

*What can X-ray observations  
tell us about:*

*The role of Gravitational Waves  
in Low Mass X-ray Binaries*

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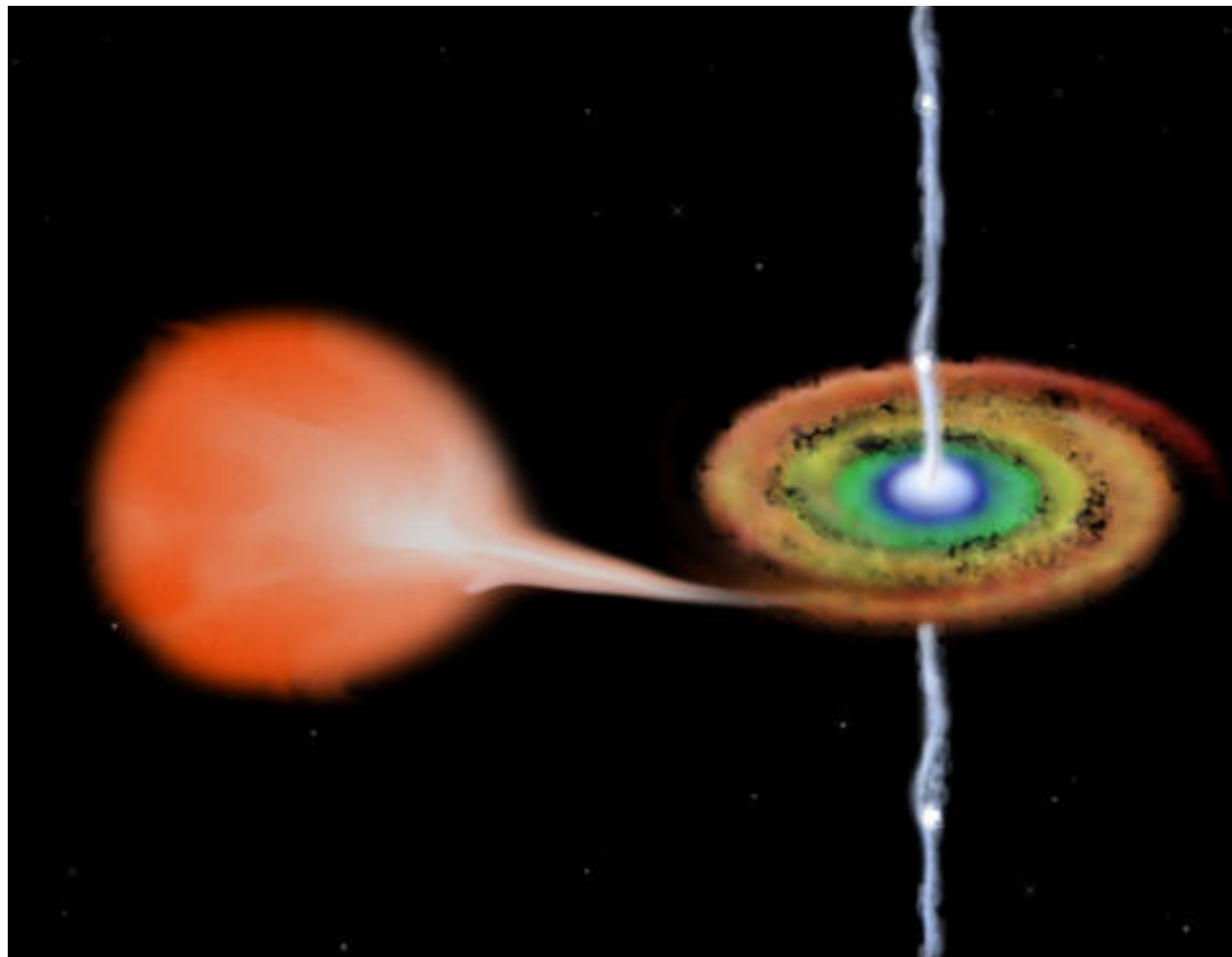
*Brynmor Haskell*

Astronomical Institute "Anton Pannekoek"



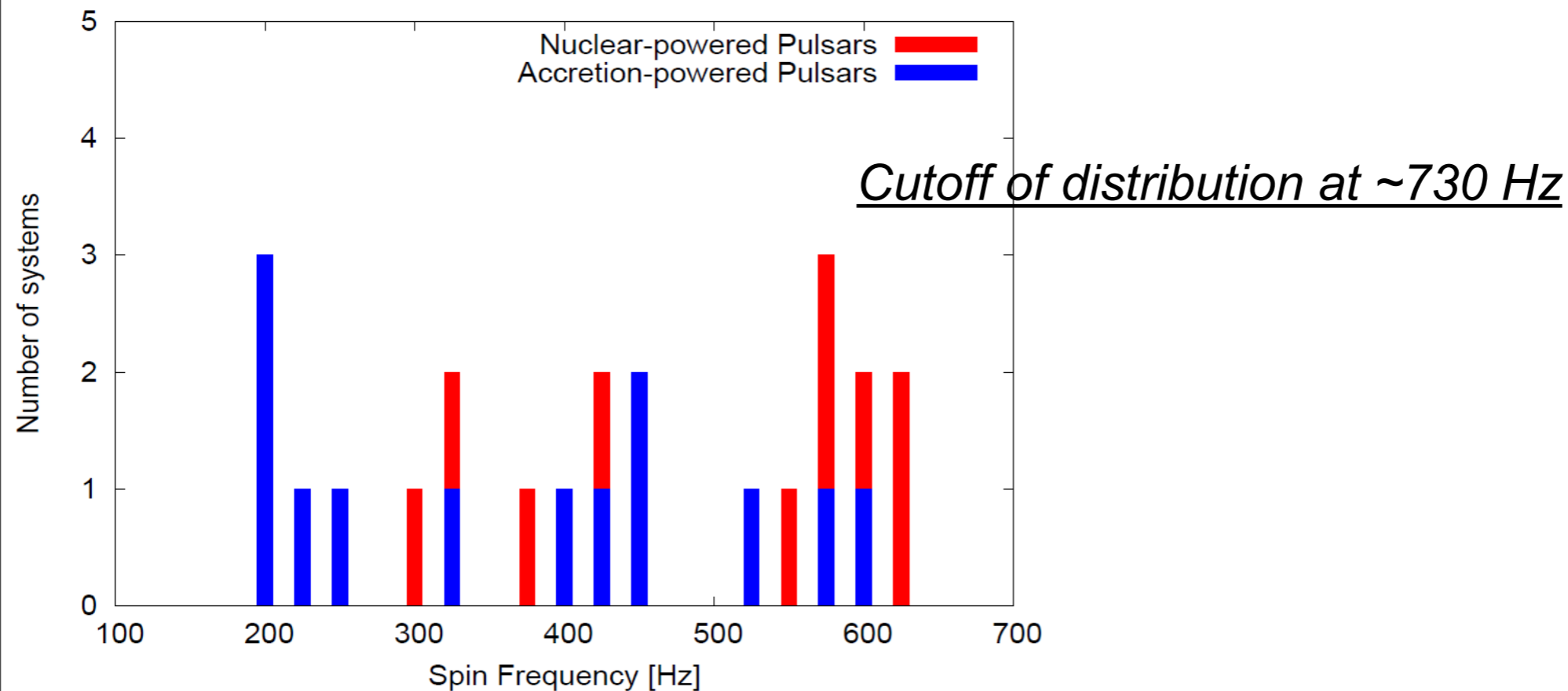
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## Low Mass X-ray Binaries



- Mass is stripped from the donor
- Forms a disc and spirals in
- Interacts with the magnetic field
- Transfers angular momentum to the central NS, spinning it up

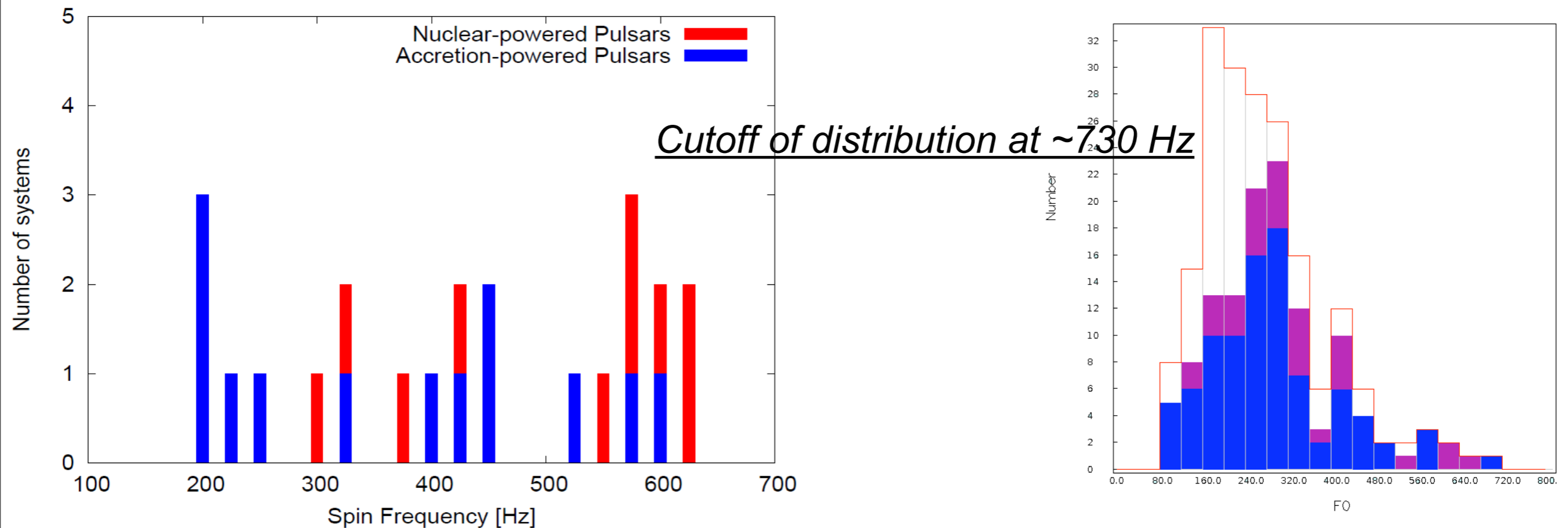
## GWs from LMXBs



(Chakrabarty et al 2003, Patruno 2010)

- LMXB spin distribution points to a mechanism that halts the spin-up before the break up limit.
- GWs!: “mountains”, unstable modes, magnetic deformations.. (Papaloizou & Pringle 1978, Wagoner 1984, Bildsten 1998)

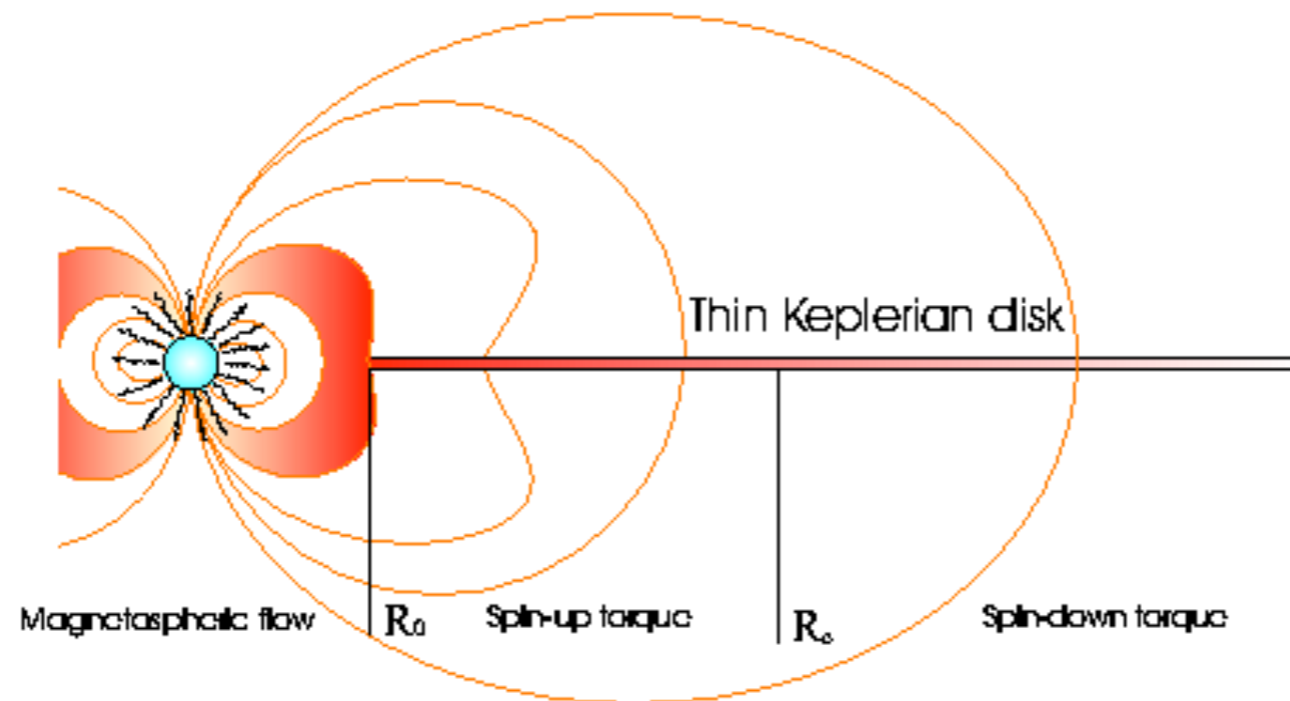
## GWs from LMXBs



(Chakrabarty et al 2003, Patruno 2010)

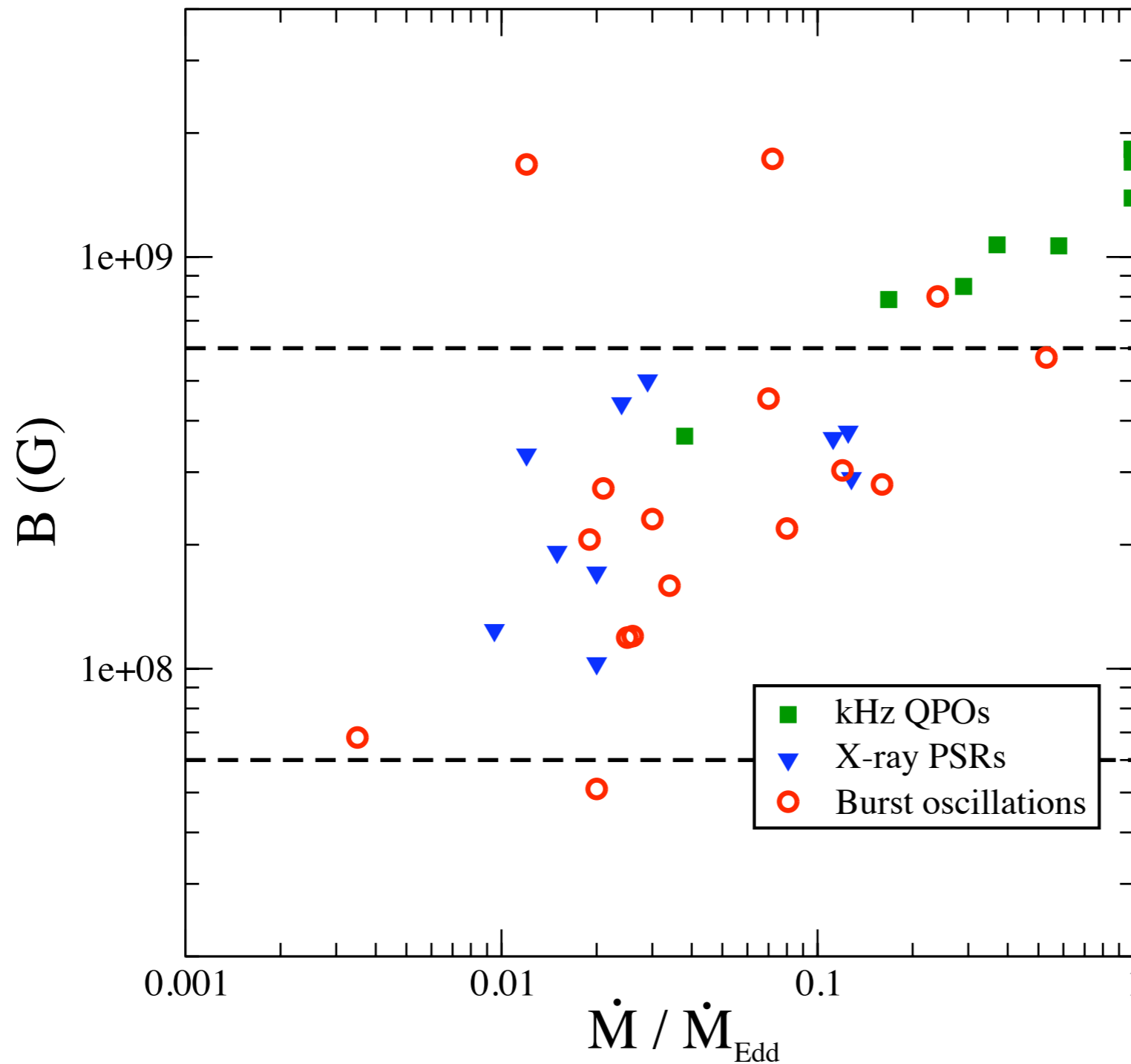
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## Spin equilibrium?



- Interaction at magnetospheric radius
- Accretion torque: spin up
- Magnetic torques and propeller: spin down

# Spin equilibrium?



(White & Zhang 1997)

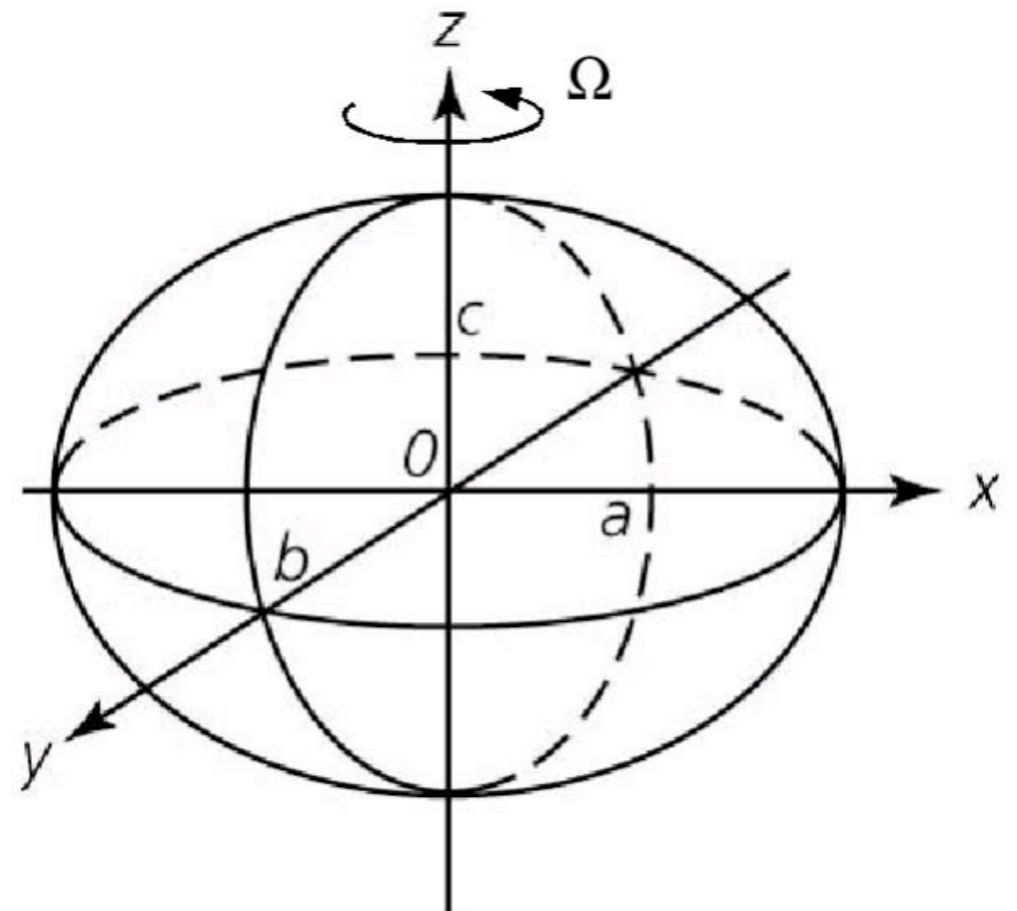
## Neutron star mountains

■ 
$$\epsilon = \frac{I_{xx} - I_{yy}}{I_{zz}}$$

■ Emission at  $\omega = 2\Omega$

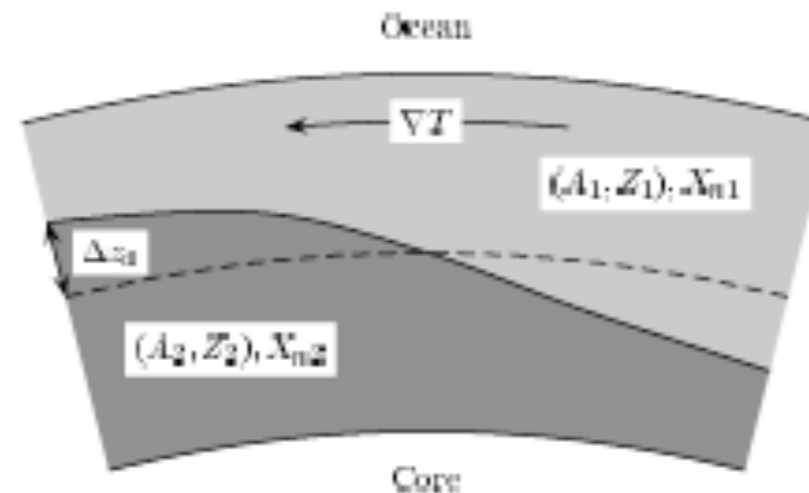
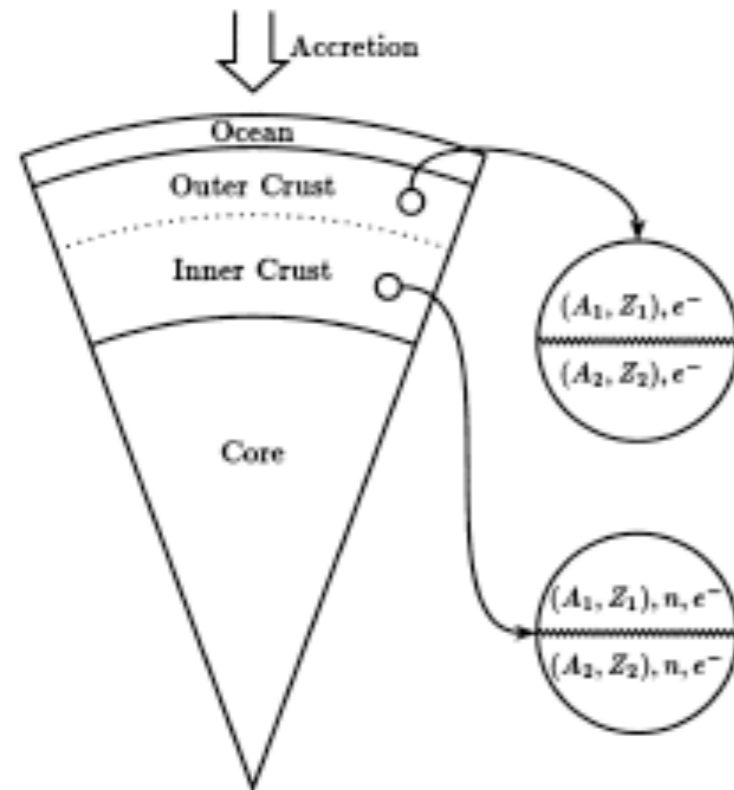
■ 
$$\frac{dE}{dt} \approx \epsilon^2 \Omega^6$$

■ Theoretical upper limit  $\epsilon \approx 10^{-6}$  (Haskell, Jones, Andersson 2006)



## Neutron star mountains-II

- Mountains from 'wavy' capture layers in crust

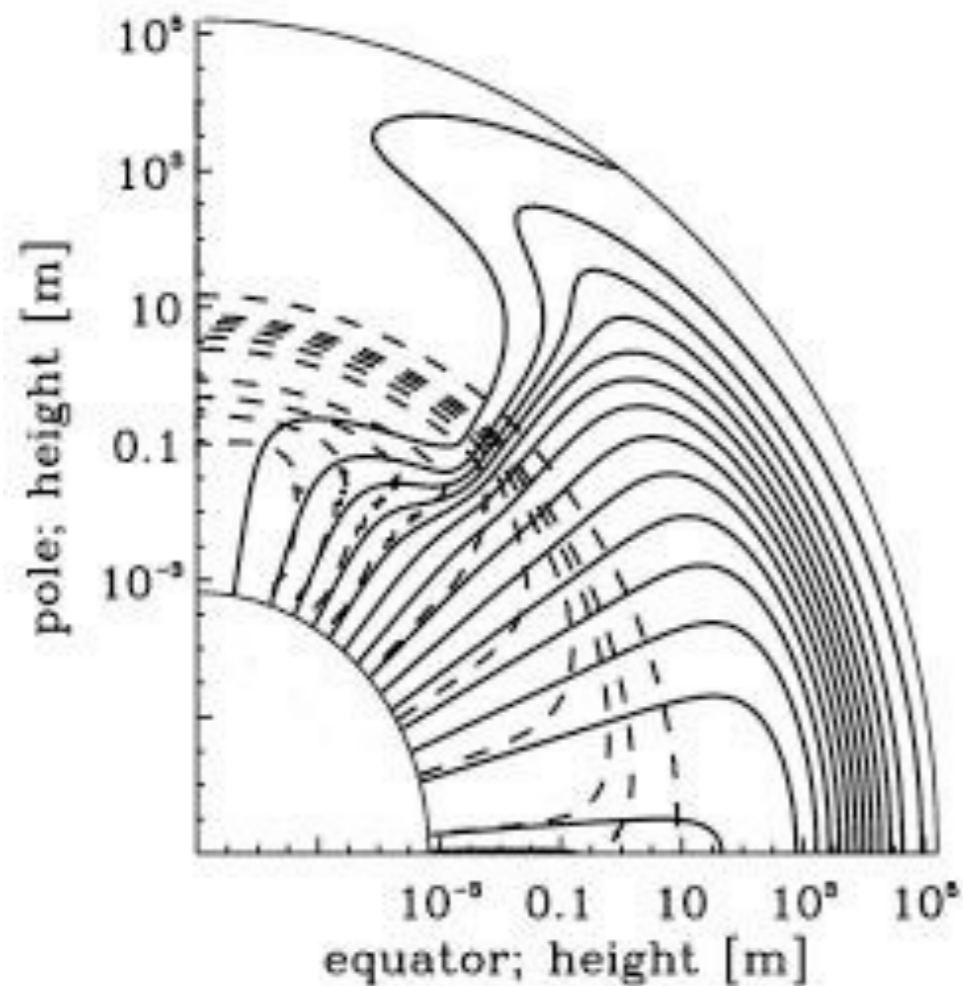


(Ushomirsky, Cutler, Bildsten 2000)

- Deep crustal heating 'consistent' with cooling observations from X-ray transients.



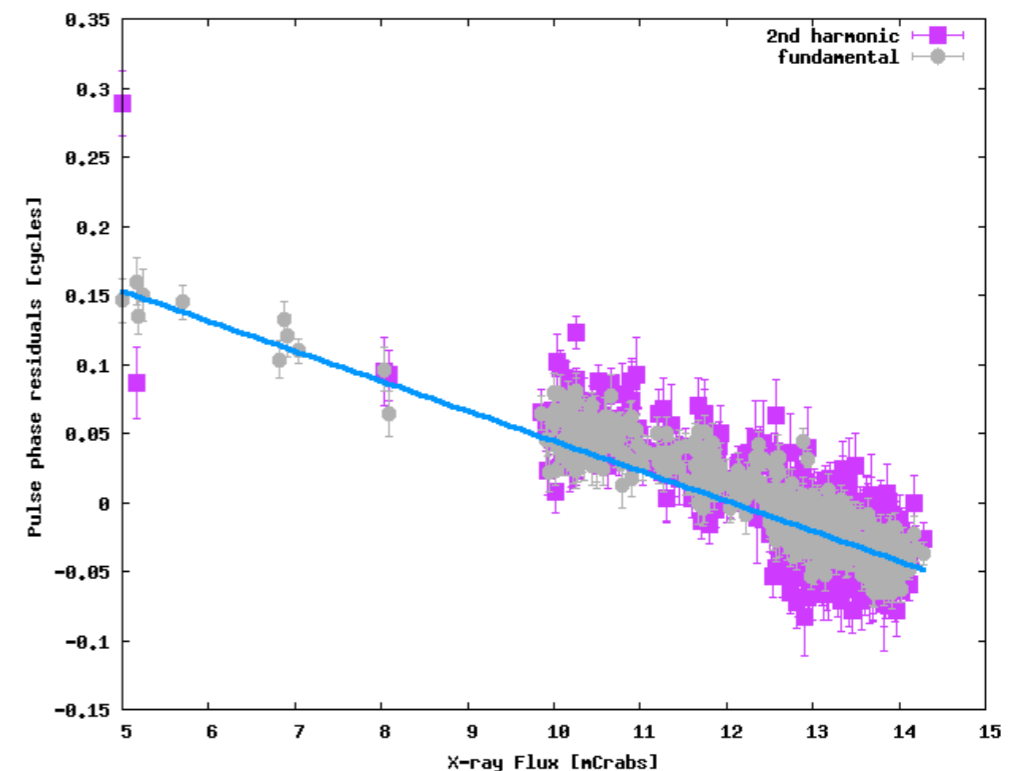
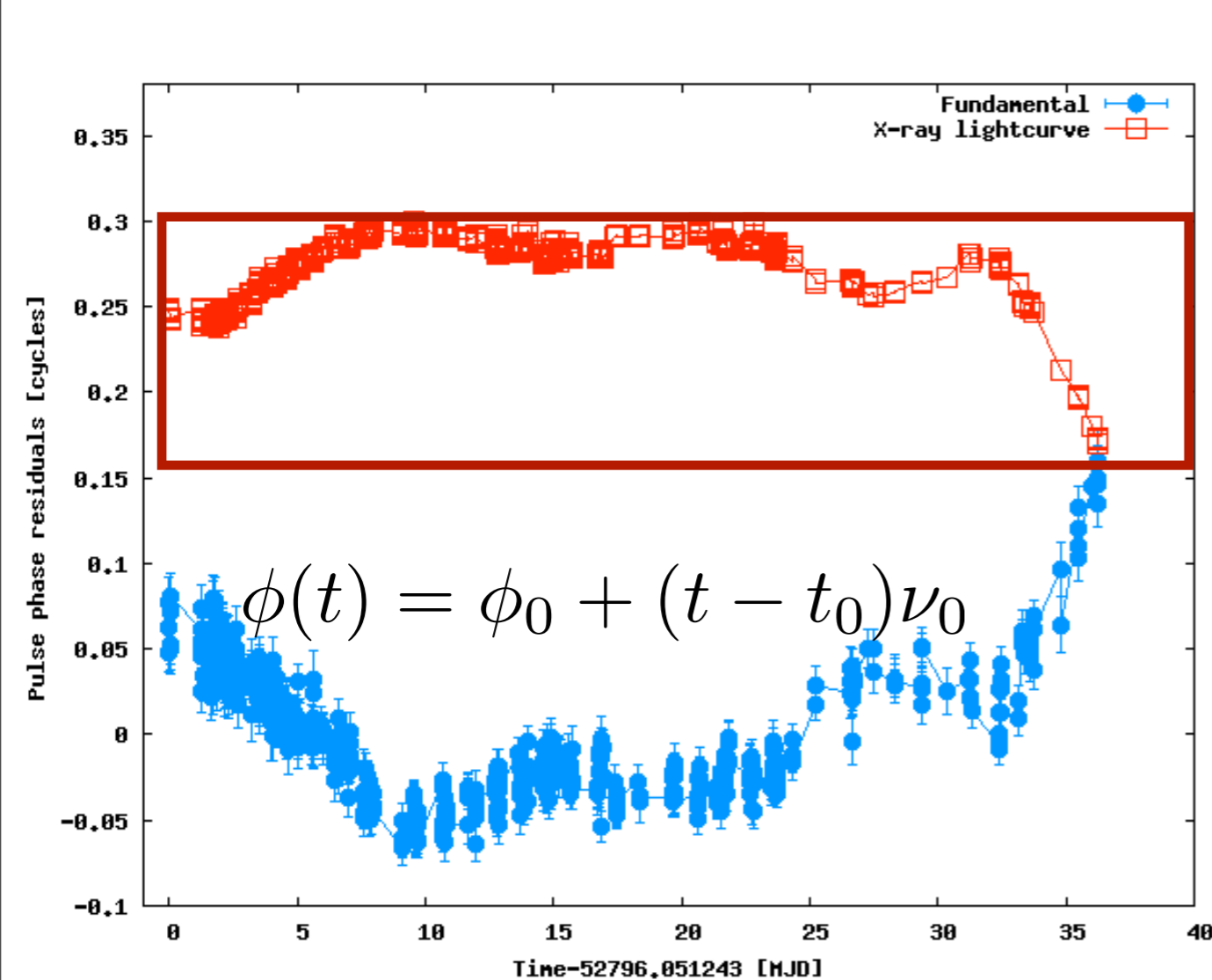
## Magnetic mountains



- Magnetic field distorted by the accretion flow
- Possibility of confining a 'mountain'

(Payne & Melatos 2004, Melatos and Payne 2005)

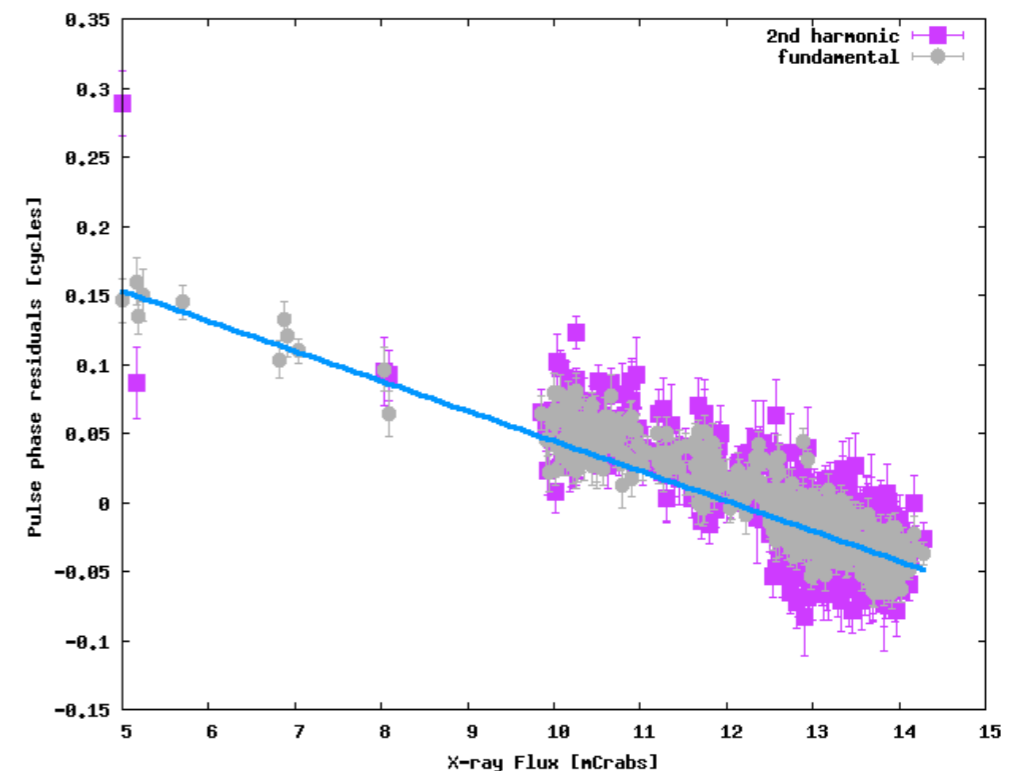
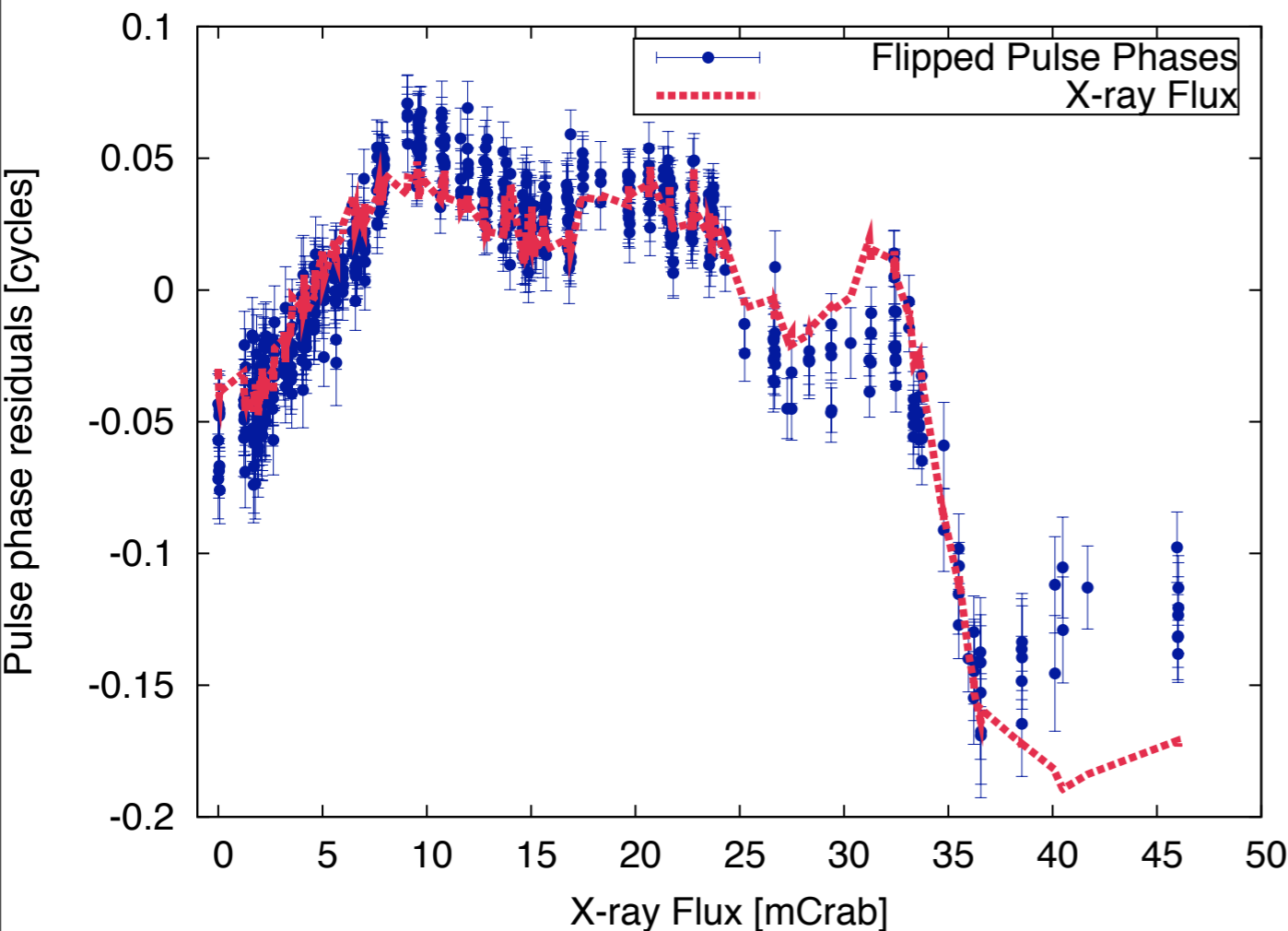
# The strange case of XTE J1814-338 (& SAX J1808.4-3658)



(Patruno, Wijnands & Van der Klis 2009)

- Flux correlated to phase?
- Constant frequency during outburst?  
 $|\dot{\nu}| < 1.5 \times 10^{-14}$
- Where is the angular momentum going? GWs?  
[Haskell & Patruno (2011)]

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## Gravitational waves?

### ■ Crustal mountains.

↓ Ushomirsky, Cutler & Bildsten (2000)

Not enough heat deposited in the crust



$$\delta T \approx 10^3 C_k^{-1} p_{30}^{-1} Q_n \Delta M_{22} \text{ K} \quad \text{Ushomirsky \& Rutledge (2001)}$$

$$Q_{22} \approx 1.3 \times 10^{35} R_6^4 \left( \frac{\delta T_q}{10^5 \text{ K}} \right) \left( \frac{Q}{30 \text{ MeV}} \right)^3 \text{ g cm}^2$$

(To balance accretion one would need  $Q_{22} \approx 10^{37}$  )

## Gravitational waves?

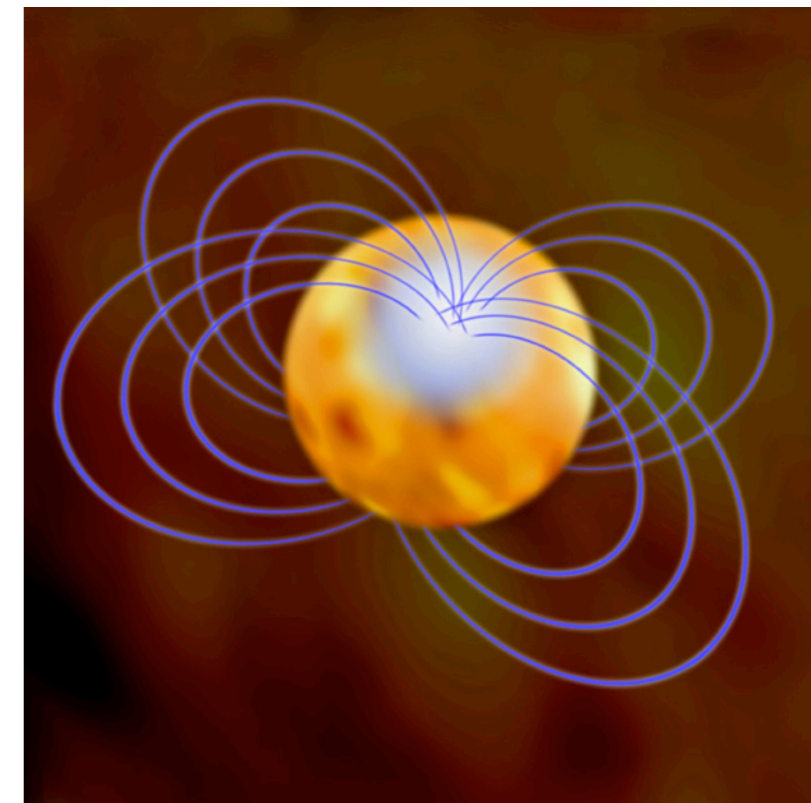
- Magnetic mountains.



Cutler (2002), Melatos & Payne (2005)

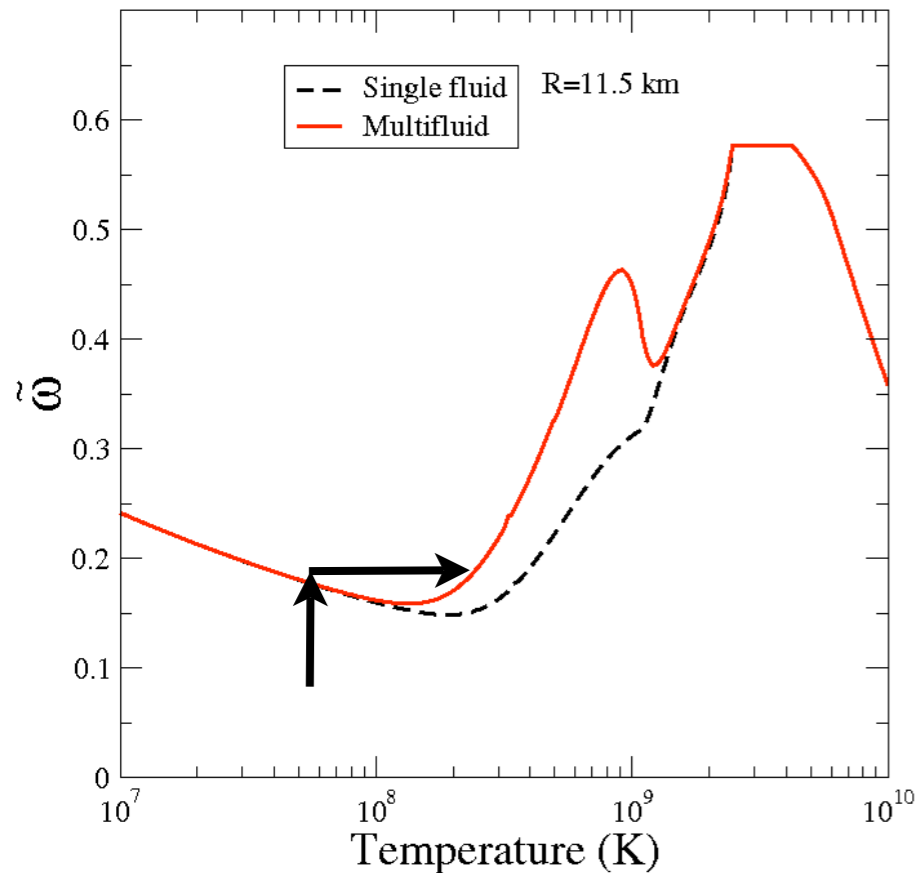
**B field too weak. Spin-down too strong?**

Requires strong internal toroidal  
(or surface higher multipole)  
component of the order  $B \approx 10^{12}$  G





## Gravitational waves?



■ Hyperon bulk viscosity (or quark bulk viscosity, mutual friction etc.) can halt the r-mode thermal run-away.

Andersson, Jones & Kokkotas (2002)

Nayyar & Owen (2006)

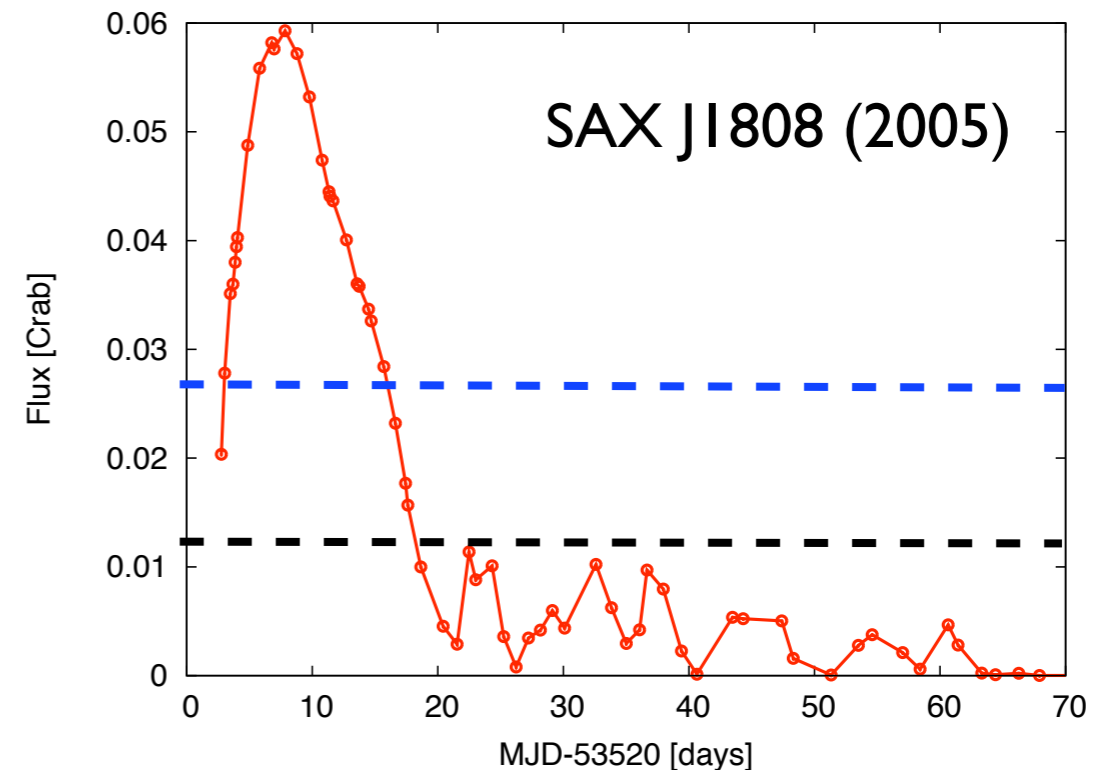
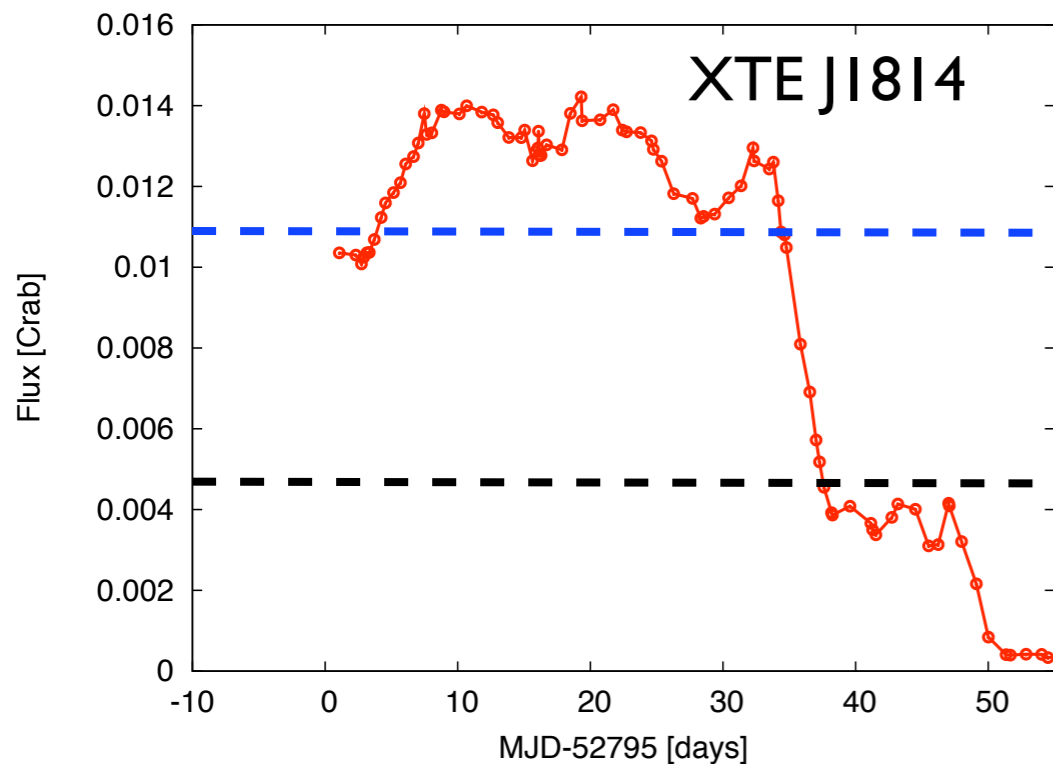
Star too hot and spin-down in quiescence too strong ?

Haskell & Andersson (2010)

Estimated core temperature:  $T_c \approx 10^7 K$

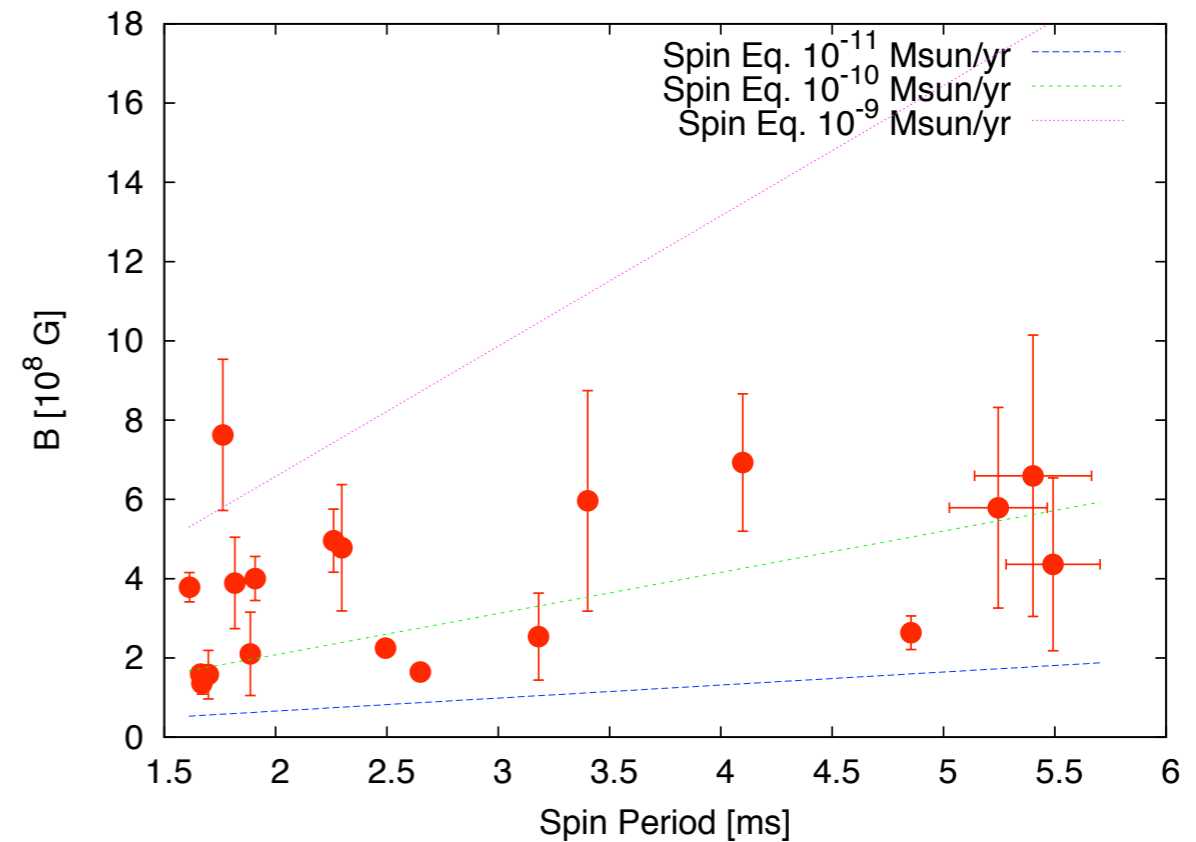
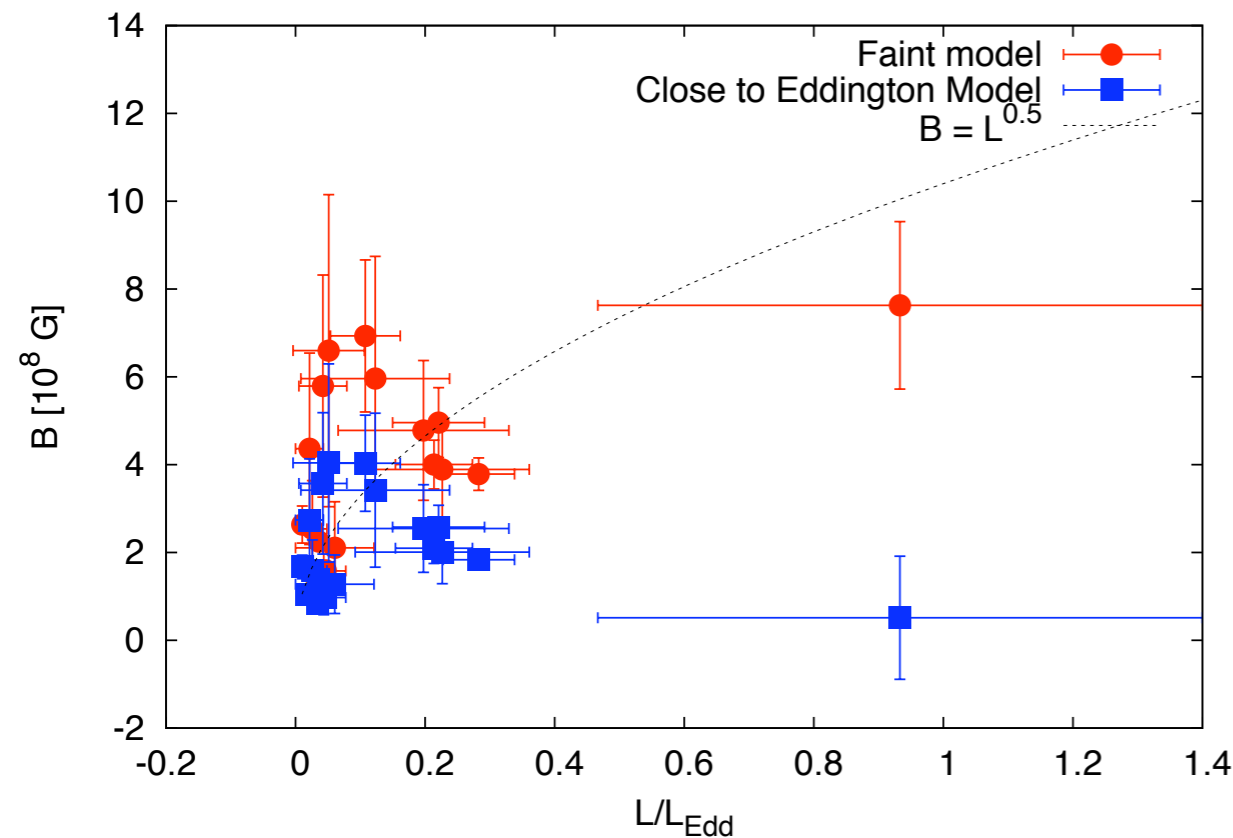
Would need strong direct URCA with no superfluidity

## Spin equilibrium?



- Assume a propeller phase with  $R_c = R_m$  [Haskell & Patruno (2011)]
- The simple model of Andersson et al. 2005 gives spin equilibrium at approximately the mean accretion rate ( $R_m = 0.8 R_c$ )
- $R_m/R_c = 0.75$  XTE J1814,  $R_m/R_c = 0.75 - 0.84$  SAX J1808

# Spin equilibrium?



■ Correlation between  $B$  and  $L$  weak

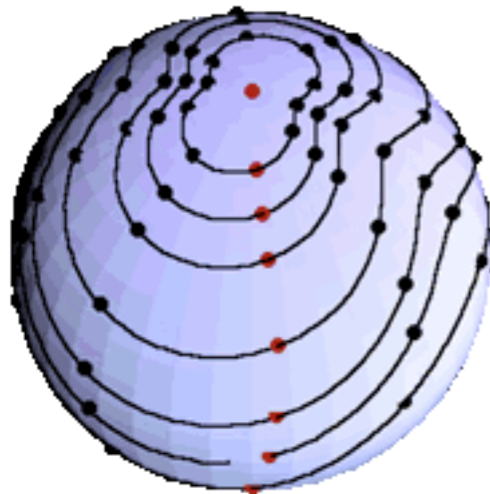
■ Many systems may be close to spin equilibrium as set by the disc/magnetosphere interaction

[ Patruno, Haskell & D'Angelo (in preparation) ]

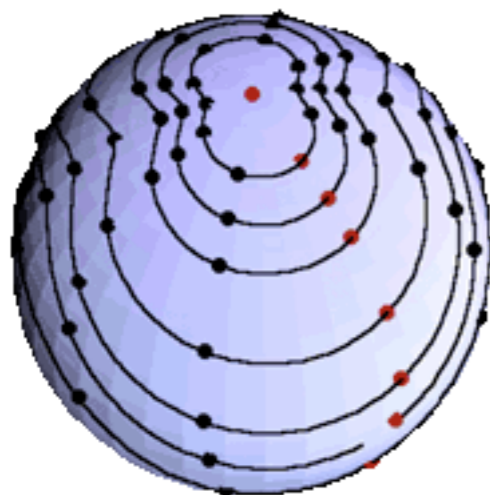


## r-mode instability

(Animation by Ben Owen)



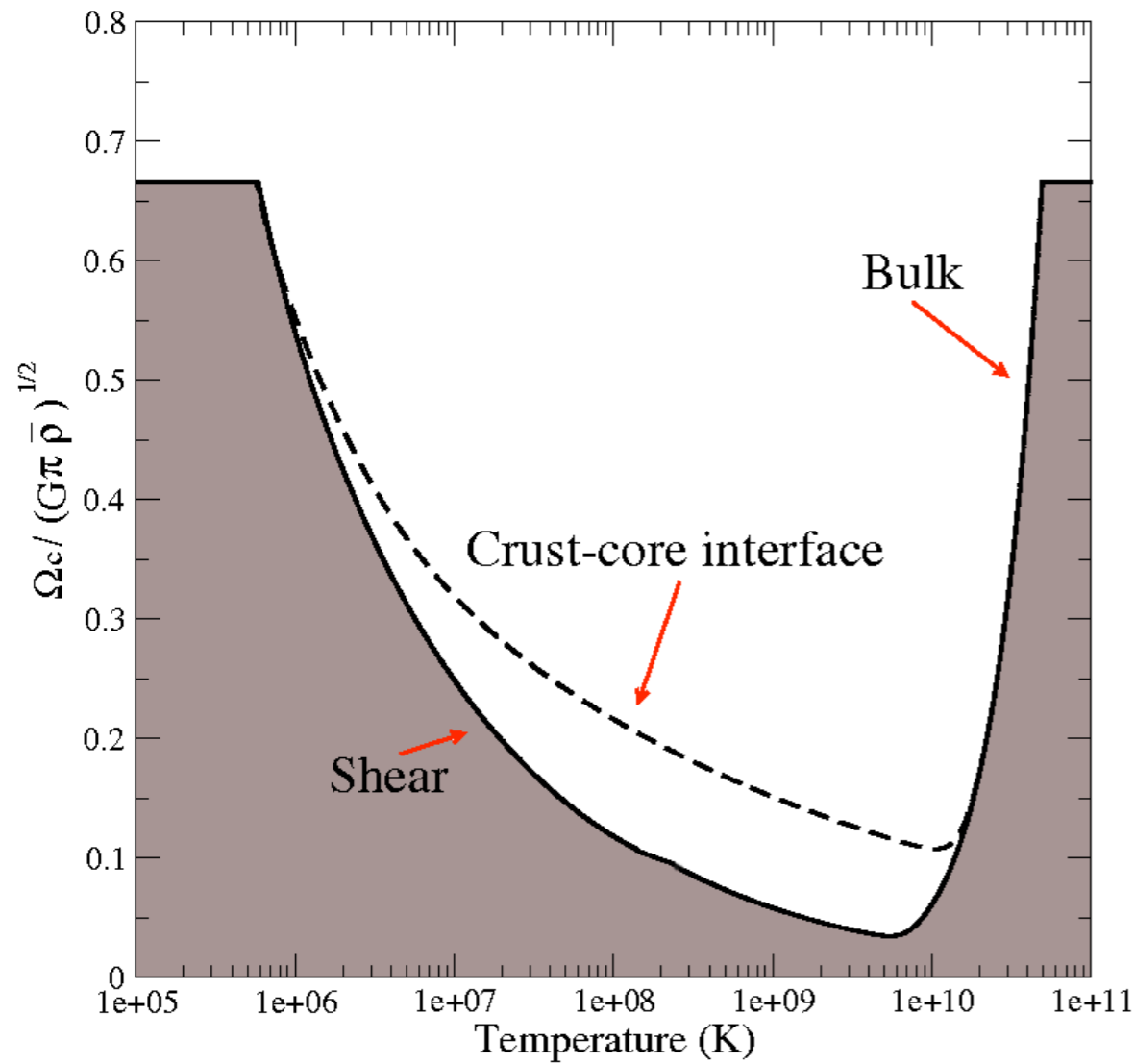
Rotating observer



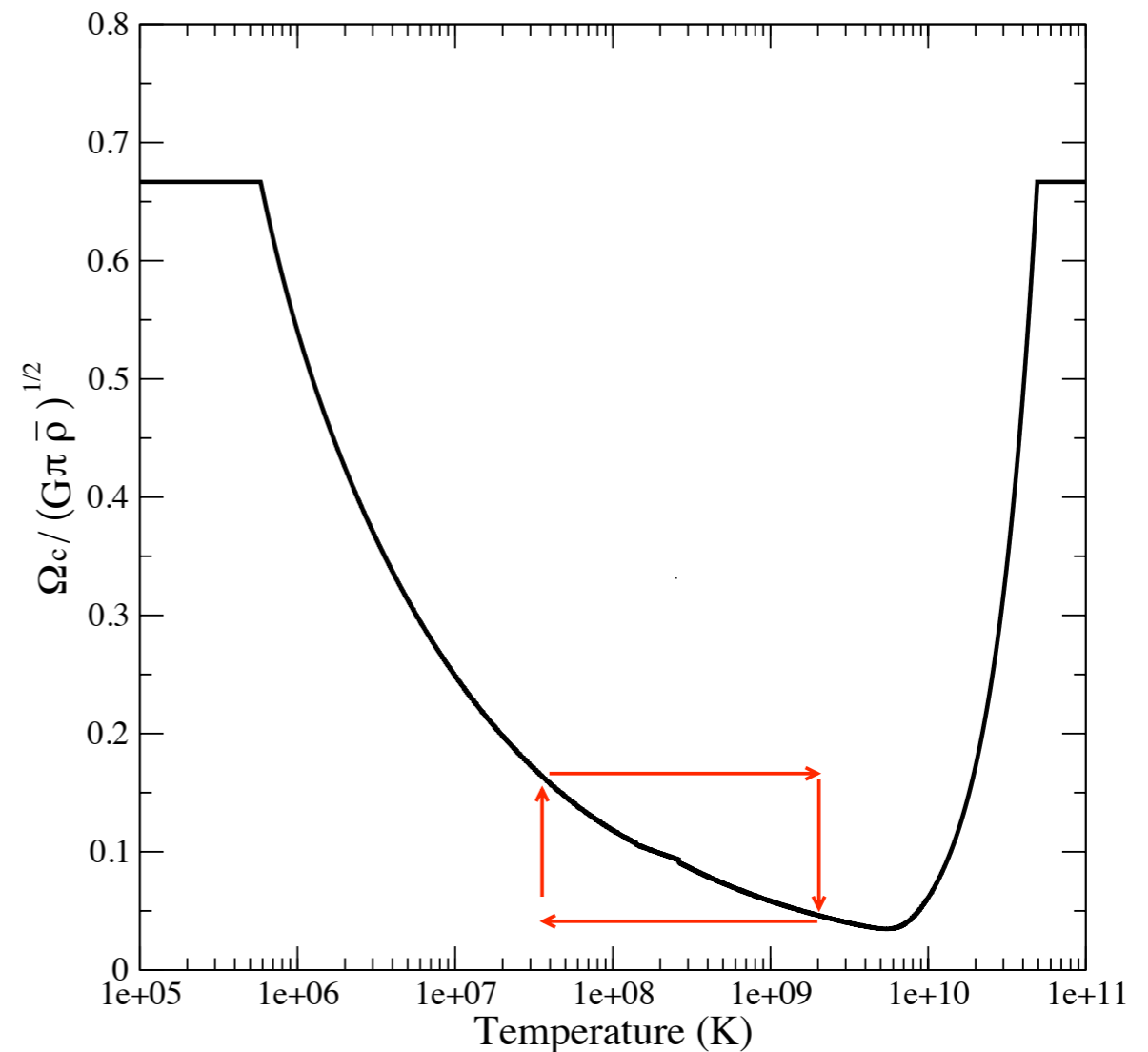
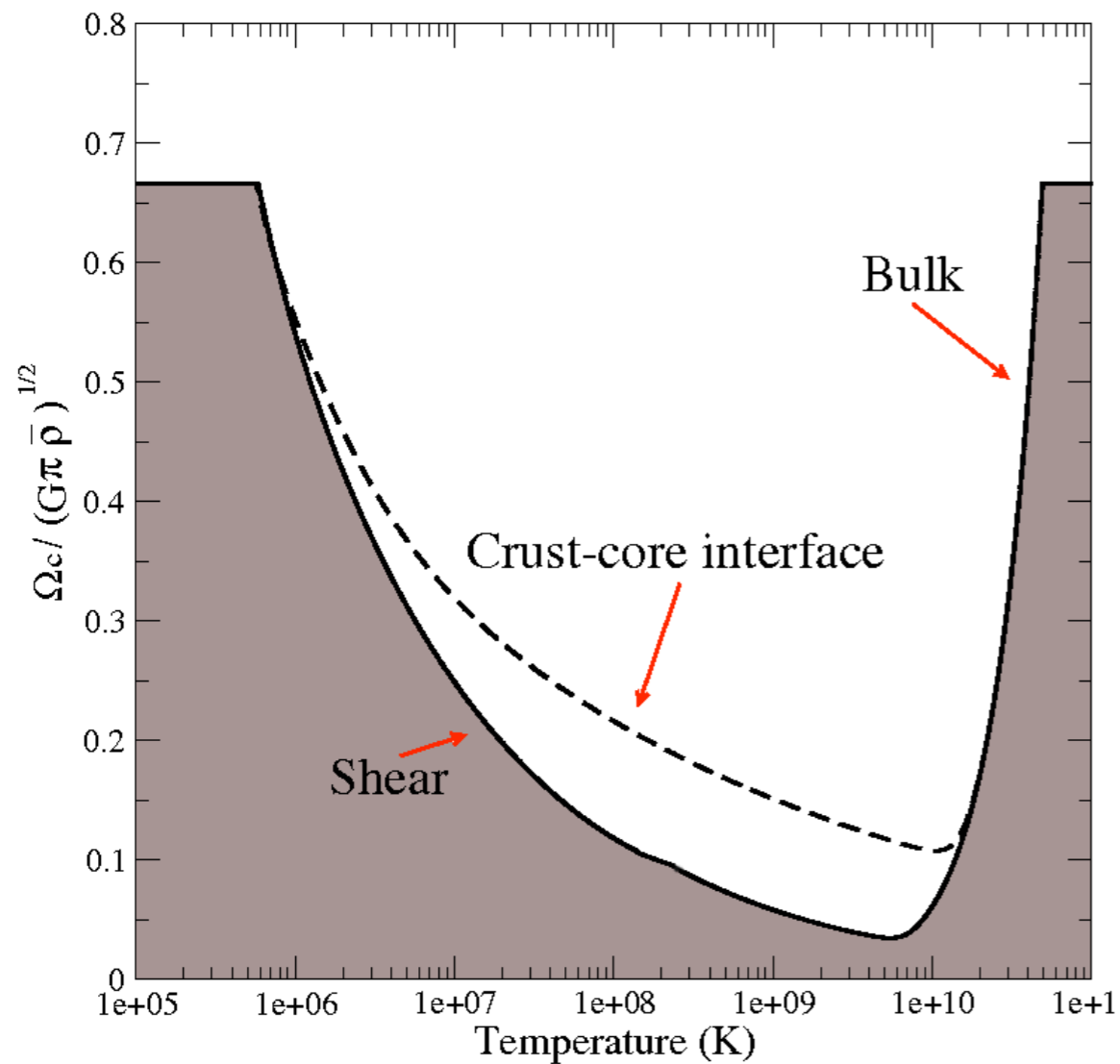
Inertial observer

- r-mode generically unstable to GW emission
- Emission at  $\omega \approx \frac{4}{3}\Omega$
- Viscosity damps the mode except in a window of temperatures and frequencies

# r-mode instability window - I

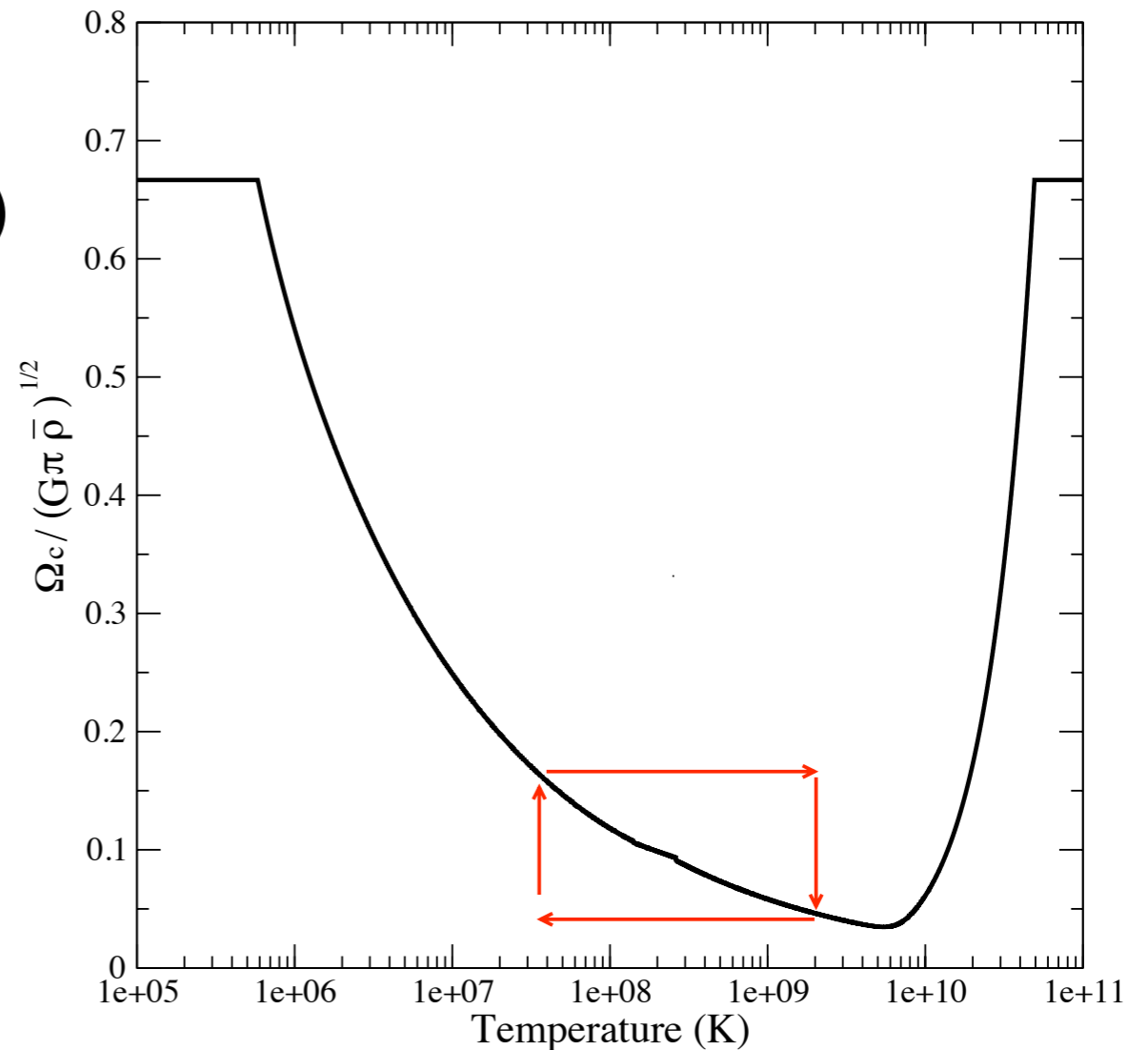


# r-mode instability window - I

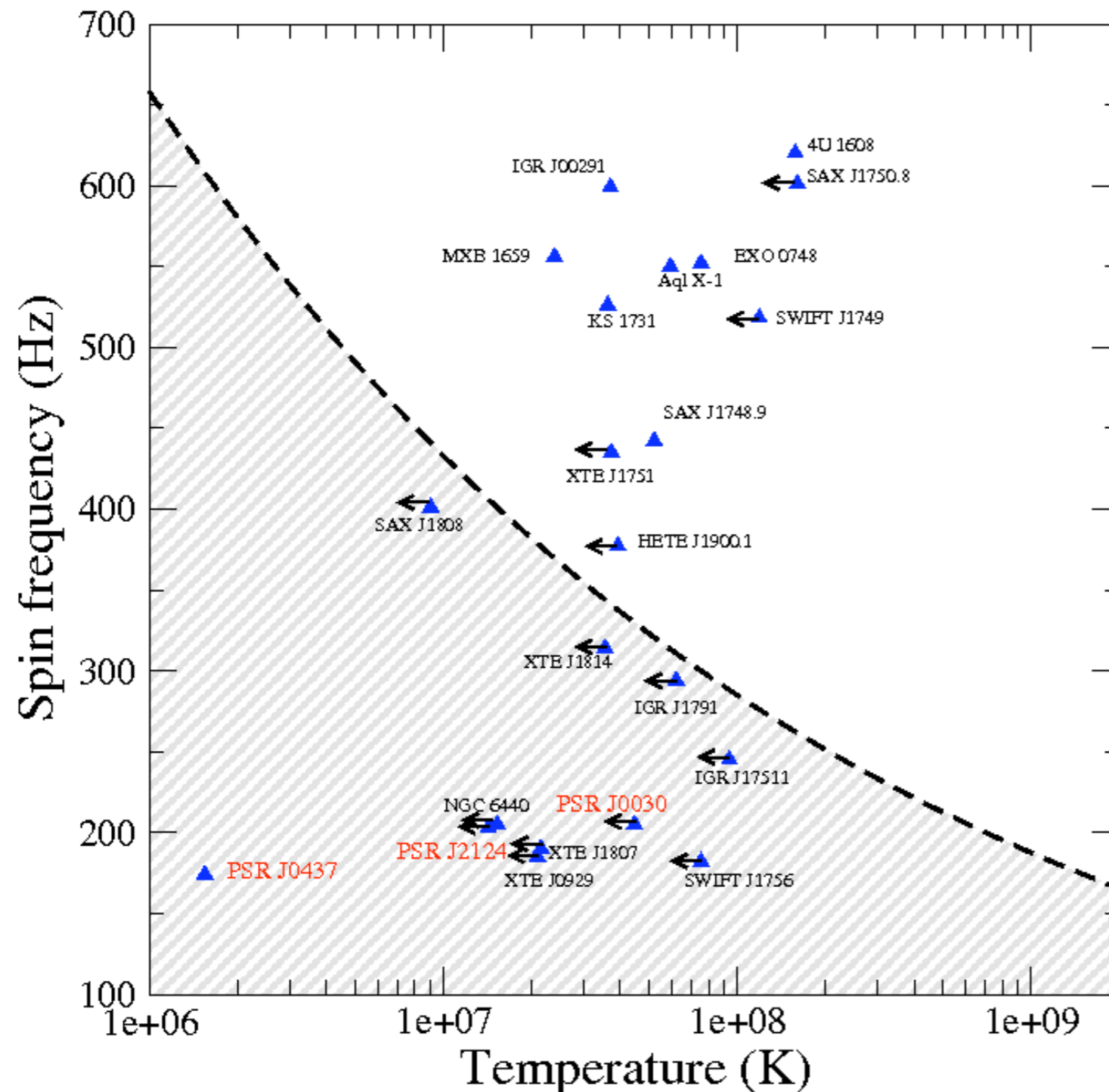


## r-mode instability window - I

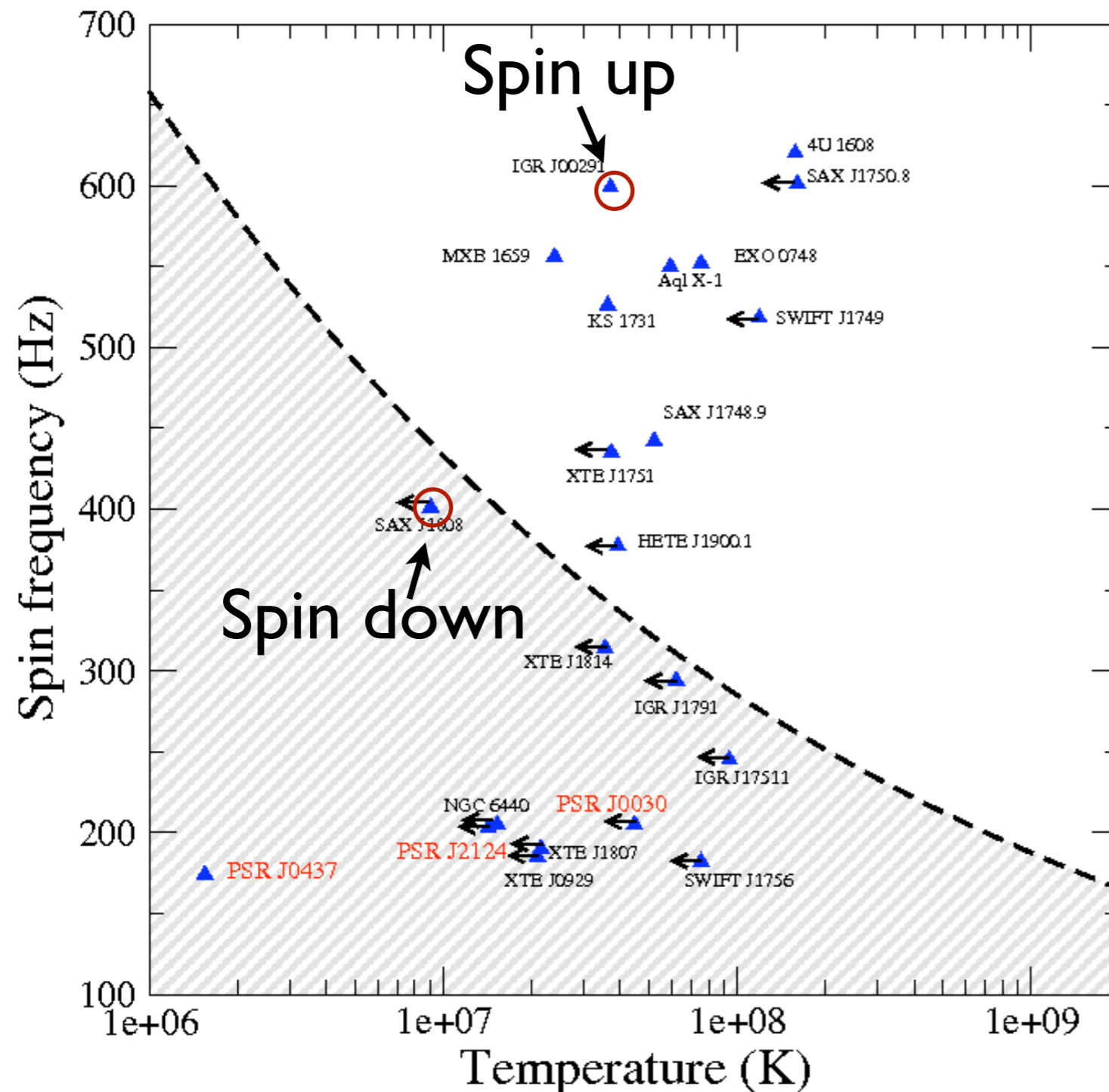
- Duty cycle short (10% or less)
- Effects of EOS? (Hyperons..)
- Effects of superfluidity?



## r-mode instability window - II



# r-mode instability window - II



## Multifluid hydrodynamics

$$\partial_t \rho_x + \nabla_i (\rho_x v_x^i) = 0$$

$$(\partial_t + v_x^j \nabla_j) (v_i^x + \varepsilon_x w_i^{yx}) + \nabla_i (\tilde{\mu}_x + \Phi) + \varepsilon_x w_{yx}^j \nabla_i v_j^x = f_i^x / \rho_x + \nabla_j D_i^j$$

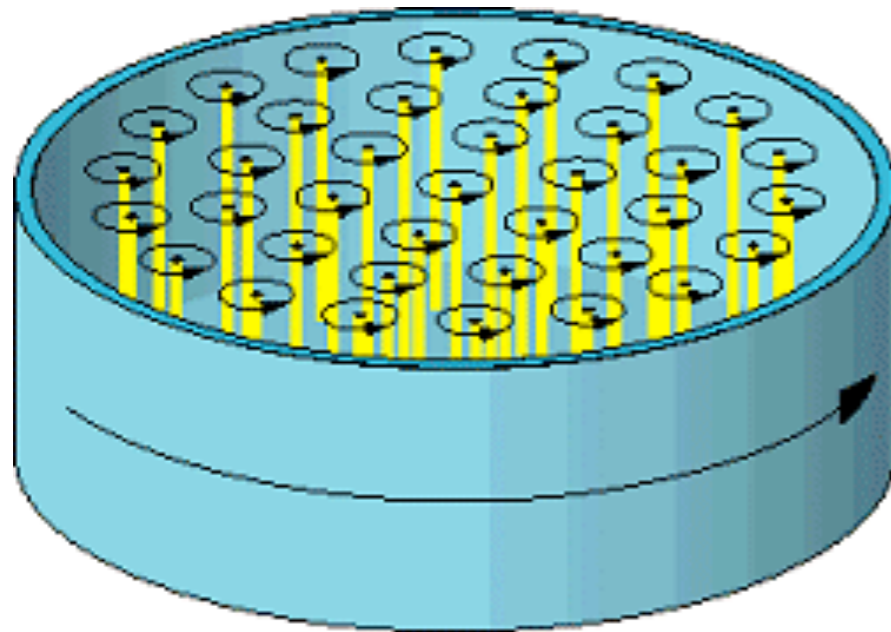
$D_i^j$

**Dissipative terms (bulk viscosity, shear viscosity, etc..)**

$$f_i^x = 2\rho_n \mathcal{B}' \epsilon_{ijk} \Omega^j w_{xy}^k + 2\rho_n \mathcal{B} \epsilon_{ijk} \hat{\Omega}^j \epsilon^{klm} \Omega_l w_m^{xy}$$

**Mutual Friction**

## Mutual friction

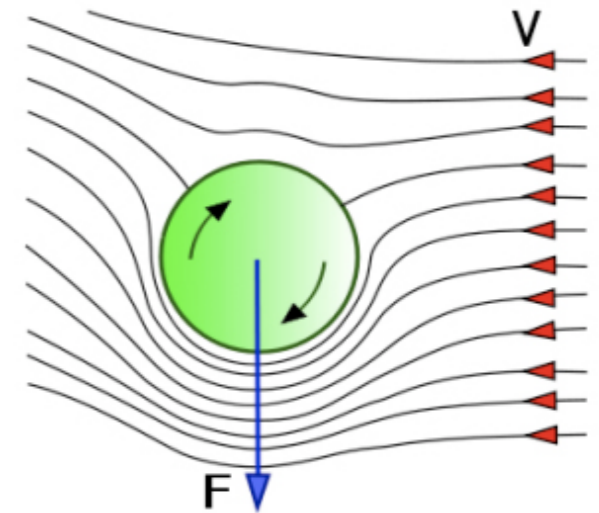


- Superfluids rotate by forming quantised vortices
- Vortex density determines spin :  
vortices must move out to spin down the fluid!
- Vortices could be strongly pinned in the crust

### Magnus Force

PINNED :  $\epsilon^{ijk} \hat{k}_j (v_k^v - v_k^n) + F_p^i = 0$

FREE :  $\epsilon^{ijk} \hat{k}_j (v_k^v - v_k^n) + \mathcal{R}(v_c^i - v_v^i) = 0$

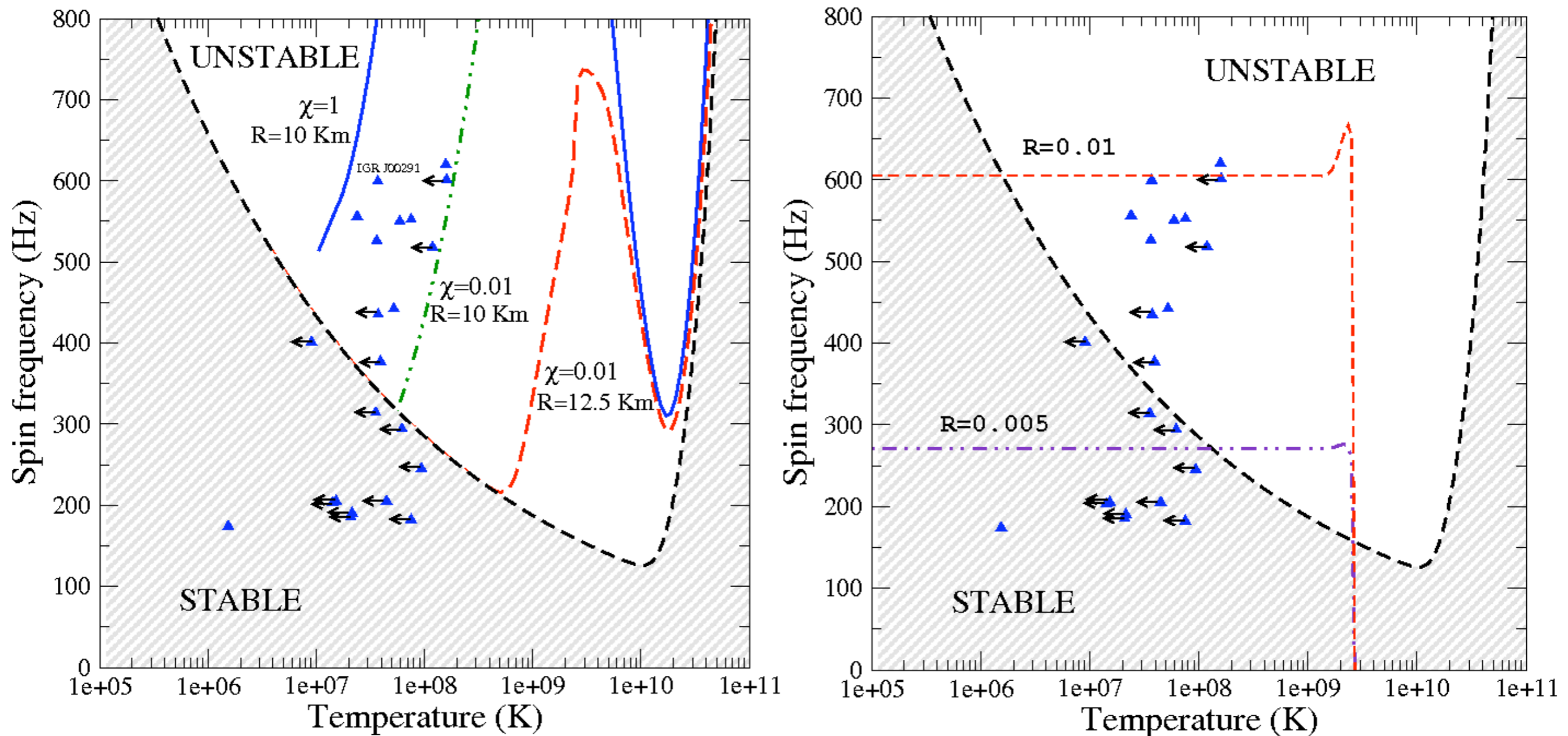






# r-mode instability window - III

# r-mode instability window - III



## Conclusions

- We have a problem with the r-mode. Can a system be inside the instability window?
- Need to model the accretion torque.
- Are some AMXPs emitting GWs? Precise timing and cooling measurements can guide us..(while we wait for GW measurements...)
- Important input for choosing targets for GW observations: persistent sources best targets?