

The non-linear Glasma

Soeren Schlichting
J.Berges

“Gauge field dynamics in and out of equilibrium”
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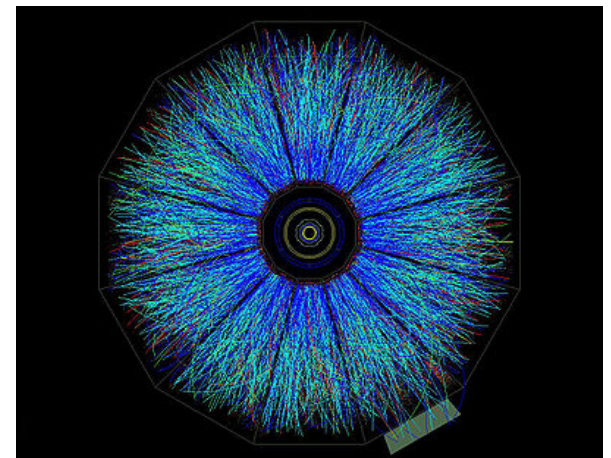
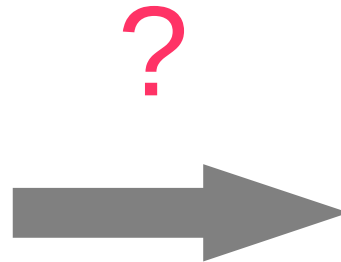
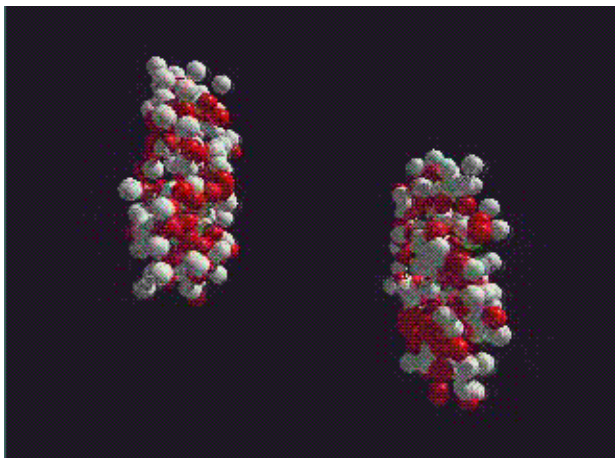


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Motivation

- understand thermalization and related questions from first principles
- use ab-initio approach to relativistic heavy-ion collisions
- weak coupling and high energies

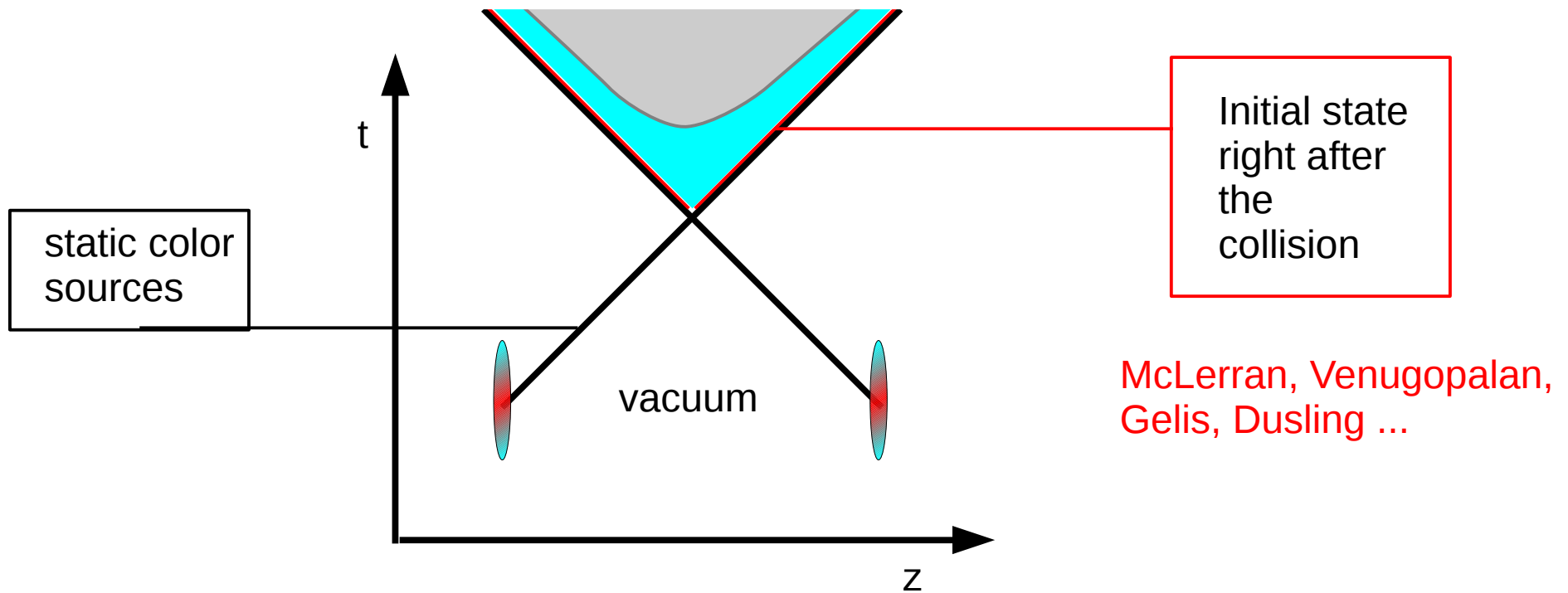


Outline

- Introduction to CGC framework and the Glasma
- Plasma instabilities and non-linear dynamics
- Summary & Outlook

CGC and the Glasma

- eikonal approximation – problem becomes particle production in the presence of strong sources



McLerran, Venugopalan, Gelis, Dusling ...

$$J_a^\mu(t, x_\perp, z) = \delta^{\mu+} \varrho_{(1)}^a(x_\perp) \delta(x^-) + \delta^{\mu-} \varrho_{(2)}^a(x_\perp) \delta(x^+)$$

CGC and the Glasma

Initial value problem in QFT

$$J_a^\mu(\mathbf{x}) \sim 1/g \quad \Rightarrow \quad A_\mu^a(\mathbf{x}) \sim 1/g \quad F_{\mu\nu}^{ab}(x, y) = \frac{1}{2} \langle \{ \hat{A}_\mu^a(x), \hat{A}_\nu^b(y) \} \rangle - \mathcal{A}_\mu^a(x) \mathcal{A}_\nu^b(y)$$

$$\rho_{\mu\nu}^{ab}(x, y) = i \langle [\hat{A}_\mu^a(x), \hat{A}_\nu^b(y)] \rangle \sim 1 \quad \sim 1 \text{ (initially)}$$

unstable in forward lightcone

Quantum evolution equations:

$$\frac{\delta S[J, A]}{\delta A_\mu^a(x)} = -J_a^\mu(x) + \text{loop corrections}$$

$$iG_0^{-1}[x; \mathcal{A}] \rho(x, y) = 0$$

$$iG_0^{-1}[x; \mathcal{A}] F(x, y) = 0 \quad + \text{loop corrections}$$

Closed set of coupled
integro-differential
equations

Glasma and Fluctuations

Weak coupling, small fluctuations

$$\frac{\delta S[J, A]}{\delta A_\mu^a(x)} = -J_a^\mu(x) + \text{loop corrections}$$

=> recovers classical field solution
(McLerran, Venugopalan, Fukushima, Gelis, Lappi,...)

$$\begin{aligned} iG_0^{-1}[x; \mathcal{A}] \rho(x, y) &= 0 \\ iG_0^{-1}[x; \mathcal{A}] F(x, y) &= 0 \end{aligned} + \text{loop corrections}$$

equivalent to linearized classical evolution equations
=> spectrum of fluctuations right after the collision

Dusling, Gelis, Venugopalan (2011)

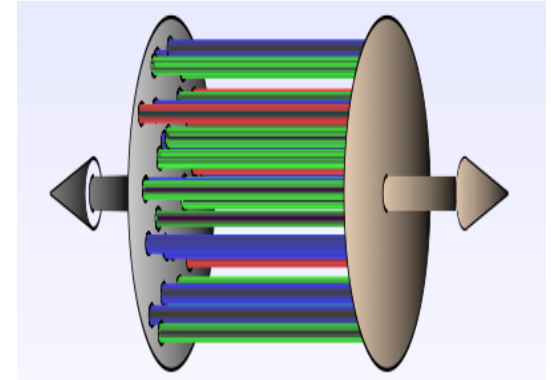
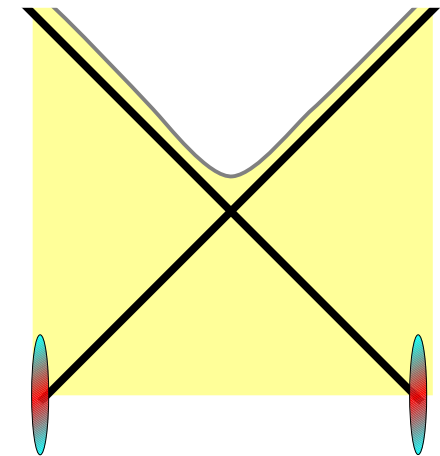


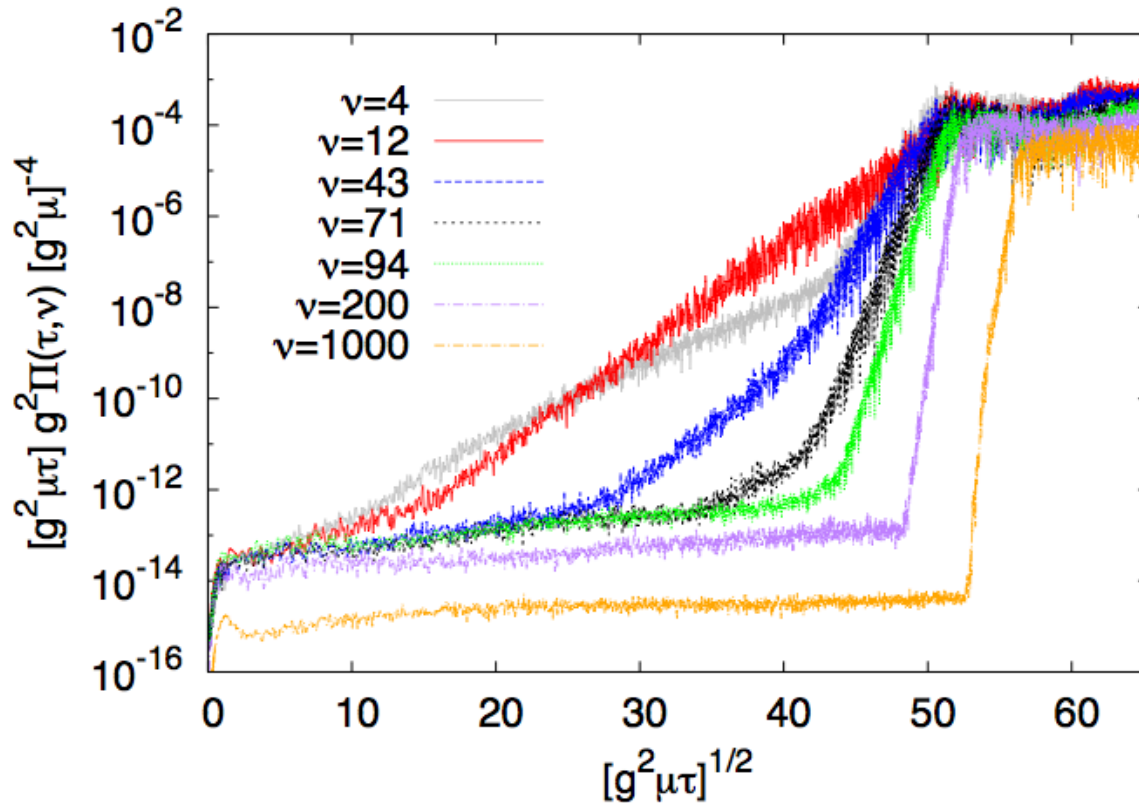
Fig. by F. Gelis



Plasma instabilities

Consider boost non-invariant fluctuations
=> Grow exponentially in the forward light-cone

Initially small fluctuations become *large*

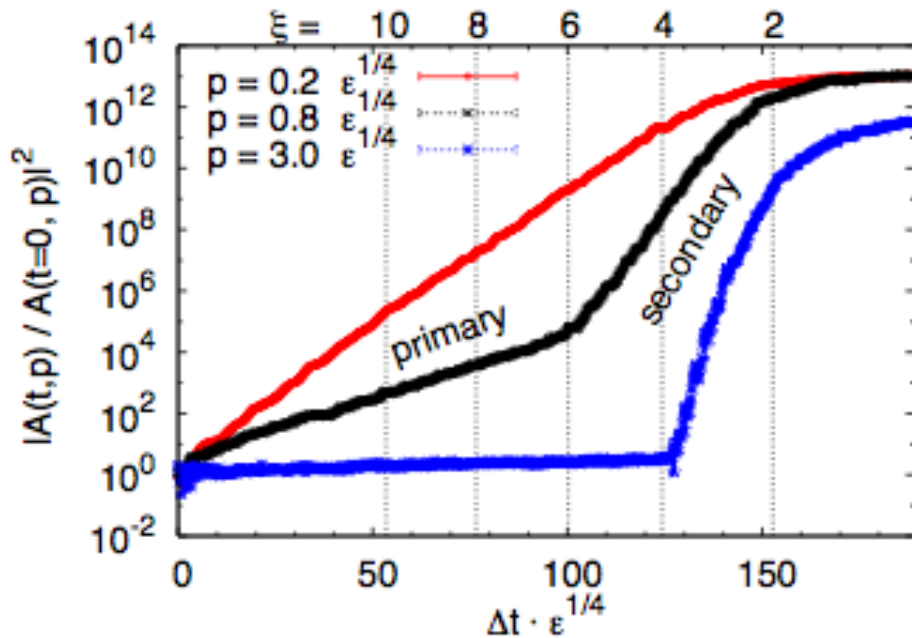


Classical statistical lattice simulation

- **CGC initial conditions** (MV model)
- **simplified fluctuations**

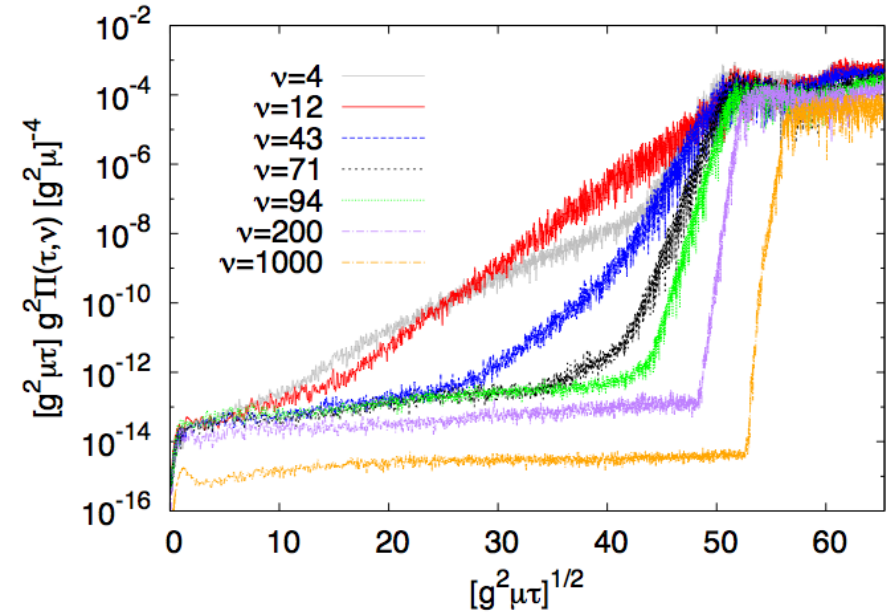
Romastschke, Venugopalan (2006); Fukushima, Gelis (2011); SS, Berges (in preparation)

Non-linear amplification



SU(2) – fixed box

Berges, Scheffler, Sexty (2008)



SU(2) – CGC expanding

Berges, SS (2012)

Non-linear amplification

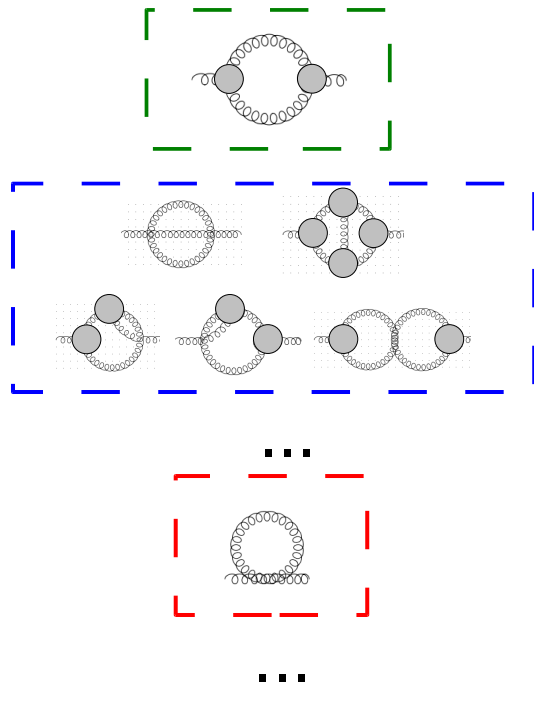
Weak coupling, fluctuations grow with time

$$\frac{\delta S[J, A]}{\delta A_\mu^a(x)} = -J_a^\mu(x) + \text{loop corrections}$$

$$\begin{aligned} iG_0^{-1}[x; \mathcal{A}] \rho(x, y) &= 0 \\ iG_0^{-1}[x; \mathcal{A}] F(x, y) &= 0 \end{aligned} + \text{loop corrections}$$

Q: What are the relevant loop corrections?

Power counting



$$F \sim 1/g$$

Take into account enhancement due to large fluctuations

$$F \sim 1/g^{4/3}$$

Can distinguish different dynamical regimes where higher order corrections are suppressed by at least a fractional power of the coupling constant

$$F \sim 1/g^2$$

Most important diagrams contained in classical statistical lattice gauge theory

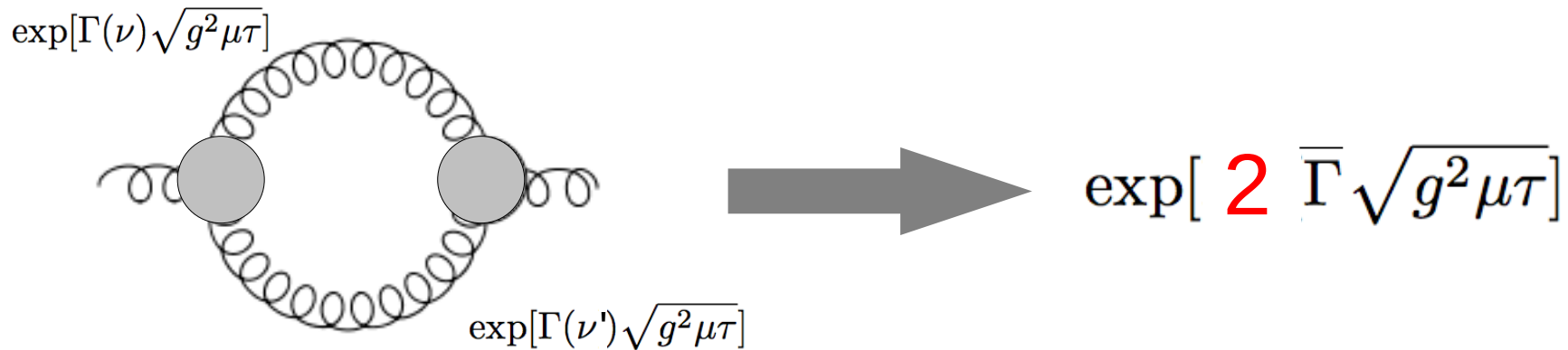
power counting

$$g^2 F^2, g^4 F^3, \dots, g^2 F, \dots$$

Non-linear amplification

dominated by *soft unstable* modes

secondary instabilities

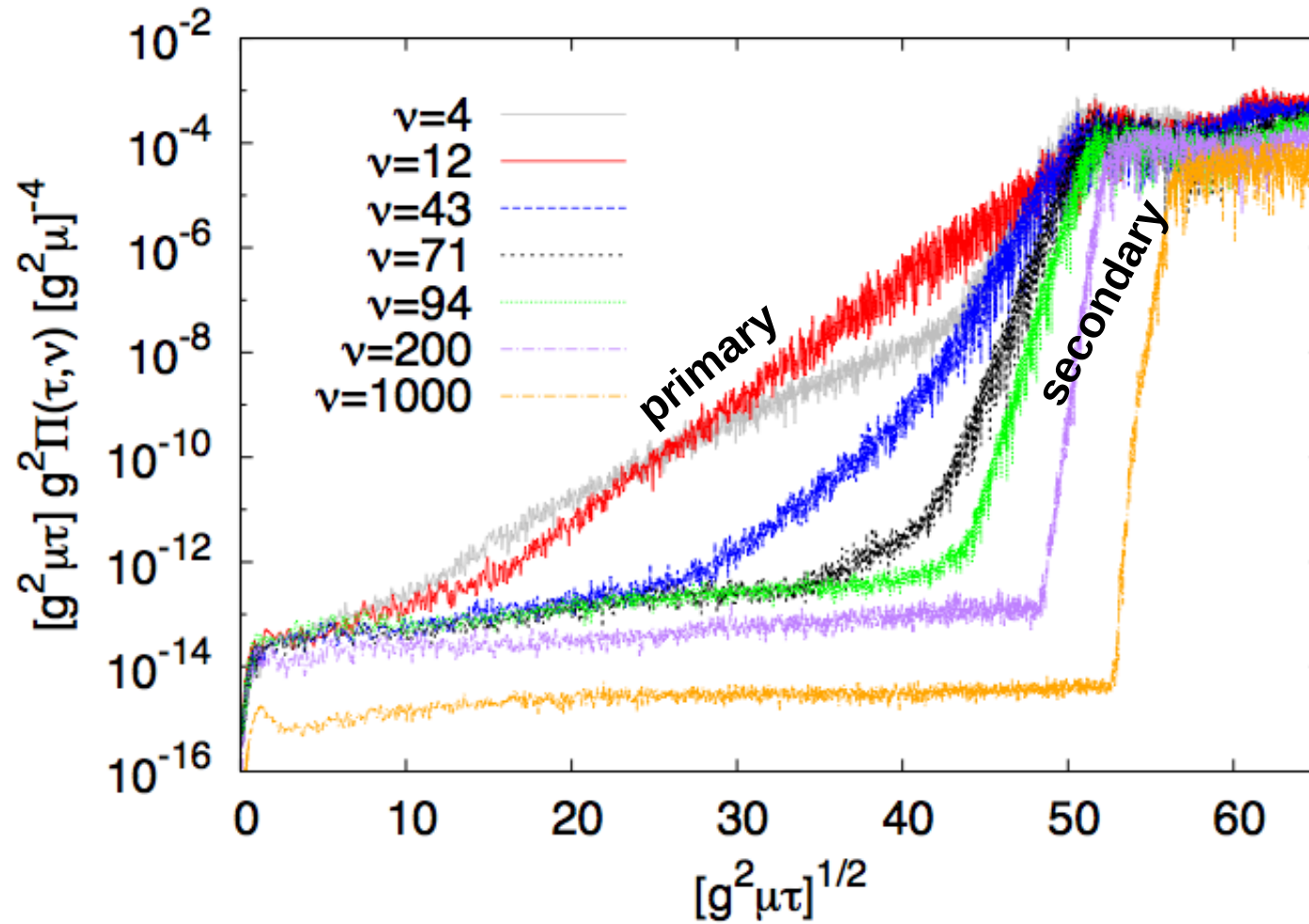


relevant time scale **parametrically**

$$\sqrt{g^2\mu\tau_{\text{Sec}}} \stackrel{g \ll 1}{\sim} \frac{1}{2\Gamma_0} \ln g^{-2}$$

Can be calculated explicitly when primary instabilities are described analytically (e.g. Berges, Serreau; Berges, Boguslavski, SS (scalar field theory))

Non-linear amplification

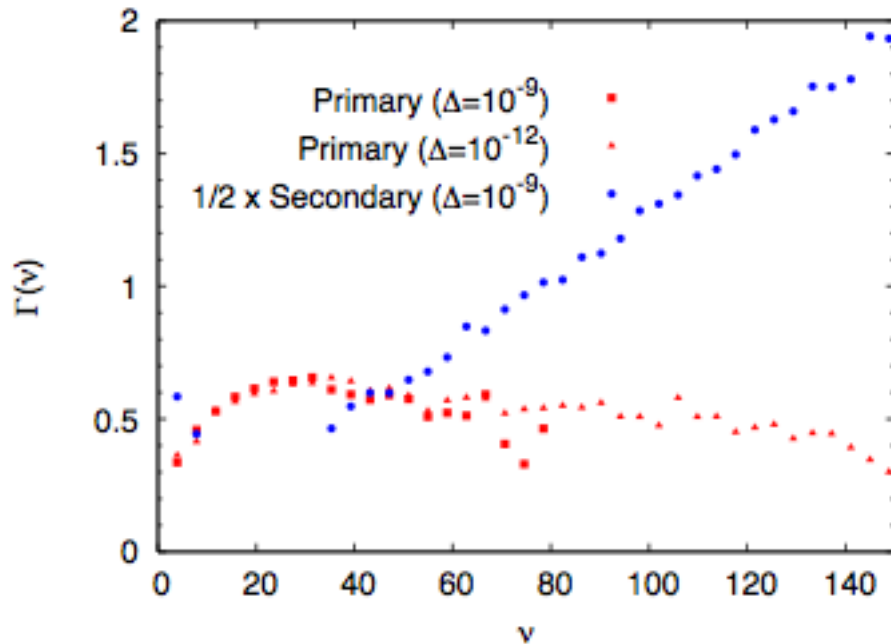


MV model, spectrum of fluctuations simplified

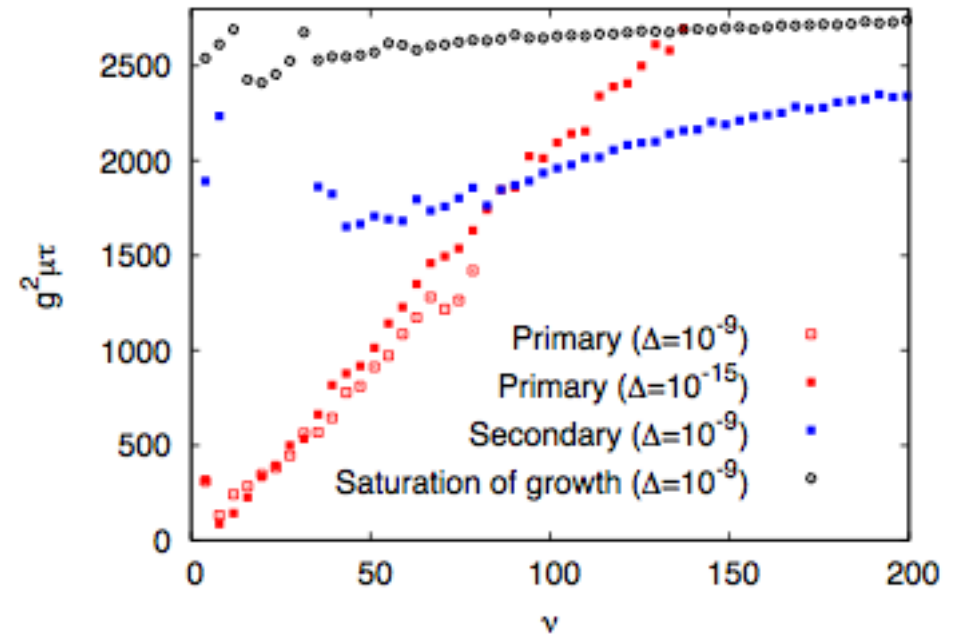
SS, Berges (in preparation)

Growth rates & time scales

Growth rates

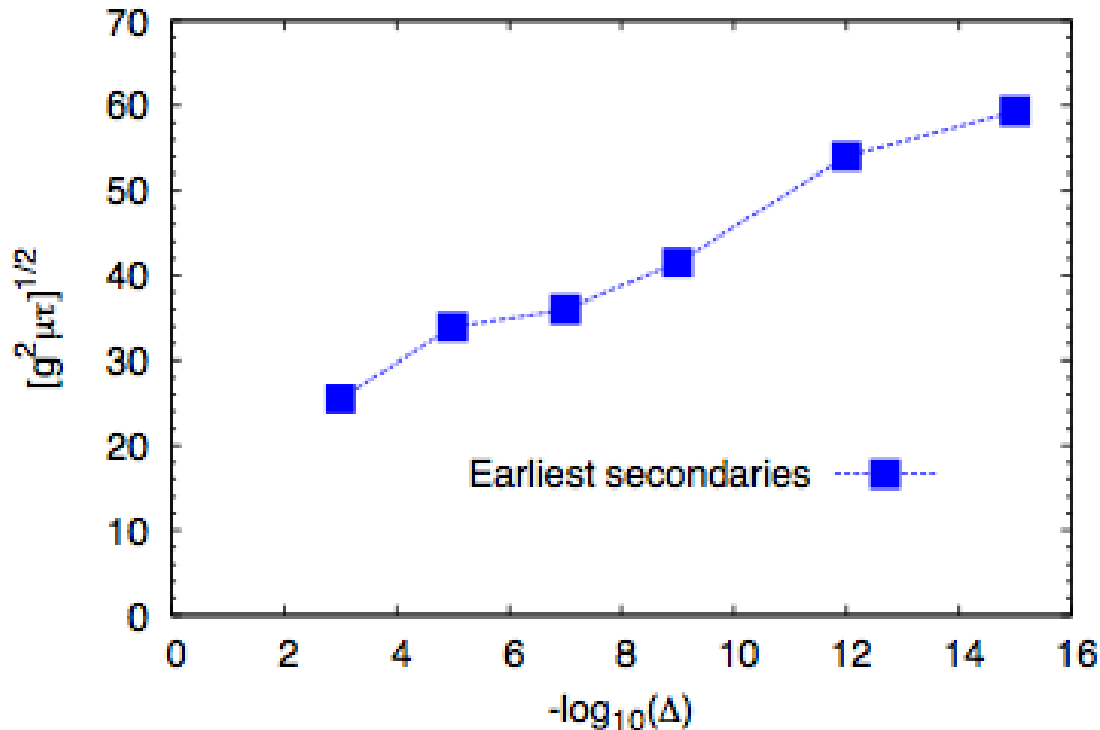


Set-in times



SS, Berges (in preparation)

Coupling dependence



Confirms logarithmic behavior

$$\sqrt{g^2 \mu \tau_{\text{Sec}}} \stackrel{g \ll 1}{\sim} \frac{1}{2\Gamma_0} \ln g^{-2}$$

Subleading corrections from:

- delayed set-in of primary instability
- Spectrum of initial fluctuations

Δ^2 : size of initial fluctuations $\sim g^2$

SS, Berges (in preparation)

Summary

Different dynamical regimes of a system undergoing instabilities:

- linear instability regime
- non-linear amplification regime
- saturation of growth

Non-linear effects occur before saturation and are dominant for high-momentum modes

Relevant time scale depends only logarithmically on g^2

$$\sqrt{g^2 \mu \tau_{\text{Sec}}} \stackrel{g \ll 1}{\sim} \frac{1}{2\Gamma_0} \ln g^{-2} ,$$