

Fusion Reactions at very low energies.

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Interface between structure and reactions for rare isotopes and astrophysics.

- The description is based on the coupled-channels technique. It includes couplings to the low-lying 2^+ and 3^- states in projectile and target, and mutual and two-phonon excitations of these states.
- The influence of transfer is also studied,
 - for example, in the fusion of $^{40}\text{Ca}+^{48}\text{Ca}$, $^{12}\text{C}+^{13}\text{C}$, and $^{13}\text{C}+^{13}\text{C}$.
- The analysis of measurements has focused on understanding
 - the hindrance of fusion at extreme subbarrier energies,
 - fusion reactions of interest to astrophysics (e. g. $^{12}\text{C}+^{12}\text{C}$),
 - the constraints on the extrapolation to very low energies.

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Coupled equations.

$$(H_n^{(0)} + E_n - E)\phi_n(r) = - \sum_{n' \neq n} \langle n | \delta V_C + \delta V_N | n' \rangle \phi_{n'}(r).$$

Nuclear interaction of the form $V_N(r - R_1 - R_2 - \delta R)$, where

$$\delta R = \sum_{\lambda\mu} \alpha_{1\lambda\mu} R_1 Y_{\lambda\mu}^*(\hat{r}) + \alpha_{2\lambda\mu} R_2 Y_{\lambda\mu}^*(-\hat{r}).$$

$\alpha_{\lambda\mu}$: dynamic (or static) surface deformation amplitudes.

Rotating frame approximation: $\hat{r} = \hat{z}$,

$$Y_{\lambda\mu}(\hat{r}) = \delta_{\mu,0} \sqrt{\frac{2\lambda + 1}{4\pi}}.$$

Implies that the magnetic quantum number is conserved.
There is only one channel for each state with spin I ,
instead of $2I + 1$ or $I + 1$ channels.

Standard approach.

Include Coulomb couplings δV_C to first order,
and nuclear couplings up to second order in δR ,

$$\delta V_N = -U'(r) \delta R + \frac{1}{2} U''(r) \left(\delta R^2 - \langle 0 | \delta R^2 | 0 \rangle \right),$$

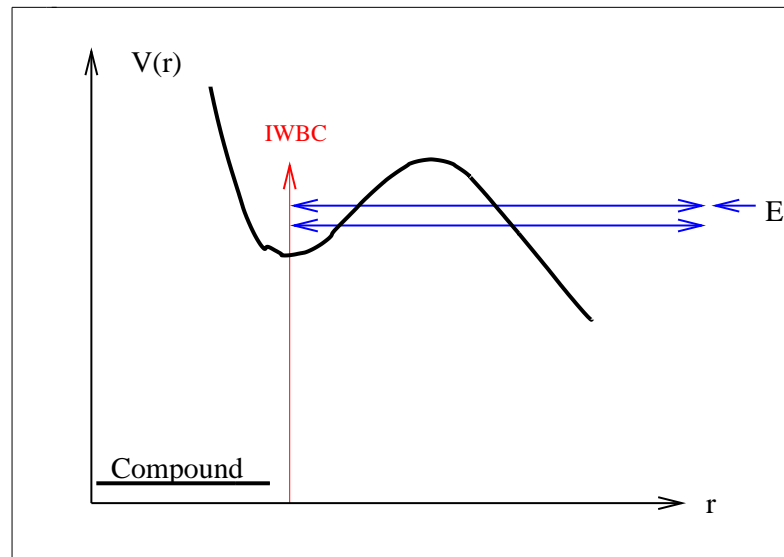
where $U(r)$ is the conventional ion-ion potential

$$U(r) = \frac{16\pi\gamma a R_1 R_2}{R_1 + R_2} \frac{1}{1 + \exp[(r - R_1 - R_2)/a]}.$$

A Proximity type interaction by Broglia and Winther.

Usual scattering boundary
conditions at large value of r .
FUSION is determined by **IWBC**
(ingoing-wave boundary conditions)
that are imposed for overlapping nuclei.

Avoid imaginary potentials if possible.



Standard two-phonon calculation.



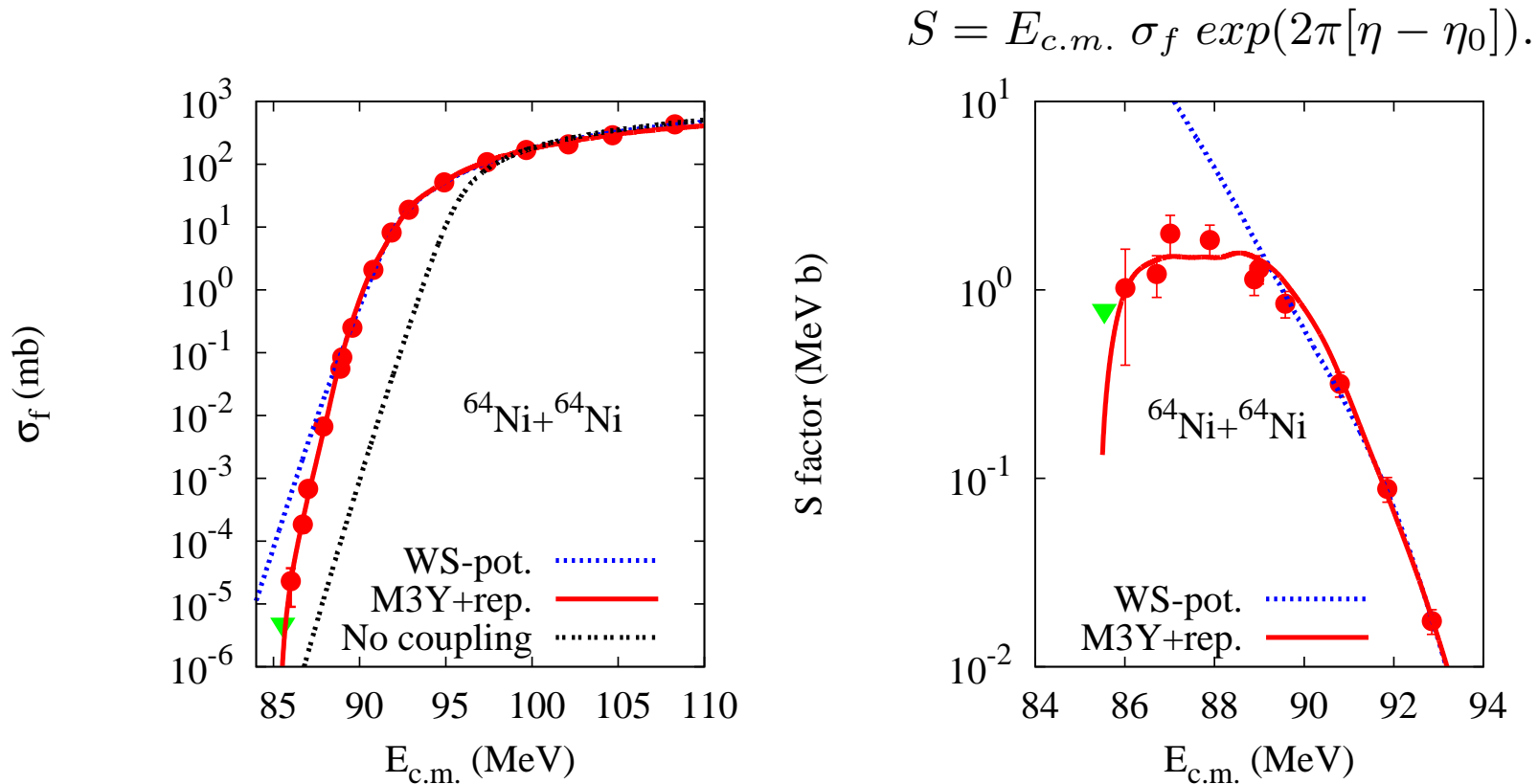
No. of channels: 1 (GS) + 4 (1PH) + 4 (2PH) + 6 (Mutual) = 15 channels.

Works quite well for lighter and medium heavy systems, with some exceptions,

- in fusion reactions where transfer plays a role,
- for heavy and soft, or strongly deformed nuclei,
 - become sensitive to high-lying states (multiphonon or high spin states.)

Hindrance of fusion far below the Coulomb barrier.

- Was first recognized at Argonne by Jiang et al., $^{60}\text{Ni}+^{89}\text{Y}$, PRL 89 (2002), $^{64}\text{Ni}+^{64}\text{Ni}$, PRL 93 (2004), $^{64}\text{Ni}+^{100}\text{Mo}$, PRC 71 (2005), $^{28}\text{Si}+^{64}\text{Ni}$, PL B640 (2006).



- The data are *hindered* compared to coupled-channels calculations that are based on *a conventional Woods-Saxon potential*.
- Right panel: *the S factor for fusion has a maximum at low energy.*

Theoretical description of the fusion hindrance,

Mișicu and Esbensen, PRL 96, 112701 (2006); PRC 75, 034606 (2007).

Conventional Woods-Saxon potential (WS), $a \approx 0.65$ fm.

M3Y double-folding potential:

$$U(\mathbf{r}) = \int d\mathbf{r}_1 d\mathbf{r}_2 \rho_1(\mathbf{r}_1) \rho_2(\mathbf{r}_2) \times v_{NN}(\mathbf{r} + \mathbf{r}_2 - \mathbf{r}_1). v_{NN} =$$

M3Y effective NN interaction.

M3Y+repulsion potential:

Supplement with a repulsive term,

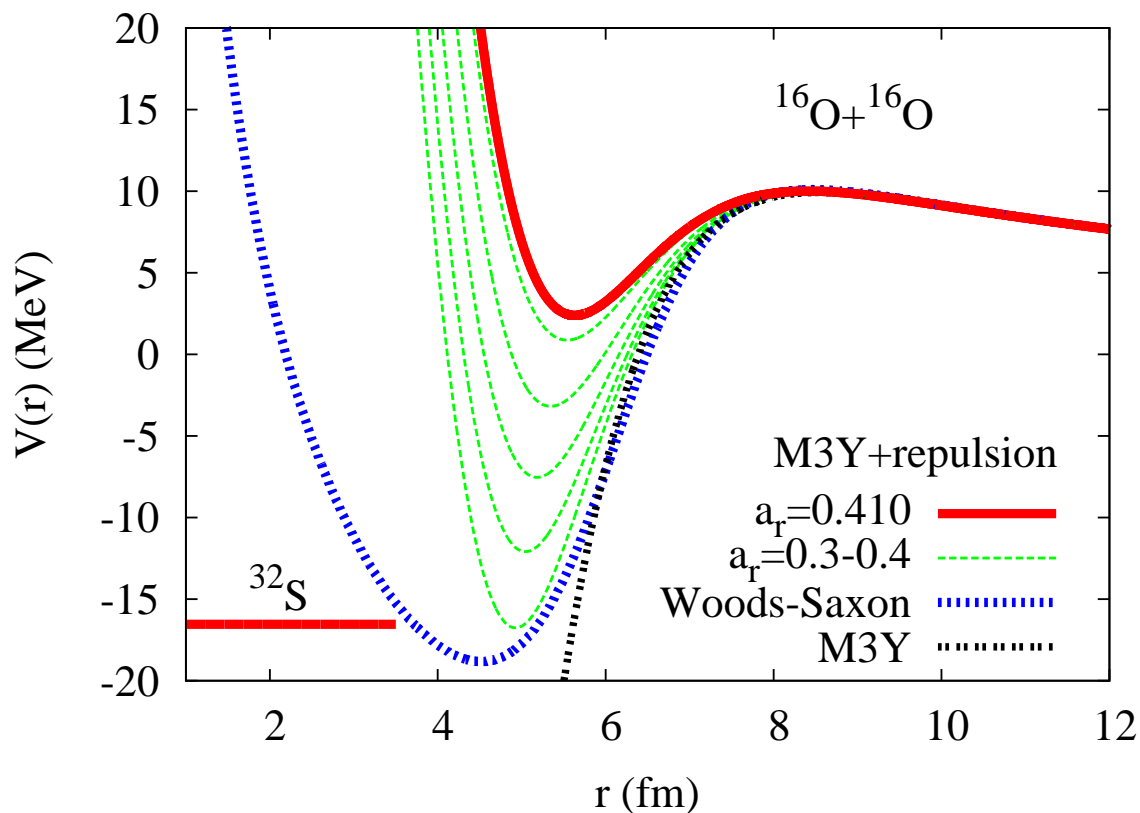
$$v_{NN}^{\text{rep}} = v_r \delta(\mathbf{r} + \mathbf{r}_2 - \mathbf{r}_1).$$

Use $\rho(r)$ with adjustable diffuseness a_r .

v_r is calibrated to give the nuclear

incompressibility $K = 234$ MeV.

Adjustable parameters: R and a_r . Can produce a shallow pocket and a thicker barrier.

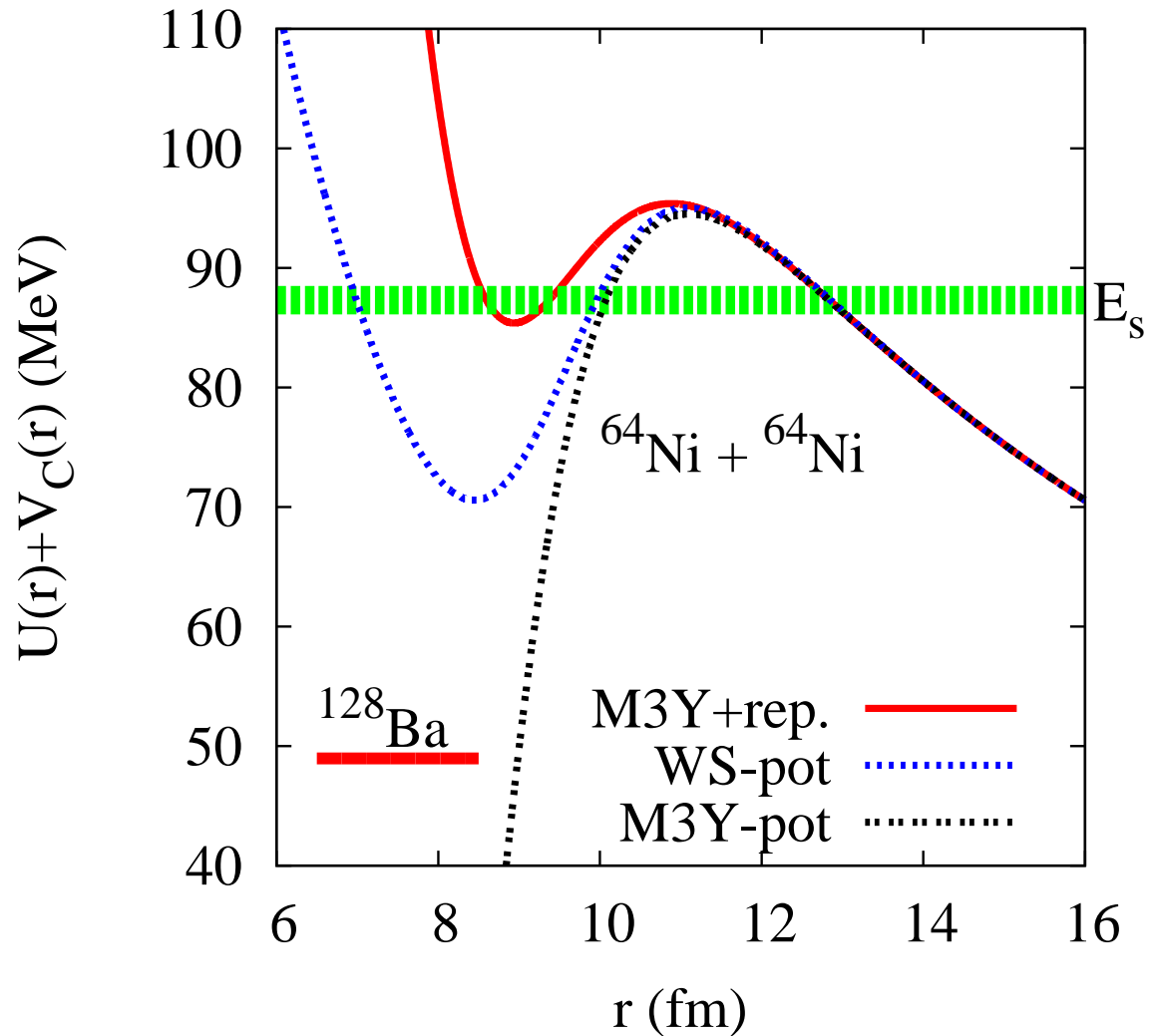


Hindrance of fusion far below the Coulomb barrier.

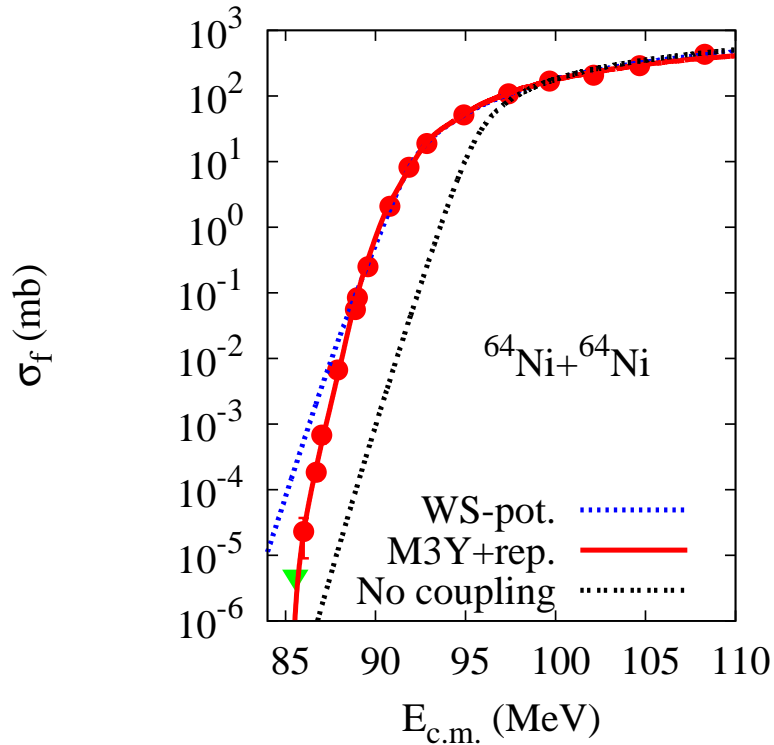
The hindrance
sets in near E_s .

Adjust a_r so pocket
is slightly deeper:
 $a_r = 0.403$ fm.

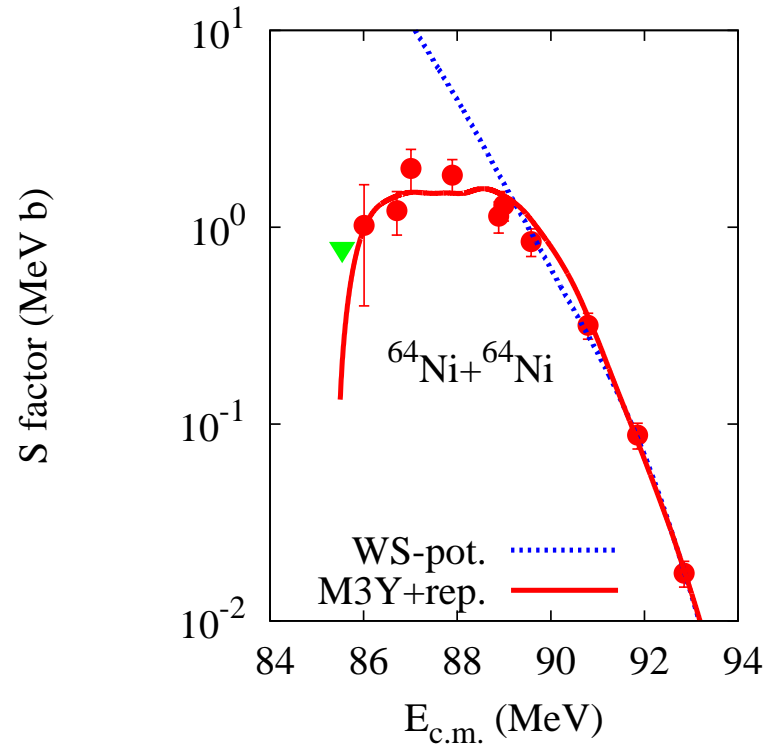
Fusion cross section
goes to zero for
 $E_{cm} \leq E_{pocket}$.



Jiang et al., $^{64}\text{Ni}+^{64}\text{Ni}$, PRL 93 (2004).



$$S = E_{c.m.} \sigma_f \exp(2\pi[\eta - \eta_0]).$$



- The data are *hindered* compared to coupled-channels calculations that are based on *a conventional Woods-Saxon potential*.
- Strong hindrance (right panel):
 - *the S factor for fusion has a maximum* at low energy.

Does fusion hindrance occur in systems with positive Q values?
Important issue for extrapolations in astrophysics.

- Confirmed in new experiments at Argonne:

$^{28}\text{Si}+^{30}\text{Si}$, PRC 78, 17601 (2008),

$^{27}\text{Al}+^{45}\text{Sc}$, PRC 81, 24611 (2010).

Almost a maximum S factor.

- $^{36}\text{S}+^{48}\text{Ca}$, $Q_{gg} = +7.6$ MeV.

Stefanini et al.,

PRC 78, 044607 (2008).

Hindrance does occur but not

always with a maximum S factor.

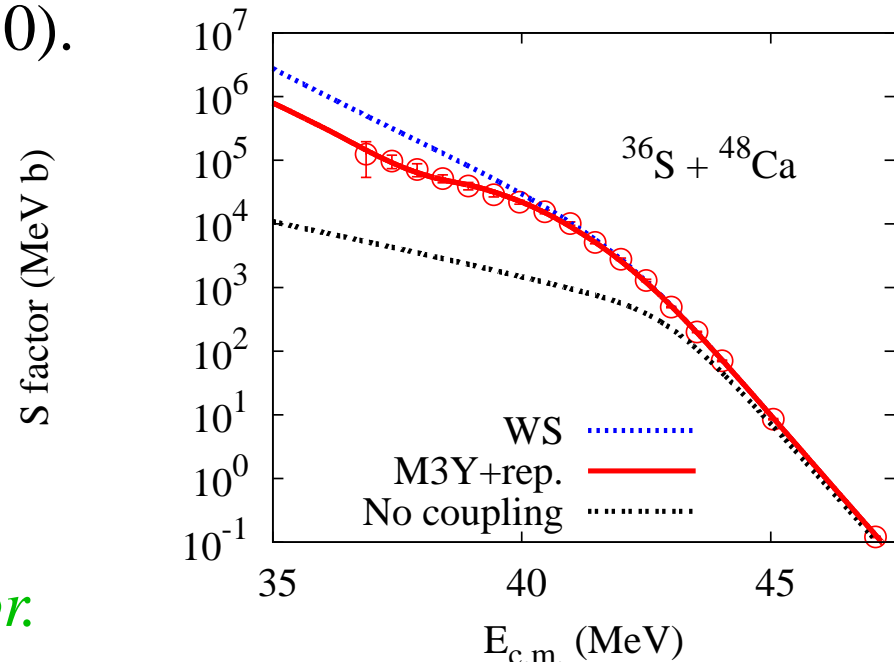
- Systematic study of fusion with calcium isotopes:

$^{48}\text{Ca}+^{48}\text{Ca}$, Stefanini et al., PLB 679, 95 (2009),

$^{40}\text{Ca}+^{48}\text{Ca}$, Jiang et al. (incl. Esbensen), PRC 82 (2010),

$^{40}\text{Ca}+^{40}\text{Ca}$, Stefanini et al. (incl. Esbensen, unpublished.)

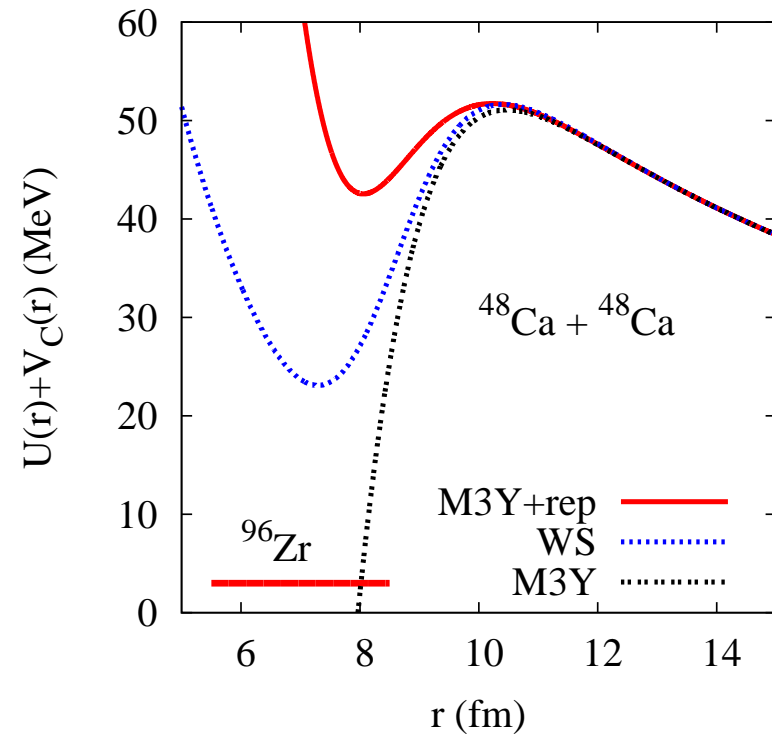
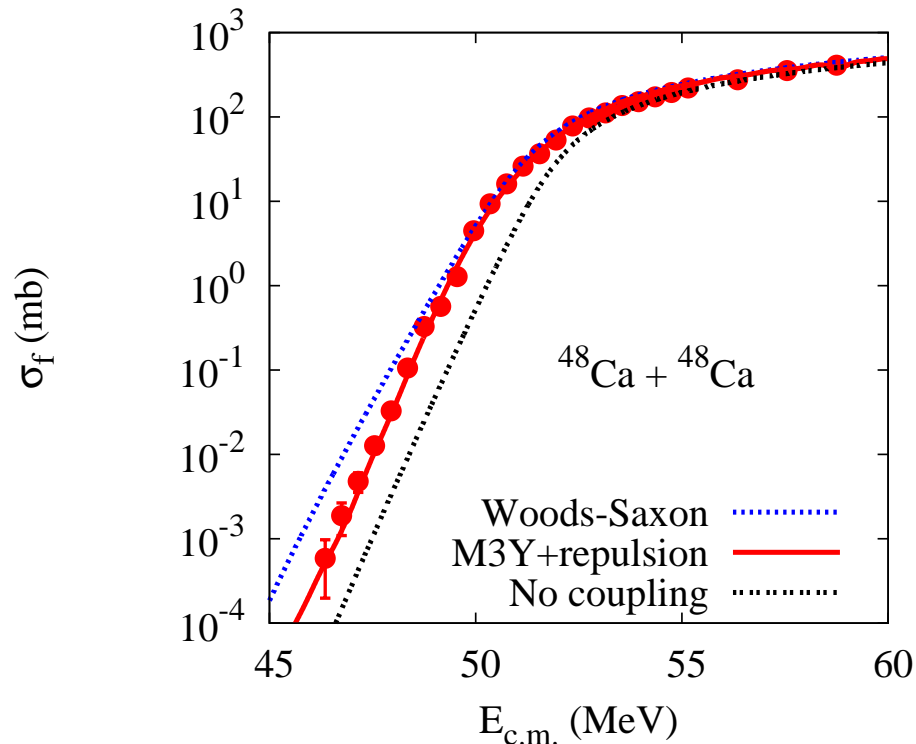
- Previously studied by Aljuwair et al. (84-Exp), and Esbensen et al. (89-Theo.)



Analysis of $^{48}\text{Ca}+^{48}\text{Ca}$ fusion experiment,

Esbensen, Jiang, and Stefanini, PRC 82, 054621 (2010).

Coupling to the low-lying 2^+ , 3^- , 5^- states, and mutual excitations.



Fusion is a sensitive probe of the surface of nuclei.

Best fit density of ^{48}Ca has the rms-radius = 3.56 fm.

Compare to point-protons: 3.39 fm, point-nucleons: 3.53 fm.

Extracted radius is model dependent (dynamic polarization effects.)

There are two solutions to the analysis.

$a_r = 0.4070$ fm:

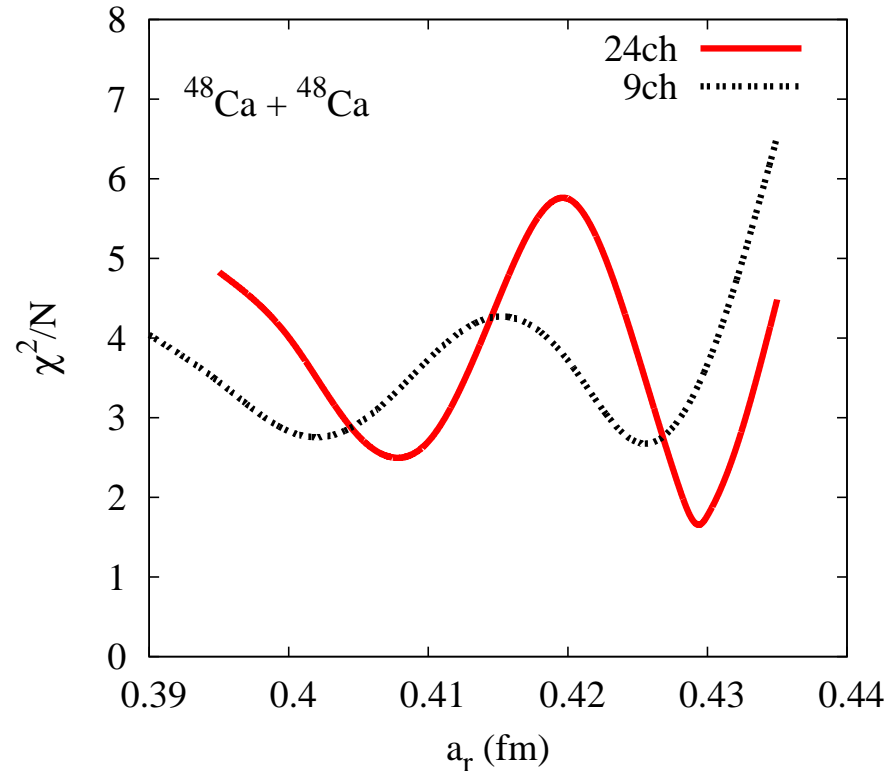
Ch9: $\langle r^2 \rangle^{1/2} = 3.55$ fm.

Ch24: $\langle r^2 \rangle^{1/2} = 3.53$ fm.

$a_r = 0.4295$ fm:

Ch9: $\langle r^2 \rangle^{1/2} = 3.57$ fm.

Ch24: $\langle r^2 \rangle^{1/2} = 3.56$ fm.



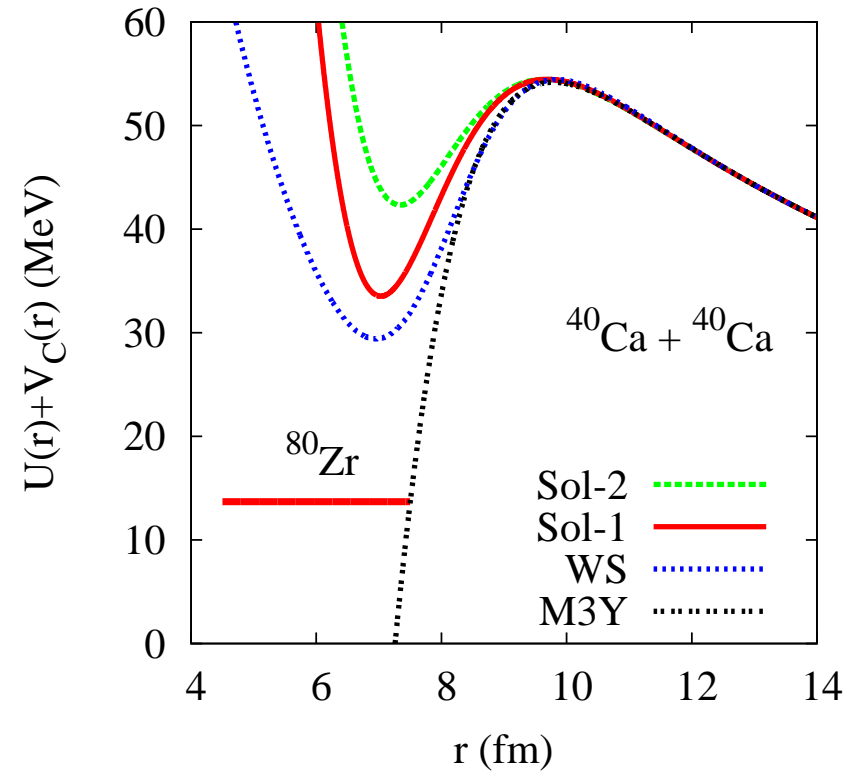
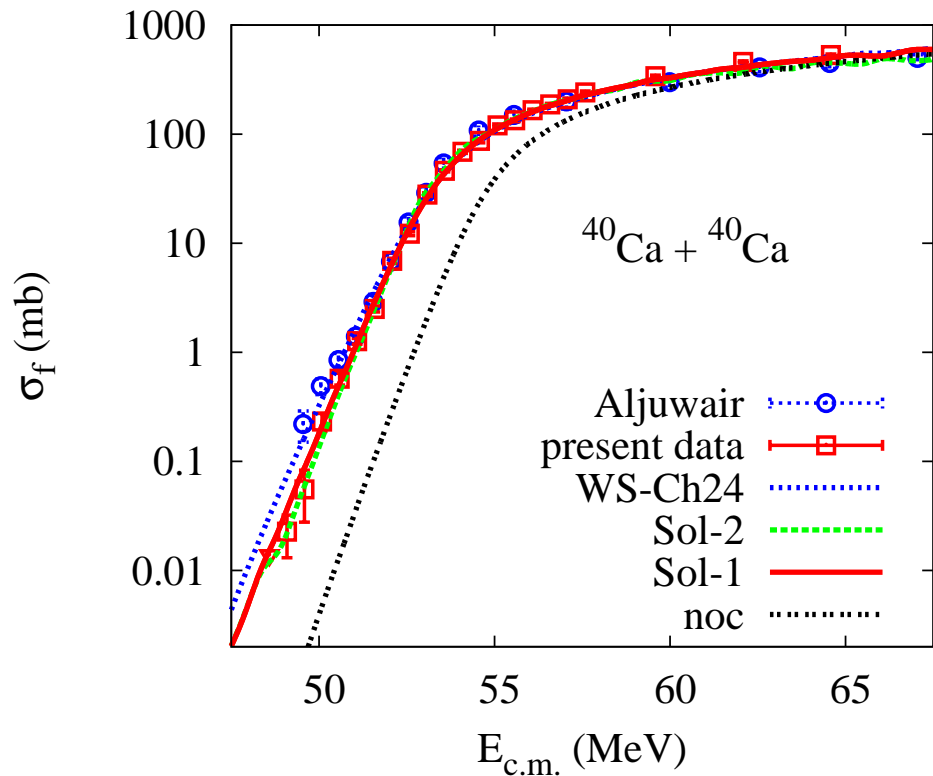
Compare: point matter rms = 3.53(3) fm.

Which solution should I choose? The one with the smallest χ^2/N .

Dynamic polarization effect: fewer channels implies a larger radius.

Similar analysis of $^{40}\text{Ca}+^{40}\text{Ca}$ experiment.

Coupling to the low-lying 2^+ , 3^- , 5^- states, and mutual excitations.



Also two solutions of the analysis.

Stefanini et al. (unpublished)

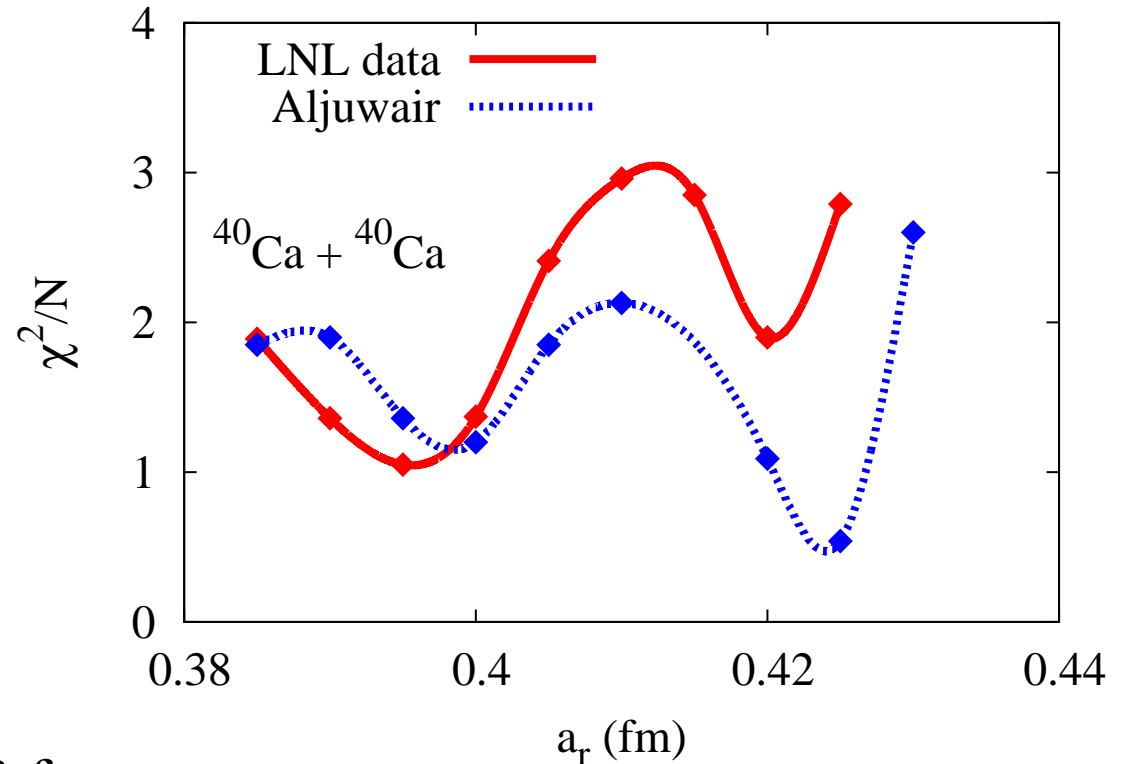
Also two solutions:

$a_r = 0.395$ fm:

$\langle r^2 \rangle^{1/2} = 3.38$ fm.

$a_r = 0.42$ fm:

$\langle r^2 \rangle^{1/2} = 3.41$ fm.



Compare point-protons: 3.38 fm,

Which solution should I choose?

A larger radius is justified by the dynamic polarization effect.

Table 1: Nuclear structure input for ^{40}Ca and ^{48}Ca . The $B(E\lambda)$ values are from ENDSF, Nat. Nuclear Data Center (BNL). Values marked with * are from Ref. [1]. Eff 2PH: effective parameters for the $(0^+, 2^+, 4^+)$ two-phonon quadrupole excitations.

Nucleus	I^π	E_x (MeV)	$B(E\lambda)$ (W.u.)	$\frac{(\beta R)_C}{\sqrt{4\pi}}$ (fm)	$\frac{(\beta R)_N}{\sqrt{4\pi}}$ (fm)
^{40}Ca	2_1^+	3.905	2.26(14)	0.138*	0.125*
	3^-	3.737	27(4)	0.465*	0.315*
	5^-	4.491		0.344*	0.175*
	Eff 2PH	5.269	41(5)	0.416	0.416
^{48}Ca	2_1^+	3.832	1.71(9)	0.126*	0.190*
	3^-	4.507	5.0(8)	0.250*	0.190*
	5^-	5.146		0.049*	0.038*
	Eff 2PH	4.849	4.7(29)	0.15	0.15

[1] Esbensen and Videbaek, PRC 40, 126 (1989), analysis of $^{16}\text{O}+^A\text{Ca}$ scattering data.

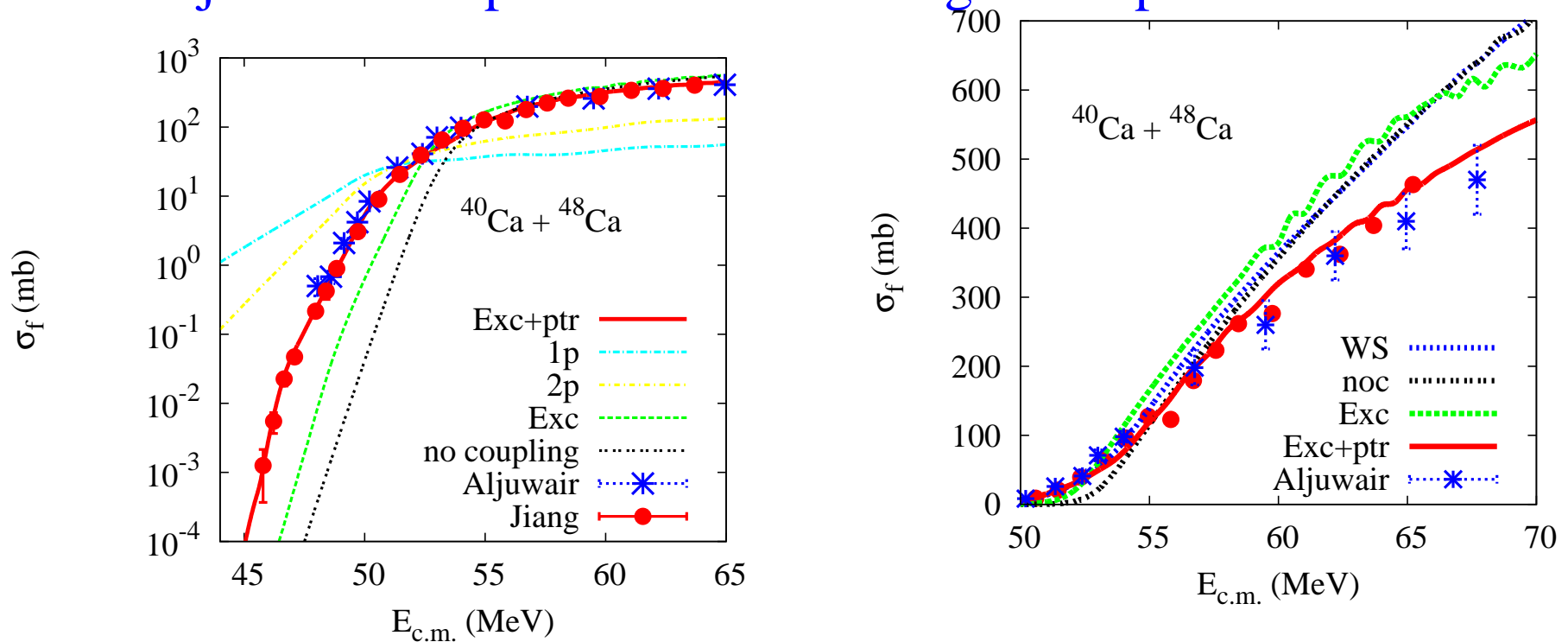
Note that ^{40}Ca is much softer than ^{48}Ca .

Coulomb and nuclear parameters are not (always) identical.

$^{40}\text{Ca} + ^{48}\text{Ca}$ experiment, Jiang et al, PRC 82, 041601 (2010).

Include one- and two-proton transfers with positive Q-values.

Adjust the two-proton transfer strength to reproduce data.



Two-proton transfer gives a large enhancement of subbarrier fusion and has "a good effect" (suppression) on fusion at high energies.

Systematics of the fusion of calcium isotopes.

$^{48}\text{Ca}+^{48}\text{Ca}$, Stefanini et al.
Phys. Lett. B679, 95 (2009).

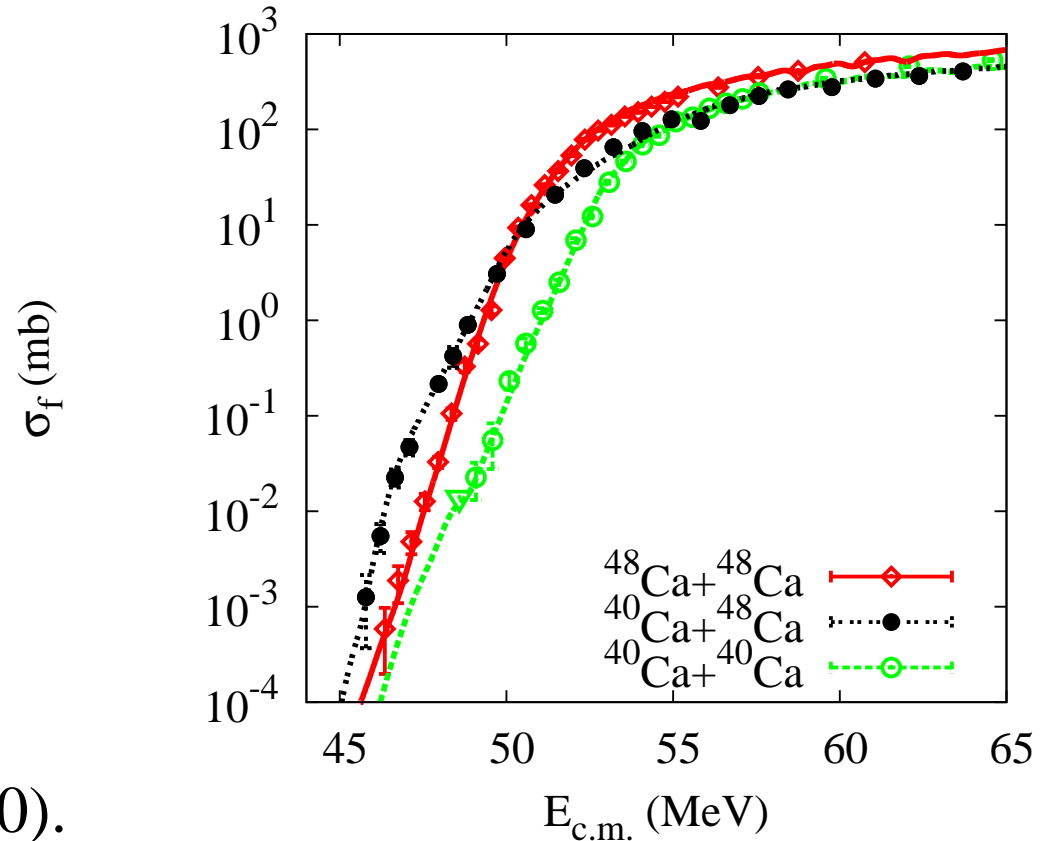
$^{40}\text{Ca}+^{40}\text{Ca}$, Stefanini et al.
(unpublished.) ^{40}Ca is soft.

$^{40}\text{Ca}+^{48}\text{Ca}$, Jiang et al,
Phys. Rev. C 82, 041601 (2010).

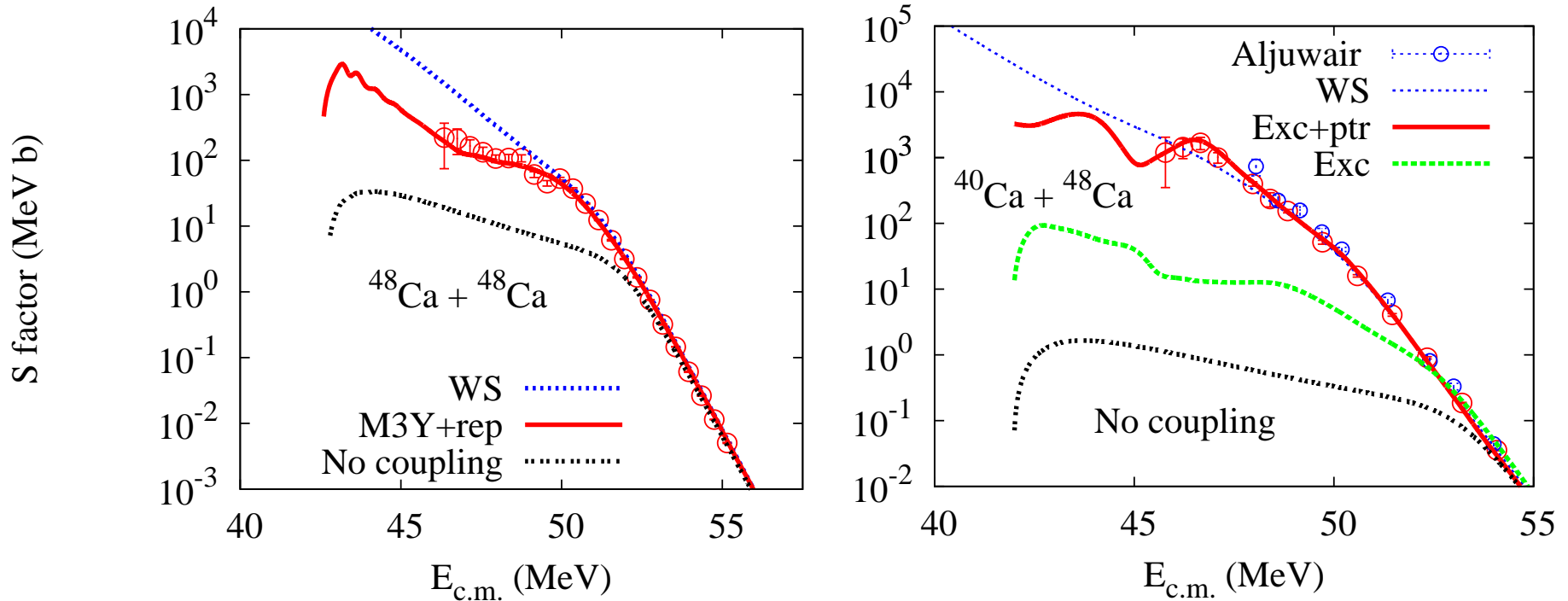
Enhanced compared to the $^{48}\text{Ca}+^{48}\text{Ca}$ data.

*Large effect of one- and two-proton transfer with positive Q -values
on the fusion of the asymmetric $^{40}\text{Ca}+^{48}\text{Ca}$ system.*

The hindrance sets in at a lower energy.



S factors for $^{48}\text{Ca}+^{48}\text{Ca}$ and $^{40}\text{Ca}+^{48}\text{Ca}$ fusion.



Relatively weak coupling for $^{48}\text{Ca}+^{48}\text{Ca}$: Hindrance occurs.

Strong coupling effects in the fusion of $^{40}\text{Ca}+^{48}\text{Ca}$
due to two-proton transfer and a softer ^{40}Ca .

The hindrance sets in at a lower energy.

Coupled-channels calculations of $^{12}\text{C}+^{12}\text{C}$ fusion,
H. Esbensen, X. Tang, and C. L. Jiang.

The goal is to put constraints on the extrapolation

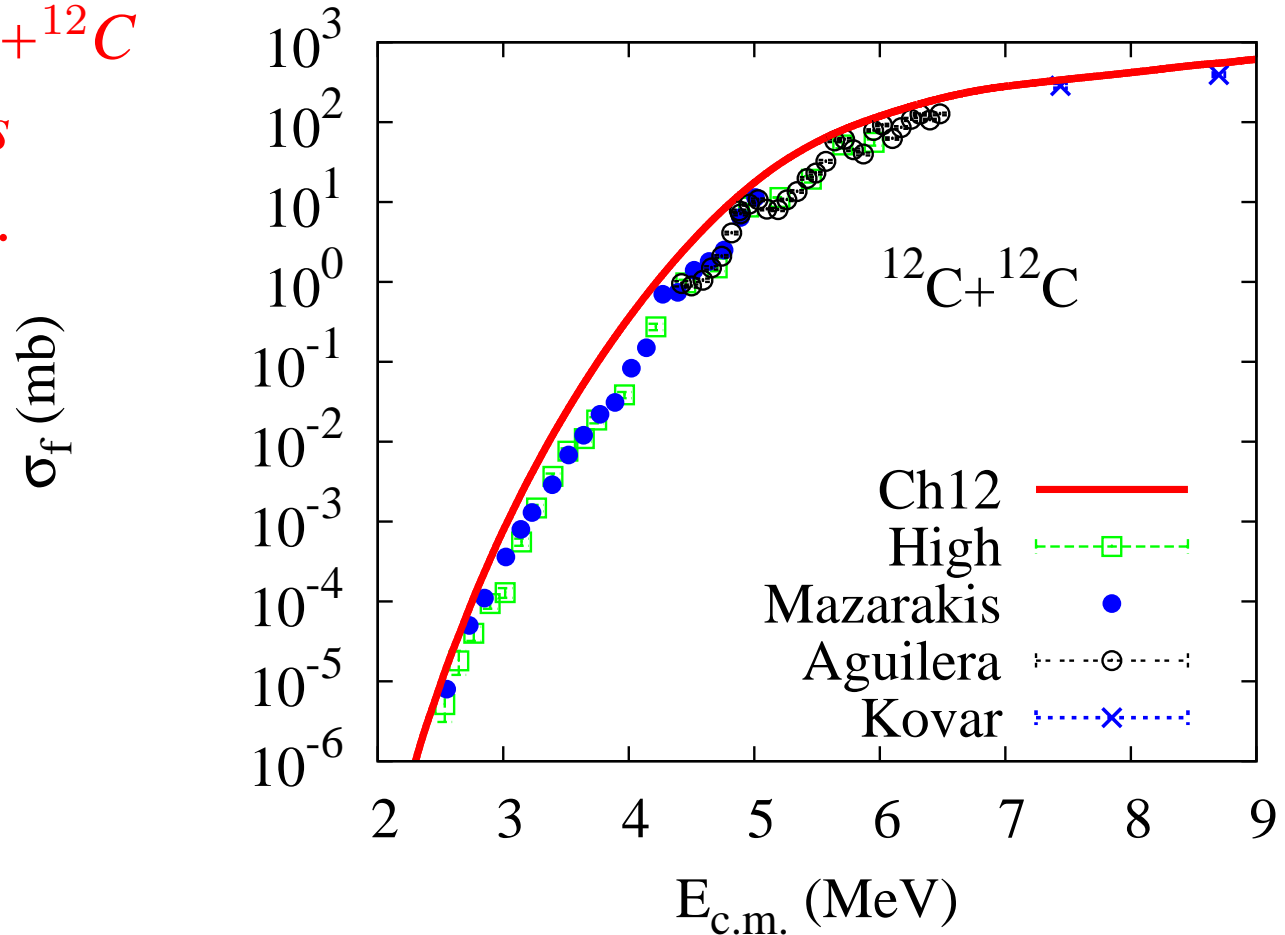
*of the measured $^{12}\text{C}+^{12}\text{C}$
fusion cross sections
to very low energies.*

Problems:
Structures
(peaks) in data.

Large experimental
uncertainties
up to 30%.

Closed channels:

use decaying states when $E_{c.m.} \leq E_x$ (Whittaker function.)



Entrance channel potential.

^{12}C density = *point-proton density*. $R = 1.696$ fm, $a = 0.52$ fm.

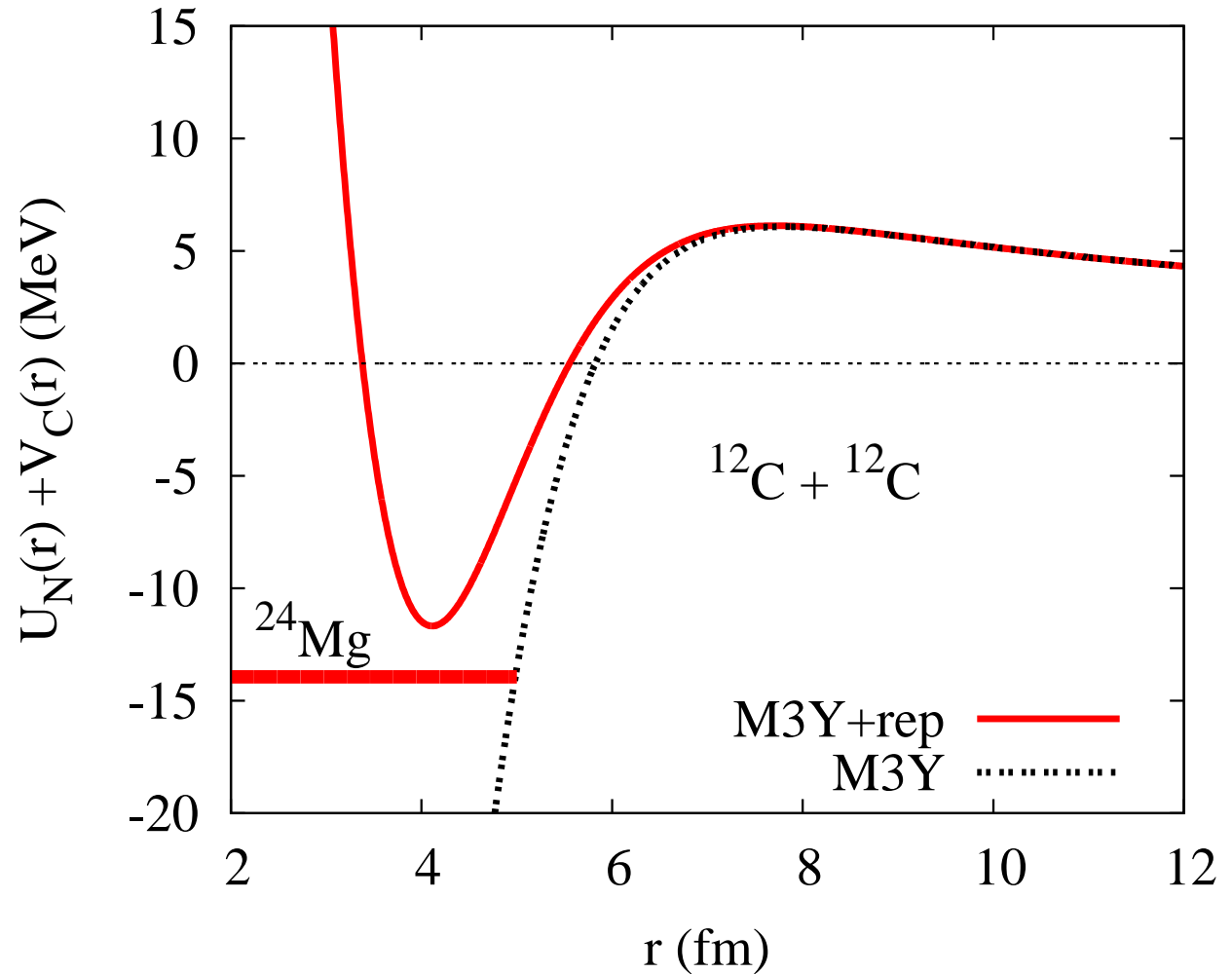
Repulsion:

$K=234$ MeV,

$a_r = 0.426$ fm.

$V_{CB} = 6.12$ MeV.

$V_{min} = -11.67$ MeV.



Nuclear structure input.

Table 2: Properties of E2 and E3 transitions to the low-lying states in ^{12}C . The intrinsic quadrupole moment was extracted from the lowest E2 transition.

Nucleus	State	E_x (MeV)	Transition	$B(E\lambda)$ (W.u.)	β_λ^C
^{12}C	0_1^+	0		$Q_0 = -19.5 \text{ fm}^2$	0.570
	2^+	4.439	E2: $0_1^+ \rightarrow 2^+$	4.65(26)	0.570
	0_2^+	7.654	E2: $2^+ \rightarrow 0_2^+$	8.0(11)	0.236
	3^-	9.641	E3: $0_1^+ \rightarrow 3^-$	12(2)	0.90(7)

β_3 extremely large. Matrix element in ^{13}C is quenched: $\beta_3^{\text{eff}} = \sqrt{\frac{3}{7}} \beta_3$.

Coupling scheme.

Assume independent modes of excitation.

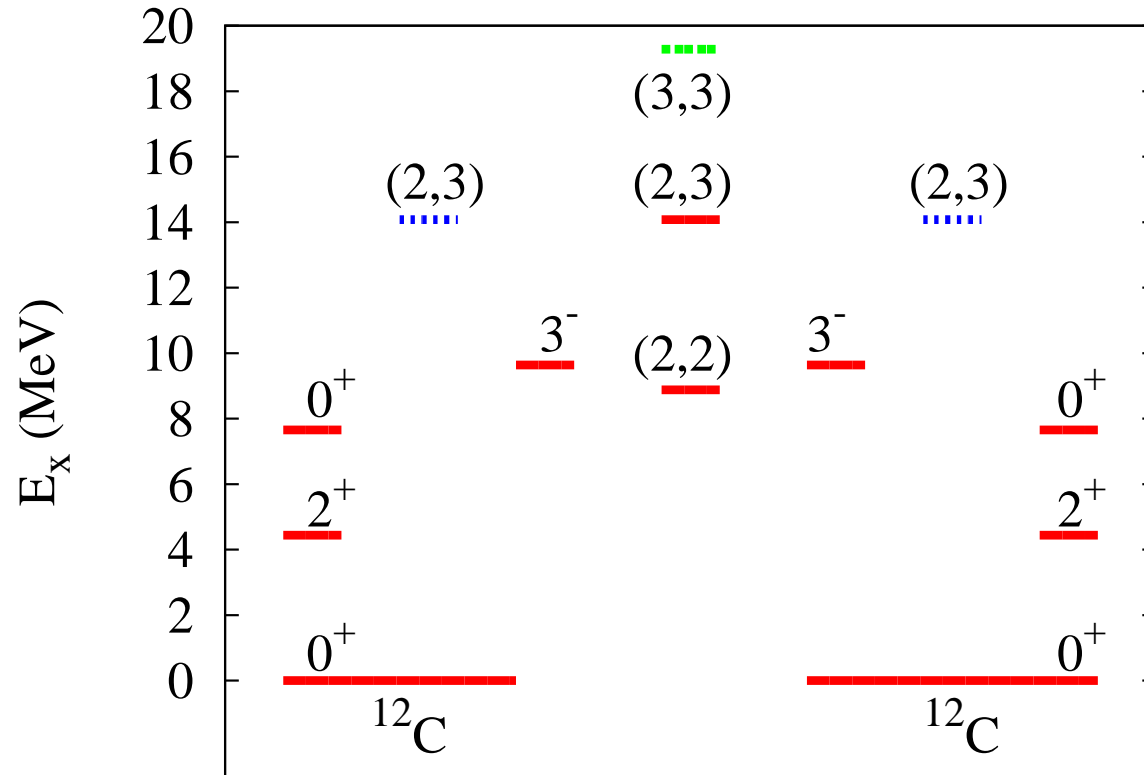
$$\beta_2 = 0.57$$

$$\beta_3 = 0.90(7)$$

Ch10: all red states

Ch12: add (2,3)

ph1: all one-phonon excitations (Ch5).



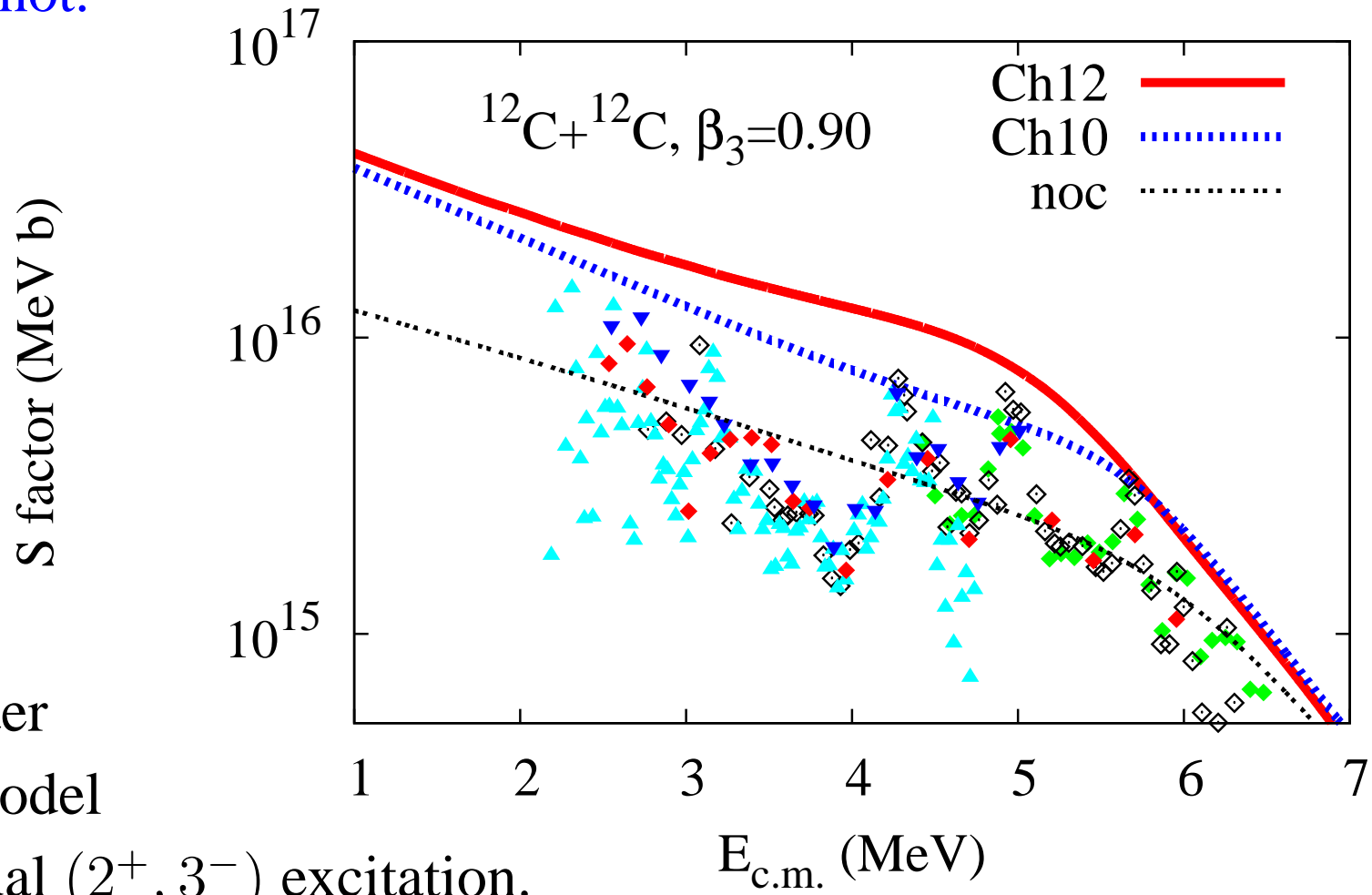
Ignore the $(3^-, 3^-)$ excitation. The energy is too high.

Model of mutual $(2^+, 3^-)$ excitation in same nucleus is questionable.

Sensitivity to the mutual $(2^+, 3^-)$ excitation in ^{12}C .

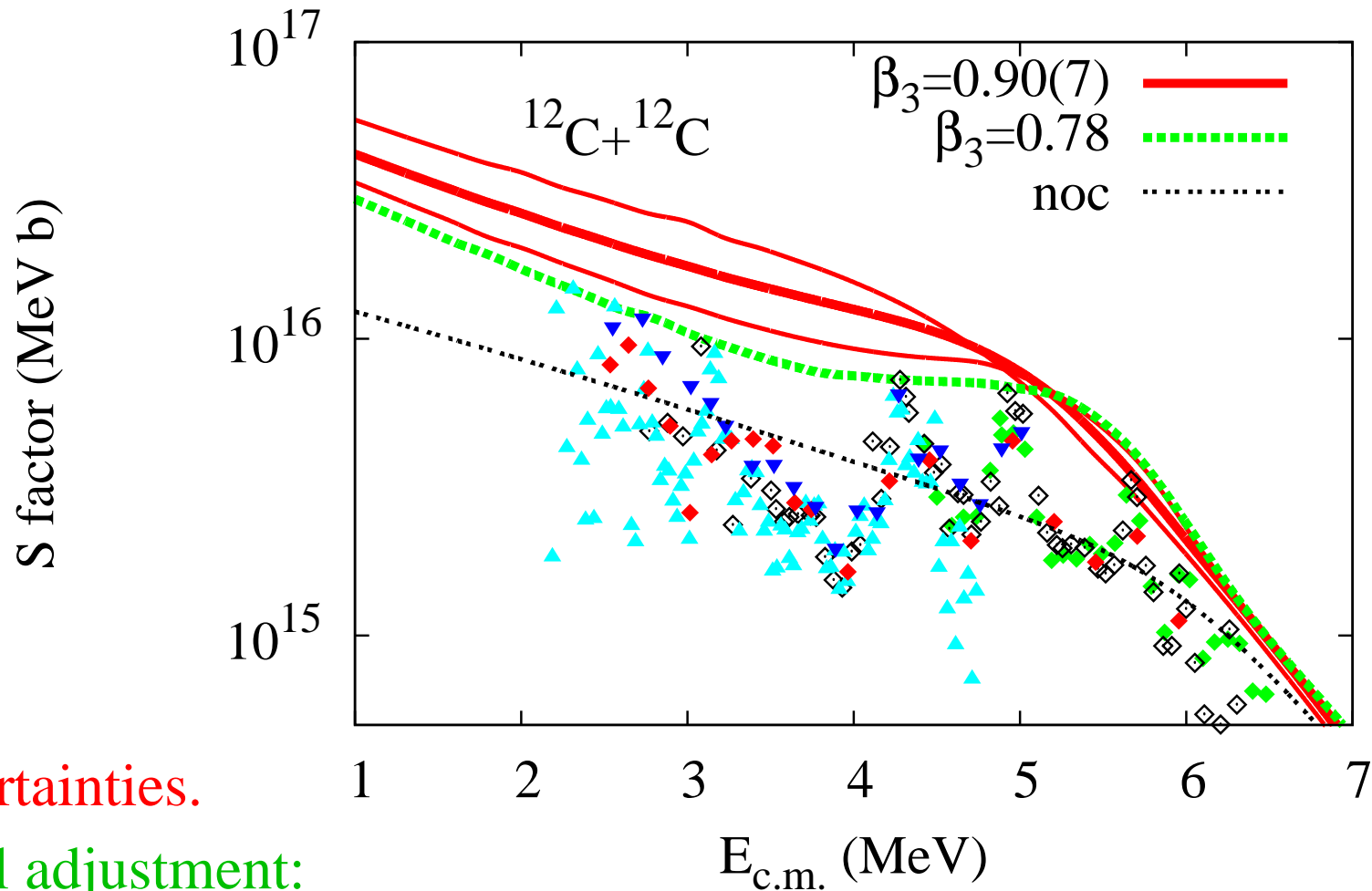
Ch12 includes this excitation, one in each nucleus.

Ch10 does not.



Need a better structure model of the mutual $(2^+, 3^-)$ excitation.

Sensitivity to $\beta_3 = 0.90 \pm 0.07$ in Ch12 calculations.



Large uncertainties.

Make small adjustment:

set $\beta_3 = 0.78$ to match the peaks.

Predict extrapolated value: $S(1 \text{ MeV}) = 3.0 \cdot 10^{16} \text{ MeV b}$.

Conclusions.

- The fusion hindrance at low energies is *a general phenomenon*
- the data are suppressed compared to CCC based on conventional WS potentials, - they have a steeper logarithmic slope.
- The data can be reproduced by adjusting the **M3Y+rep potential**.
Apart from structure, essentially only two parameters: R and a_r .
- **Extracted radii** can be slightly larger than expected
- due to dynamic polarization of states not included.
- **Large uncertainty** in the prediction of the $^{12}\text{C}+^{12}\text{C}$ fusion rate
- due to large β_3 -value and large sensitivity to $(2^+, 3^-)$ excitation.
- Conjecture: Coupled-channels calculations with **IWBC** should provide an **upper limit** for the $^{12}\text{C}+^{12}\text{C}$ fusion. **Structures in the data are due to the low level density of the compound system.**