Multi-dimensional Progenitors of Core Collapse Supernovae

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Topics:

•Pre-collapse dynamics: O+O and Si shell interaction (a major issue!)

- •A Turbulence Model to replace mixinglength theory (MLT) for evolutionary time scales
- Rotation and convection: a bifurcation in progenitor evolution (first results)

C, Ne and O burning versus C, Ne, O and Si burning

- Map same 1D model onto a 2D grid
- Use the same microphysics
- Allow time derivatives to be nonzero
- What happens? Are 1D models OK?

- We need 3D simulations to confirm these 2D results (a major computational challenge; Meakin and Arnett, in progress)
- We need full 4 pi steradian geometry to get the lowest order modes
- Contrary to conventional wisdom, progenitor models are a MAJOR uncertainty for core collapse! (e.g., the compactness parameter of O'Conner & Ott 2012; see also Ugliano et al 2012)

Toward a 3D Turbulence theory:

- \bullet numerically simulate turbulent convection in a realistic stellar model (O burning shell in a collapse progenitor)
- \bullet theoretically analyze the numerical data
- \bullet synthesize a theory which is useable in a stellar evolution code (replaces MLT) for long term evolution
- anelastic and low-mach solvers give an increase in time step of only 10-100 for realistic progenitors due to accuracy limitations (good but not big enough for evolution)
- •implicit solvers (Viallet) are also limited
- \bullet evolutionary time steps require reformulation to "step over" turbulent variations (Lorenz attractor)

New results:

- Kolmogorov damping balances buoyancy (to fix velocity scale and remove adjustable alpha parameter)
- **turbulent braking layers enclose Schwarzschild convection** zones (overshoot)
- Fluctuations, cells and Lorenz strange attractor (dynamic systems theory)
- Turbulent heating (thermodynamic consistency)

Rotation and Convection: a beginning

Case 1: Deep interior

triangles denote top and bottom convective boundary (and shear layer)

Case 2: Surface

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Schou, et al., 1998, ApJ

FIG. 5.—As a companion to Fig. 3, inversions for rotation rate $\Omega/2\pi$ with radius and latitude for four inversion methods: (a) 2dRLS; (b) 2dSOLA; (c) $1d \times 1dSOLA$; (d) 1.5dRLS.

lying angular velocity (cf. The error bars indicate only α). The error bars indicate only α

considerable degree of consistency between the results from

•Balbus & Weiss (2009, MNRAS) show that this behavior does NOT require a magnetic field, but may be a result of purely hydrodynamic, turbulent flow.

Case 3: Red Giant

Figure 11. We display at mid-convection zone for cases *RG1* (left) and *RG2* (right), the angular velocity as a function of latitude for four different temporal intervals along with the average over the full temporal range sampled (solid line). For *RG1* we display four consecutives 1 rotation period averages, and for *RG2* we have chosen the temporal intervals corresponding to Figure 14 with either one or two meridional circulation cells. We clearly note that the angular velocity profile for *RG1* is stable

(A color version of this figure is available in the online journal.)

slope for the specific angular momentum *j*. Moving toward the top of the domain, the mean radial angular velocity dramatically *4.2. Structure of Meridional Flows* Another important large-scale flow established in rotating It is unclear how the ASH energy non-conservation in stratified environments affects this result (Brown et al, 2012).

convective envelopes is the mean (axisymmetric) meridional

The profiles achieved by models *RG1* and *RG1t* are very similar

both in shape and amplitude. Similarly, the profiles obtained for

models *RG2* and *RG2t* are also almost identical. This indicates

that the distribution of specific angular momentum (and of the

mean radial angular velocity) achieved in the simulations is

not influenced by the turbulent state (e.g., Reynolds number)

of the simulations up to the values that we have been able

to compute. This is a quite different behavior compared to the

solar convection, where lowering the diffusivities while keeping

Pr constant in order to increase *Re* by a factor of 2 leads to

some modification of the angular velocity profile (see Figure 4

of Brun & Toomre 2002). On the other hand, the specific

angular momentum distribution achieved in our convective shell

strongly depends on the bulk rotation considered. For models

- We see a variety of behavior!
- It seems to depend upon Rossby number (ratio of inertial to Coriolis force), and perhaps stratification
- above $j \sim 10$ ^14 cm^2/sec, rotation will dominate convective mixing, giving a bifurcation in behavior (and a new class of progenitors)
- The insights from 3D simulations should be applied to calculations of stellar evolution
- We must regard our understanding of the evolution to core collapse as incomplete at best