

Renormalization Group Techniques and Applications

Achim Schwenk



CANADA'S NATIONAL LABORATORY FOR PARTICLE AND NUCLEAR PHYSICS

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TECHNISCHE
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INT Weakly-Bound Systems Workshop, March 10, 2010



**NSERC
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NRC-CMRC



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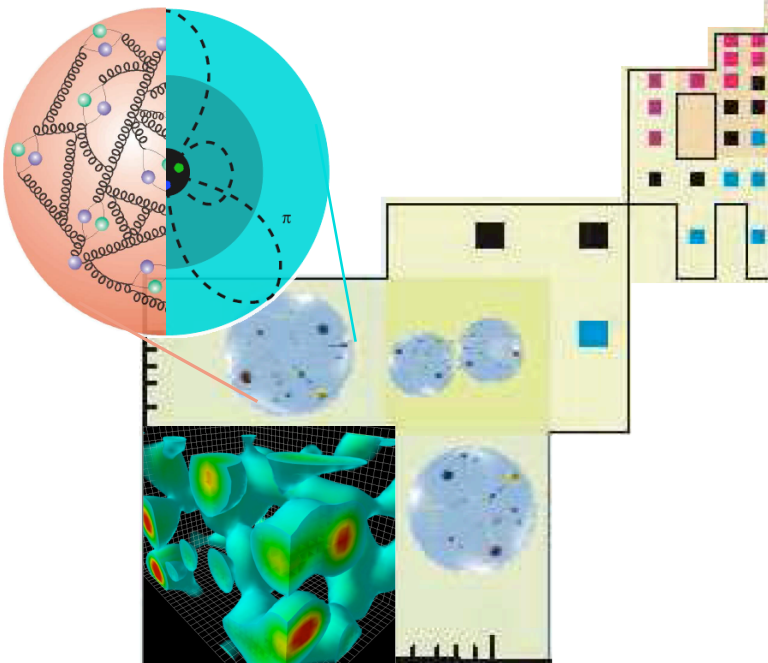
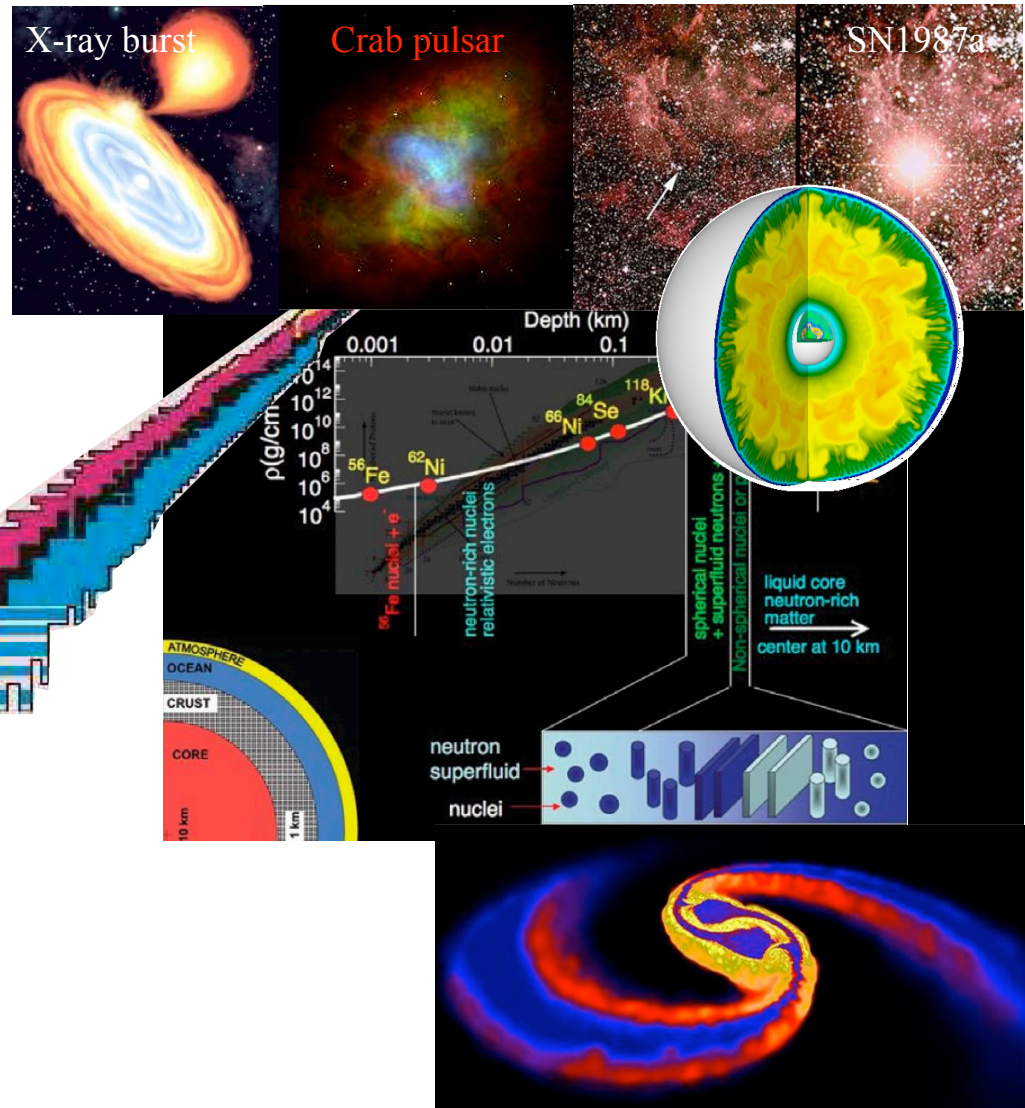
Strong interaction physics in the lab and cosmos

Matter at the extremes:

density $\rho \sim \dots 10^{15} \text{ g/cm}^3$

proton-rich, neutron-rich,
 ${}^8\text{He}$ to $Z/N \sim 0.05$

temperatures $T \sim \dots 100 \text{ MeV}$



Outline

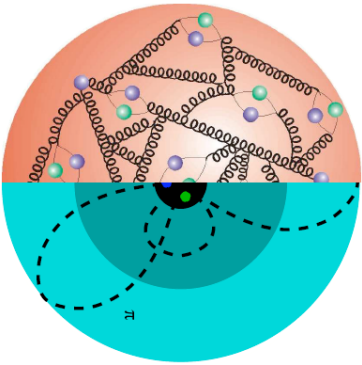
Effective field theory and renormalization group
for nuclear forces

Applications to weakly-bound and neutron-rich nuclei

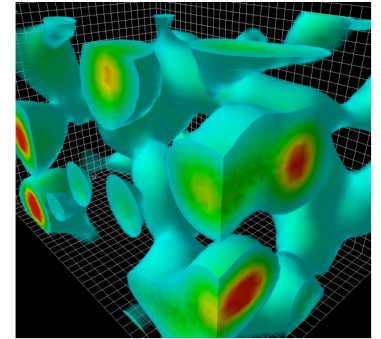
Similarity renormalization group for nuclei

Λ / Resolution dependence of nuclear forces

with high-energy probes:
quarks+gluons



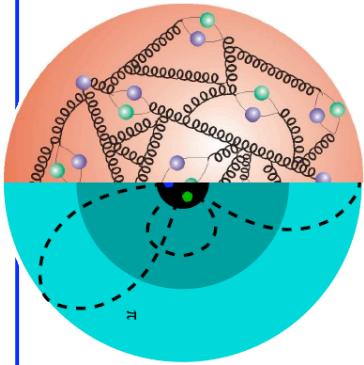
at low energies:
complex QCD vacuum



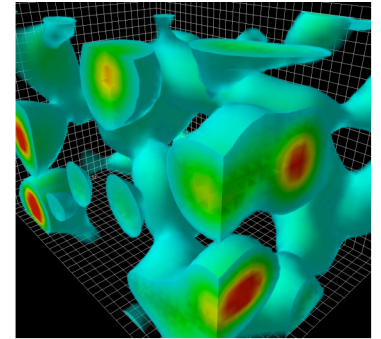
lowest energy excitations:
pions, nearly massless, $m_\pi=140$ MeV
“phonons” of QCD vacuum

Λ / Resolution dependence of nuclear forces

with high-energy probes:
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at low energies:
complex QCD vacuum



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Λ_{chiral}

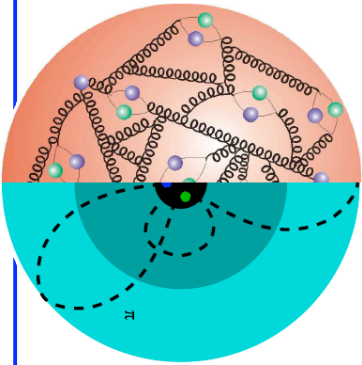
momenta $Q \sim \lambda^{-1} \sim m_\pi$

$\Lambda_{\text{pionless}}$

$Q \ll m_\pi$

Λ / Resolution dependence of nuclear forces

with high-energy probes:
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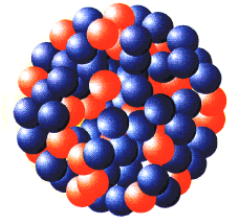
Effective theory for NN, 3N, many-N interactions and electroweak operators: resolution scale/ Λ -dependent

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$

Λ_{chiral}

momenta $Q \sim \lambda^{-1} \sim m_{\pi}$: chiral effective field theory (EFT)

neutrons and protons interacting via pion exchanges
and shorter-range contact interactions



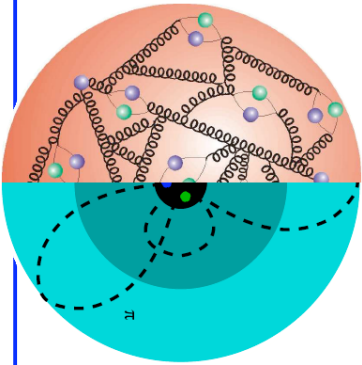
typical momenta in nuclei $\sim m_{\pi}$

$\Lambda_{\text{pionless}}$

$Q \ll m_{\pi}$

Λ / Resolution dependence of nuclear forces

with high-energy probes:
quarks+gluons



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$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$

Λ_{chiral}

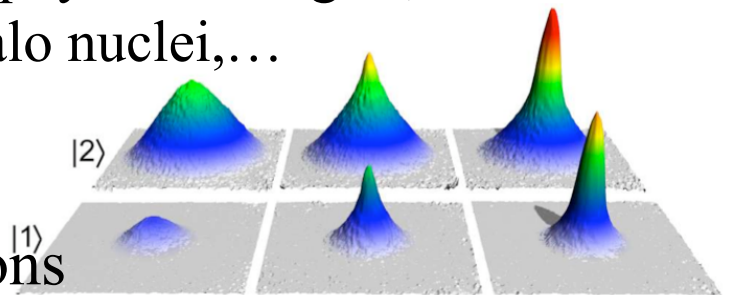
momenta $Q \sim \lambda^{-1} \sim m_{\pi}$

universal properties of neutrons and cold atoms,
reactions at astrophysical energies,
loosely-bound halo nuclei,...

$\Lambda_{\text{pionless}}$

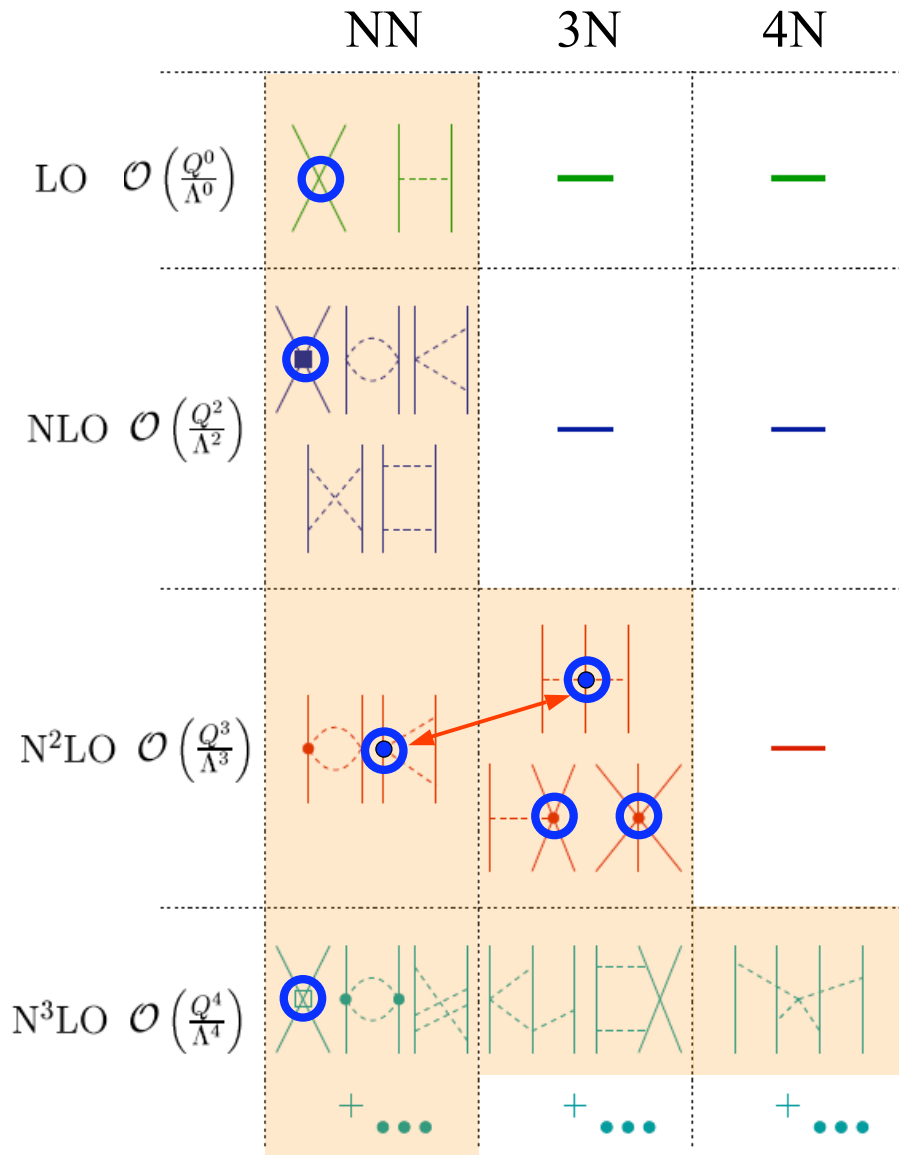
$Q \ll m_{\pi}$: pionless effective field theory

large scattering length physics and corrections



Chiral EFT for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



limited resolution at low energies,
can expand in powers Q/Λ_b

include long-range pion physics

details at short-distance not resolved

capture in few **short-range couplings**,
fit to experiment once, Λ -dependent

systematic: can work to desired
accuracy and obtain error estimates
from truncation order and Λ variation

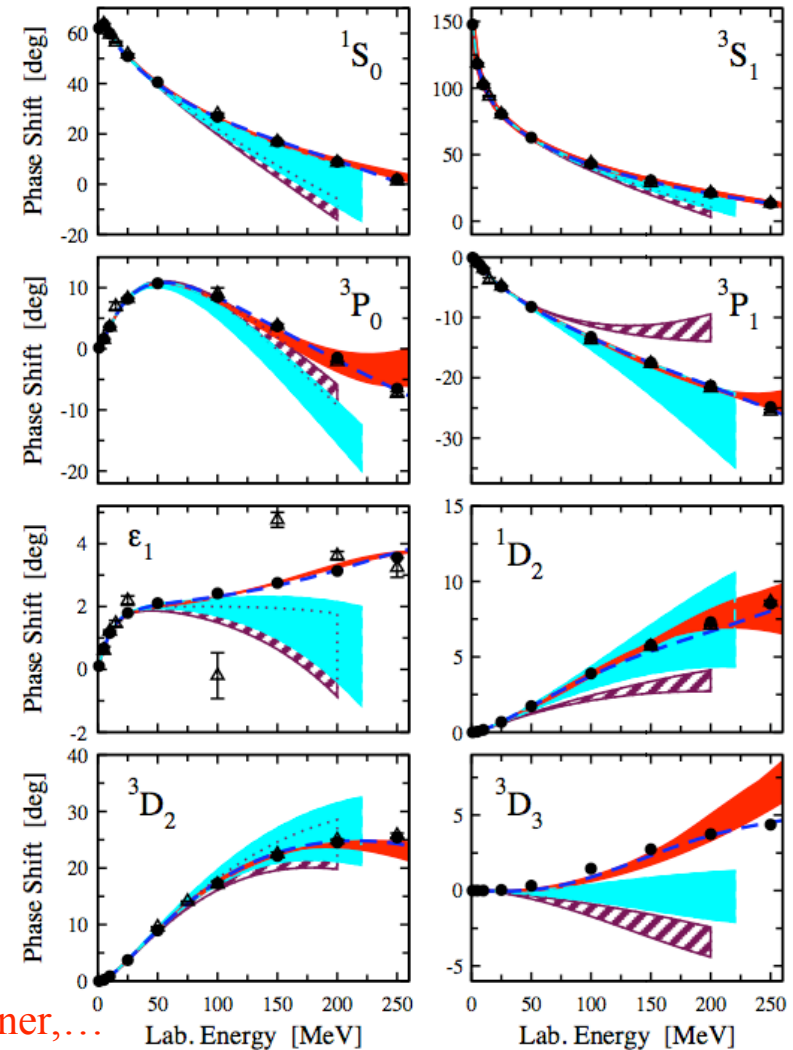
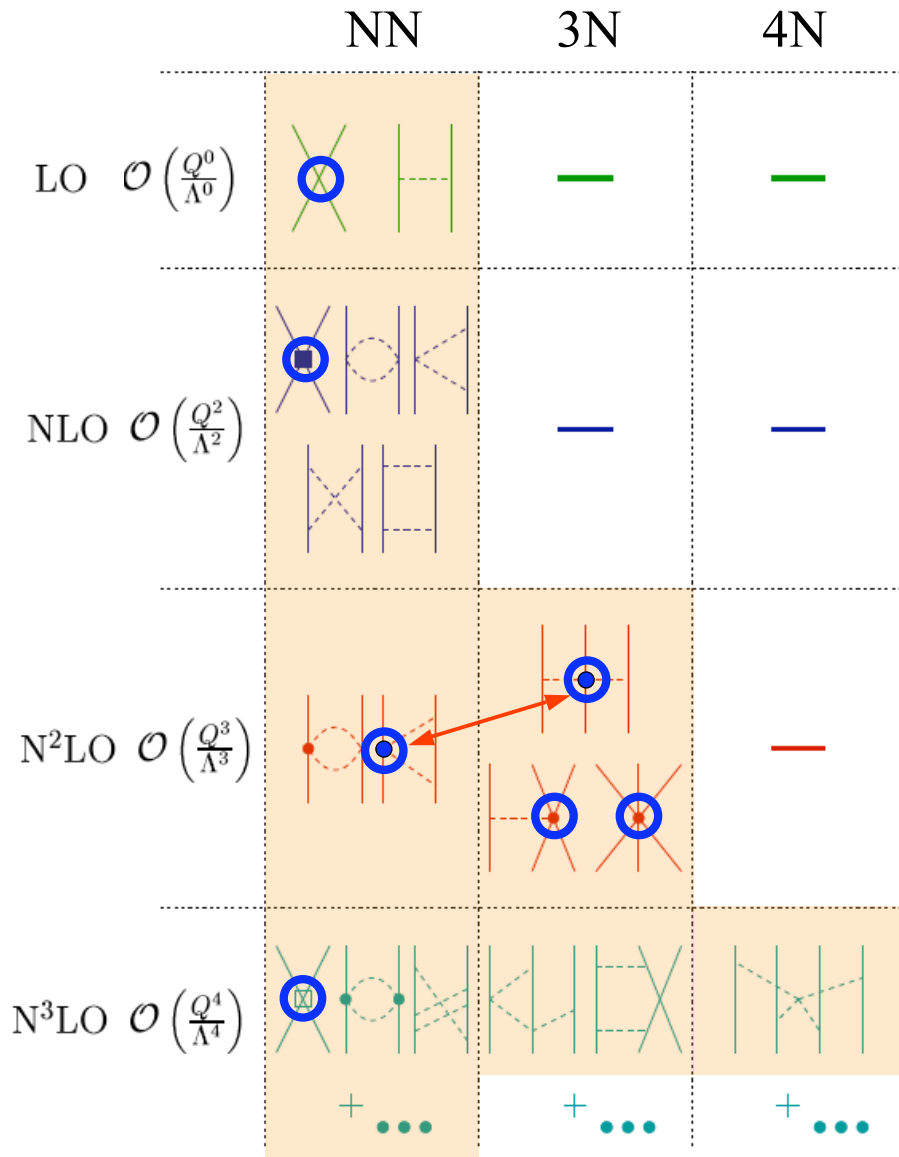
several open problems

Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner,...

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Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

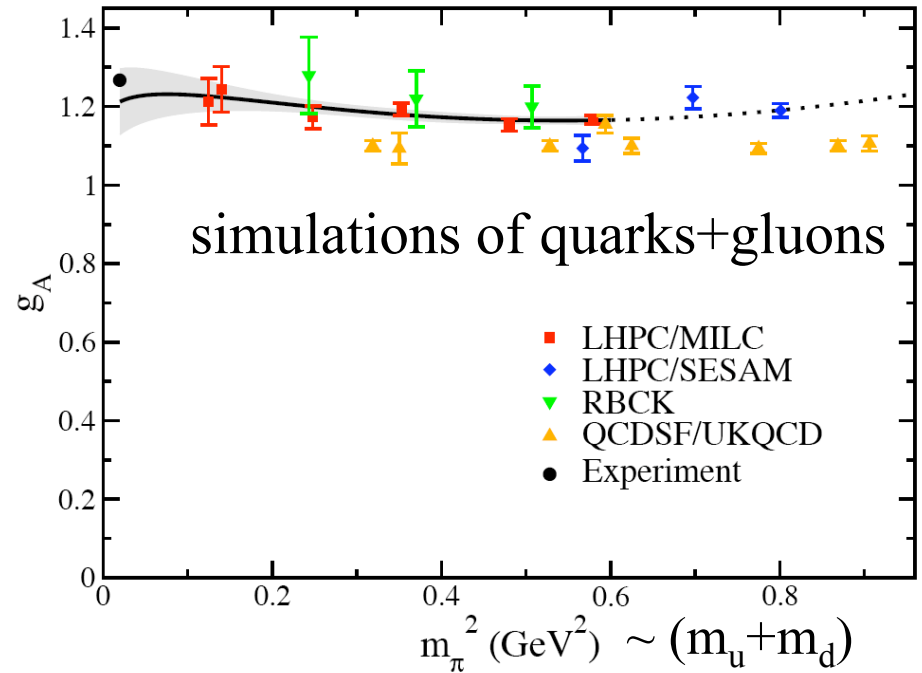
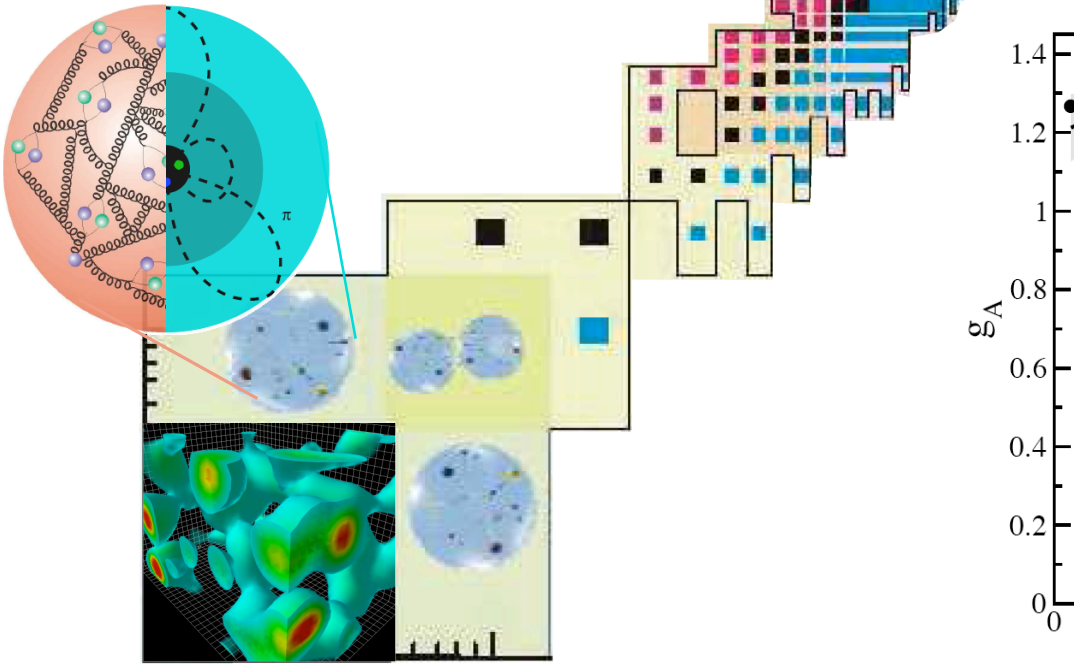
accurate reproduction of low-energy NN scattering at N³LO



Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner,...

Lattice QCD

and nuclear forces



First, exciting efforts to connect nuclear forces to underlying QCD

Long-range couplings: pion-nucleon coupling g_a [Edwards et al. \(2006\)](#)

chiral EFT extrapolation to physical pion mass agrees with experiment

variation of nuclear forces with quark masses [Beane et al. \(2006\)](#)

Future possibility: access/constrain 3-neutron forces (3n exp difficult)

Nuclear forces and the Renormalization Group

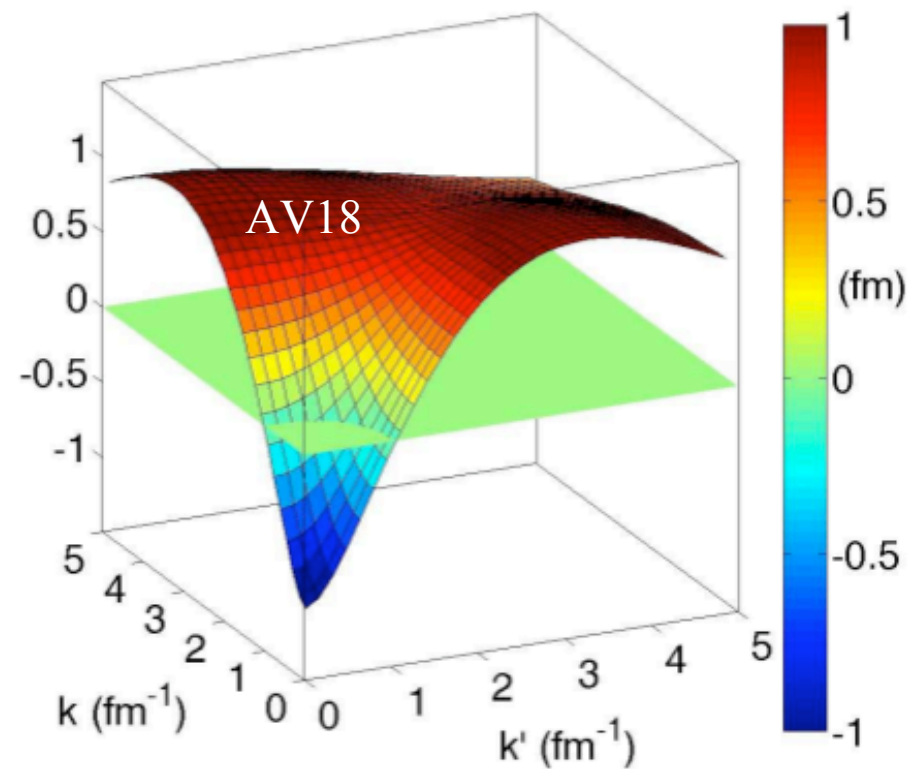
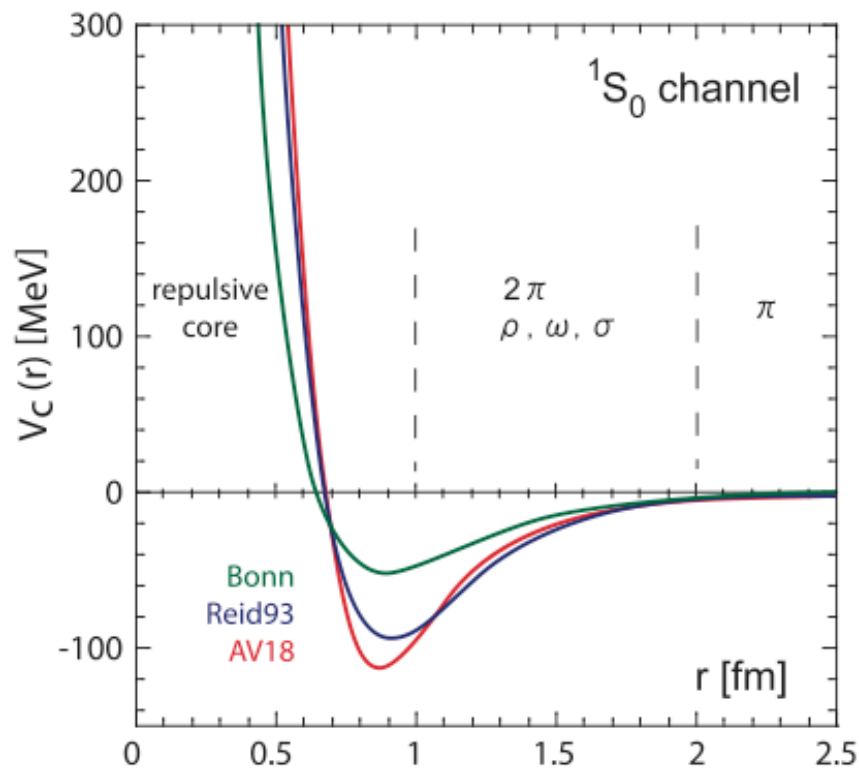
RG evolution to lower resolution/cutoffs

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$

exact RG for NN interactions

Bogner, Kuo, AS, Furnstahl, ...

$$\frac{d}{d\Lambda} V_{\text{low } k}^{\Lambda}(k', k) = \frac{2}{\pi} \frac{V_{\text{low } k}^{\Lambda}(k', \Lambda) T^{\Lambda}(\Lambda, k; \Lambda^2)}{1 - (k/\Lambda)^2}$$



red = short-range repulsion

Nuclear forces and the Renormalization Group

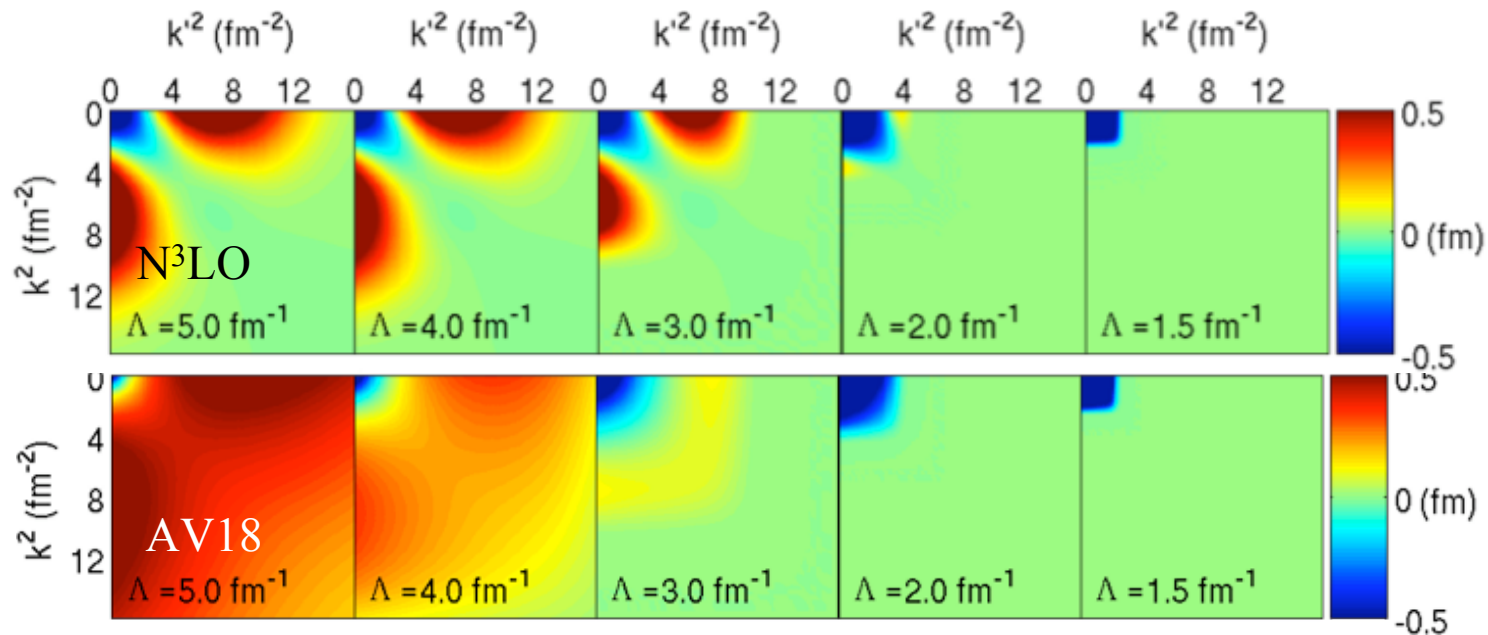
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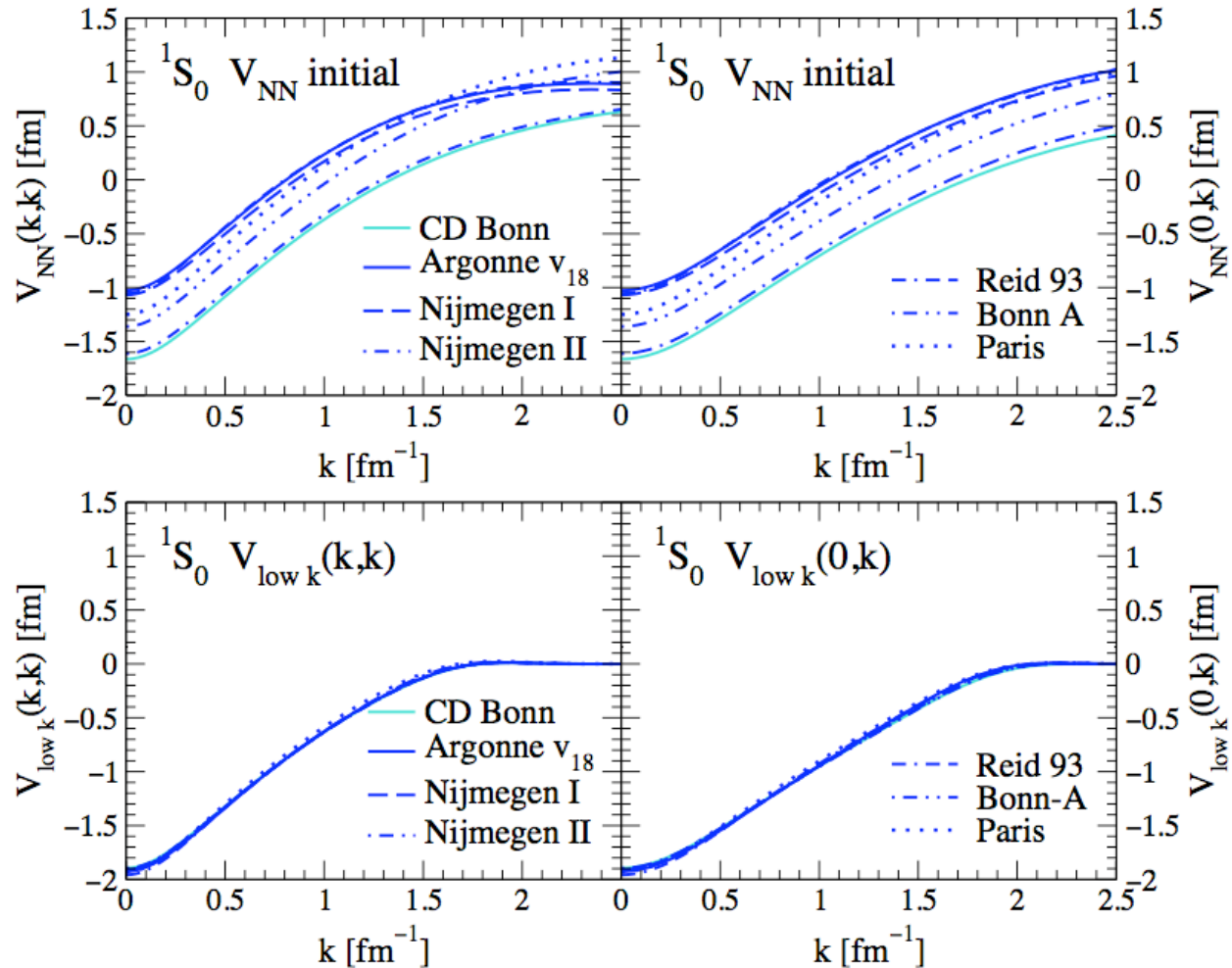


low-momentum interactions $V_{\text{low } k}(\Lambda)$ with sharp or smooth regulators

decouples low-momentum physics from high momenta

red = short-range repulsion and short-range tensor parts

Low-momentum universality

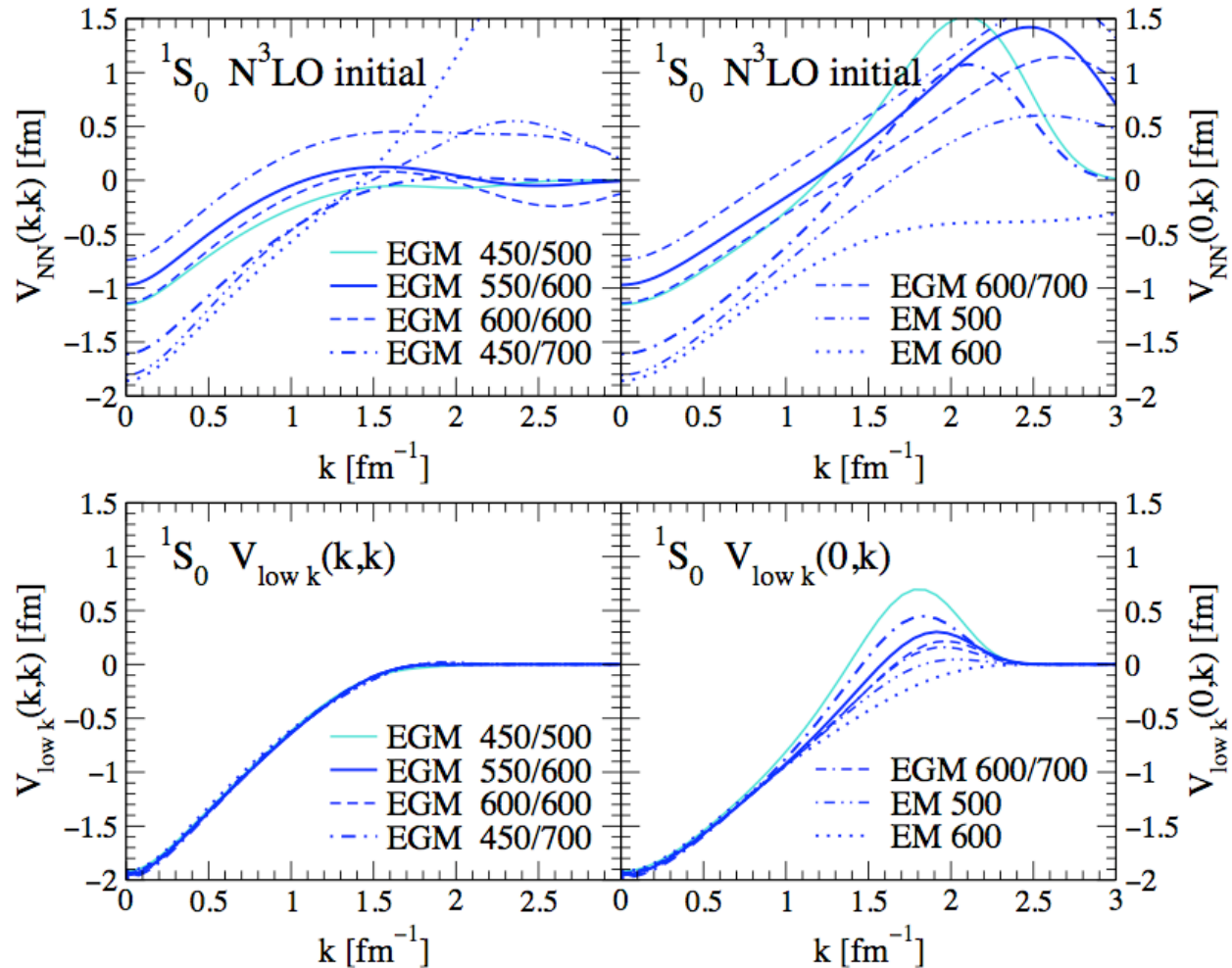


\approx universality from different phenomenological potentials

RG preserves NN observables and long-range parts

decouples low-momentum physics from high momenta

Low-momentum universality



\approx **universality** from different chiral $N^3\text{LO}$ potentials

RG preserves NN observables and long-range parts

decouples low-momentum physics from high momenta

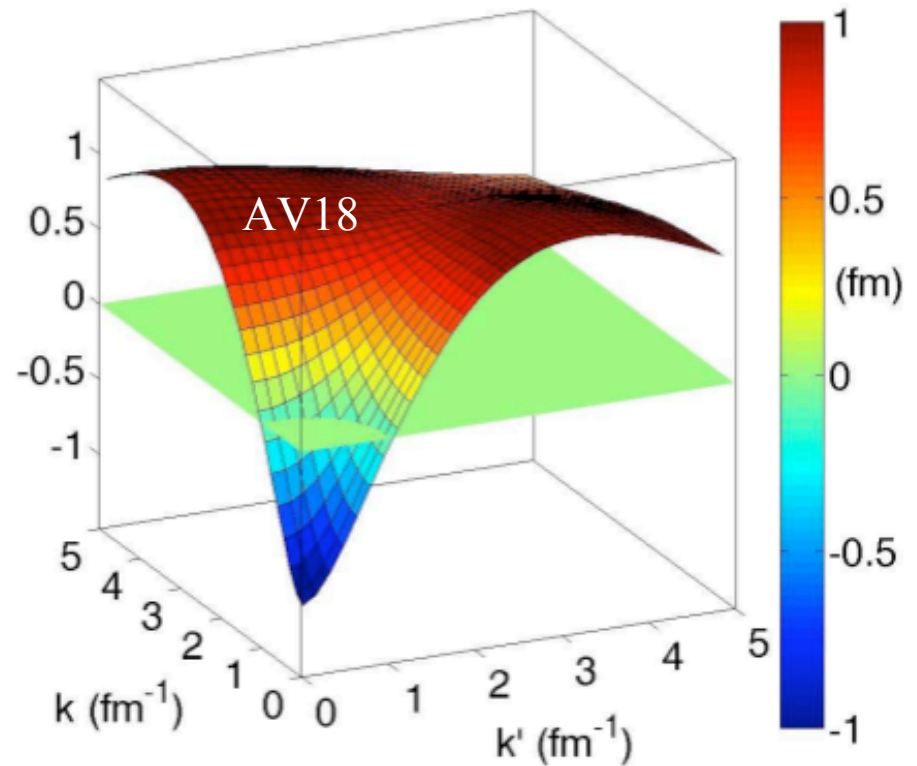
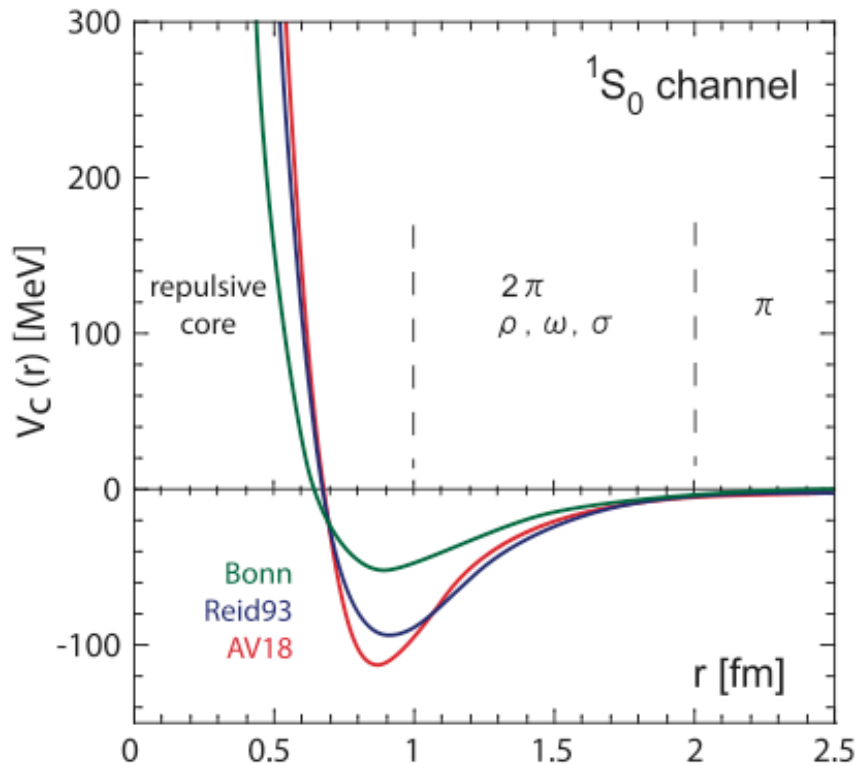
Weinberg eigenvalue diagnostic

study spectrum of $G_0(z)V |\Psi_\nu(z)\rangle = \eta_\nu(z) |\Psi_\nu(z)\rangle$ at fixed energy z

governs convergence $T(z) |\Psi_\nu(z)\rangle = (1 + \eta_\nu(z) + \eta_\nu(z)^2 + \dots) V |\Psi_\nu(z)\rangle$

can write as Schrödinger equation $(H_0 + \frac{1}{\eta_\nu(z)} V) |\Psi_\nu(z)\rangle = z |\Psi_\nu(z)\rangle$

large cutoffs lead to flipped-potential bound states of $-\lambda V$ for small λ



Weinberg eigenvalue diagnostic

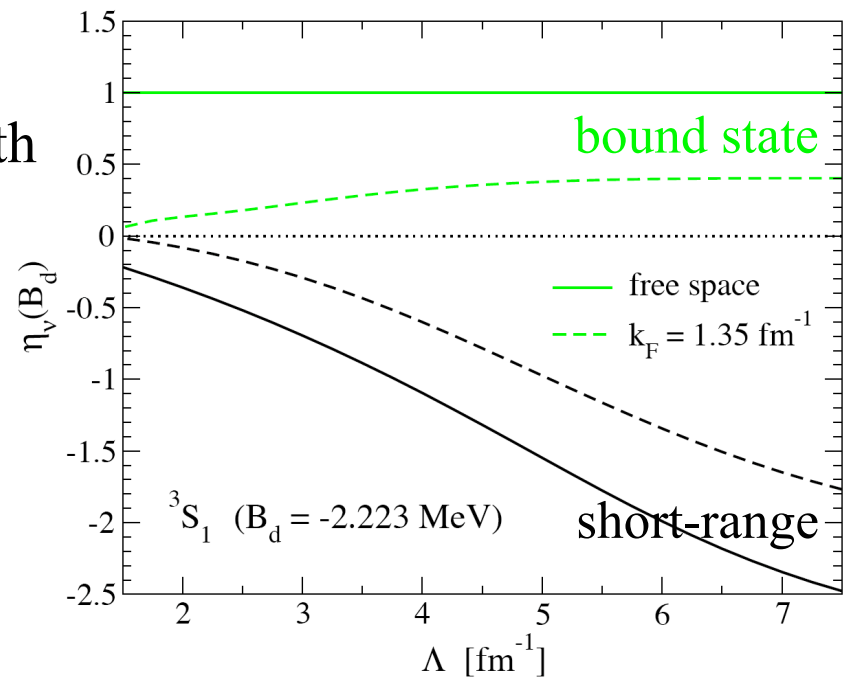
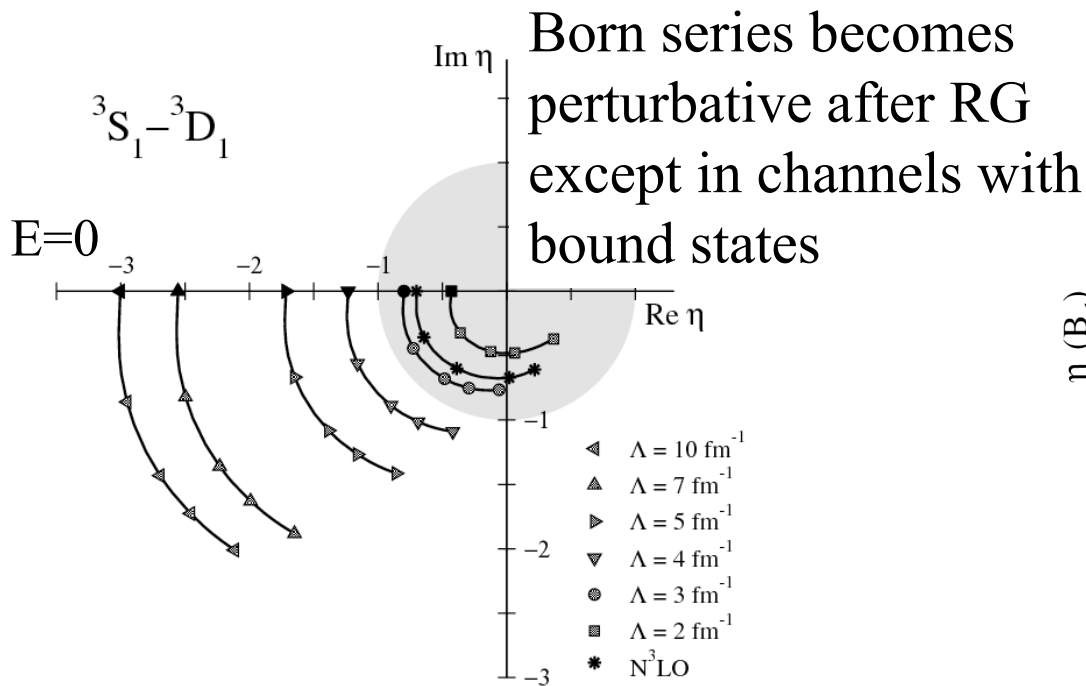
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large cutoffs lead to flipped-potential bound states of $-\lambda V$ for small λ
 \rightarrow large $\eta \rightarrow$ strong coupling to high momenta and Born series nonpert.

leads to slow convergence for all nuclei



Outline

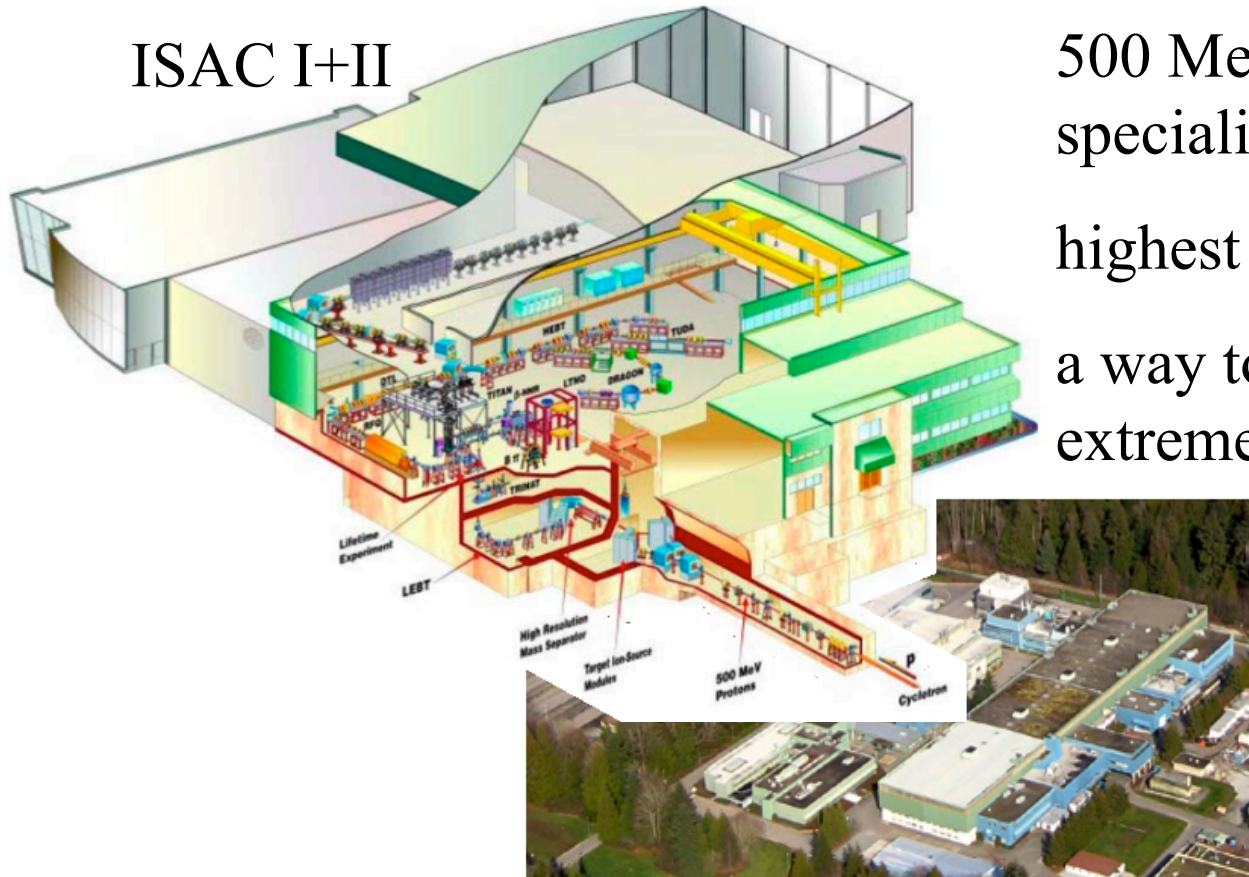
Effective field theory and renormalization group
for nuclear forces

Applications to weakly-bound and neutron-rich nuclei

Similarity renormalization group for nuclei

Neutron-rich nuclei in the laboratory

ISAC I+II



500 MeV proton beam on specialized targets

highest power at TRIUMF-ISAC

a way to add neutrons to reach extreme neutron-rich nuclei

^3He



bound

^4He



bound

^5He



unbound

^6He



bound
halo

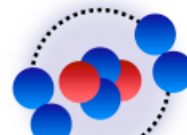
806ms

^7He



unbound

^8He



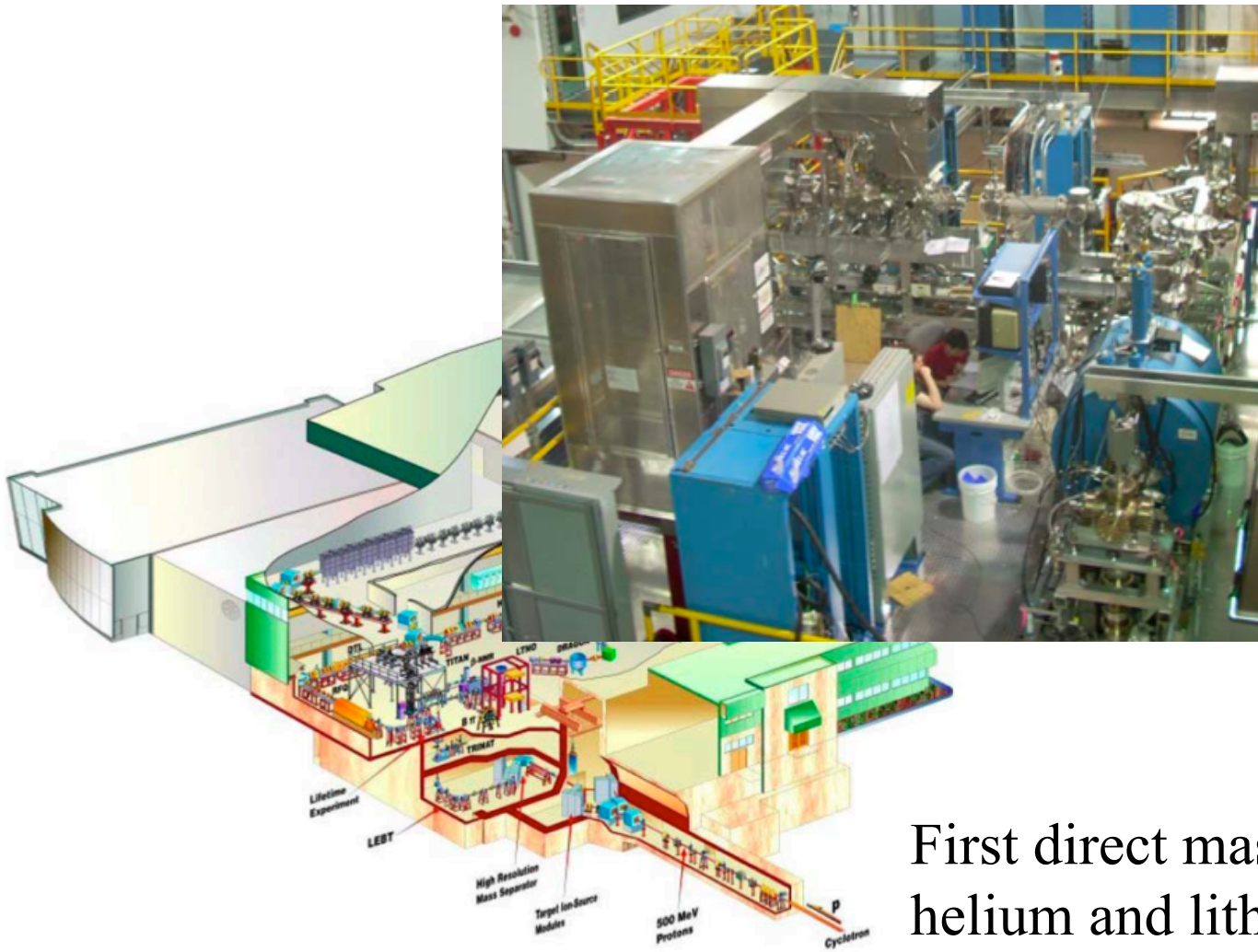
bound
halo

104ms

Novel forms of matter:
halo nuclei

^8He : most neutron-rich nucleus in the lab

TITAN Penning trap see talk by S. Ettenauer

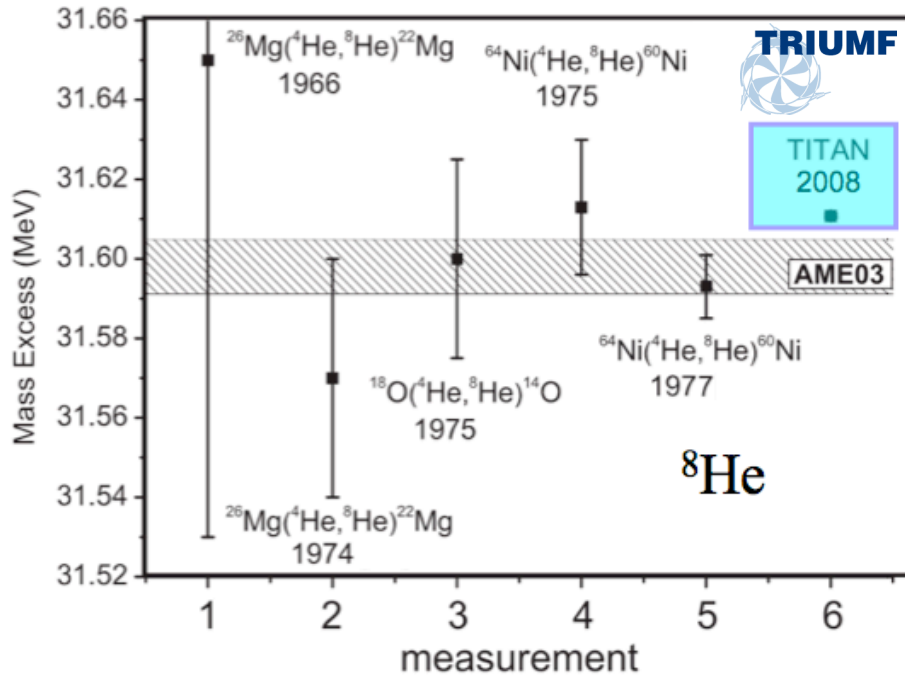


First direct mass measurement of helium and lithium halo nuclei

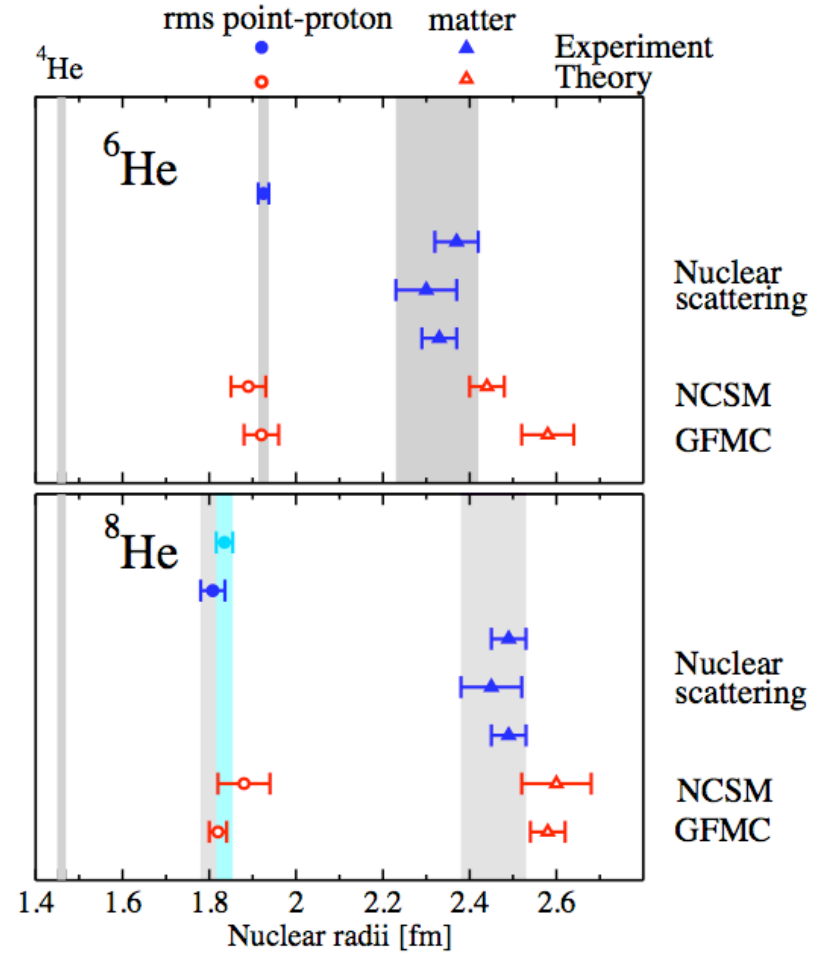
high precision $\delta m/m \sim 10^{-8}$

Neutron halos

New precision era for masses and charge radii (from isotope shifts)



TRIUMF: Ryjkov et al., PRL (2006)



ANL: Wang et al., PRL (2004)

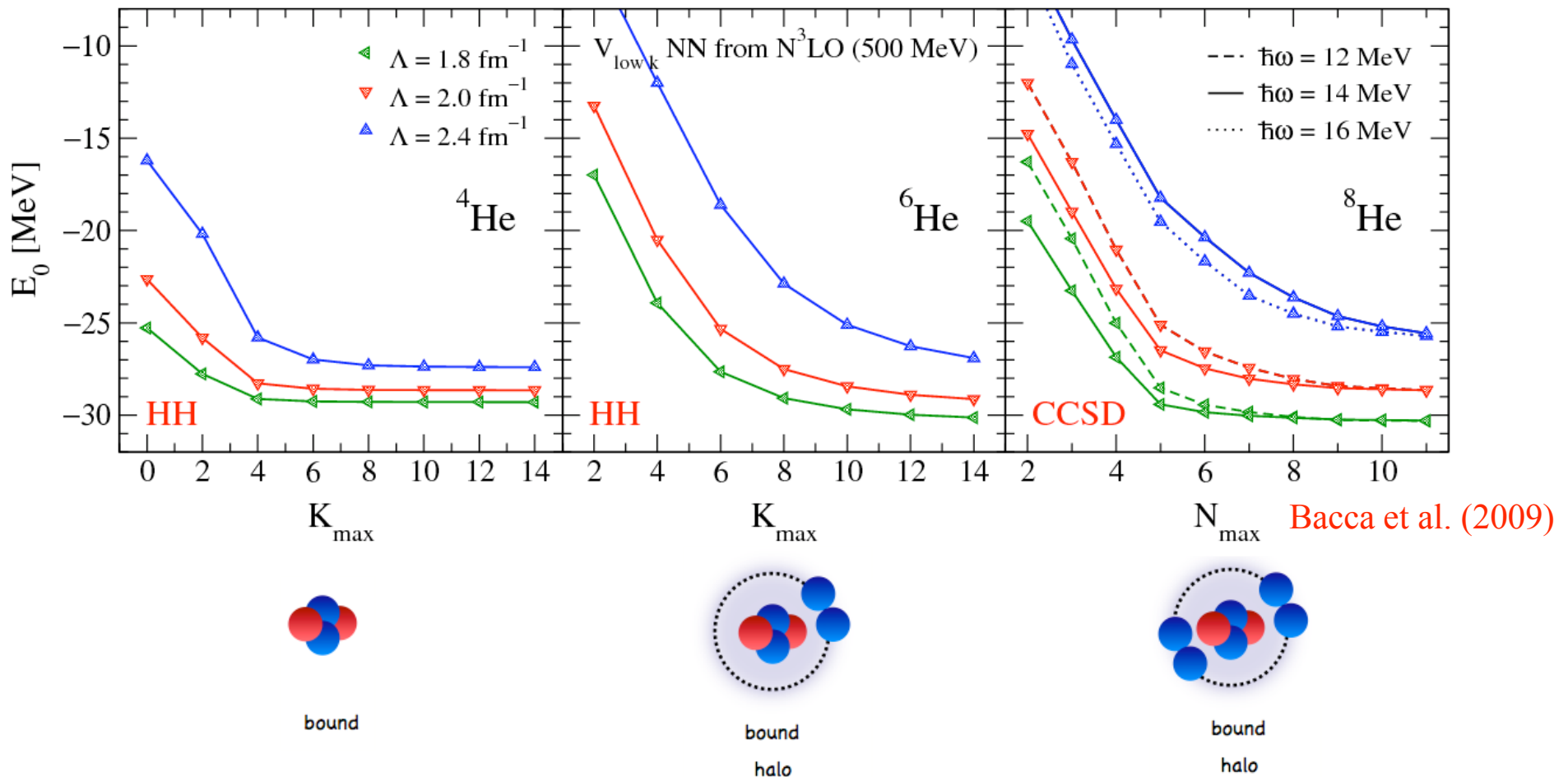
GANIL: Mueller et al., PRL (2007)

poses extraordinary challenges for theory

Neutron halos

Hyperspherical Harmonics for ${}^6\text{He}$ and Coupled-Cluster theory for ${}^8\text{He}$ describe weakly-bound nuclei with correct asymptotics

but compare convergence to stable ${}^4\text{He}$!

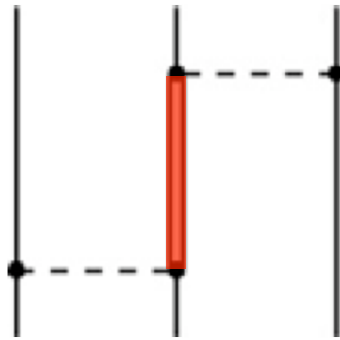


based on $N^3\text{LO}$ NN potential, RG cutoff variation \rightarrow need 3N forces

Why are there three-nucleon (3N) forces?

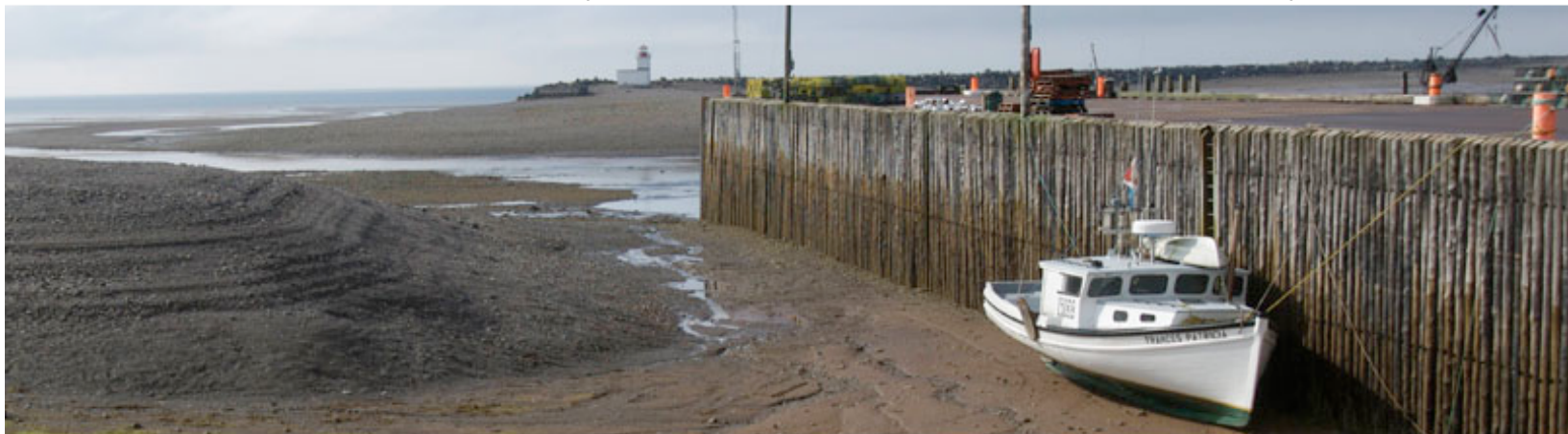
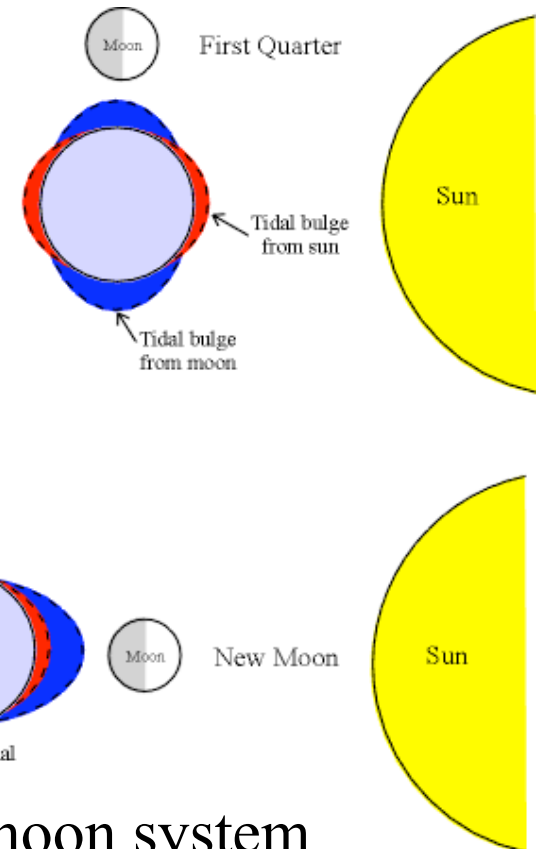
Nucleons are finite-mass composite particles,
can be excited to resonances

dominant contribution from $\Delta(1232 \text{ MeV})$



+ shorter-range parts

tidal effects leads to 3-body forces in earth-sun-moon system

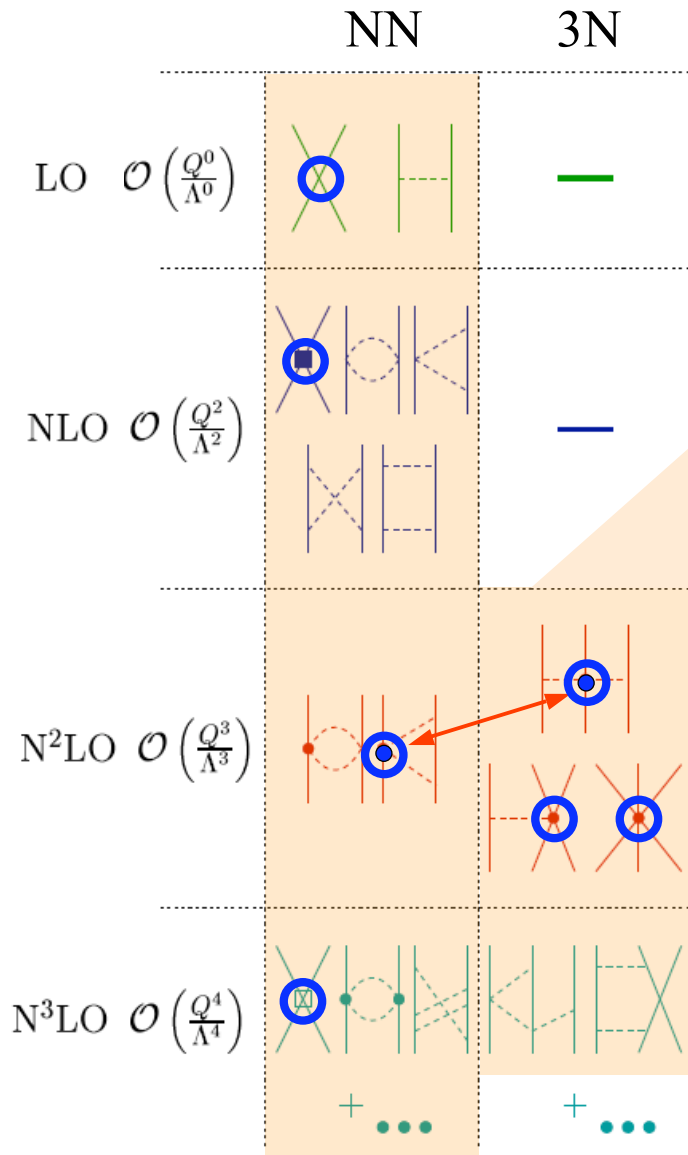


Chiral EFT for 3N forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

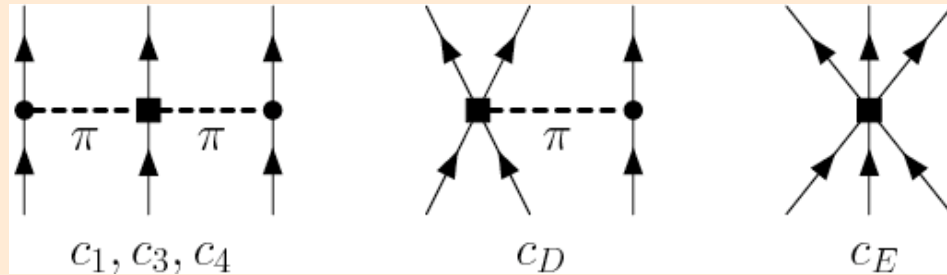
consistent NN-3N interactions

3N,4N: only 2 new couplings to N³LO



leading 3N: N²LO

van Kolck (1994), Epelbaum et al. (2002)



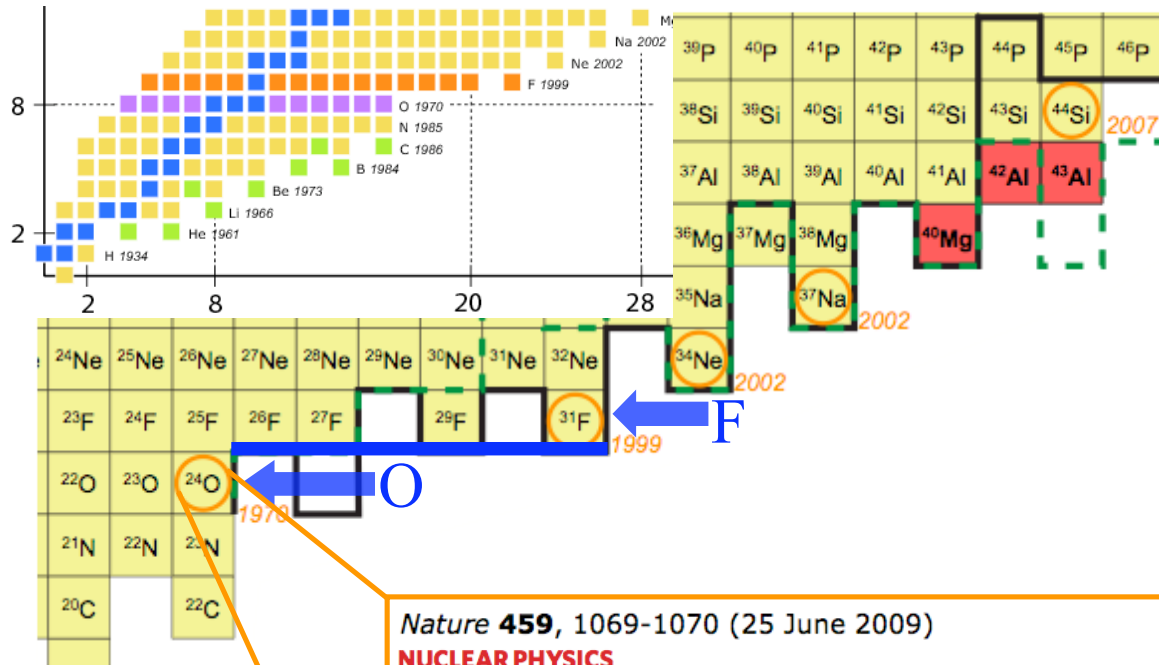
c_i from π N and NN from Meissner (2007)

$$c_1 = -0.9_{-0.5}^{+0.2}, \quad c_3 = -4.7_{-1.0}^{+1.2}, \quad c_4 = 3.5_{-0.2}^{+0.5}$$

single- Δ excitation = particular c_i

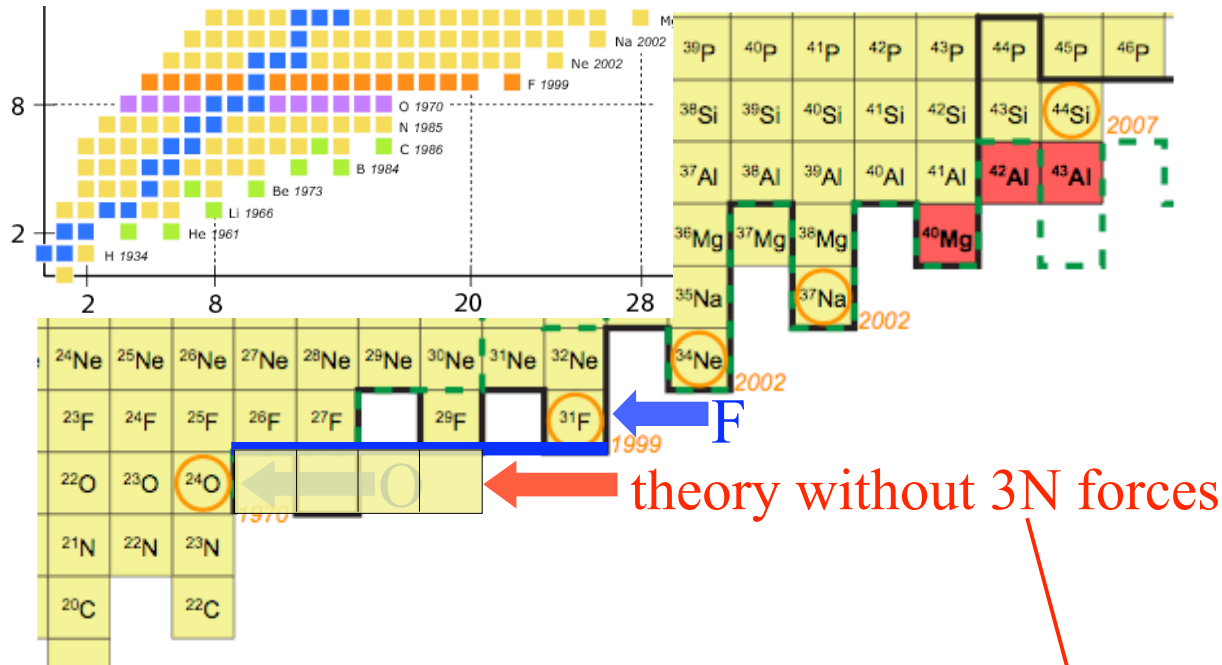
c_D, c_E fit to ${}^3\text{H}$ binding energy and ${}^4\text{He}$ radius (or ${}^3\text{H}$ beta decay half-life)

The oxygen anomaly

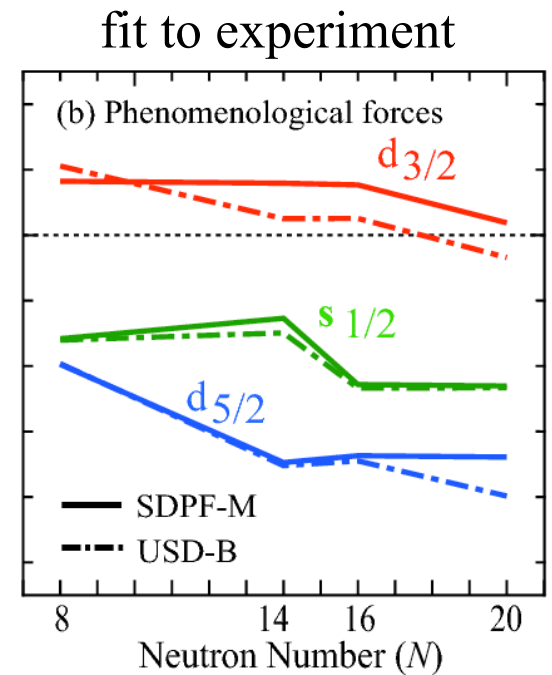
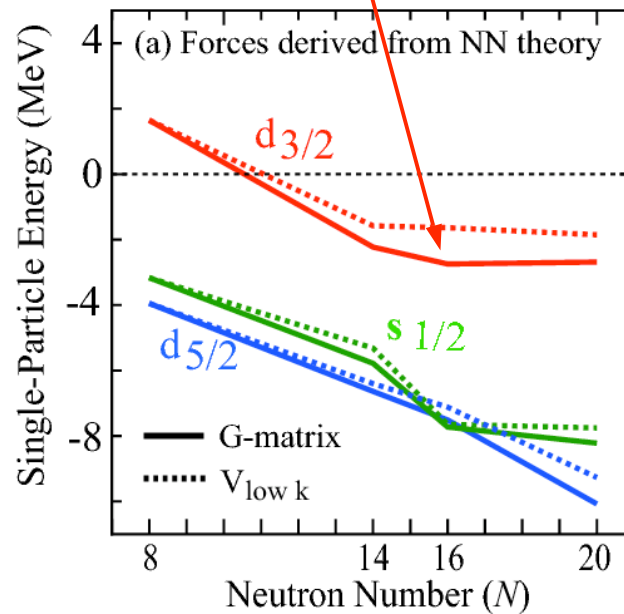
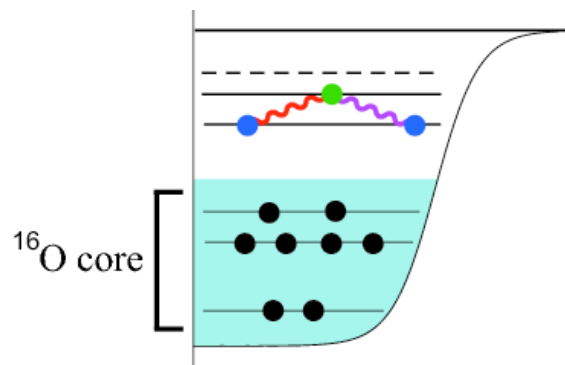


Nature **459**, 1069-1070 (25 June 2009)
NUCLEAR PHYSICS
Unexpected doubly magic nucleus
 Robert V. F. Janssens
 Nuclei with a 'magic' number of both protons and neutrons, dubbed doubly magic, are particularly stable. The oxygen isotope ^{24}O has been found to be one such nucleus — yet it lies just at the limit of stability.

The oxygen anomaly - not reproduced without 3N forces



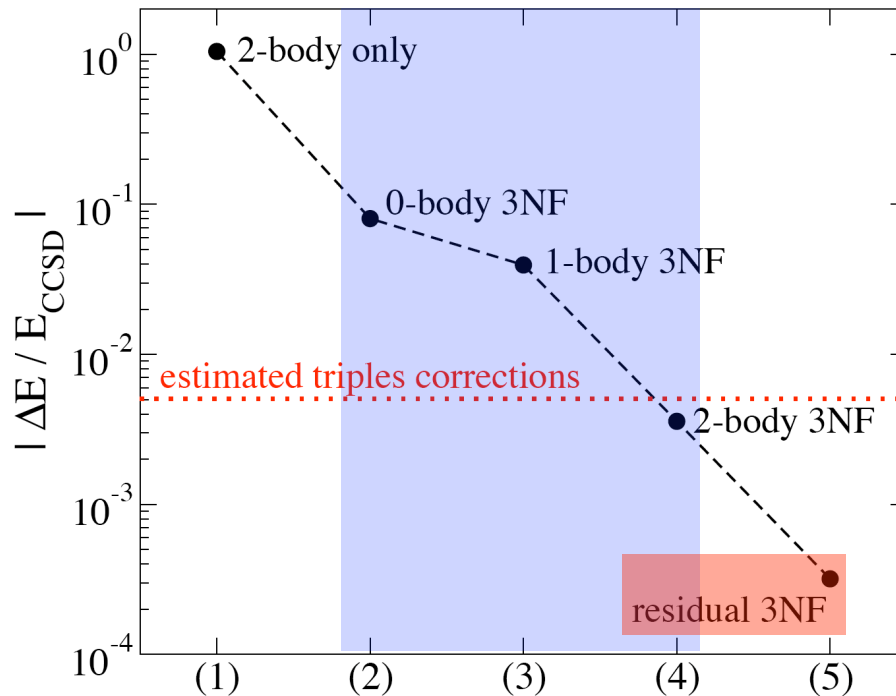
many-body theory based on two-nucleon forces:
 drip-line incorrect at ^{28}O



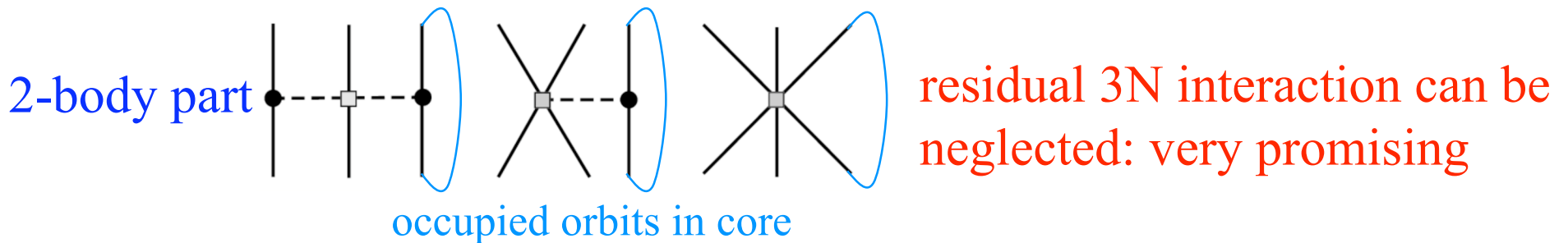
Towards 3N forces in medium-mass nuclei

Coupled-Cluster theory with 3N forces Hagen et al. (2007)

first benchmark for ^4He , work on ^{16}O in progress



normal-ordered 0-, 1- and 2-body parts of 3N forces dominate

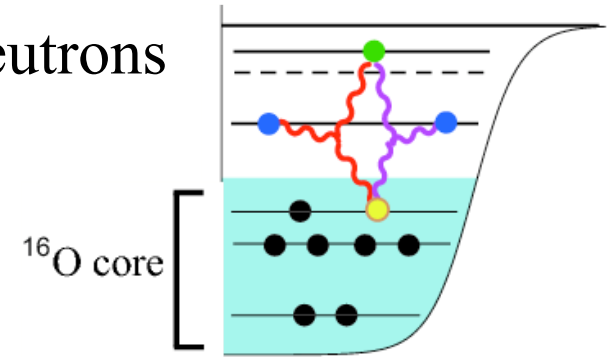


The oxygen anomaly - impact of 3N forces

include normal-ordered 2-body part of 3N forces (enhanced by core A)

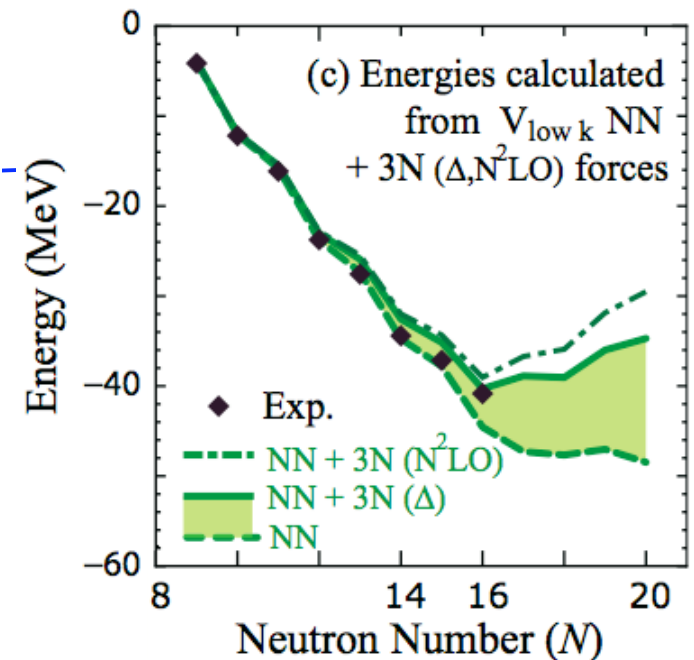
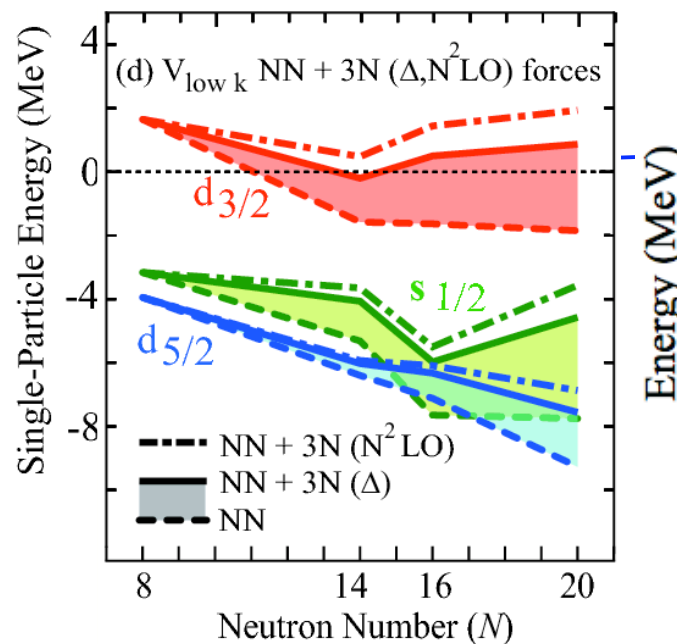
leads to repulsive interactions between valence neutrons (repulsive based on the Pauli principle)

$d_{3/2}$ orbital remains unbound



first microscopic explanation of the oxygen anomaly

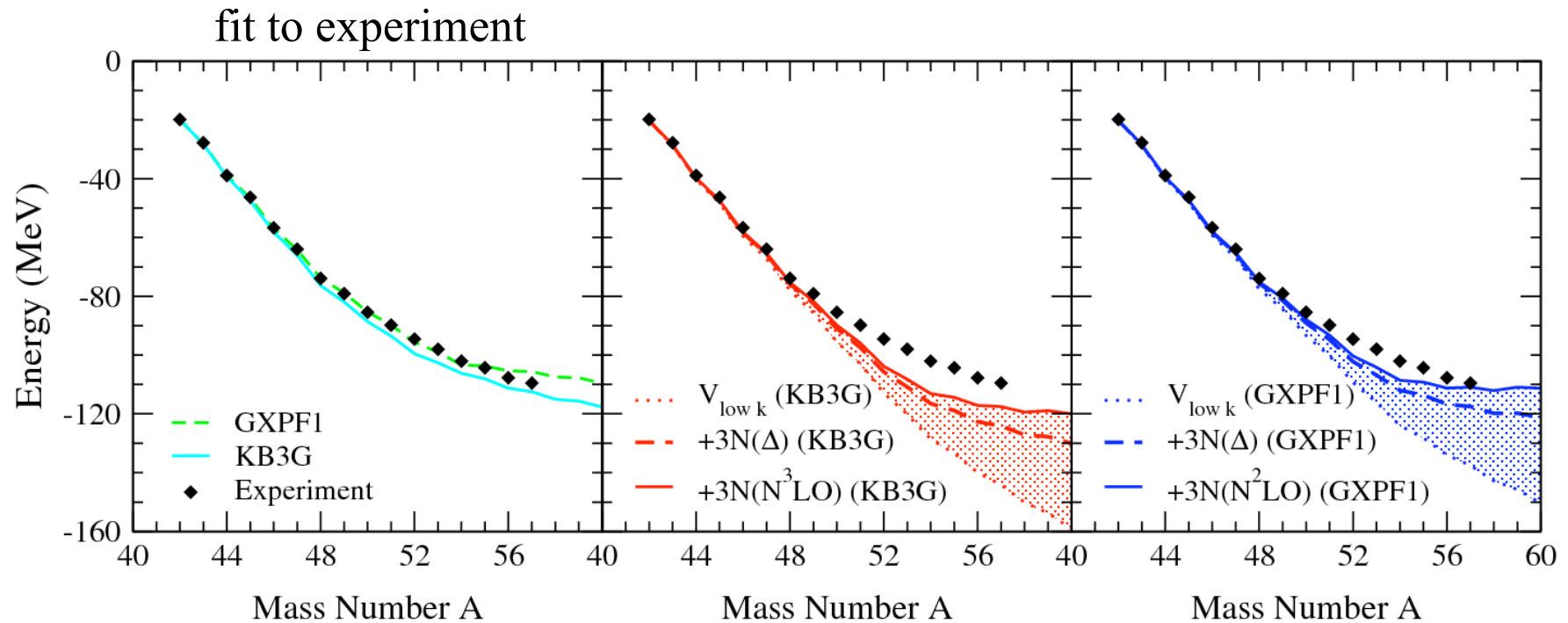
Otsuka, Suzuki, Holt, AS, Akaishi (2009)



Evolution to neutron-rich calcium isotopes

repulsive 3N contributions also key for calcium ground-state energies

Holt, Otsuka, AS, Suzuki, in prep.



Evolution to neutron-rich calcium isotopes

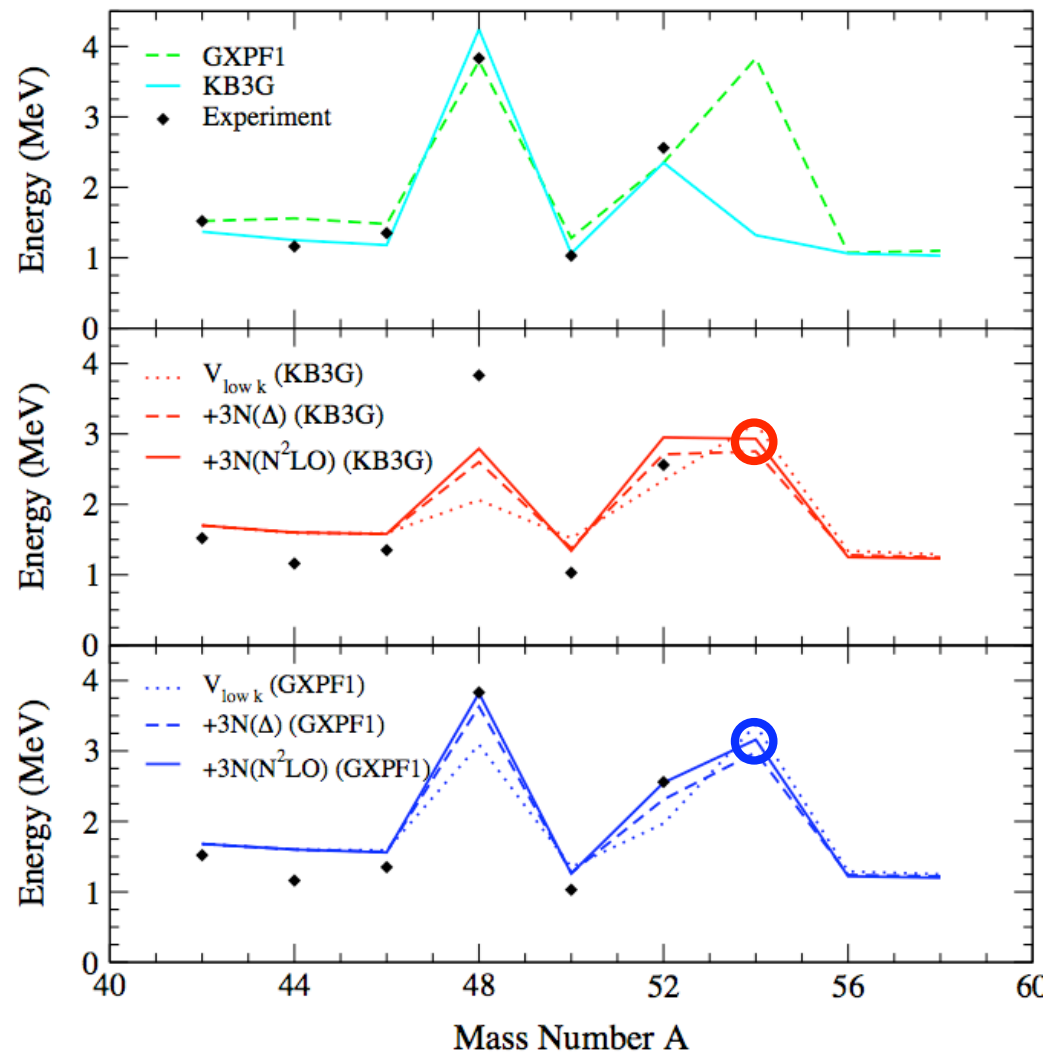
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Holt, Otsuka, AS, Suzuki, in prep.

3N mechanism important for shell structure: 2^+ excitation energy in ^{48}Ca

N=28 shell closure due to 3N and single-particle effects (^{41}Ca)

predict 2^+ excitation energy in ^{54}Ca at ~ 3 MeV



Weinberg eigenvalue diagnostic

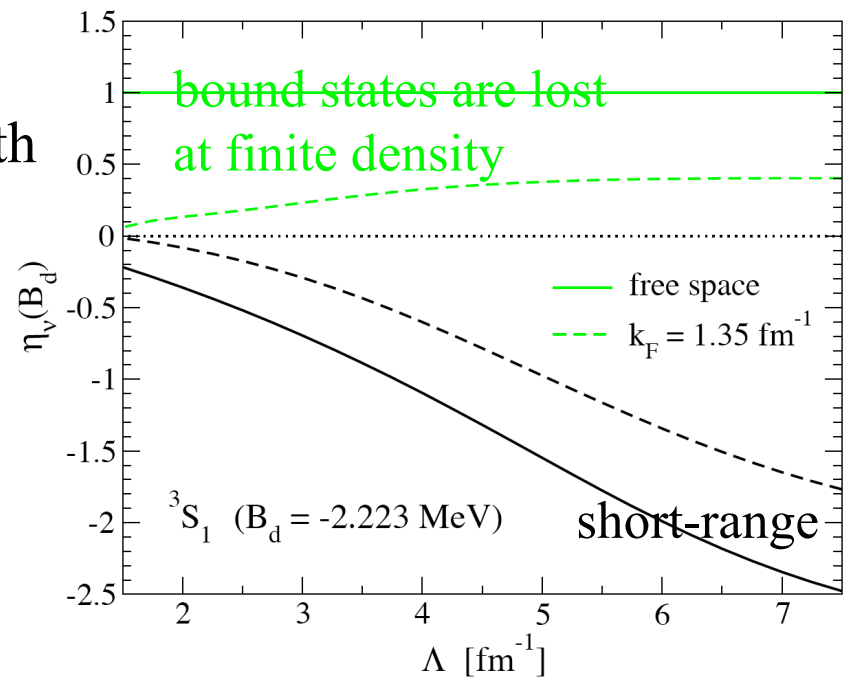
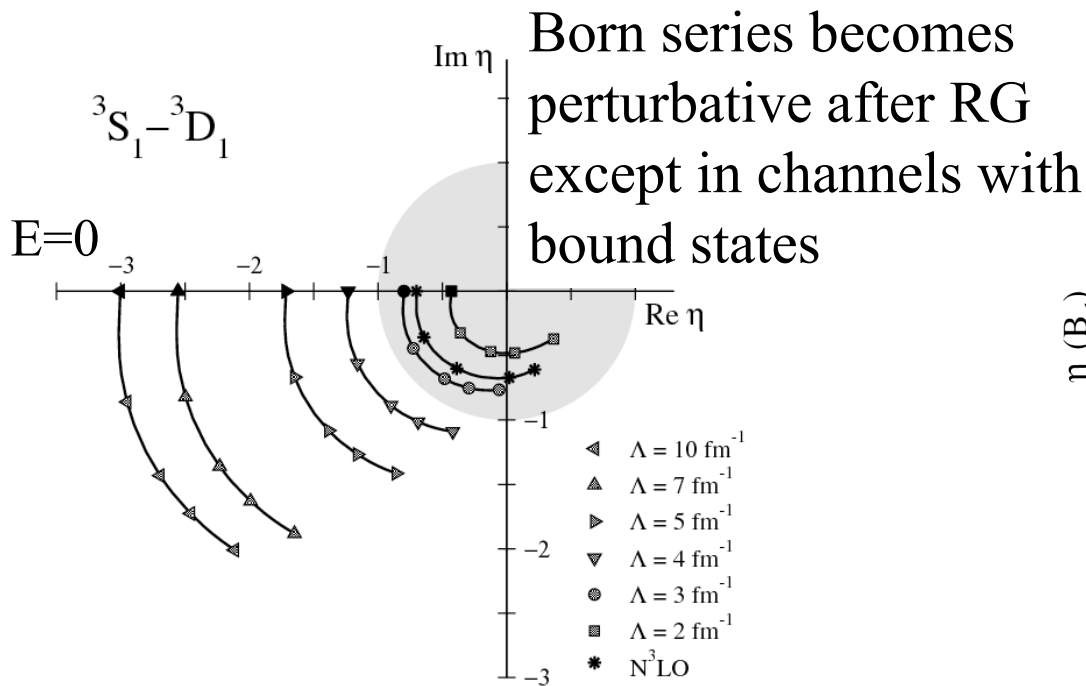
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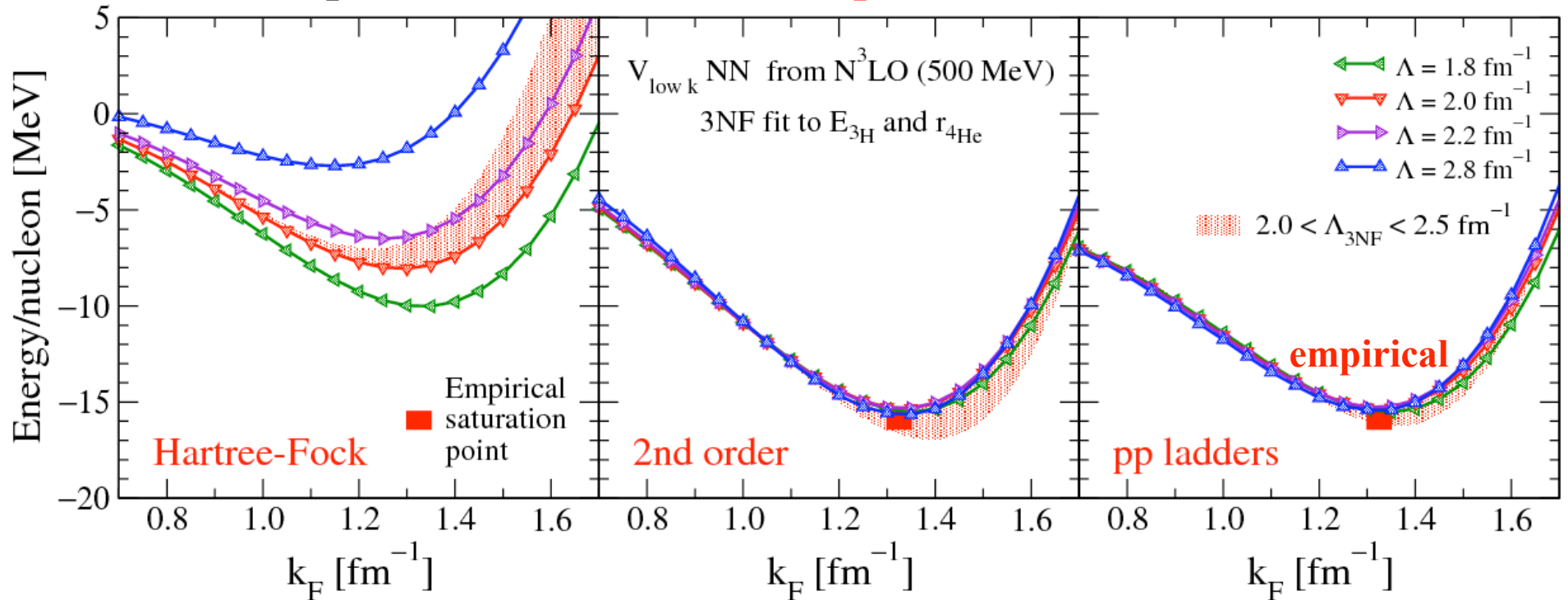
Is nuclear matter perturbative with chiral EFT and RG?

conventional Bethe-Brueckner-Goldstone expansion (sums ladders):
no, due to nonpert. cores (flipped-V bound states) and off-diag coupling

start from chiral EFT and RG evolution:

nuclear matter converged at \approx 2nd order, 3N drives saturation

weak cutoff dependence, **but need to improve 3N treatment**



Bogner, AS, Furnstahl, Nogga (2009)

see talk by K. Hebeler for neutron matter and more

Outline

Effective field theory and renormalization group
for nuclear forces

Applications to weakly-bound and neutron-rich nuclei

Similarity renormalization group for nuclei

Similarity RG

unitary transformations to band-diagonal $V_{\text{srg}}(\lambda)$ from flow equations

Glazek, Wilson (1993), Wegner (1994)

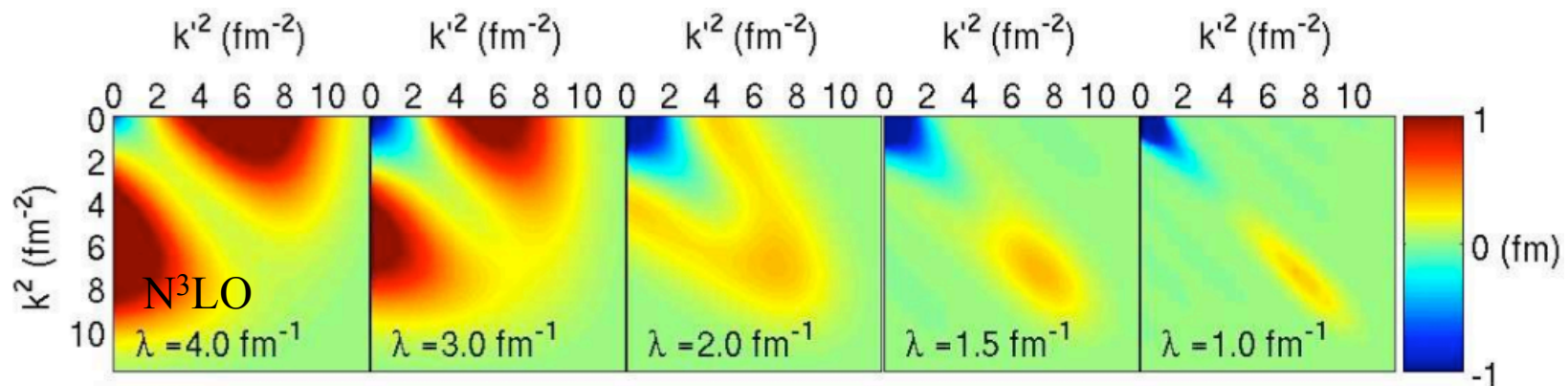
$$\frac{dH_s}{ds} = [\eta_s, H_s] = [[G_s, H_s], H_s]$$

evolution driven towards nonzero part of generator G_s

with flow operator $G_s = T_{\text{rel}}$ and resolution $\lambda = s^{-1/4}$

Bogner, Furnstahl, Perry, ...

$$\frac{dV_s(k, k')}{ds} = -(k^2 - k'^2)^2 V_s(k, k') + \frac{2}{\pi} \int_0^\infty q^2 dq (k^2 + k'^2 - 2q^2) V_s(k, q) V_s(q, k')$$

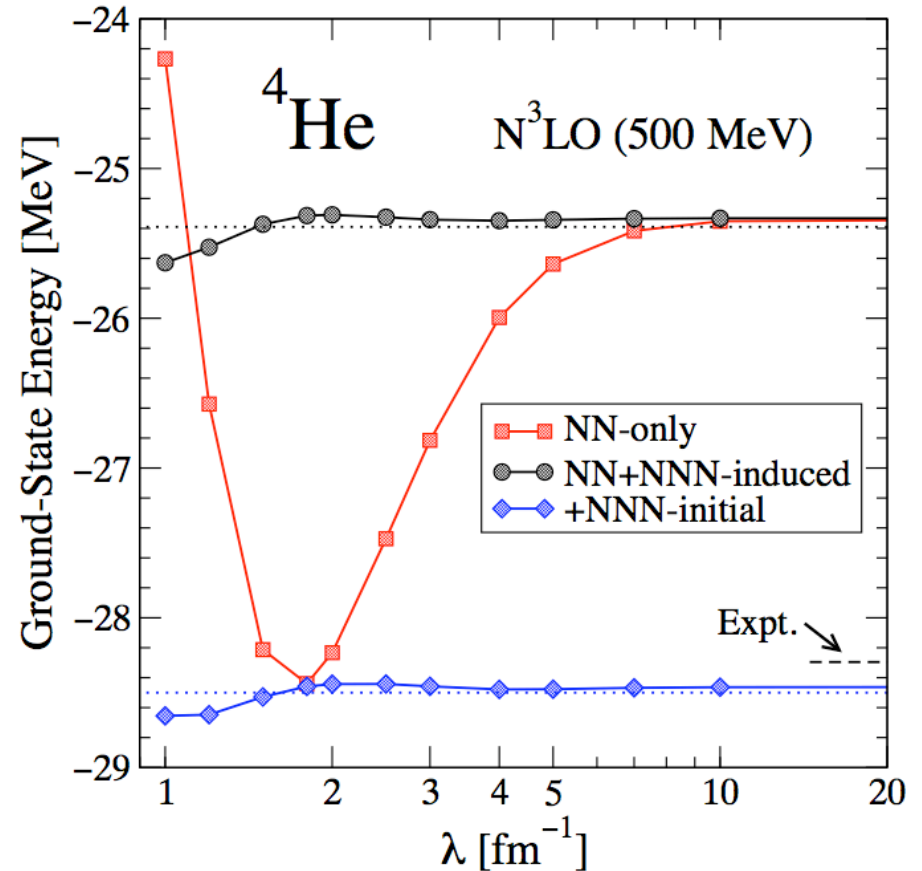
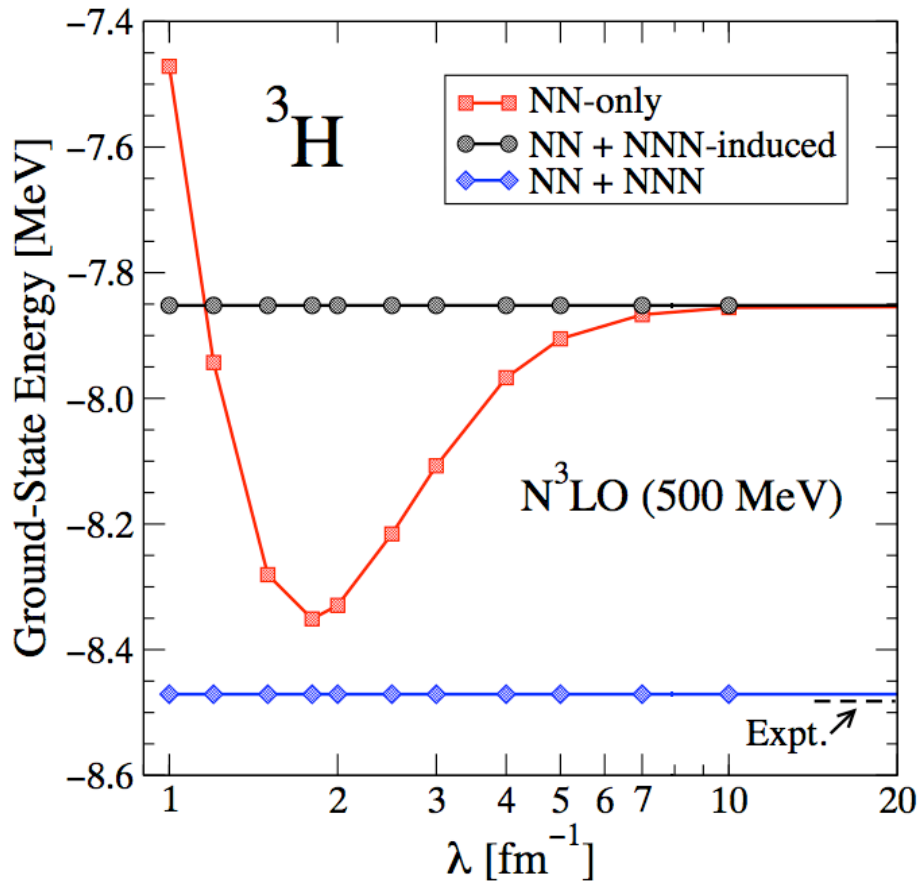


SRG decouples high momenta with similar low-momentum universality

SRG evolution of 3N forces

start from chiral NN and 3N interactions,
SRG evolution in harmonic oscillator basis

Jurgenson, Navratil, Furnstahl (2009)



induced many-body interactions consistent with truncation error in EFT

In-medium SRG for nuclei

$$H = \sum_{12} T_{12} a_1^\dagger a_2 + \frac{1}{(2!)^2} \sum_{1234} \langle 12|V|34 \rangle a_1^\dagger a_2^\dagger a_4 a_3 + \frac{1}{(3!)^2} \sum_{123456} \langle 123|V^{(3)}|456 \rangle a_1^\dagger a_2^\dagger a_3^\dagger a_6 a_5 a_4$$

normal-order Hamiltonian with respect to reference state
(e.g., Hartree-Fock ground state)

$$H = E_0 + \sum_{12} f_{12} \{a_1^\dagger a_2\} + \frac{1}{(2!)^2} \sum_{1234} \langle 12|\Gamma|34 \rangle \{a_1^\dagger a_2^\dagger a_4 a_3\} + \frac{1}{(3!)^2} \sum_{123456} \langle 123|\Gamma^{(3)}|456 \rangle \{a_1^\dagger a_2^\dagger a_3^\dagger a_6 a_5 a_4\}$$

with 0-, 1- and 2-body normal-ordered parts

$$E_0 = \langle \Phi|H|\Phi \rangle = \sum_1 T_{11} n_1 + \frac{1}{2} \sum_{12} \langle 12|V|12 \rangle n_1 n_2 + \frac{1}{3!} \sum_{123} \langle 123|V^{(3)}|123 \rangle n_1 n_2 n_3$$

$$f_{12} = T_{12} + \sum_i \langle 1i|V|2i \rangle n_i + \frac{1}{2} \sum_{ij} \langle 1ij|W|2ij \rangle n_i n_j ,$$

$$\langle 12|\Gamma|34 \rangle = \langle 12|V|34 \rangle + \sum_i \langle 12i|V^{(3)}|34i \rangle n_i ,$$

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with 0-, 1- and 2-body normal-ordered parts and **in-medium SRG eqns**
e.g., for nuclear matter with $\eta=[f,\Gamma]$ [see Bogner et al., Kehrein \(2006\)](#)

$$\frac{dE_0}{ds} = \frac{1}{2} \sum_{1234} (f_{12} - f_{34}) |\Gamma_{1234}|^2 n_1 n_2 \bar{n}_3 \bar{n}_4,$$

$$\frac{df_1}{ds} = \sum_{234} (f_{41} - f_{23}) |\Gamma_{4123}|^2 (\bar{n}_2 \bar{n}_3 n_4 + n_2 n_3 \bar{n}_4),$$

$$\begin{aligned} \frac{d\Gamma_{1234}}{ds} = & -(f_{12} - f_{34})^2 \Gamma_{1234} + \frac{1}{2} \sum_{ab} (f_{12} + f_{34} - 2f_{ab}) \Gamma_{12ab} \Gamma_{ab34} (1 - n_a - n_b) + \sum_{ab} (n_a - n_b) \\ & \times \left\{ \Gamma_{a1b3} \Gamma_{b2a4} [(f_{a1} - f_{b3}) - (f_{b2} - f_{a4})] - \Gamma_{a2b3} \Gamma_{b1a4} [(f_{a2} - f_{b3}) - (f_{b1} - f_{a4})] \right\}, \end{aligned}$$

approx. includes many-body forces and sums pp, hh, ph diagrams

In-medium SRG for nuclei Tsukiyama, Bogner, AS, in prep.

decouple 1p1h, 2p2h,... ApAh sectors from reference state

want to suppress pph and ph couplings,

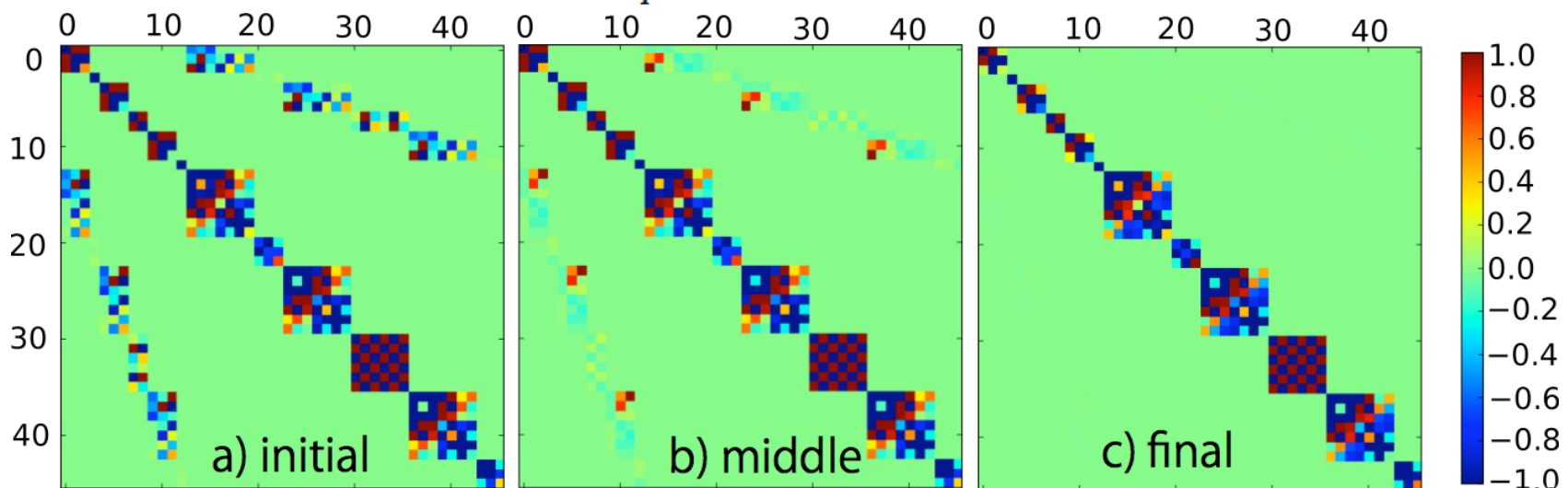
all other (normal-ordered) couplings annihilate reference state

minimal choice: $\eta(s) = [H^d(s), H(s)] = [H^d(s), H^{od}(s)]$

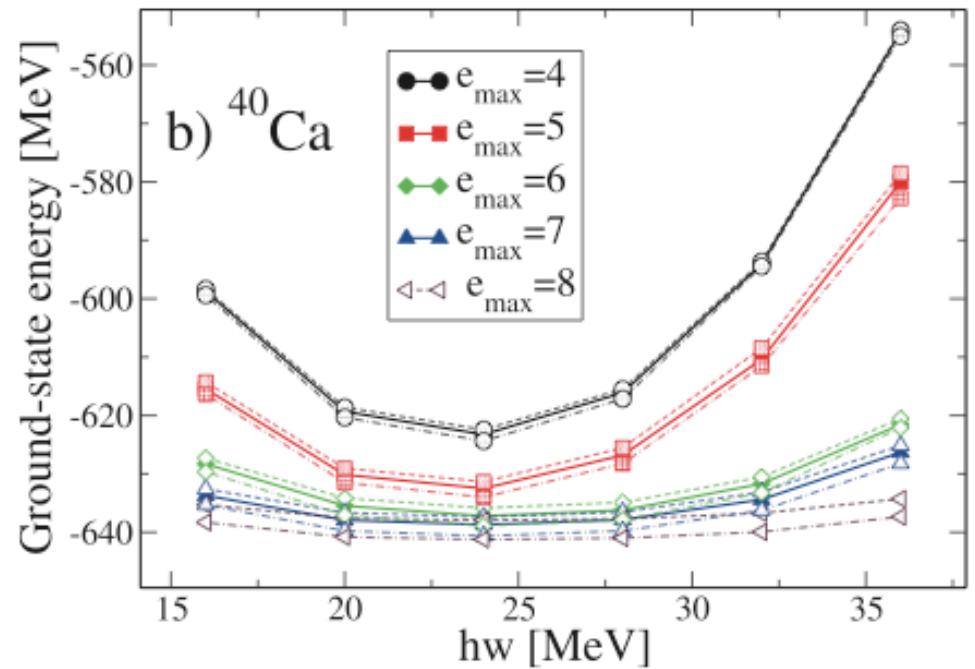
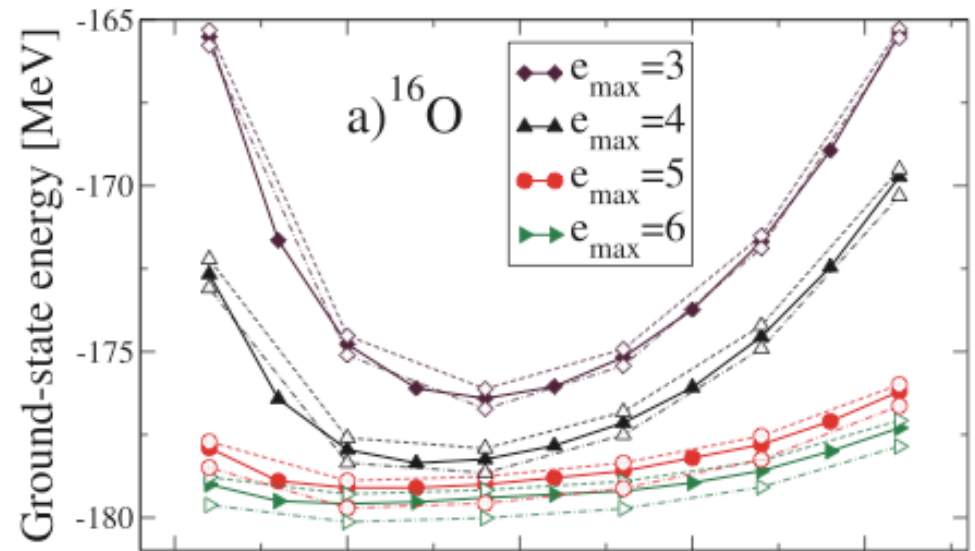
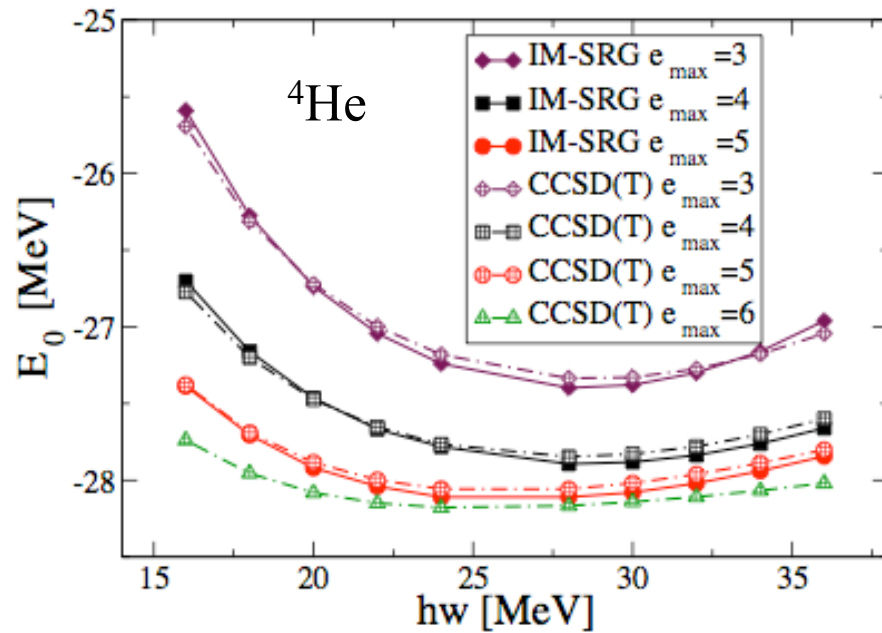
$$H^{od}(s) = g^{od}(s) + \Gamma^{od}(s)$$

$$\Gamma^{od}(s) = \sum_{pp'hh'} \Gamma_{pp'hh'}(s) a_p^\dagger a_{p'}^\dagger a_h a_{h'} + h.c.$$

$$g^{od}(s) \sum_{ph} g_{ph}(s) a_p^\dagger a_h + h.c.$$



In-medium SRG for nuclei Tsukiya, Bogner, AS, in prep.



first results for closed-shell nuclei
very promising convergence,
results comparable to CCSD(T)

can be used to derive nonperturbative valence-shell effective interactions

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Summary

development of effective field theory and the renormalization group
with approaches from light to heavy nuclei to matter in astrophysics
based on the same interactions

3N forces are a frontier for weakly-bound and neutron-rich nuclei

exciting intersections with problems in many related areas

exciting experiments to study neutron-rich matter in the laboratory
at rare isotope beam (RIB) facilities worldwide

