

Amplitudes in $D \neq 4$

Yu-tin Huang

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(Z. Bern, J. J. Carrasco, W-M Cheng, B. Czech, T. Dennen, D. Gang, H. Ita, E. Koh, S. Lee, A. Lipstein, M. Rozali, W. Siegel)

Sept-29 2011 INT

(Lorentz)
 λ
(Little)

$$P = \lambda\lambda \quad \text{Lorentz} \times \text{Little} - \text{LittleGroup} = D - 1$$

$$D = 3, 4, 6$$

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$$D = 3, 4, 6$$

1 3-D Chern-Simons matter theory (ABJM)

2 D=6

- D=6 Maximal SYM
- $\mathcal{N} = (2, 0)$ Self-Dual Tensor

3D spinor helicity

$$\lambda \xrightarrow{\alpha} \text{SL}(2\mathbb{R})$$
$$Z_2$$

Field Content

Aharony, Bergman, Jafferis, and Maldacena (ABJM)

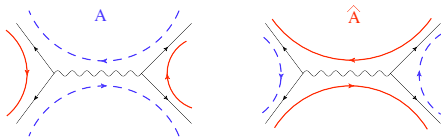
Gauge fields:

$$(A)^a{}_b, (\hat{A})^{\hat{a}}{}_{\hat{b}} \in \text{SU}(N) \times \text{SU}(N)$$

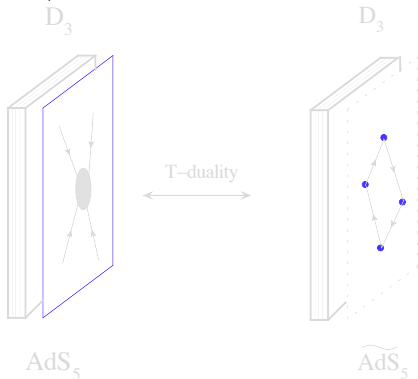
$$\mathcal{L}_{CS} = A \wedge dA + \frac{1}{3} A \wedge A \wedge A$$

Matter fields:

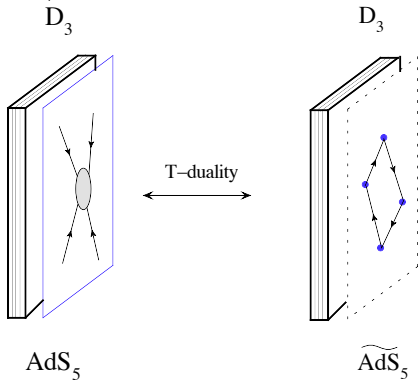
$$(\phi^I, \psi^I)^a{}_{\hat{a}}, (\bar{\phi}_I, \bar{\psi}_I)^{\hat{a}}{}_b, \quad I \in \text{SU}(4)$$



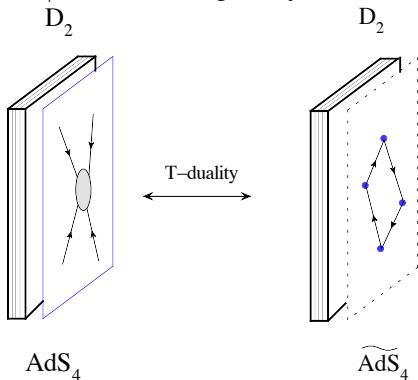
- Non-maximal, $\mathcal{N} = 6$
- SCFT $\text{OSp}(6|4) \begin{cases} \text{SO}(6) \sim \text{SU}(4) \text{ R symmetry} \\ \text{Sp}(4) \text{ Conformal symmetry} \end{cases}$
- AdS/CFT: IIA String theory $\text{AdS}_4 \times \text{CP}_3$ /ABJM



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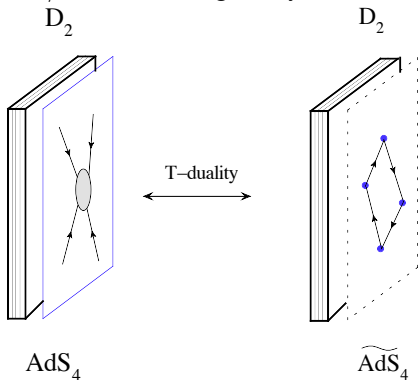
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T-duality invariant?

NO I. Adam, A. Dekel, Y. Oz (09)(10), P. A. Grassi, D. Sorokin, L. Wulff,(09),
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The On-shell Variables

Three-dimensional kinematics

$$SL(2, \mathbb{R}) : \quad p_i^{(\alpha\beta)} = \lambda_i^\alpha \lambda_i^\beta$$

$$\langle ij \rangle \equiv \lambda_i^\alpha \lambda_{j\alpha}, \quad \langle ij \rangle^2 = -2p_i \cdot p_j$$

Supersymmetry: $\mathcal{N} = 6$, η^I , $I \in SU(3)$

$$\mathcal{A}_n(\lambda_i^\alpha, \eta^I)$$

$$\Phi(\eta) = \phi^4 + \eta^I \psi_I + \frac{1}{2} \epsilon_{IJK} \eta^I \eta^J \phi^K + \frac{1}{3!} \epsilon_{IJK} \eta^I \eta^J \eta^K \psi_4$$

$$\Psi(\eta) = \bar{\psi}^4 + \eta^I \bar{\phi}_I + \frac{1}{2} \epsilon_{IJK} \eta^I \eta^J \bar{\psi}^K + \frac{1}{3!} \epsilon_{IJK} \eta^I \eta^J \eta^K \bar{\phi}_4$$

General n -point amplitude

- \mathcal{A}_n : n =even
- \mathcal{A}_n : R-symmetry invariance $\rightarrow \eta^{\frac{3}{2}n}$
- \mathcal{A}_n : Supersymmetry invariance.
 $\rightarrow \mathcal{A}_n = \delta^3(P)\delta^3(Q^\alpha)\delta^3(Q_\alpha)f_n(\lambda, \eta)$

$$\mathcal{A}_4 = \frac{\delta^3(P)\delta^3(Q^\alpha)\delta^3(Q_\alpha)}{\langle 12 \rangle \langle 23 \rangle}, \quad f_4 = \frac{1}{\langle 12 \rangle \langle 23 \rangle}$$

Is it Dual conformal?

$$p_i = x_i - x_{i+1}, \quad I[x_i] = \frac{x_i}{x_i^2}$$

$$\rightarrow I[f_4] = \sqrt{x_1^2 x_2^2 x_3^2 x_4^2} f_4$$

Explicit computation $\mathcal{A}_4 \mathcal{A}_6 \rightarrow \text{OSp}(6|4)$ Yangian invariance [T. Bargheer, F. Loebert, C. Meneghelli (10)]

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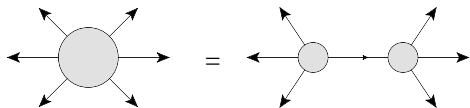
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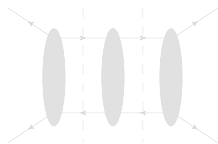
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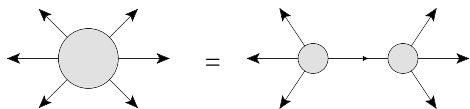


$$\rightarrow I[I_n^L] = (x_{\ell_1}^2 \cdots x_{\ell_L}^2)^3 \prod_i^n \sqrt{x_i^2} I_n^L$$

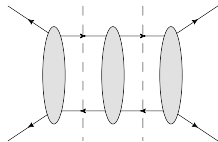
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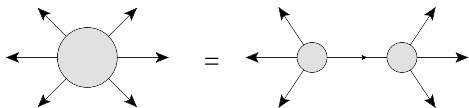


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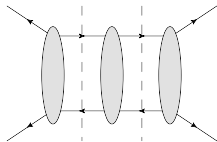
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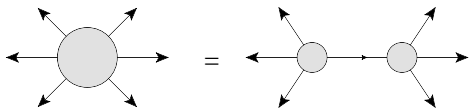


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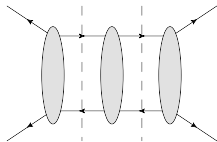
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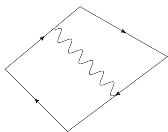
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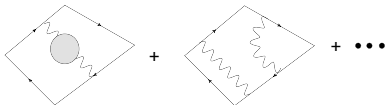
ABJM Wilson-loops

1-loop:



$$\langle W_4^1 \rangle = 0$$

2-loop:



$$\langle W_4^2 \rangle = 1 - \left(\frac{N}{K} \right)^2 \left[\frac{(-\mu'^2 x_{13}^2)^{-2\epsilon}}{(2\epsilon)^2} + \frac{(-\mu'^2 x_{24}^2)^{-2\epsilon}}{(2\epsilon)^2} - \frac{1}{2} \log^2 \left(\frac{-x_{13}^2}{-x_{24}^2} \right) + \text{const} + \mathcal{O}(\epsilon) \right]$$

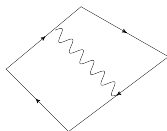
N=4 SYM

$$\langle W_4^1 \rangle = 1 + \left(\frac{g^2 N}{8\pi^2} \right)^2 \left[\frac{(-\mu'^2 x_{13}^2)^{-\epsilon}}{(\epsilon)^2} + \frac{(-\mu'^2 x_{24}^2)^{-\epsilon}}{(\epsilon)^2} - \frac{1}{2} \log^2 \left(\frac{-x_{13}^2}{-x_{24}^2} \right) + \text{const} + \mathcal{O}(\epsilon) \right]$$

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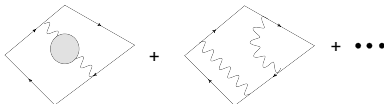
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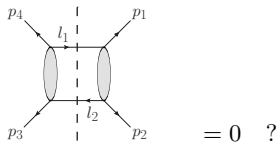
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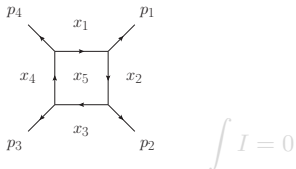
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Amplitude (W-M Chen, Y. Huang [11])

1-loop:

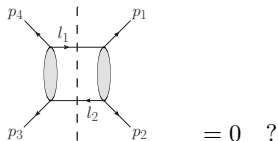


$$I = \frac{2x_{51}^2 \epsilon_{\mu\nu\rho} x_{21}^\mu x_{31}^\nu x_{41}^\rho + 2x_{31}^2 \epsilon_{\mu\nu\rho} x_{51}^\mu x_{21}^\nu x_{41}^\rho}{x_{15}^2 x_{25}^2 x_{35}^2 x_{45}^2}$$

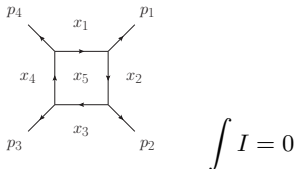


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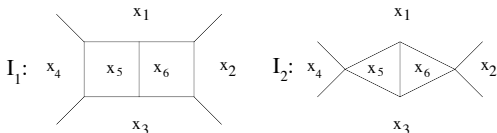


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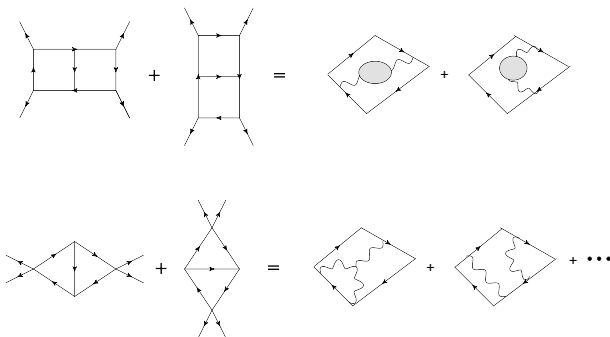
2-loop:



$$I_1 = \frac{4\xi_s}{x_{51}^2 x_{53}^2 x_{54}^2 x_{56}^2 x_{61}^2 x_{62}^2 x_{63}^2 x_{24}^2}$$

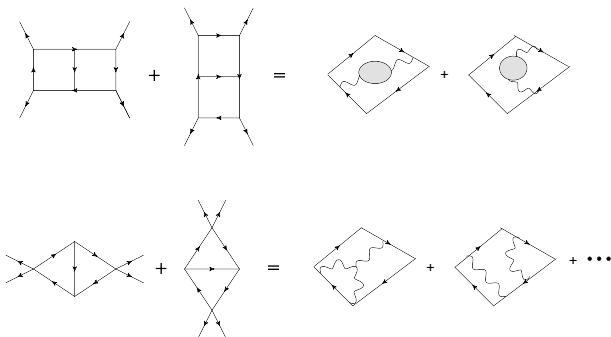
$$I_2 = \frac{x_{13}^4}{x_{51}^2 x_{53}^2 x_{56}^2 x_{61}^2 x_{63}^2}$$

$$\xi_s \equiv \epsilon_{\mu\nu\rho} (x_{51}^2 x_{21}^\mu x_{31}^\nu x_{41}^\rho + x_{31}^2 x_{51}^\mu x_{21}^\nu x_{41}^\rho) \\ \epsilon_{\gamma\sigma\eta} (x_{61}^2 x_{21}^\gamma x_{31}^\sigma x_{41}^\eta + x_{31}^2 x_{61}^\gamma x_{21}^\sigma x_{41}^\eta)$$



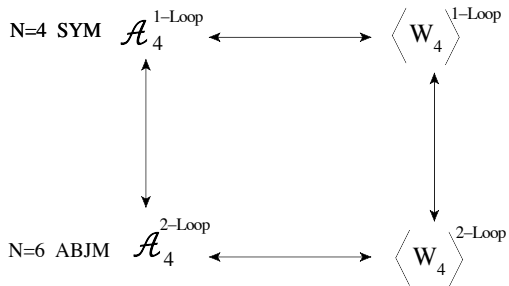
$$\text{ABJM: } \mathcal{A}_4^{2-loop} = \frac{1}{16\pi^2} \left(\frac{N}{K}\right)^2 \mathcal{A}_4^{tree} \left[-\frac{(-s/\tilde{\mu}^2)^{-2\epsilon}}{(2\epsilon)^2} - \frac{(-t/\tilde{\mu}^2)^{-2\epsilon}}{(2\epsilon)^2} + \frac{1}{2} \log^2 \left(\frac{-s}{-t}\right) + a + \mathcal{O}(\epsilon) \right]$$

$$\mathcal{N} = 4 \quad \mathcal{A}_4^{1-loop} = \frac{ig^2 N}{8\pi^2} \mathcal{A}_4^{tree} \left[-\frac{(-s/\mu^2)^{-\epsilon}}{\epsilon^2} - \frac{(-t/\mu^2)^{-\epsilon}}{\epsilon^2} + \frac{1}{2} \log^2 \left(\frac{-s}{-t}\right) + 4\zeta_2 + \mathcal{O}(\epsilon) \right]$$

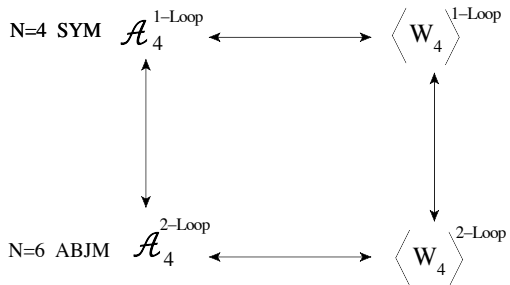


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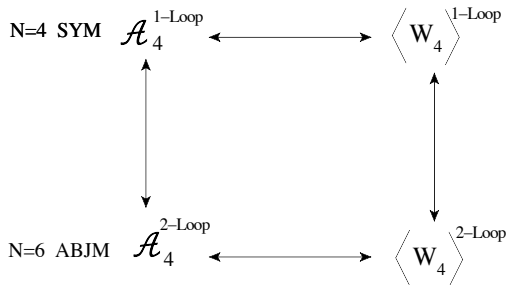


- (1) Higher loops ? (BDS)
- (2) Higher points ? (N^n MHV)

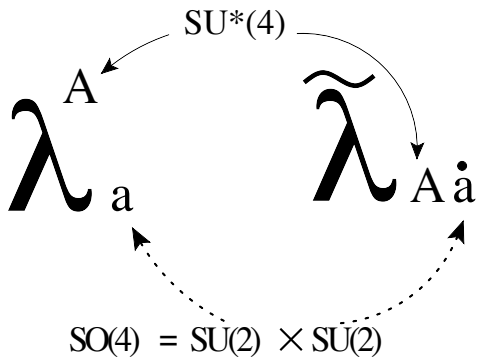


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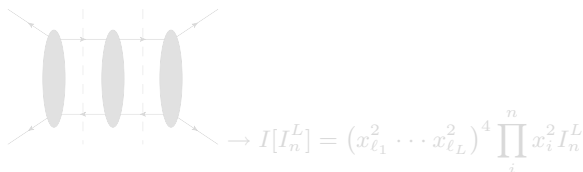
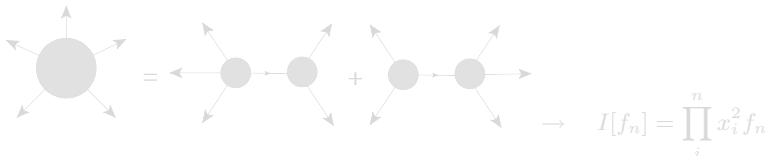


$\mathcal{N} = (1, 1) \text{ SYM}$ [J. J. Carrasco, T. Dennen, Y. Huang, H. Ita, W. Siegel, (10)(11)]

$$\mathcal{A}_4 = \delta^6 \left(\sum_{i=1}^4 p_i \right) \delta^4 \left(\sum_{i=1}^4 q_i^A \right) \delta^4 \left(\sum_{i=1}^4 \tilde{q}_{iA} \right) \frac{i}{st}$$

$$f_4 = \frac{i}{st} = \frac{i}{x_{13}^2 x_{24}^2}$$

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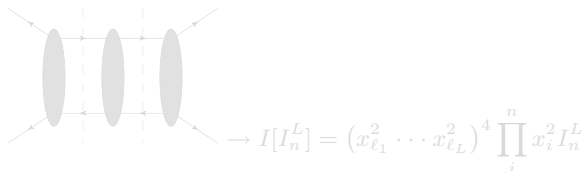
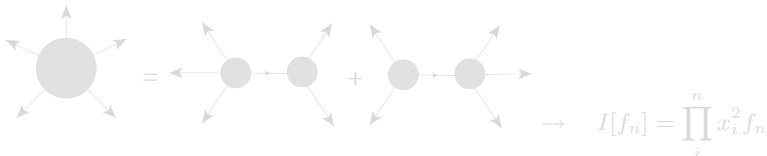


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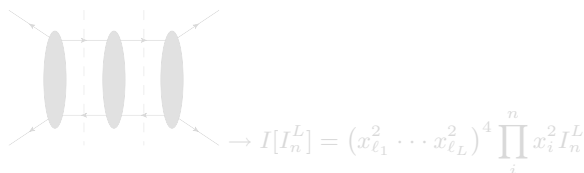
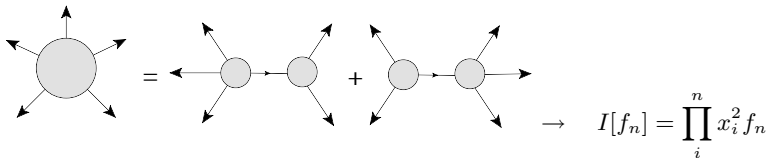


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$$\mathcal{A}_4 = \delta^6 \left(\sum_{i=1}^4 p_i \right) \delta^4 \left(\sum_{i=1}^4 q_i^A \right) \delta^4 \left(\sum_{i=1}^4 \tilde{q}_{iA} \right) \frac{i}{st}$$

$$f_4 = \frac{i}{st} = \frac{i}{x_{13}^2 x_{24}^2}$$

$$I[f_4] = x_1^2 x_2^2 x_3^2 x_4^2 f_4$$

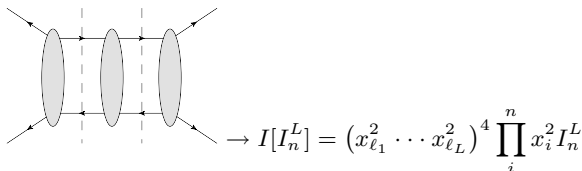
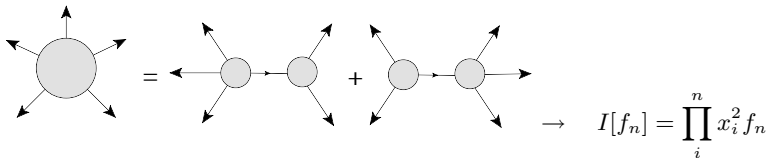


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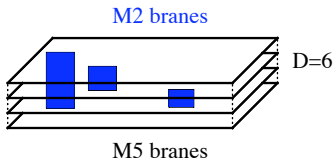
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$\mathcal{N} = (2, 0)$ Self-Dual Tensor [B. Czech, Y. Huang, M. Rozali(11)]

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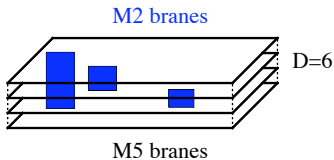
$$\rightarrow (B_{\mu\nu}, \phi^I), \quad I \in SO(5)$$

$$H_{\mu\nu\rho} = \partial_{[\mu} B_{\nu\rho]}, \quad \delta B_{\mu\nu} = \partial_{[\mu} \Lambda_{\nu]}$$

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- A SCFT $OSp(8|4)$
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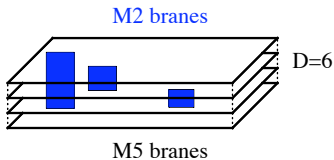
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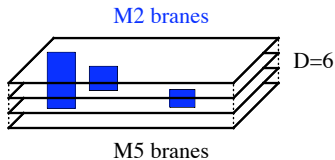
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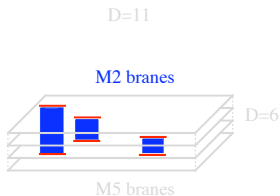
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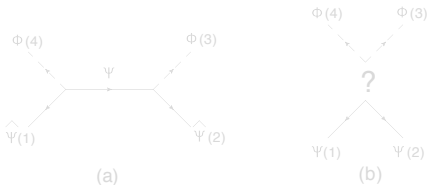
The search for $\mathcal{A}_3(1)$

- $\mathcal{N} = (2, 0), (1, 0)$ SUSY \rightarrow No pure tensor-interaction Wrong degrees of freedom?



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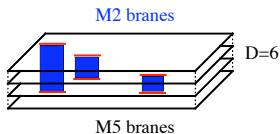
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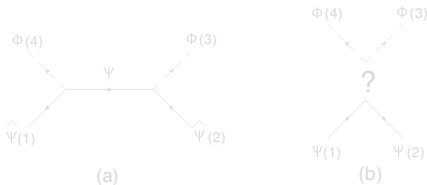
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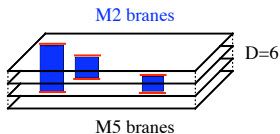
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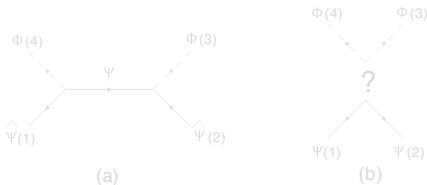
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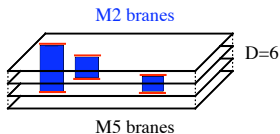
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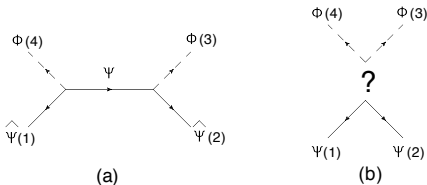
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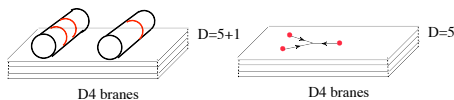
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The search for $\mathcal{A}_3(\text{II})$

Lets go to D=5: $SU(4) \rightarrow USp(2,2)$

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We find

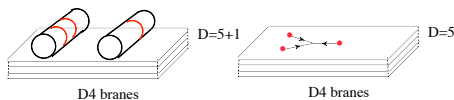
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- $\mathcal{A}_3(123) = -\mathcal{A}_3(132) \rightarrow f^{aij} = -f^{aji}$
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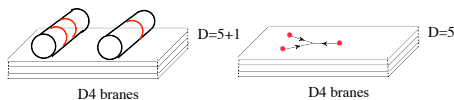
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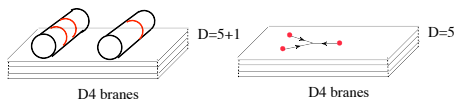
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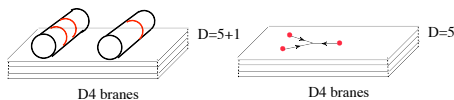
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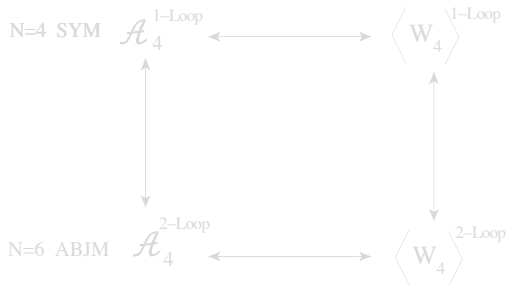
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Conclusion

- Dual superconformal symmetry of ABJM

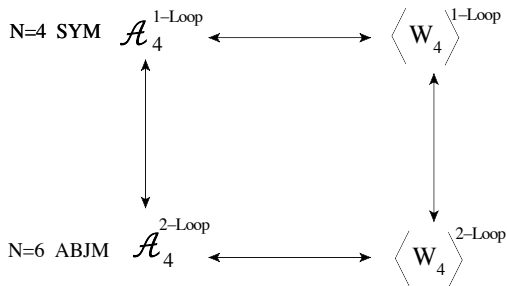


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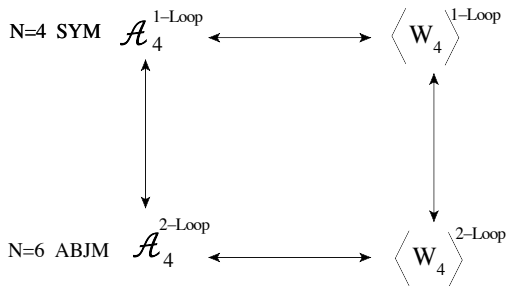


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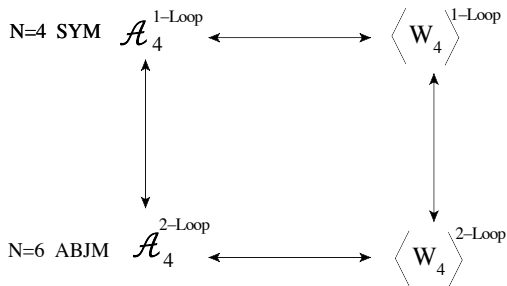


- Dual conformal symmetry of non-conformal theory (6D SYM)
- Gone where no action has gone before. First $\mathcal{A}_3, \mathcal{A}_4$ for M -theory.

Conclusion

- Dual superconformal symmetry of ABJM

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- Dual conformal symmetry of non-conformal theory (6D SYM)
- Gone where no action has gone before. First $\mathcal{A}_3, \mathcal{A}_4$ for M -theory.