

Ab initio reaction theory for light nuclei

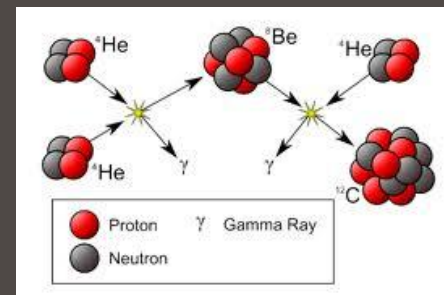
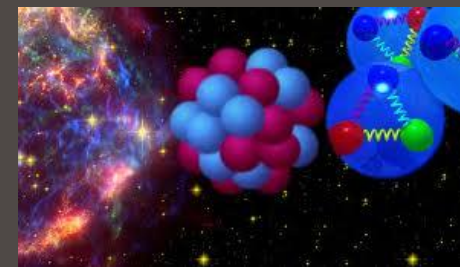
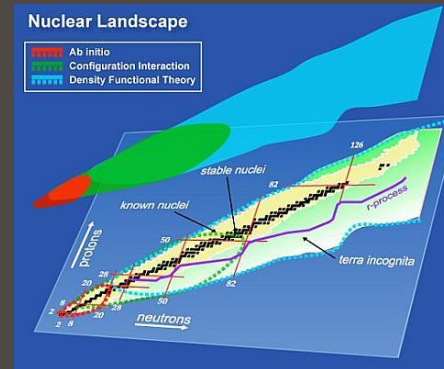
7th ANL/INT/JINA/MSU annual FRIB workshop

INT Program INT-11-2d: Interfaces between structure and reactions
 for rare isotopes and nuclear astrophysics

Seattle, August 8, 2011

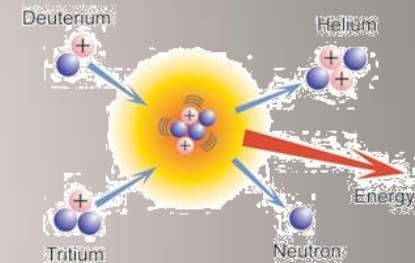
Petr Navratil | TRIUMF

Collaborators: Sofia Quaglioni (LLNL),
 Robert Roth (TU Darmstadt), W. Horiuchi (RIKEN),
 E. Jurgenson (LLNL), M. Kruse (U of A), S. Baroni (TRIUMF)

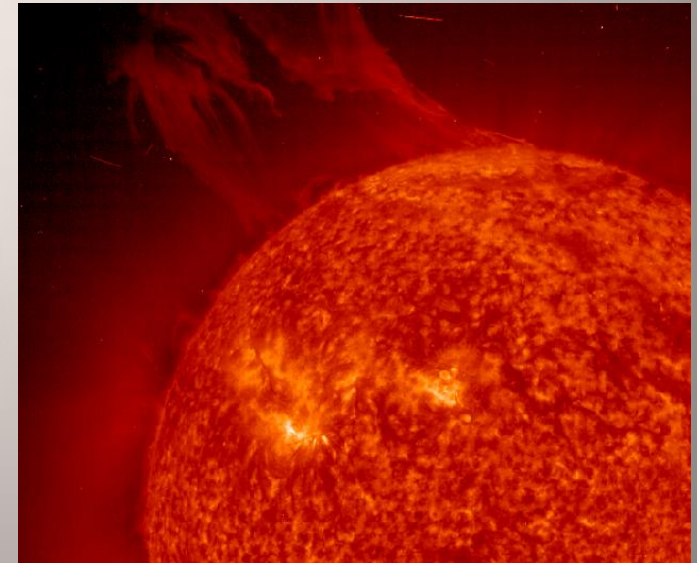
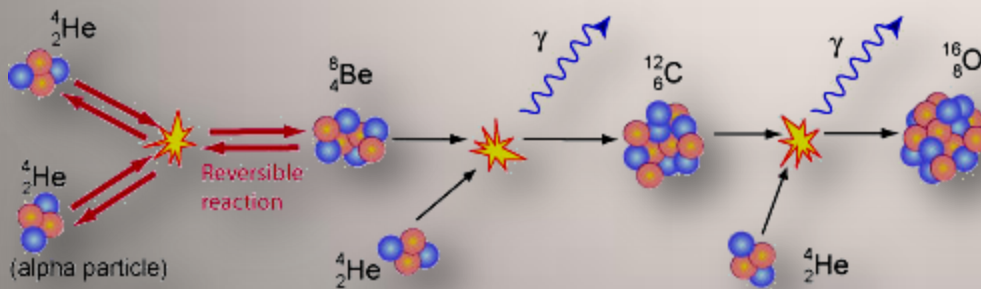


Light nuclei from first principles

- **Goal:** Predictive theory of structure and reactions of light nuclei
- Needed for
 - Physics of exotic nuclei, tests of fundamental symmetries
 - **Understanding of nuclear reactions important for astrophysics**
 - Understanding of reactions important for energy generation
- **From first principles or *ab initio*:**
- ✓ Nuclei as systems of nucleons interacting by nucleon-nucleon (and three-nucleon) forces that describe accurately nucleon-nucleon (and three-nucleon) systems



Understanding our Sun

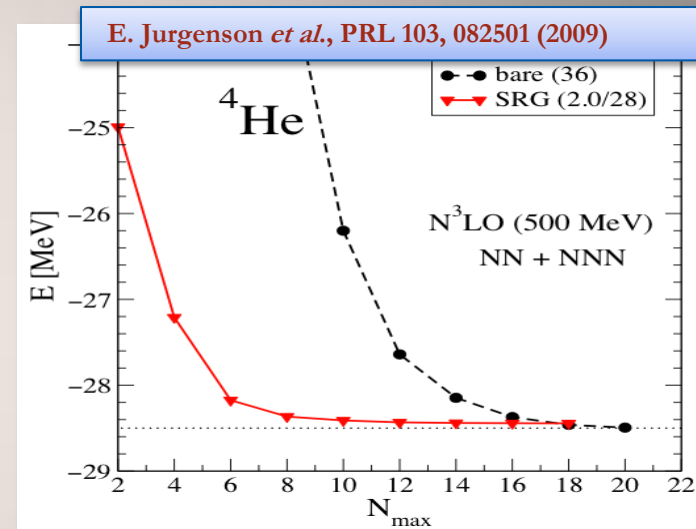


Our many-body technique:

- **Combine** the *ab initio* no-core shell model (NCSM) with the resonating group method (RGM)

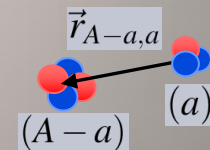
- **The NCSM:** An approach to the solution of the A -nucleon bound-state problem

- Accurate nuclear Hamiltonian
- Finite harmonic oscillator (HO) basis
 - Complete N_{max} model space
- Effective interaction due to the model space truncation
 - Similarity-Renormalization-Group evolved NN(+NNN) potential
- Short & medium range correlations
- No continuum



- **The RGM:** A microscopic approach to the A -nucleon scattering of clusters

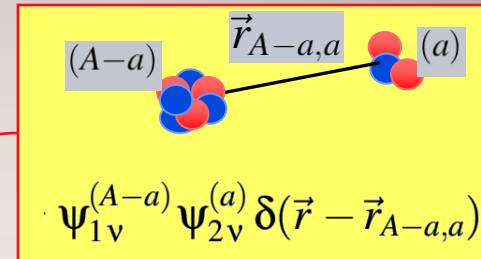
- Nuclear Hamiltonian may be simplistic
- Cluster wave functions may be simplified and inconsistent with the nuclear Hamiltonian
- Long range correlations, relative motion of clusters



Ab initio NCSM/RGM: Combines the best of both approaches
Accurate nuclear Hamiltonian, **consistent** cluster wave functions
Correct asymptotic expansion, **Pauli principle** and **translational invariance**

The *ab initio* NCSM/RGM in a snapshot

- Ansatz: $\Psi^{(A)} = \sum_{\mathbf{v}} \int d\vec{r} \phi_{\mathbf{v}}(\vec{r}) \hat{\mathcal{A}} \Phi_{\mathbf{v}\vec{r}}^{(A-a,a)}$



eigenstates of $H_{(A-a)}$ and $H_{(a)}$ in the *ab initio* NCSM basis

- Many-body Schrödinger equation:

$$H\Psi^{(A)} = E\Psi^{(A)}$$

$$T_{\text{rel}}(r) + \mathcal{V}_{\text{rel}} + \bar{V}_{\text{Coul}}(r) + H_{(A-a)} + H_{(a)}$$

$$\sum_{\mathbf{v}} \int d\vec{r} \left[\mathcal{H}_{\mu\mathbf{v}}^{(A-a,a)}(\vec{r}', \vec{r}) - E \mathcal{N}_{\mu\mathbf{v}}^{(A-a,a)}(\vec{r}', \vec{r}) \right] \phi_{\mathbf{v}}(\vec{r}) = 0$$

$$\langle \Phi_{\mu\vec{r}'}^{(A-a,a)} | \hat{\mathcal{A}} H \hat{\mathcal{A}} | \Phi_{\mathbf{v}\vec{r}}^{(A-a,a)} \rangle$$

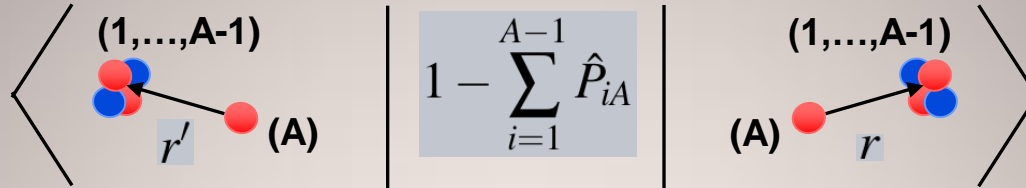
Hamiltonian kernel

$$\langle \Phi_{\mu\vec{r}'}^{(A-a,a)} | \hat{\mathcal{A}}^2 | \Phi_{\mathbf{v}\vec{r}}^{(A-a,a)} \rangle$$

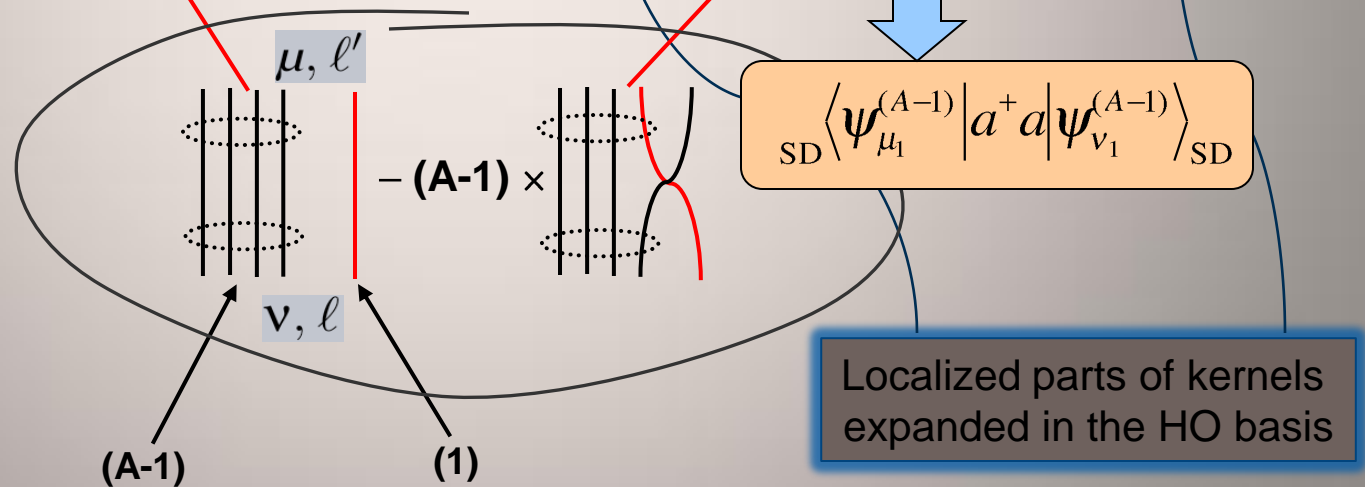
Norm kernel

realistic nuclear Hamiltonian

Single-nucleon projectile: the norm kernel



$$\mathcal{N}_{\mu\ell, \nu\ell}^{(A-1,1)}(r', r) = \delta_{\mu\nu} \delta_{\ell\ell} \frac{\delta(r' - r)}{r'r} - (A-1) \sum_{n'n} R_{n'\ell}(r') \langle \Phi_{\mu n'\ell'}^{(A-1,1)JT} | P_{A,A-1} | \Phi_{\nu n\ell}^{(A-1,1)JT} \rangle R_{n\ell}(r)$$



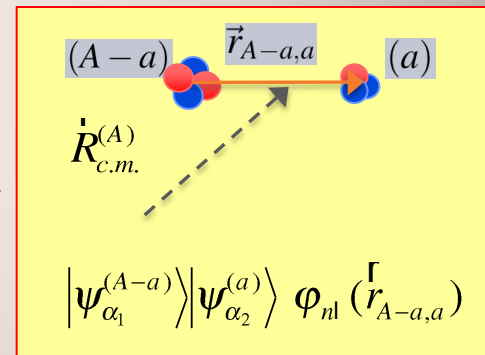
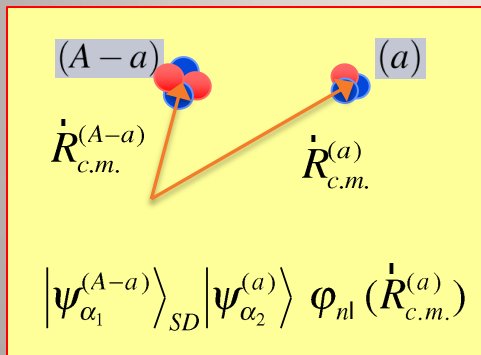
Matrix elements of translationally invariant operators

- Translational invariance is preserved (exactly!) also with SD cluster basis

$${}_{SD} \left\langle \Phi_{f_{SD}}^{(A-a',a')} \left| \hat{O}_{t.i.} \right| \Phi_{i_{SD}}^{(A-a,a)} \right\rangle_{SD} = \sum_{i_R f_R} M_{i_{SD} f_{SD}, i_R f_R} \left\langle \Phi_{f_R}^{(A-a',a')} \left| \hat{O}_{t.i.} \right| \Phi_{i_R}^{(A-a,a)} \right\rangle$$

Calculate these

Interested in these



- Advantage: can use powerful second quantization techniques

$${}_{SD} \left\langle \Phi_{v'n'}^{(A-a',a')} \left| \hat{O}_{t.i.} \right| \Phi_{vn}^{(A-a,a)} \right\rangle_{SD} \propto {}_{SD} \left\langle \psi_{\alpha_1'}^{(A-a')} \left| a^+ a \right| \psi_{\alpha_1}^{(A-a)} \right\rangle_{SD}, \quad {}_{SD} \left\langle \psi_{\alpha_1'}^{(A-a')} \left| a^+ a^+ a a \right| \psi_{\alpha_1}^{(A-a)} \right\rangle_{SD}, \quad \text{L}$$

Solving the RGM equations

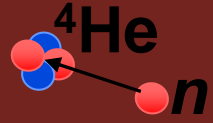
- The many-body problem has been reduced to a **two-body problem!**
 - **Macroscopic degrees of freedom:** nucleon clusters
 - **Unknowns:** relative wave function between the two clusters
- Non-local integral-differential coupled-channel equations:

$$\left[T_{rel}(r) + V_C(r) + E_{\alpha_1}^{(A-a)} + E_{\alpha_2}^{(a)} \right] u_{\nu}^{(A-a,a)}(r) + \sum_{a'\nu'} \int dr' r' W_{av,a'\nu'}(r,r') u_{\nu'}^{(A-a',a')}(r') = 0$$

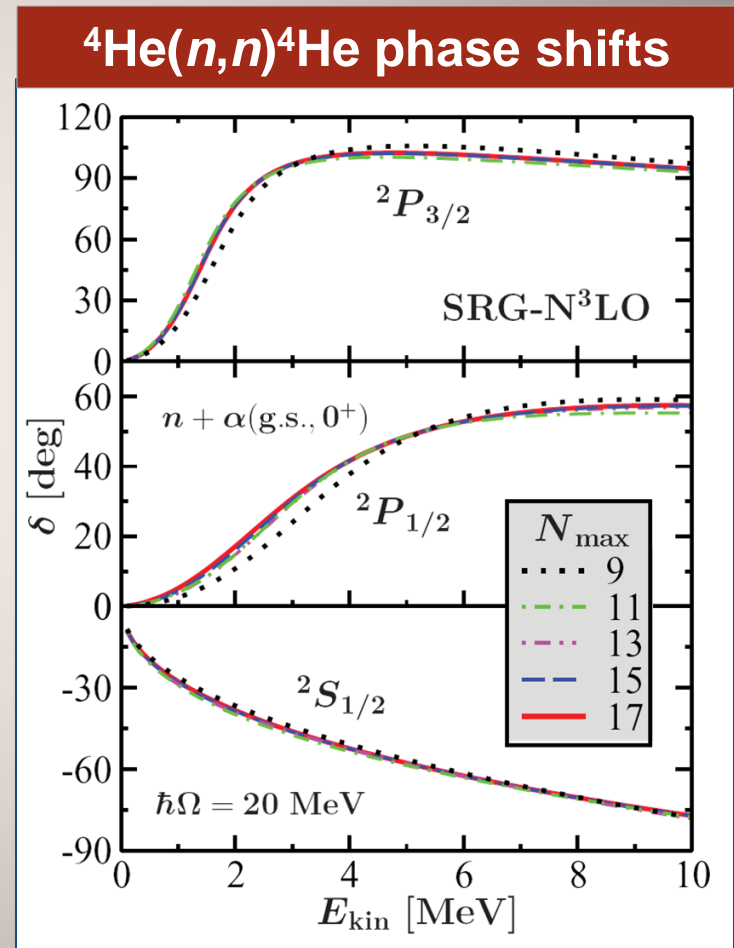
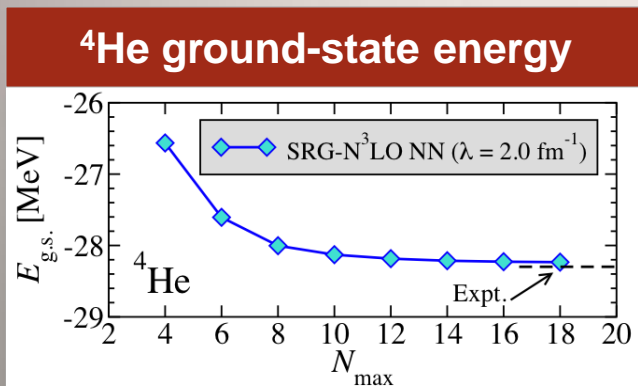
- Solve with R-matrix theory on Lagrange mesh imposing
 - **Bound state boundary conditions** → eigenenergy + eigenfunction
 - **Scattering state boundary conditions** → Scattering matrix
 - Phase shifts
 - Cross sections
 - ...

The R-matrix theory on Lagrange mesh is an elegant and powerful technique, particularly for calculations with non-local potentials

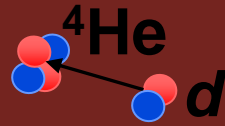
Convergence with respect to HO basis expansion



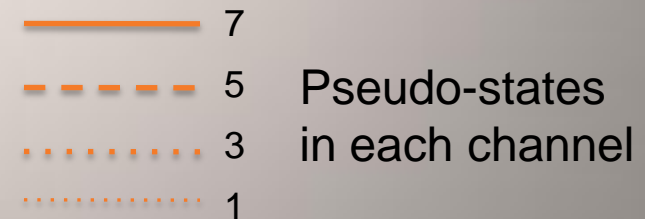
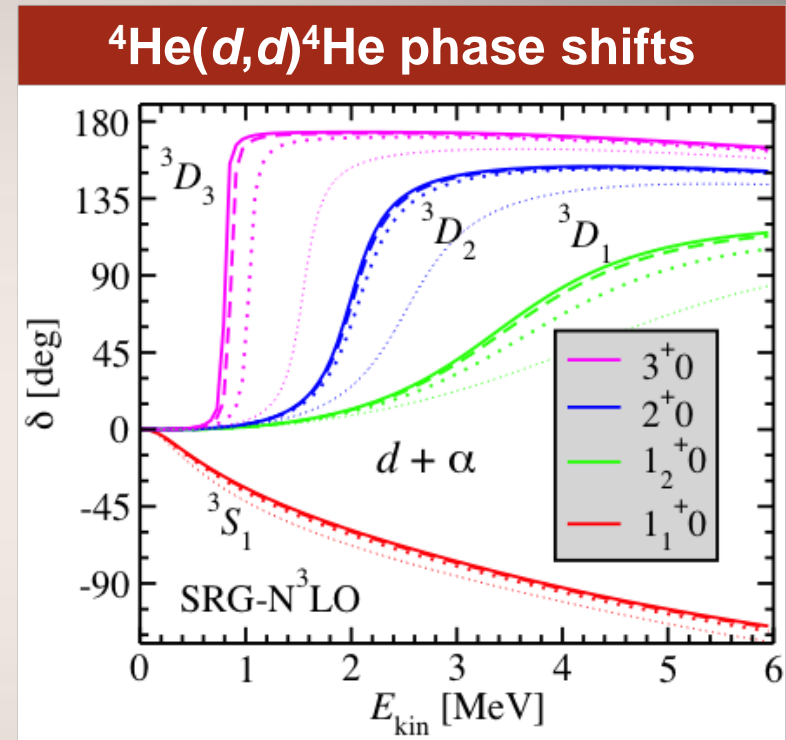
- Influenced by:
 - 1) Convergence of target and projectile wave functions
 - 2) Convergence of localized parts of the integration kernels
- Here:
 - $n + 4\text{He}(\text{g.s.}, 0^+)$ phase shifts
 - SRG-N³LO NN potential ($\lambda = 2 \text{ fm}^{-1}$)



Convergence with respect to RGM model space

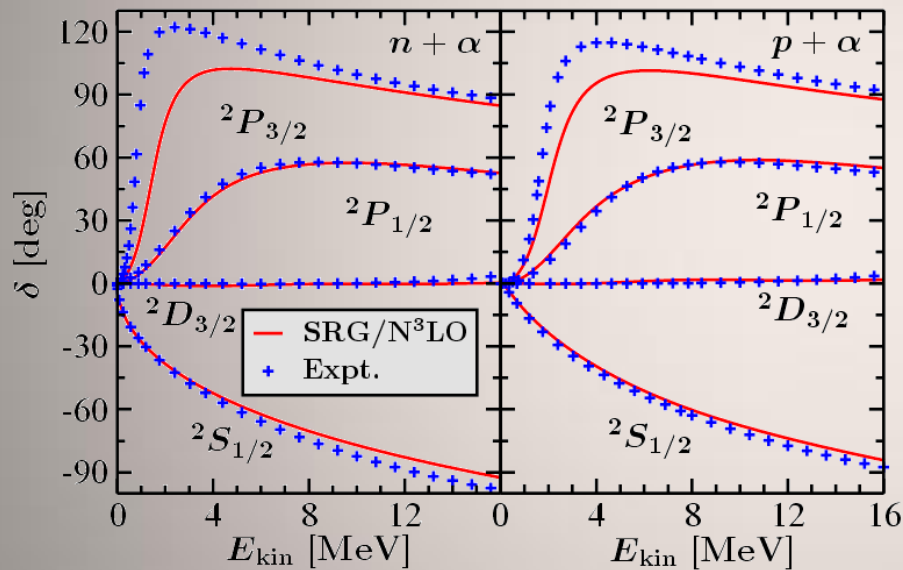
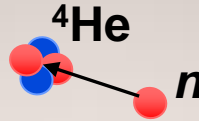


- NCSM/RGM describes binary reactions (below three-body breakup threshold)
- If projectile (or target) can be easily deformed or broken apart
 - Need to account for virtual breakup
 - Approximate treatment:
 - Include multiple excited (pseudo-) states of the clusters
 - Exact treatment:
 - 1) Inclusion of three-body clusters
 - 2) Solution of three-body scattering
- Here:
 - $d(\text{g.s.}, {}^3S_1\text{-}{}^3D_1, {}^3D_2, {}^3D_3\text{-}{}^3G_3) + {}^4\text{He}(\text{g.s.})$
 - SRG-N³LO NN potential ($\lambda = 1.5 \text{ fm}^{-1}$)

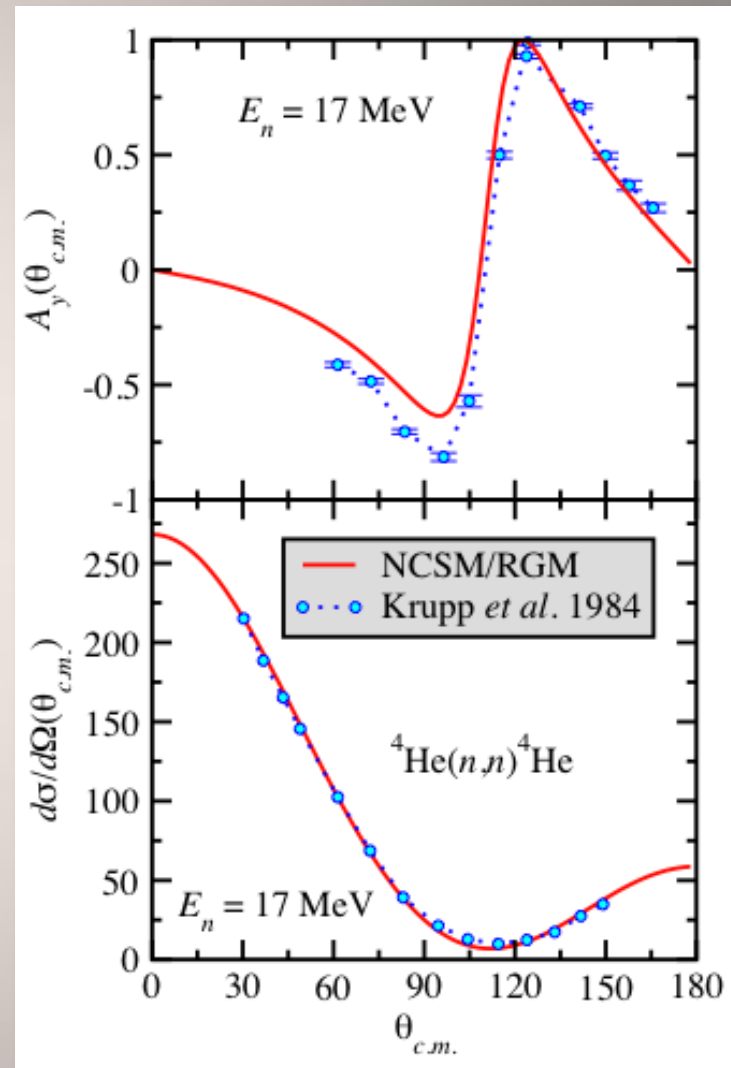


The best system to start with: $n+{}^4\text{He}$, $p+{}^4\text{He}$

- NCSM/RGM calculations with
 - $N + {}^4\text{He}(\text{g.s.}, 0^+0)$
 - SRG- $N^3\text{LO}$ NN potential with $\Lambda=2.02 \text{ fm}^{-1}$



- Differential cross section and analyzing power @ 17 MeV neutron energy
 - Polarized neutron experiment at Karlsruhe



NNN missing: Good agreement only for energies beyond low-lying $3/2^-$ resonance

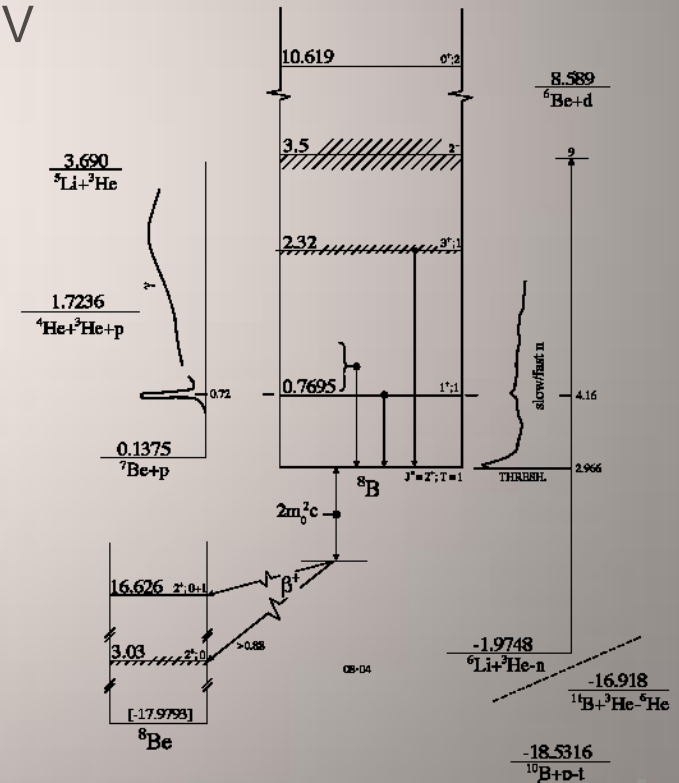
${}^7\text{Be}(p,\gamma){}^8\text{B}$ S-factor

- S_{17} one of the main inputs for understanding the solar neutrino flux
 - Needs to be known with a precision better than 9 %
- Current evaluation has uncertainty $\sim 10\%$
 - Theory needed for extrapolation to ~ 10 keV

$$S(E) = E\sigma(E) \exp[2\pi\eta(E)]$$

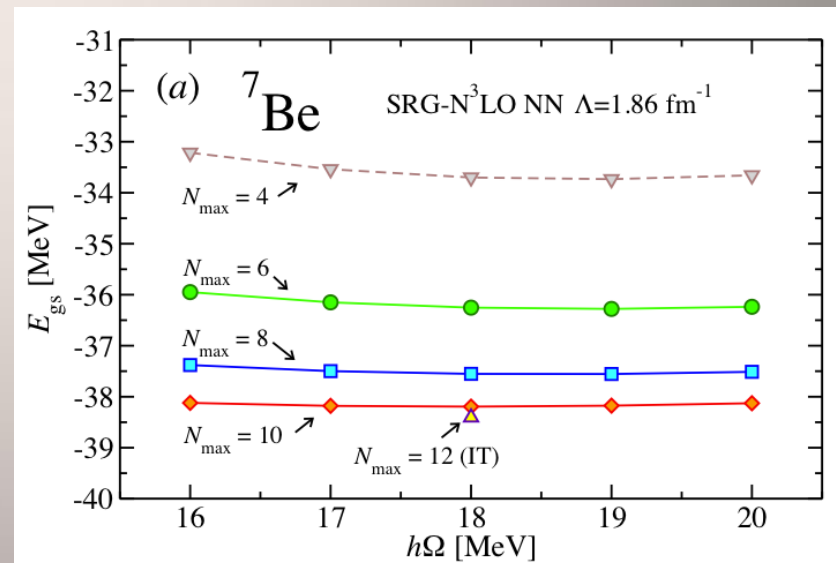
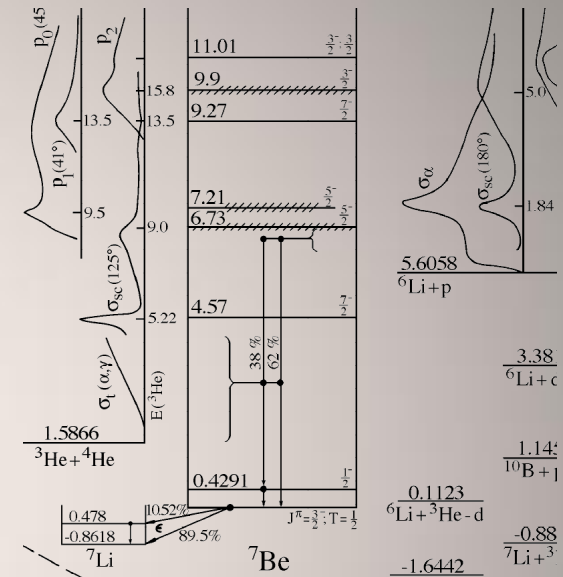
$$\eta(E) = Z_{A-a}Z_a e^2 / \hbar v_{A-a,a}$$

$$\langle {}^8\text{B}_{\text{g.s.}} | E1 | {}^7\text{Be}_{\text{g.s.}} + p \rangle$$



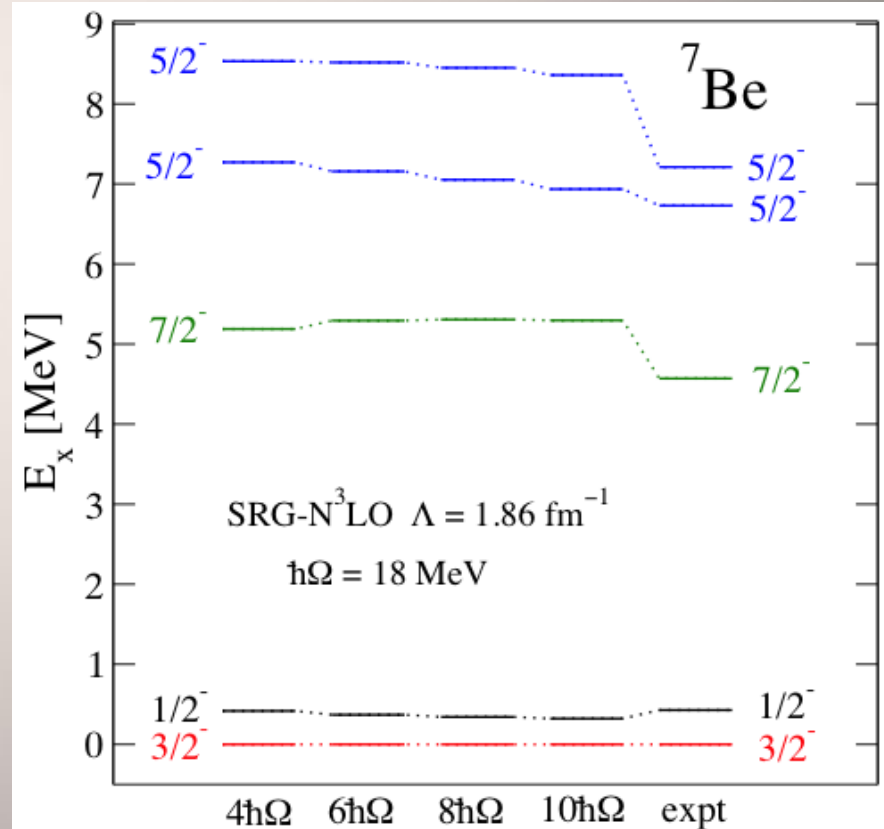
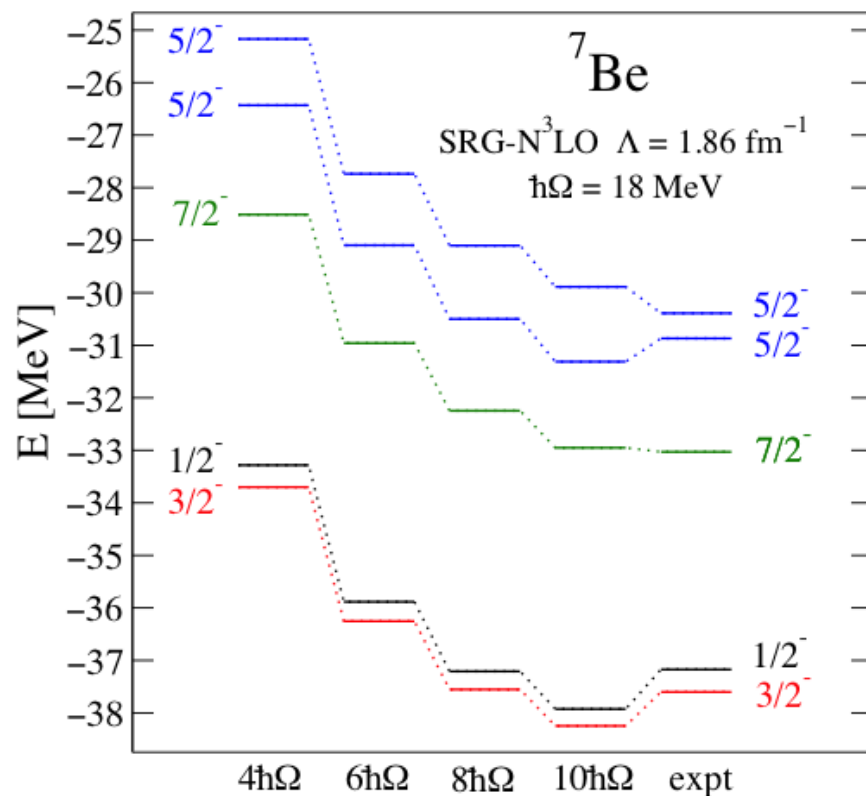
Input: NN interaction, ${}^7\text{Be}$ eigenstates

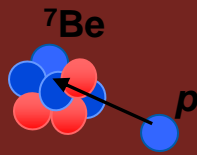
- Similarity-Renormalization-Group (SRG) evolved chiral $N^3\text{LO}$ NN interaction
 - Accurate
 - Soft: Evolution parameter Λ
 - Study dependence on Λ
- ${}^7\text{Be}$
 - NCSM up to $N_{\text{max}}=10$, Importance Truncated NCSM up to $N_{\text{max}}=14$
 - Variational calculation
 - optimal HO frequency from the ground-state minimum
 - For the selected NN potential with $\Lambda=1.86 \text{ fm}^{-1}$: $\hbar\Omega=18 \text{ MeV}$



Input: ${}^7\text{Be}$ eigenstates

- Ground- and excited states at the optimal HO frequency, $\hbar\Omega=18$ MeV





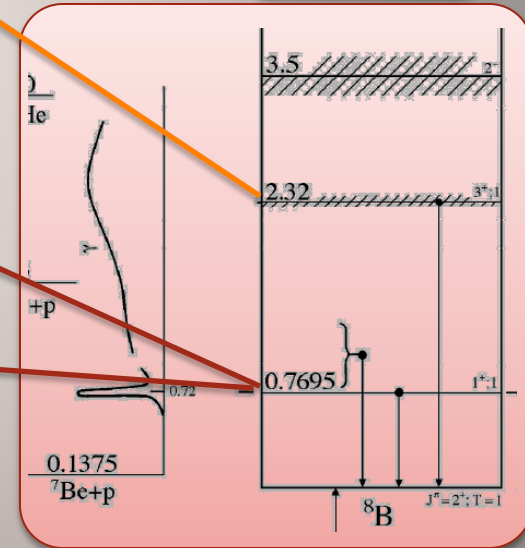
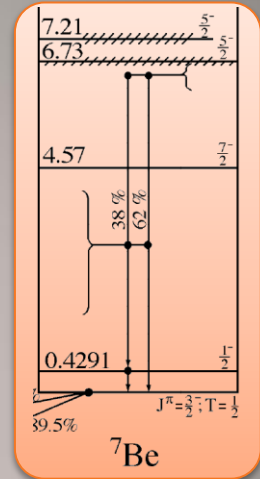
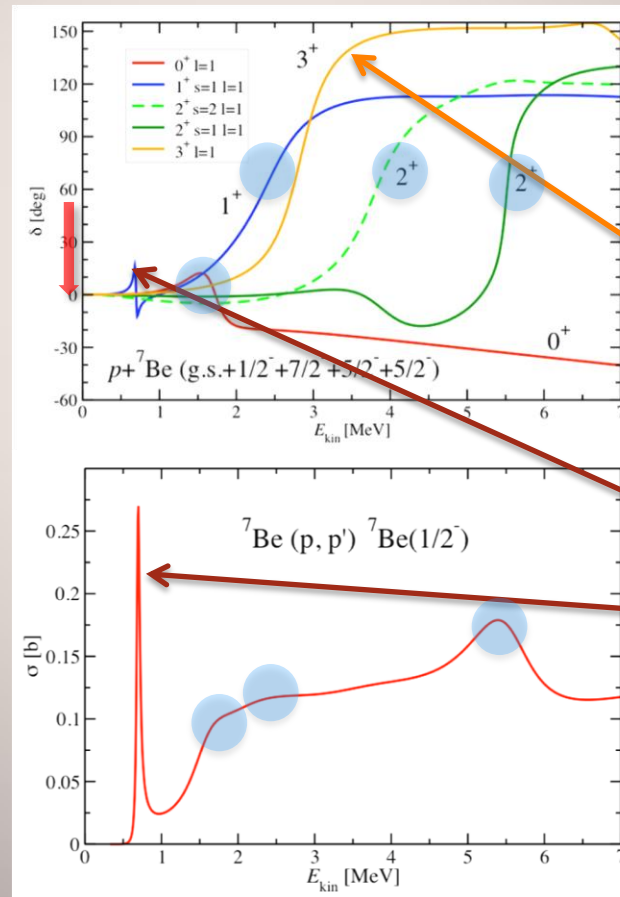
p - ${}^7\text{Be}$ scattering

- NCSM/RGM calculation of p - ${}^7\text{Be}$ scattering
 - ${}^7\text{Be}$ states $3/2^-$, $1/2^-$, $7/2^-$, $5/2^-_1$, $5/2^-_2$
 - Soft NN potential (SRG- N^3LO with $\Lambda = 1.86 \text{ fm}^{-1}$)

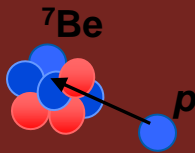
${}^8\text{B}$ 2^+ g.s. bound by 136 keV
(expt. bound by 137 keV)

New 0^+ , 1^+ , and two 2^+ resonances predicted

$s=1$ $l=1$ 2^+ clearly visible in (p, p') cross sections



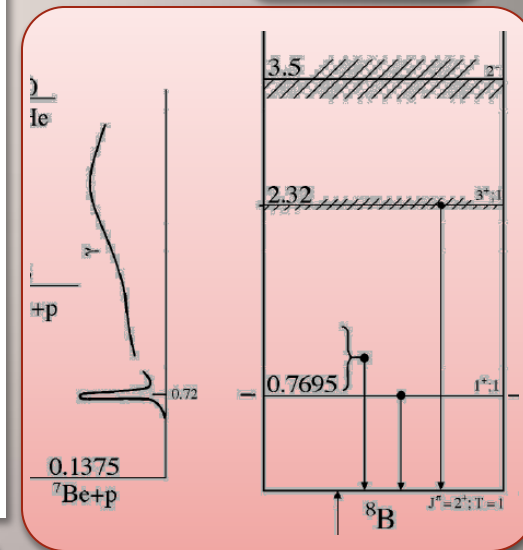
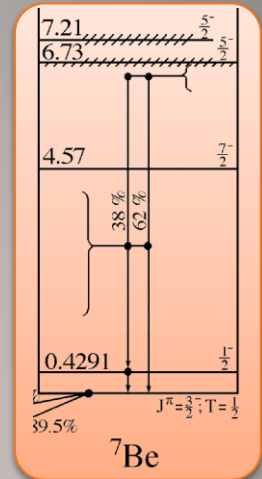
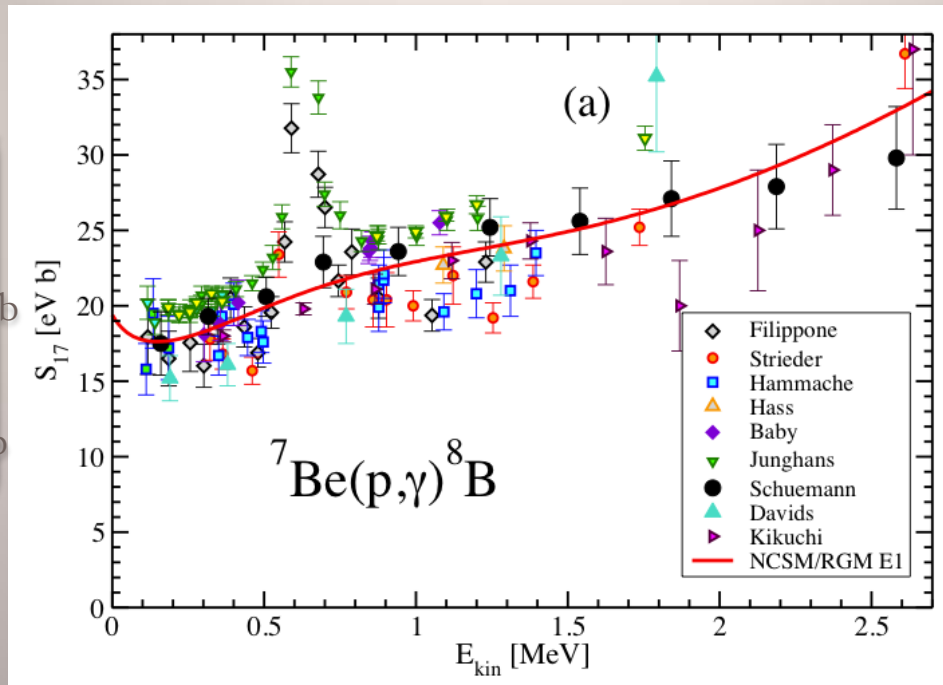
P. Navratil et al., PRC **82**, 034609 (2010)



${}^7\text{Be}(p,\gamma){}^8\text{B}$ radiative capture

- NCSM/RGM calculation of ${}^7\text{Be}(p,\gamma){}^8\text{B}$ radiative capture
 - ${}^7\text{Be}$ states $3/2^-, 1/2^-, 7/2^-, 5/2^-_1, 5/2^-_2$
 - Soft NN potential (SRG- N^3LO with $\Lambda = 1.86 \text{ fm}^{-1}$)

${}^8\text{B}$ 2^+ g.s. bound by
 136 keV
 (expt. 137 keV)
 $S(0) \sim 19.4(0.7) \text{ eV b}$
 Data evaluation:
 $S(0) = 20.8(2.1) \text{ eV b}$



The first ever *ab initio* calculations of ${}^7\text{Be}(p,\gamma){}^8\text{B}$

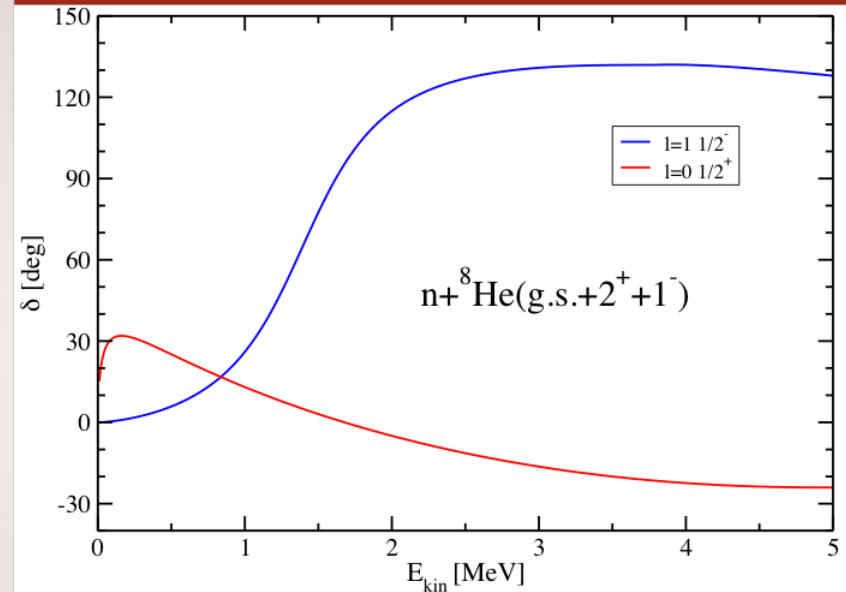
arXiv:1105.5977 [nucl-th]



Structure of the unbound ${}^9\text{He}$ nucleus

- ${}^9\text{He}$ offers the opportunity to study the evolution of nuclear structure as a function of increasing number of neutrons
- Does the ground state of ${}^9\text{He}$ present the same parity inversion observed in the neighboring ${}^{11}\text{Be}$ and ${}^{10}\text{Li}$?
- Disappearance of the $N = 8$ magic number with increasing N/Z ratio
- Controversy on the nature of $S_{1/2}$ contribution to the ${}^9\text{He}$ spectrum
- Here:
 - $n + {}^8\text{He}(\text{g.s.}, 2^+, 1^-)$, $N_{\text{max}} = 13$
 - SRG- $N^3\text{LO}$ NN pot. ($\lambda = 2.02 \text{ fm}^{-1}$)

n - ${}^8\text{He}$ scattering phase shifts

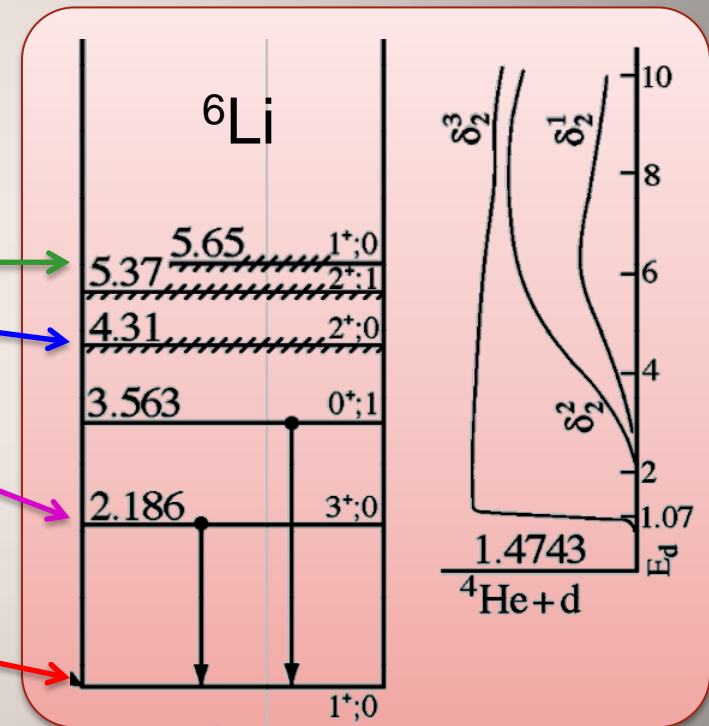
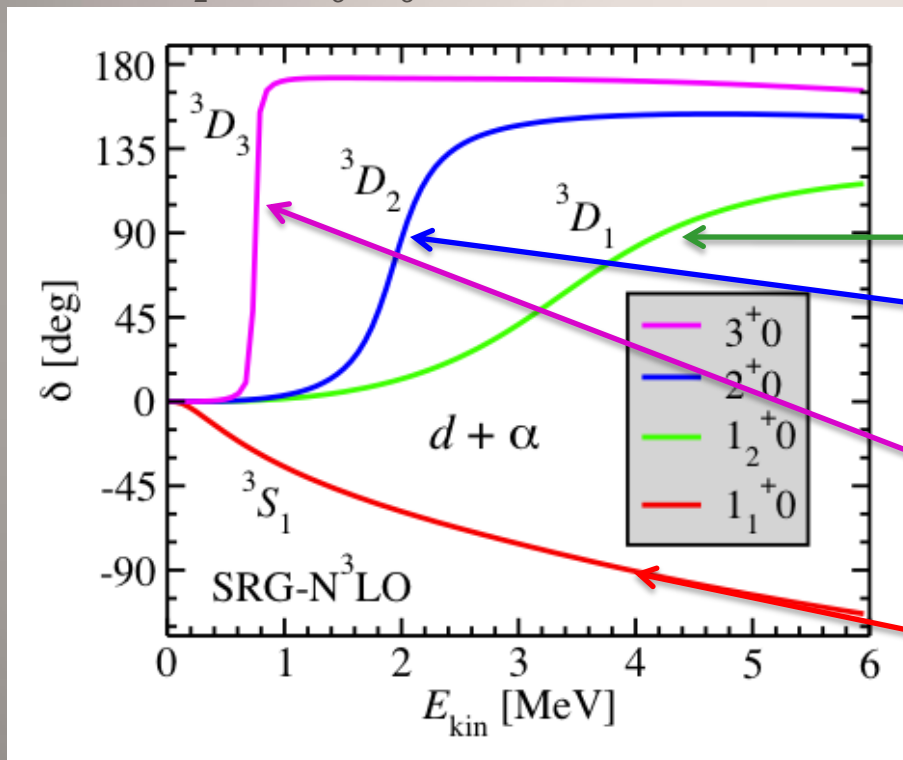


NCSM/RGM results for the S - and P -wave diagonal phase shifts. Need to study N_{max} dependence for an unambiguous answer.

g.s. parity inv. for exotic $N=7$ nuclei, well established in ${}^{11}\text{Be}$ and ${}^{10}\text{Li}$, disappears for ${}^9\text{He}$?

NCSM/RGM *ab initio* calculation of d - ^4He scattering

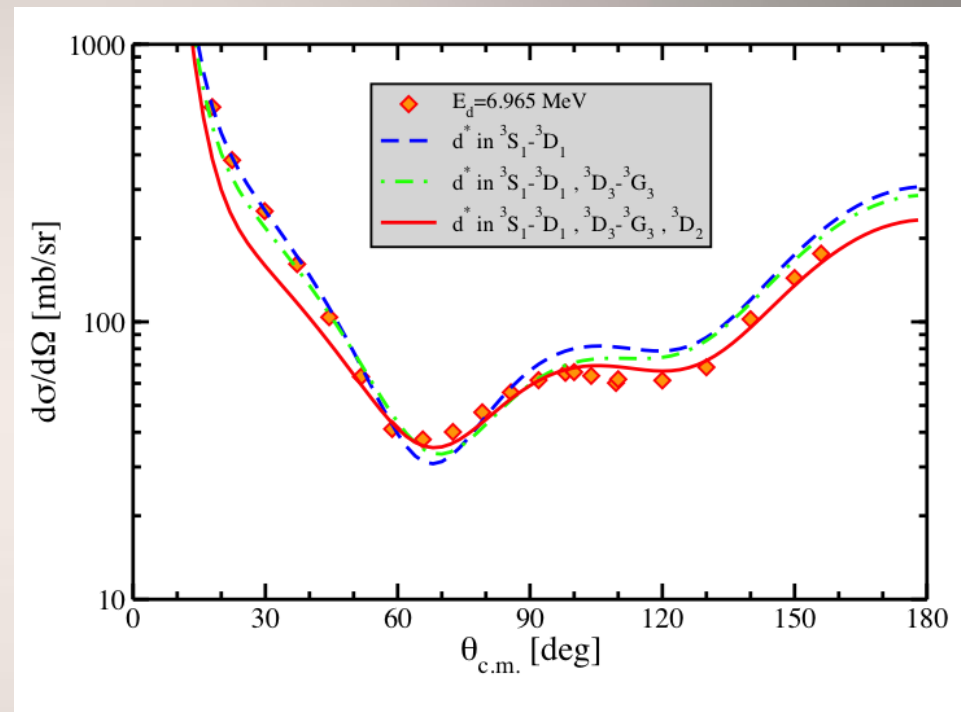
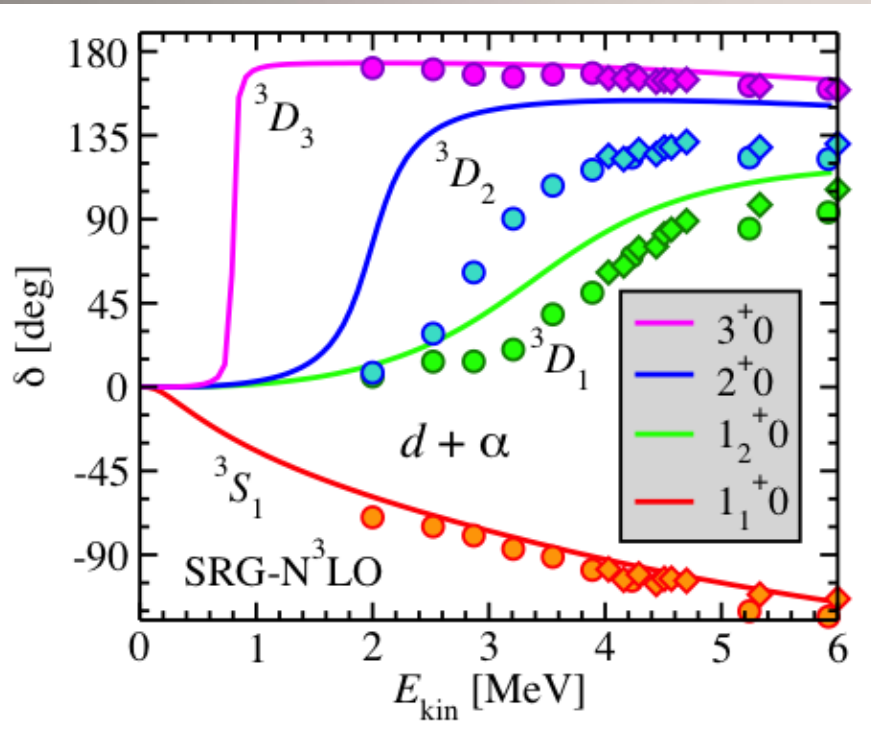
- NCSM/RGM calculation with $d + ^4\text{He}(\text{g.s.})$ up to $N_{\text{max}} = 12$
 - SRG- N^3LO potential with $\Lambda = 1.5 \text{ fm}^{-1}$
 - Deuteron breakup effects included by continuum discretized by pseudo states in $^3\text{S}_1$ - $^3\text{D}_1$, $^3\text{D}_2$ and $^3\text{D}_3$ - $^3\text{G}_3$ channels



- The 1^+0 ground state bound by 1.9 MeV (expt. 1.47 MeV)
- Calculated $T=0$ resonances: 3^+ , 2^+ and 1^+ in correct order close to expt. energies

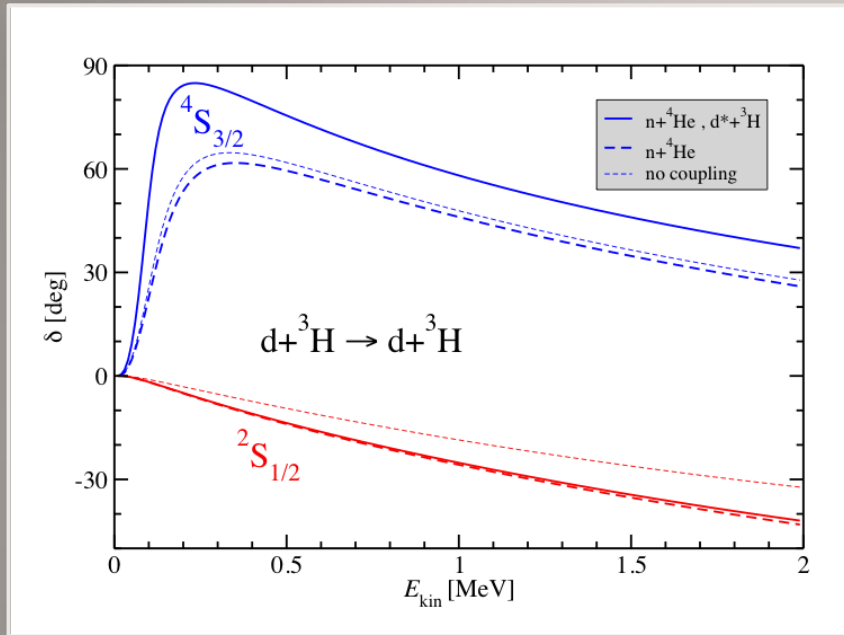
NCSM/RGM *ab initio* calculation of d - ^4He scattering

PHYS. REV. C **83**, 044609 (2011)

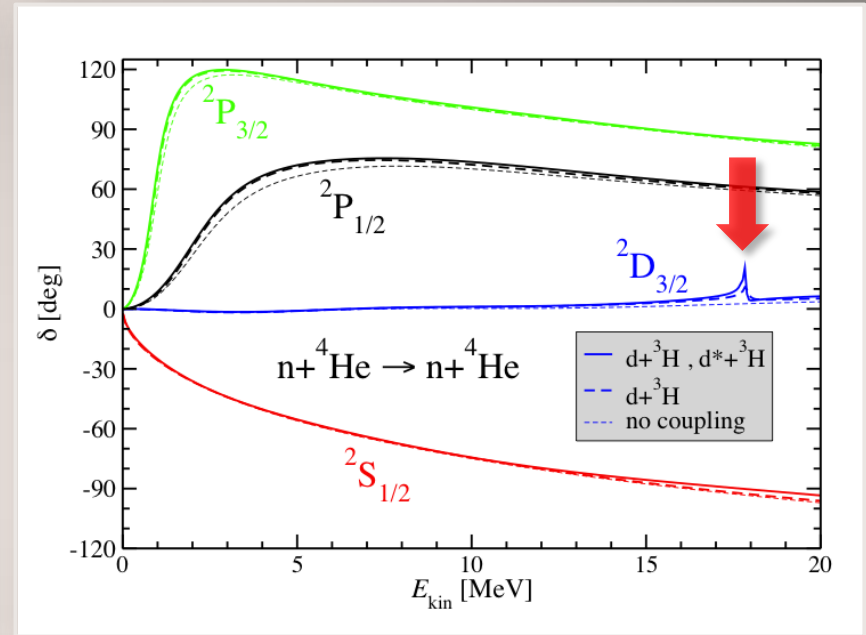


- NCSM/RGM a superior theory: Bound states, resonances, scattering
- NCSM efficiently accounts for many-nucleon correlations: **Coupling of the NCSM and the NCSM/RGM basis desirable**
- Scattering provides a strict test of NN and NNN forces

$d+^3\text{H}$ and $n+^4\text{He}$ elastic scattering: phase shifts



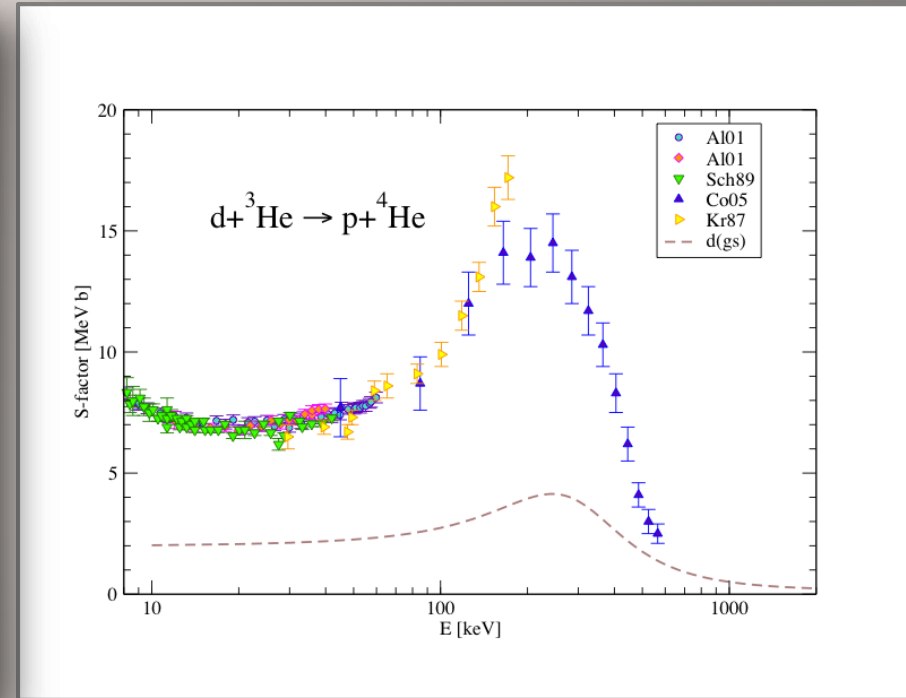
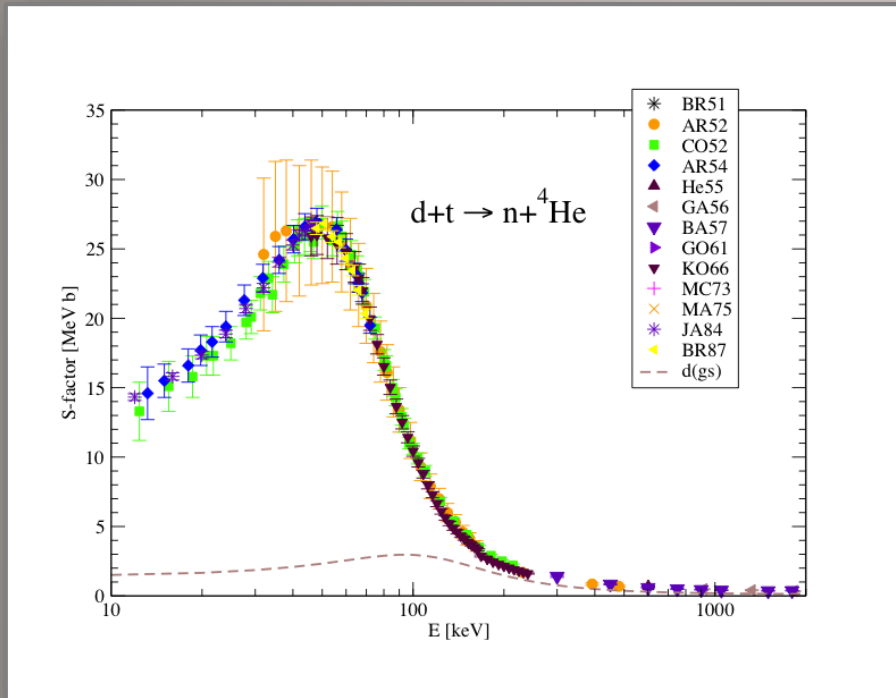
- $d+^3\text{H}$ elastic phase shifts:
 - Resonance in the $^4\text{S}_{3/2}$ channel
 - Repulsive behavior in the $^2\text{S}_{1/2}$ channel \rightarrow Pauli principle



- $n+^4\text{He}$ elastic phase shifts:
 - $d+^3\text{H}$ channels produces slight increase of the P phase shifts
 - Appearance of resonance in the $3/2^+$ D -wave, just above $d-^3\text{H}$ threshold

The D-T fusion takes place through a transition of $d+^3\text{H}$ is S -wave to $n+^4\text{He}$ in D -wave

${}^3\text{H}(d,n){}^4\text{He}$ and ${}^3\text{He}(d,p){}^4\text{He}$ cross sections

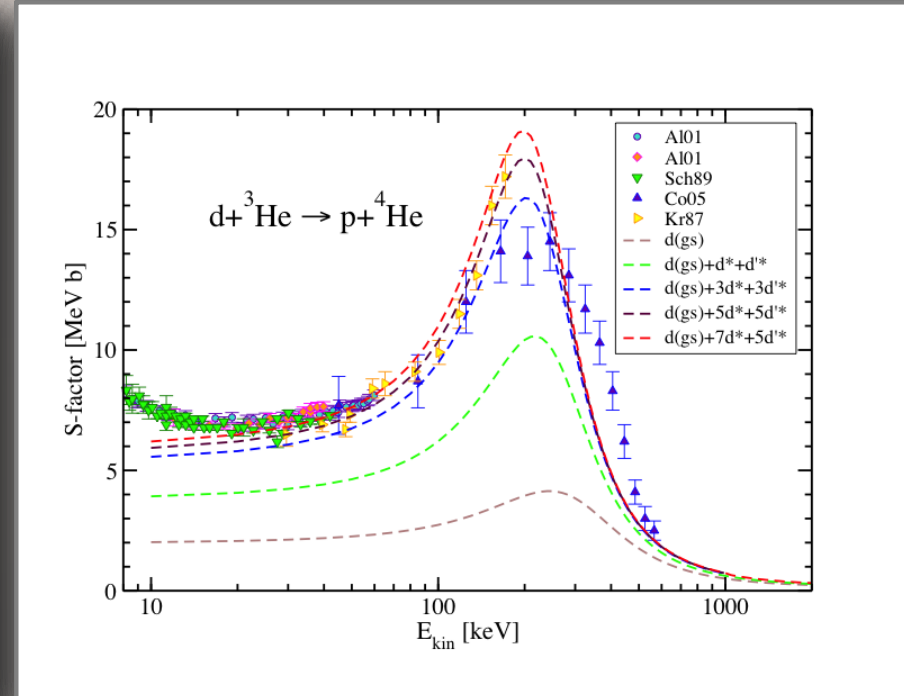
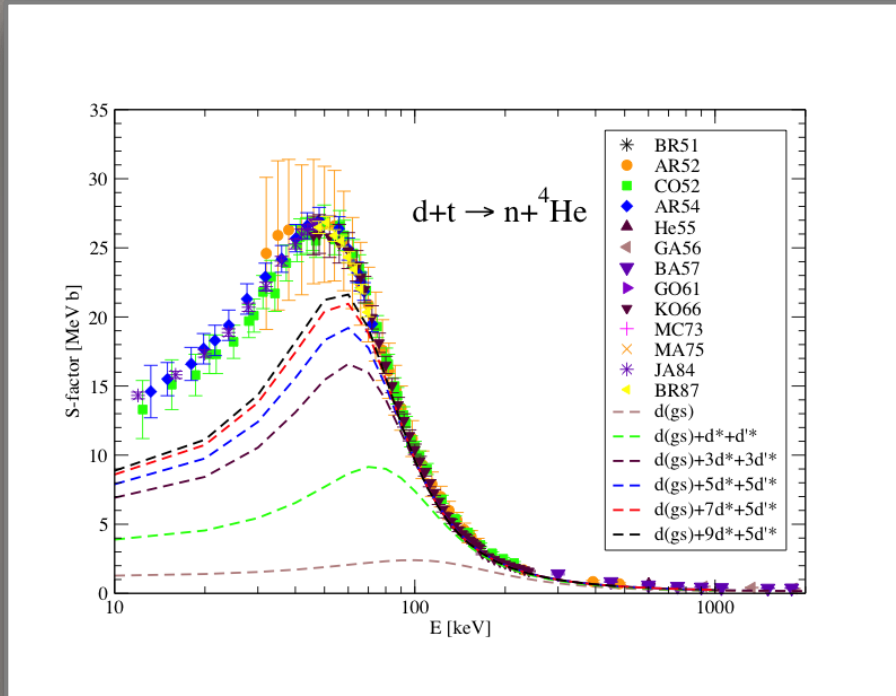


- NCSM/RGM:

- $N_{\text{max}} = 13$
- SRG- $N^3\text{LO}$ NN ($\Lambda = 1.5 \text{ fm}^{-1}$) potential
- NNN interaction interaction effects for $A=3,4,5$ partly included by the choice of Λ
- Only **g.s.** of d , ${}^3\text{H}$, ${}^4\text{He}$ included above

$$S(E) = E\sigma(E) \exp\left(\frac{2\pi Z_1 Z_2 e^2}{h\sqrt{2mE}}\right)$$

${}^3\text{H}(d,n){}^4\text{He}$ and ${}^3\text{He}(d,p){}^4\text{He}$ cross sections



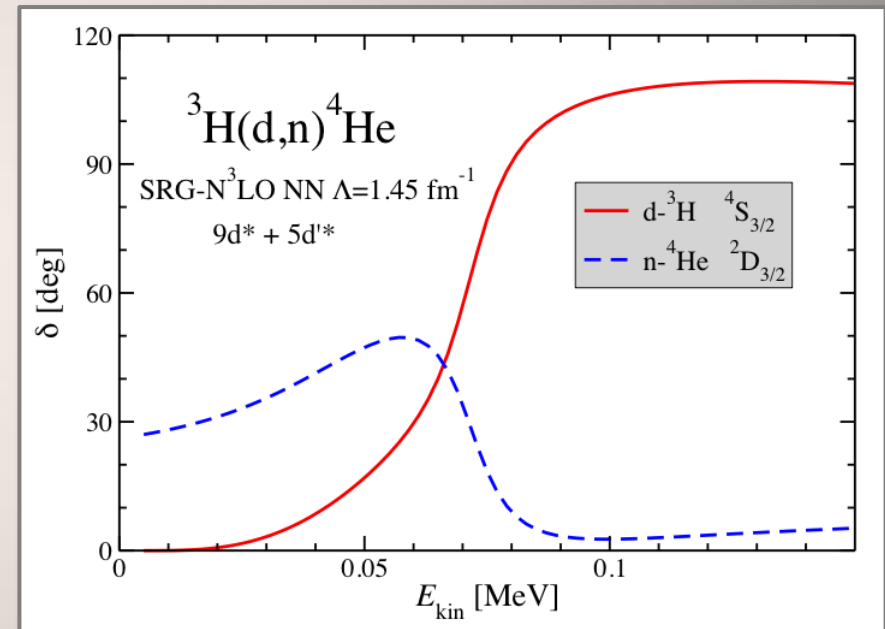
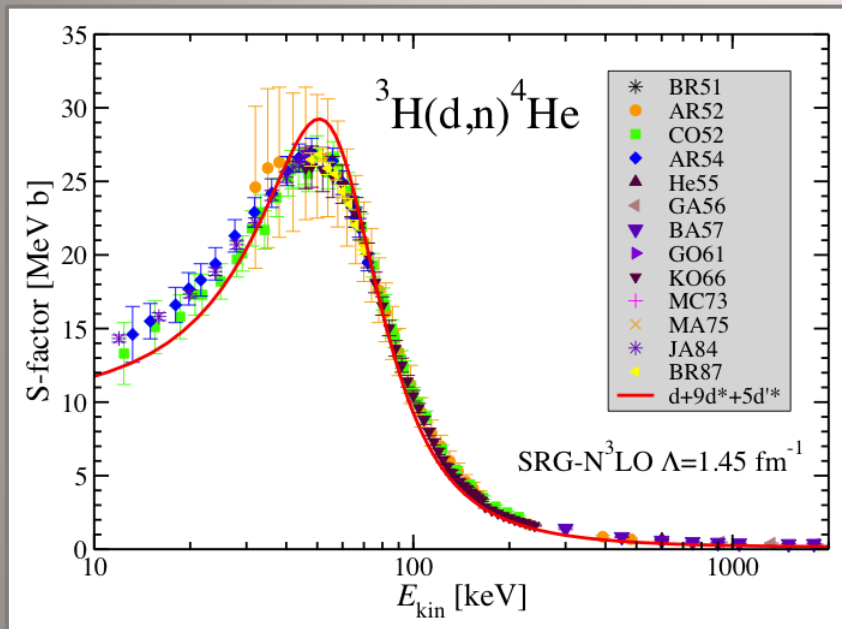
- The cross section improves with the inclusion of virtual breakup of the deuteron
 - Deuteron weakly bound: easily gets polarized and easily breaks
 - These effects included below the breakup threshold with continuum discretized by excited deuteron pseudo-states

First *ab initio* results for $d\text{-T}$ and $d\text{-}{}^3\text{He}$ fusion:

Very promising, correct physics, can become competitive with fitted evaluations ...

${}^3\text{H}(d,n){}^4\text{He}$ cross section

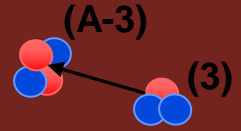
- SRG- N^3LO ($\Lambda=1.45 \text{ fm}^{-1}$) NN potential
 - Position of the resonance matches experiment



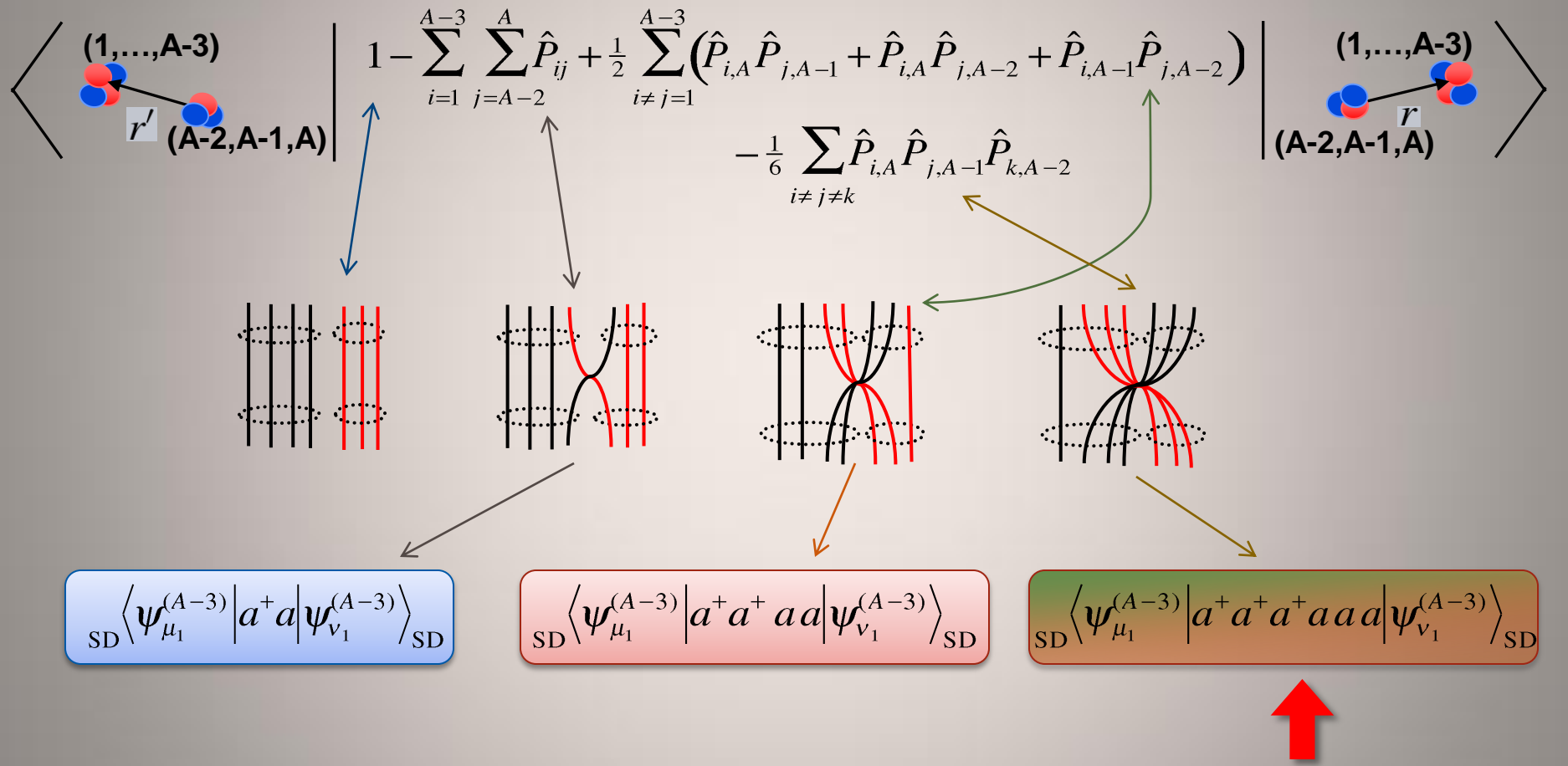
Narrower than the evaluation

Improvements:

Excitations of ${}^4\text{He}$; n-p- ${}^3\text{H}$ rather than d*, d*
 Polarization of ${}^3\text{H}$; NNN interaction; Increase $N_{\text{max}} (=15)$



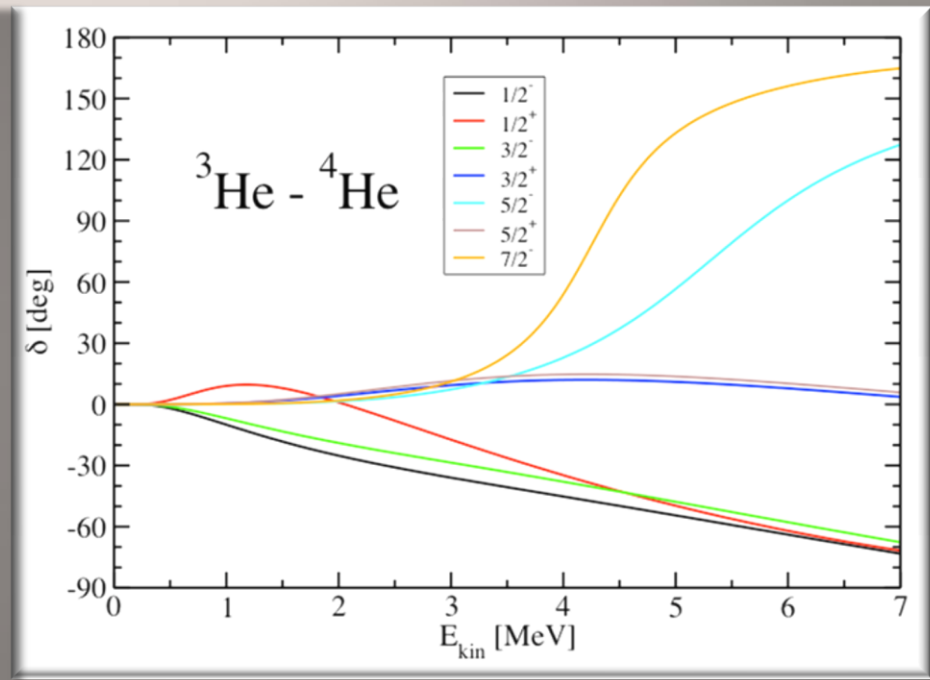
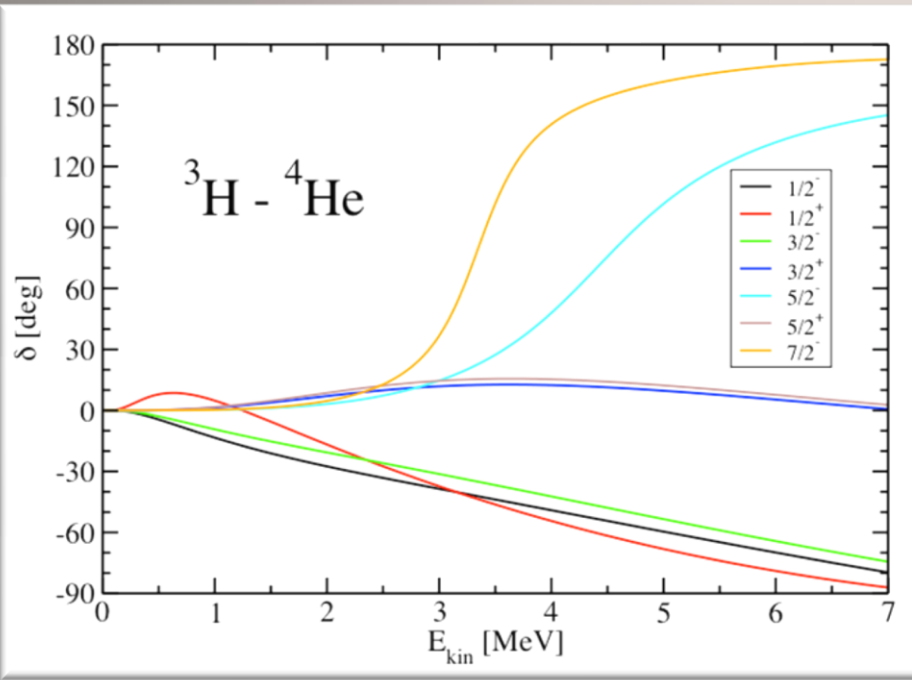
Towards an *ab initio* calculation of the $^3\text{He}+\alpha$ scattering



↑
For A=7 use completeness

Calculations for $a=3$ projectile under way. Kernels completed...

Ab initio calculations of ${}^3\text{H}+\alpha$ and ${}^3\text{He}+\alpha$ scattering: First results (preliminary)



Calculations for $a=3$ projectile under way:
 Soft SRG interactions ($\Lambda=1.5 \text{ fm}^{-1}$), codes working up to $N_{\text{max}}=11$

Addressing the program goals

- Needs of reaction theory
 - More efficient coupling of the *ab initio* reaction theory with the *ab initio* structure calculations
 - No-core shell model with continuum (NCSMC)
 - Coupling of the *ab initio* reaction theory with traditional approaches
 - Breakup reactions on heavy targets (^{11}Be , ^8B ...), fusion, ... with the projectile described *ab initio*
 - Extension of *ab initio* reaction theory to heavier nuclei
 - Higher body-density calculations, $3N$ interactions, importance truncation...
- Making codes available
 - Some of the codes developed at LLNL
 - Proper release procedure must be followed
 - Multiple codes involved, large-scale computation
 - Sharing the codes for collaborations possible now
 - Later a full release possible

Ab initio No-Core Shell Model with continuum

- Original idea:



$$|\Psi_A^J\rangle = \sum c_\lambda |A\lambda J\rangle + \sum \int dr \int dr' \hat{\mathcal{A}} \mathcal{D}_{vr}^{(A-a,a)} \mathcal{N}_{vv'}^{-1/2}(r, r') \chi_w(r')$$

- A better idea:

$$\begin{pmatrix} H & \bar{h} \\ \bar{h} & \mathcal{N}^{-1/2} \mathcal{H} \mathcal{N}^{-1/2} \end{pmatrix} \begin{pmatrix} c \\ \chi \end{pmatrix} = E \begin{pmatrix} 1 & \bar{g} \\ \bar{g} & 1 \end{pmatrix} \begin{pmatrix} c \\ \chi \end{pmatrix}$$

- Test case: ${}^9\text{Li} \leftrightarrow {}^8\text{Li} + n$

– SRG-N³LO NN ($\Lambda=1.9 \text{ fm}^{-1}$), ${}^8\text{Li}(2^+, 1^+, 3^+, 0^+)$, $N_{\text{max}}=6$

– Ground state energy [MeV]:

- ${}^8\text{Li}(2^+)$: (NCSM) -39.27; (Expt.) -41.28
- ${}^9\text{Li}(3/2^-)$: (NCSM/RGM) -42.36; (NCSM) -43.03
(NCSMC-HO) -43.27; (Expt.) -45.34

${}^7\text{Be}(3/2^-)$: (NCSM) -38.19; ${}^8\text{B}(2^+)$: (NCSM/RGM) -38.32; (NCSM) -38.27

Conclusions and Outlook

- With the NCSM/RGM approach we are extending the *ab initio* effort to describe low-energy reactions and weakly-bound systems
- The first ${}^7\text{Be}(p,\gamma){}^8\text{B}$ *ab initio* S-factor calculation arXiv:1105.5977 [nucl-th]
 - Both the bound and the scattering states from first principles
 - SRG-N³LO NN potential selected to match closely the experimental threshold ($\Lambda \approx 1.8 \sim 2 \text{ fm}^{-1}$)
 - No fit: Both normalization and shape predicted
 - Prediction of new ${}^8\text{B}$ resonances
- New results with SRG-N³LO NN potentials: PRC **83**, 044609 (2011)
 - d - ${}^4\text{He}$ scattering
 - Initial results for ${}^3\text{H}(d,n){}^4\text{He}$ & ${}^3\text{He}(d,p){}^4\text{He}$ fusion arXiv:1009.3965 [nucl-th]
- Under way:
 - ${}^3\text{He}+{}^4\text{He}$ scattering calculations
 - *Ab initio* NCSM with continuum (NCSMC)
 - Three-cluster NCSM/RGM and treatment of three-body continuum
- To do:
 - Inclusion of NNN force
 - Alpha clustering: ${}^4\text{He}$ projectile