



New approaches to multi-dimensional radiation transport

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Monte Carlo

E. Abdkamalov, A. Burrows, C. D. Ott, F. Löffler, E. O'Connor, J. Dolence, E. Schnetter, 2012, ApJ

Filtered Spherical Harmonics (FP_N)

D. Radice, E. Abdkamalov, L. Rezzolla, and C. D. Ott, *in prep*

Monte Carlo

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Transport equation

$$\frac{1}{c} \frac{\partial I}{\partial t} + \nabla I = \eta - \kappa I$$

Transport equation

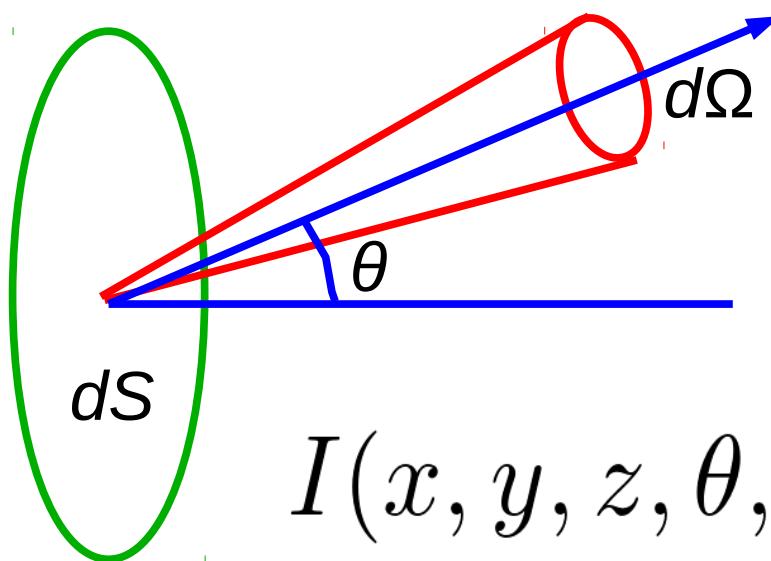
$$\frac{1}{c} \frac{\partial I}{\partial t} + \nabla I = \eta - \kappa I$$

The diagram illustrates the components of the transport equation. It features three blue-outlined circles containing mathematical symbols: ∇I , η , and κI . Blue arrows point from these circles to red text labels below them: 'flux' under ∇I , 'absorption' under η , and 'extinction' under κI .

Transport equation

$$\frac{1}{c} \frac{\partial I}{\partial t} + \nabla I = \eta - \kappa I$$

flux absorption extinction



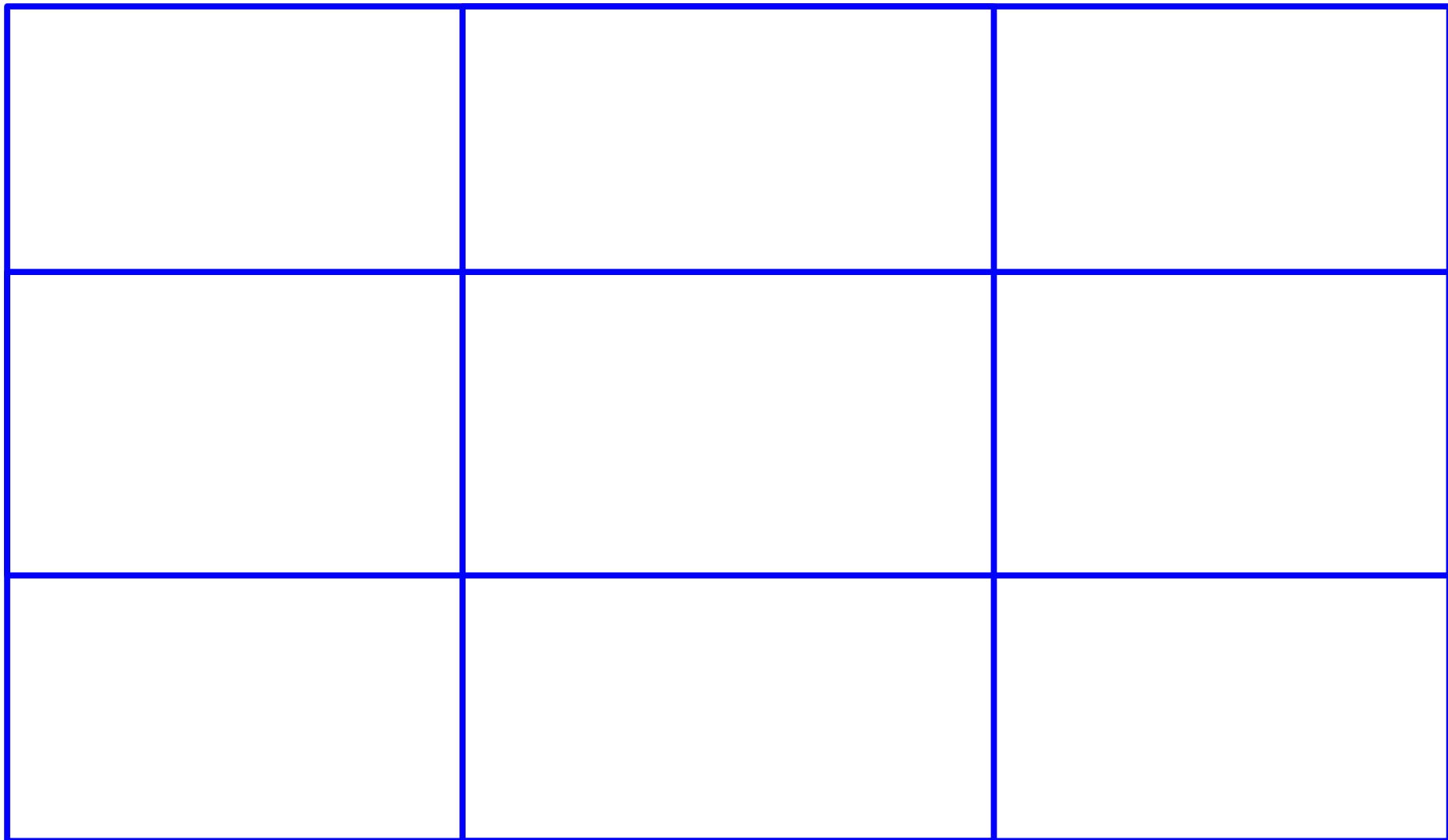
$$I(x, y, z, \theta, \phi, \nu, t) = \frac{dE}{d\Omega dS d\nu \cos \theta dt}$$

Deterministic **and** Monte Carlo radiation transport

Monte Carlo Transport



Monte Carlo Transport



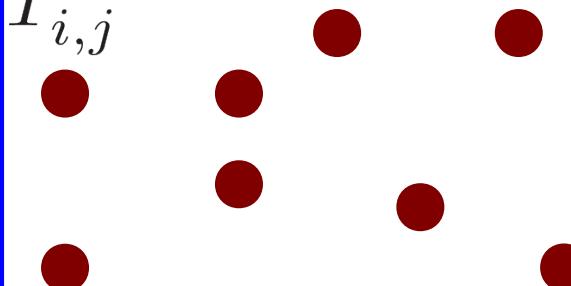
Monte Carlo Transport

$T_{i-1,j+1}^n$	$T_{i,j+1}^n$	$T_{i+1,j+1}^n$
$T_{i-1,j}^n$	$T_{i,j}^n$	$T_{i+1,j}^n$
$T_{i-1,j-1}^n$	$T_{i,j-1}^n$	$T_{i+1,j-1}^n$

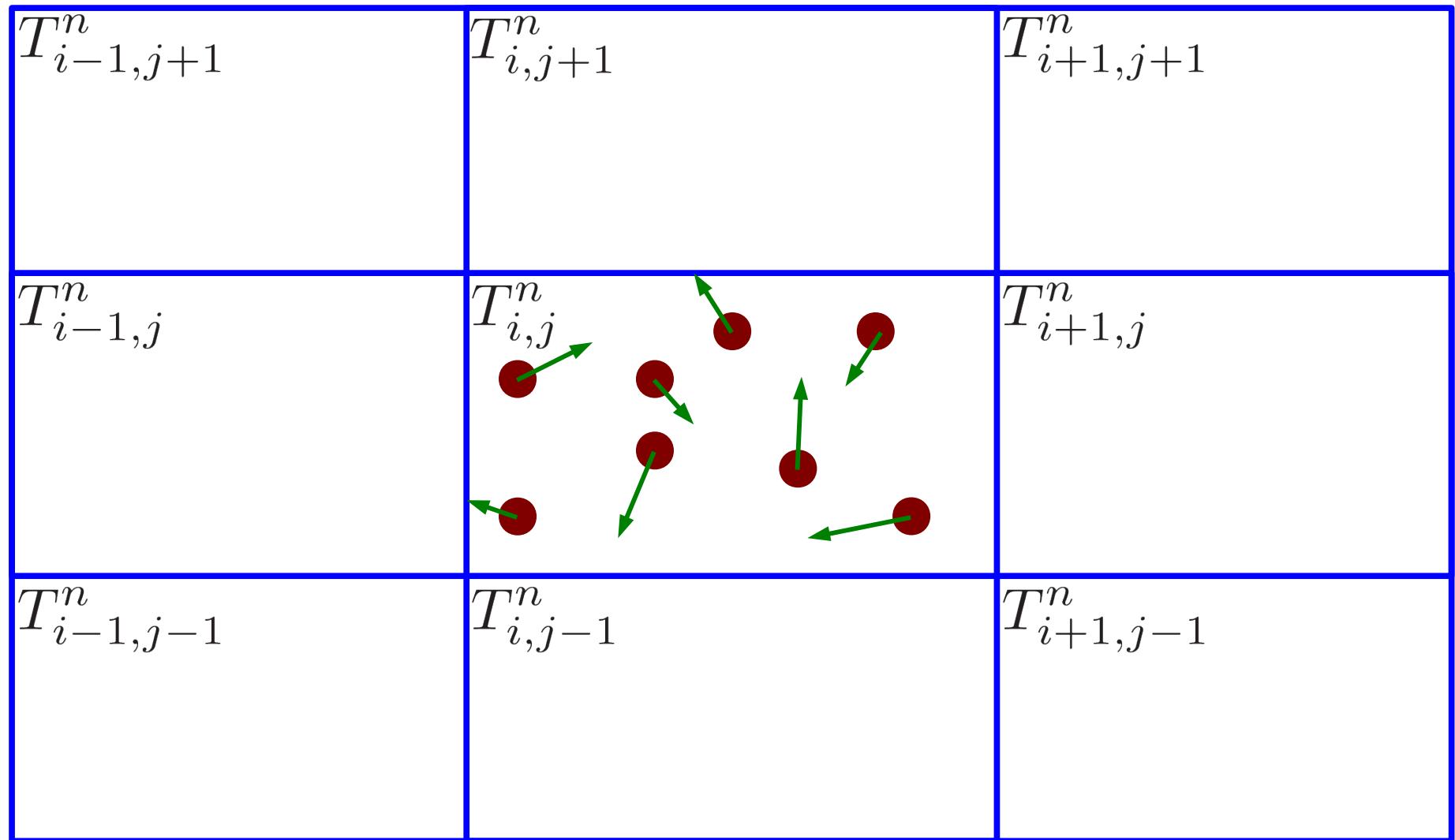
Monte Carlo Transport

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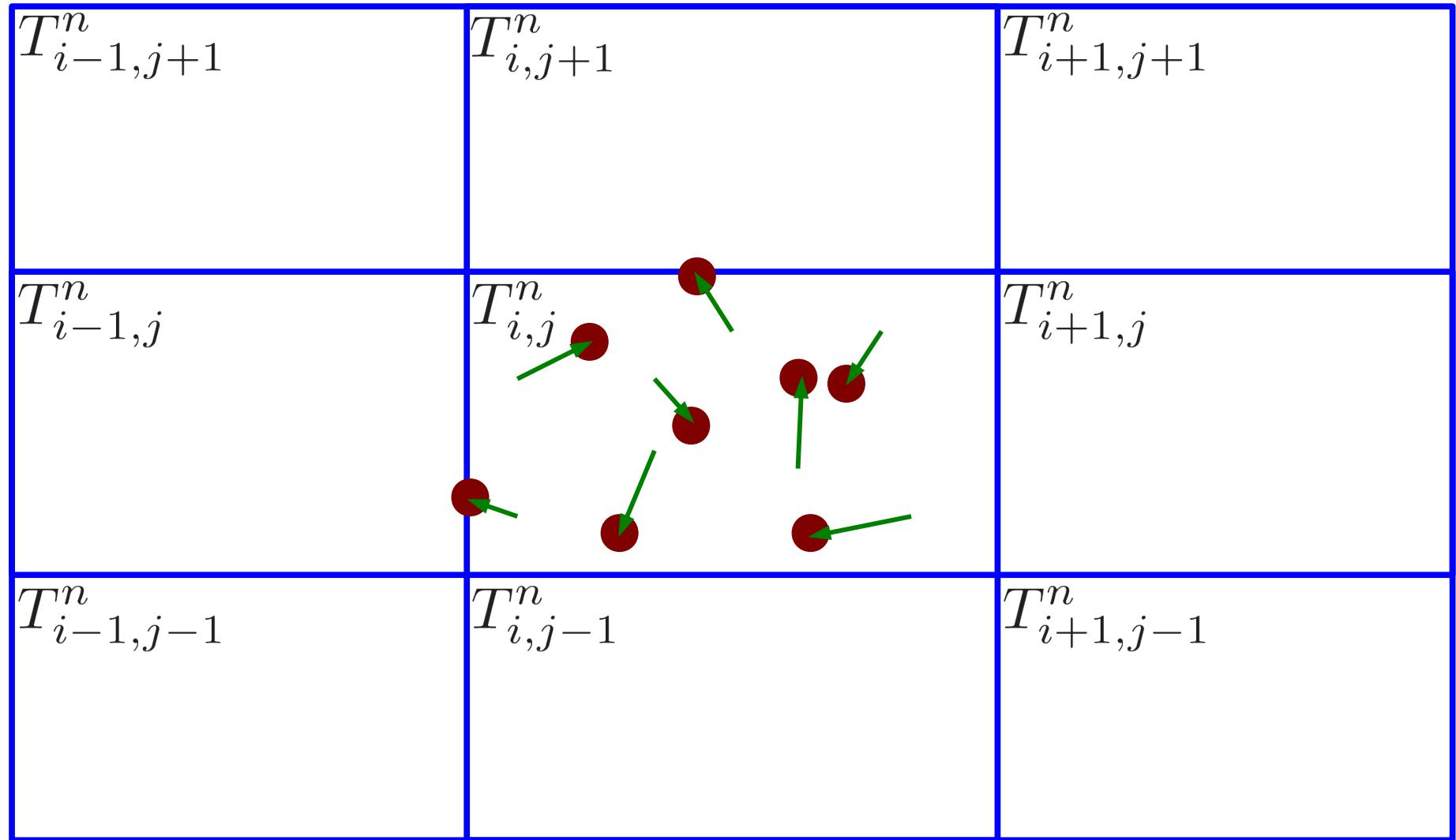
Monte Carlo Transport

$T_{i-1,j+1}^n$	$T_{i,j+1}^n$	$T_{i+1,j+1}^n$
$T_{i-1,j}^n$	$T_{i,j}^n$ 	$T_{i+1,j}^n$
$T_{i-1,j-1}^n$	$T_{i,j-1}^n$	$T_{i+1,j-1}^n$

Monte Carlo Transport



Monte Carlo Transport



Monte Carlo Transport

- Simple
- Easy extension to multi-D
- Parallelization
- Random noise
- Expensive

Monte Carlo transport scheme goal:

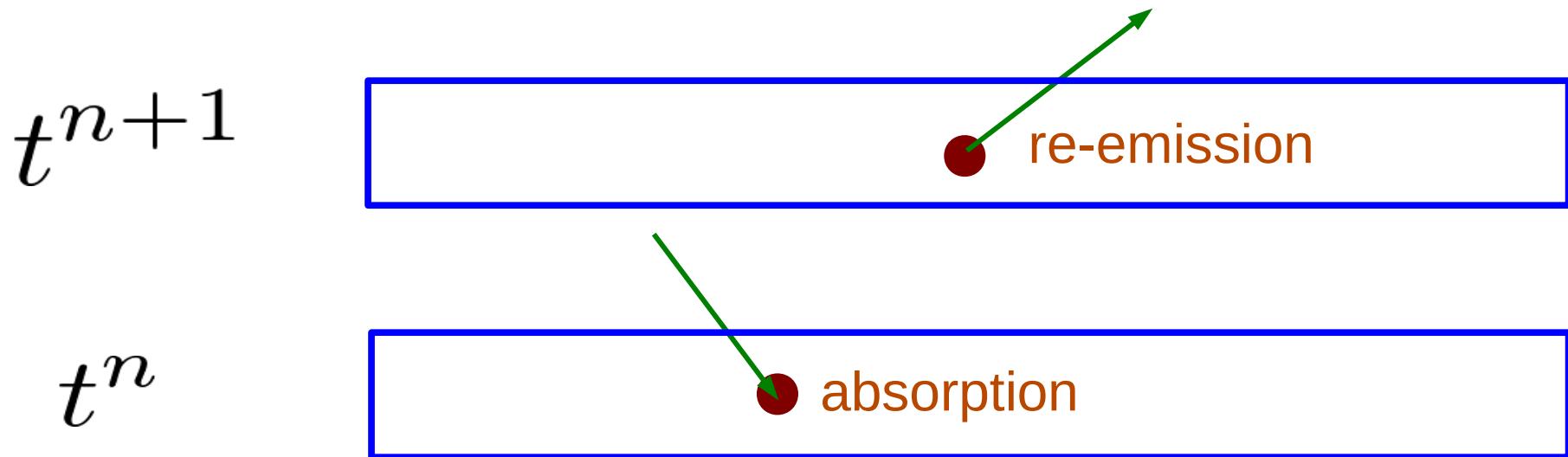
- Time-dependence
(coupling to matter energy and lepton number)
- Efficient treatment of high optical depth.
- Energy dependence
- Velocity dependence

A simple explicit discretization does not work!

$$\text{Timestep} \lesssim \frac{\rho c_v}{a T^3 c \kappa}$$

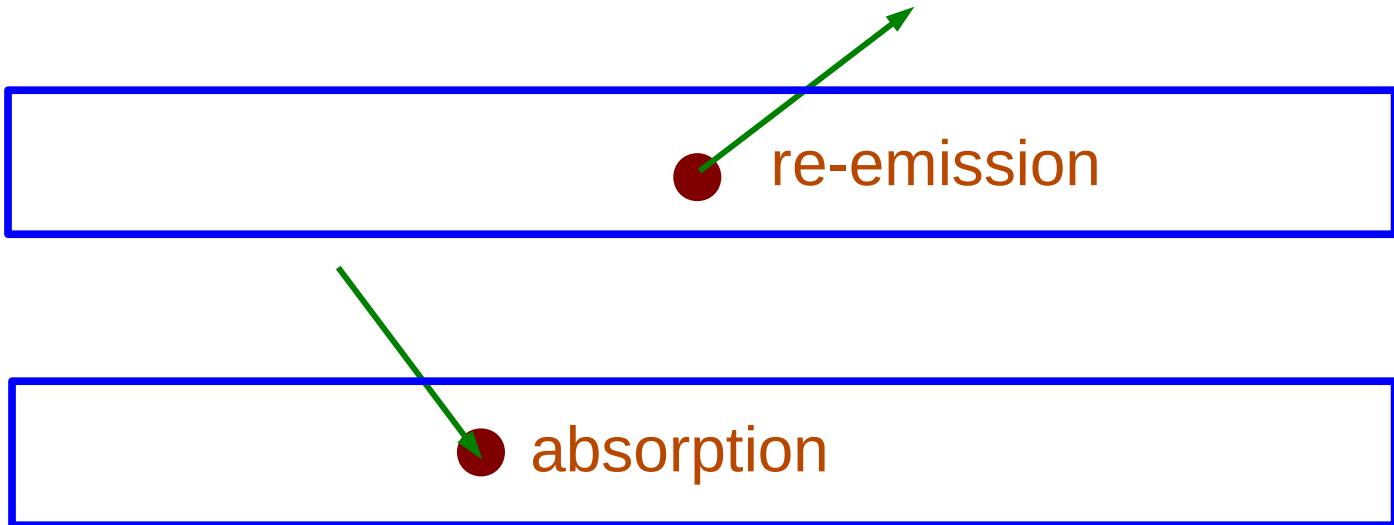
A popular solution:
Implicit Monte Carlo method
(Fleck & Cummings 1971)

A simple explicit scheme:



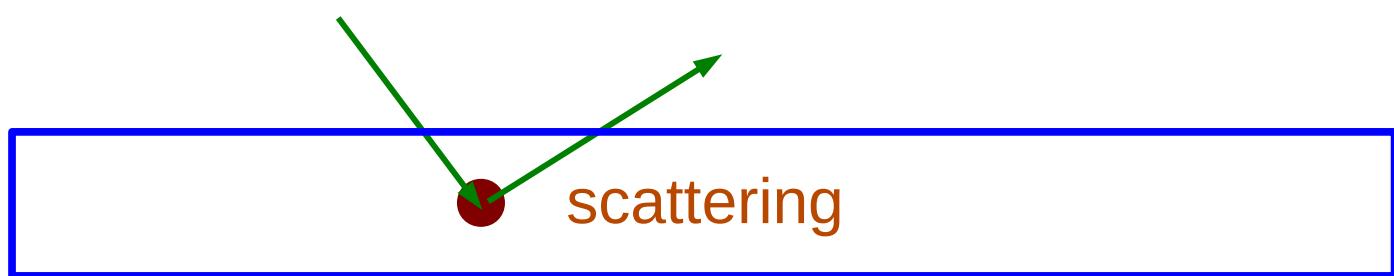
A simple explicit scheme:

t^{n+1}



Implicit Monte Carlo scheme:

t^n



Implicit Monte Carlo for Neutrinos

- Energy and lepton number coupling
- The same benefits as for photons
(large timestep, unconditional stability, accuracy)
- Both energy and lepton number conservation

Implicit Monte Carlo

at high optical depth

$T_{i-1,j+1}^n$	$T_{i,j+1}^n$	$T_{i+1,j+1}^n$
$T_{i-1,j}^n$	$T_{i,j}^n$ 	$T_{i+1,j}^n$
$T_{i-1,j-1}^n$	$T_{i,j-1}^n$	$T_{i+1,j-1}^n$

Implicit Monte Carlo

at high optical depth

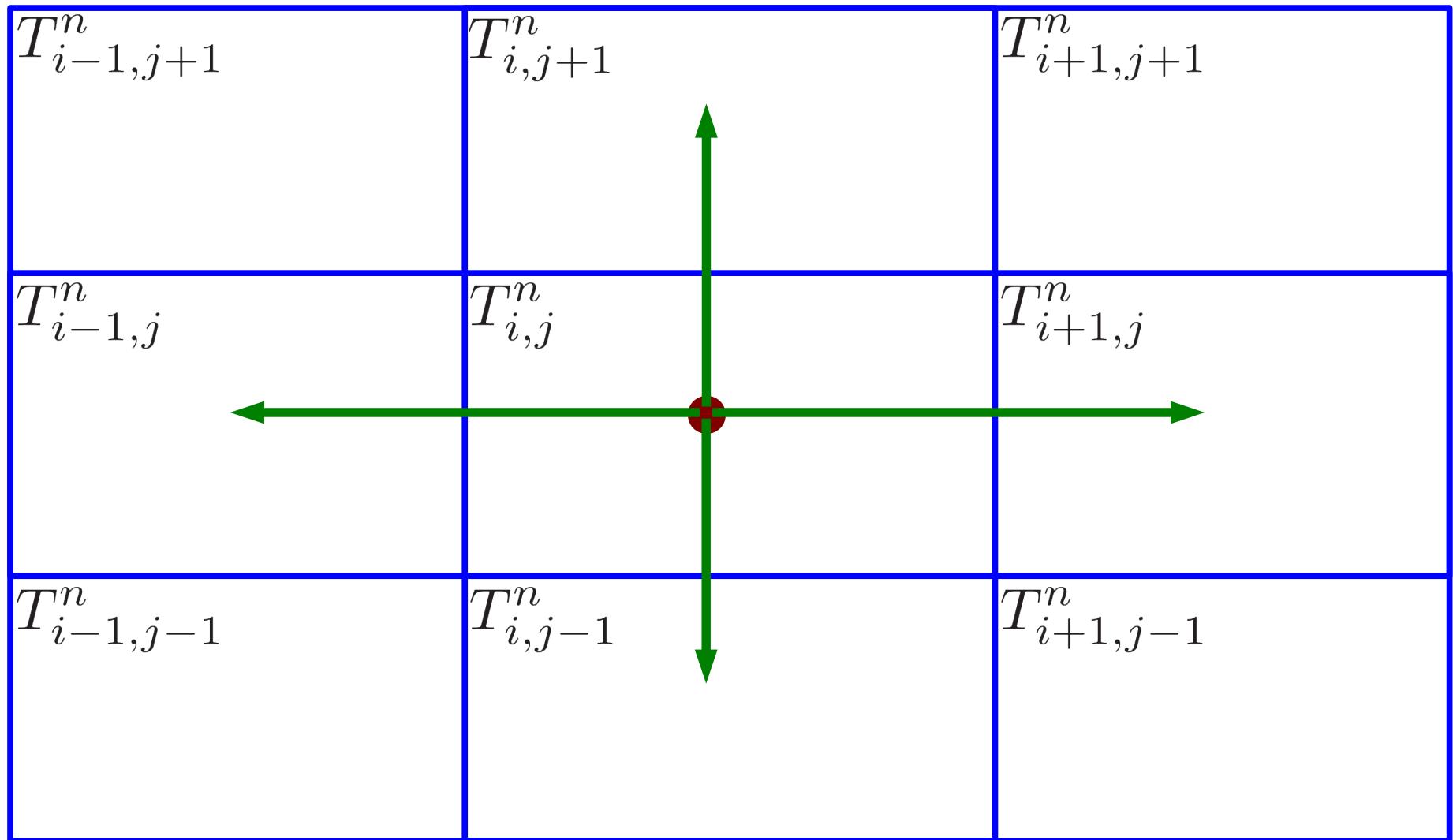
$T_{i-1,j+1}^n$	$T_{i,j+1}^n$	$T_{i+1,j+1}^n$
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$T_{i-1,j-1}^n$	$T_{i,j-1}^n$	$T_{i+1,j-1}^n$

Inefficient!

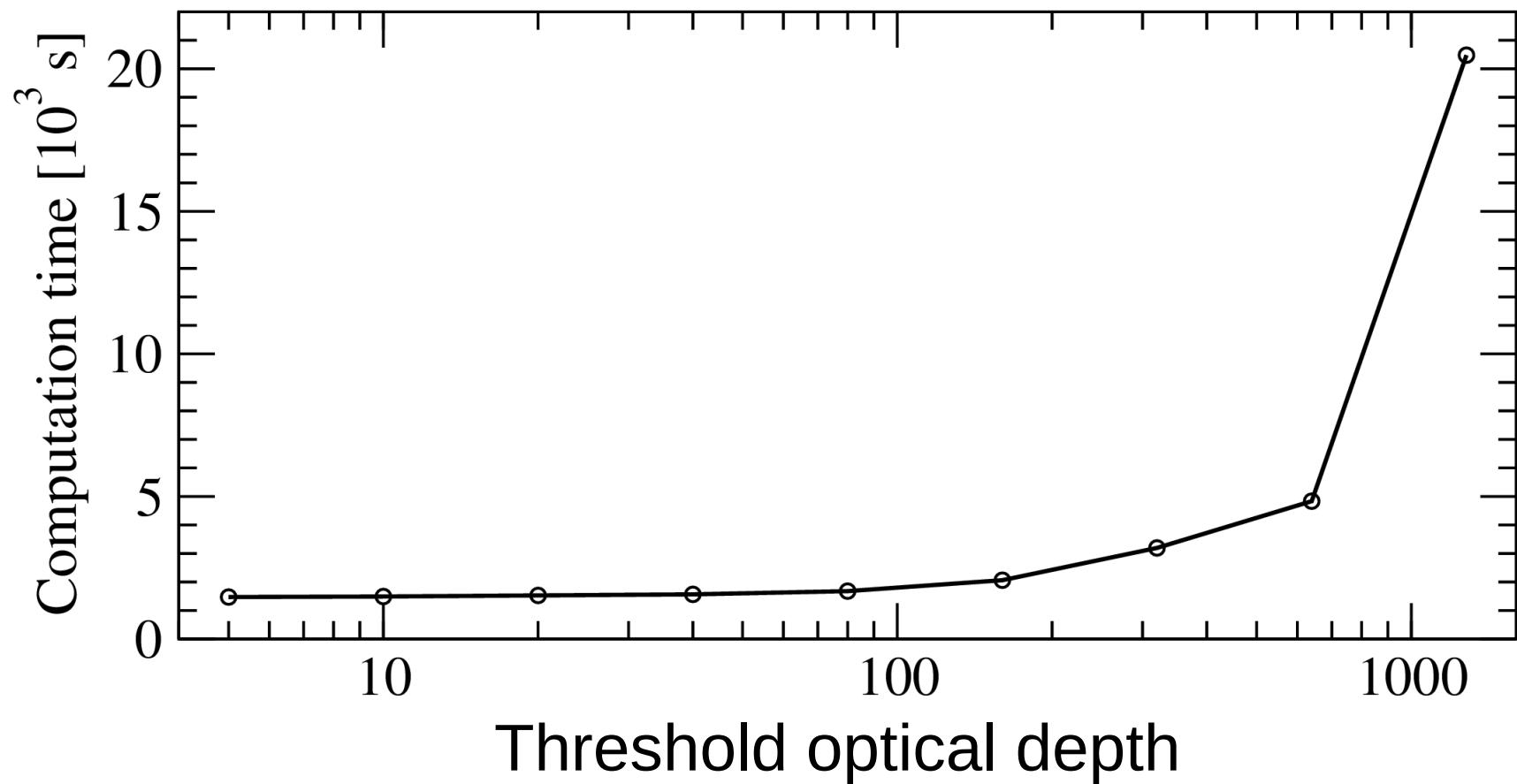


Discrete-Diffusion Monte Carlo

by Densmore+ '07 for gray transport for non-moving matter



Discrete-diffusion speed-up for proto-neutron star cooling



Velocity dependence

Velocity-dependent Monte Carlo

Mixed frame formalism

Eulerian frame: Transport

Comoving frame: emission, absorption, scattering

Velocity-dependent discrete-diffusion

Transport is performed in comoving frame
with $O(\lambda v/Lc)$ accuracy:

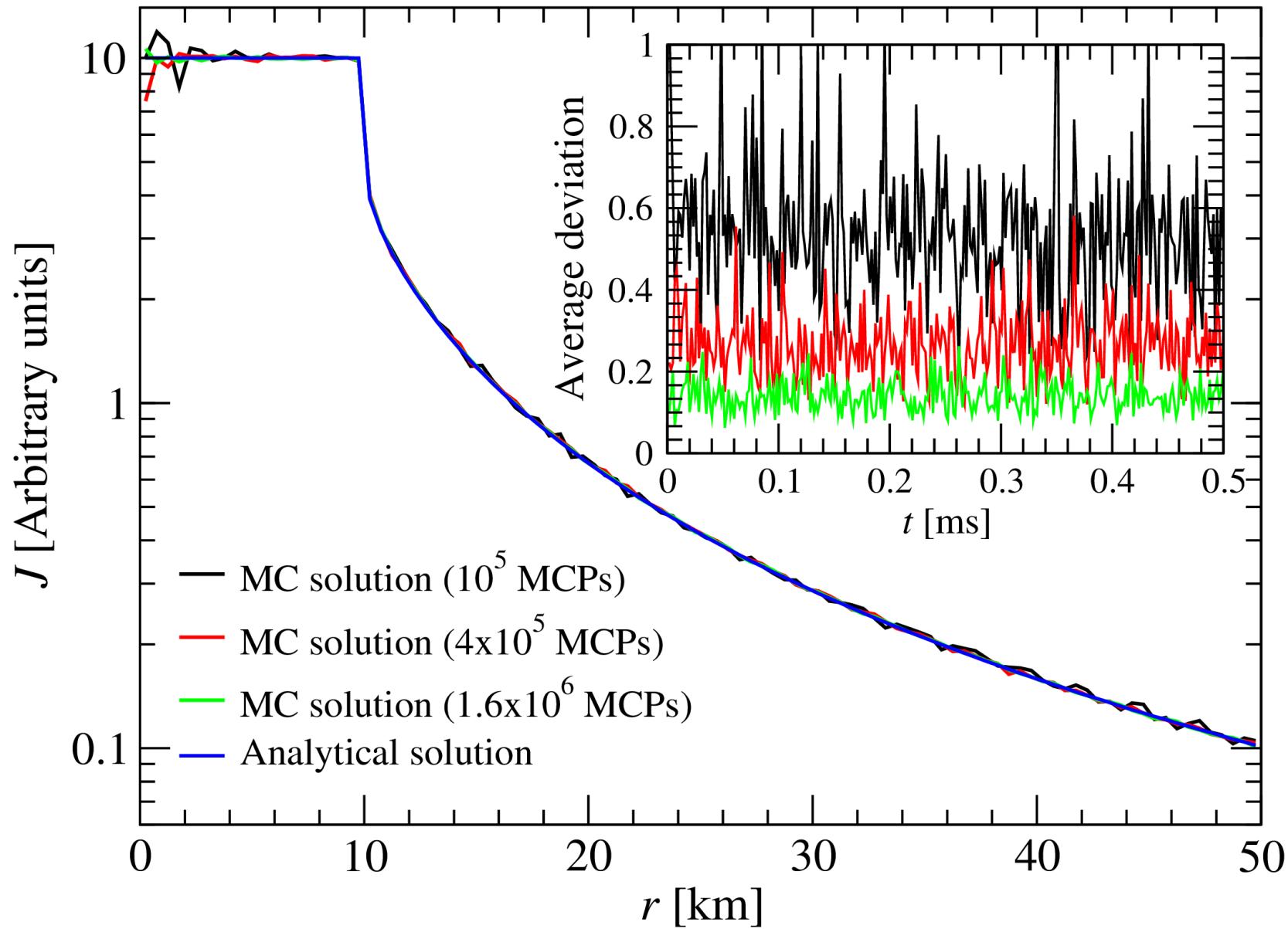
$$\frac{1}{c} \frac{dJ_0}{dt} + \frac{v}{c} \frac{\partial J_0}{\partial r} + \frac{J_0}{c} \frac{\partial v}{\partial r} + \frac{\varepsilon_0}{3c} \frac{\partial J_0}{\partial \varepsilon_0} \frac{D \ln \rho}{Dt} + \frac{\partial H_0}{\partial x} = \kappa_0 (B - J_0)$$

Three effects: advection,
compression/expansion, and Doppler shift.

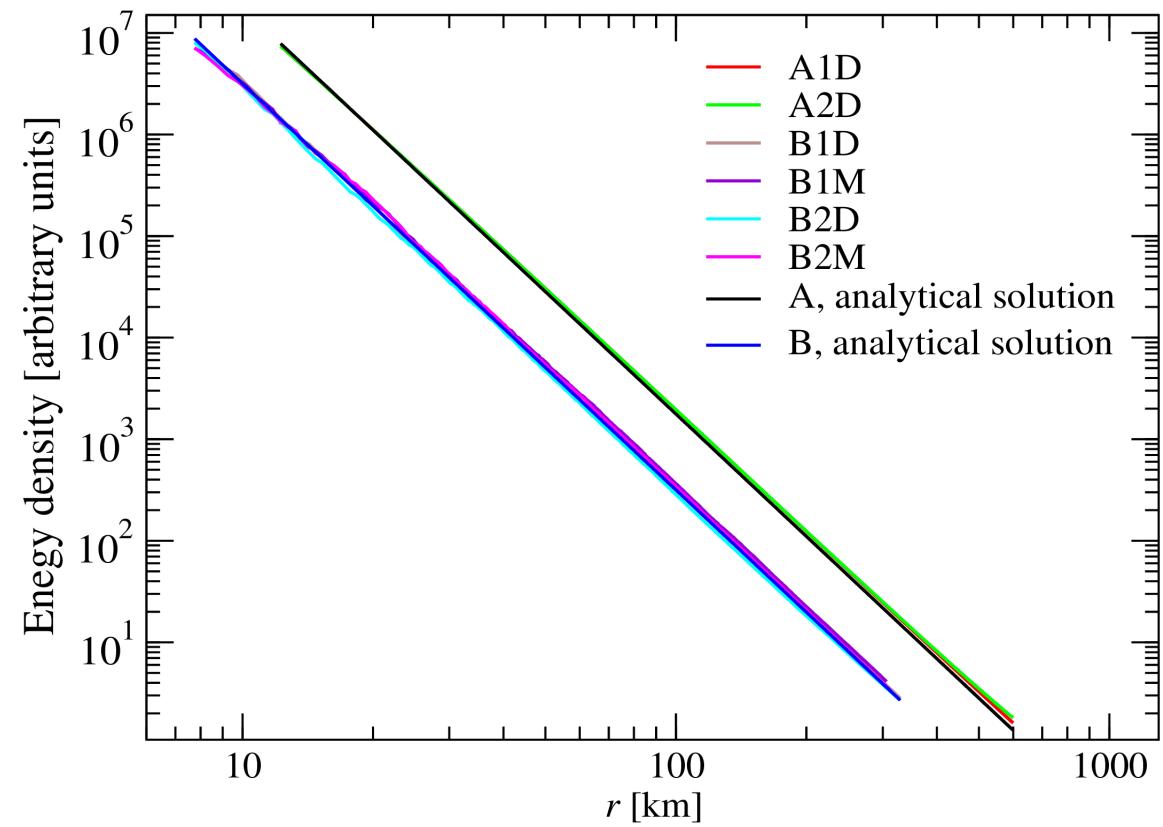
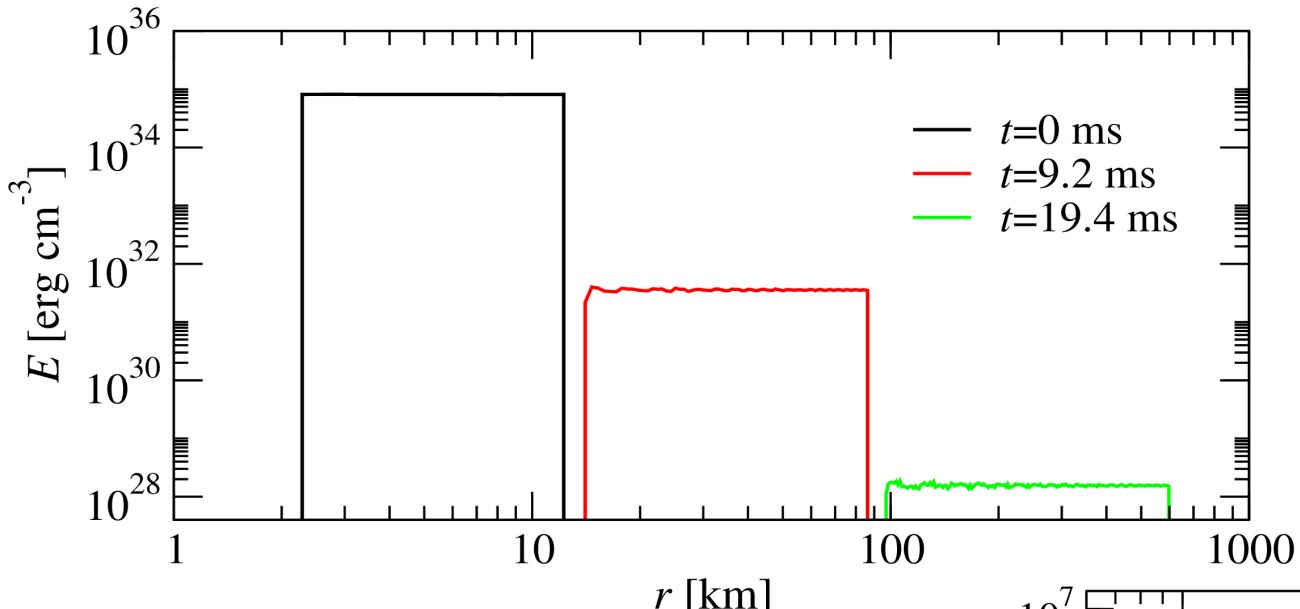
Solution method: Operator-splitting

Tests

Homogeneous sphere

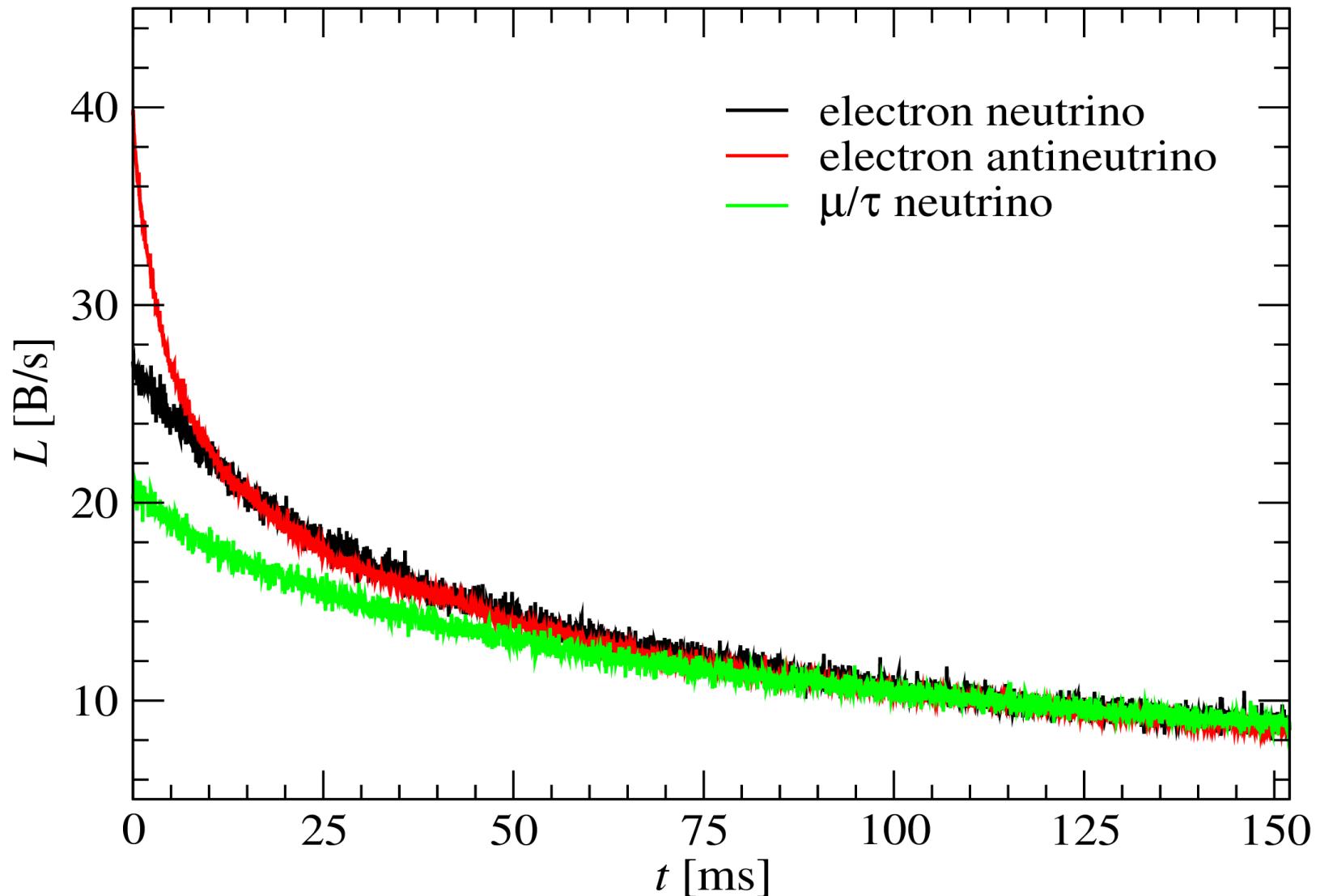


Homologously expanding shell



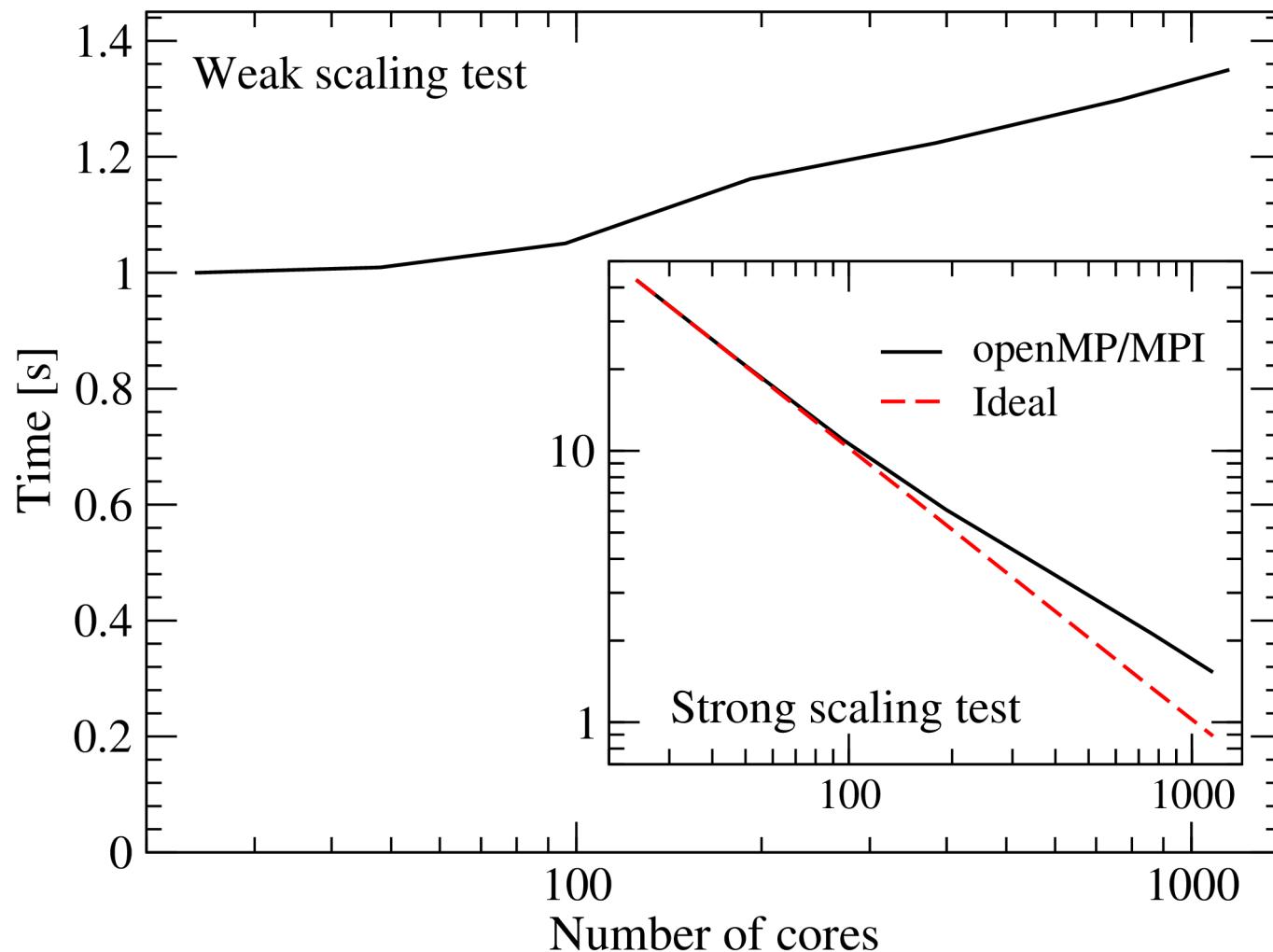
Proto-neutron star cooling

using Ott et al. (2008) PNS model



Parallel scaling

mesh replication method



Monte Carlo Summary

- Implicit MC for neutrino transport
- Multi-group discrete-diffusion
- Velocity-dependence
- Applicable to both neutrinos and photons
- Parallel scaling (in 1D)

Monte Carlo

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Filtered Spherical Harmonics (FP_N)

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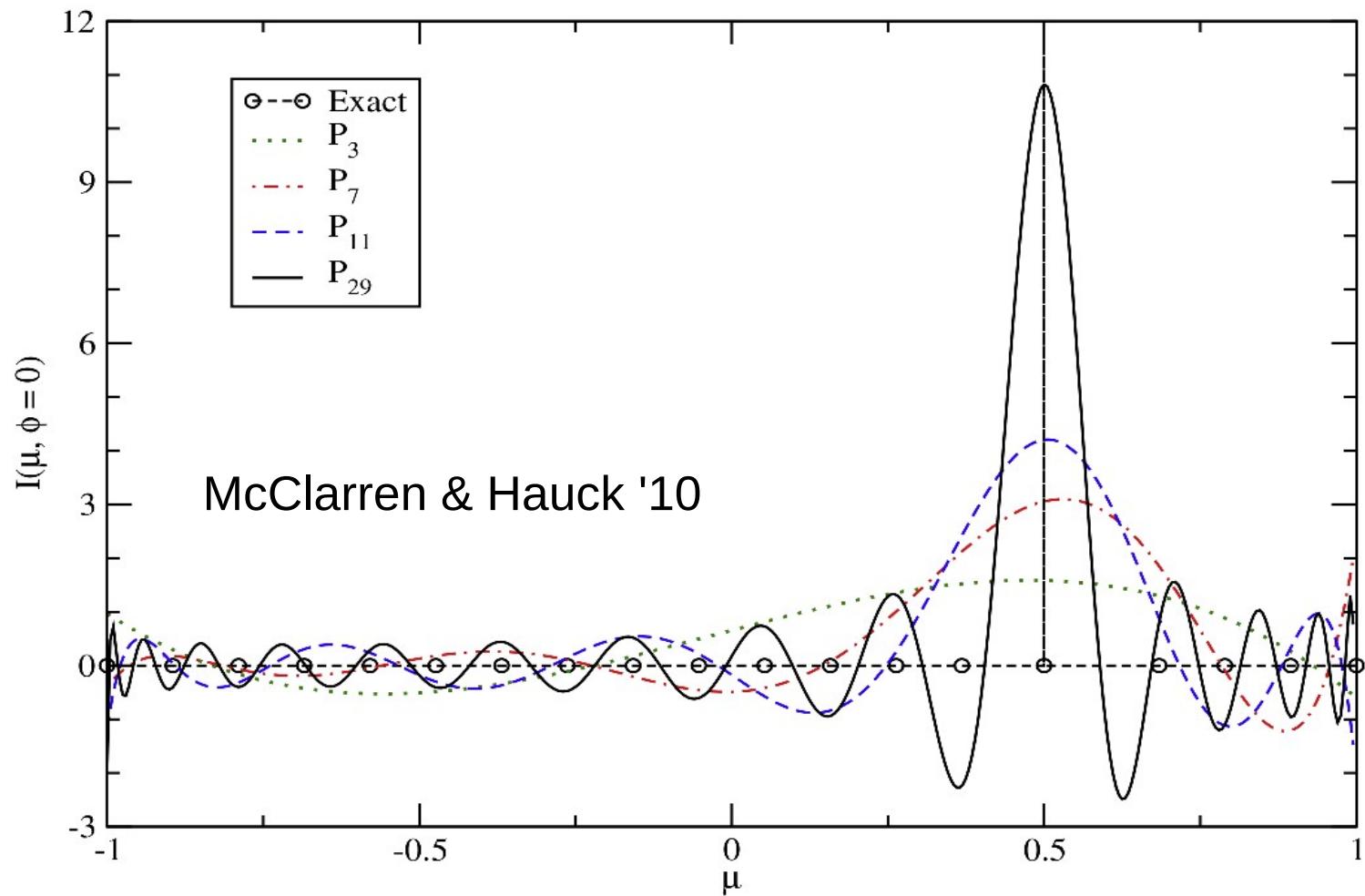
P_N scheme

$$I(x, \nu, \Omega, t) \simeq \sum_{l=0}^N \sum_{m=-l}^l E^{ml}(x, \nu, t) Y_{ml}(\Omega)$$

$$E^{lm}(x, \nu, t) = \int_{4\pi} I(x, \nu, \Omega, t) Y^{lm}(\Omega) d\Omega$$

- Hyperbolic system ($v \leq c$)
- Rotationally invariant (no ray-effects as in S_N)
- Less memory ($P_{N-1} \sim S_N$)

Oscillations in P_N



(a) Beam: $I(\mu^c \varphi) = \delta(\mu - \frac{1}{2})$.

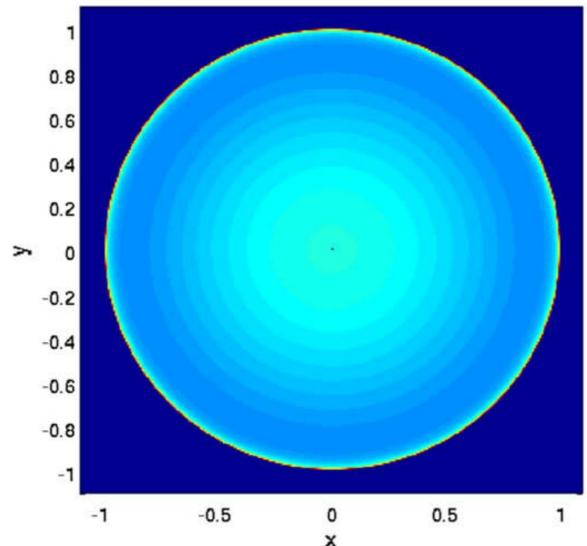
Filtered P_N

McClarren & Hauck 2010

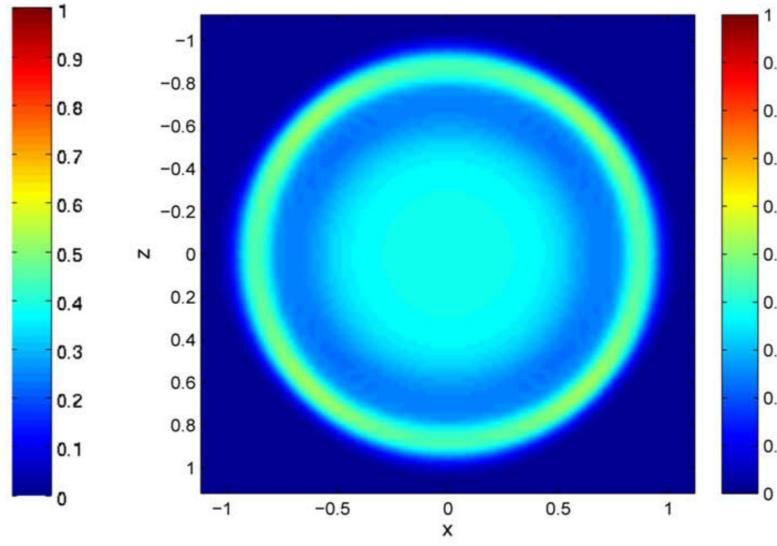
- Use filters to remove oscillations
- Preserves rotational invariance
- Converges to the transport solution
- Efficient and accurate

Filtered P_N : 2D line problem

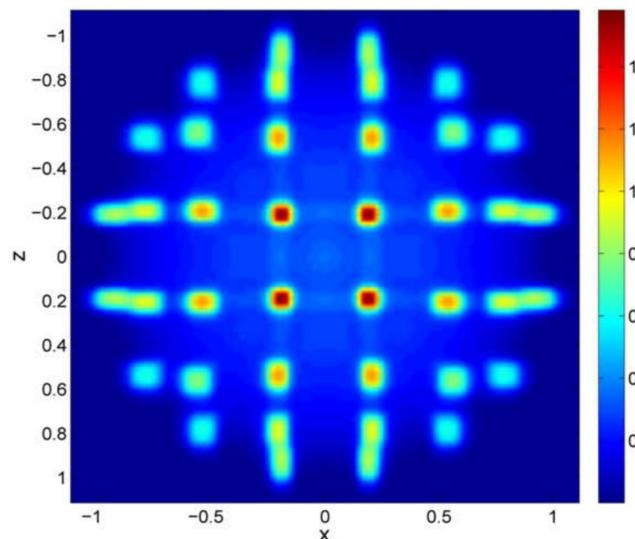
McClarren & Hauck 2010



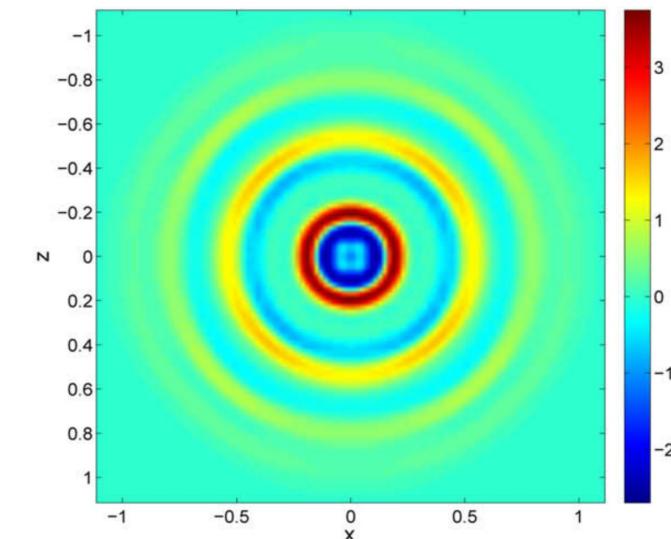
(a) Transport



(b) FP_7



(c) S_8



(d) P_7

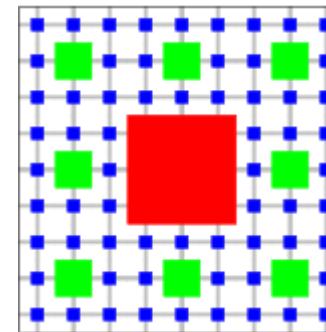
Filtered P_N

Open issues

- Other filters?
- What about 3D?
- Filtering as a continuum operation

Charon code

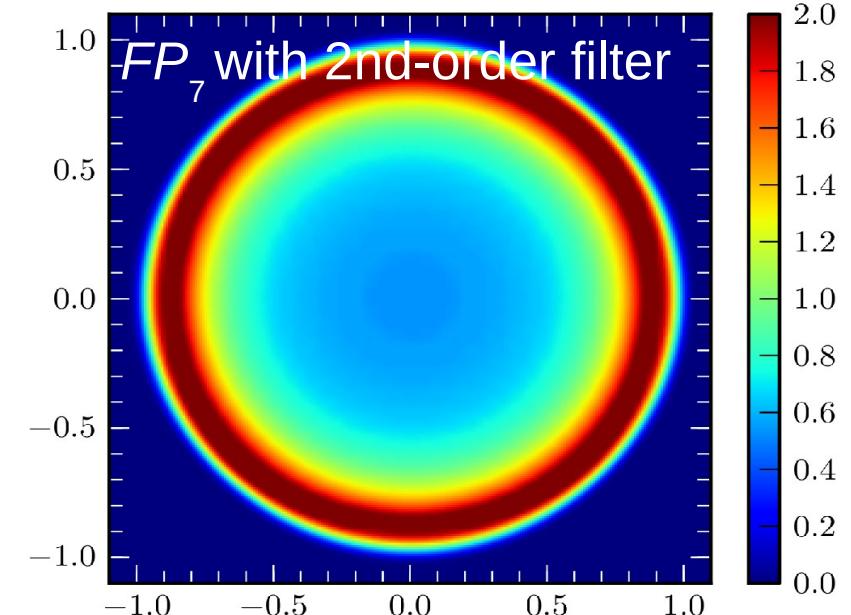
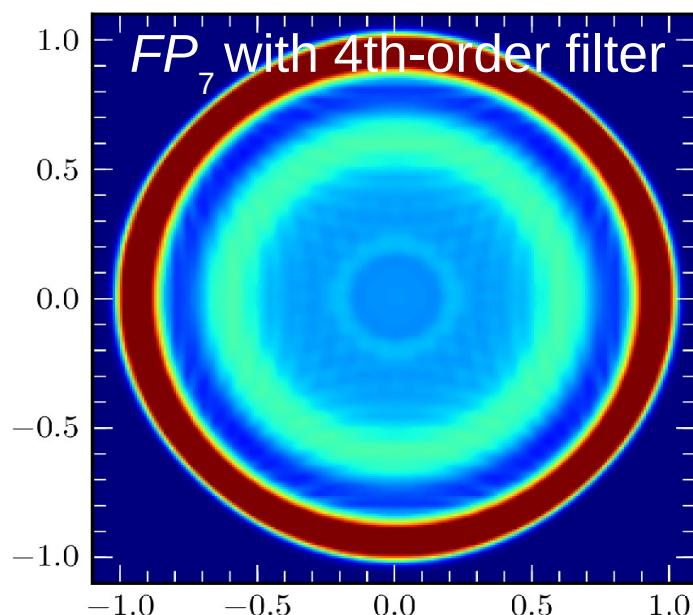
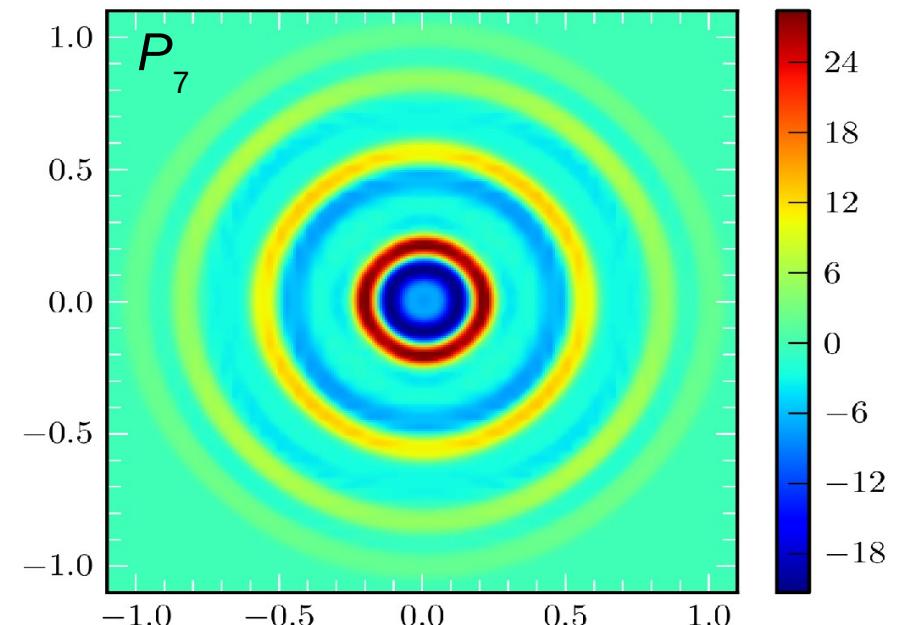
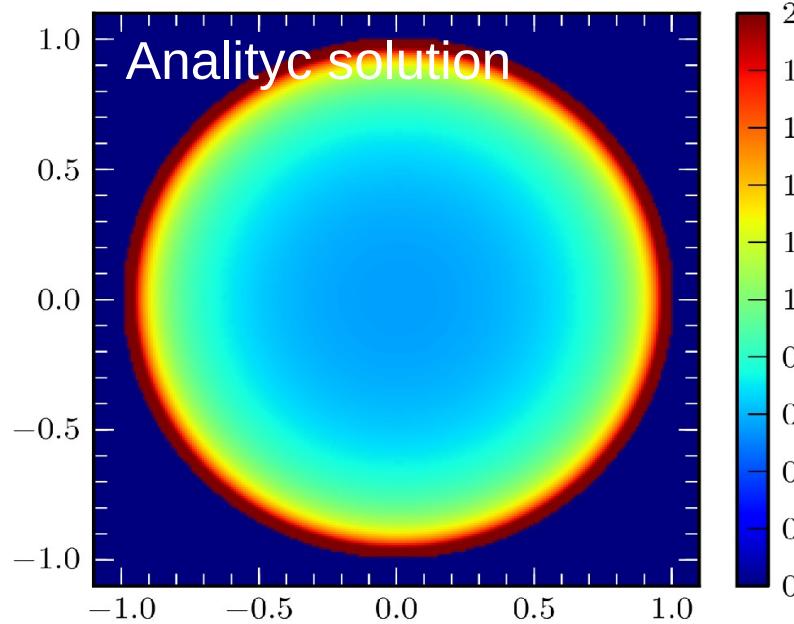
- 3D (using Cactus and Carpet)
- Space discretization: AP DG scheme
(Mcclarren & Lowrie '08)
- Semi-implicit time integration
(McClarren+ '06)
- 2nd- and 4th-order filters
(continuum limit)
- Special relativity
- No velocity dependence



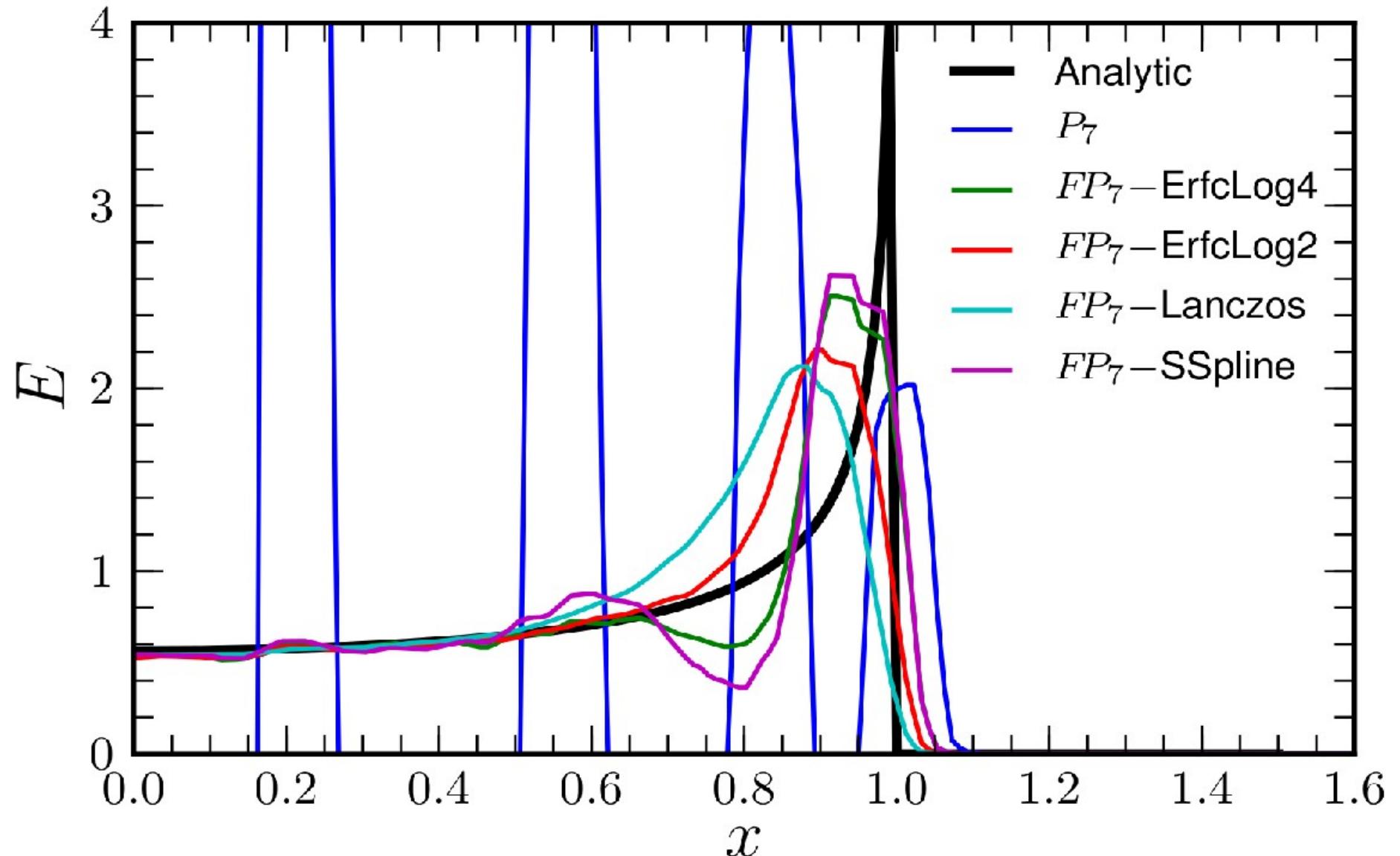
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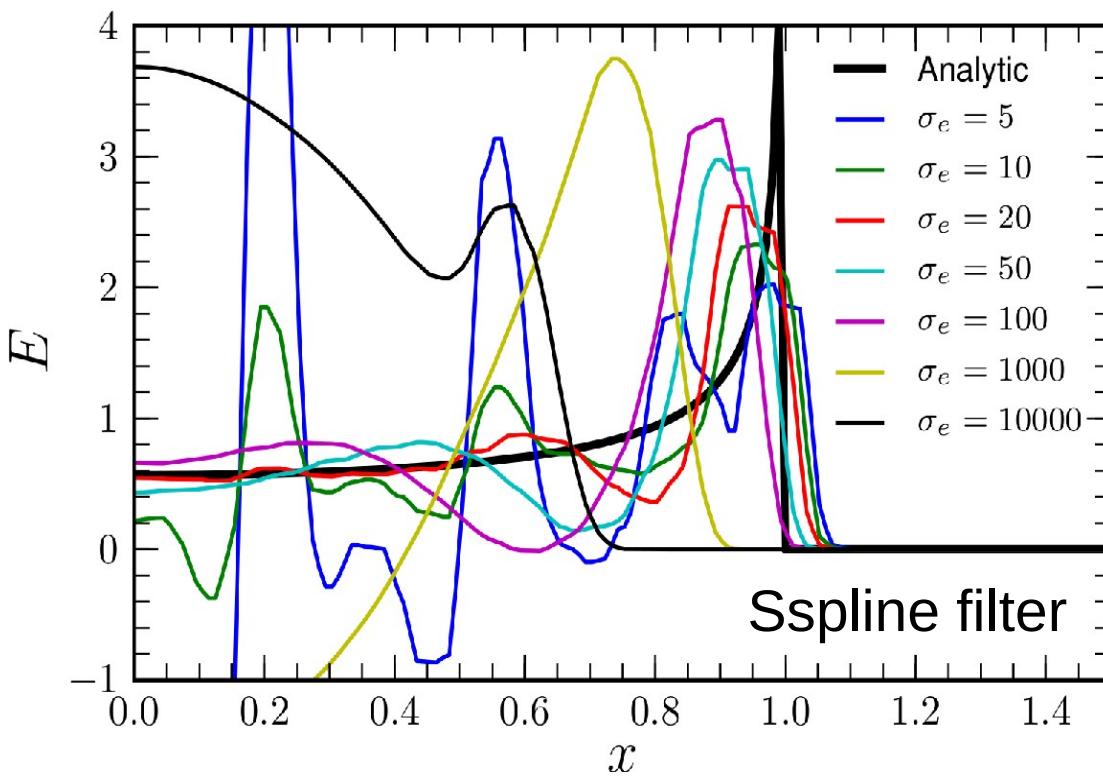
Line problem



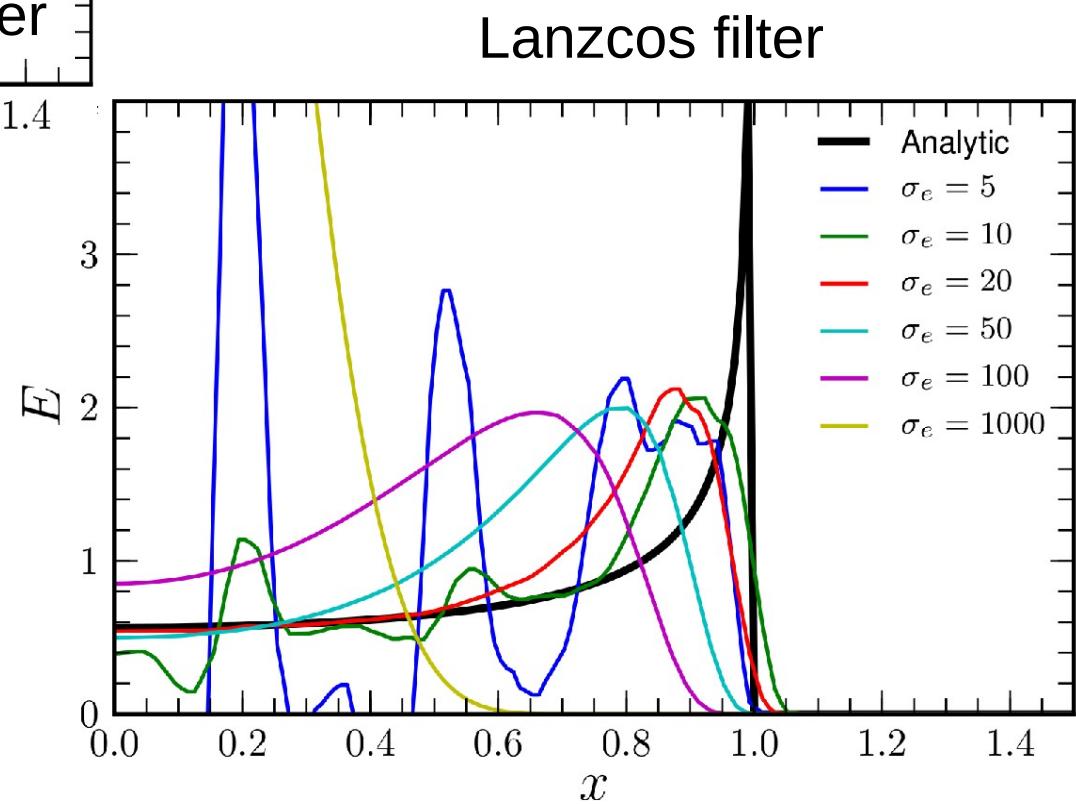
Line problem



Line problem

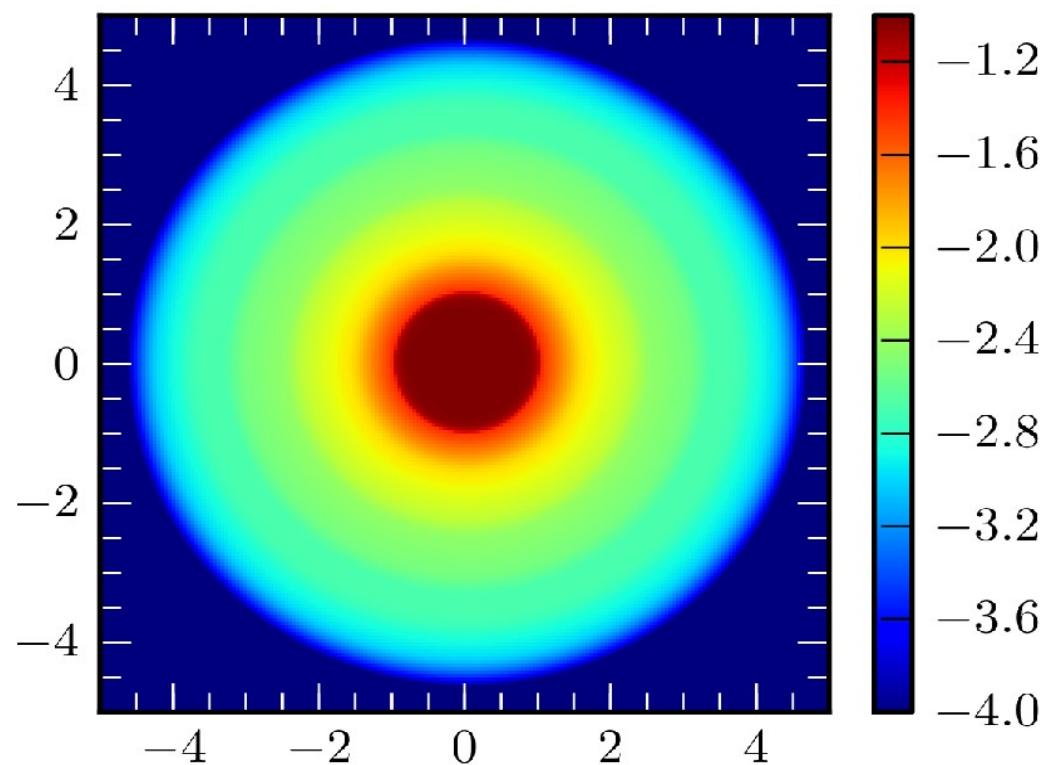
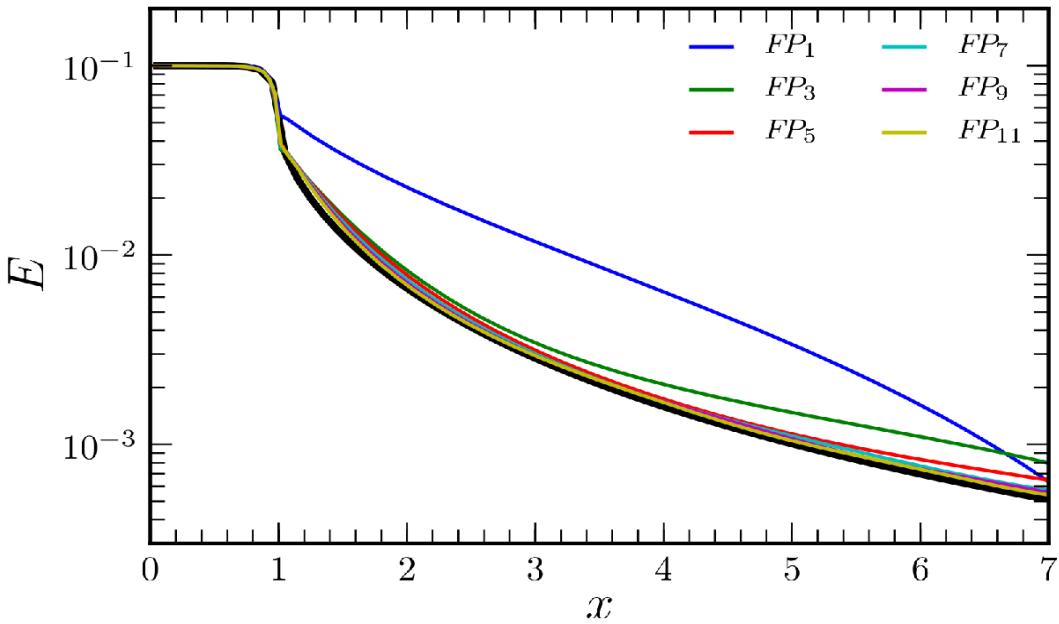


Spline filter



Lanczos filter

Homogeneous sphere test



Conclusion

- 3D filtered spherical harmonics
- Continuum filter formulation
- 2nd-order filters are better
- Overall, FP_N is a promising approach to 3D radiation transport

Supplemental Material

Velocity-dependent discrete-diffusion: operator splitting

$$\frac{1}{c} \frac{dJ_0}{dt} + \frac{\partial H_0}{\partial x} = \kappa_0 (B - J_0)$$

$$\frac{1}{c} \frac{dJ_0}{dt} + \frac{\varepsilon_0}{3c} \frac{\partial J_0}{\partial \varepsilon_0} \frac{D \ln \rho}{Dt} = 0$$

$$\frac{1}{c} \frac{dJ_0}{dt} + \frac{v}{c} \frac{\partial J_0}{\partial r} + \frac{J_0}{c} \frac{\partial v}{\partial r} = 0$$

Implicit Monte Carlo for Photons

[Fleck & Cummings '71]

$$\frac{1}{c} \frac{\partial I}{\partial t} + \mathbf{n} \cdot \nabla I = \kappa(B - I)$$



$$\frac{1}{c} \frac{\partial I}{\partial t} + \mathbf{n} \cdot \nabla I = \kappa_{ea,n}(B_n - I)$$

$$+ \chi_n \int \int \kappa_{es,n} I d\Omega d\varepsilon - \kappa_{es,n} I$$

Much larger timesteps!