

Reactions at Intermediate Energies

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Rutherford is (was) wrong

Aguiar, Aleixo, Bertulani, PRC 42, 2180 (1990)

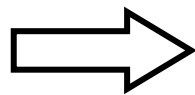
$$L = L^{LO} + L^{NLO} + L^{N^2LO} + \dots$$

$$L^{LO} = \frac{1}{2} \mu c^2 \left(\frac{v}{c} \right)^2 - \frac{Z_1 Z_2 e^2}{r}$$

$$L^{NLO} = \frac{\mu^4 c^2}{8} \left[\frac{1}{m_1^3} - \frac{1}{m_2^3} \right] \left(\frac{v}{c} \right)^4 - \frac{\mu^2 Z_1 Z_2 e^2}{2m_1 m_2 r} \left[\left(\frac{v}{c} \right)^2 + \left(\frac{\mathbf{v} \cdot \mathbf{r}}{cr} \right)^2 \right]$$

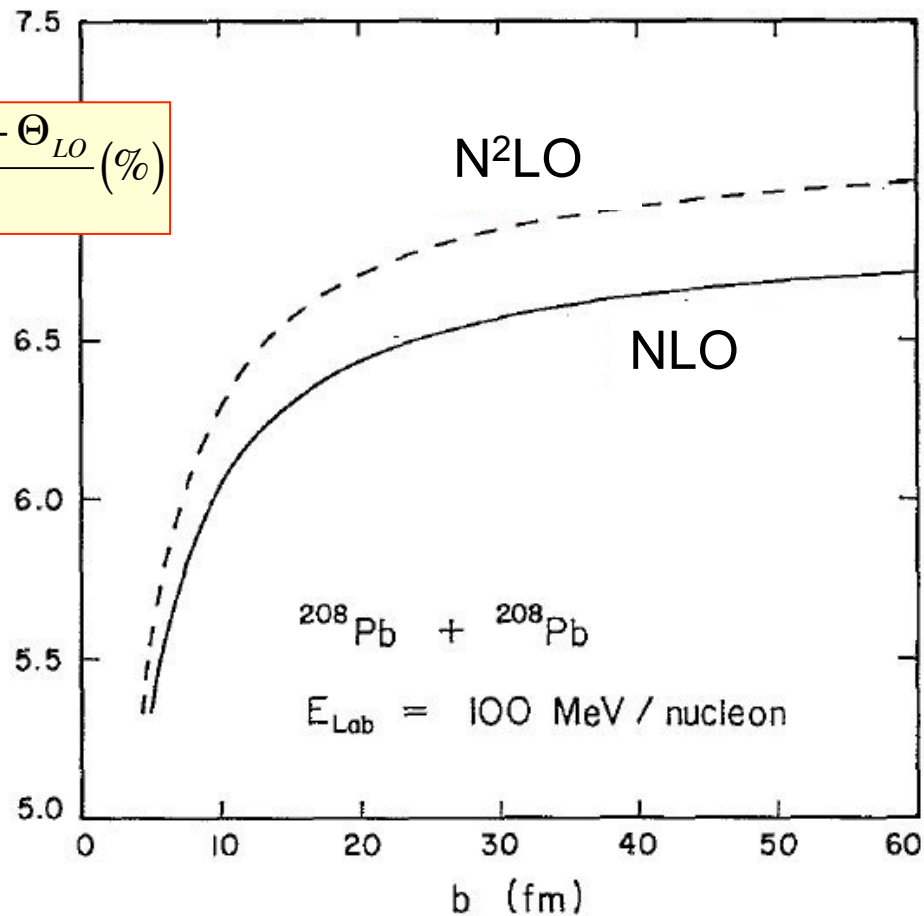
$$L^{N^2LO} = \frac{\mu c^2}{512} \left(\frac{v}{c} \right)^6 + \frac{Z_1 Z_2 e^2}{16r}$$

$$\times \left[\frac{1}{8} \left\{ \left(\frac{v}{c} \right)^4 - 3 \left(\frac{v_r}{c} \right)^4 + 2 \left(\frac{v_r v}{c} \right)^2 \right\} + \frac{Z_1 Z_2 e^2}{\mu c^2 r} \left\{ 3 \left(\frac{v_r}{c} \right)^2 - \left(\frac{v}{c} \right)^2 \right\} + \frac{4 Z_1^2 Z_2^2 e^4}{\mu^2 c^4 r^2} \right]$$



$$\frac{d\sigma}{d\Omega}$$

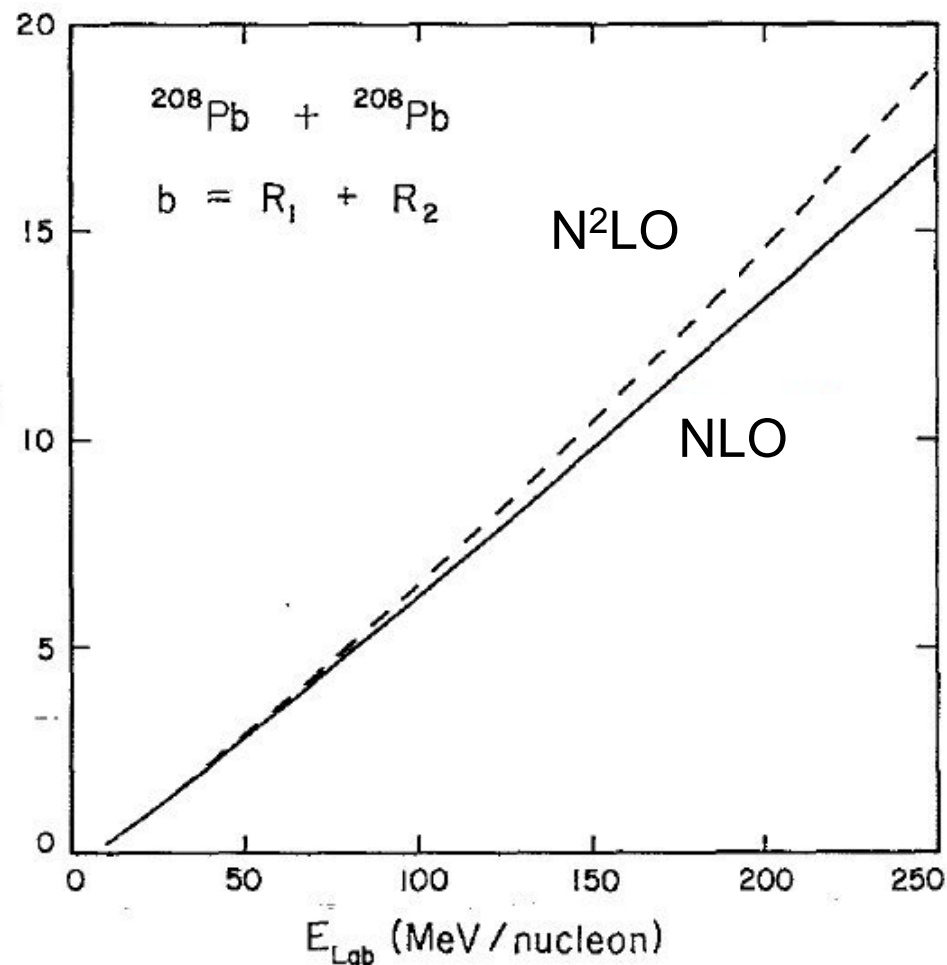
$$\frac{\Theta(E,b) - \Theta_{LO}}{\Theta_{LO}} (\%)$$

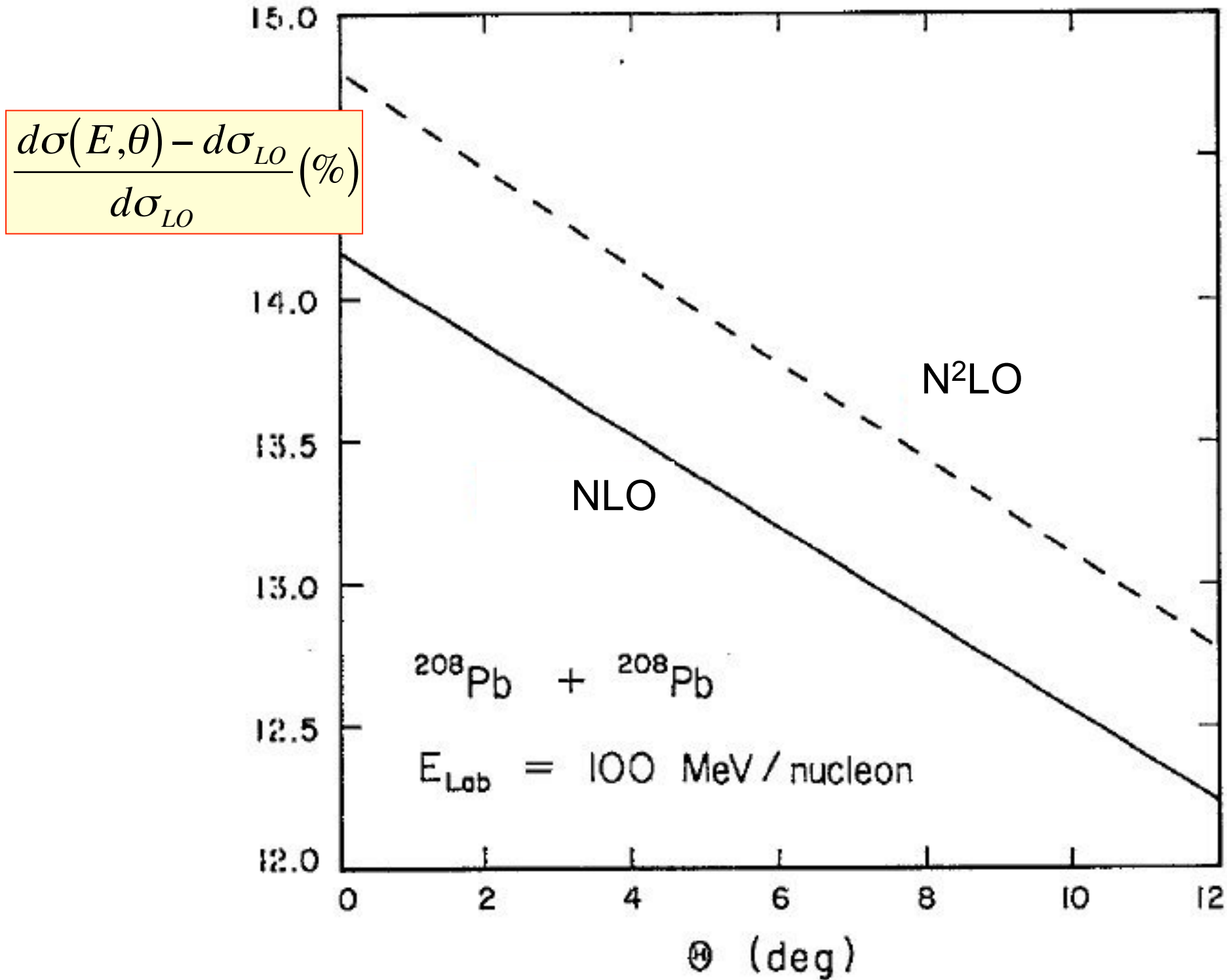


Deviations from Rutherford

important for elastic scattering:
 experimental data often reported
 as

$$\frac{d\sigma_{\text{elast}}}{d\sigma_{\text{Ruth}}}$$





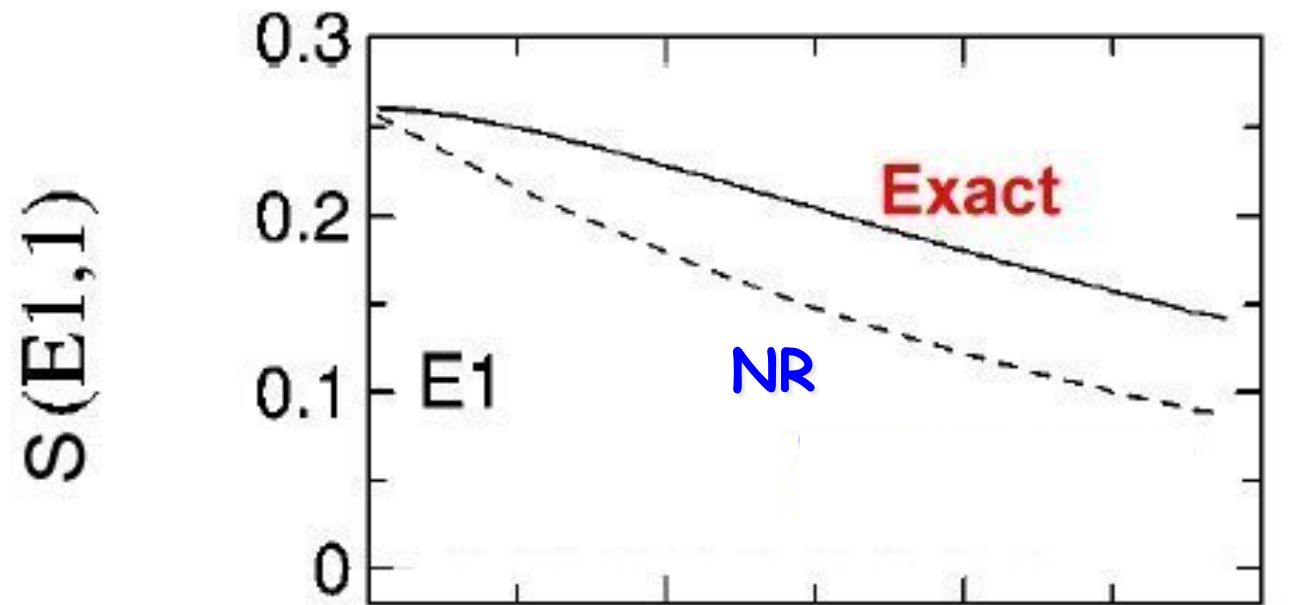
Coulomb excitation: orbital integrals with retardation

Aleixo, Bertulani, NPA 505, 448 (1989)

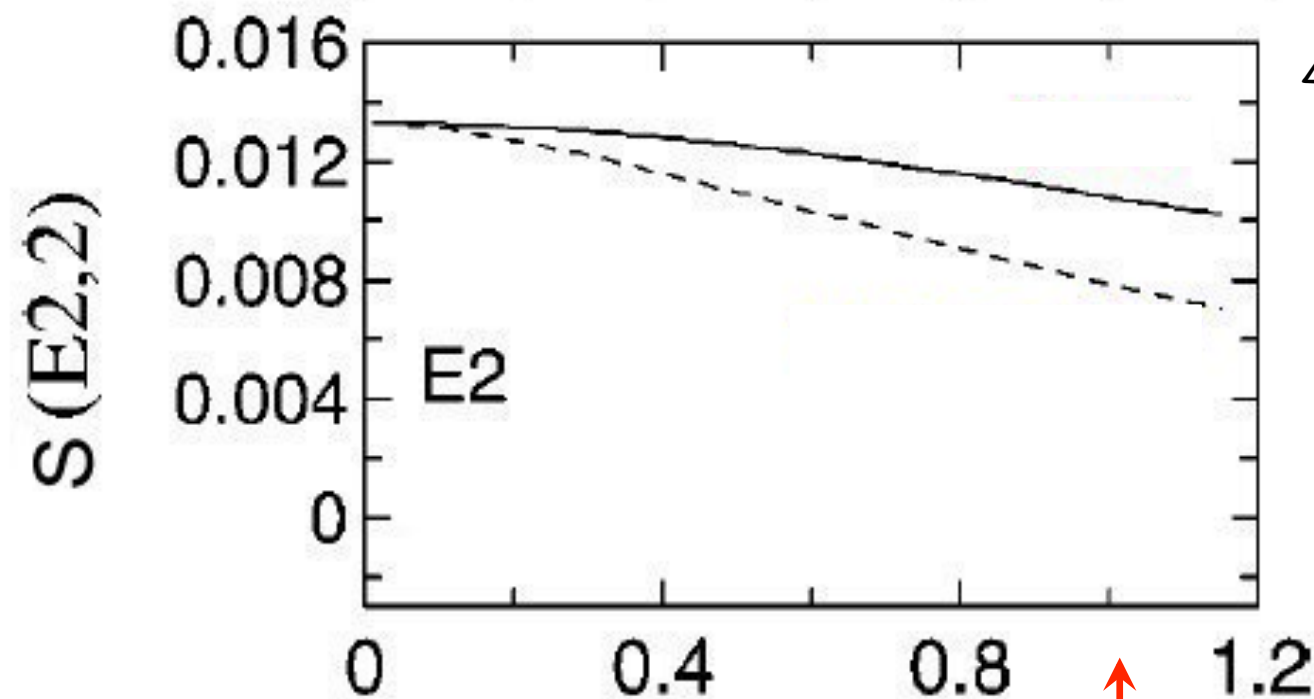
$$\begin{aligned}\frac{d\sigma}{d\Omega} &= \left(\frac{d\sigma}{d\Omega} \right)_{elast} \times P_{exc} \\ &= \sum_{\pi\lambda\mu} |S(\pi\lambda\mu)|^2 |M_{fi}(\pi\lambda, -\mu)|^2\end{aligned}$$

$$M_{fi}(\pi\lambda\mu) = \langle f | \text{EM Operator}(\lambda\mu) | i \rangle$$

$$S(\pi\lambda\mu) = \text{orbital integrals}$$



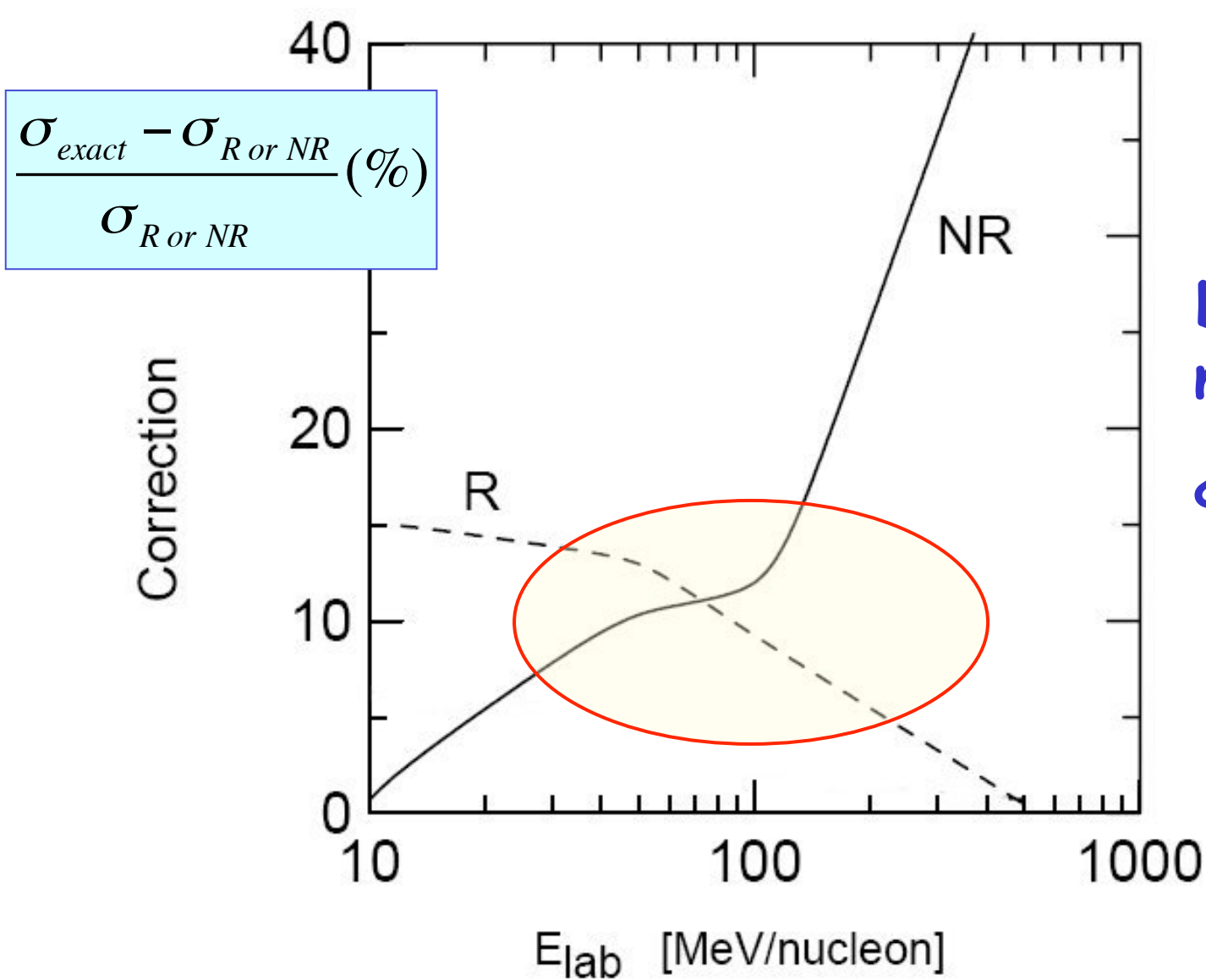
Deviations from
non-relativistic



^{40}S (100 MeV/nucleon) + Au

$$\xi = \frac{E_x b}{\gamma \hbar v}$$

Corrections important
large b 's, large E_x 's



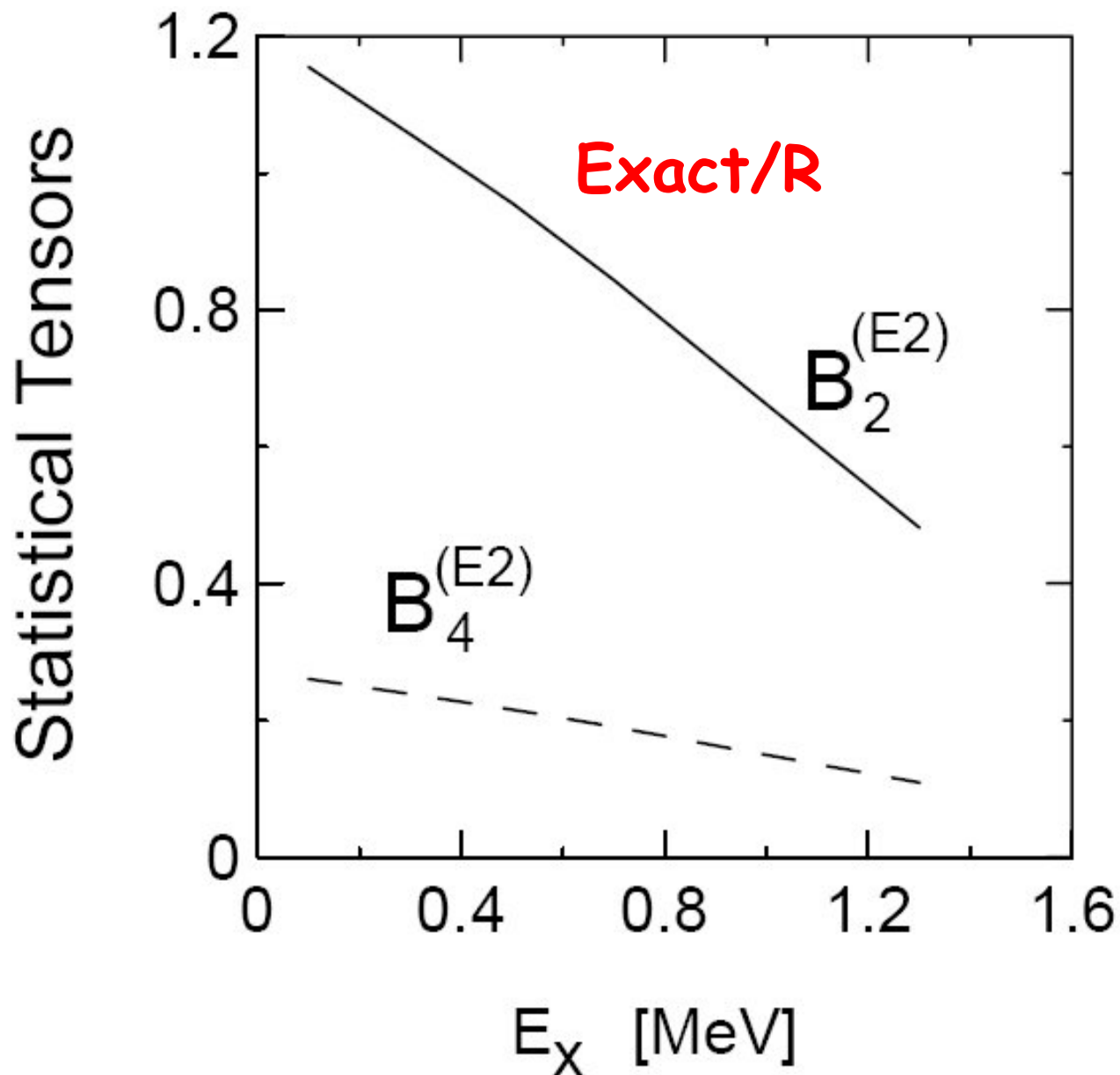
Deviations from
non-relativistic
& from relativistic

^{40}S (100 MeV/nucleon) + Au

$E_x = 0.89$ MeV

De-excitation by γ -ray emission

$$W_\gamma(\theta_\gamma) = 1 + \sum_{\kappa=2,4} B_\kappa Q_\kappa(E_\gamma) P_\kappa(\cos\theta_\gamma)$$



Deviations from
from relativistic
theory

^{38}S (100 MeV/nucleon) + Au

We all know that:

- Relativity obviously important at GANIL, GSI, MSU and RIKEN
(we are talking dynamics)
- ‘Rather’ easy to include for Coulomb interaction
(transformation properties of E/M fields well known)

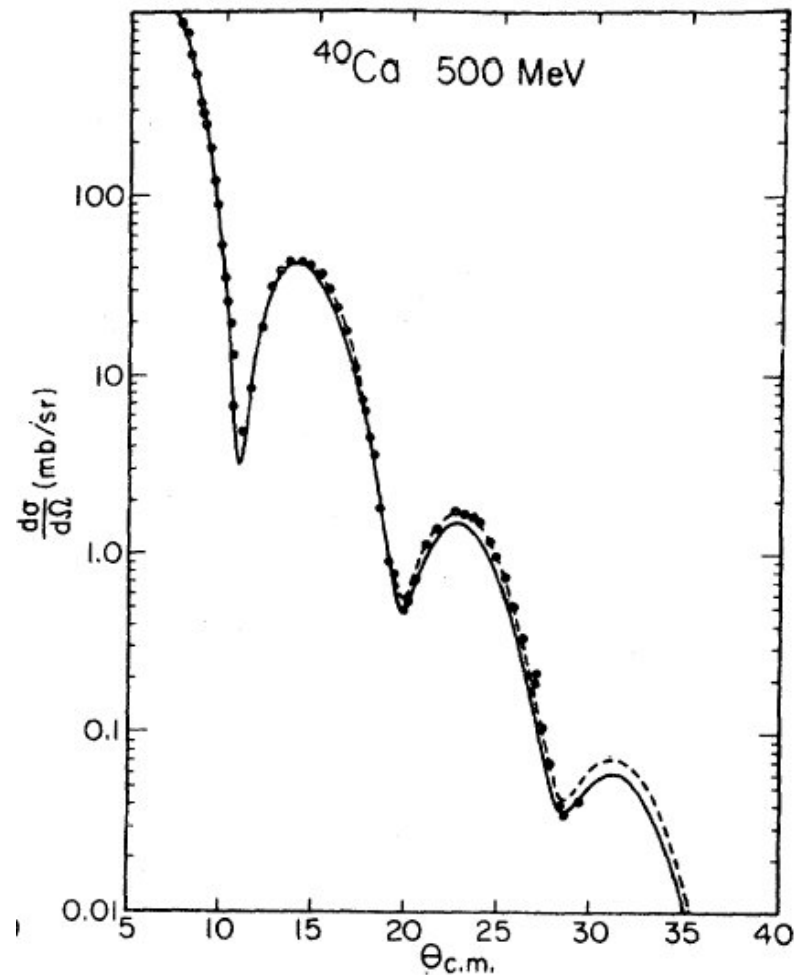
How about nuclear interaction?

- Transformation properties of nucleus-nucleus potentials not exactly known
- Solution has to be based on QFT (QM + relativity)
- Can we save our DWBA, CC, or CDCC knowledge for something practical?

Clue: Proton-nucleus scattering at intermediate energies¹⁰

- meson exchange, two-nucleon interaction
- mean field approximation, U_0 (ω exchange), U_S (2π exchange)

$$\left[E - V_C - U_0 - \beta (mc^2 + U_S) \right] \Psi = -i\hbar c \alpha \cdot \nabla \Psi$$



non-relativistic reduction

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + U_{cent} + \left(\frac{\hbar}{2mc} \right)^2 \frac{1}{r} \frac{d}{dr} U_{SO} \sigma \cdot \mathbf{L} \right] \phi = E \phi$$

$$U_{cent} = m^* (U_0 + U_S) + \dots$$

$$m^* = 1 - \frac{U_0 - U_S}{2mc^2} + \dots$$

$$U_{SO} = U_0 - U_S + \dots$$

Arnold, Clark, PLB 84, 46 (1979)

Continuum (CDCC)

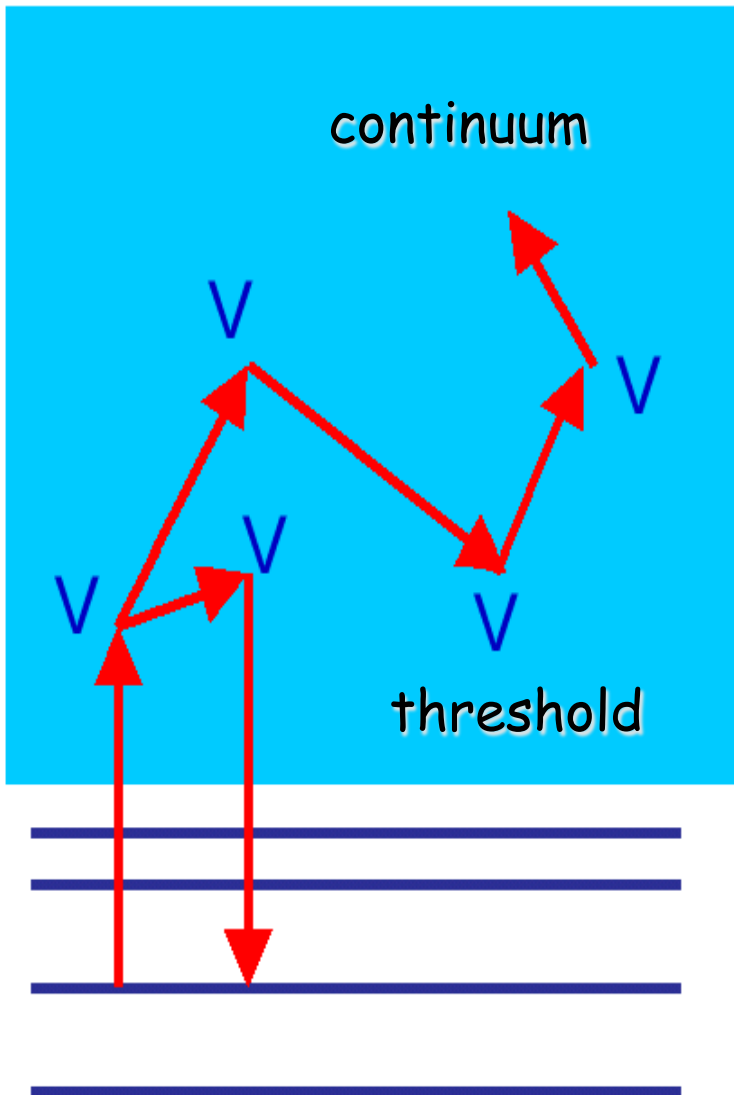
$$|\varphi_b\rangle = |E_b, J_b M_b\rangle$$

$$|\varphi_{jJM}^{(c)}\rangle = \int \Gamma_j(E) |E, JM\rangle dE$$

$$\int \Gamma_i(E) \Gamma_j(E) dE = \delta_{ij}$$

continuum discretization

$$V_{\alpha\beta}(\mathbf{R}) = \langle \phi_\alpha(\mathbf{r}) | U(\mathbf{R}, \mathbf{r}) | \phi_\beta(\mathbf{r}) \rangle$$



$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 - E \right] \chi_\alpha(\mathbf{R}) = - \sum_{\beta=0}^N V_{\alpha\beta}(\mathbf{R}) \chi_\beta(\mathbf{R})$$

Coupled-channels

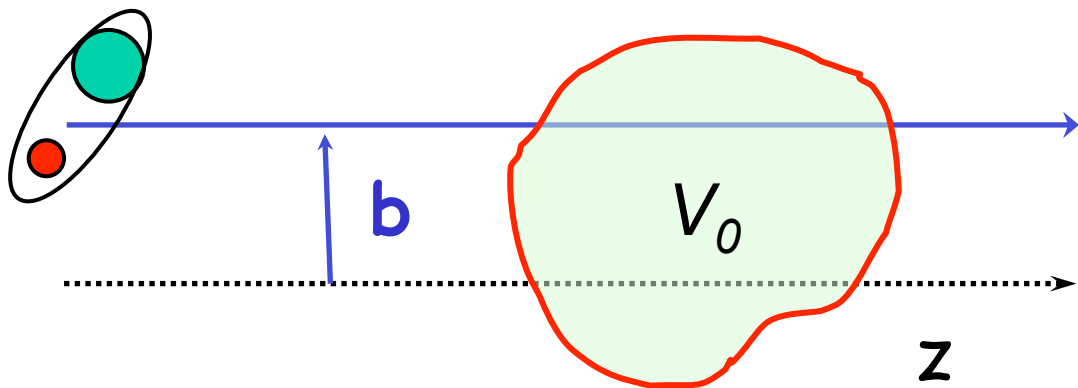
Bertulani, Canto, NPA 539, 163 (1992)
 $^{11}\text{Li} + ^{208}\text{Pb}$ (100 MeV/nucleon)

Relativistic-CDCC

Bertulani, PRL 94, 072701 (2005)

$$[\nabla^2 + k^2 - U] \Psi(\mathbf{R}, \mathbf{r}) = 0$$

$$U = V_0(2E - V_0)$$



$$\Psi(\mathbf{R}, \mathbf{r}) = \sum_{\alpha} S_{\alpha}(\mathbf{b}, z) e^{ik_{\alpha}z} \phi_{\alpha}(\mathbf{r}),$$

$$\mathbf{R} = (\mathbf{b}, z)$$

$$U \approx 2V_0E$$

$$\nabla^2 S \ll ik_z \partial_z S$$



$$iv \partial_z S_{\alpha}(\mathbf{b}, z) = \sum_{\beta} V_{\alpha\beta}(\mathbf{b}, z) S_{\beta}(\mathbf{b}, z) e^{i(k_{\beta} - k_{\alpha})z}$$

$$f_{\alpha}(\mathbf{Q}) = -\frac{ik}{2\pi} \int d\mathbf{b} e^{i\mathbf{Q} \cdot \mathbf{b}} [S_{\alpha}(\mathbf{b}, z = \infty) - \delta_{\alpha,0}]$$

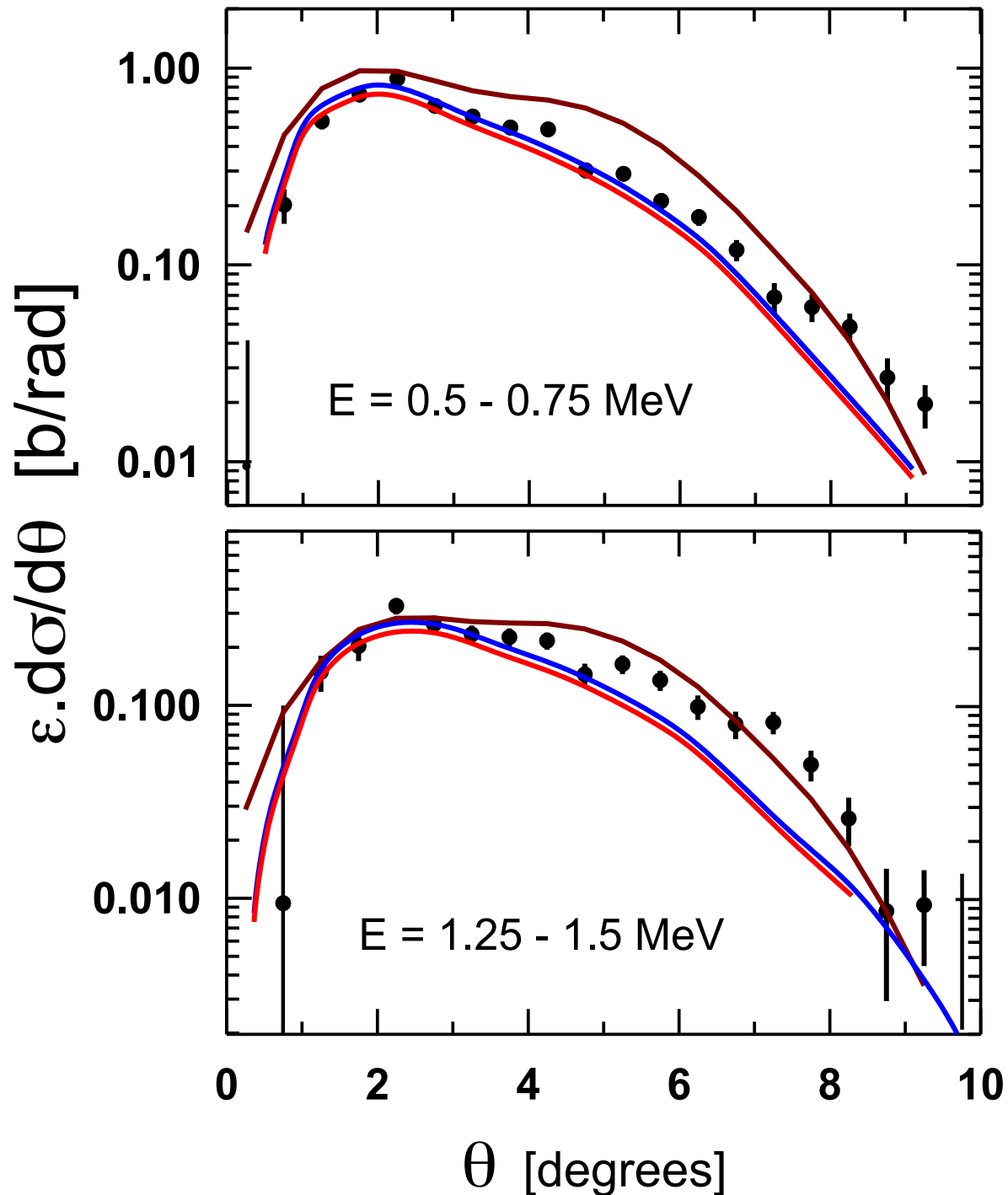
$$\mathbf{Q} = \mathbf{K}'_{\perp} - \mathbf{K}_{\perp} \quad \alpha = j l J M$$

V_0 = time-like
component of 4-vector



Relativistic CDCC
= Lorentz invariant

Pb(^8B , $p^7\text{Be}$) at 50 MeV/nucleon



DATA: Kikuchi et al, 1997

— LO

Bertulani, Gai, NPA 626, 227 (1998)

— All orders

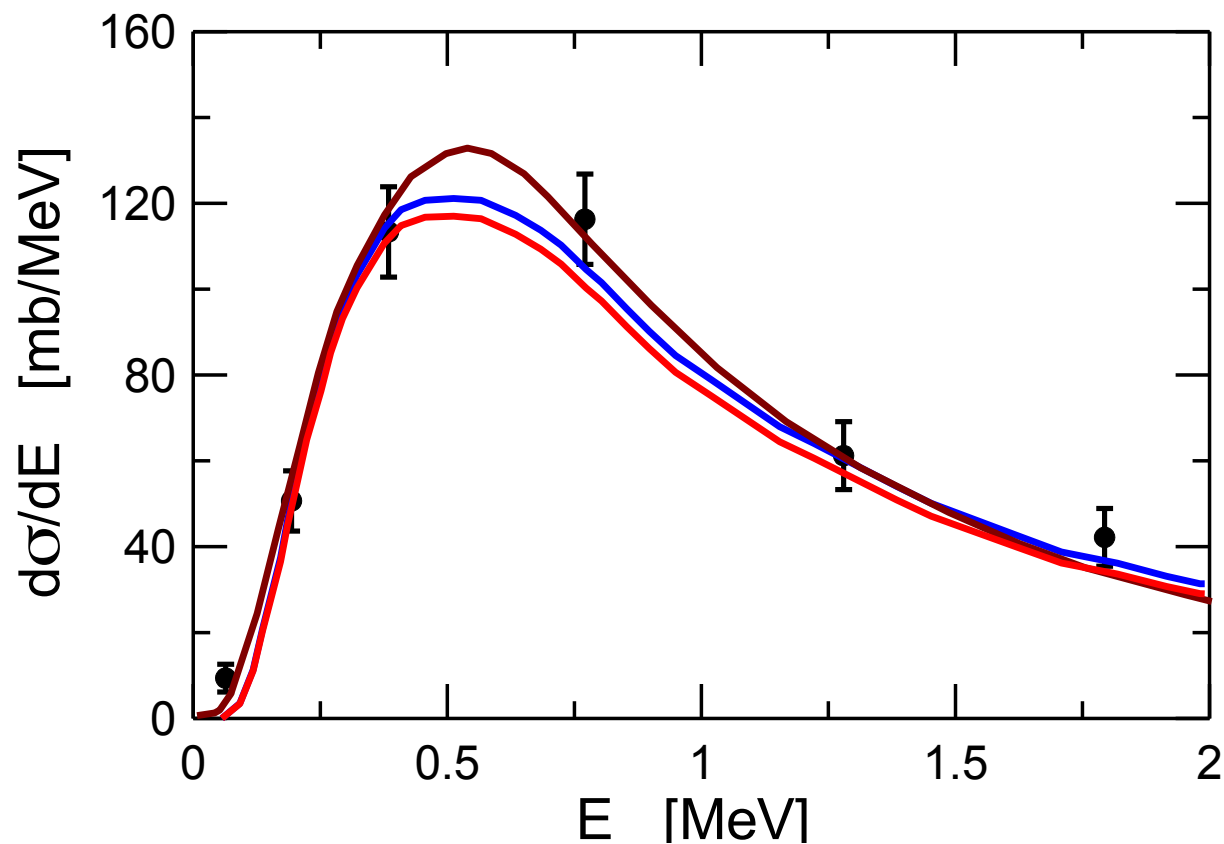
— All orders
relativistic

Bertulani, PRL 94, 072701 (2005)

$V_0 = \text{Coulomb} + \text{nuclear}$
with relativistic
corrections

5-7% effect

Pb(^8B , $p^7\text{Be}$) at 83 MeV/nucleon



DATA: Davids et al, 2002

— LO
— all orders
— all orders
— relativistic

*$V_0 = \text{Coulomb} + \text{nuclear}$
with relativistic
corrections*

4-10% effect

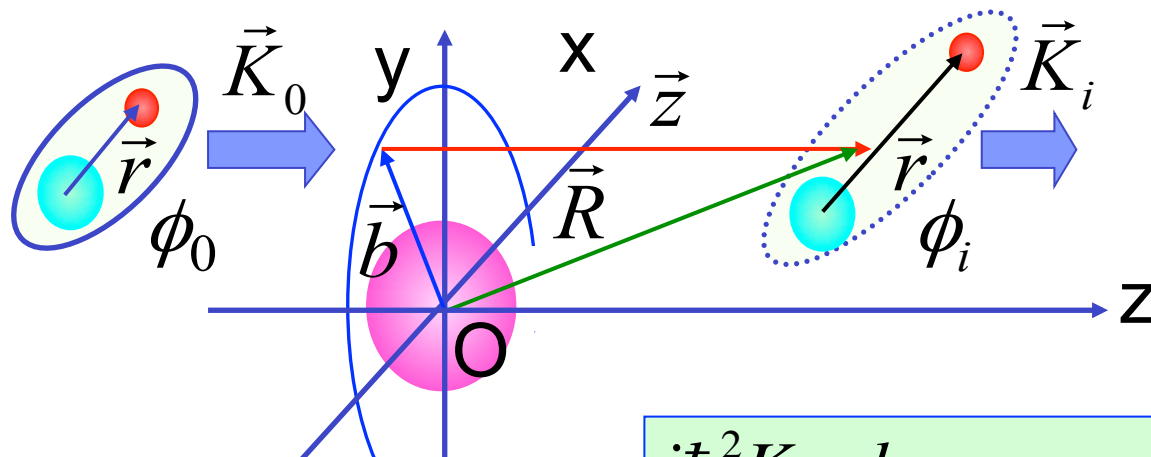
Transition: low to high energies

Eikonal scattering waves $\hat{S}_i(K_i, \vec{R})$

$$\psi^{E-CDCC} = \sum_i \hat{\phi}_i(\vec{r}) \hat{S}_i(b, z) \exp(i\vec{K}_i \cdot \vec{R})$$

$$K_i = \sqrt{2\mu_R(E - \varepsilon_i)} / \hbar,$$

Energy conservation



● Boundary condition

$$\hat{S}_i(b, z) \xrightarrow{z \rightarrow -\infty} \delta_{i,0}$$

$$\Delta \hat{S}_i(b, z) \cong 0 \quad \longrightarrow$$

$$\frac{i\hbar^2 K_i}{\mu_R} \frac{d}{dz} \hat{S}_i^{(b)}(z) = \sum_{i'} F_{ii'}^{(b)}(z) \hat{S}_{i'}^{(b)}(z) e^{i(K_{i'} - K_i)z}$$

Eikonal scattering amplitude transformed into QM form

$$f_{i,0}^E = \sum_L f_L^E \equiv \sum_L \frac{2\pi}{iK_i} \sqrt{\frac{2L+1}{4\pi}} i^m Y_{Lm}(\Omega) [S_{i,0}^{b(L;i)} - \delta_{i,0}]$$

Hybrid scattering amplitude is given by

$$f_{i,0}^H \equiv \sum_{L=0}^{L_C} f_L^Q + \sum_{L=L_C+1}^{L_{\max}} f_L^E$$

Ogata., et al, PRC68, 064609 (2003)

Relativistic CDCC

Form factor of non-rel. E-CDCC

$$F_{c'c}^{(b)}(Z) = \langle \Phi_{c'} | U_{1A} + U_{2A} | \Phi_c \rangle_r e^{-i(m-m')\phi} = \sum_{\lambda} F_{c'c}^{(b);\lambda}(Z)$$

Lorentz transform of form factor and coordinates

$$F_{c'c}^{(b);\lambda}(Z) \rightarrow f_{\lambda, m'-m} \gamma F_{c'c}^{(b)\lambda}(\gamma Z)$$

$$f_{\lambda, m'-m}^{\text{Coul}} = \begin{cases} 1/\gamma & (\lambda=1, m'-m=0) \\ \gamma & (\lambda=2, m'-m=\pm 1) \\ 1 & (\text{otherwise}) \end{cases}$$

$$f_{\lambda, m'-m}^{\text{nucl}} = 1$$

Assumptions

- ✓ Point charges for 1, 2 and A
- ✓ Neglecting far-field ($r_i > R$) contribution
- ✓ Correction to nuclear form factor

Ogata, Bertulani, PTP 121 (2009), 1399
PTP, 123 (2010) 701

Reaction

$^{208}\text{Pb}(^8\text{B}, ^7\text{Be}+p)$ at 250 A MeV and 100 A MeV

$^{208}\text{Pb}(^{11}\text{Be}, ^{10}\text{Be}+n)$ at 250 A MeV and 100 A MeV

Projectile wave function and distorting potential

Standard Woods-Saxon

Modelspace

^8B

$$l_{\max} = 3$$

$$N_s = 20, N_{p-d} = 10,$$

$$N_f = 5$$

$$\varepsilon_{\max} = 10 \text{ MeV}$$

$$r_{\max} = 200 \text{ fm}$$

$$R_{\max} = 500 \text{ fm}$$

$$N_{\text{ch}} = 138$$

^{11}Be

$$l_{\max} = 3$$

$$N_{s,p} = 20, N_d = 10,$$

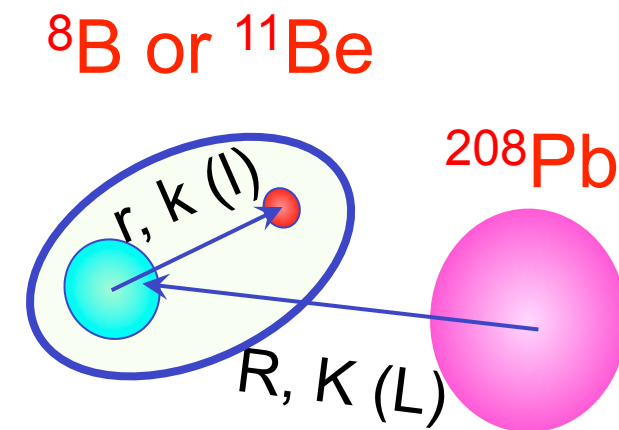
$$N_f = 5$$

$$\varepsilon_{\max} = 10 \text{ MeV}$$

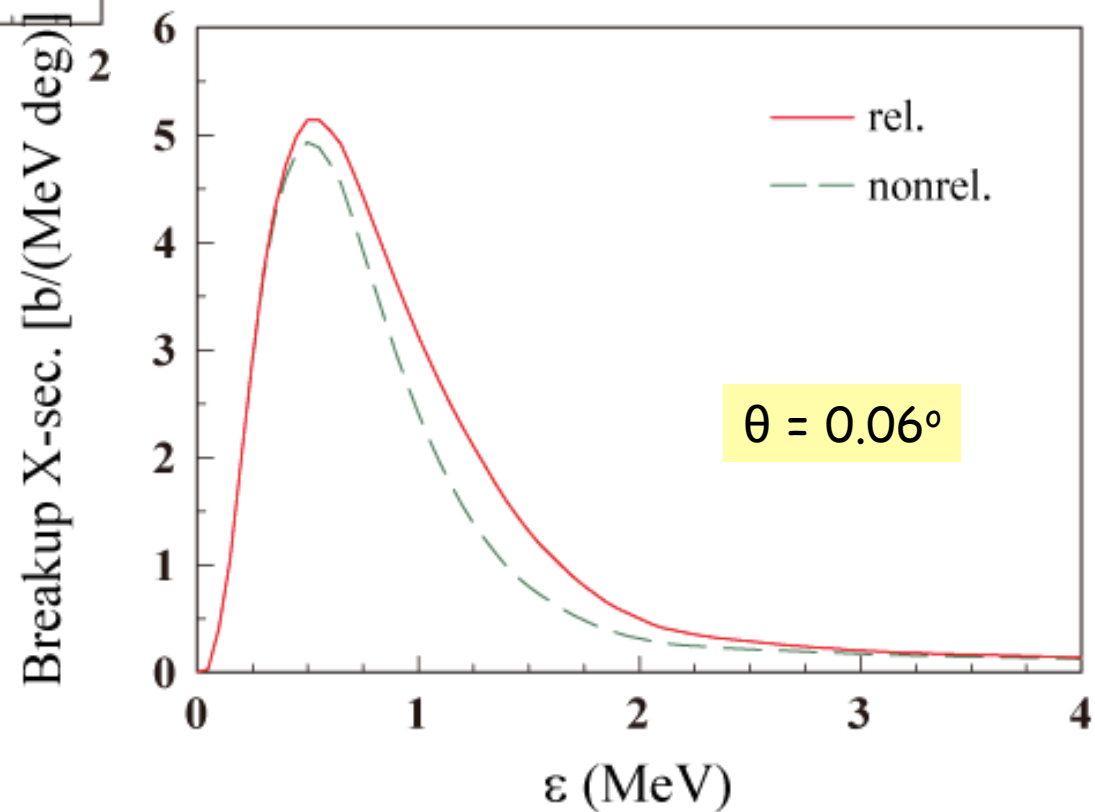
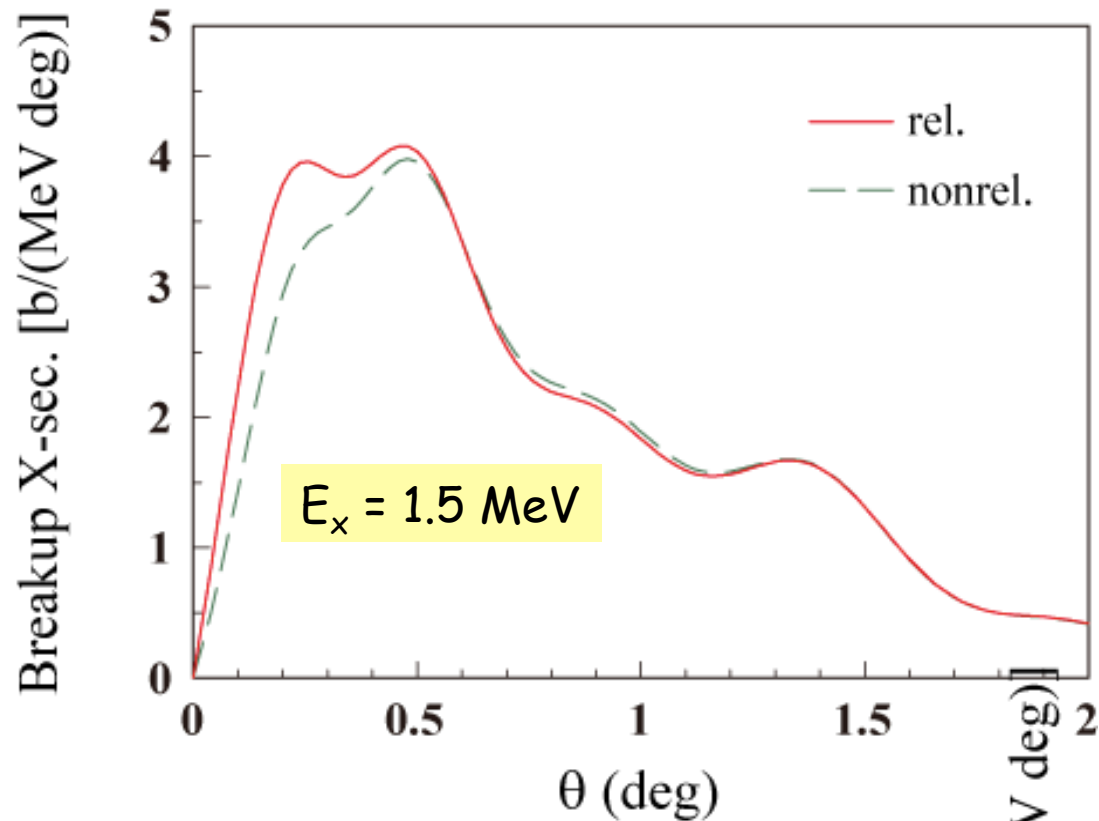
$$r_{\max} = 200 \text{ fm}$$

$$R_{\max} = 450 \text{ fm}$$

$$N_{\text{ch}} = 166$$



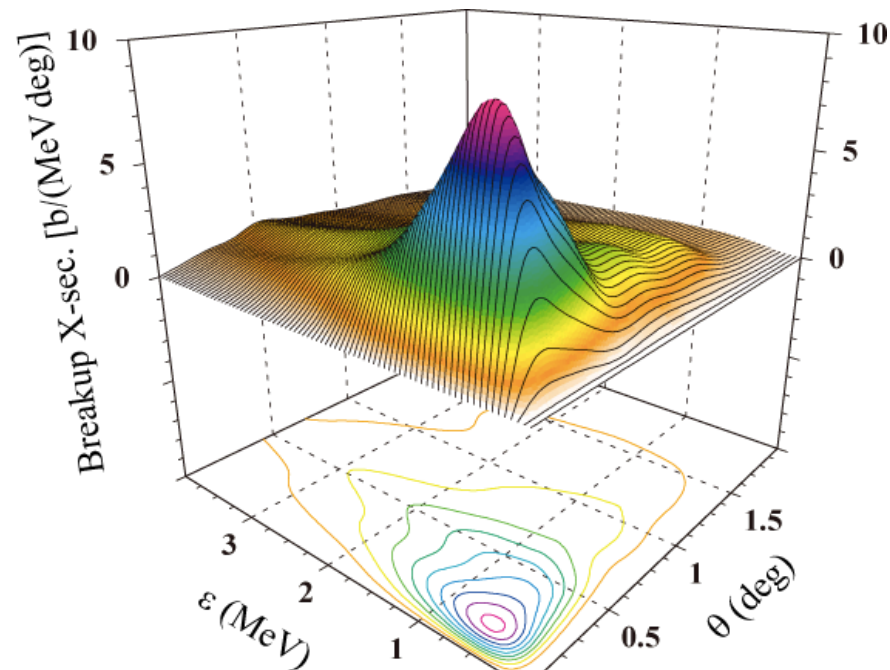
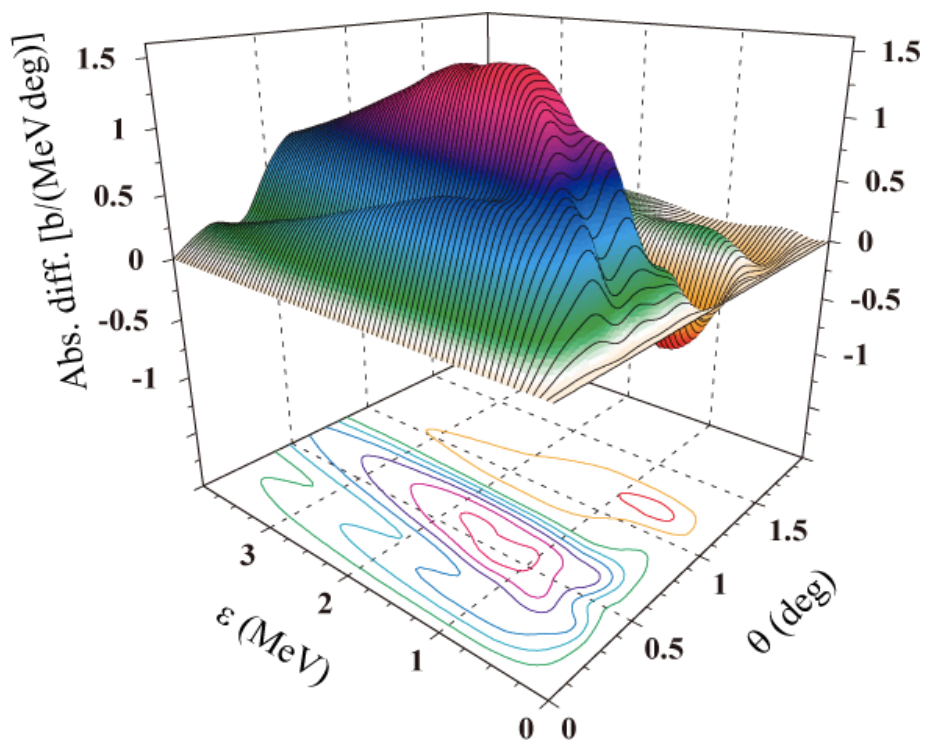
Pb($^8\text{B}, p^7\text{Be}$) at 250 MeV/nucleon



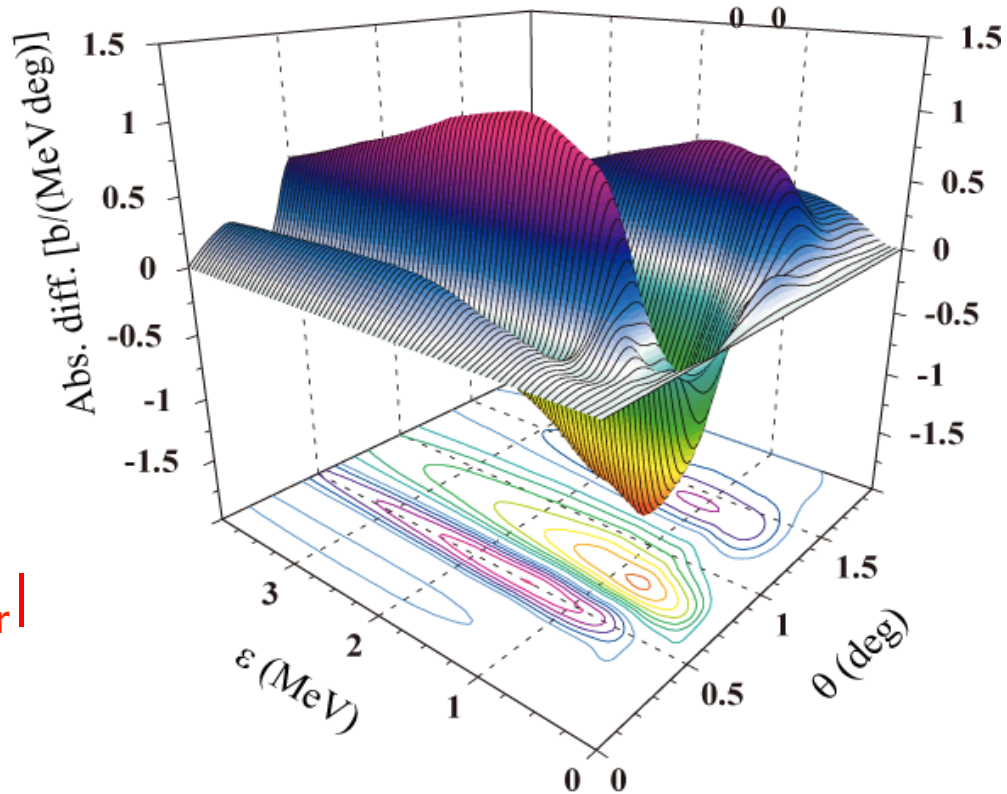
Pb(^8B , p ^7Be) at 250 MeV/nucleon

all orders

$$|\sigma_{\text{all}} - \sigma_{\text{NR}}|$$



$$|\sigma_{\text{all}} - \sigma_{\text{no-nuclear}}|$$



Relativistic MF nucleus-nucleus potential

Long, Bertulani, PRC 83, 024907 (2011).

σ, ω, ρ and γ exchange

$$E = \int d^3r \sum_a \bar{\psi}_a (-i\boldsymbol{\gamma} \cdot \nabla + M) \psi_a$$

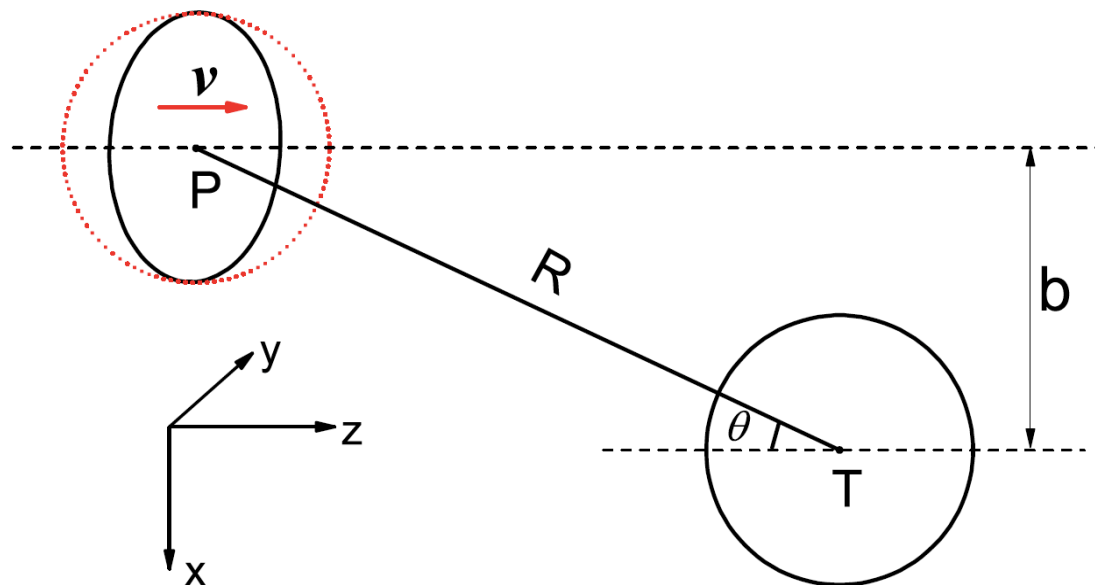
$$+ \frac{1}{2} \sum_{\phi=\sigma,\omega,\rho,\gamma} \int d^3r d^3r' \sum_{ab} \bar{\psi}_a(\mathbf{r}) \bar{\psi}_b(\mathbf{r}') \Gamma_\phi(\mathbf{r}, \mathbf{r}') D_\phi(\mathbf{r} - \mathbf{r}') \psi_a(\mathbf{r}) \psi_b(\mathbf{r}')$$

$$\Gamma_\phi(\mathbf{r}, \mathbf{r}') = -g_\sigma(\mathbf{r}) g_\sigma(\mathbf{r}')$$

$$\Gamma_\omega(\mathbf{r}, \mathbf{r}') = -\left(g_\omega \gamma^\mu\right)_\mathbf{r} \cdot \left(g_\omega \gamma_\mu\right)_{\mathbf{r}'}$$

$$\Gamma_\rho(\mathbf{r}, \mathbf{r}') = -\left(g_\rho \gamma^\mu \vec{\tau}\right)_\mathbf{r} \cdot \left(g_\rho \gamma_\mu \vec{\tau}\right)_{\mathbf{r}'}$$

$$\Gamma_\gamma(\mathbf{r}, \mathbf{r}') = \frac{e^2}{4} \left[\gamma^\mu (1 - \tau_z)\right]_\mathbf{r} \cdot \left[\gamma_\mu (1 - \tau_z)\right]_{\mathbf{r}'}$$



$$D_\phi = \frac{1}{4\pi} \frac{e^{m_\phi |\mathbf{r} - \mathbf{r}'|}}{|\mathbf{r} - \mathbf{r}'|}$$

$$D_\gamma = \frac{1}{|\mathbf{r} - \mathbf{r}'|}$$

Lorentz transform

$$x_p = x_t + b, \quad y_p = y_t$$

$$z_p = \gamma(z_t + R \cos \theta)$$

$$E(A_t, A_p, v) = E(A_t) + E(A_p, v) + \mathbf{E}(A_t, A_p, v)$$

$$\mathbf{E}(A_t, A_p, v) = \sum_{\phi=\sigma, \omega, \rho, \gamma} \int d^3 r \int d^3 r' \sum_{ab} \bar{\psi}_{t,a}(\mathbf{r}) \bar{\psi}_{p,b}(\mathbf{r}') \Gamma_{\phi}(\mathbf{r}, \mathbf{r}') D_{\phi}(\mathbf{r} - \mathbf{r}') \psi_{t,a}(\mathbf{r}) \psi_{p,b}(\mathbf{r}')$$

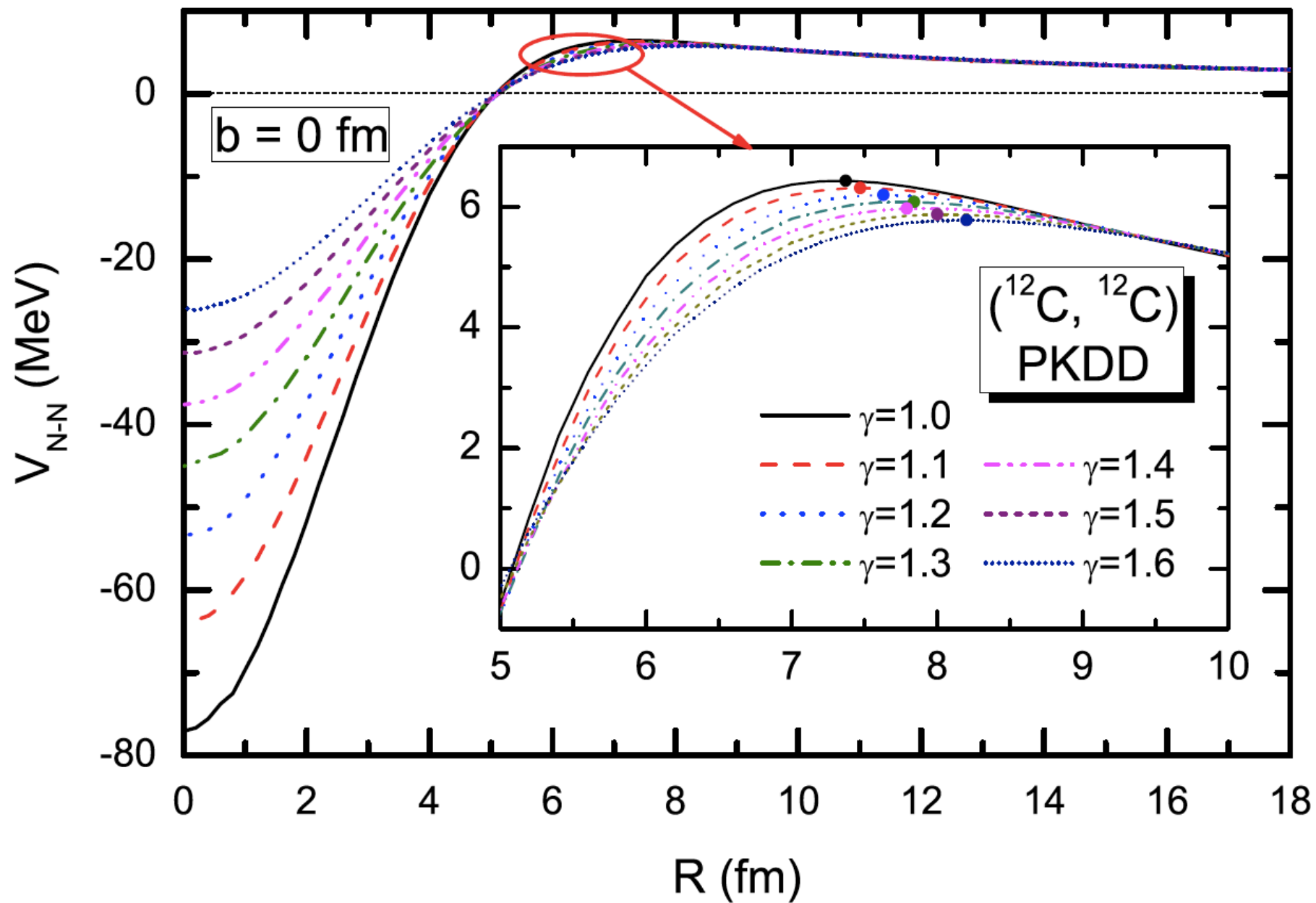
Ex: σ and ω contributions

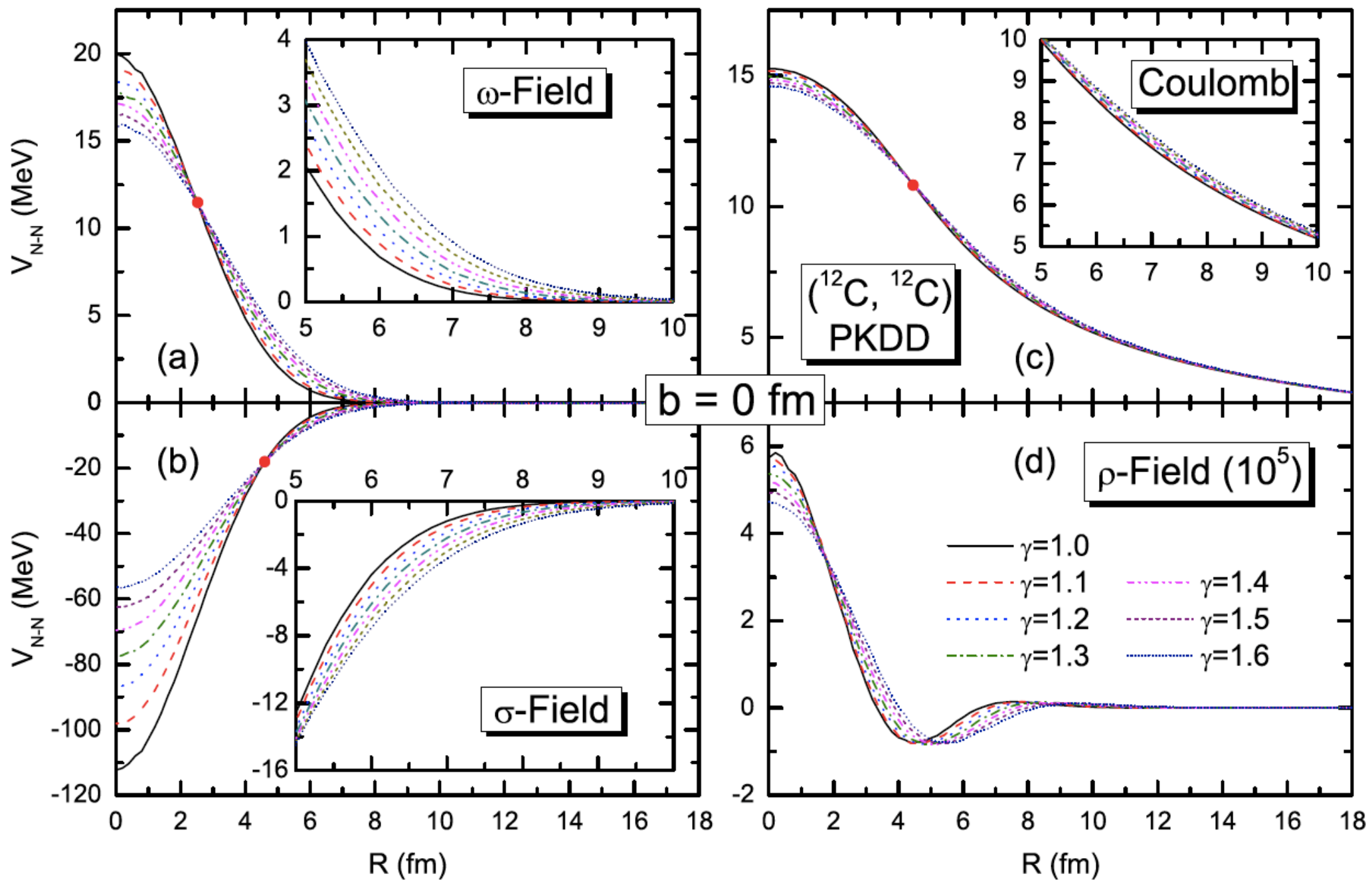
$$\mathbf{E}_{\sigma} = -\frac{1}{\gamma} \int d^3 r_t \int d^3 r'_p g_{\sigma}(\mathbf{r}_t) \rho_{s,t}(\mathbf{r}_t) D_{\sigma}(\mathbf{r} - \mathbf{r}') \rho_{s,p}(\mathbf{r}'_p) g_{\sigma}(\mathbf{r}'_p)$$

$$\mathbf{E}_{\omega} = \int d^3 r_t \int d^3 r'_p g_{\omega}(\mathbf{r}_t) \rho_{b,t}(\mathbf{r}_t) D_{\omega}(\mathbf{r} - \mathbf{r}') \rho_{b,p}(\mathbf{r}'_p) g_{\omega}(\mathbf{r}'_p)$$

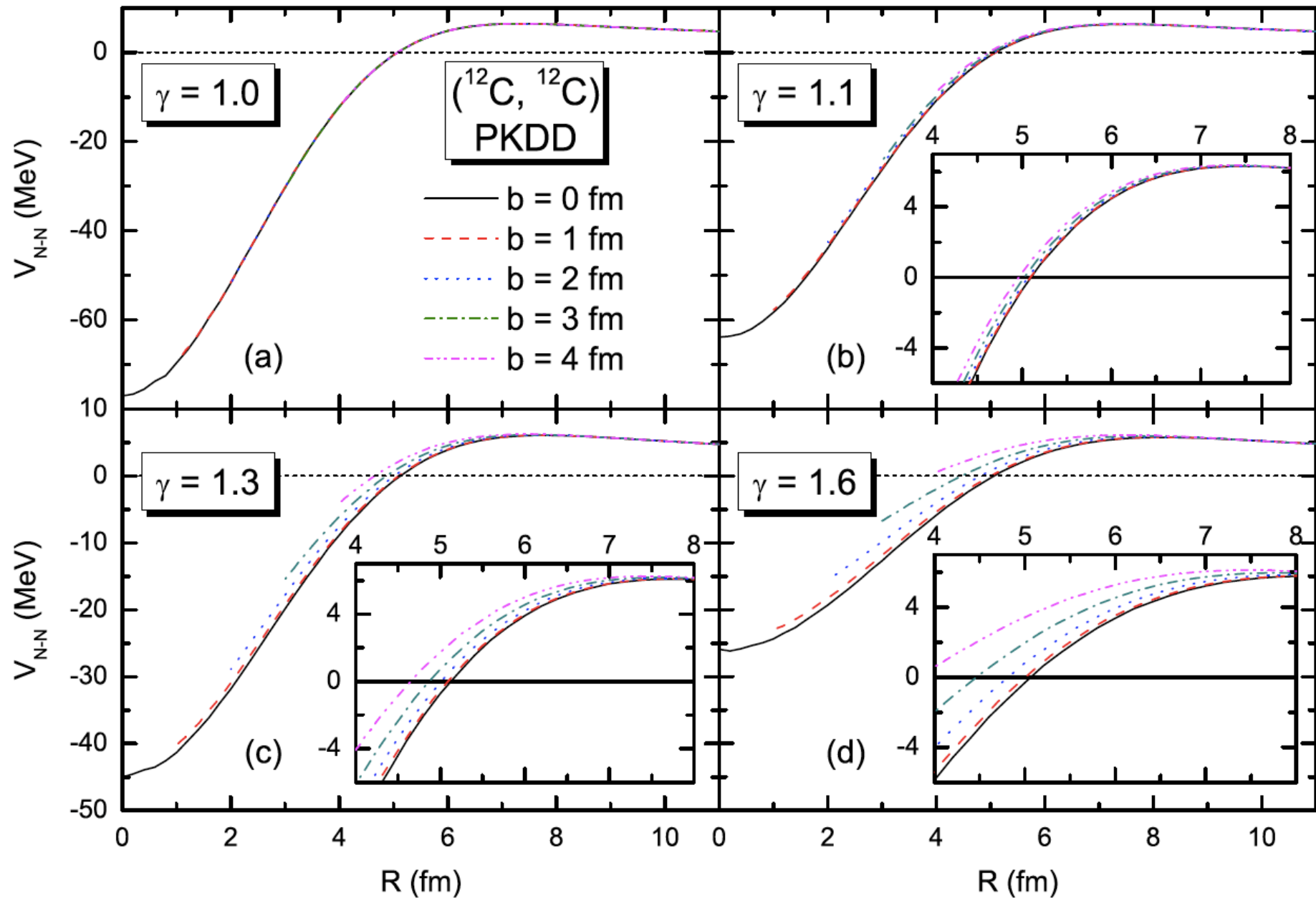
$$\rho_s(\mathbf{r}) = \sum_a \bar{\psi}_a(\mathbf{r}) \psi_a(\mathbf{r}), \quad \rho_b(\mathbf{r}) = \sum_a \bar{\psi}_a(\mathbf{r}) \gamma^0 \psi_a(\mathbf{r})$$

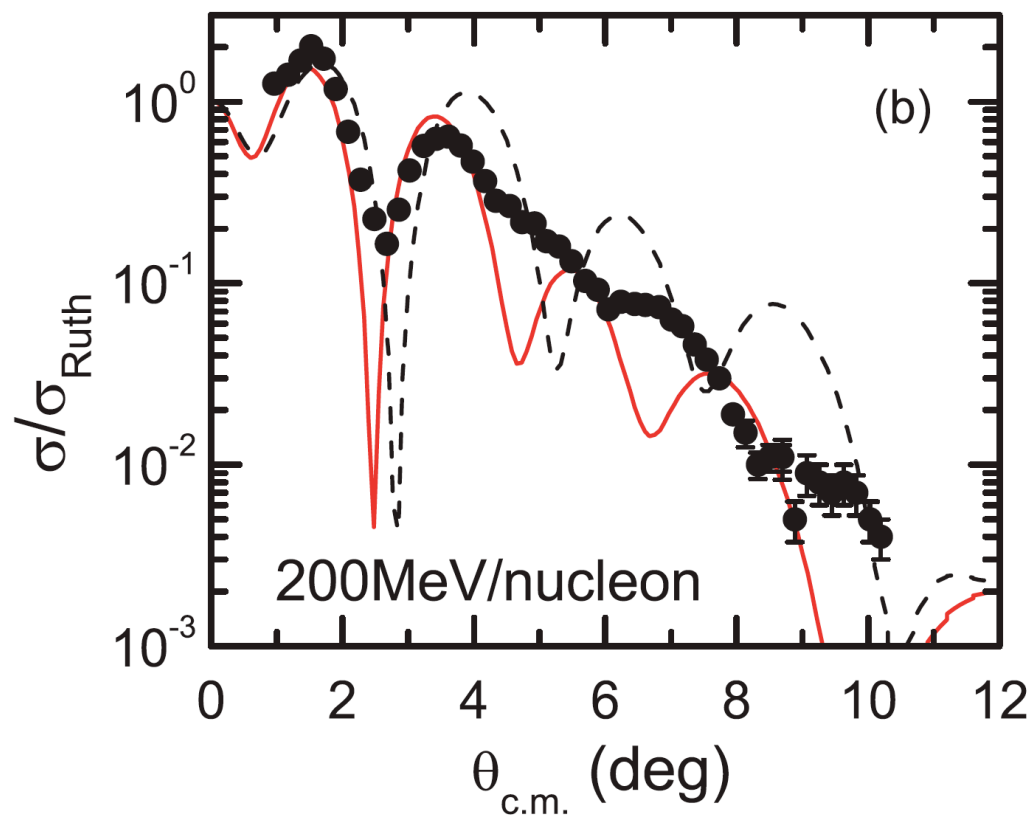
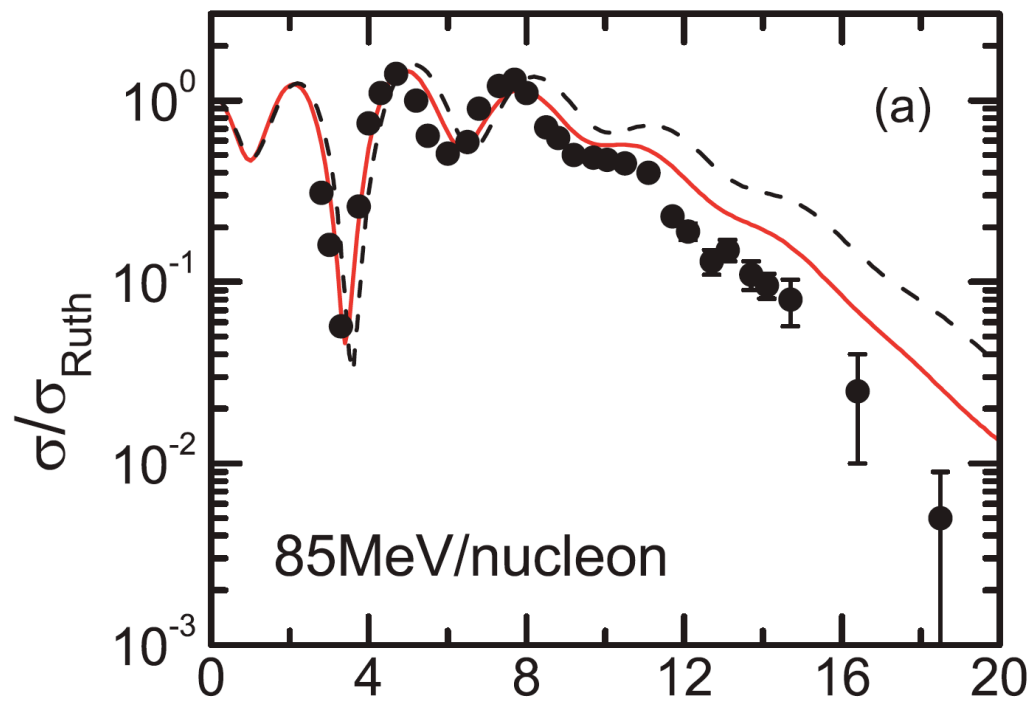
Projectile densities boosted to the target frame





Dependence on energy and impact parameter





Medium effects in σ_{NN}

$$\langle \mathbf{k} | G | \mathbf{k}_0 \rangle = \langle \mathbf{k} | V_{NN} | \mathbf{k}_0 \rangle - \int \frac{d^3 k'}{(2\pi)^3} \frac{\langle \mathbf{k} | V_{NN} | \mathbf{k}' \rangle Q(\mathbf{k}') \langle \mathbf{k}' | G | \mathbf{k}_0 \rangle}{E(\mathbf{k}') - E_0 - i\epsilon}$$

$$E(\mathbf{P}, \mathbf{k}) = e(\mathbf{P} + \mathbf{k}) + e(\mathbf{P} - \mathbf{k})$$

e = single-particle energies

$E_0 = E$ on-shell

$$Q(\mathbf{P}, \mathbf{k}) = 1, \quad \text{if } k_{1,2} > k_F \\ = 0, \quad \text{otherwise}$$

$$\mathbf{k}_{1,2} = \mathbf{P} \pm \mathbf{k}$$

In real calculations:

$$\bar{Q}(P, k) = \frac{\int d\Omega Q(\mathbf{P}, \mathbf{k})}{\int d\Omega}$$

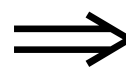
$$e(p) = T(p) + v(p)$$

$$v(p) = \langle p | v | p \rangle = \text{Re} \sum_{q \leq k_F} \langle pq | G | pq - qp \rangle$$

- e depends on v

- v depends on G

- G depends on v



Solve self-consistently
(Brueckner theory)

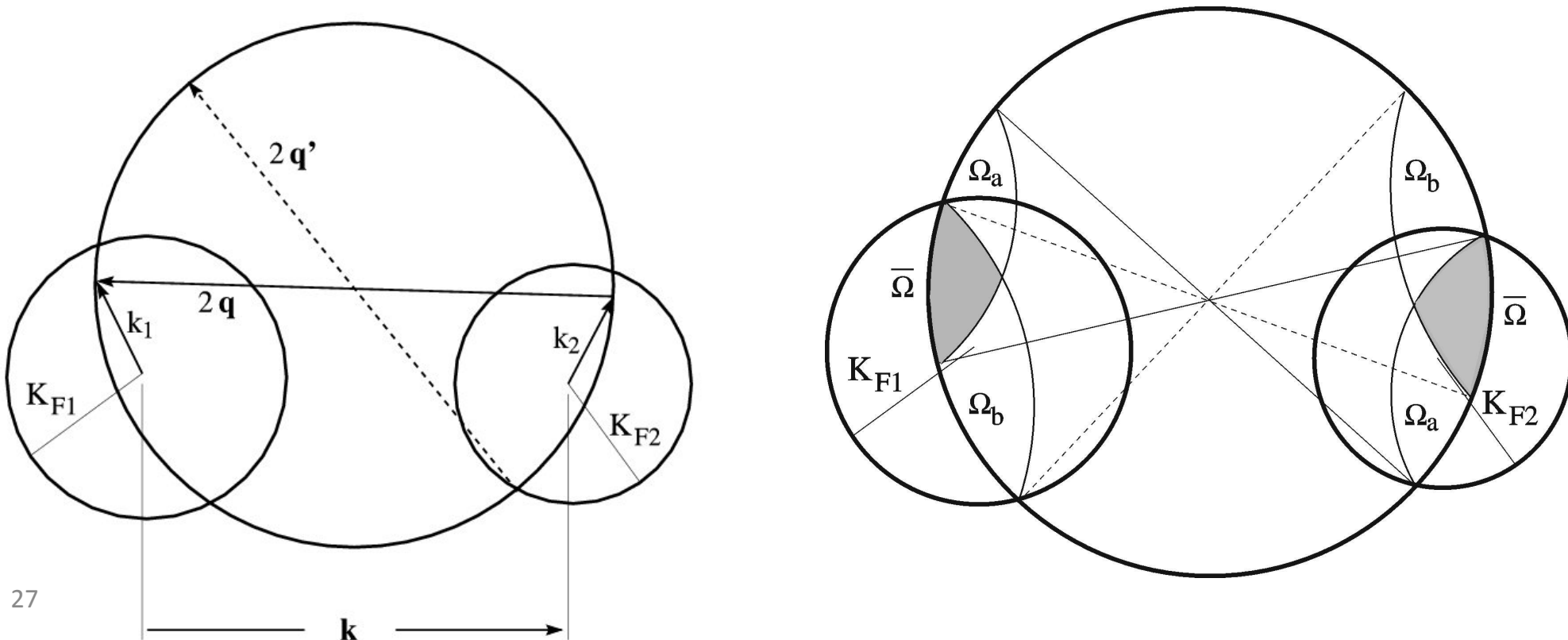
Geometric approximation + LDA

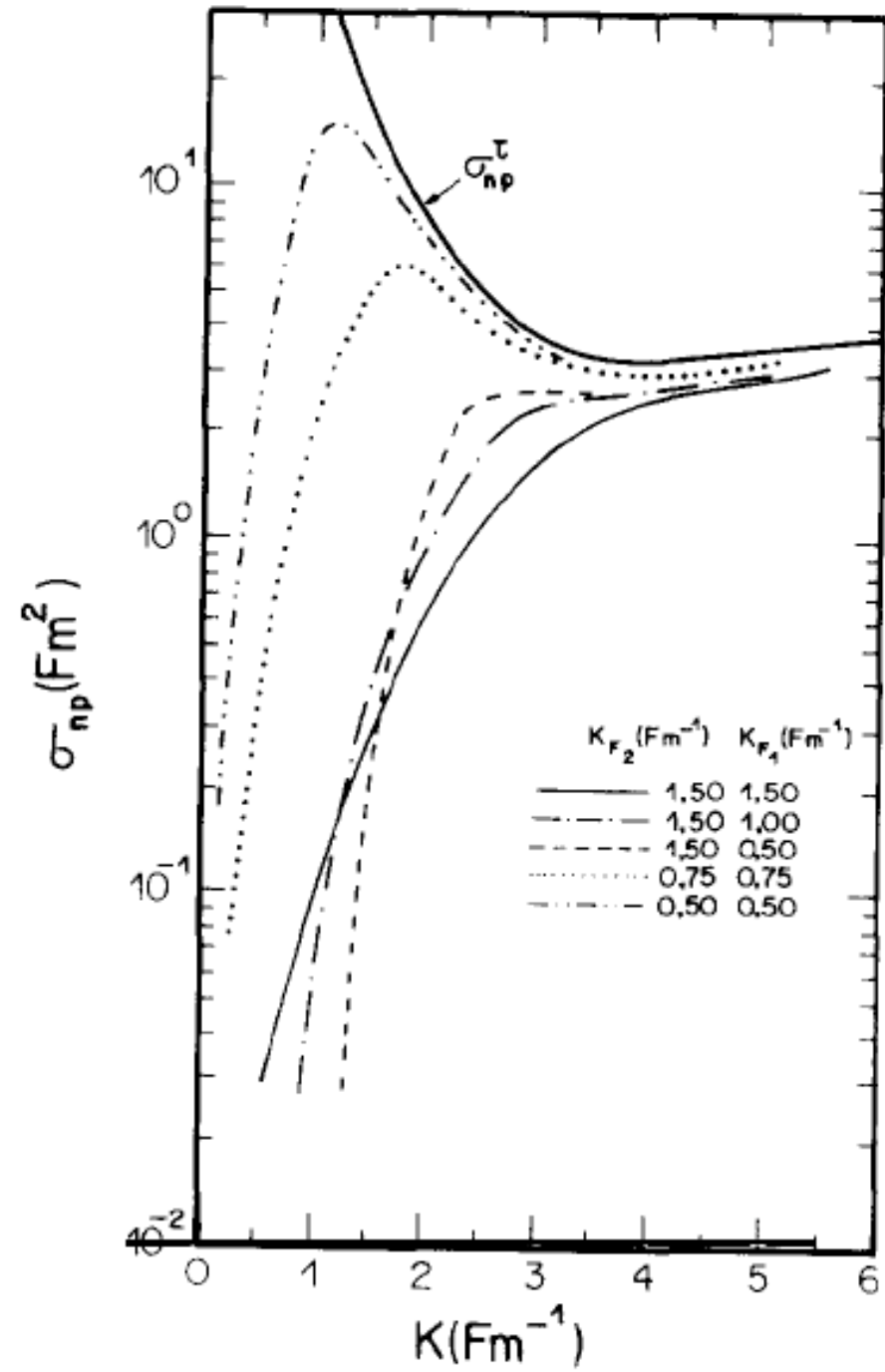
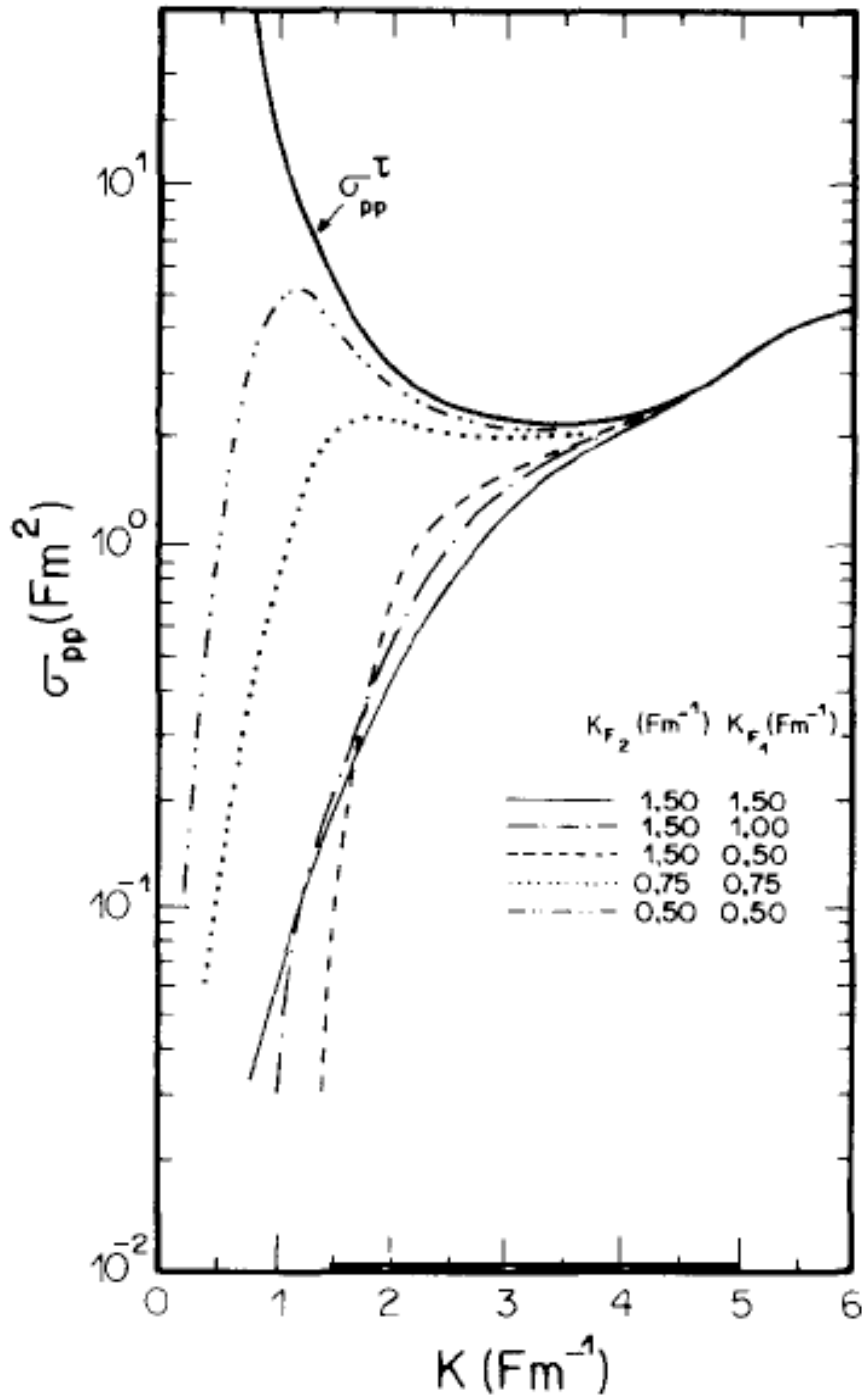
CB, Phys. Rep. (1991), JPG 27, L67 (2001)

CB, De Conti, PRC C 81, 064603 (2010)

$$\bar{\sigma}_{NN}(E) = \int \frac{d^3 k_1 d^3 k_2}{(4\pi k_{1F}^3 / 3)(4\pi k_{2F}^3 / 3)} \frac{2q}{k} \sigma_{NN}(q) \frac{\Omega_{Pauli}}{4\pi}$$

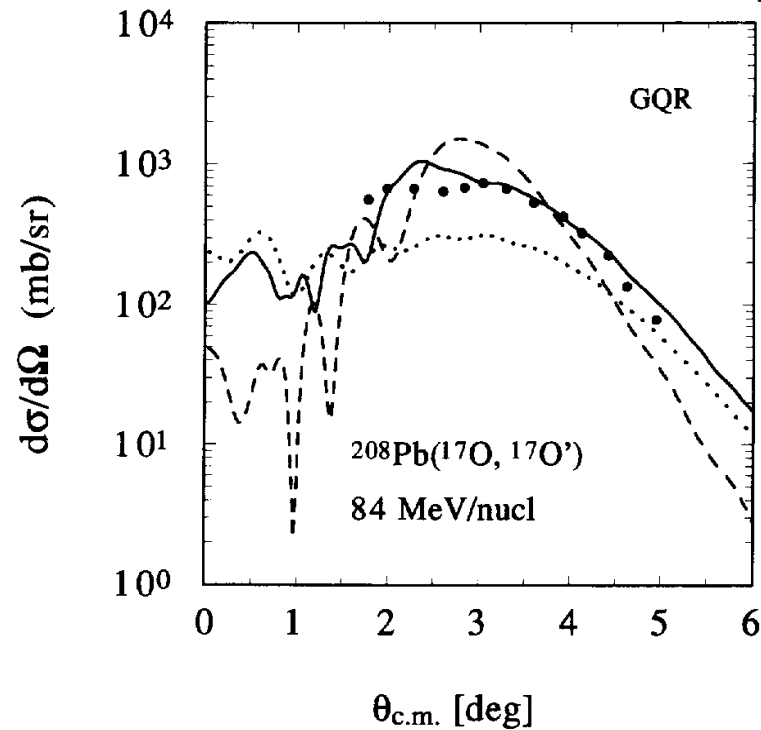
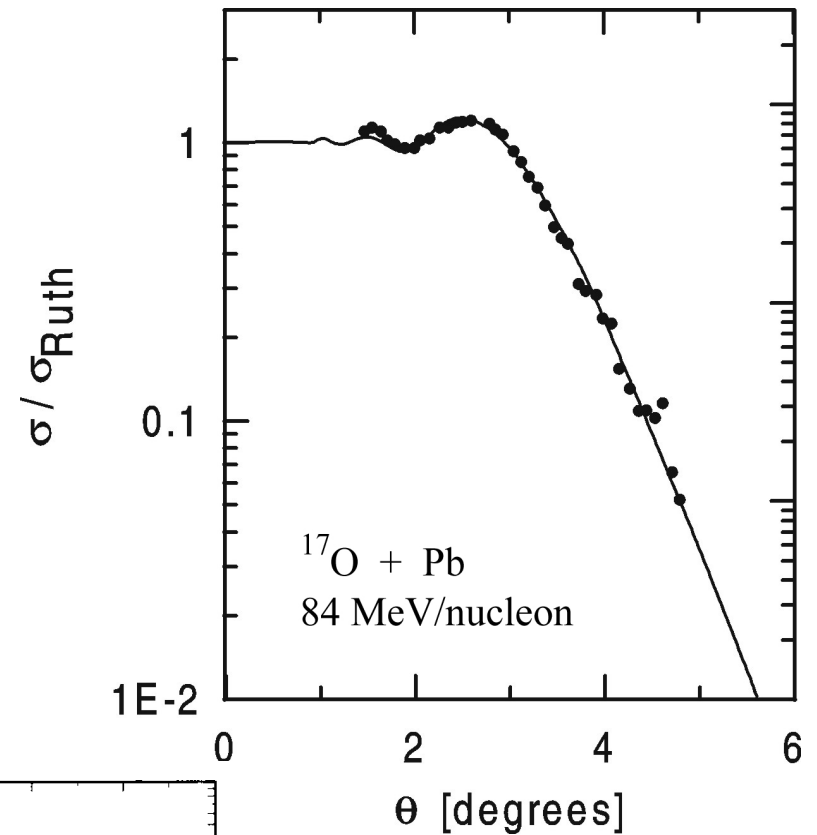
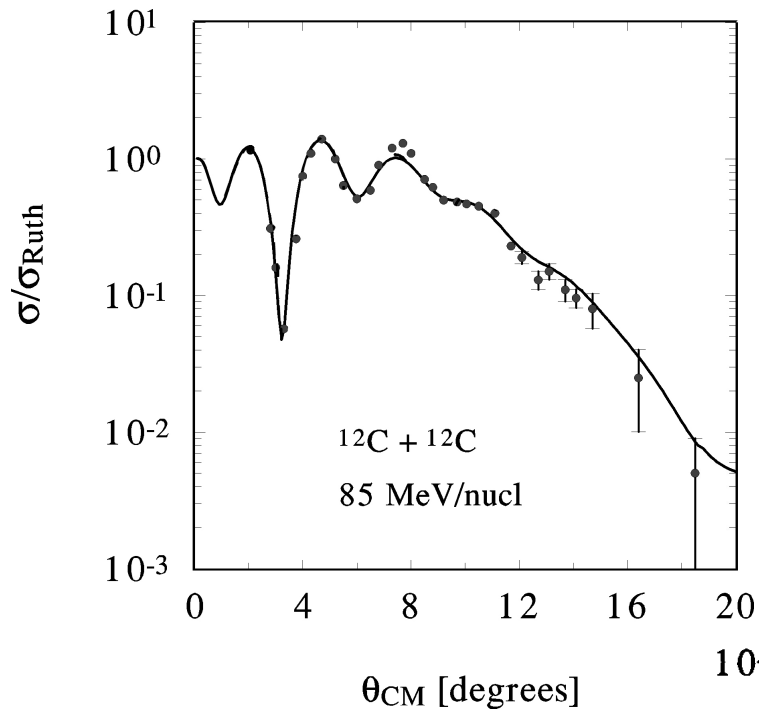
$$\Omega_{Pauli} = 4\pi - 2(\Omega_a + \Omega_b - \bar{\Omega}) = \text{analytic}$$

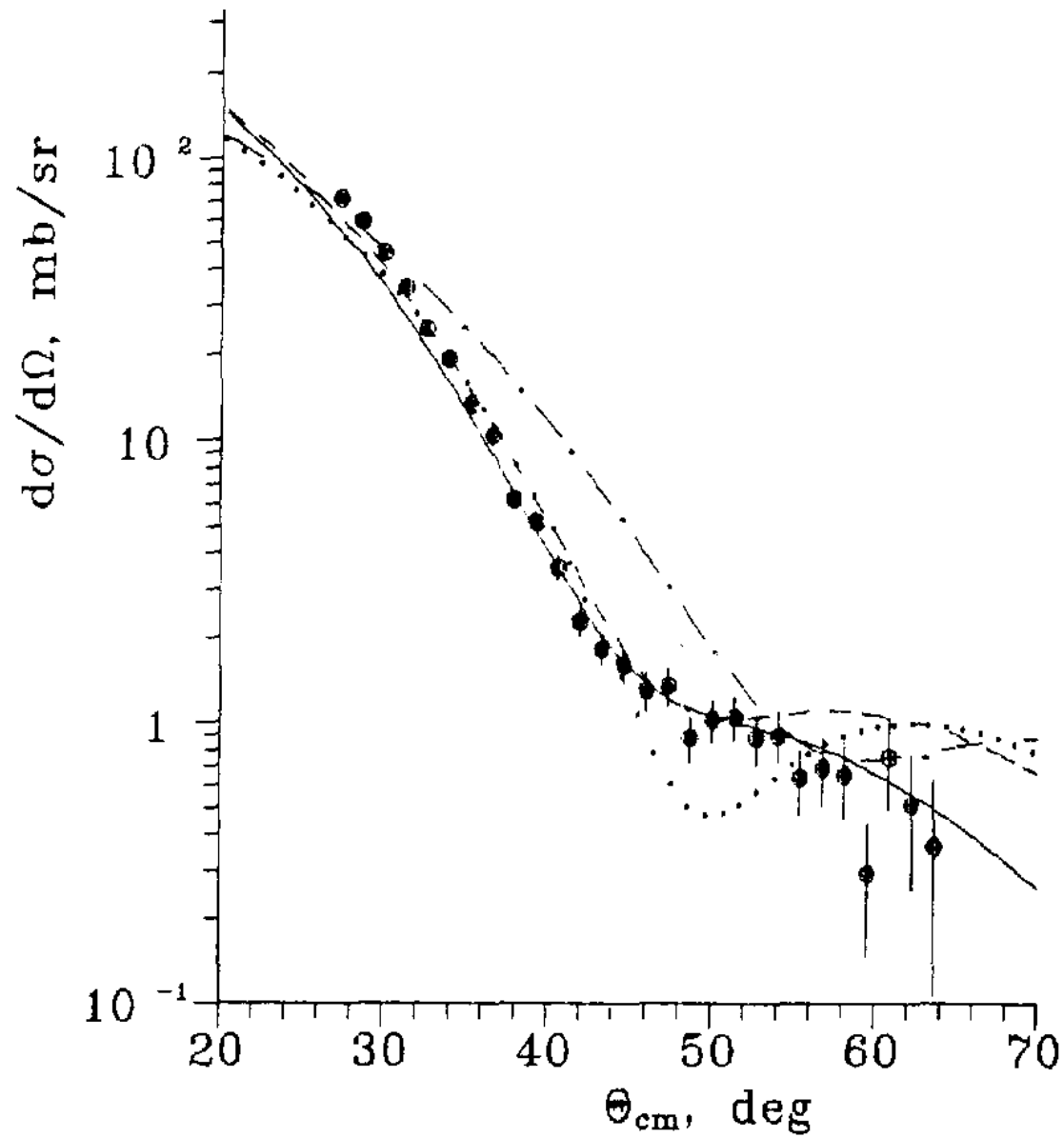




Nucleus-nucleus elastic and inelastic scattering

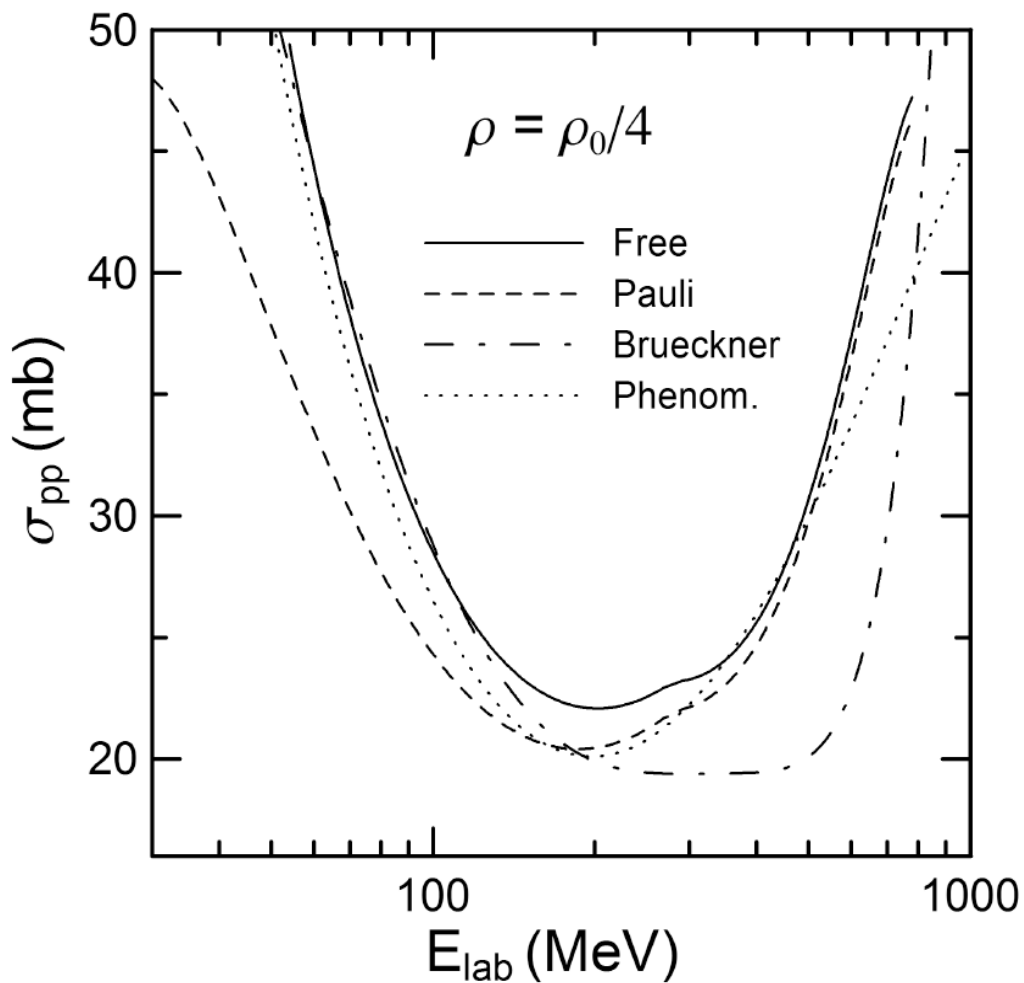
Bertulani, Sagawa, PLB 300 (1993) 205
 NPA 588 (1995) 667



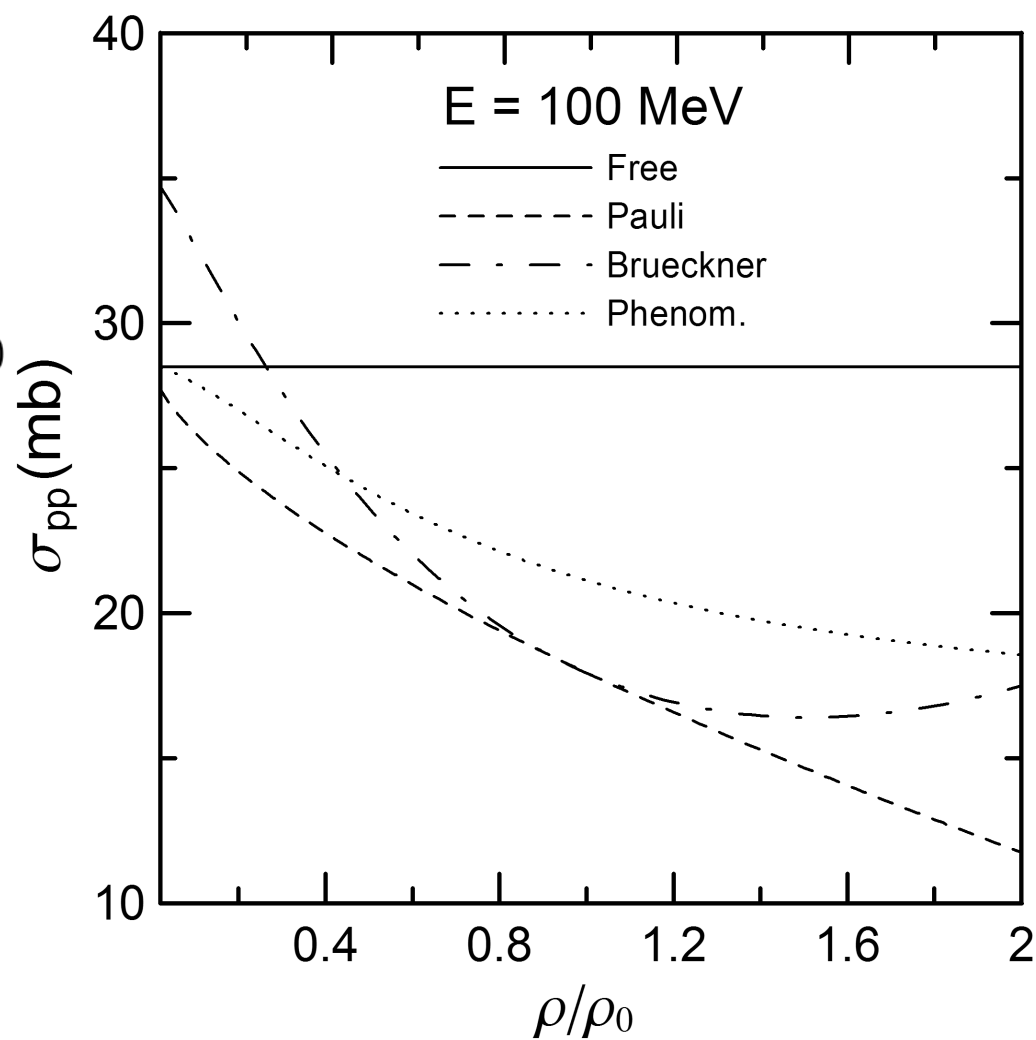
Elastic Scattering with ME in σ_{NN} 

Chulkov, Bertulani, Korshenninikov NPA 587, 291 (1995)

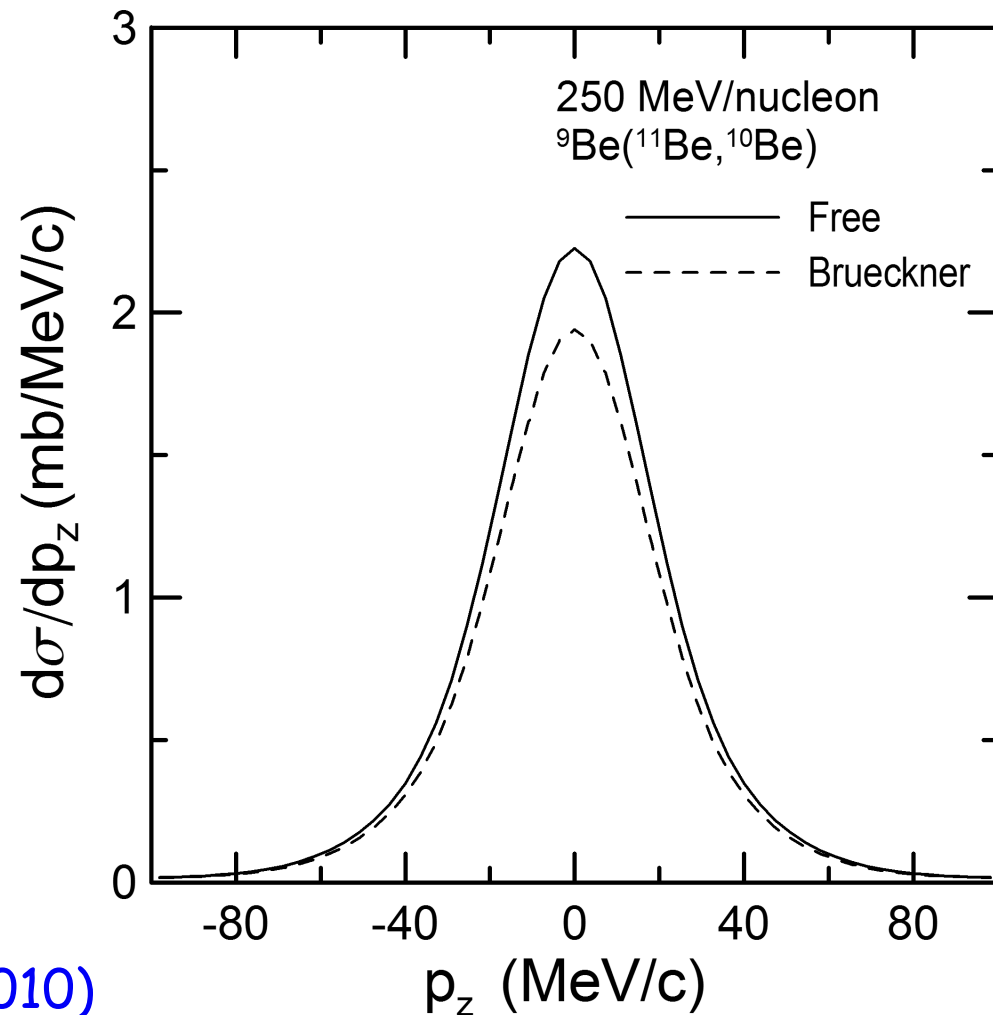
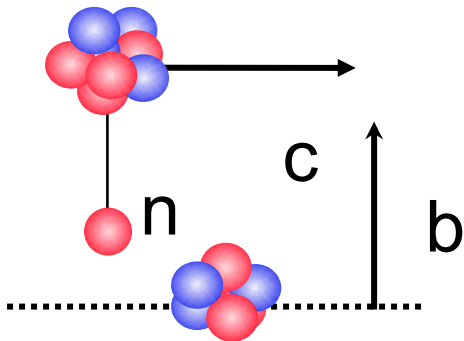
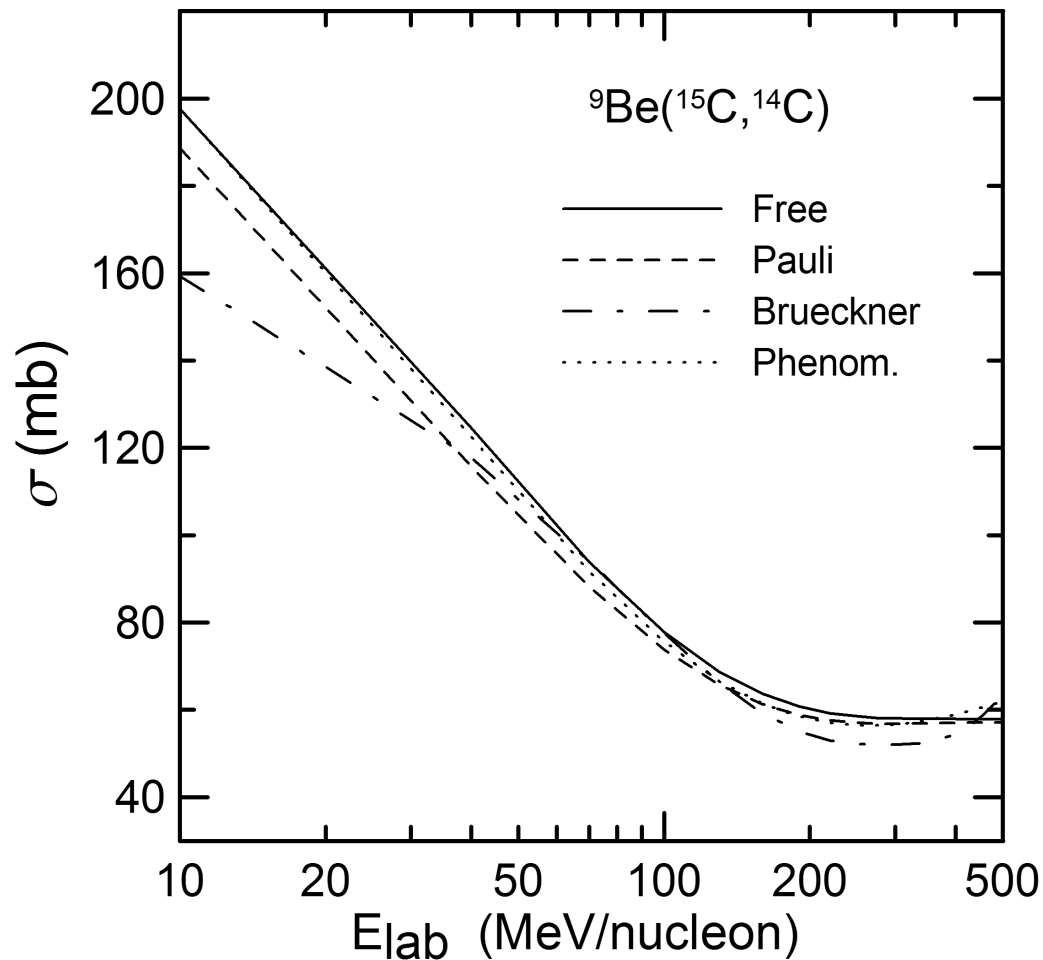
Medium effects in σ_{NN}



Bertulani, De Conti,
PRC C 81, 064603 (2010)

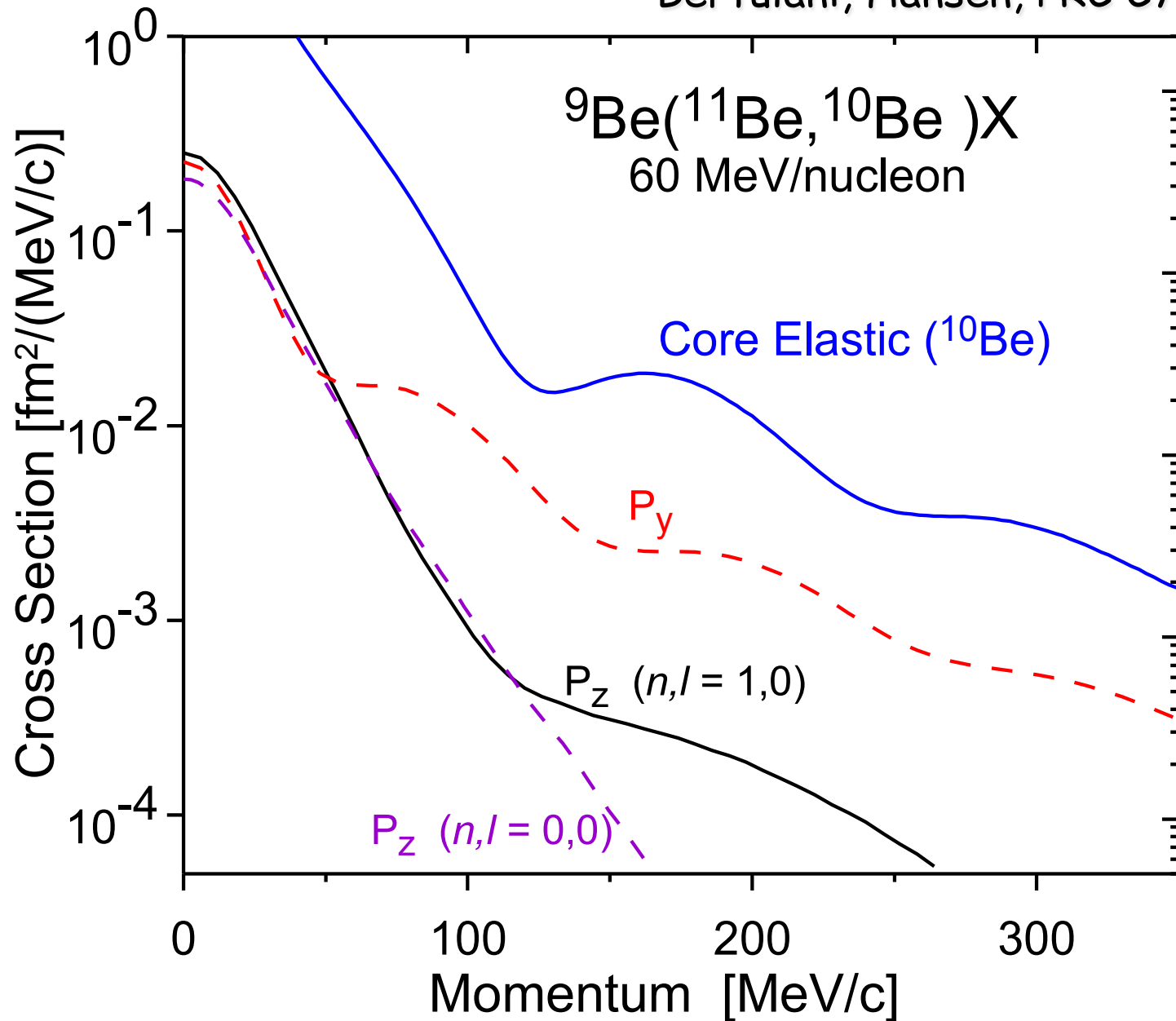


Medium effects in knockout reactions and momentum distributions



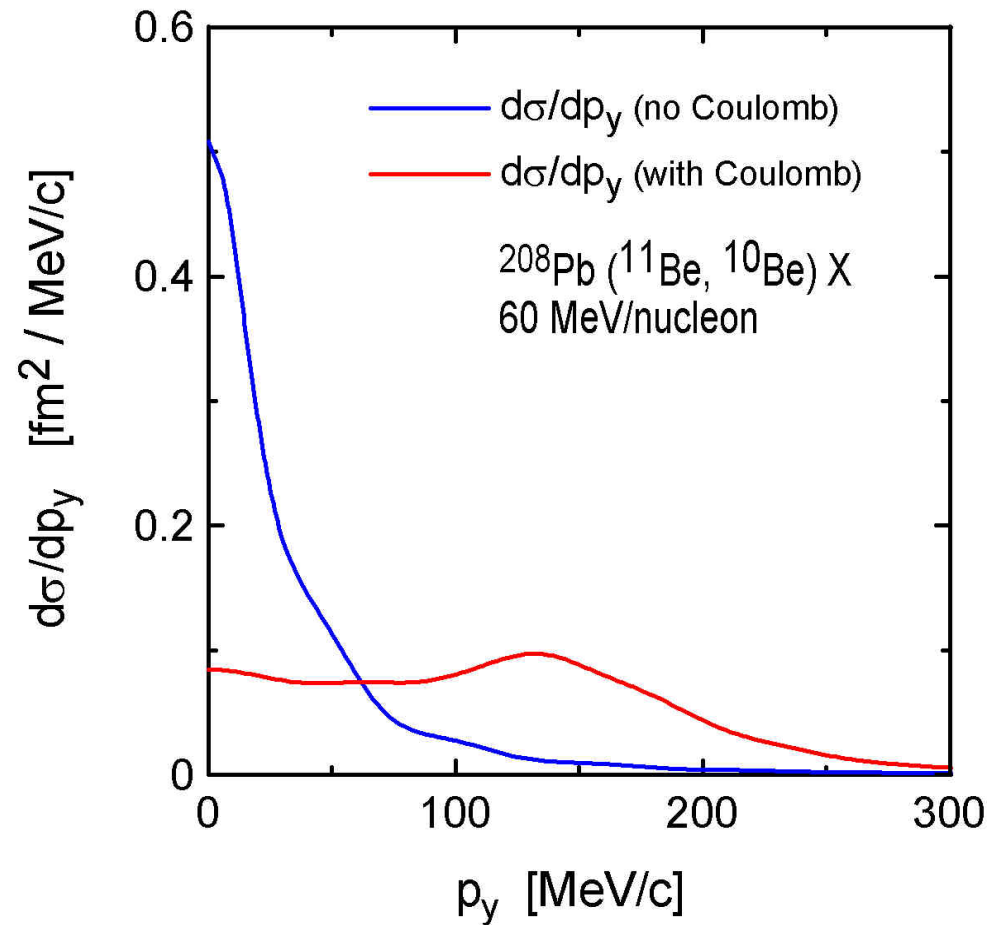
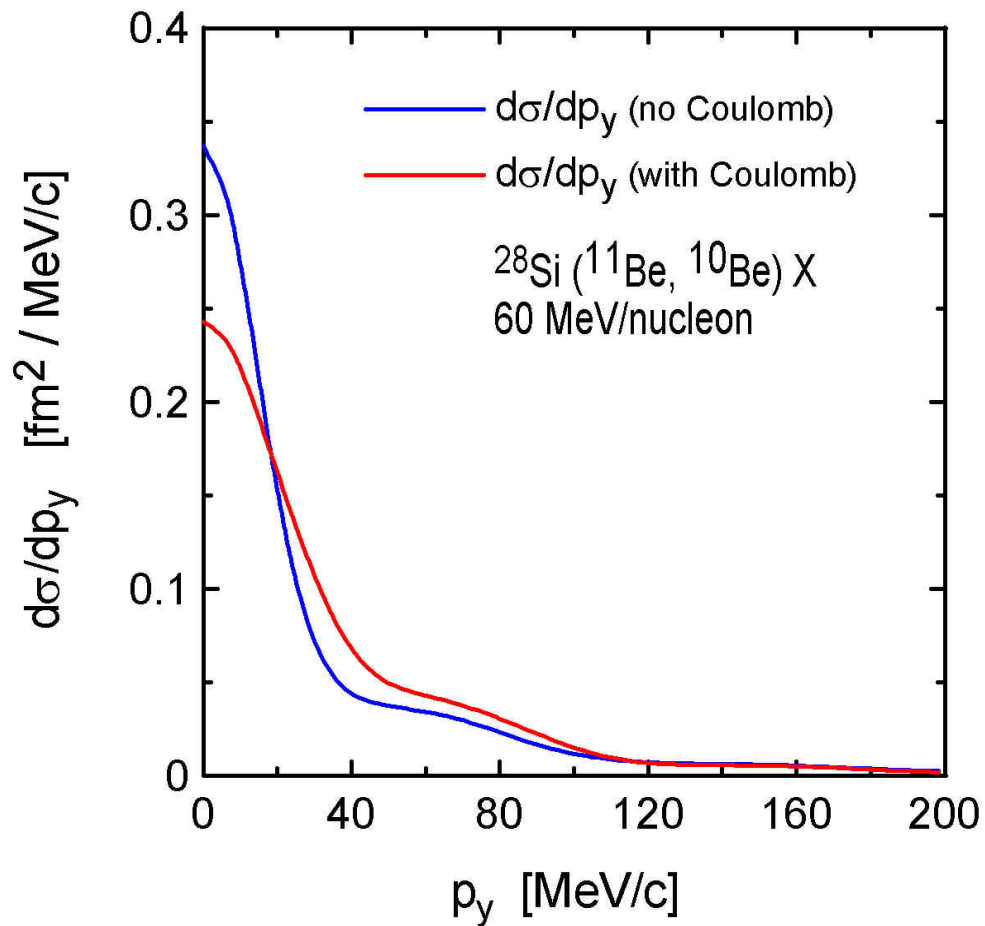
Final state core-target scattering

Bertulani, Hansen, PRC C70, 034609 (2004)

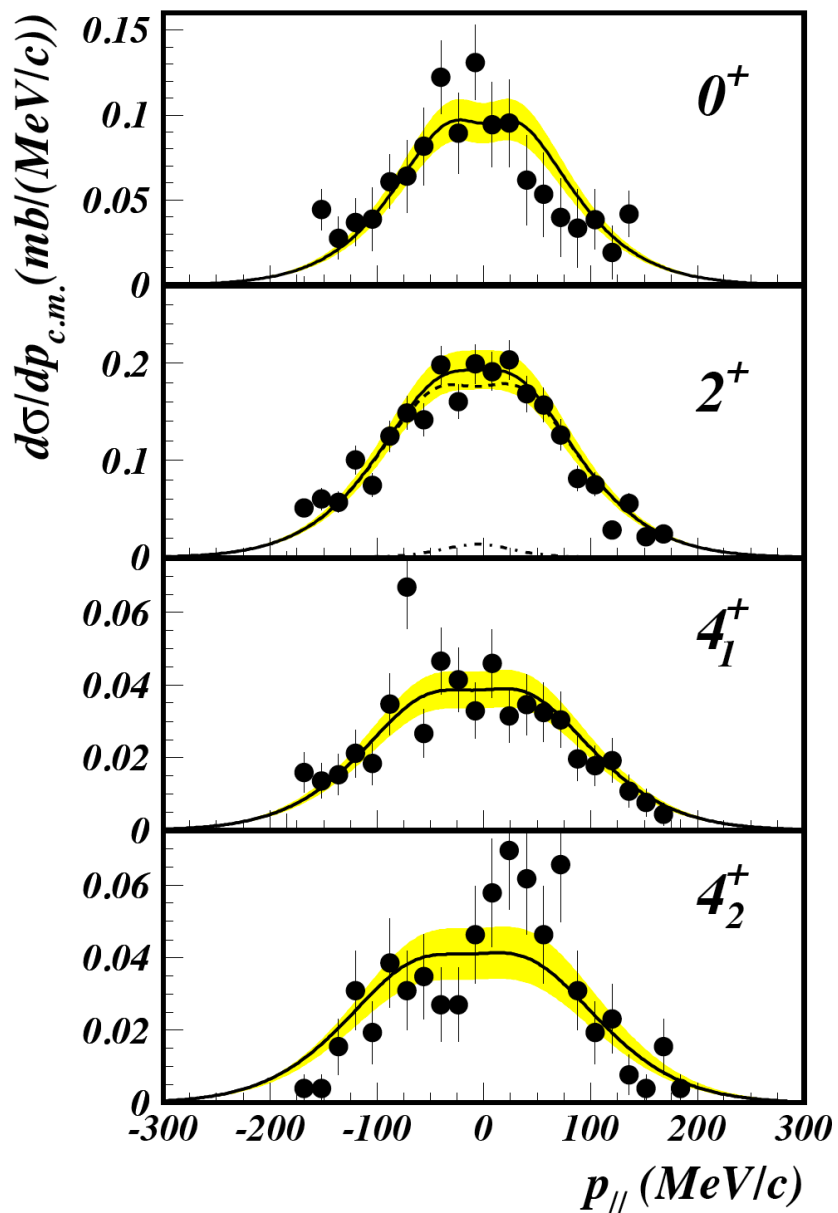


Coulomb reacceleration

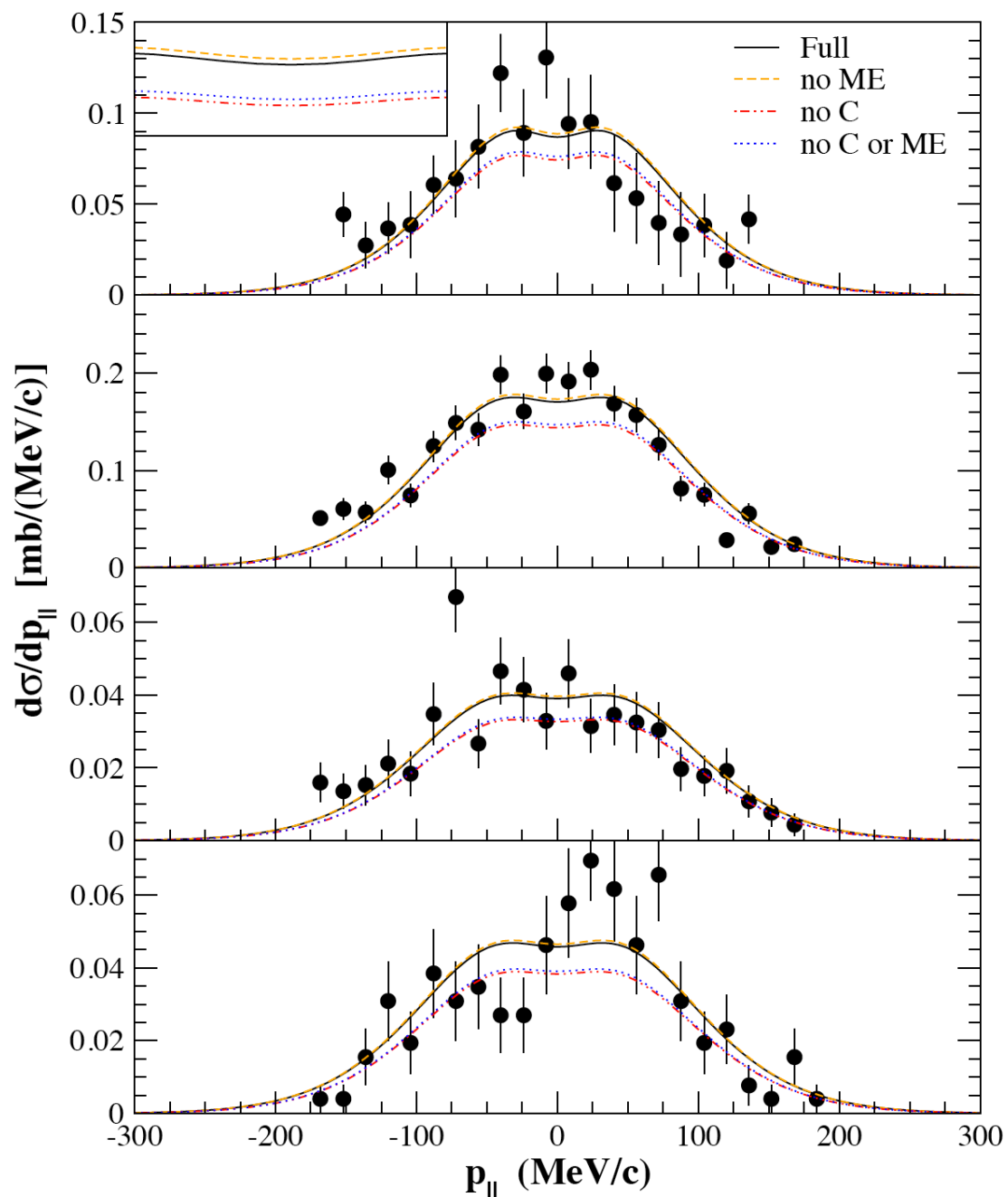
Bertulani, Hansen, PRC C70, 034609 (2004)



Data: Banu et al, PRC 84, 015803 (2011)



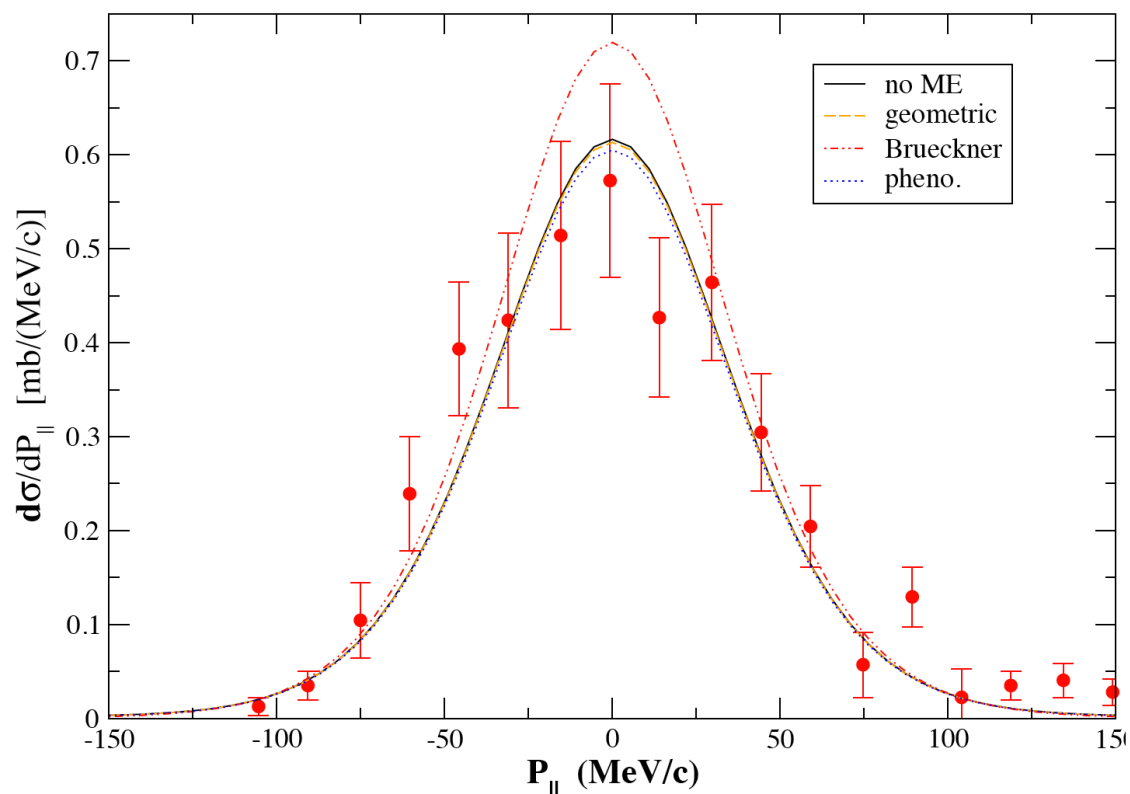
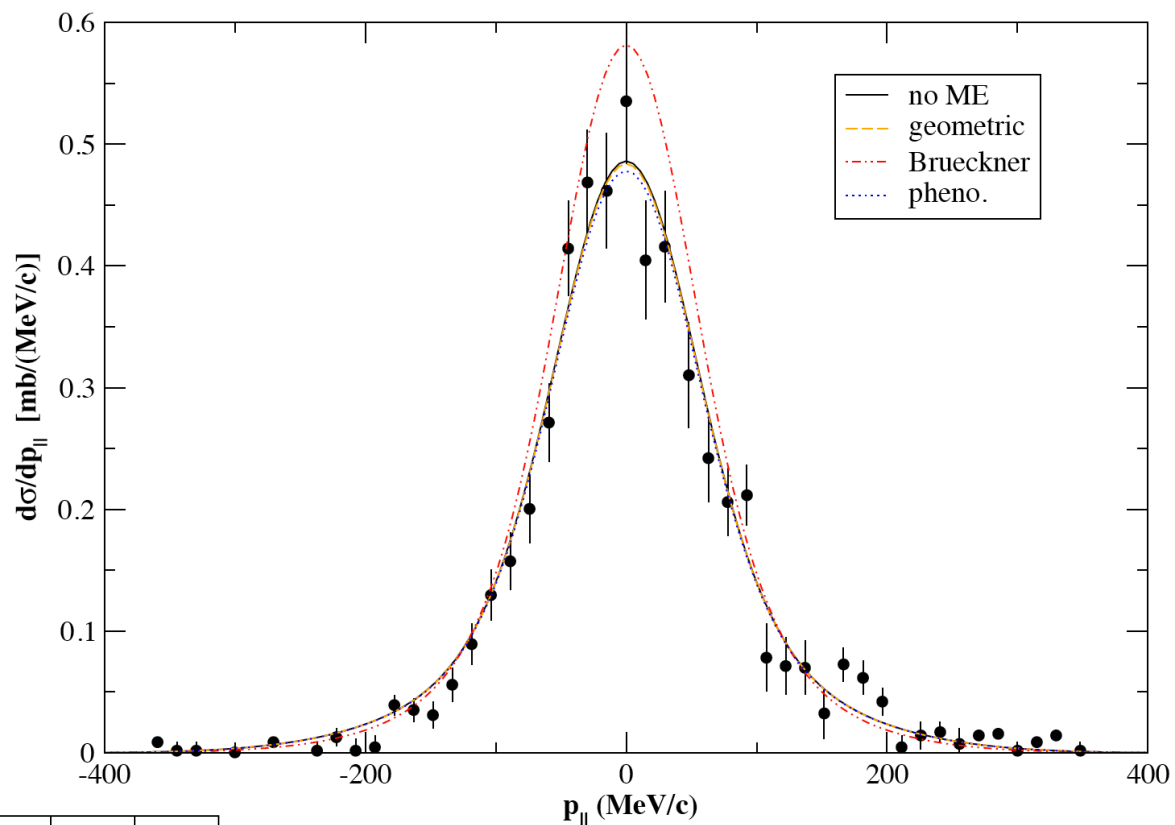
$^{12}\text{C}(^{23}\text{Al}, ^{22}\text{Mg})\text{X}$ @ 50 MeV/u



Karakoc, Bertulani, to be published

Data: Kanungo et al,
PRL 102, 152501 (2009)

$^{12}\text{C}(^{24}\text{O}, ^{23}\text{O})\text{X}$ @ 920 MeV/u



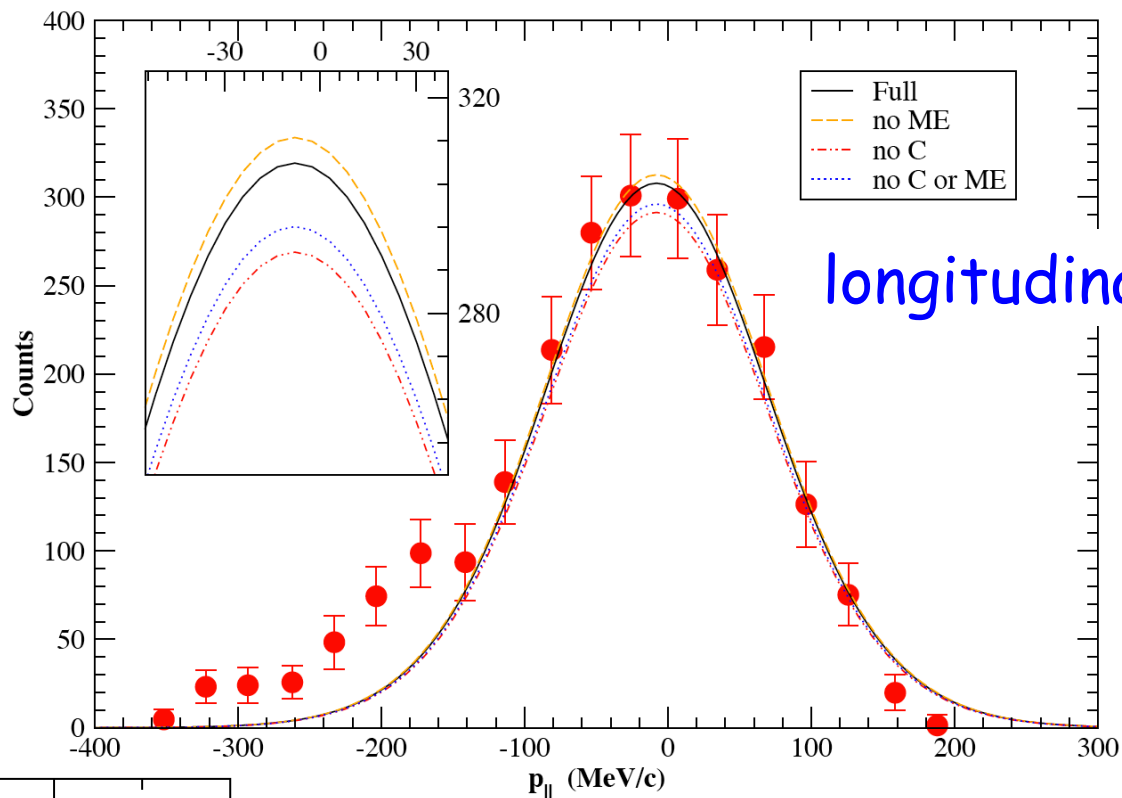
Data: Kanungo et al,
PLB 685, 253 (2010)

$^{12}\text{C}(^{33}\text{Mg}, ^{32}\text{Mg})\text{X}$ @ 898 MeV/u

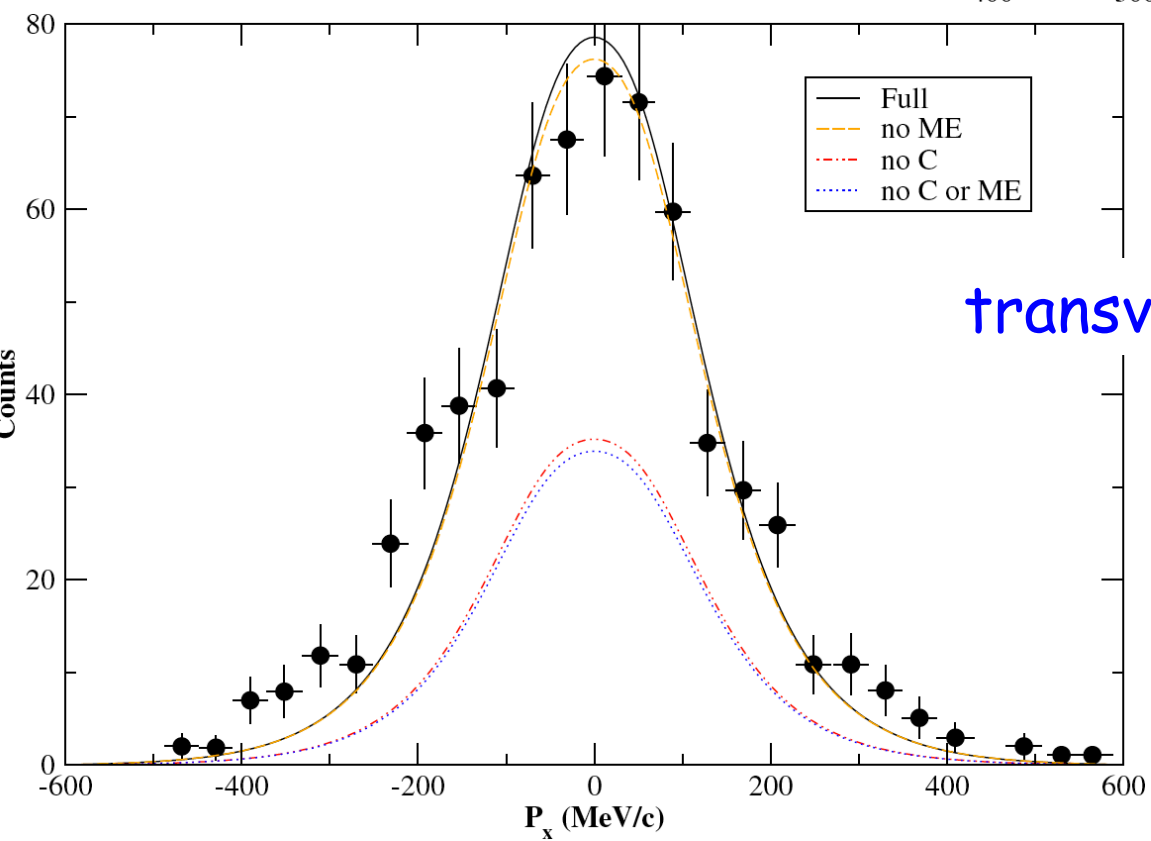
Karakoc, Bertulani, to be published

Data: Jeppesen et al,
NPA 739, 57 (2004)

${}^9\text{Be}({}^{15}\text{O}, {}^{14}\text{N})\text{X}$ @ 56 MeV/u



longitudinal



transverse

Data: Lecouey et al,
PLB 672, 6 (2009)

${}^{12}\text{C}({}^{17}\text{C}, {}^{16}\text{B})\text{X}$ @ 35 MeV/u

Karakoc, Bertulani, to be published

OUT OF TIME