The Challenges of Modeling Explosive Phenomena

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in collaboration with

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Convection in Astrophysics

- Evolution of many stellar systems dominated by convective transport of energy
	- Supernovae (both thermonuclear and gravitational)
	- X-ray bursts and novae (thermonuclear explosion of accreted material on a surface of compact object)
	- General stellar evolution, including post main-sequence evolution of massive stars
- Often the convection is highly subsonic
	- Challenging for traditional astrophysical hydrodynamics codes
- New algorithms are needed for efficient simulation of convective astrophysical flows

Type Ia Supernovae Observations

- Bright as host galaxy, L \sim 10⁴³ erg s⁻¹
- Large amounts of 56 Ni produced
	- Radioactivity powers the lightcurve

SN 1994D (High-Z SN Search team)

- No H seen in spectra, but strong Si, Ca, and Fe lines
- Occur in old stellar populations
- Lightcurve is robust
	- Variations can be corrected for via a single parameter function.
	- SNe Ia act as standard candles.

Type Ia Supernovae (single-degenerate scenario)

(David A. Hardy & PPARC)

1 Accretion from binary companion. Grows to M_{ch}

> 2 "Smoldering" phase—central T $rises \rightarrow$ flame born

SN 1994D (High-Z SN Search team)

3 Flame propagation. Initially subsonic, but detonation transition?

4 Explosion! Lightcurve powered by Ni decay. Width / luminosity relation.

(Roepke and Hillebrandt 2005)

Outstanding Questions in SNe Ia

- What is the progenitor?
	- Alternate models exist, including the merger of two white dwarfs
- Does the burning front remain subsonic or does it transition to a detonation?
	- A late time transition to a supersonic burning front (detonation) appears to give the best match to observations. Turbulence likely plays a key role in this transition.
- What are the initial conditions?
	- Variations in the spatial and temporal distribution of hot spots gives different explosion outcomes (Gracia-Senz & Bravo 2005, Plewa et al. 2004, Roepke et al. 2007, ...)
- What is the physical basis for the width-luminosity relationship in the lightcurve?
	- Some variation in the explosion is needed to account for the diversity in explosions.

Each of these questions requires a unique code.

▼ Roepke and Hillebrandt: ignition seeds in many points distributed around the center.

Why Study Ignition?

- **Explosion outcome very** sensitive to spatial and temporal distribution of initial flames (ignition points)
	- Single point on/off-center vs. multi-point explored by various groups
- Majority of explosion calculations begin with no initial velocity field

◀ Jordan et al. 2007: Single off-centered ignition point leads to very asymmetric explosion. Also discussed in Plewa et al. 2004, Roepke and Woosley 2006.

... what does nature do?

Multidimensional Simulations

- Nature is 3-d
	- Convection driven by nuclear energy release
	- Fluid instabilities / turbulence
	- Localized burning/runaway
	- Rotation
- Challenging simulations
	- Large computing / memory requirements
	- Making sense of enormous amounts of data
- SNe Ia begin with periods of low speed convection driven by nuclear energy release
	- Requires ability to model the domain for long timescales
- Requires a different algorithmic approach than those traditionally used in astrophysics

Simulating Low Mach Phenomena

With explicit timestepping, information cannot propagate more than one zone per step:

$$
\Delta t = \min \left\{ \frac{\Delta x}{|u| + c} \right\}
$$

For $M \equiv |u|/c \ll 1$ this is

$$
\Delta t \approx \frac{\Delta x}{c}
$$

We would like to have

$$
\Delta t \approx \frac{\Delta x}{|u|}
$$

For very low Mach number flows, it takes $\sim 1/M$ timesteps for a fluid element to move more than one zone. Can't we do better?

 λ

▶ A Mach 0.01 front moving to the right (a) initially, (b) after 1 step, (c) after 100 steps.

MAESTRO: Low Mach Number Hydrodynamics

Almgren, Bell, Rendleman, & Zingale 2006 ApJ, 637, 922 Almgren, Bell, Rendleman, & Zingale 2006, ApJ, 649, 927 Almgren, Bell, Nonaka, & Zingale 2008, ApJ, 684, 449

- Reformulation of compressible Euler equations
	- Retain compressibility effects due to heating and stratification
	- Asymptotic expansion in Mach number decomposes pressure into thermodynamic and dynamic parts
	- Analytically enforce hydrostatic equilibrium through base state:

$$
\nabla p_0 = \rho_0 g
$$

Elliptic constraint on velocity field

$$
\nabla \cdot (\beta_0 \mathbf{U}) = \beta_0 \left(S - \frac{1}{\bar{\Gamma}_1 p_0} \frac{\partial p_0}{\partial t} \right)
$$

- β_0 is a density-like variable
- $-$ S represents heating sources
- Timestep based on bulk fluid velocity, not sound speed
- Weak scaling to \sim 100,000 processors

Previous SNe Ia Convection Calculations

▶ Hoflich and Stein modeled a 2-d wedge using an implicit code. Found flow caused compression near the center. Suggested ignition near the center.

(Hoflich and Stein 2002)

◀ Kuhlen et al. modeled the convectively unstable region, with the very center cut out. The observed a characteristic dipole feature and suggested that off-center ignition was likely.

(Kuhlen et al. 2006)

No previous calculations have modeled the entire star.

Computational Demands

- **Computer time measured in CPU-hours**
- Single 384^3 run: $~1$ million CPU-hours
- Single 576^3 run: \sim 7 million CPU-hours
	- 10368 processors (1728 MPI tasks \times 6 OpenMP threads/MPI task)
	- 2100 plotfiles, each 18 GB in size $=$ ~40 TB of data for a single run

► The OLCF Cray XT5 jaguarpf machine at ORNL. This machine has 224,000 cores and is currently ranked #2 on the Top500 list.

Pre-SNe Ia Convection: Dipole Convection

Zingale, Almgren, Bell, Nonaka, & Woosley 2009,, ApJ, 704, 196. Nonaka, Aspden, Zingale, Almgren, Bell, & Woosley 2011, in preparation

- Dipole feature seen in previous calculations better described as a jet
	- Asymmetry in radial velocity field
- Direction changes rapidly

Radial velocity field (red = outflow; blue = inflow) in an 1152³ non-rotating WD simulation.

Pre-SNe Ia Convection: Runaway

Zingale, Almgren, Bell, Nonaka, & Woosley 2009, ApJ, 704, 196.

 $\times 10^8$ 8.0 • Temperature increase $\times10^8$ nonlinear 8.0 7.8 – Ignition occurs as T 7.6 crosses 8 x 10 $^{\rm 8}$ K 7.5 7.4 – "Failed" hotspots seen 7.2 toward the end. 7.0 $T_{\rm peak}~~({\rm K})$ 6.8 7.0 6.6 6900 7000 7100 6.5 6.0 2000 3000 4000 5000 6000 7000 1000 t (s)

Pre-SNe Ia Convection: Shear Layer

Zingale, Nonaka, Almgren, Bell, Malone, Woosley 2011, accepted to ApJ

- Clear separation between the convecting and stable regions
- Persists up to ignition
	- Strong shearing here will greatly affect the flame evolution

Vorticity field in a non-rotating model.

Pre-SNe Ia Convection: Ignition Radius Likelihood

Zingale, Nonaka, Almgren, Bell, Malone, Woosley 2011, accepted to ApJ

- Distribution of likely ignition locations
	- Average hotspot radius over 1 s intervals
	- Consider final 200 s of evolution
- Vast majority of hotspots are moving outward from the center
- Off-center ignition likely

 \blacktriangleright Histogram of likely ignition radii from 576³ non-rotating model. Hotspot radii are averaged into 1 s intervals and colored by radial velocity.

Pre-SNe Ia Convection: Multiple Ignition?

Nonaka, Aspden, Zingale, Almgren, Bell, Woosley 2011, in preparation

- Disable burning in a hot spot once it ignites to allow further evolution
- Second hot spot is not present over a short timescale
- Single-point, off-center ignition most likely.

user: ajnonaka Wed May 11 17:37:37 2011

Current Work: Turbulent Properties

- High resolution simulations show a well-resolved turbulent cascade
- Integral scale large/turbulent intensity small \rightarrow turbulence unlikely to affect flame propagation.
- Velocities in stable region are much higher/shearing.

X-ray Bursts

- Thermonuclear runaway in thin accreted H/He layer on surface of a neutron star
- Accretion timescale \sim hours to days
- Runaway timescale \sim seconds
- \cdot > 70 sources known, some with 10s or more individual bursts.
- Potential site for rpprocess nucleosynthesis

Strohmayer et al., 1996, ApJ, 469:L9

Outstanding Questions in XRBs

- How does the fuel spread over the surface?
- How does the ignition begin?
- Is the burning localized?
- Does convection modify the nucleosynthesis?
- What are the effects of rotation?
- Does convection bring ash to the surface?

Localization

- Accreted layer not degenerate enough to localize a hot spot
	- Fizzles out, maybe stirs
	- Similar to nova (Shankar & Arnett 1994)

Fig. 3.13. Initial evolution of a burning hot spot ignited off the equator as seen in a frame rotating with the neutron star. Velocity vectors show the circulation of the fluid induced by the Coriolis forces. The hot spot expands due to burning and drifts west-southwest because of the latitude dependence of the Coriolis force (after Spitkovsky, Levin & Ushomirsky $2002)$

MAESTRO simulation (L = 3.84 m, Δx = 0.5 cm) evolved to $t = 1.5 \times 10^{-5}$ s showing dissiapation of a hot spot (initial $T = 10^9$ K) in a He NS accreted layer. Temperature is shown on the left, nuclear energy generation rate is shown on the right.

- Spitkovsky et al. (2001)
	- Shallow-water calculations of spreading on NS
	- Coriolis force balances lateral spreading of burning front
	- Simplified vertical structure

- Pure He layers
	- Convective energy transport develops quickly
	- Resolution requirements much higher than previously thought
	- 2-d simulations challenging. 3-d?
- Mixed H/He bursts underway
	- Nuclear physics more time consuming
	- Steep composition change challenging

What Can't We Do (Now)?

- Lateral flame propagation w/ resolved nuclear physics
	- Low Mach methods cannot (currently) describe two different scale heights (fuel and ash)
	- Lengthscale for Coriolis force to balance pressure gradient (Rossby length): $L = \sqrt{gH_0/f} \sim$ few km
		- Much bigger domain that we currently use

Novae

- Explosion of surface H layer on white dwarf
	- Similar progenitor system to SNe Ia (but much fainter)
	- Recur when new layer of fuel is accreted (recurrance times of decades to 1000s of years, depending on white dwarf mass)
	- CNO burning required to explain luminosity
	- Core material seen in ejecta

(Young, Corwin, Bryan, and De Vaucouleurs)

Novae

- Biggest issue: dredge-up
	- Underlying C/O needed to catalyze the reaction—enrichment by atleast order-of-magnitude
	- Convection? Shear instabilities during accretion? Mass diffusion? (Livio and Truran 1990)
- Does the WD gain or lose mass as a result of the explosion?
- Can novae be SNe la progenitors?

Novae

- Numerical challenges
	- Numerical mixing at base of accretion layer can artifically enhance burning
	- Expansion is significant—needs to be tracked to accurately model dynamics.

Velocity magnitude ^{15}N abundance

sub-Chandra SNe Ia Models

- Numerical challenges
	- Convective region potentially extends to the "top" of the star—difficult to capture steep gradient.

Slice through 3-d domain of vorticity in He layer on surface of 1.0 $\textsf{M}_{_{\mathbb{O}}}$ WD.

Summary / What's Next?

- Modern algorithms / supercomputers can model convective astrophysical flows for many turnover times in 3-d.
	- Requires involvement of many different disciplines: mathematics, computational science, application scientists.
- Convection in pre-SNe Ia white dwarfs:
	- Rapidly changing convective field
	- Range of ignition locations: between central and \sim 80 km offcenter
	- Details of ignition distribution can be learn by looking at late time fluctuations
	- Single-point, off-center ignition likely

Summary / What's Next?

- X-ray bursts
	- Low Mach number hydrodynamics provides an efficient means for modeling convection on neutron stars
	- Exploring mixed H/He burning underway.
- Applications to nova, sub-Chandra SNe Ia ignition, and H core convection in massive stars underway.
- Future Work
	- Addition of long wavelength acoustics to extend range of validity of low Mach number model
	- Work on understanding limits to method
	- Restart ignition calculation in compressible code to follow the subsequent flame evolution.