Multidimensional CCSN Simulations with FLASH

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FLASH as a tool for studying CCSNe
EOS effects in simulations with parametric neutrino heating in 1D and 2D

Beyond 2D

Magnetorotational CCSN



FLASH Capabilities Span a Broad Range...







experiment



Laser-driven shock instabilities



Gravitationally confined detonation



Rayleigh-Taylor instability



Orzag/Tang MHD vortex



Intracluster interactions



Richtmyer-Meshkov instability



Nova outbursts on white dwarfs



Cellular detonation



Flash Center for Computational Science The University of Chicago



FLASH Capabilities Span a Broad Range...





FLASH v4.0

Fryxell et al. (2000) – My FLASH CCSN application shares no lines of code in common with FOO.

We are preparing a new 'FLASH' paper: Lee, SMC, et al. (in prep). See also A. Dubey et al. (2009).

- Ø Open-source. Get it at: flash.uchicago.edu
- Subset Contributions accepted!

FLASH v4.0

- Directionally-unsplit staggered mesh MHD solver with constrained transport
- Ideal and non-ideal MHD
- Reconstruction orders: 1 (Godunov), 2 (MUSCL-Hancock), 3 (PPM), 5 (WENO)
- Multiple slope limiters and Riemann solvers (HLL, HLLC, HLLD, Roe, Lax-Friedrichs,...)
- New multipole Poisson solver. Significantly faster, more accurate and efficient.
- Multigroup FLD with HYPRE linear algebra

FLASH v4.0 for CCSNe

- Extension of unsplit solvers to spherical and cylindrical geometries
- Addition of 1.5D and 2.5D rotation
- Finite temperature EOS (via E. O'Connor, stellarcollapse.org)
- Neutrino `lightbulb' heating/cooling
- Deleptonization a la Liebendorfer (2005)
- And new this week: Ray-by-Ray Neutrino Leakage (from GR1D, E. O'Connor)

3D CCSNe Simulations Require Petascale Computing!

- About 70 million zones in 3D (1% angular/radial resolution)
- Argonne Leadership Computing Facility: Intrepid (557 Tera-FLOP BG/P) and next year, Mira (10 Peta-FLOP BG/Q)







*MHD perfect scaling achieved with bigger AMR blocks

First Steps: The influence of EOS in parameterized simulations (Couch 2012)

$$\mathcal{H} = 1.544 \times 10^{20} \left(\frac{L_{\nu_e}}{10^{52} \text{ erg s}^{-1}} \right) \left(\frac{T_{\nu_e}}{4 \text{ MeV}} \right)^2 \\ \times \left(\frac{100 \text{ km}}{r} \right)^2 (Y_p + Y_n) e^{-\tau_{\nu_e}}$$

$$C = 1.399 \times 10^{20} \left(\frac{T}{2 \text{ MeV}}\right)^6 (Y_p + Y_n) e^{-\tau_{\nu_e}}$$

3 EOS in 1D and 2D with s15s7b2: HShen and Lattimer & Swesty (K=180,220 MeV)

HShen, L=1.7e52



L=1.3e52, 600 ms post-bounce:







- 2D is significantly easier than
 1D
- STOS
 LS curves are lower than
- Solution LS180 is lower than LS220





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Why do the explosion times depend on EOS?

 Difference in alpha-particle abundance causing difference in (Y_p+Y_n)? Tried (1.-Y_H)² in heating/cooling terms. No qualitative change.

The Difference in buoyant convection/turbulence? Look at optical depth through gain region (Murphy, Burrows, & Dolence 2012): $E_k \sim L_{\nu_o} \tau$

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 Could be due to different development of the SASI. The softer LS EOS produce more compact, denser PNS and advective deceleration regions. Thus, acoustic wave generation is more efficient. Denser PNS are better sonic transducers.

Looking into 3D

Simulations not done, so sorry, no conclusions yet.





What I'd like to do next...

Magnetic, rotating progenitors!

How important could MR be for normal supernovae? Aiding neutrinos, shaping explosions, driving explosions?

How does MR affect the SASI, convection, and turbulence in a range of progenitors?

What is the threshold in initial P/B above which MR is important?

Can we better capture the important dynamics of unresolved amplification mechanisms?

Rotation and Magnetic Fields Are Important!

- Important implications for initial spin and B-field of neutron stars
- Possible r-process site (e.g.
 Winteler et al. 2012)
- Connection to GRBs
- Possible explanation for observational evidence of bipolar SNe (but see, e.g., Hammer et al. 2010)



Winteler et al. 2012

$$\begin{split} & \textbf{Magnetorotational CCSNe} \\ & E_{\rm rot,PNS} \approx \frac{1}{2} I_{\rm PNS} \Omega_{\rm PNS}^2 \\ & \approx 9 \times 10^{50} \text{ergs} \left(\frac{M_{\rm PNS}}{1.5 \ M_{\odot}} \right) \left(\frac{\Omega_{\rm PNS}}{250 \ \text{s}^{-1}} \right)^2 \left(\frac{R_{\rm PNS}}{50 \ \text{km}} \right)^2 \end{split}$$

Explosions by MHD jets first suggested by LeBlanc & Wilson (1970). Explored further in Meier et al. (1976)

Possible importance revisited by Wheeler et al. (2000,2002).

Recent work by Burrows et al. (2007) and Takiwaki et al. (2009) in 2D, Winteler et al (2012) in 3D.

Magnetorotational SNe





Burrows et al. 2007

B-field Amplification in CCSNe

See, e.g., Endeve et al. (2010,2012), Obergaulinger & Janka (2011)

Field compression: field carried along with collapsing plasma: "flux-freezing"

- Field winding: linear process, wraps up field lines. $B_{\phi} \approx 2\pi n_{\phi} B_p$
- Possible (small-scale) dynamo

 Magnetorotational Instability (MRI): exponential growth of initial field.
 Saturation field strengths as high as 10¹⁵ – 10¹⁶ G. Akiyama et al. (2003). MRI in the Linear Regime E.g., Balbus & Hawley (1992), Obergaulinger et al. (2009) Considering only radial gradients, instability criterion:

$$\begin{split} C_r &= \begin{bmatrix} \frac{1}{\rho} \frac{\partial P}{\partial r} \left(\frac{1}{\rho} \frac{\partial \rho}{\partial r} - \frac{1}{\Gamma_1 P} \frac{\partial P}{\partial r} \right) + r \frac{\partial \Omega^2}{\partial r} \end{bmatrix} / \Omega^2 < 0 \\ \text{gravity} \qquad \text{buoyancy} \qquad \text{rotation} \end{split}$$

 $k_{\rm FGM} \approx [-C_r (C_r + 8)]^{1/2} \Omega / 4 v_A \qquad \omega_{\rm FGM} \approx (-C_r^2)^{1/2} \Omega / 4$

 $B_f \sim B_0 e^{i\omega_{\rm FGM}t}$ $t_{\rm MRI} \sim \ln(B_f/B_0)/i\omega_{\rm FGM}$

 $B_{\rm sat} \approx \alpha r \Omega (4\pi\rho)^{1/2}, \ \alpha < 1$

Magnetic, Rotating Progenitors Heger, Woosley, & Spruit (2005)

Magnetic torques slow rotation of cores by a factor of 30–50 relative to non-magnetic progenitors

standard 15 M_{sun} model: 0.2 s⁻¹ (v. 8.0 s⁻¹)
 1.5D collapse simulations for 0.2, 1.0, 2.0 s⁻¹



Magnetic, Rotating Progenitors

Magnetic, Rotating Progenitors

 Is (unresolved) MRI turbulence important?
 Angular momentum transport, viscous
 heating (e.g.
 Thompson et al.
 2003).

Progenitor structure is still a big uncertainty. Mass loss, binarity, etc.

Conclusions

FLASH is a great tool for astrophysical simulation, and now CCSNe as well

In parametric-neutrino simulations, the timeto-explosion depends on EOS

Possibly due to enhanced acoustic flux from PNS accelerating SASI growth

FLASH 3D parametric simulations on the way!

Magnetorotational effects need further study. Could be very important for certain progenitors

3D Parameterized Models





$$F_{100}^{-} = \dot{M}c_{S}^{2} \frac{1 - \mathcal{M}}{\mathcal{M}} \left(\frac{\delta P}{\Gamma_{1}P}\right).$$



SN 1987A











Cas A



SN Polarization

- ALL core-collapse SNe are polarized
- Higher asymmetries in the cores of explosions
- Often show a "dominant axis" in Q/U plane – indicates an elongated explosion
- Loops in Q/U plane indicate non-axisymmetry



Wang & Wheeler 2008

Type IIP Polarization

SN 2004dj



Leonard et al. 2006

What Do the Observations Tell Us?

- Massive stars explode all the time, with energies around 10⁵¹ erg!
- They are NOT spherically-symmetric
- They often show general 'bi-polarity' with significant non-axisymmetry and time-dependent polarization.
- They leave remnants that often have high kick velocities and strong magnetic fields.
- Some CCSNe are associated with GRBs.
- Mixing & overturn commonly indicated.





Nordhaus et al. 2010



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