



# Symbology and Scattering Amplitudes

Anastasia Volovich  
Brown University

based on work with Goncharov, Spradlin and Vergu



# Scattering Amplitudes

- The last few years have seen a lot of progress in our understanding of the mathematical structure of gauge theories and in our ability to do computations both for theoretical and phenomenological purposes.
- Remarkable results range from **precision predictions in QCD** which are important for understanding the LHC data to the discovery of new miraculous structures in **N=4 Yang-Mills** and **N=8 supergravity**.

**The goals of this research program are:**

- to explore the rich hidden mathematical structure of amplitudes
- to exploit this structure as much as possible to make previously impossible computations trivial

Many have contributed to these recent developments including:

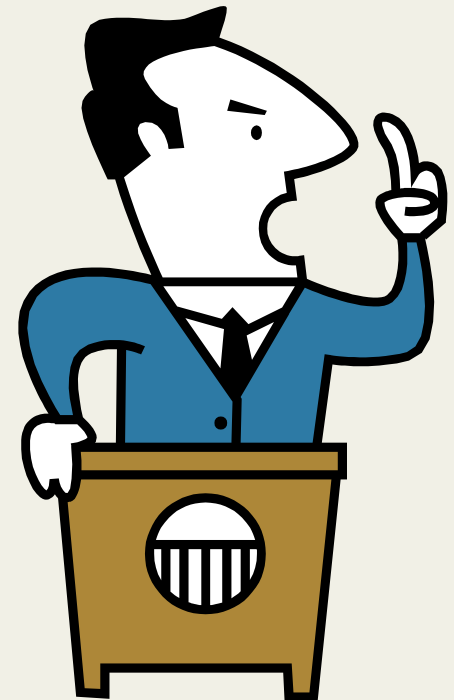
Alday, Arkani-Hamed, Bargheer, Beisert, Belitski, Bern, Berkovits, Boels, Bourjaily, Brandhuber, Broedel, Bjerrum-Bohr, Bullimore, Britto, Cachazo, Caron-Huot, Carrasco, Cheung, Damgaard, Del Duca, Duhr, Dixon, Dolan, Duhr, Drummons, Eden, Elvang, Fend, Ferro, Forde, Freedman, Gaiotto, Goddard, Goncharov, Green, Henn, Heslop, Hodges, Huang, Ita, Johannson, Kallosh, Kaplan, Khoze, Kiermaier, Korchemsky, Kosower, Lipatov, Loebbert, Maitre, Mafra, Maldacena, Mason, Naculich, Nastase, Plefka, Prygarin, Roiban, Sabio Vera, Schnizer, Sever, Skinner, Smirnov, Sokatchev, Spence, Spradlin, Staudacher, Stelle, Steiberger, Svrcek, Taylor, Travaglini, Trnka, Tye, Vanhove, Vergu, Vieira, Wen, Witten

# In My Talk

- I will present **a new very powerful tool** for analysis of multiloop amplitudes = **symbol** of an amplitude = which comes from modern mathematics called **theory of motives**.
- Symbol of an amplitude
  - a very general powerful tool which leads to the compact expression for the amplitude
  - captures the essential physical (and “motivic”) content of the amplitude

# Plan

- Introduction
- Definition of Symbol
- Why is Symbol very useful
- Results: 2-loop 6-points;  
1-loop 2m-gons;
- Outlook and Conclusions



# N=4 Yang-Mills Amplitudeology

We have discovered a lot about N=4 amplitudes

## 1. Dual conformal and Yangian symmetries

Beisert, Henn, Drummond, Korchemsky, McLoughlin, Plefka, Sokatchev, Smirnov

## 2. Amplitude/Wilson loop/Correlation funcs triality

Adamo, Bullimore, Caron-Huot, Mason, Skinner

Alday, Eden, Heslop, Henn, Drummond, Korchemsky, Maldacena, Sokachev,

## 3. Color-kinematic relations

Bern, Carrasco, Johansson; Dennen, Huang, Kiermaier

## 4. Grassmanian & Integrand & Polytopes & Twistors

Arkani-Hamed, Bourjaily, Cachazo, Cheung, Kaplan, Trnka;

Mason, Hodges, Skinner; Spradlin, Wen, AV

## 5. Strong Coupling Y-system

Alday, Gaiotto, Maldacena, Sever, Vieira

# N=4 Yang-Mills Amplitudeology

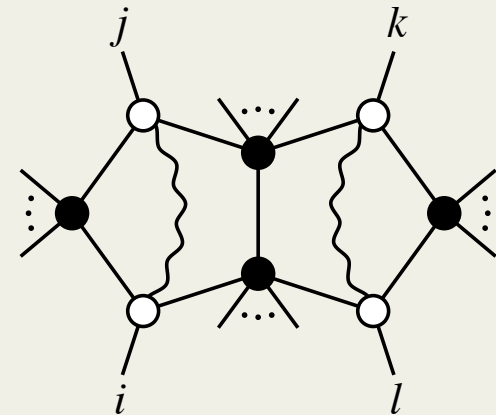
1. Dual conformal and Yangian symmetries
2. Amplitude/Wilson loop/Correlation funcs triality
3. Color-kinematic relations
4. Grassmanian & Integrands & Polytopes & Twistors

## Important Lessons

All of these are completely invisible in the Lagrangian, and were discovered only after doing very hard calculations and then simplifying and analyzing the answers. Simplifications don't happen by accident. This is an **experimental science**, which requires collecting and analyzing lots of "data" to see the hidden structure.

Simplicity has to be believed to be seen!

# Amplitudes and Integrand



- Standard Lore

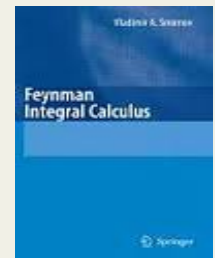
$$\textit{Amplitude} = \int d(\textit{loop momenta}) (\textit{Integrand})$$

- Integrand = solved for all loops and all legs in N=4

Bern, Dixon, et al, Arkani-Hamed et al

- Integral = **still very hard to evaluate!** This step requires a lot of blood and tears begging for some new technology!

e.g. Smirnov's book



- Symbol = “half way between integrand and amplitude”

Goncharov, Spradlin, Vergu, AV



# MHV Remainder function

Simplest nontrivial multi-loop amplitude in N=4 Yang-Mills is the **2-loop 6-particle MHV remainder function** = what is left after subtracting IR divergences, leaving a finite and dual conformal invariant quantity

$$R(u_1, u_2, u_2)$$

$$u_1 = \frac{s_{12}s_{45}}{s_{123}s_{345}} \quad u_2 = \frac{s_{23}s_{56}}{s_{234}s_{456}} \quad u_3 = \frac{s_{34}s_{61}}{s_{345}s_{234}} \quad s_{123} = (k_1 + k_2 + k_3)^2$$

Numerically known since '08 Bern, Dixon, Kosower, Roiban, Vergu, Spradlin, AV Henn, Drummond, Korchemsky, Sokatchev

In a heroic effort, Del Duca, Duhr, Smirnov '09 found a manageable way to evaluate the relevant diagrams and obtained **an analytic formula** for it:













$$\begin{aligned}
& \frac{3}{4}\mathcal{G}\left(v_{231}, 1, \frac{1}{1-u_2}; 1\right)H(0; u_3) + \frac{3}{4}\mathcal{G}\left(v_{231}, \frac{1}{1-u_2}, 1; 1\right)H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(v_{312}, 1, \frac{1}{1-u_3}; 1\right)H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{312}, \frac{1}{1-u_3}, 1; 1\right)H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(v_{321}, 1, \frac{1}{1-u_3}; 1\right)H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{321}, \frac{1}{1-u_3}, 1; 1\right)H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right)H(0; u_1)H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}; 1\right)H(0; u_1)H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right)H(0; u_1)H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, u_{231}; 1\right)H(0; u_1)H(0; u_3) - \\
& \frac{1}{4}G\left(\frac{1}{1-u_2}, v_{213}; 1\right)H(0; u_1)H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right)H(0; u_1)H(0; u_3) + \\
& \frac{5}{24}\pi^2H(0; u_1)H(0; u_3) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}; 1\right)H(0; u_2)H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right)H(0; u_2)H(0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right)H(0; u_2)H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, u_{123}; 1\right)H(0; u_2)H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right)H(0; u_2)H(0; u_3) - \\
& \frac{1}{4}G\left(\frac{1}{1-u_1}, v_{132}; 1\right)H(0; u_2)H(0; u_3) + \frac{5}{24}\pi^2H(0; u_2)H(0; u_3) + \\
& 3H(0; u_2)H(0; u_1)H(0; u_3) + 3H(0; u_1)H(0; u_2)H(0; u_3) + \\
& \frac{1}{4}H(0; u_2)H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right)H(0; u_3) + \frac{1}{2}H(0; u_1)H(0, 1; (u_1+u_3))H(0; u_3) + \\
& \frac{1}{4}H(0; u_1)H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right)H(0; u_3) + \frac{1}{2}H(0; u_2)H(0, 1; (u_2+u_3))H(0; u_3) + \\
& \frac{3}{4}H(0; u_2)H(1, 0; u_1)H(0; u_3) + \frac{3}{4}H(0; u_1)H(1, 0; u_2)H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{213}; 1\right)H(0, 0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right)H(0, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, v_{312}; 1\right)H(0, 0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right)H(0, 0; u_1) - \frac{23}{24}\pi^2H(0, 0; u_1) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right)H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right)H(0, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, v_{312}; 1\right)H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right)H(0, 0; u_2) - \\
& \frac{25}{4}H(0, 0; u_1)H(0, 0; u_2) - \frac{23}{24}\pi^2H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right)H(0, 0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right)H(0, 0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{213}; 1\right)H(0, 0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right)H(0, 0; u_3) + 3H(0; u_1)H(0; u_2)H(0, 0; u_3) - \\
& \frac{25}{4}H(0, 0; u_1)H(0, 0; u_3) - \frac{25}{4}H(0; u_2)H(0, 0; u_3) - \frac{23}{24}\pi^2H(0, 0; u_3) + \frac{1}{12}\pi^2H(0, 1; u_1) + \\
& \frac{1}{12}\pi^2H(0, 1; u_2) - \frac{1}{24}\pi^2H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right) + \frac{1}{2}H(0; u_1)H(0; u_2)H(0, 1; (u_1+u_2)) +
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{12}\pi^2H(0, 1; (u_1+u_2)) + \frac{1}{12}\pi^2H(0, 1; u_3) + \frac{1}{4}H(0; u_1)H(0; u_2)H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) - \\
& \frac{1}{24}\pi^2H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) + \frac{1}{12}\pi^2H(0, 1; (u_1+u_3)) - \frac{1}{24}\pi^2H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) + \\
& \frac{1}{12}\pi^2H(0, 1; (u_2+u_3)) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_1) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, \frac{u_1-1}{u_1+u_3-1}; 1\right)H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_1) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, u_{312}; 1\right)H(1, 0; u_1) - \frac{3}{4}H(0, 0; u_2)H(1, 0; u_1) - \frac{3}{4}H(0, 0; u_3)H(1, 0; u_1) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right)H(1, 0; u_1) - \frac{1}{3}\pi^2H(1, 0; u_1) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}; 1\right)H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_2) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, u_{123}; 1\right)H(1, 0; u_2) - \frac{3}{4}H(0, 0; u_1)H(1, 0; u_2) - \frac{3}{4}H(0, 0; u_3)H(1, 0; u_2) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right)H(1, 0; u_2) - \frac{1}{4}H(1, 0; u_1)H(1, 0; u_2) - \frac{1}{3}\pi^2H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_3) - \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}; 1\right)H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_3) - \frac{1}{3}\pi^2H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, u_{231}; 1\right)H(1, 0; u_3) + \\
& \frac{3}{4}H(0; u_1)H(0; u_2)H(1, 0; u_3) - \frac{3}{4}H(0, 0; u_1)H(1, 0; u_3) - \frac{3}{4}H(0, 0; u_2)H(1, 0; u_3) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right)H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_1)H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_2)H(1, 0; u_3) + \\
& \frac{1}{24}\pi^2H(1, 1; u_1) + \frac{1}{24}\pi^2H(1, 1; u_2) + \frac{1}{24}\pi^2H(1, 1; u_3) + \frac{1}{2}H(0; u_2)H(0, 0, 0; u_1) + \\
& \frac{1}{2}H(0; u_3)H(0, 0, 0; u_2) + \frac{1}{2}H(0; u_1)H(0, 0, 0; u_3) - \frac{1}{2}H(0; u_2)H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - \\
& \frac{1}{2}H(0; u_3)H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - H(0; u_1)H(0, 0, 1; (u_1+u_2)) - \\
& H(0; u_2)H(0, 0, 1; (u_1+u_2)) - \frac{1}{2}H(0; u_1)H\left(0, 0, 1; \frac{u_1+u_3-1}{u_1-1}\right) -
\end{aligned}$$





$$\begin{aligned}
& \frac{1}{4}H(0; u_2) \mathcal{H} \left( 0, 1, 1; \frac{1}{u_{312}} \right) + \frac{1}{4}H(0; u_2) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{123}} \right) - \frac{1}{4}H(0; u_3) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{123}} \right) - \\
& \frac{1}{4}H(0; u_2) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{132}} \right) + \frac{1}{4}H(0; u_3) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{132}} \right) + \frac{1}{4}H(0; u_1) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{213}} \right) - \\
& \frac{1}{4}H(0; u_3) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{213}} \right) - \frac{1}{4}H(0; u_1) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{231}} \right) + \frac{1}{4}H(0; u_3) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{231}} \right) + \\
& \frac{1}{4}H(0; u_1) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{312}} \right) - \frac{1}{4}H(0; u_2) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{312}} \right) - \frac{1}{4}H(0; u_1) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{321}} \right) + \\
& \frac{1}{4}H(0; u_2) \mathcal{H} \left( 0, 1, 1; \frac{1}{v_{321}} \right) + \frac{1}{4}H(0; u_3) \mathcal{H} \left( 1, 0, 1; \frac{1}{u_{123}} \right) + \frac{1}{4}H(0; u_1) \mathcal{H} \left( 1, 0, 1; \frac{1}{u_{231}} \right) + \\
& \frac{1}{4}H(0; u_2) \mathcal{H} \left( 1, 0, 1; \frac{1}{u_{312}} \right) + \frac{1}{4}H(0; u_2) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{123}} \right) - \frac{1}{4}H(0; u_3) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{123}} \right) - \\
& \frac{1}{4}H(0; u_2) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{132}} \right) + \frac{1}{4}H(0; u_3) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{132}} \right) + \frac{1}{4}H(0; u_1) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{213}} \right) - \\
& \frac{1}{4}H(0; u_3) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{213}} \right) - \frac{1}{4}H(0; u_1) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{231}} \right) + \frac{1}{4}H(0; u_3) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{231}} \right) + \\
& \frac{1}{4}H(0; u_1) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{312}} \right) - \frac{1}{4}H(0; u_2) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{312}} \right) - \frac{1}{4}H(0; u_1) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{321}} \right) + \\
& \frac{1}{4}H(0; u_2) \mathcal{H} \left( 1, 0, 1; \frac{1}{v_{321}} \right) + H(0; u_2) \mathcal{H} \left( 1, 1, 1; \frac{1}{v_{123}} \right) - H(0; u_3) \mathcal{H} \left( 1, 1, 1; \frac{1}{v_{123}} \right) - \\
& H(0; u_1) \mathcal{H} \left( 1, 1, 1; \frac{1}{v_{231}} \right) + H(0; u_3) \mathcal{H} \left( 1, 1, 1; \frac{1}{v_{231}} \right) + H(0; u_1) \mathcal{H} \left( 1, 1, 1; \frac{1}{v_{312}} \right) - \\
& H(0; u_2) \mathcal{H} \left( 1, 1, 1; \frac{1}{v_{312}} \right) - \frac{3}{2}\mathcal{H} \left( 0, 0, 0, 1; \frac{1}{u_{123}} \right) - \frac{3}{2}\mathcal{H} \left( 0, 0, 0, 1; \frac{1}{u_{231}} \right) - \\
& \frac{3}{2}\mathcal{H} \left( 0, 0, 0, 1; \frac{1}{u_{312}} \right) - 3\mathcal{H} \left( 0, 0, 0, 1; \frac{1}{v_{132}} \right) - 3\mathcal{H} \left( 0, 0, 0, 1; \frac{1}{v_{213}} \right) - 3\mathcal{H} \left( 0, 0, 0, 1; \frac{1}{v_{321}} \right) - \\
& \frac{1}{2}\mathcal{H} \left( 0, 0, 1, 1; \frac{1}{u_{123}} \right) - \frac{1}{2}\mathcal{H} \left( 0, 0, 1, 1; \frac{1}{u_{231}} \right) - \frac{1}{2}\mathcal{H} \left( 0, 0, 1, 1; \frac{1}{u_{312}} \right) - \\
& \frac{1}{2}\mathcal{H} \left( 0, 1, 0, 1; \frac{1}{u_{123}} \right) - \frac{1}{2}\mathcal{H} \left( 0, 1, 0, 1; \frac{1}{u_{231}} \right) - \frac{1}{2}\mathcal{H} \left( 0, 1, 0, 1; \frac{1}{u_{312}} \right) + \\
& \frac{1}{4}\mathcal{H} \left( 0, 1, 1, 1; \frac{1}{v_{123}} \right) + \frac{1}{4}\mathcal{H} \left( 0, 1, 1, 1; \frac{1}{v_{132}} \right) + \zeta_3 H(0; u_1) + \zeta_3 H(0; u_2) + \zeta_3 H(0; u_3) + \\
& \frac{5}{2}\zeta_3 H(1; u_1) + \frac{5}{2}\zeta_3 H(1; u_2) + \frac{5}{2}\zeta_3 H(1; u_3) + \frac{1}{2}\zeta_3 \mathcal{H} \left( 1; \frac{1}{u_{123}} \right) + \frac{1}{2}\zeta_3 \mathcal{H} \left( 1; \frac{1}{u_{231}} \right) + \\
& \frac{1}{2}\zeta_3 \mathcal{H} \left( 1; \frac{1}{u_{312}} \right) - \frac{1}{2}\mathcal{H} \left( 1, 0, 0, 1; \frac{1}{u_{123}} \right) - \frac{1}{2}\mathcal{H} \left( 1, 0, 0, 1; \frac{1}{u_{231}} \right) - \frac{1}{2}\mathcal{H} \left( 1, 0, 0, 1; \frac{1}{u_{312}} \right) + \\
& \frac{1}{4}\zeta_3 \mathcal{H} \left( 1; \frac{1}{v_{123}} \right) + \frac{1}{4}\zeta_3 \mathcal{H} \left( 1; \frac{1}{v_{132}} \right) + \frac{1}{4}\zeta_3 \mathcal{H} \left( 1; \frac{1}{v_{213}} \right) + \frac{1}{4}\zeta_3 \mathcal{H} \left( 1; \frac{1}{v_{231}} \right) + \frac{1}{4}\zeta_3 \mathcal{H} \left( 1; \frac{1}{v_{312}} \right) + \\
& \frac{1}{4}\zeta_3 \mathcal{H} \left( 1; \frac{1}{v_{321}} \right) + \frac{1}{4}\mathcal{H} \left( 0, 1, 1, 1; \frac{1}{v_{213}} \right) + \frac{1}{4}\mathcal{H} \left( 0, 1, 1, 1; \frac{1}{v_{231}} \right) + \frac{1}{4}\mathcal{H} \left( 0, 1, 1, 1; \frac{1}{v_{312}} \right) + \\
& \frac{1}{4}\mathcal{H} \left( 0, 1, 1, 1; \frac{1}{v_{321}} \right) + \frac{1}{4}\mathcal{H} \left( 1, 0, 1, 1; \frac{1}{v_{123}} \right) + \frac{1}{4}\mathcal{H} \left( 1, 0, 1, 1; \frac{1}{v_{132}} \right) + \frac{1}{4}\mathcal{H} \left( 1, 0, 1, 1; \frac{1}{v_{213}} \right) + \\
& \frac{1}{4}\mathcal{H} \left( 1, 0, 1, 1; \frac{1}{v_{231}} \right) + \frac{1}{4}\mathcal{H} \left( 1, 0, 1, 1; \frac{1}{v_{312}} \right) + \frac{1}{4}\mathcal{H} \left( 1, 1, 0, 1; \frac{1}{v_{123}} \right) + \frac{1}{4}\mathcal{H} \left( 1, 1, 0, 1; \frac{1}{v_{132}} \right) + \\
& \frac{1}{4}\mathcal{H} \left( 1, 1, 0, 1; \frac{1}{v_{213}} \right) + \frac{1}{4}\mathcal{H} \left( 1, 1, 0, 1; \frac{1}{v_{231}} \right) + \frac{1}{4}\mathcal{H} \left( 1, 1, 0, 1; \frac{1}{v_{312}} \right) + \\
& \frac{1}{4}\mathcal{H} \left( 1, 1, 0, 1; \frac{1}{v_{321}} \right)
\end{aligned}$$

There should be a better presentation!

Otherwise we should abandon N=4 Yang Mills!



FullSimplify[ ] is not enough, we need a serious weapon to battle this 17-page



# How do we proceed?

Classical Polylogs

$$Li_k(z) = \int_0^z Li_{k-1}(t) d \log t \quad Li_1(z) = -\log(1-z)$$

Satisfy various identities

$$-Li_2\left(1 - \frac{1}{x}\right) = Li_2(1-x) + \frac{1}{2} \log(x)^2$$

Goncharov polylogs satisfy huge number of identities

$$G(a_k, a_{k-1}, \dots; z) = \int_0^z G(a_{k-1}, \dots; t) \frac{dt}{t - a_k}, \quad G(z) \equiv 1$$

How can we bring the identities under control?

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How can we bring the identities under control?

**We will use motivic high tech!!!**

# Symbol of Transcendental Function

Goncharov

$$T_k \rightarrow S(T_k) = R_1 \otimes \cdots \otimes R_k$$

Symbol is an element of the k-fold tensor product of the multiplicative group of rational functions, it can be defined recursively

$$dT_k = \sum_i T_{k-1}^i d \log R_i \rightarrow S(T_k) = \sum_i S(T_{k-1}^i) \otimes R_i$$

$$R_1 \otimes (R_2 R_3) = R_1 \otimes R_2 + R_1 \otimes R_3$$

Properties

$$R_1 \otimes (cR_2) = R_1 \otimes R_2 \quad c = \text{const}$$

## Examples

Function	Differential	Symbol
$\log R$	$d \log R$	$R$
$\log R_1 \log R_2$	$\log R_1 d \log R_2 + \log R_2 d \log R_1$	$R_1 \otimes R_2 + R_2 \otimes R_1$
$Li_2(R)$	$-\log(1 - R) d \log R$	$-(1 - R) \otimes R$
$Li_k(R)$	$Li_{k-1}(R) d \log R$	$-(1 - R) \otimes R \cdots \otimes R$

# Why is the Symbol Very Useful?

Symbol converts polylog functional equations into rational function identities



◆  $Li_2(x) + Li_2(-x) = \frac{1}{2}Li_2(x^2)$

$-(1-x) \otimes x - (1+x) \otimes (-x) = -(1-x^2) \otimes x = -\frac{1}{2}(1-x^2) \otimes x^2$

◆  $Li_2(z) + Li_2(1-z) + \log(z) \log(1-z) = \frac{\pi^2}{6}$   
 $-z \otimes (1-z) - (1-z) \otimes z + z \otimes (1-z) + (1-z) \otimes z = 0$

Symbol fixes only the leading functional transcendental piece



# Our Strategy

- Compute symbol of DDS = **very simple!** It has many nice properties and symmetries Goncharov

$$A \wedge A = 0 \rightarrow \text{classical Li's only}$$

- Construct the function which has the same symbol, its nontrivial, need luck/guessing

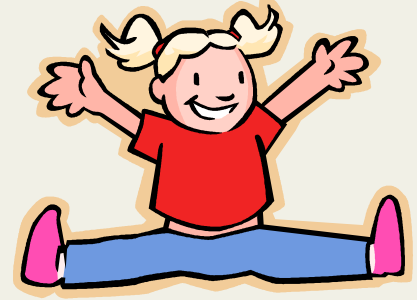
$$(1 + x + x^2) \otimes x = \frac{1-x^3}{1-x} \otimes x = \frac{1}{3}(1 - x^3) \otimes x^3 - (1 - x) \otimes x$$



$$-\frac{1}{3}Li_2(x^3) + Li_2(x)$$

- Fix the lower transcendental terms by numerical fit or collinear limits

# Final Result



Goncharov, Spradlin, Vergu, AV '10

Two-loop six-point MHV remainder function is

$$R_6^{(2)}(u_1, u_2, u_3) = \sum_{i=1}^3 (L_4(x_i^+, x_i^-) - \frac{1}{2} Li_4(1-1/u_i)) - \frac{1}{8} \left( \sum_{i=1}^3 Li_2(1-1/u_i) \right)^2 + \frac{1}{24} J^4 + \frac{\pi^2}{12} J^2 + \frac{\pi^4}{72}.$$

$$L_4(x^+, x^-) = \sum_{m=0}^3 \frac{(-1)^m}{(2m)!!} \log(x^+ x^-)^m (\ell_{4-m}(x^+) + \ell_{4-m}(x^-)) + \frac{1}{8!!} \log(x^+ x^-)^4$$

$$\ell_n(x) = \frac{1}{2} (Li_n(x) - (-1)^n Li_n(1/x))$$

$$J = \sum_{i=1}^3 (\ell_1(x_i^+) - \ell_1(x_i^-))$$

$$x_i^\pm = u_i x^\pm$$

$$x^\pm = \frac{u_1 + u_2 + u_3 - 1 \pm \sqrt{\Delta}}{2u_1 u_2 u_3}$$

$$\Delta = (u_1 + u_2 + u_3 - 1)^2 - 4u_1 u_2 u_3$$

Bartels, Prygarin, Lipatov '10: analyzed Regge limit  
Gaiotto, Maldacena, Sever, Veira '11: OPE approach

# What's next?

- Our formula provides hope to the idea that we might be able to really unlock the secrets of multi-loop SYM amplitudes, and connect to strong coupling **Alday, Gaiotto, Maldacena, Sever, Veira**
- To do: higher points  
**Gaiotto, Maldacena, Sever, Veiera OPE approach:** one integral  
**Caron-Huot:** Symbol for all n, need to integrate.  
non-MHV ratio function  
higher loops **Bartels, Lipatov, Prygarin**  
**Dixon, Drummond, Henn:** 3-loops 6-pts symbol
- It would be nice to have a derivation from integrals: work on two-loops is in progress, but solved for one-loop higher dimensional polygons

# More Motivic Magic

Spradlin, AV 05/11

- Consider **one-loop  $2m$ -gon in  $2m$ -dimensions**:  
appear in dim regulated one-loop MHV & can also be related to higher-loop four-dimensional integrals
- One-loop hexagon in six dimensions have been evaluated for massless & three-mass case:  
surprisingly the same  $L_3$  function appears.....

Dixon, Drummond, Henn; Del Duca, Duhr, Smirnov 05/11

- It turns out its symbol for **any  $m$**  can be read off from theorem by **Goncharov'96** on mixed Tate motives

# Recursions for Symbols

- After Feynman parameterization, any one-loop  $2m$ -gon integral in  $D=2m$  takes the form

$$F(Q) = \int_{\mathbb{CP}^{2m-1}} \frac{D^{2m-1} W}{(W \cdot Q \cdot W)^m} \quad Q_{ij} = (p_i + \cdots + p_j)^2$$

in terms of a quadric  $Q$  in  $\mathbb{CP}^{2m-1}$

- The symbol of the integral is given recursively by application of a result on mixed Tate motives [Goncharov 96](#)

$$S_m(Q) = \sum_{i < j} S_{m-1}(\overline{Q}_{ij}) \otimes \frac{Q_{ij}^{-1} - \sqrt{(Q_{ij}^{-1})^2 - Q_{ii}^{-1} Q_{jj}^{-1}}}{Q_{ij}^{-1} + \sqrt{(Q_{ij}^{-1})^2 - Q_{ii}^{-1} Q_{jj}^{-1}}}$$

( $\overline{Q} = Q$  with rows and columns  $i, j$  deleted) [Spradlin, AV 05/11](#)

# Conclusions

- If a problem is too hard, look for something simpler. The symbol is a useful stepping stone half way between an integrand and its integral.
- We need a technology for writing down symbols without first evaluating integrals at all. *in progress*
- Feynman integrals are natural projective objects, loop amplitude should also have a direct relation to the Grassmannian formulation for  $N=4$ . *in progress*
- Symbols are not the ultimate goal, but serve the simplest incarnation of a deeper **motivic structure of amplitudes** which we are only beginning to see.

**A LOT OF WORK TO DO!!!**