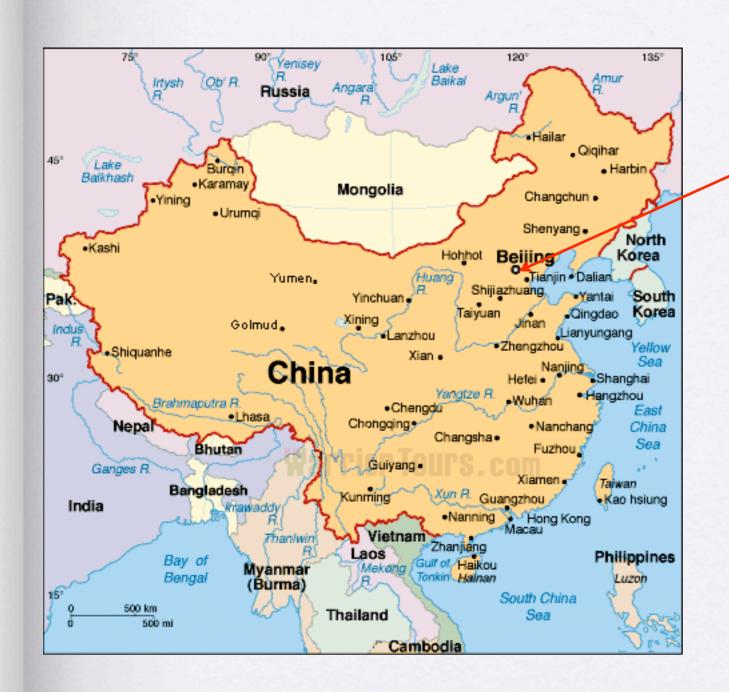
# **Spatially Modulated Interaction and Dipolar Interaction Induced Resonances**

Hui Zhai Institute for Advanced Study Tsinghua University Beijing, China



INT Seattle Fermion Workshop 2011





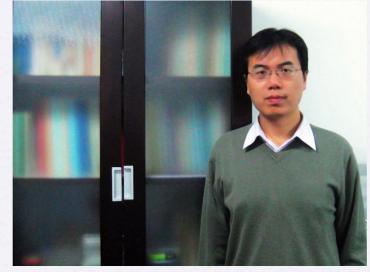
## **Three major recent topics:**



- 1. Quantum gases in synthetic gauge field (with spin-orbit coupling)
- 2. New strongly interacting systems in quantum gases
- 3. Magnetism in quantum gases

# **Main Group Members**





Zeng-Qiang Yu



**Xiaoling Cui** 



Ran Qi

**Chao Gao** 

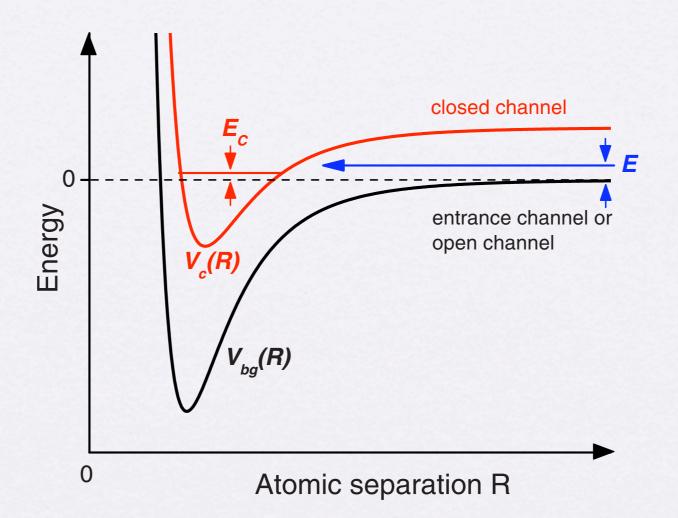






**Chao-Ming Jian** 

# Feshbach resonance is an important tool to achieve strong interactions in ultracold Fermi gases



magnetic Feshbach resonance; optical Feshbach resonance; confinement induced resonance New Way to achieve scattering resonance in New cold atom systems with Novel features

A: Spatially modulated interaction induced resonance Alkali-earth-(like) atomic gases

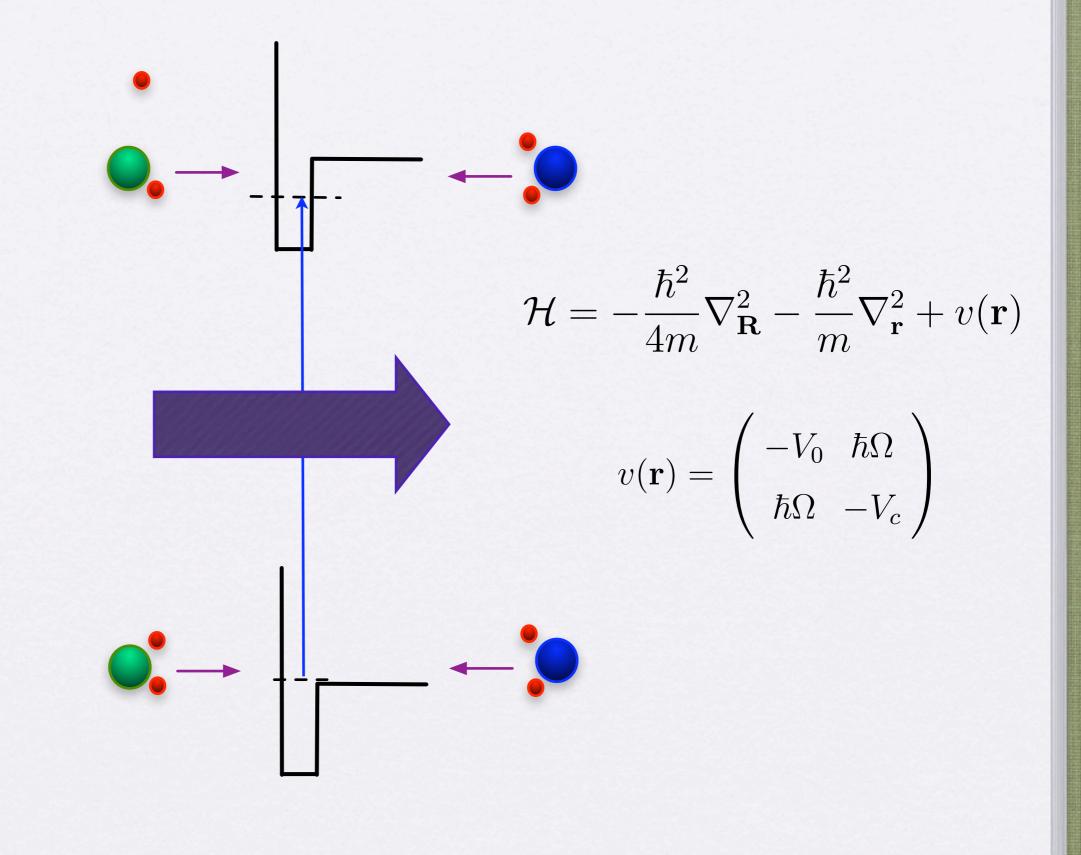
**B:** Strong dipolar interaction induced resonace

**Polar molecular gases** 

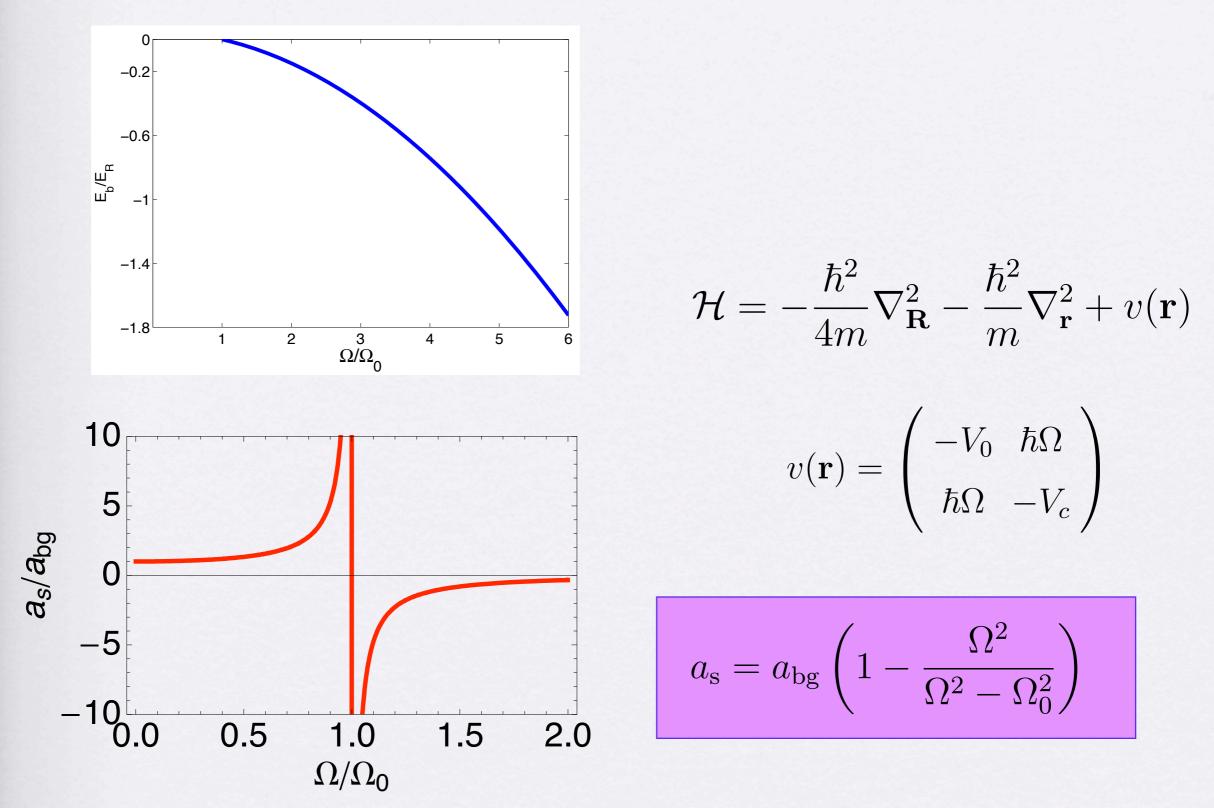
New Way to achieve scattering resonance in New systems with New features

A: Spatially modulated interaction induced resonance Alkali-earth-(like) atomic gases

# **Idea of Optical Feshbach resonance**

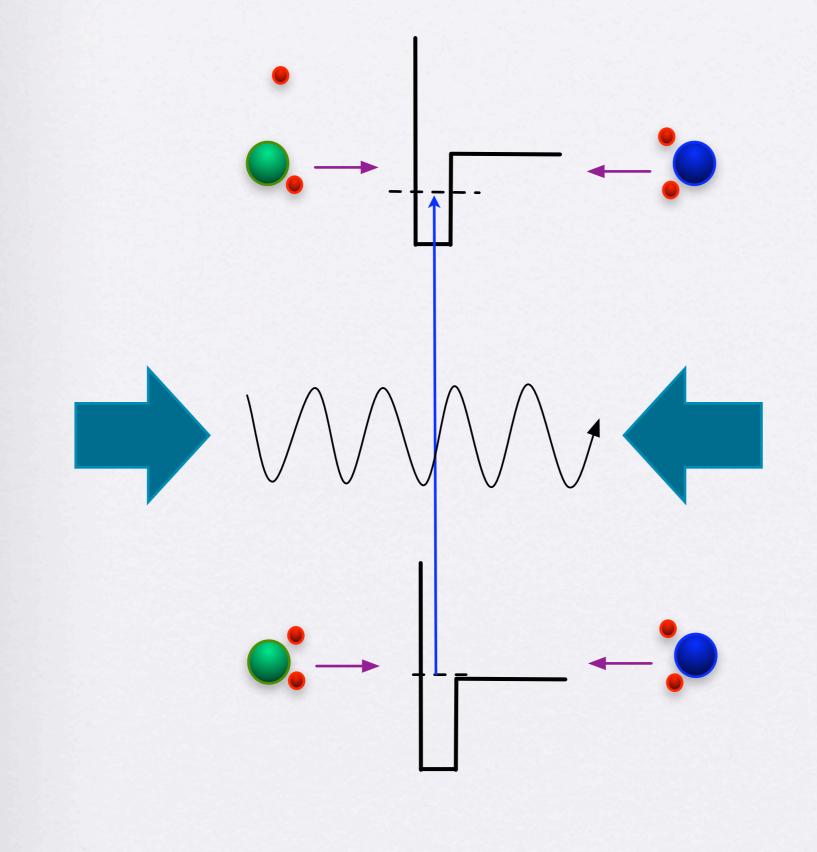


**Optical Feshbach resonance** 



C. Chin, R. Grimm, P. Julienne, and E. Tiesinga, Rev. Mod. Phys. 82, 1225 (2010).

# **Optical Feshbach resonance with <b>Standing** wave

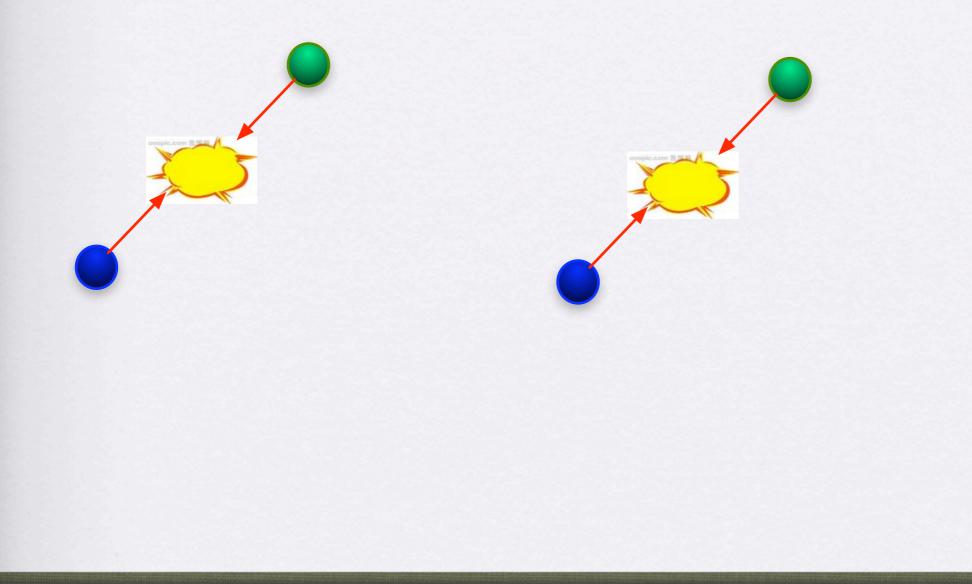


# **Spatial dependent interaction**

**Two-body interaction potential:** 

$$V(\mathbf{r_1}, \mathbf{r_2}) = V(\mathbf{r_1} - \mathbf{r_2})$$

**Spatial independent** 



#### **Spatial dependent interaction**

**Two-body interaction potential:** 

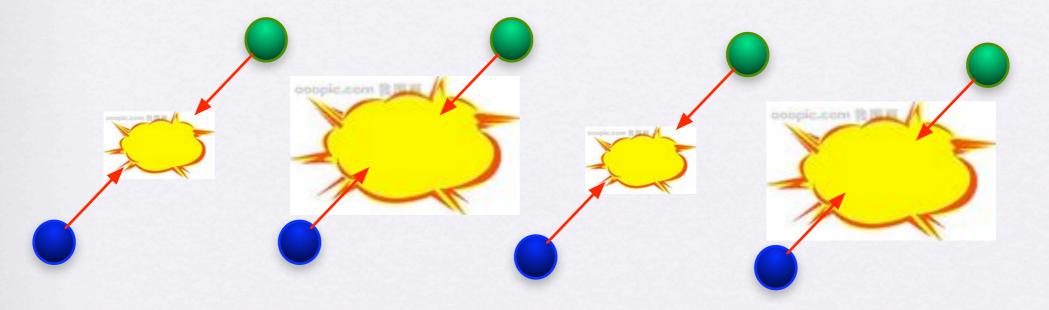
$$V(\mathbf{r_1}, \mathbf{r_2}) = V\left(\mathbf{r_1} - \mathbf{r_2}, \frac{\mathbf{r_1} + \mathbf{r_2}}{2}\right)$$

Spatially periodically modulated

 $\mathcal{H} = -\frac{\hbar^2}{4m}\nabla_{\mathbf{R}}^2 - \frac{\hbar^2}{m}\nabla_{\mathbf{r}}^2 + v(\mathbf{r}, \mathbf{R})$ 

$$v(\mathbf{r}) = \begin{pmatrix} -V_0 & \hbar\Omega(\mathbf{R}) \\ \hbar\Omega(\mathbf{R}) & -V_c \end{pmatrix}$$

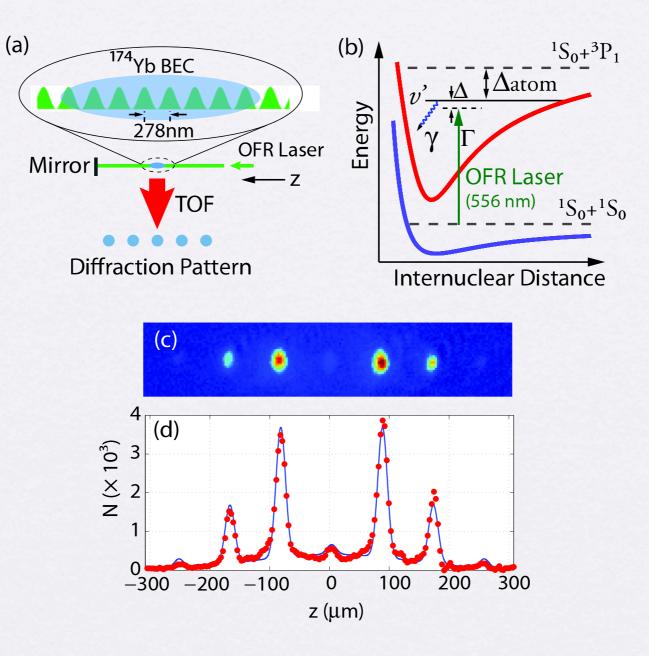
 $\Omega(\mathbf{R}) = \Omega \cos(Kx)$ 



 $a_{s}(x)$  is spatially dependent and modulates periodically in space

#### **Experimental Realization**

Submicron spatial modulation of an interatomic interaction in a Bose-Einstein condensate, PRL, 105, 050405 (2010) Kyoto group



How  $a_s(x)$  modulates in space?

$$a_{\rm s} = a_{\rm bg} \left( 1 - \frac{\Omega^2}{\Omega^2 - \Omega_0^2} \right) \qquad \Omega(\mathbf{R}) = \Omega \cos(Kx)$$

$$a_{\rm s}(x) = a_{\rm bg} \left( 1 - \frac{\Omega^2 \cos^2(Kx)}{\Omega^2 \cos^2(Kx) - \Omega_0^2} \right)$$

Any other physics effects?

# What we have done: Solve two-body problem of this Hamiltonian

0

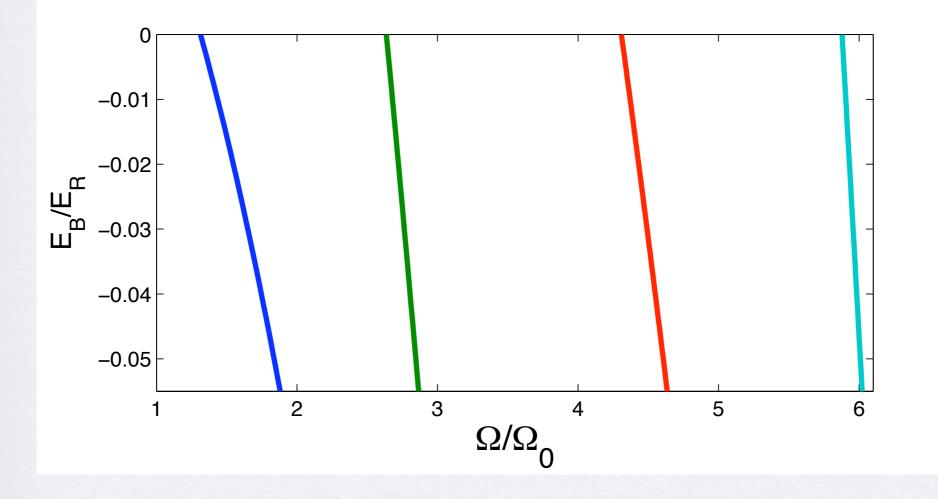
$$\mathcal{H} = -\frac{\hbar^2}{4m} \nabla_{\mathbf{R}}^2 - \frac{\hbar^2}{m} \nabla_{\mathbf{r}}^2 + v(\mathbf{r})$$
$$v(\mathbf{r}) = \begin{pmatrix} -V_0 & \hbar\Omega(\mathbf{R}) \\ \hbar\Omega(\mathbf{R}) & -V_c \end{pmatrix}$$

0

$$\Omega(\mathbf{R}) = \Omega \cos(Kx)$$

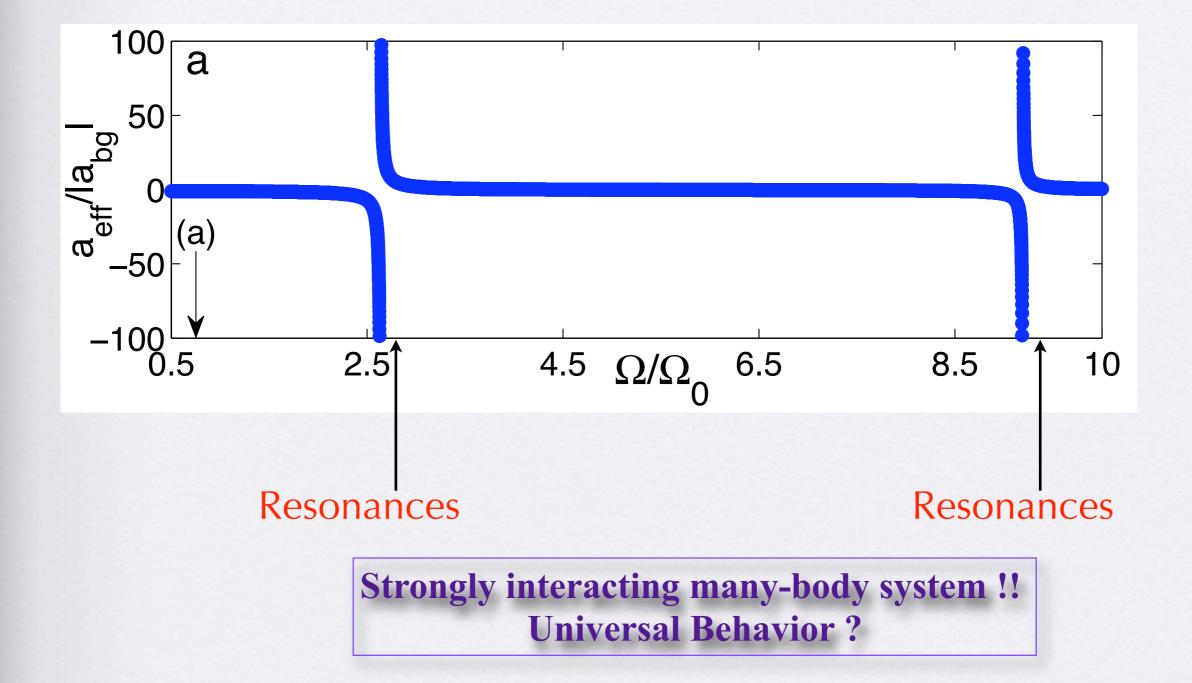
Qi Ran and HZ, arXiv: 1101.4464

#### **Results I: Bound States**



**Results II: Scattering Resonances** 

$$a_{\text{eff}} = \lim_{k \to 0} \frac{\tan \delta(k)}{k}$$



#### **Results III: Local Scattering Length** --- related to local interaction energy

**Bethe-Peierls condition:** 

$$\lim_{r \to 0} \psi(r, x) = \frac{1}{r} - \frac{1}{a_{loc}(x)}$$

Local scattering length

$$a_{\rm loc}(x) = -\lim_{r \to r_0} \frac{r\psi_{\rm o}(x,r)}{\partial_r(r\psi_{\rm o}(x,r))}$$

The mean-field energy for a BEC:

$$\mathcal{E} = \int dx \left[ -\frac{\hbar^2}{2m} \varphi^* \nabla^2 \varphi + \frac{4\pi\hbar^2}{m} a_{loc}(x) n^2(x) \right]$$

#### **Results III: Local Scattering Length**

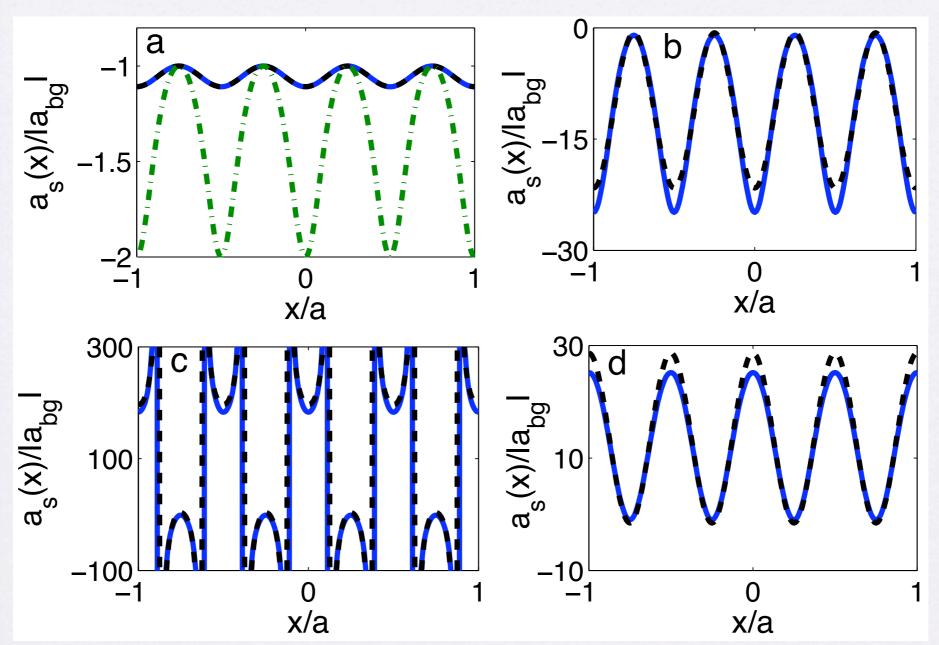
Exact formula:

$$a_{\rm loc}(x) = \frac{1 - \sum_{m \neq 0} U_m \cos(mKx) / U_0}{a_{\rm eff}^{-1} - \sum_{m \neq 0} U_m |m| K \cos(mKx) / (2U_0)}$$

$$Ka_{\text{eff}} \ll 1$$
  $a_{\text{loc}}(x) = a_{\text{eff}} \left[ 1 - \frac{2U_2}{U_0} \cos(2Kx) \right]$   
 $Ka_{\text{eff}} \gg 1$   $a_{\text{loc}}(x) = \frac{1}{K} \left[ 1 - \frac{U_0}{2U_2 \cos(2Kx)} \right]$ 

Simplified formula

#### **Results III: Local Scattering Length**



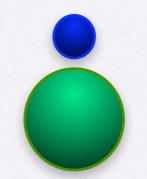
## Summary: Take Home Message

New Mechanism	New System	New Features
Two-body interaction	Alkali-earth-(like)	Spatially dependent
potential has center-	atomic gases: Sr,	local scattering
of-mass dependence	Ca, Yb	length

New Way to achieve scattering resonance in New systems with New features

**B:** Strong dipolar interaction induced resonace

**Polar molecular gases** 



large dipole moment: d

#### Hard to cool it directly !!

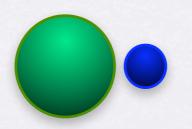




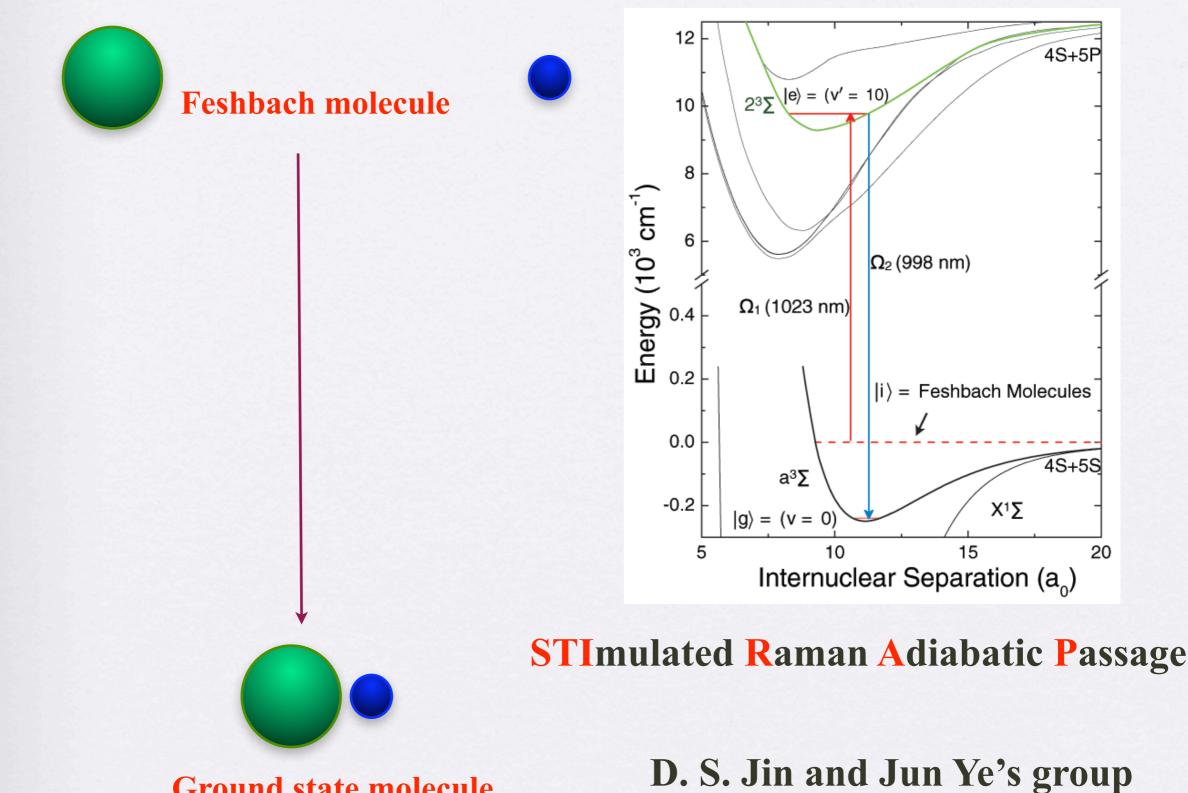
#### **Difficulties:**

1. Large energy detuning: 10-100 THz

2. Small transition matrix element



**Ground state molecule** 



**Ground state molecule** 

#### KRb+KRb --> K2+Rb2

**Chemically unstable !!** 

	Na	K	Rb	Cs
Li Na	-328(2)	-533.9(3) 74.3(3)	-618(200) 45.5(5)	-415.38(2) 236.75(20)
K Rb			-8.7(9)	37.81(13) 29.1(1.5)

#### Zuchowski and Hutson, PRA (2010)

$$E_{\rm D} \sim \frac{D}{\langle r \rangle^3} \sim k_{\rm F}^3 D \qquad E_{\rm F} \sim k_{\rm F}^2$$

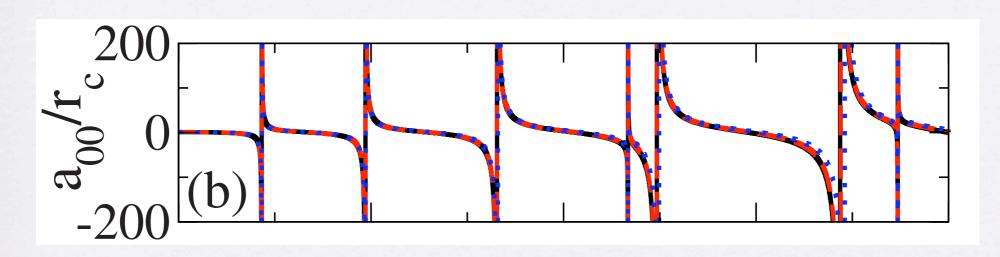
$$k_{\rm F}D \sim 1$$
  
 $E_{\rm D} \sim E_{\rm F}$ 

Dipole moments of vibrational ground state:  $\upsilon$ =0

Mixture	[Debeye]
Li-Na	0.56
Li-K	3.6
Li-Rb	4.2
Li-Cs	5.5
Na-K	2.8 ←
Na-Rb	3.3 🔶
Na-Cs	4.6 ←
K-Rb	0.6
K-Cs	1.9 🔶
Rb-Cs	1.2 ←

M. Aymar and O. Dulieu, J. Chem. Phys., 122, 204302 (2005)

s-wave resonances with strong dipolar interactions

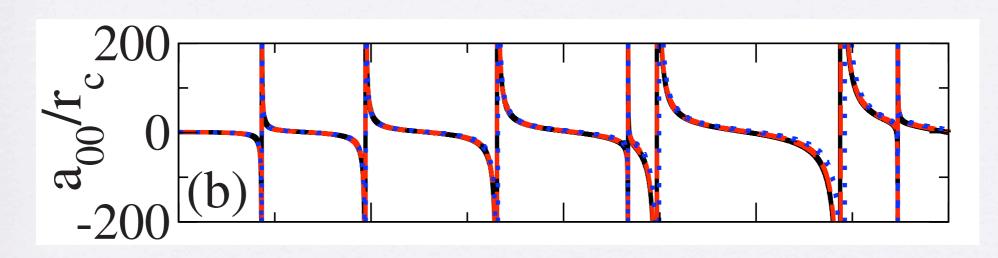


Kanjilal and Blume PRA (2008)

$$V_{\rm D} = \frac{D(1 - 3\cos^2\theta)}{r^3}$$

 $\langle Y_{00}|V_{\rm D}|Y_{00}\rangle = 0$ 

s-wave resonances with strong dipolar interactions

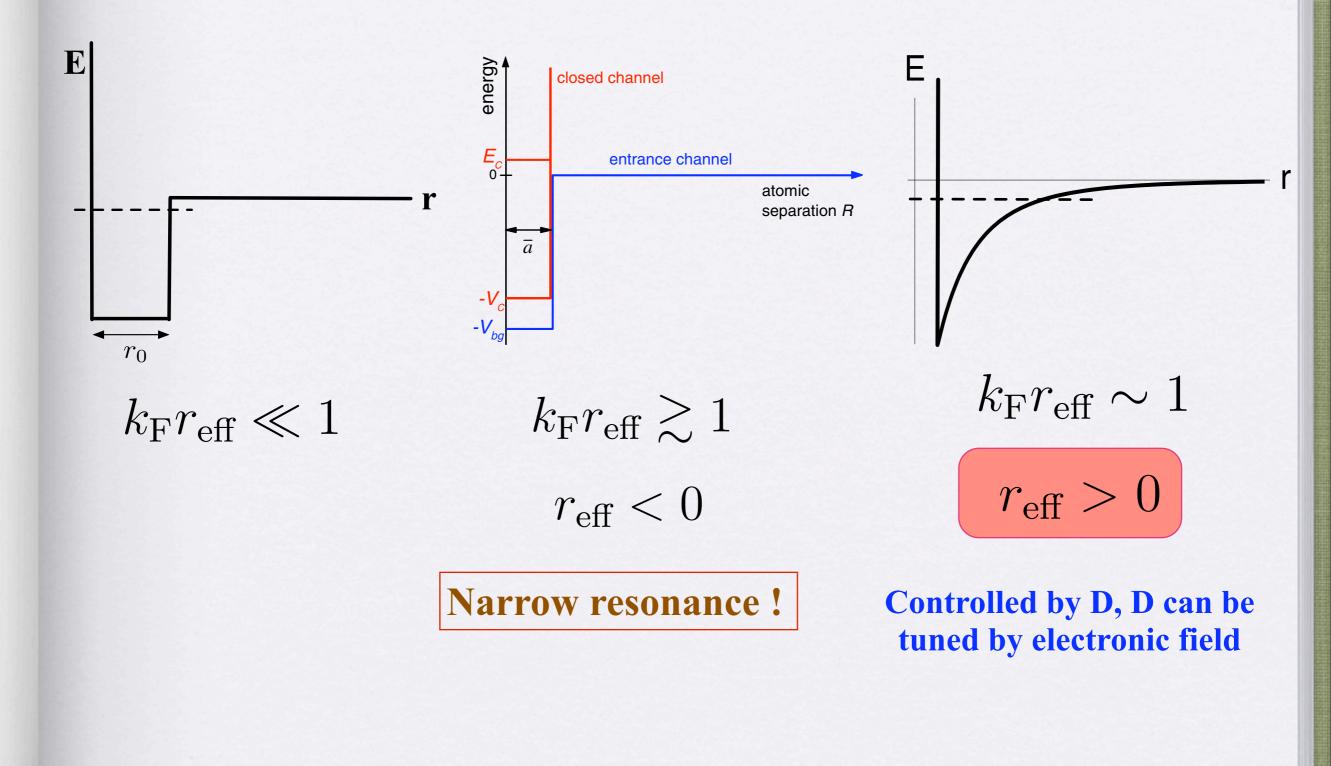


Kanjilal and Blume PRA (2008)

 $\langle Y_{20}|V_{\rm D}|Y_{00}\rangle \neq 0$ 

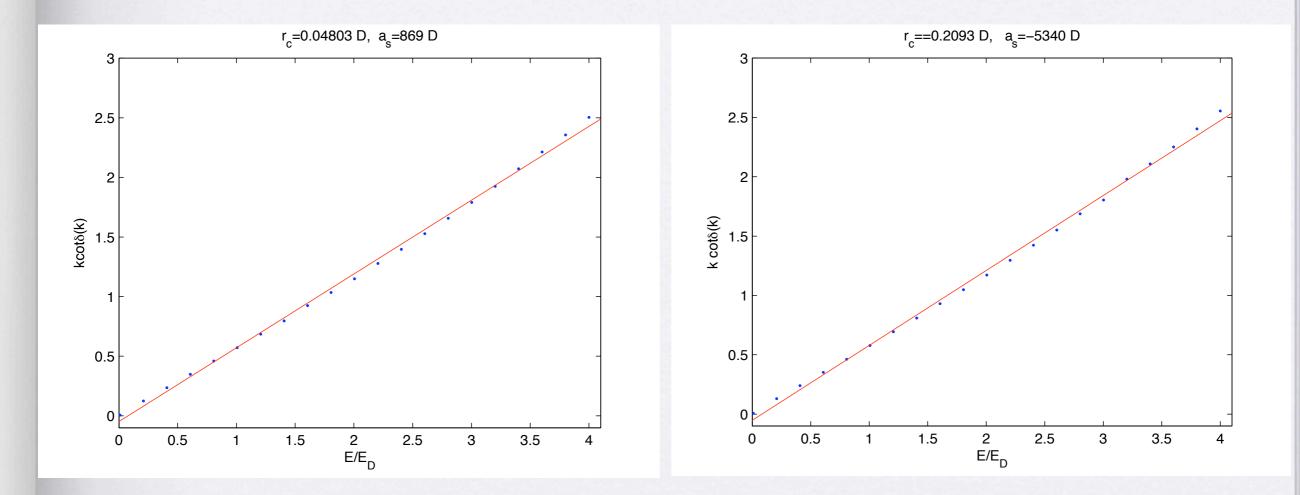
$$V_{\rm eff} = -\frac{\langle Y_{00} | V_{\rm D} | Y_{20} \rangle \langle Y_{20} | V_{\rm D} | Y_{00} \rangle}{l(l+1)/r^2} \propto -\frac{D^2}{r^4}$$

#### How well can we tune a (positive) effective range in cold atoms ? -- from the wiki



# **Positive effective range from a dipole induced resonance**

$$\frac{k}{\tan\delta(k)} = -\frac{1}{a_{\rm s}} + \frac{1}{2}r_{\rm eff}k^2$$

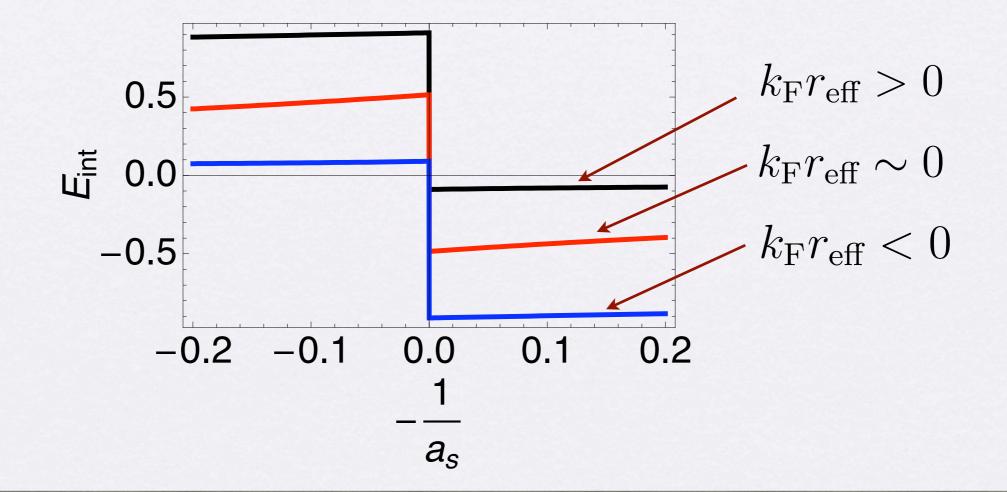


#### What is the physical effect of effective range being positive or negative

**High temperature regime:** 

$$b_2 = \int_0^{+\infty} \frac{dk}{\pi} \frac{d\delta(k)}{dk} e^{-\lambda^2 k^2/(2\pi)}$$

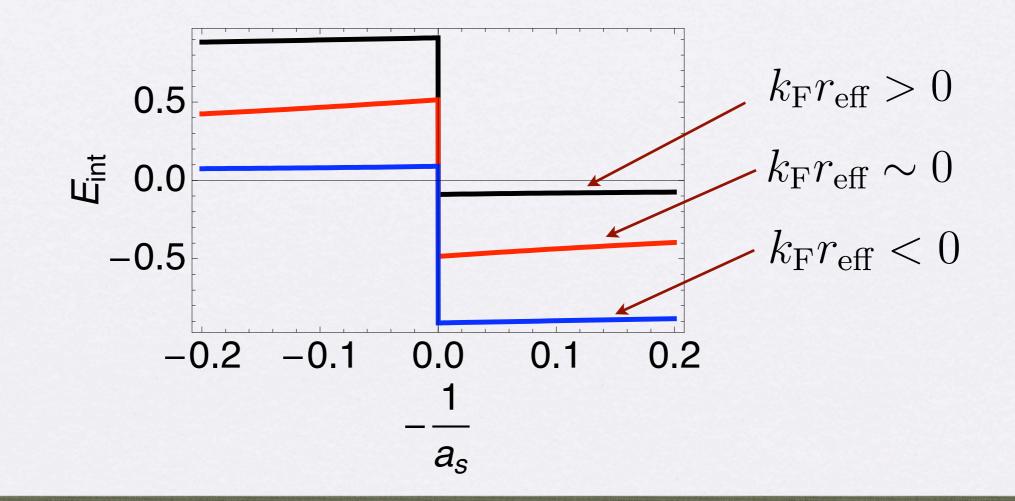
$$E_{\rm int} = \frac{3k_{\rm B}Tn}{2}(n\lambda^3) \left[ -\frac{b_2}{\sqrt{2}} + \frac{\sqrt{2}}{3}T\frac{\partial b_2}{\partial T} \right]$$



What is the physical effect of effective range being positive or negative

$$-\frac{1}{a_{\rm avg}} = -\frac{1}{a_{\rm s}} + \frac{1}{2} r_{\rm eff} \langle k^2 \rangle$$

 $\rightarrow$ 



#### Summary: Take Home Message

New Mechanism	New System	<b>New Features</b>
Two-body interaction potential has center- of-mass dependence	Alkali-earth-(like) atomic gases: Sr, Ca, Yb	Spatially dependent local scattering length
Strong long range dipolar interactions	Polar molecules	positive and sizable effective range

Thank you very much for your attention !