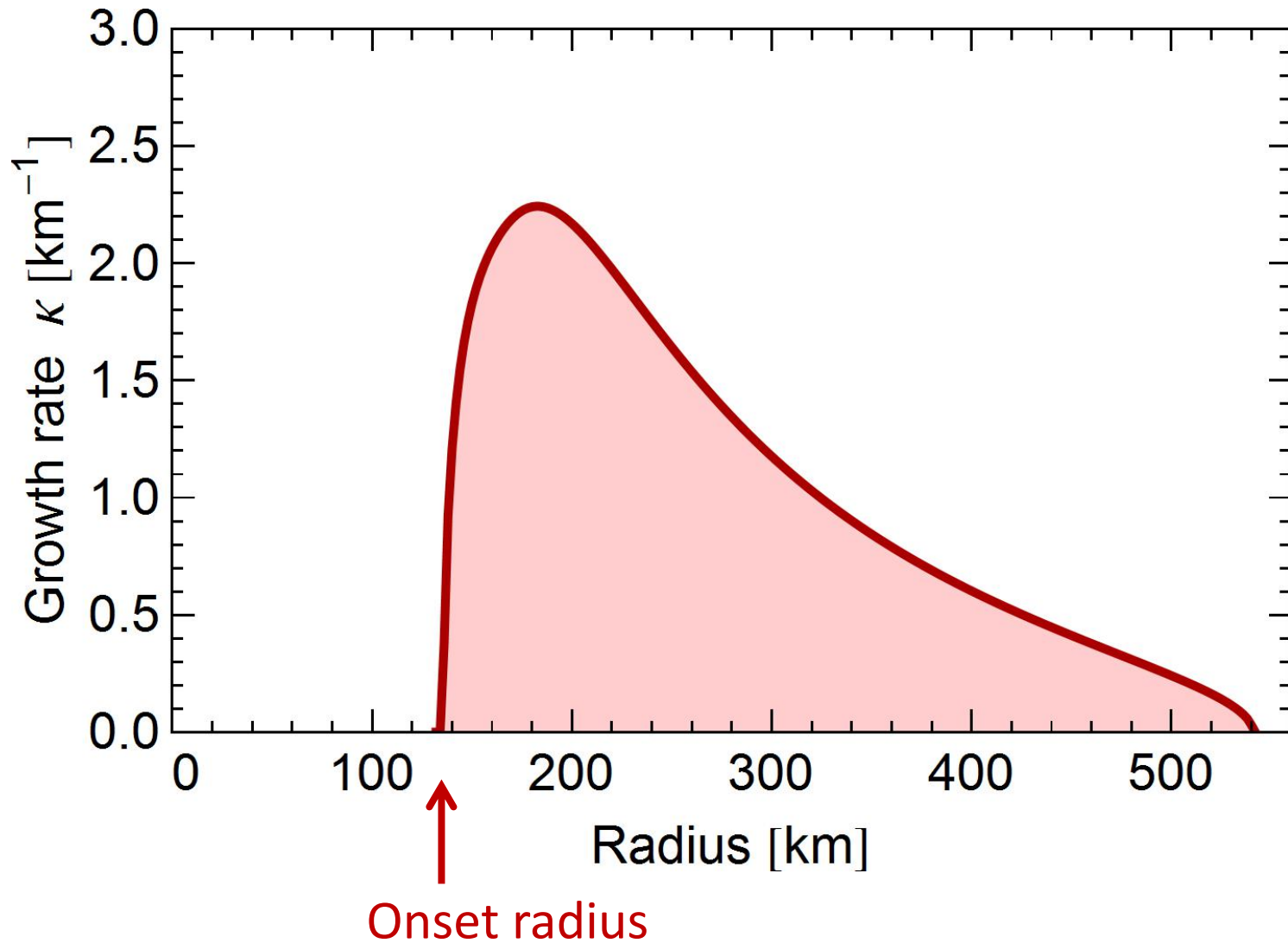
Eigenvalue equation for $Q_{\omega,u}$ in terms of eigenfrequency $\Omega = \gamma + i\kappa$, where κ is the exponential growth rate

$$\left[\omega + u \left(\lambda + \int d\omega' du' g_{\omega',u'} \right) - \Omega \right] Q_{\omega,u} = \mu \int d\omega' du' (u + u') g_{\omega',u'} Q_{\omega',u'}$$

Straightforward to solve for eigenvalue Ω and eigenfunction $Q_{\omega,u}$

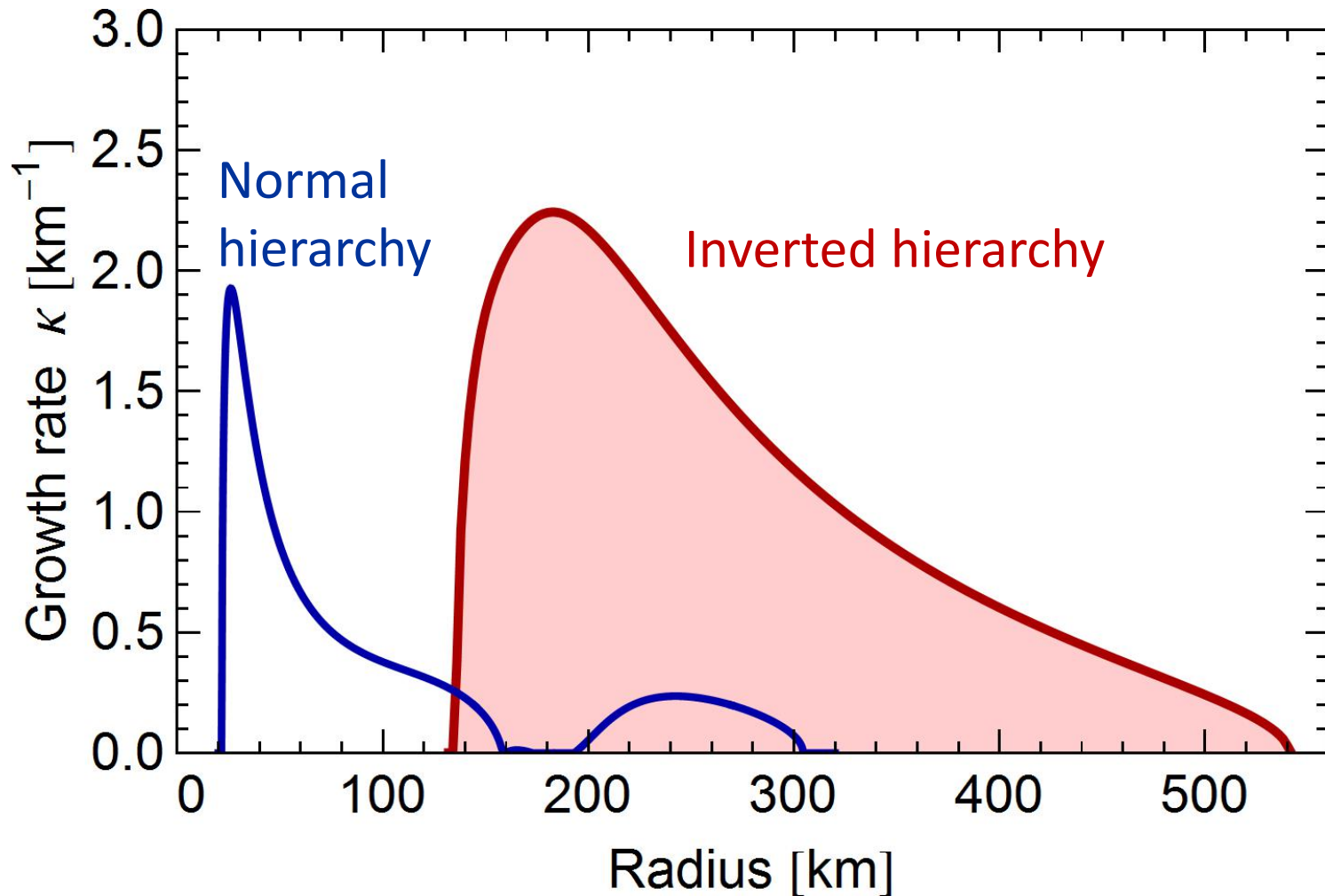
Banerjee, Dighe & Raffelt, arXiv:1107.2308

Stability Analysis for Simple SN Example



Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Normal vs Inverted Hierarchy

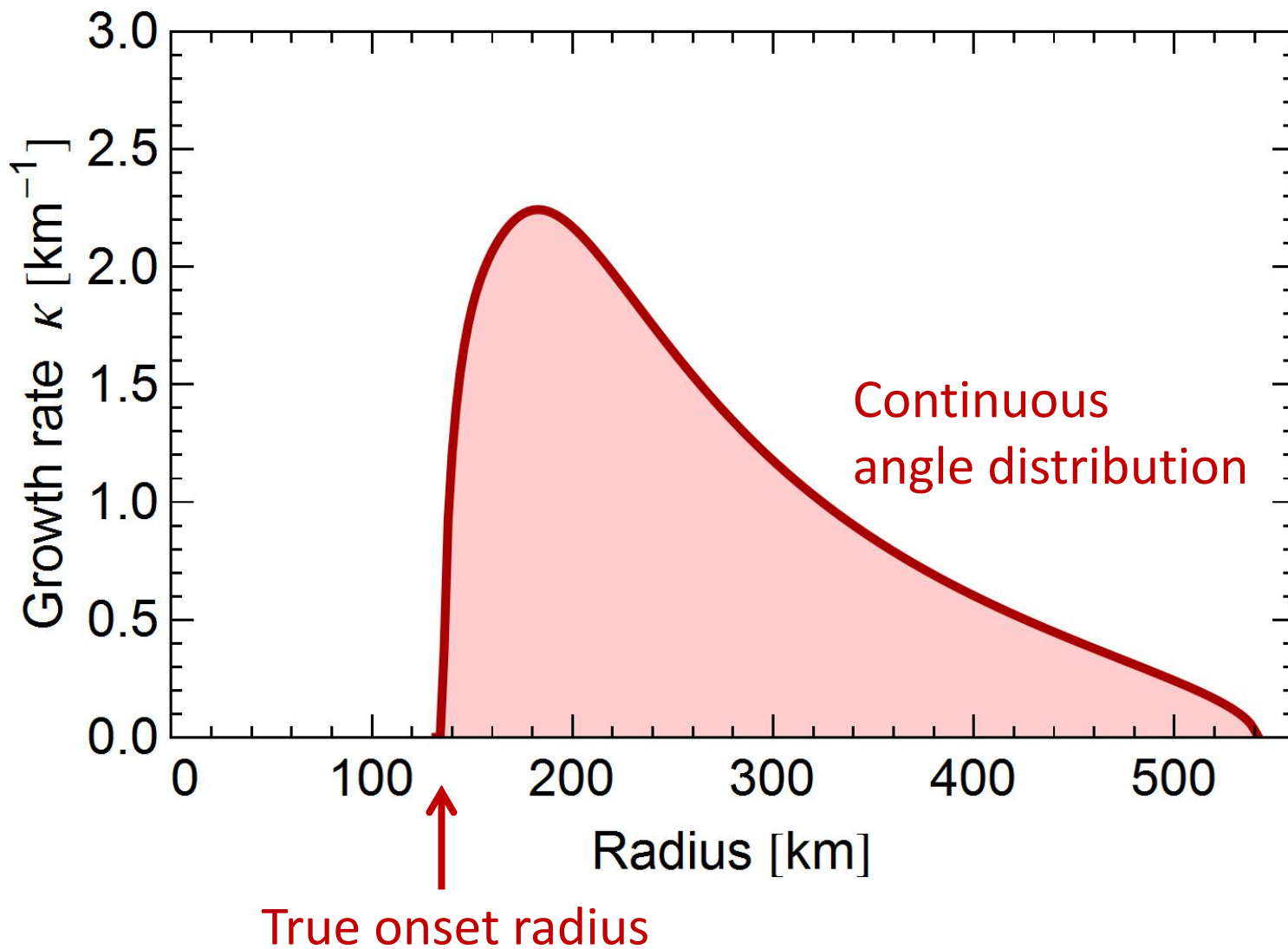


Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

**Represent neutrino field by
discrete energy and angle modes**

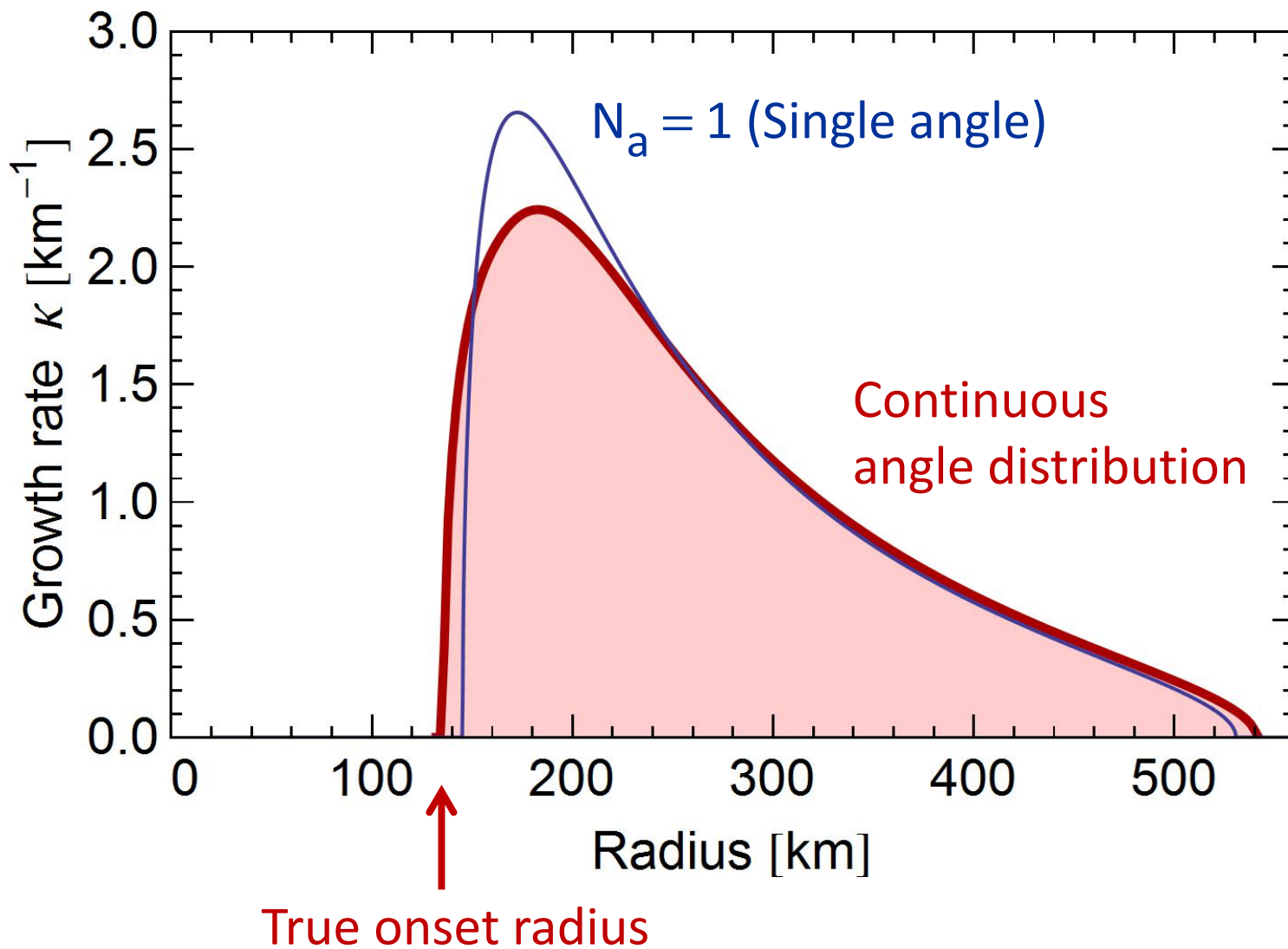
- **Number of energy modes
chosen to fit desired precision**
- **$N_a \gg 1$ of angle modes required**
 N_a too small: Unphysical solutions

New Instabilities for Discrete Angle Modes



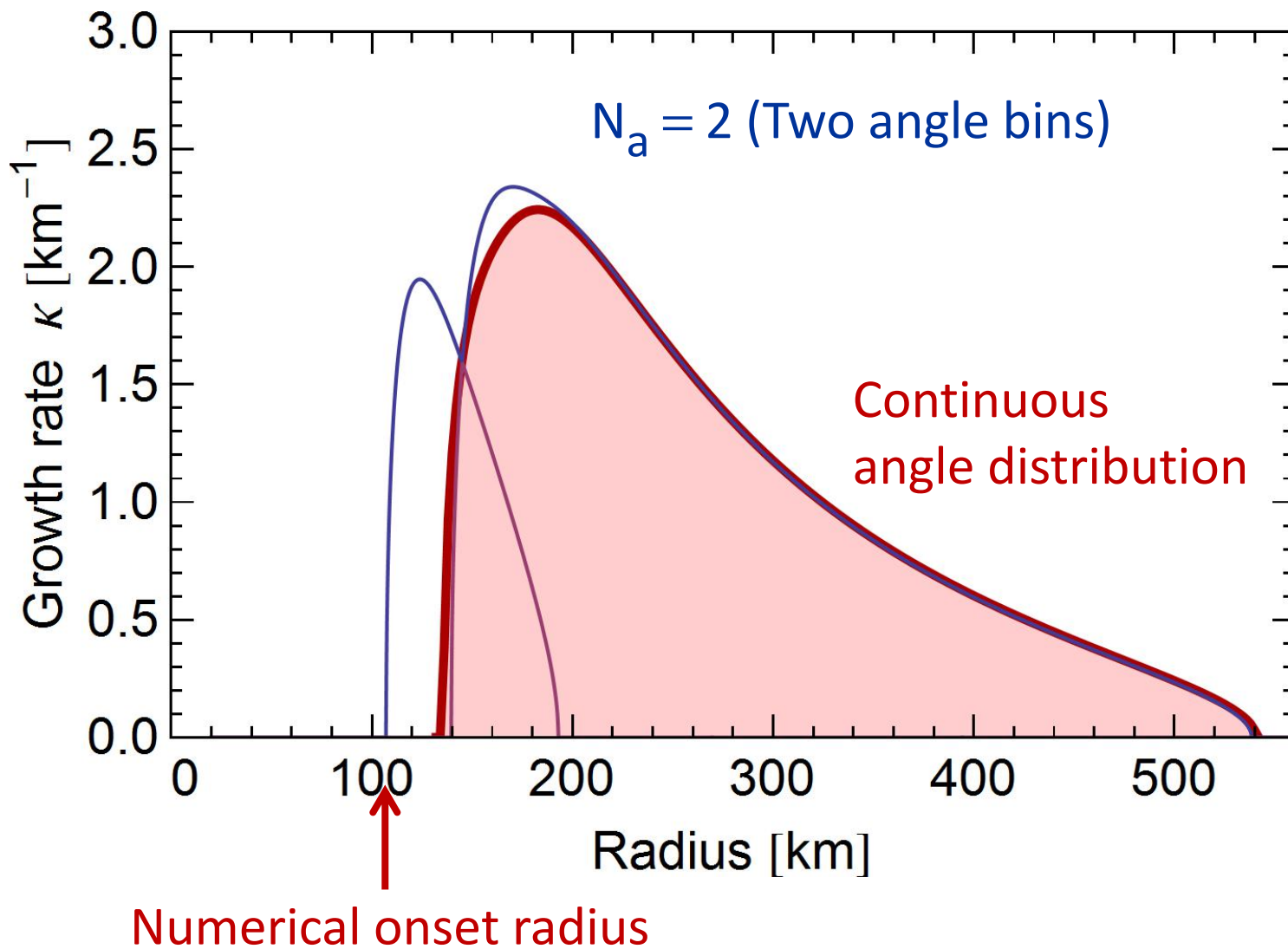
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

New Instabilities for Discrete Angle Modes



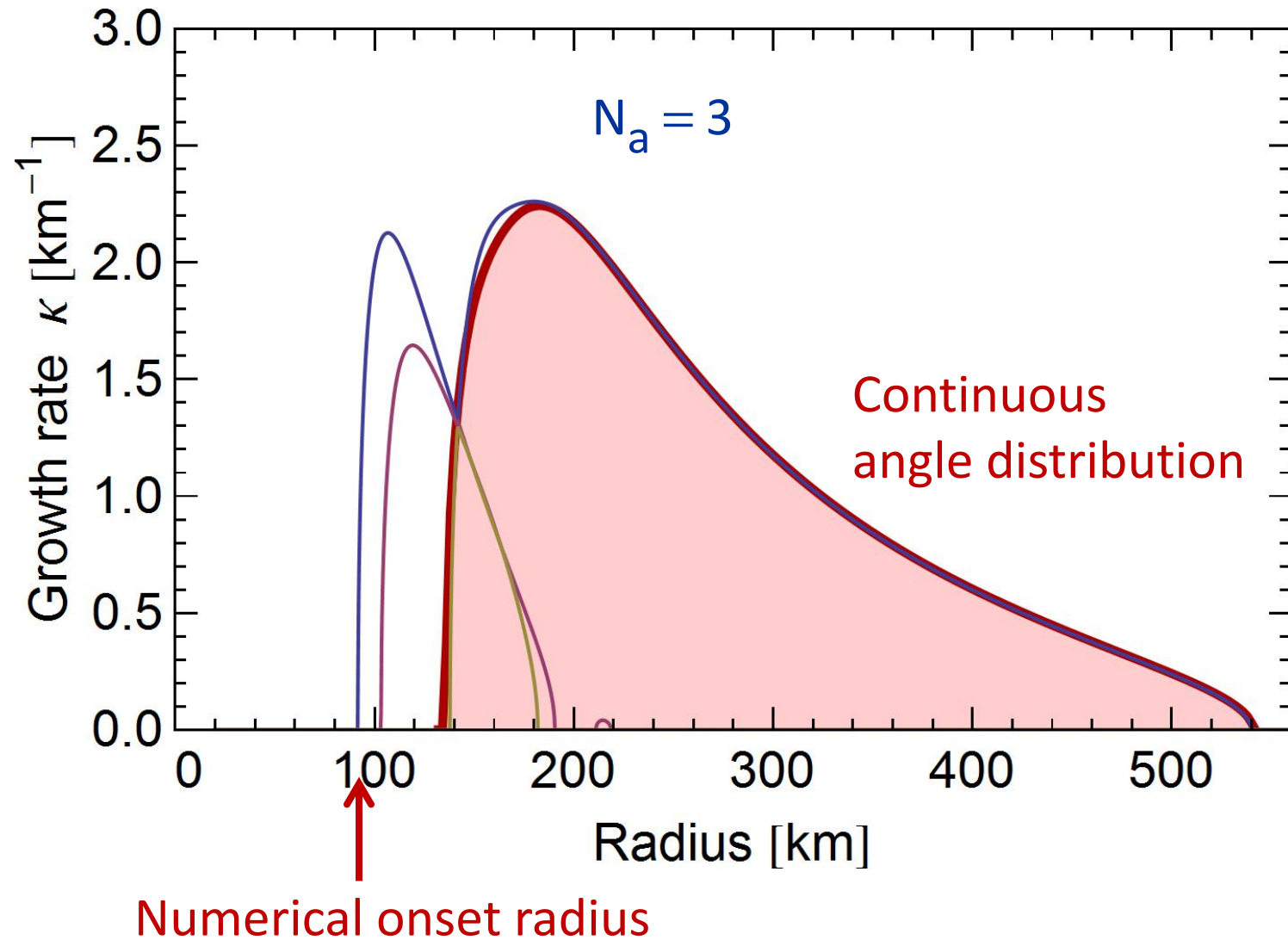
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

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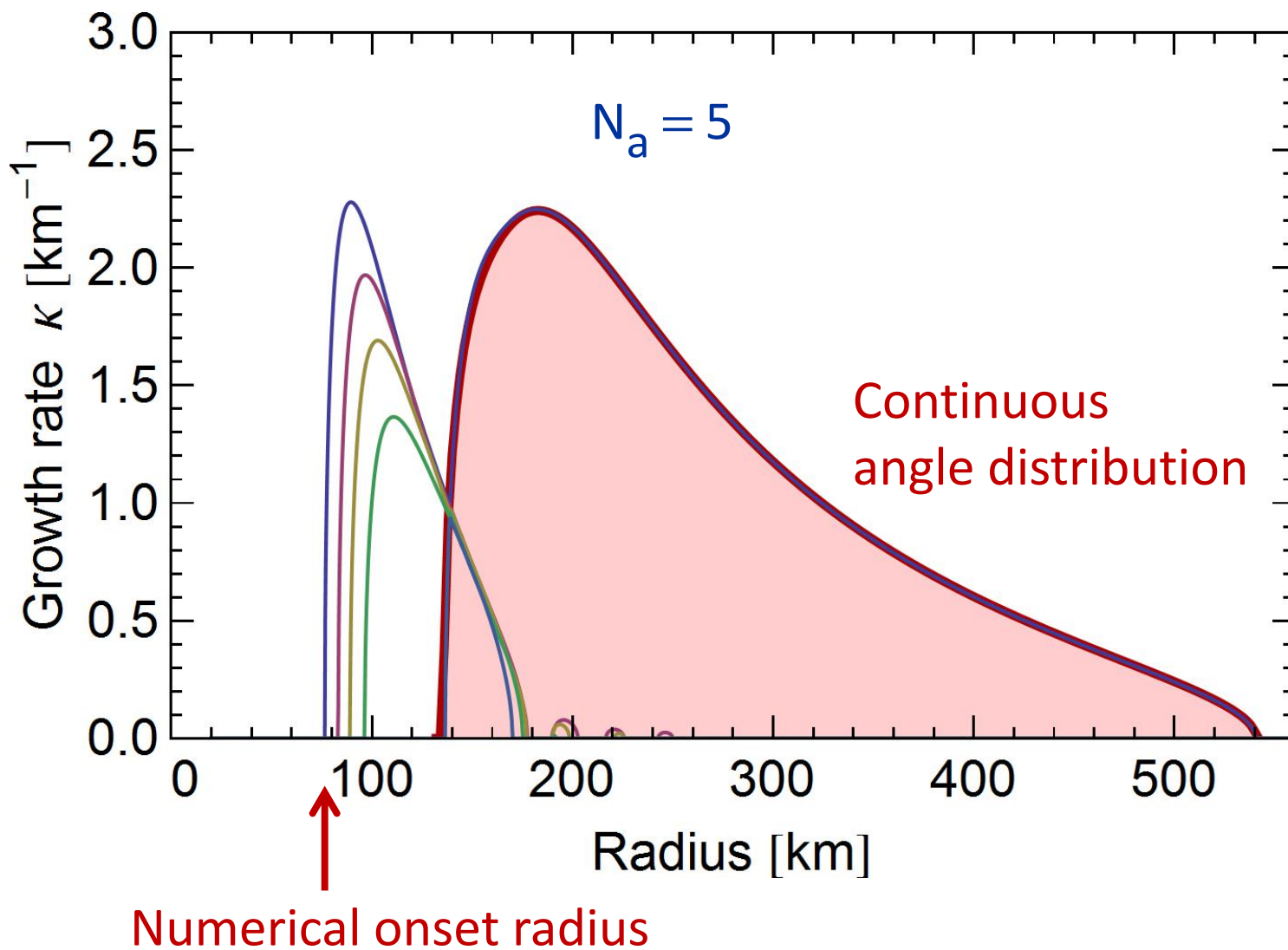
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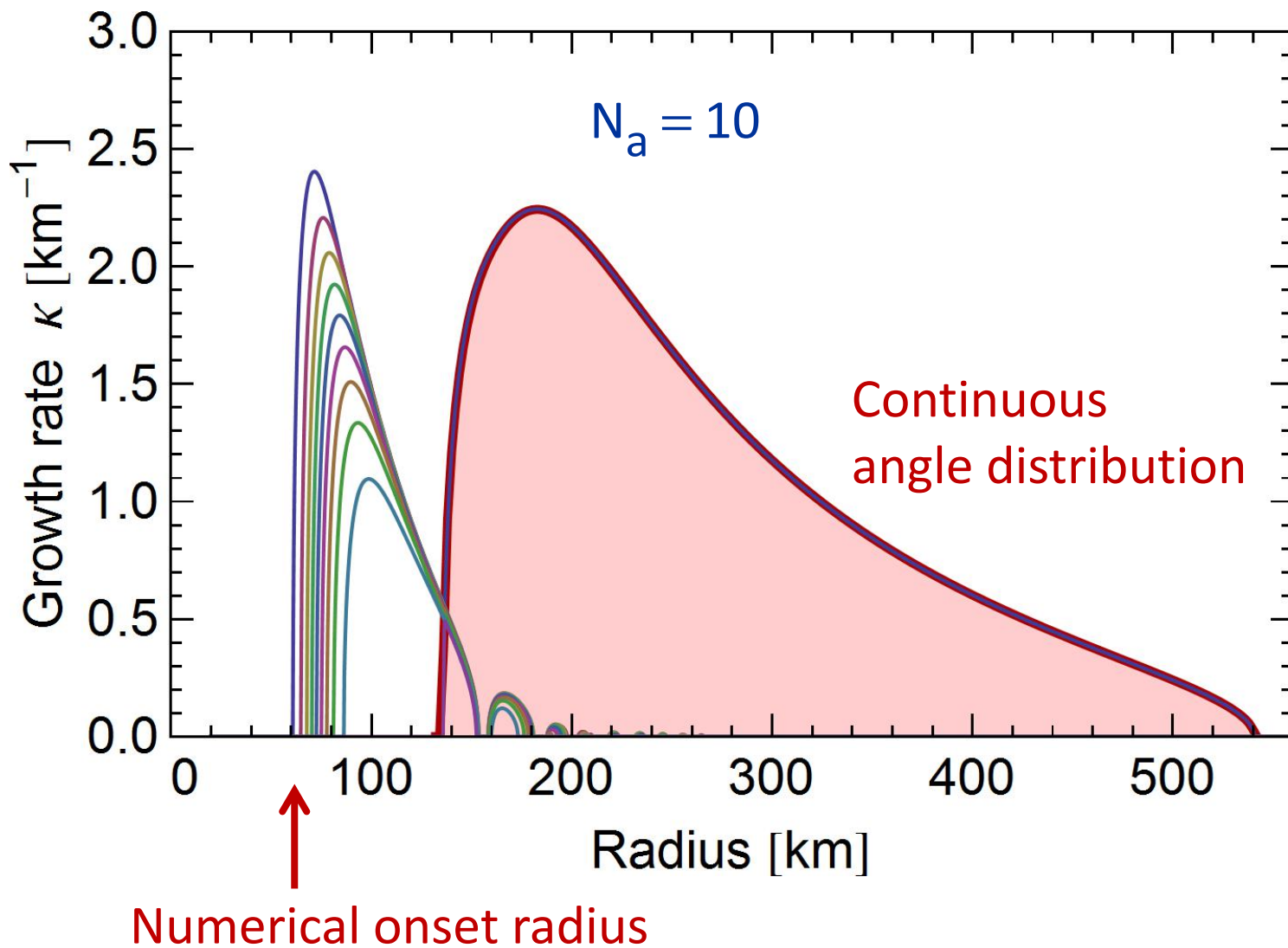
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

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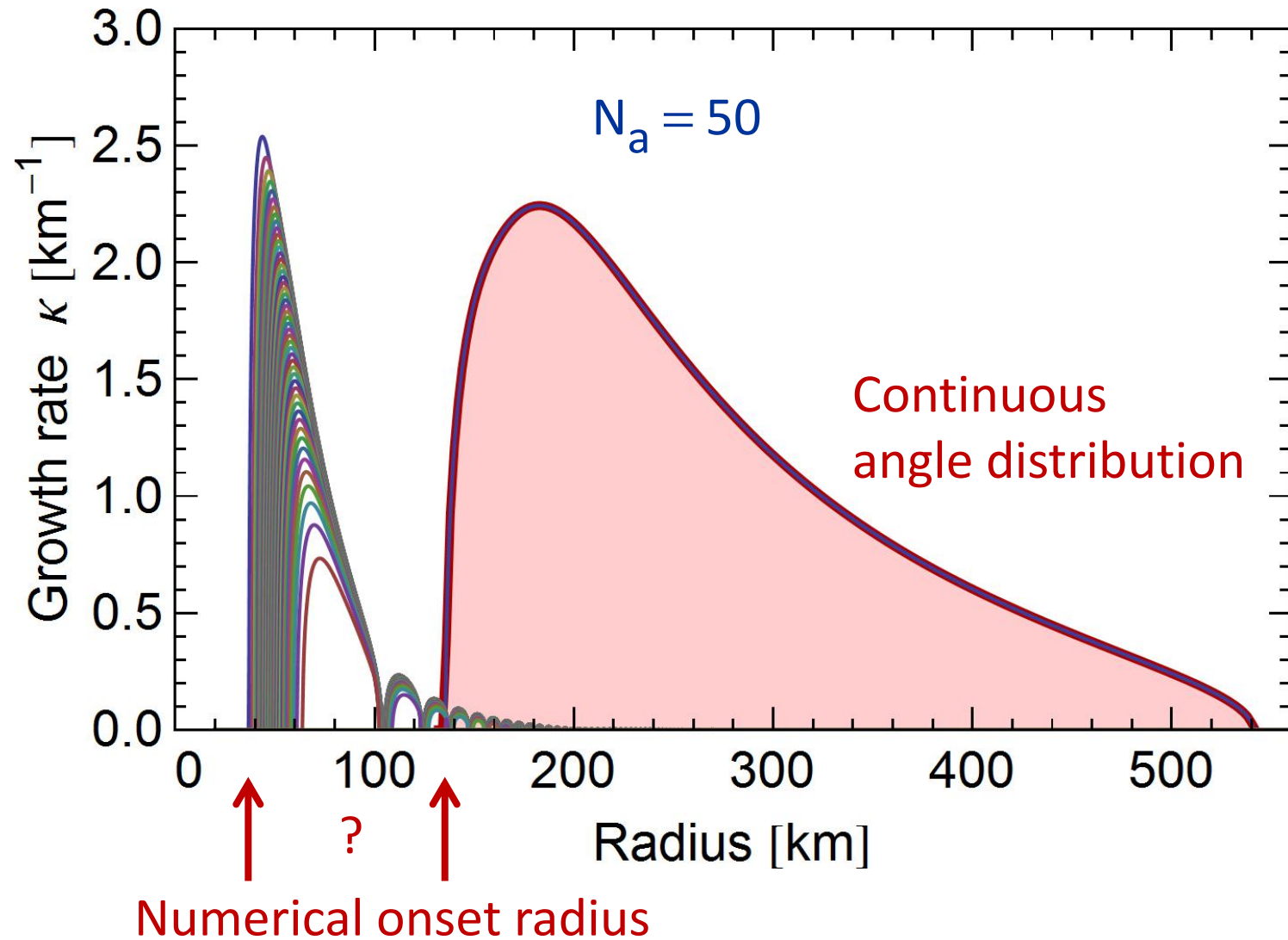
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

New Instabilities for Discrete Angle Modes



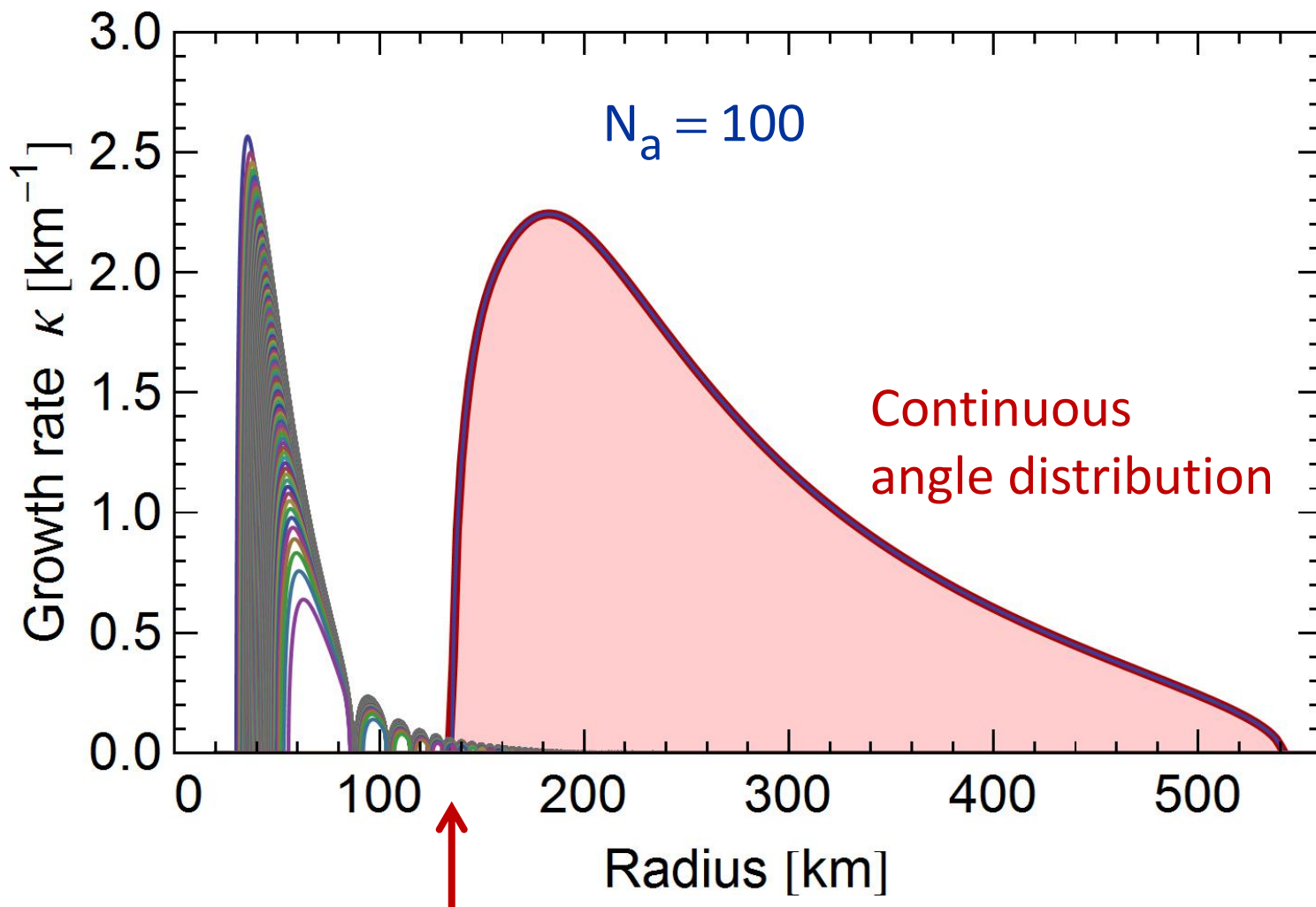
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

New Instabilities for Discrete Angle Modes



Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

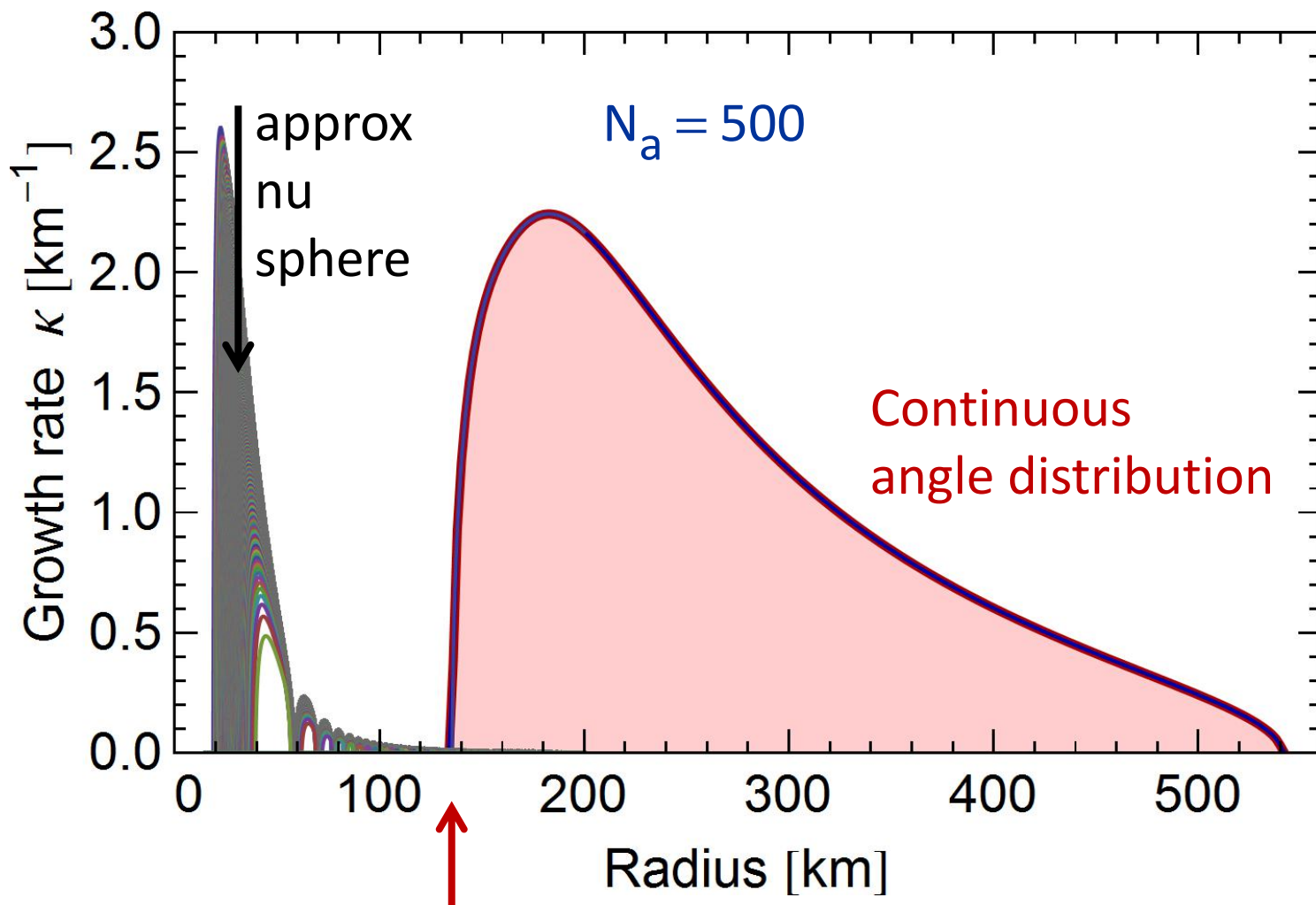
New Instabilities for Discrete Angle Modes



Numerical onset radius (start integration not too deep)

Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

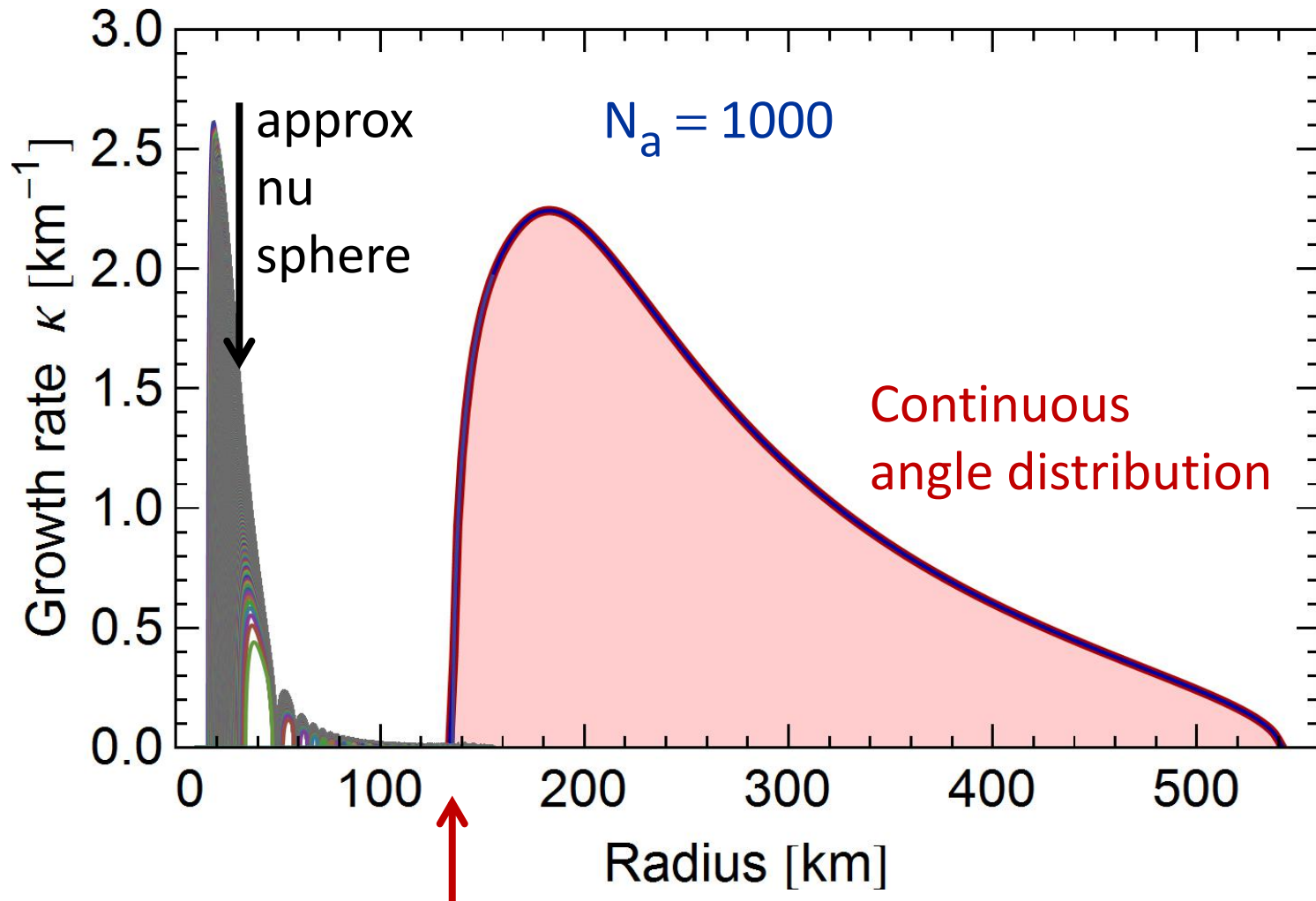
New Instabilities for Discrete Angle Modes



Numerical onset radius (start integration at nu sphere?)

Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

New Instabilities for Discrete Angle Modes



Numerical onset radius (start integration at nu sphere?)

Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Effect

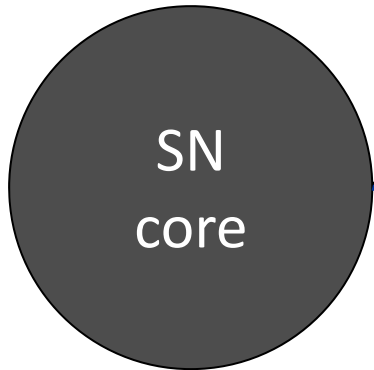
Liouville form of oscillation equation

$$\dot{\mathbf{P}}_{\omega, \mathbf{v}} + (\mathbf{v} \cdot \nabla_r) \mathbf{P}_{\omega, \mathbf{v}} = (\omega \mathbf{B} + \lambda \mathbf{L} + \mu \mathbf{P}) \times \mathbf{P}_{\omega, \mathbf{v}}$$

Drops out for stationary solutions

$$\begin{array}{cc} \uparrow & \uparrow \quad \uparrow \\ \sqrt{2} G_F N_e & \sqrt{2} G_F N_\nu \end{array}$$

$$\partial_r \mathbf{P}_{\omega, \mathbf{v}} = \frac{\omega \mathbf{B} + \lambda \mathbf{L} + \mu \mathbf{P}}{v_r} \times \mathbf{P}_{\omega, \mathbf{v}}$$

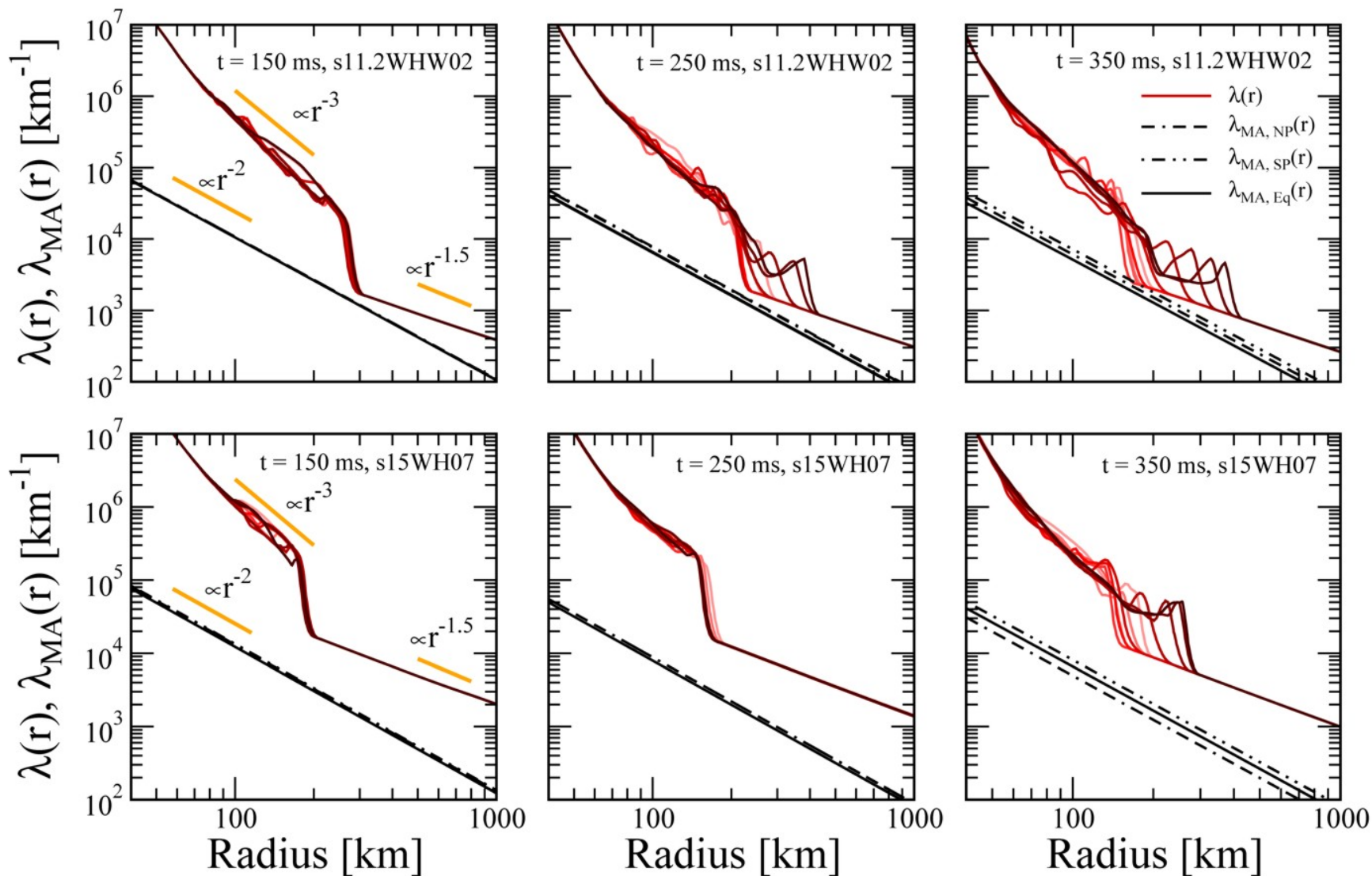


Longer path to
same radial distance:
Larger matter effect

Self-induced conversion suppressed for $N_e \gtrsim N_\nu$

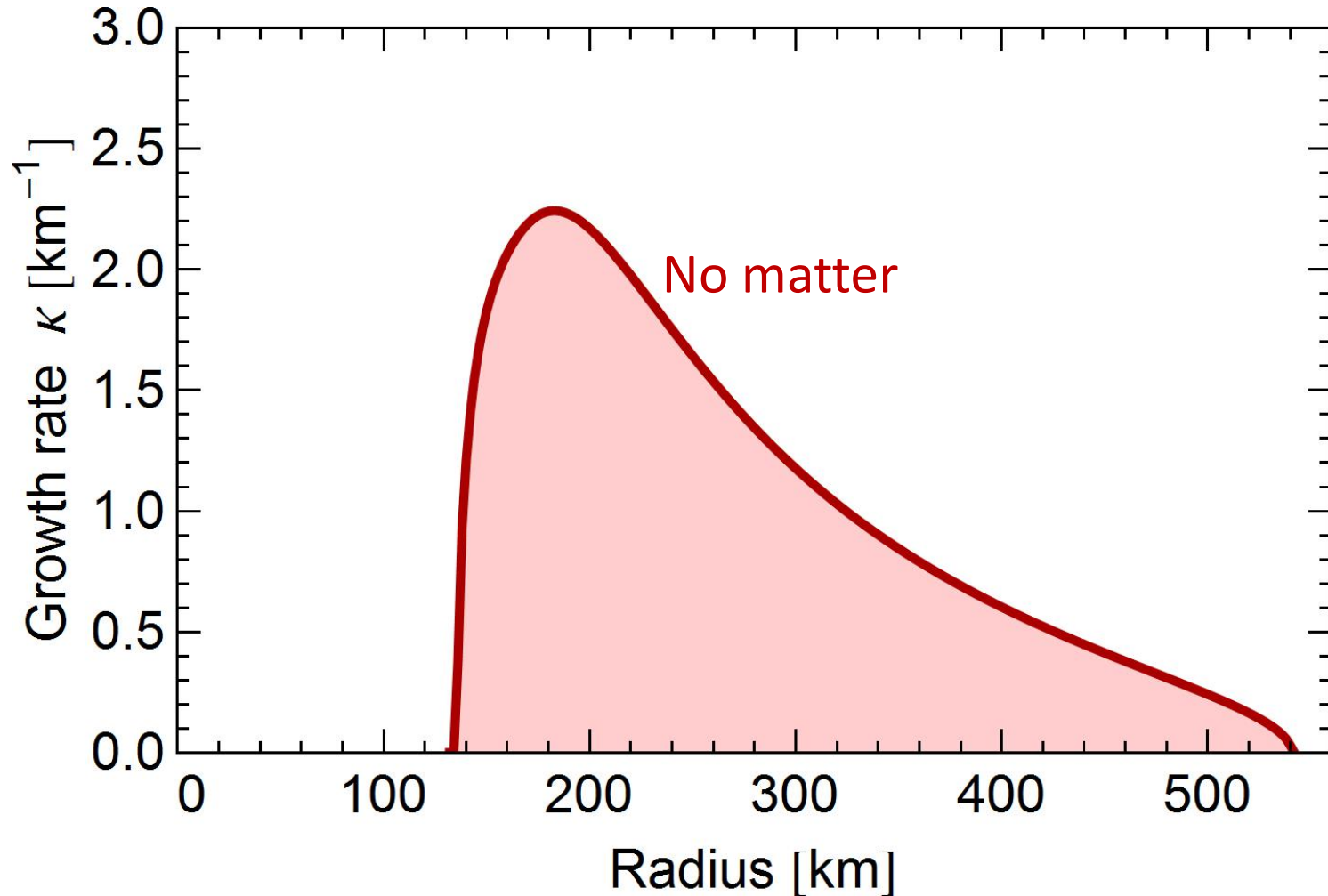
Esteban-Pretel, Mirizzi, Pastor, Tomàs, Raffelt, Serpico & Sigl, arXiv:0807.0659

Accretion-Phase Matter Profiles



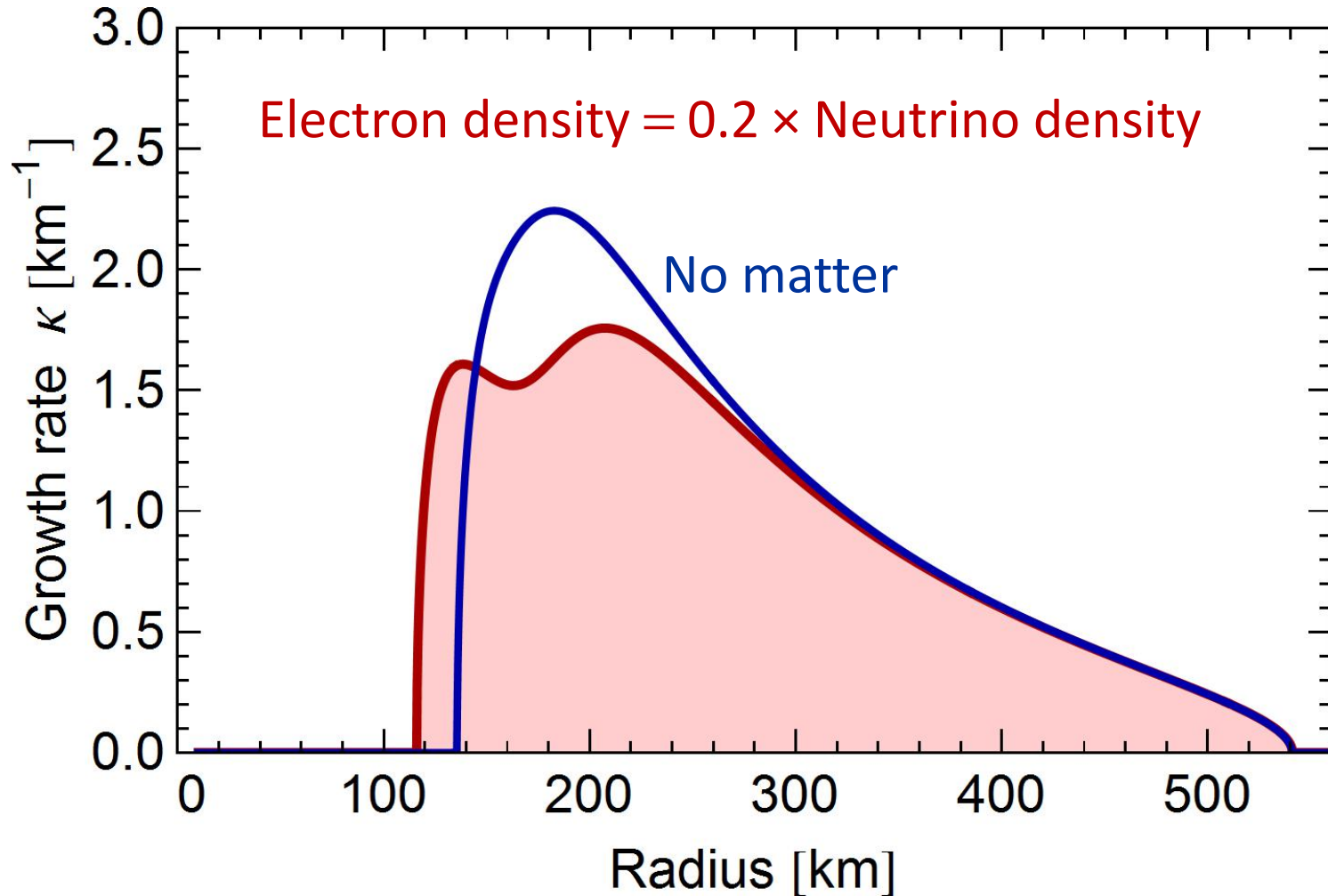
Dasgupta, O'Connor & Ott, arXiv:1106.1167

Multi-Angle Matter Suppression



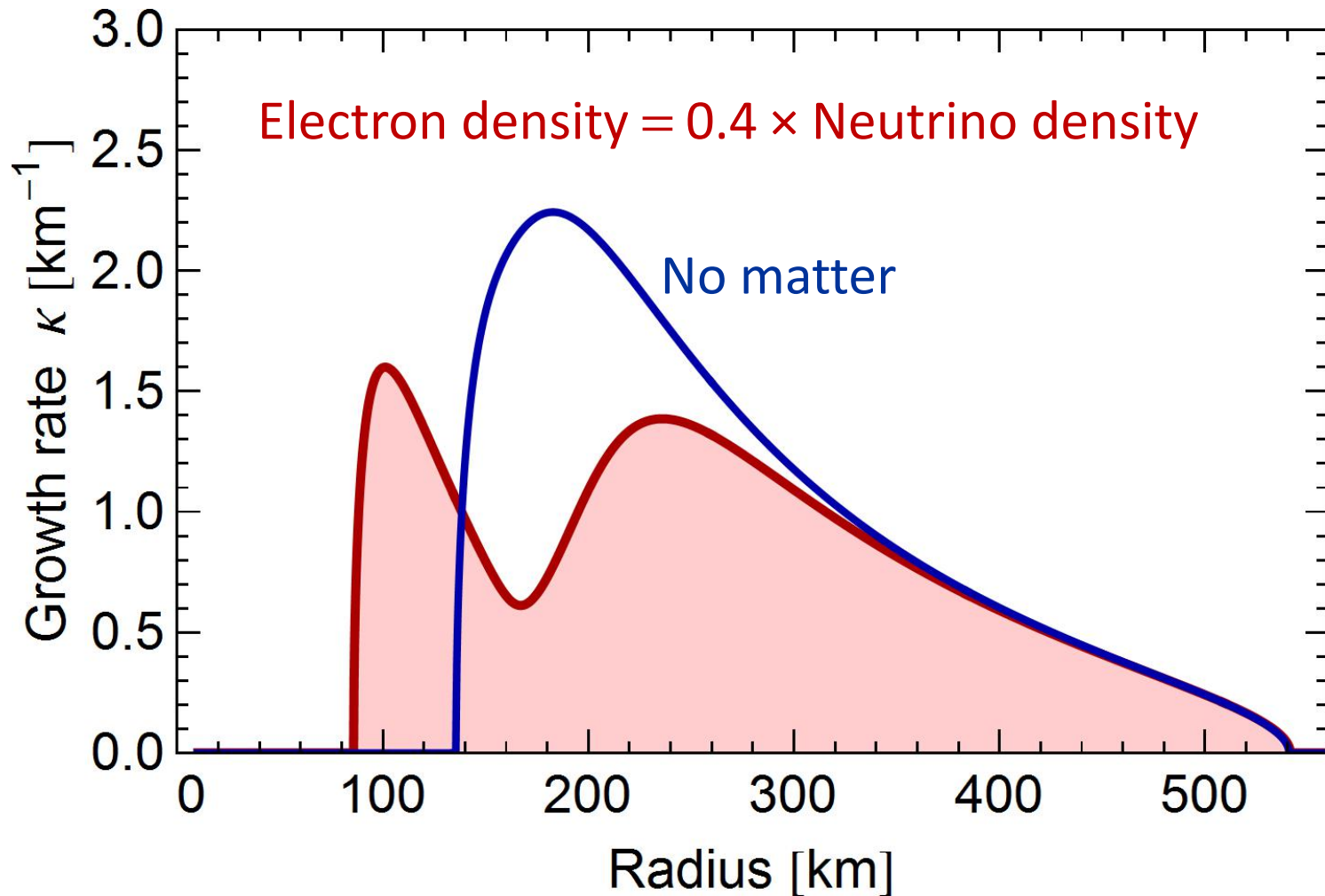
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression



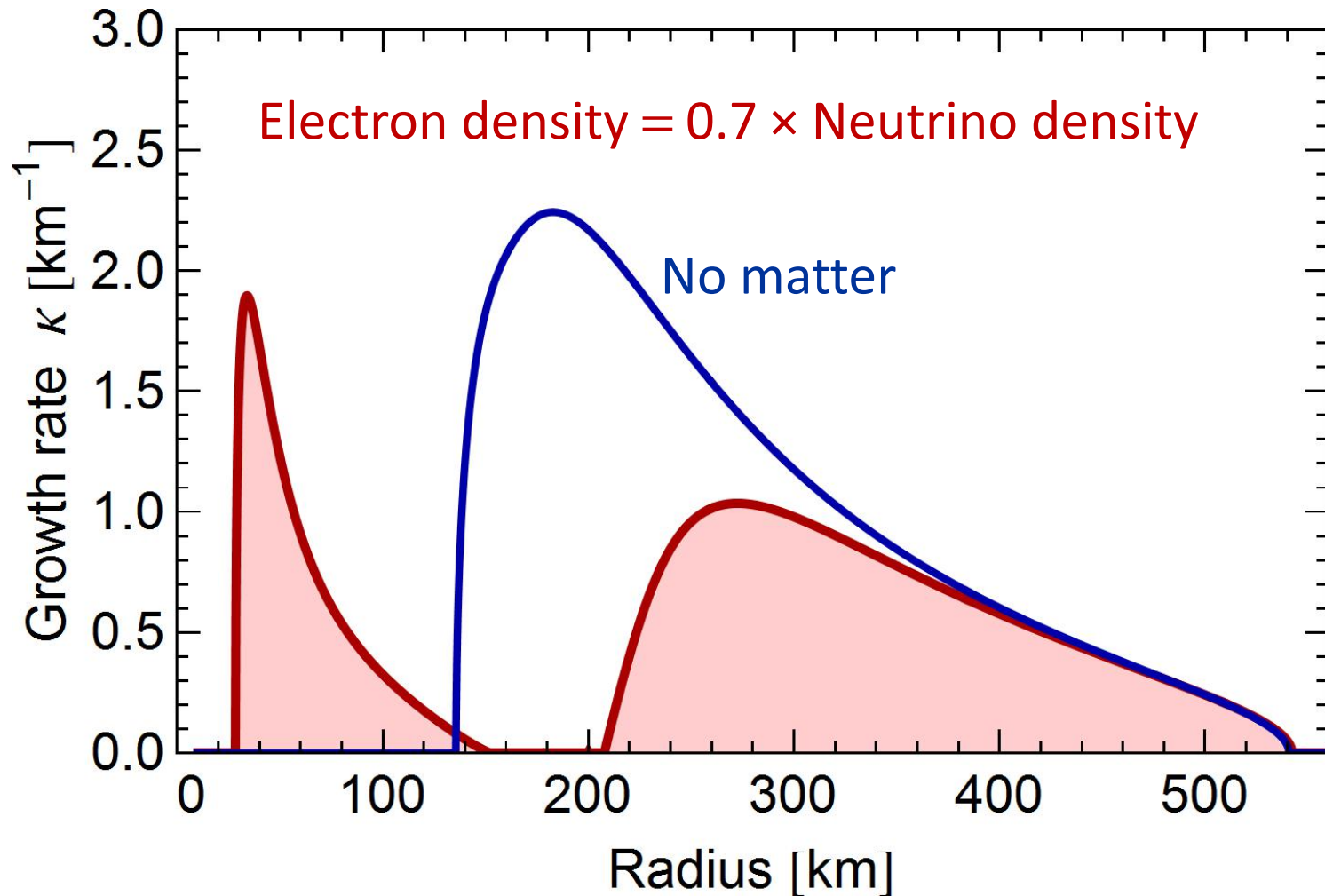
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression



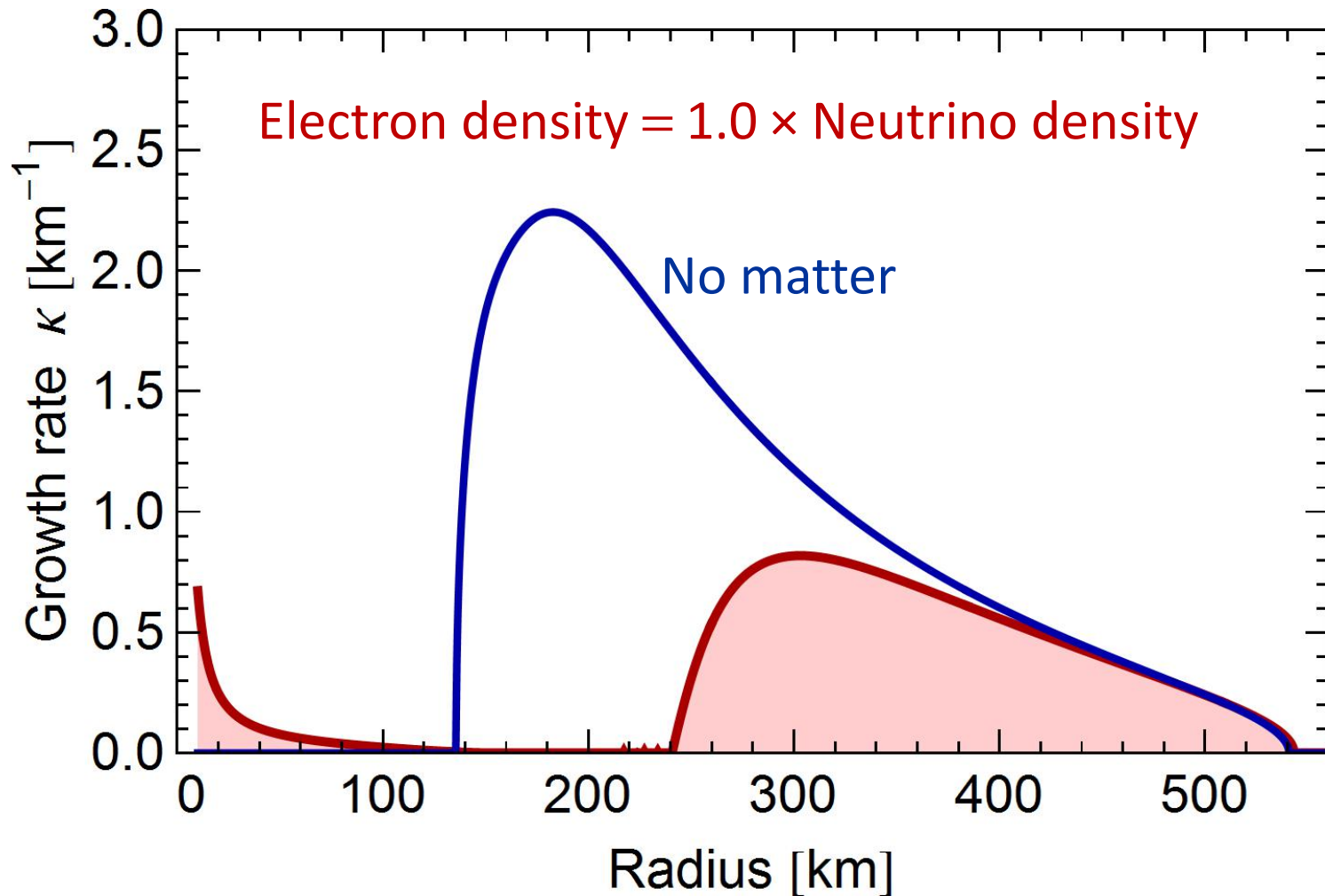
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression



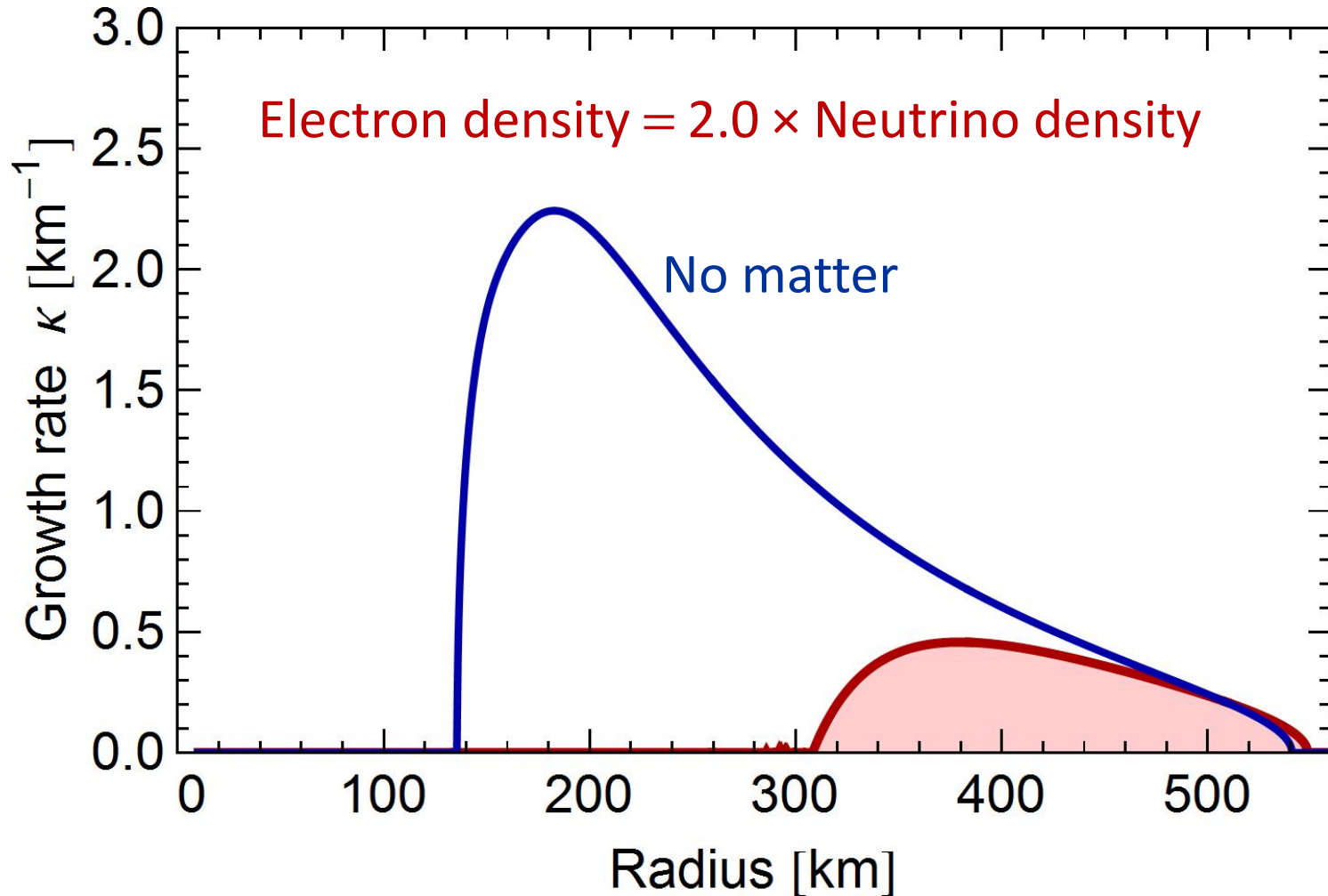
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression



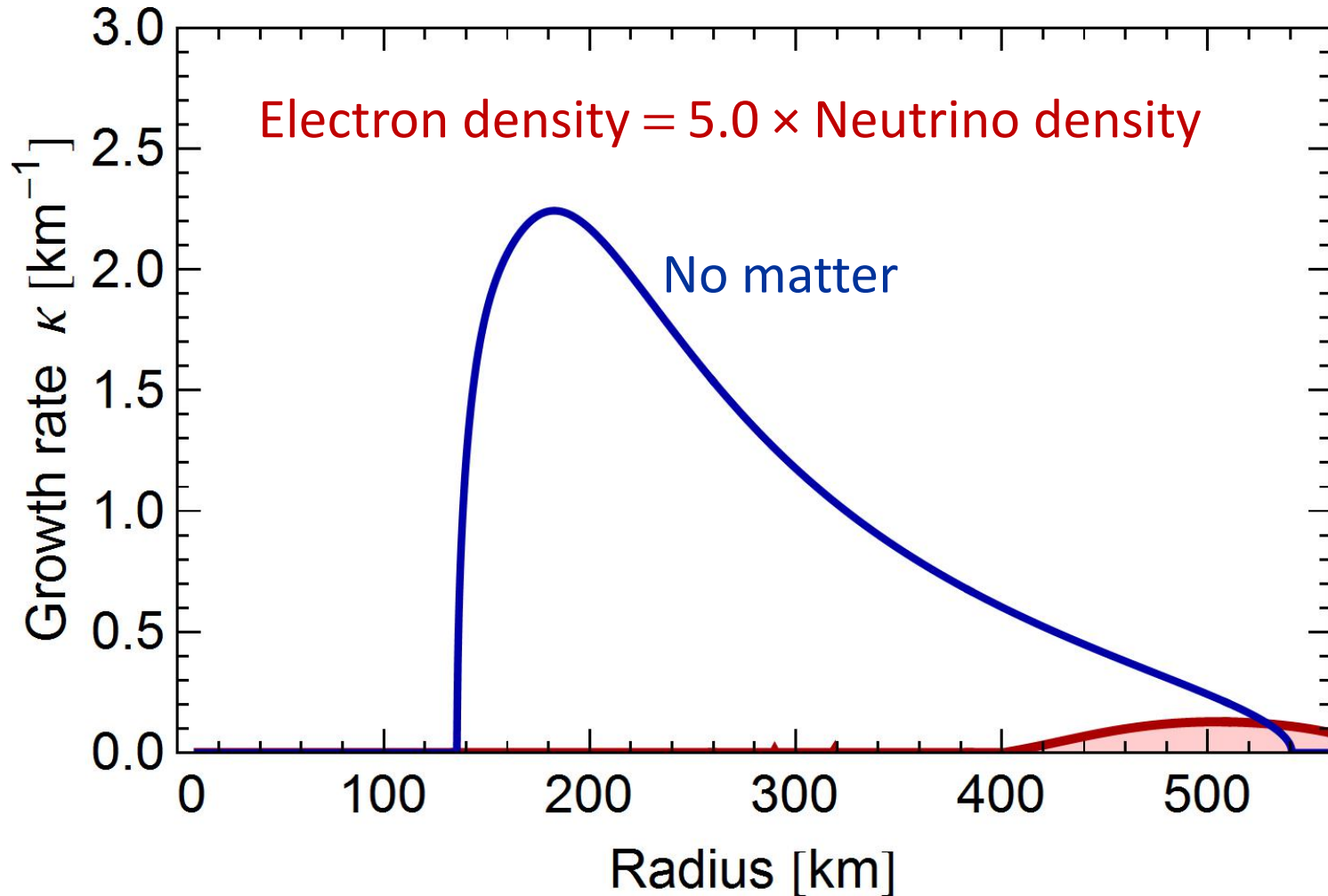
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression



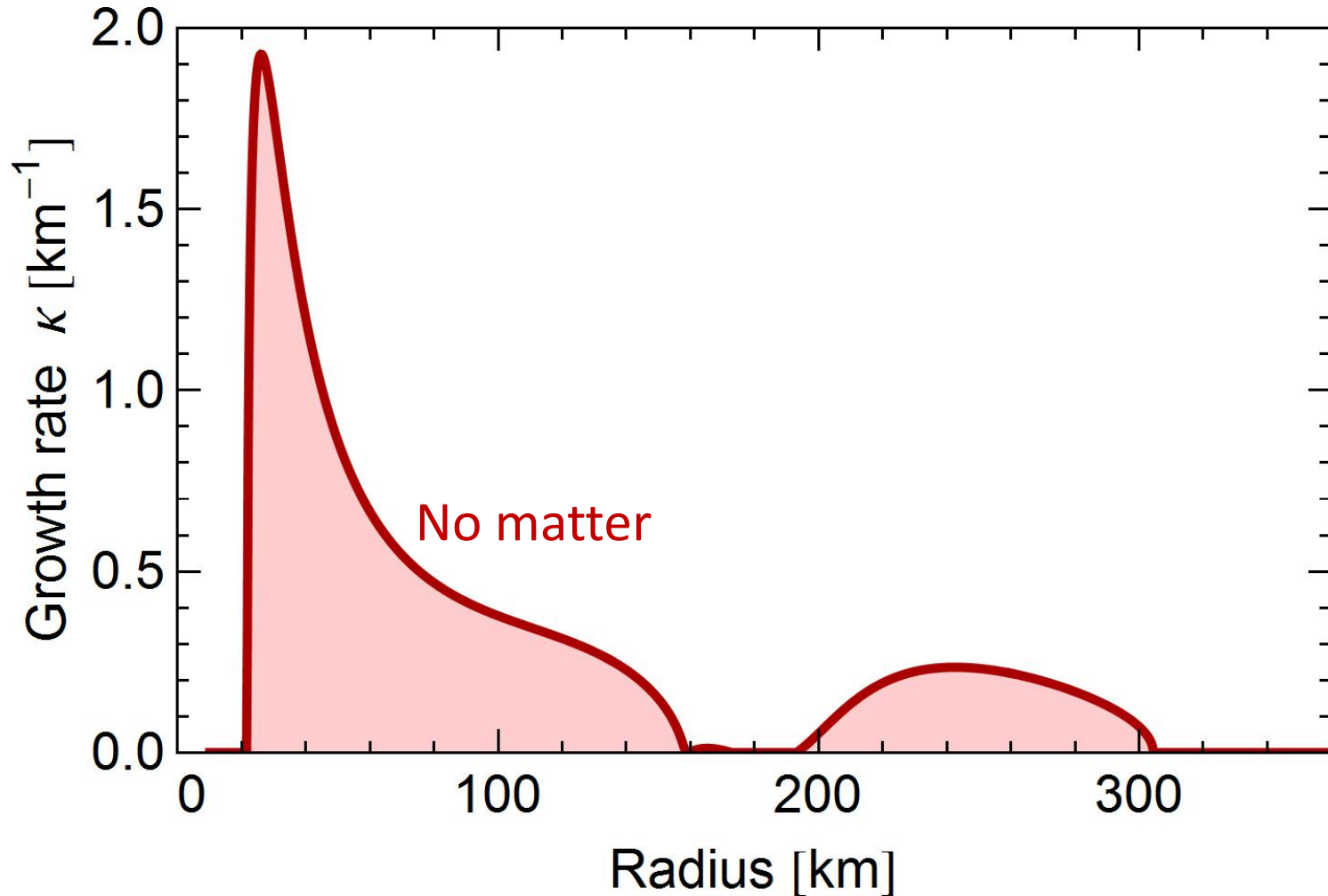
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression



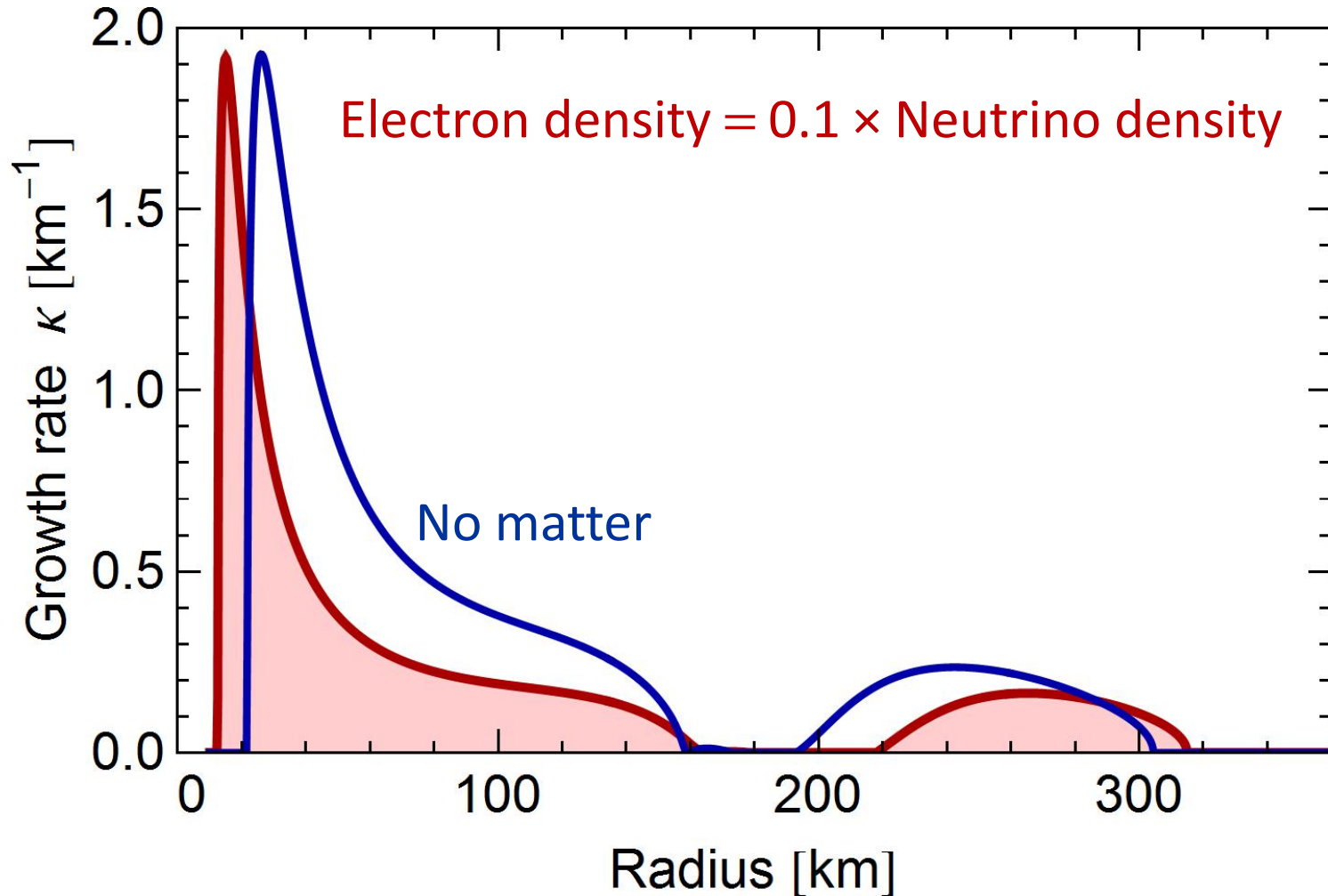
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression (NH)



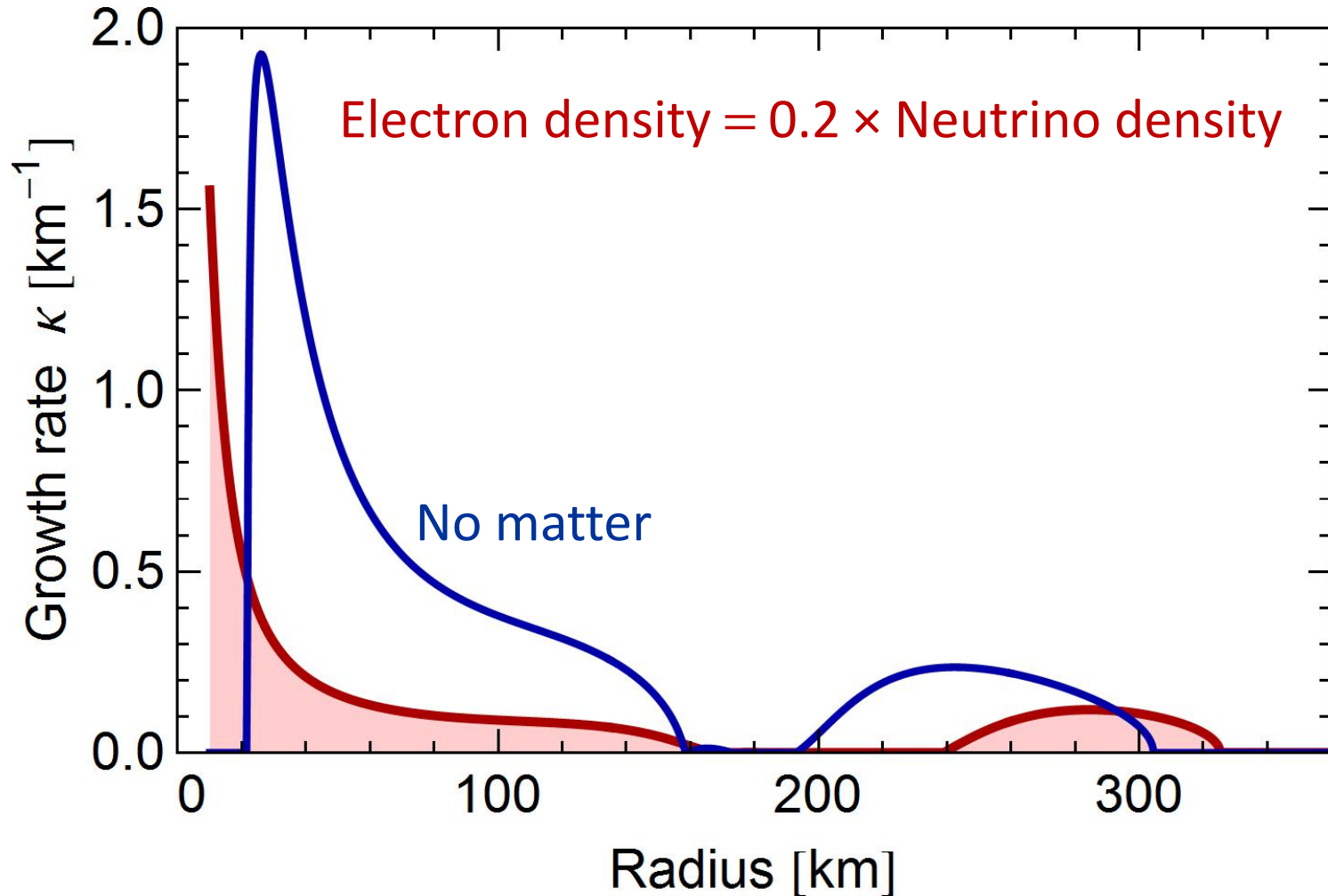
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression (NH)



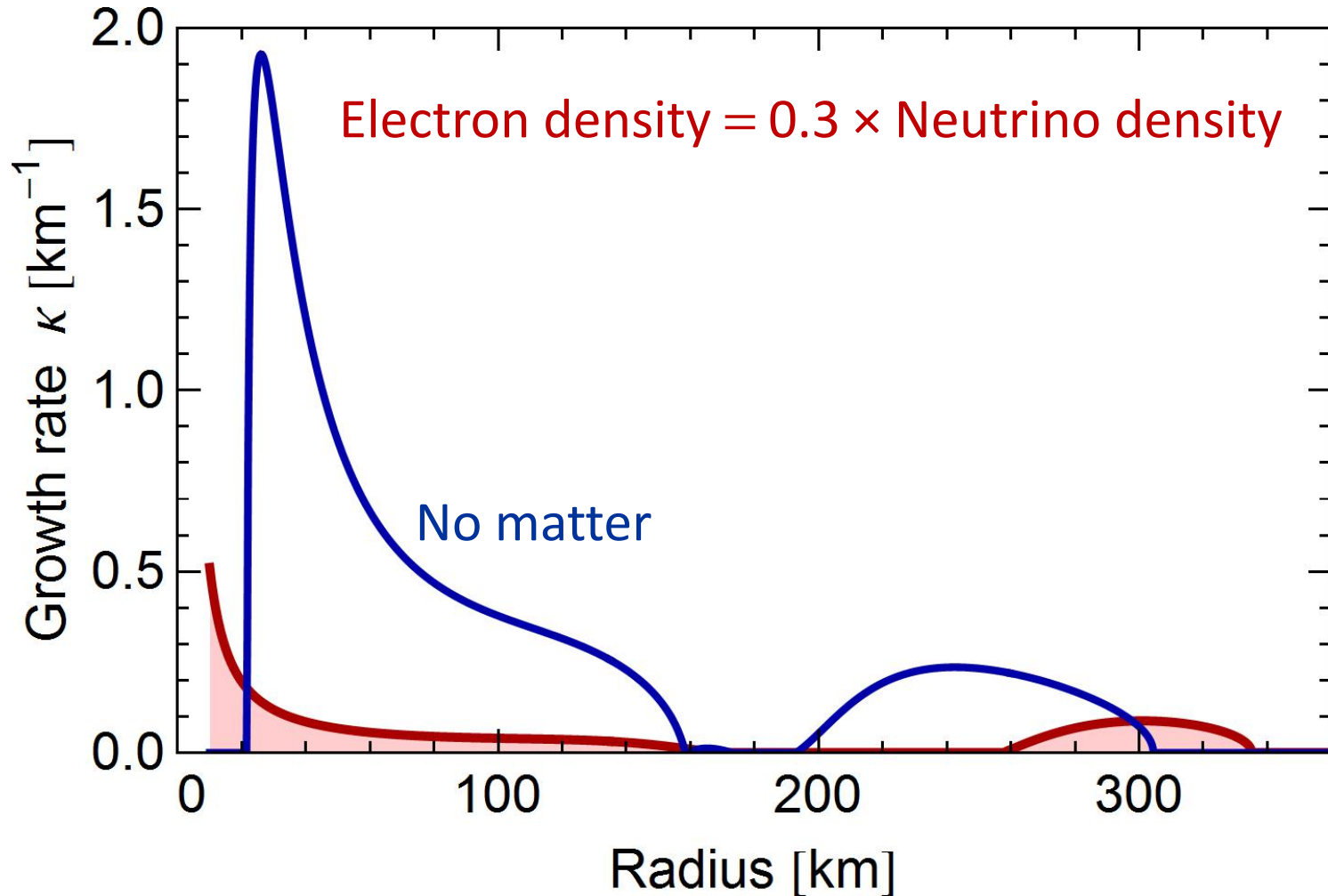
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression (NH)



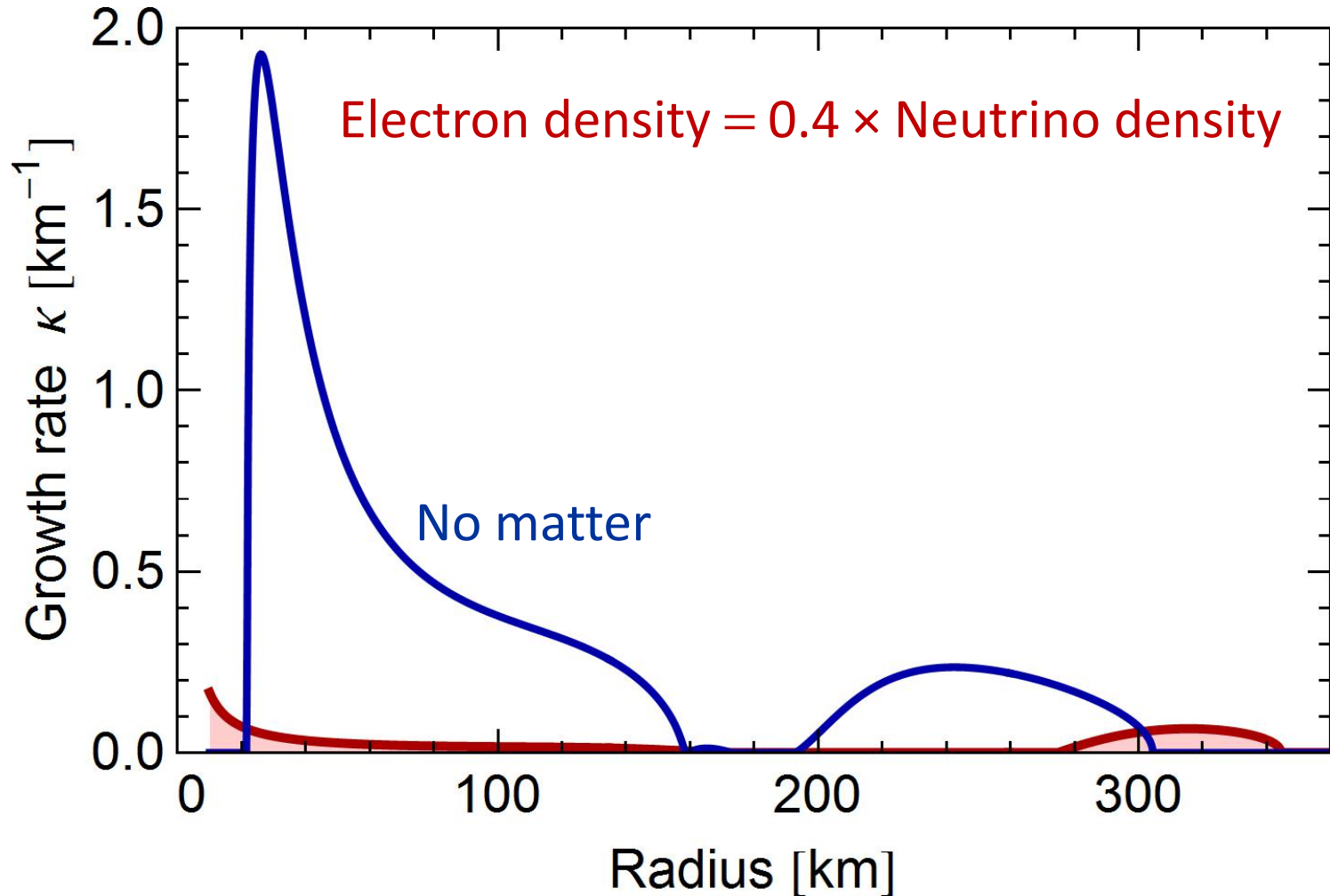
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression (NH)



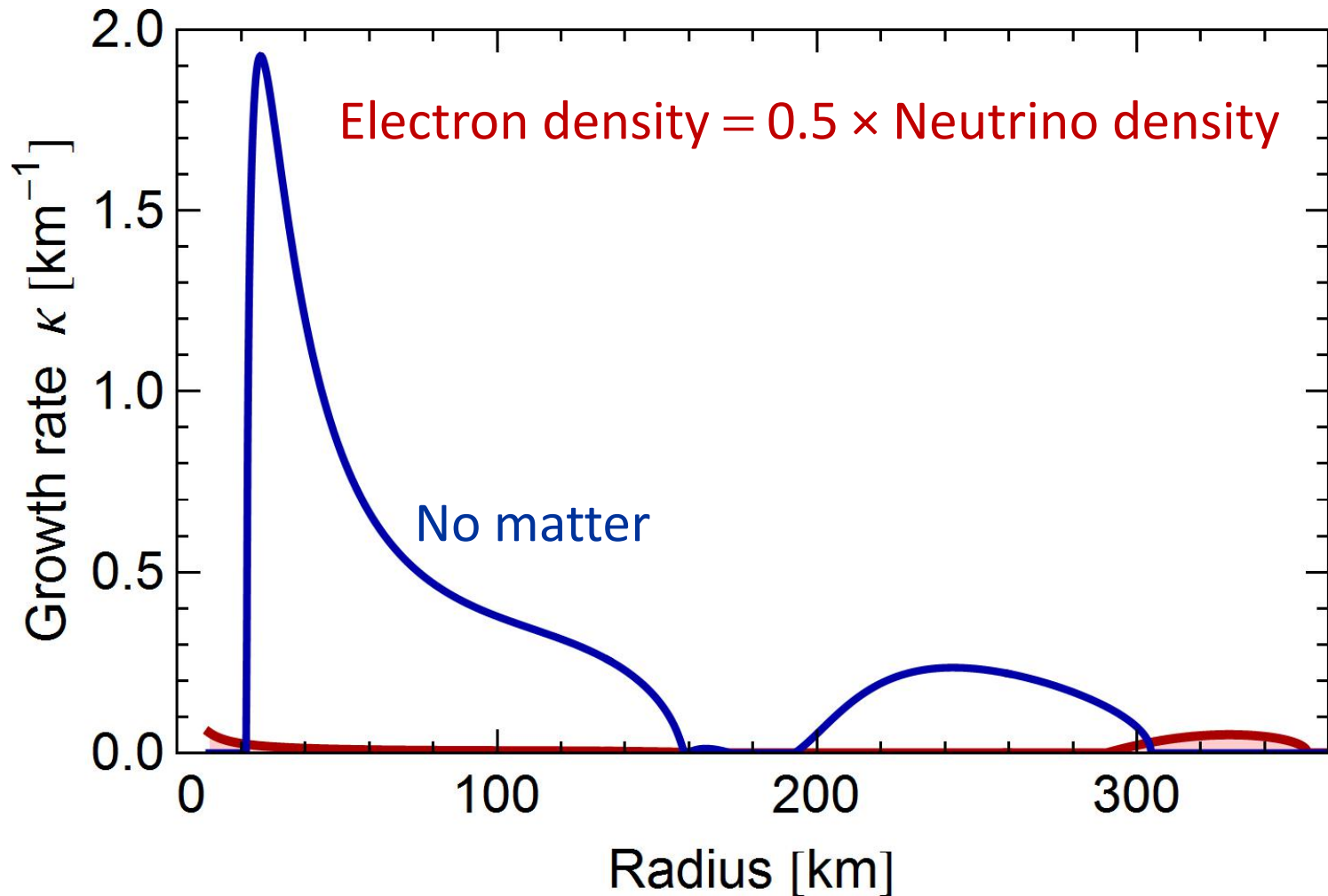
Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

Multi-Angle Matter Suppression (NH)



Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

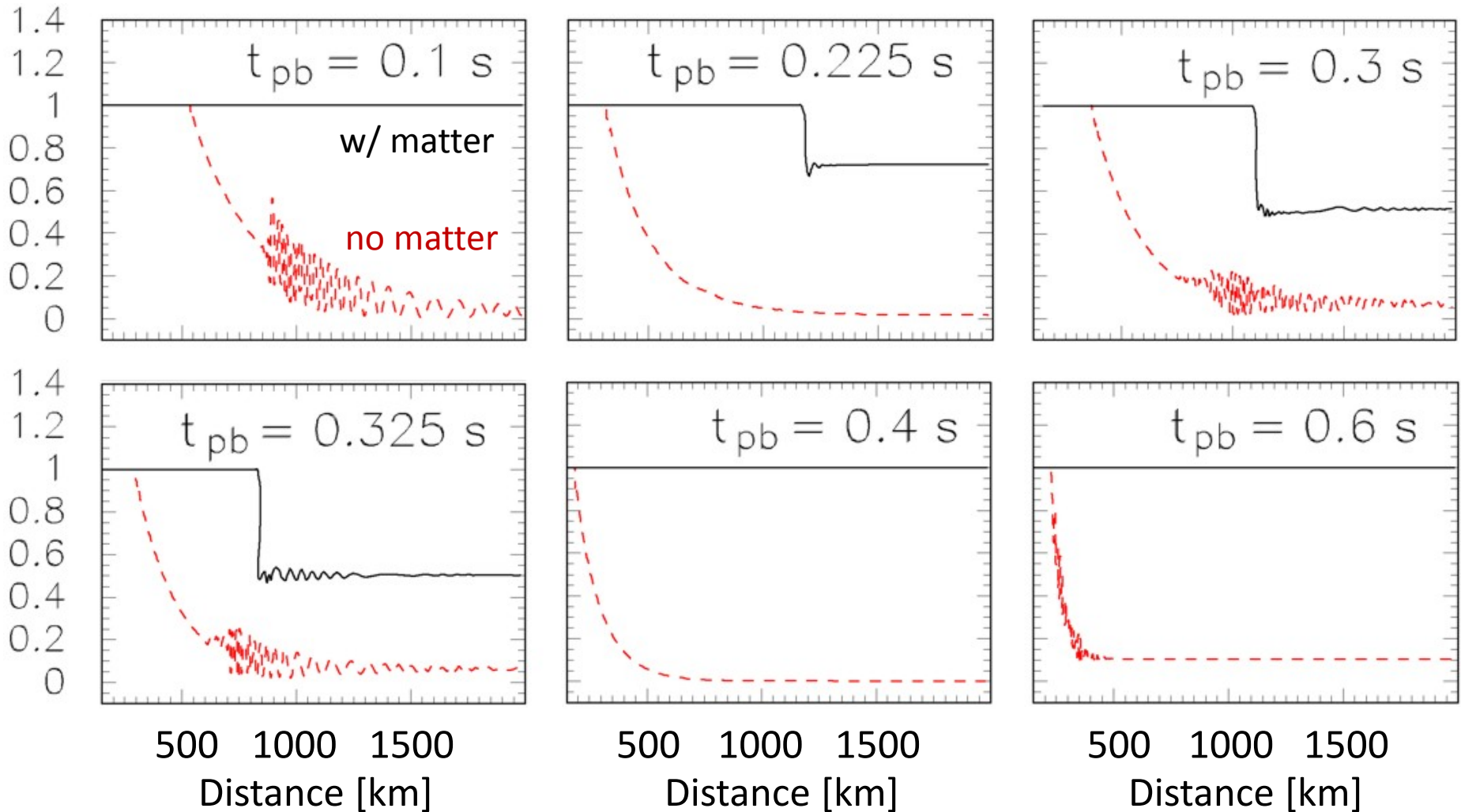
Multi-Angle Matter Suppression (NH)



Sarikas, Seixas, Tamborra & Raffelt, work in progress (2012)

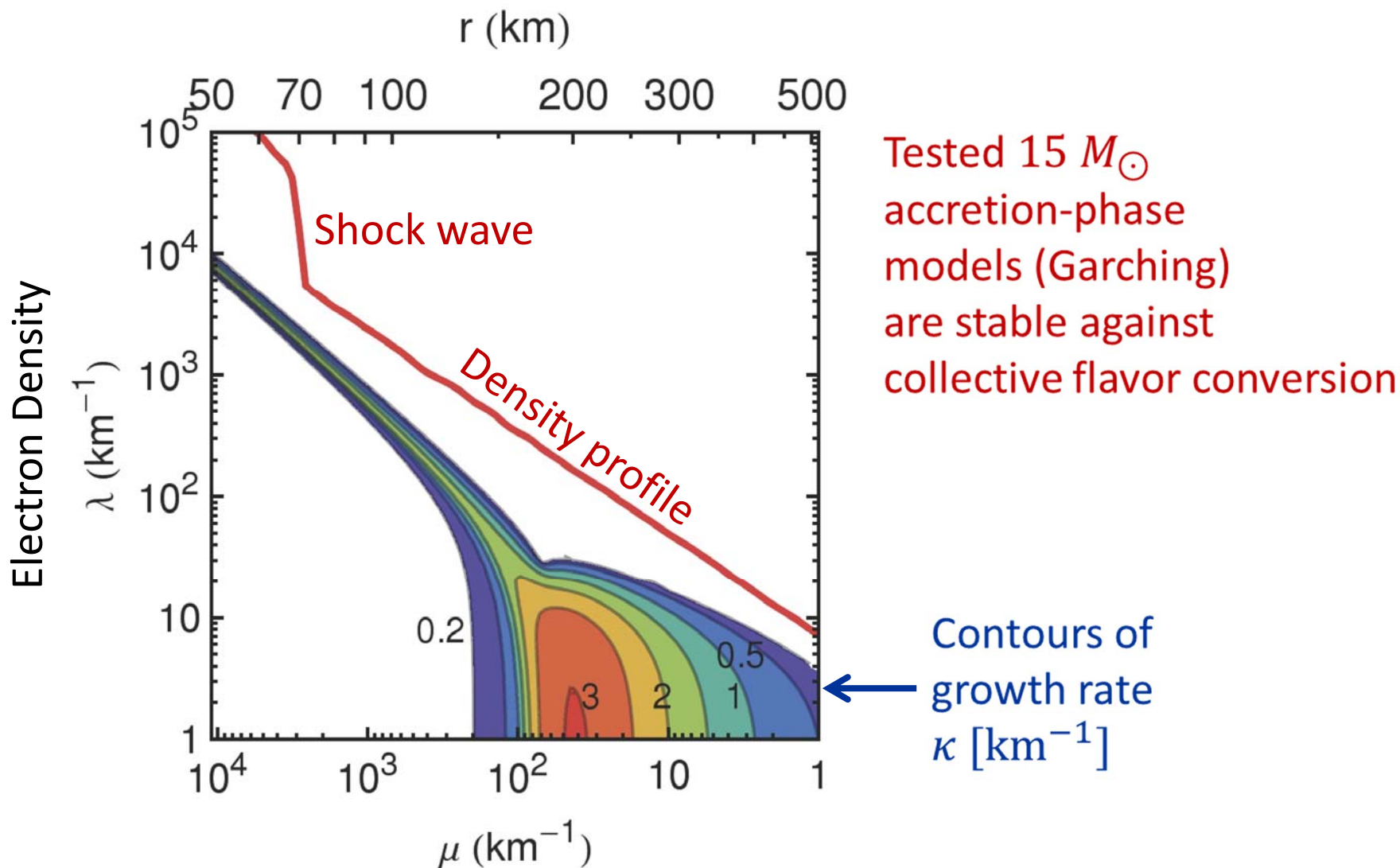
Multi-Angle Matter Effect (Basel Model $10.8 M_{\text{sun}}$)

Schematic single-energy, multi-angle simulations with realistic density profile



Chakraborty, Fischer, Mirizzi, Saviano & Tomàs, arXiv:1105.1130

Multi-Angle Multi-Energy Stability Analysis



Sarikas, Raffelt, Hüdepohl & Janka, arXiv:1109.3601

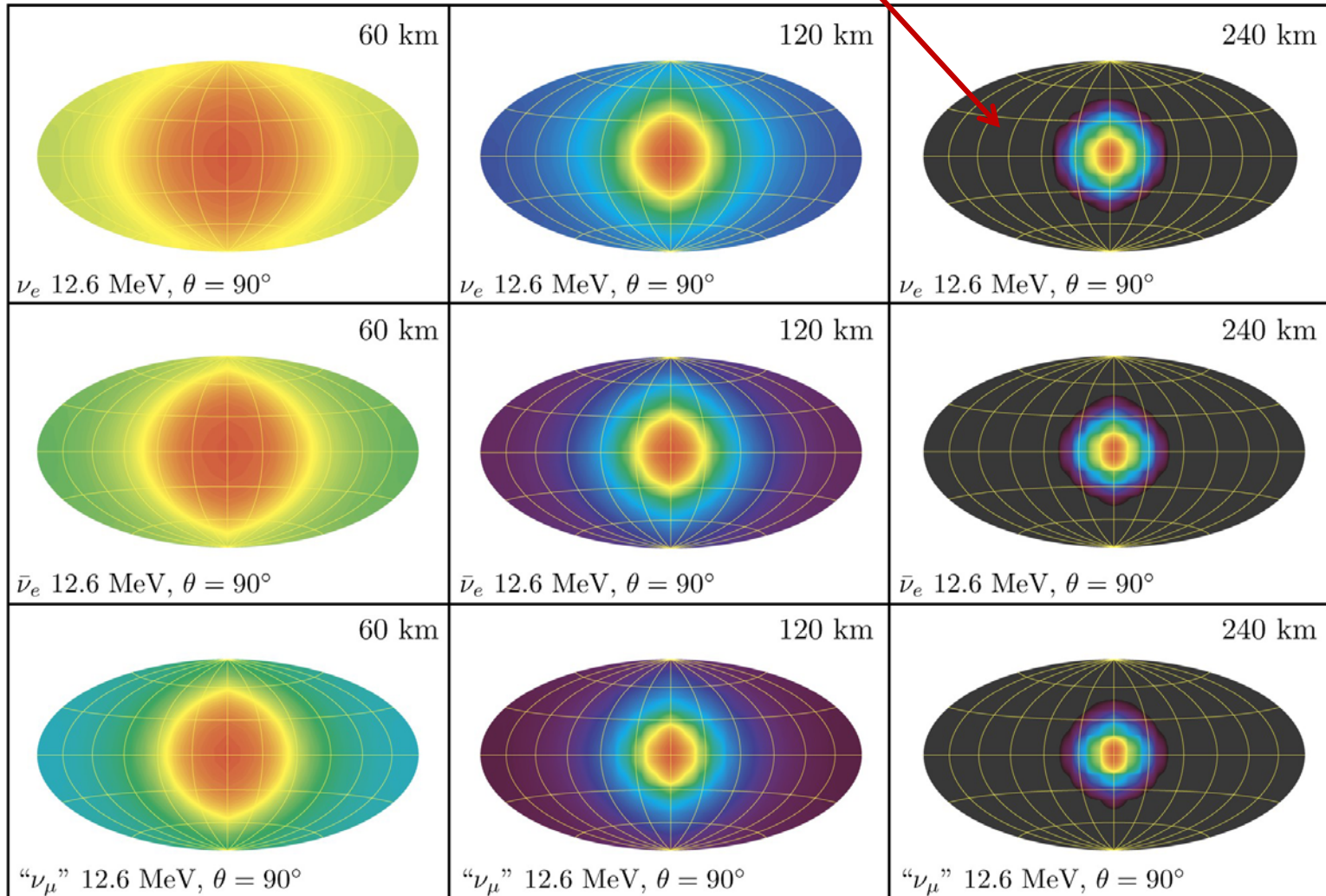
Required Neutrino Information

To investigate neutrino oscillations in SNe need

- Profile of electron density
- Flavor-dependent neutrino flux spectra
- Flavor-dependent neutrino angular distribution

Neutrino Radiation Field

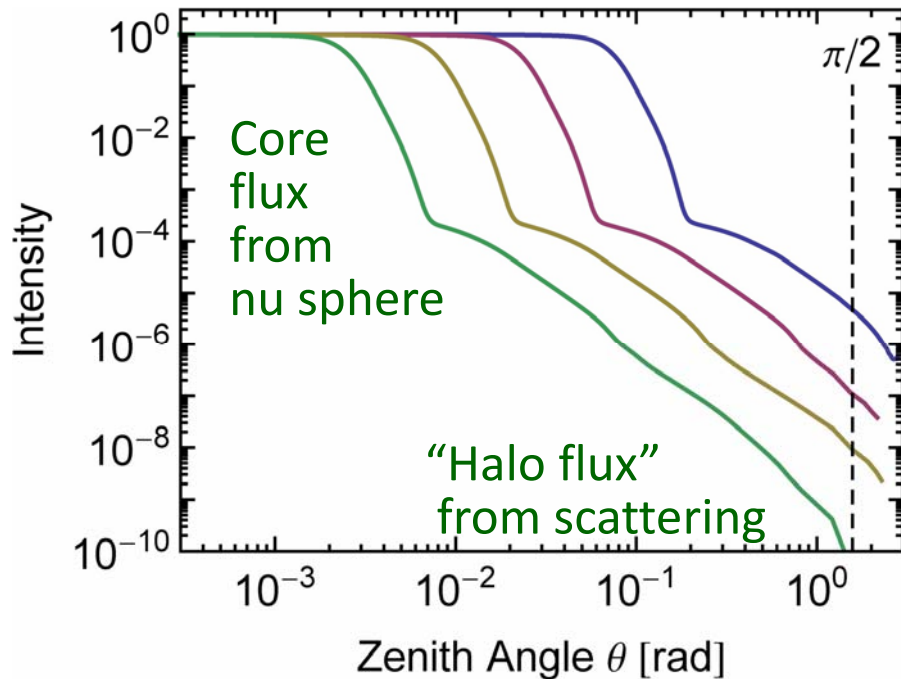
Small “scattering halo” important for nu-nu refraction? (Cherry et al., arXiv:1203.1607)



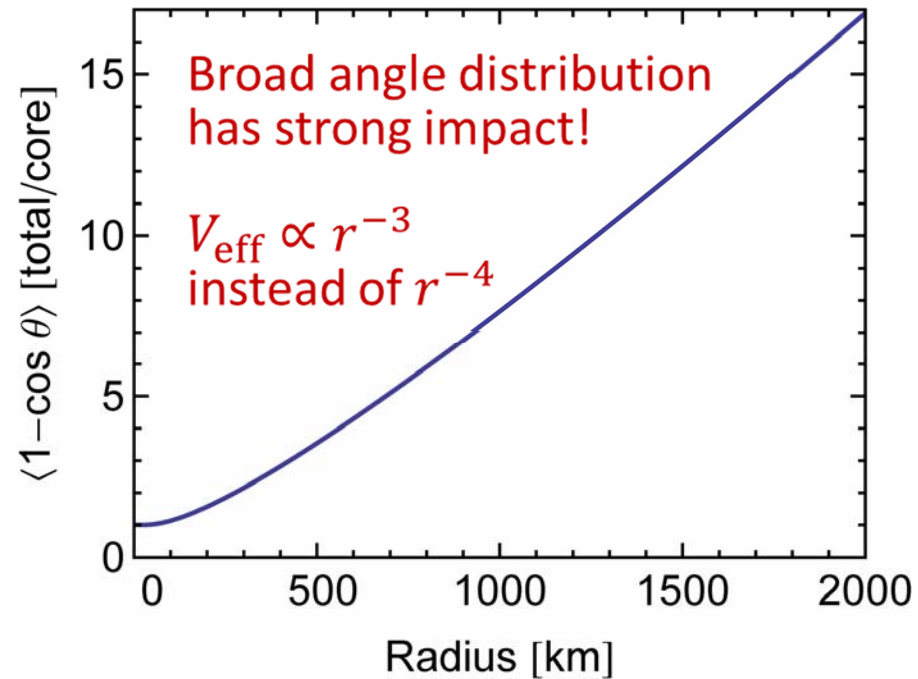
Picture from Ott, Burrows, Dessart & Livne, arXiv:0804.0239

Scattered Neutrinos as a Source of Refraction

SN neutrino angle distribution seen at 10^4 , 3000, 1000 and 300 km

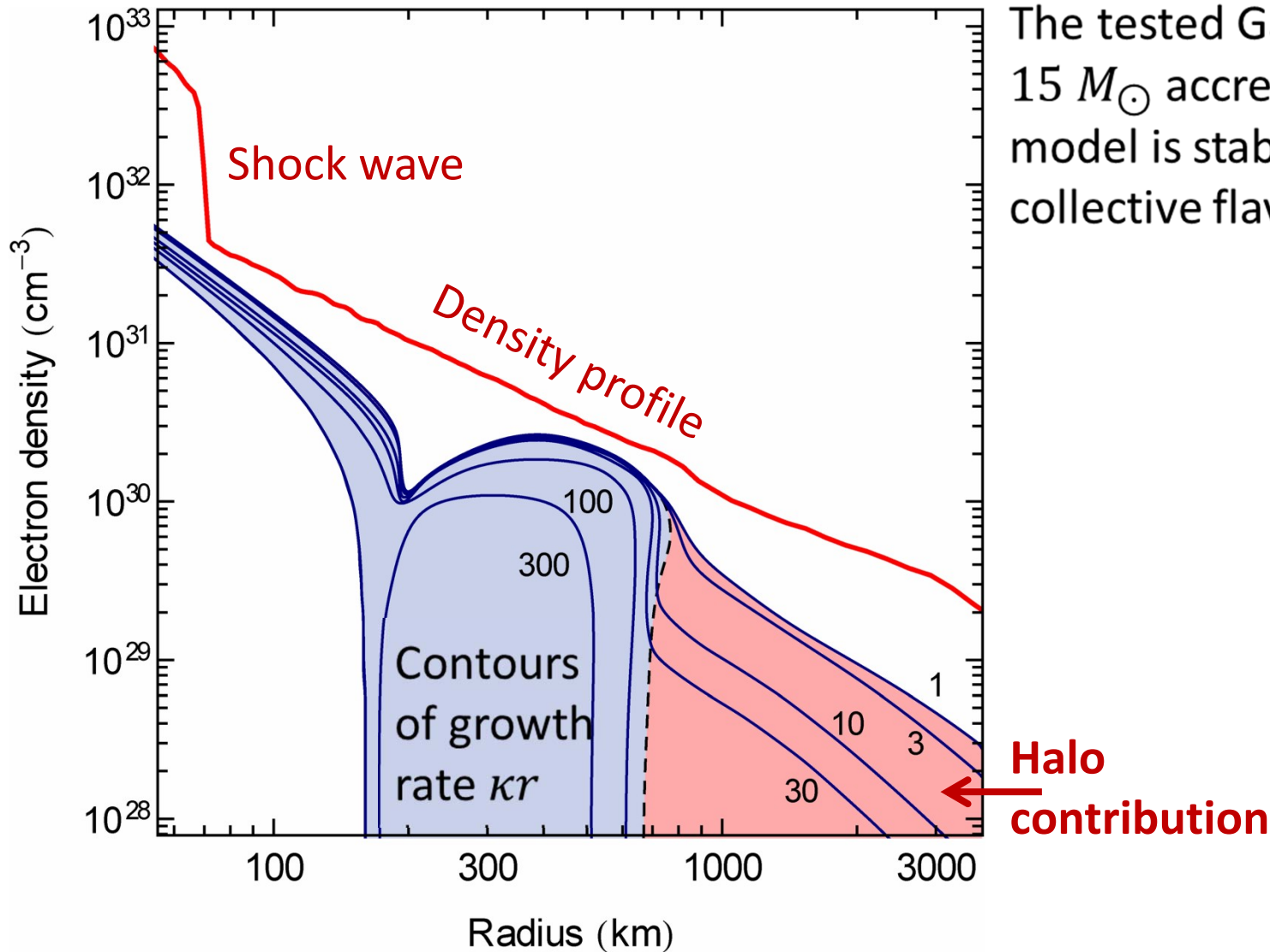


Relative importance of halo for nu-nu refractive potential



Cherry, Carlson, Friedland, Fuller & Vlasenko, arXiv:1203.1607
Sarikas, Tamborra, Raffelt, Hüdepohl & Janka, arXiv:1204.0971

Multi-Angle Matter Suppression in Realistic Model



The tested Garching $15 M_{\odot}$ accretion-phase model is stable against collective flavor conversion

Sarikas, Tamborra, Raffelt, Hüdepohl & Janka, arXiv:1109.3601, 1204.0971

Summary

Supernova neutrino flavor evolution remains a complicated subject

- Axial symmetry was always assumed – too symmetric?
- Numerical treatments challenging
- Novel role for neutrino “scattering halo”?
- Simultaneous space and time dependence important?

Theoretical developments

- Analogy to BCS theory
- Linearized stability analysis provides many conceptual insights
- And practical results

Working hypothesis for SN neutrinos

- Multi-angle matter effect can prevent instability
- No collective conversion during early accretion phase
- Can test for nu mass hierarchy (because θ_{13} is large)



**More theory progress is needed to reliably interpret
neutrino signal of next galactic supernova!**