

Effective Field Theory and Efimov Physics

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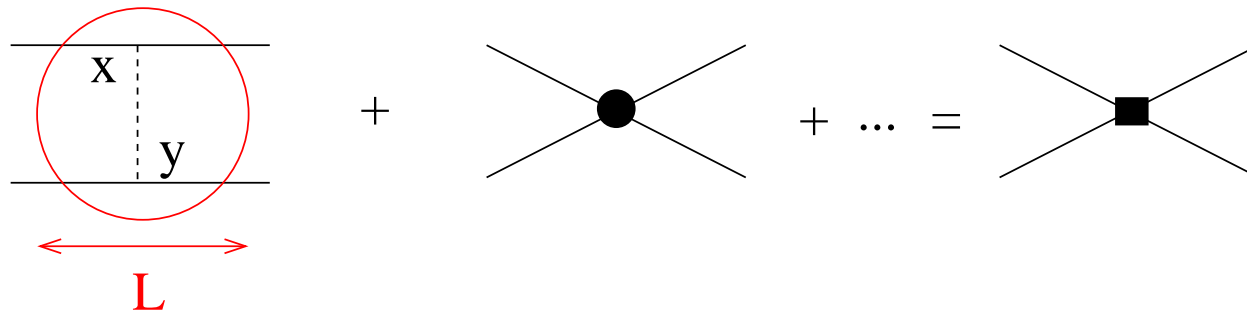
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DFG

Collaborators: E. Braaten, D. Canham, K. Helfrich, D. Kang, S. Kreuzer, D. Lee, L. Platter, R. Springer, ...

- Introduction
- Resonant Interactions and Efimov Physics
- Effective Field Theory for Large Scattering Length
- Applications
 - Ultracold atoms
 - Hadronic Molecules
- Summary and Outlook

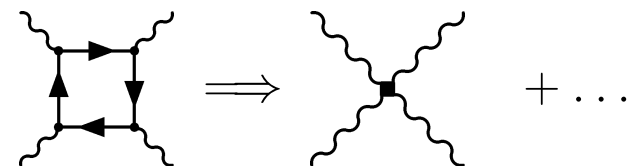
- Limited resolution in low-energy processes



- Typical momentum: $p = \hbar/L \Rightarrow$ resolution: $|\mathbf{x} - \mathbf{y}| \gtrsim L$
- Short-distance physics not resolved
 - capture in low-energy constants using renormalization
 - include long-range physics explicitly
- Systematic, model independent → universal properties
- Classic example: light-light-scattering (Euler, Heisenberg, 1936)

Simpler theory for $\omega \ll m_e$:

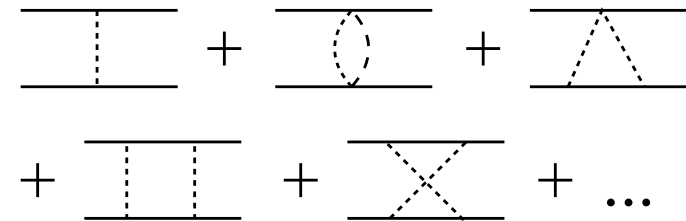
$$\mathcal{L}_{QED}[\psi, \bar{\psi}, A_\mu] \rightarrow \mathcal{L}_{eff}[A_\mu]$$



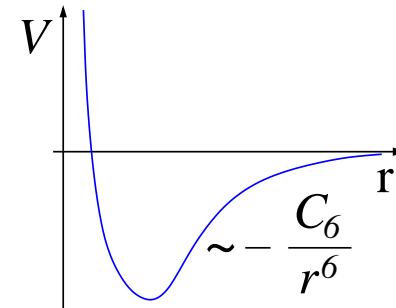
- **Renormalization Group:** Systems with very different fundamental interactions can behave similarly at low energies

⇒ **Universal properties**

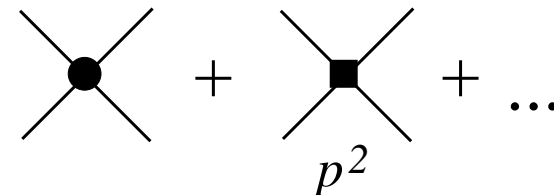
- **Nuclear Physics:**
Nucleons and pions



- **Atomic Physics:**
Born-Oppenheimer
plus van der Waals



- ⇒ **At sufficiently low energy:**
contact interactions



- Large scattering length: $|a| \gg \ell \sim r_e, l_{vdW}, \dots$
- Natural expansion parameter: $\ell/|a|, k\ell, \dots$

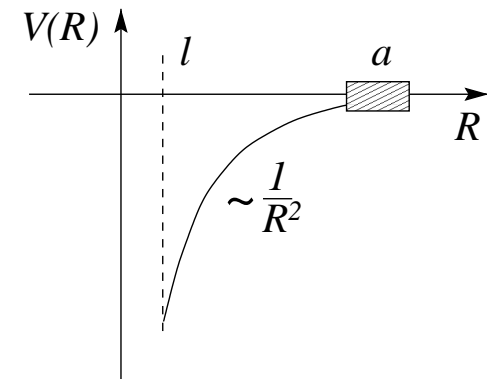
$$a > 0 \quad \Longrightarrow \quad B_d = \frac{1}{2\mu a^2} + \mathcal{O}(\ell/a)$$

- Atomic physics:
 - ^4He : $a \approx 104 \text{ \AA} \gg r_e \approx 7 \text{ \AA} \sim l_{vdW} \longrightarrow B_d \approx 100 \text{ neV}$
 - Feshbach resonances \Longrightarrow variable scattering length
- Nuclear physics: S -wave NN -scattering, halo nuclei, ...
 - $^1S_0, ^3S_1$: $|a| \gg r_e \sim 1/m_\pi \longrightarrow B_d \approx 2.2 \text{ MeV}$
 - $^6\text{He} \Longrightarrow P$ -wave universality?
- Particle physics:
 - $X(3872)$ as a $D^0\bar{D}^{0*}$ molecule? ($J^{PC} = 1^{++}$)
 $B_X = m_{D^0} + m_{D^{0*}} - m_X = (0.3 \pm 0.4) \text{ MeV}$

(V. Efimov, Phys. Lett. **33B** (1970) 563)

- Three-body system with large scattering length a
- Hyperspherical coordinates: $R^2 = (r_{12}^2 + r_{13}^2 + r_{23}^2)/3$
- Schrödinger equation simplifies for $|a| \gg R \gg l$:

$$-\frac{\hbar^2}{2m} \left[\frac{\partial^2}{\partial R^2} + \frac{s_0^2 + 1/4}{R^2} \right] f(R) = \underbrace{-\frac{\hbar^2 \kappa^2}{m}}_E f(R)$$

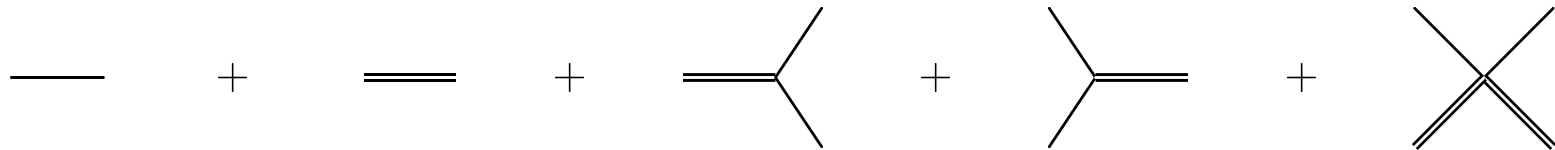


- **Singular Potential:** renormalization required
- **Boundary condition at small R :** breaks scale invariance
 \implies **dependence of observables on 3-body parameter (and a)**
- **EFT formulation:** boundary condition \implies 3-body interaction

- Effective Lagrangian

(Kaplan, 1997; Bedaque, HWH, van Kolck, 1999)

$$\mathcal{L}_d = \psi^\dagger \left(i\partial_t + \frac{\vec{\nabla}^2}{2m} \right) \psi + \frac{g_2}{4} d^\dagger d - \frac{g_2}{4} (d^\dagger \psi^2 + (\psi^\dagger)^2 d) - \frac{g_3}{36} d^\dagger d \psi^\dagger \psi + \dots$$



- Interacting dimeron propagator \longrightarrow sum bubbles



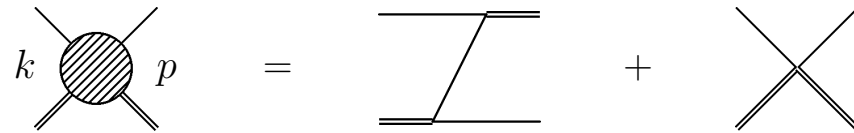
- Two-body amplitude $\mathcal{T}_2(k, k)$:

$$\propto \left[\underbrace{8\pi/g_2 + 2\Lambda/\pi}_{1/a} + ik \right]^{-1} + \dots$$

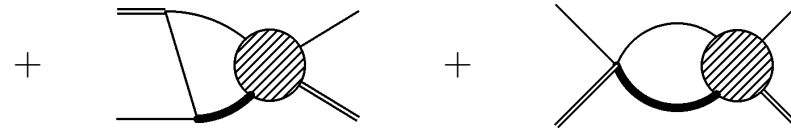
- Matching: $g_2 \longleftarrow a, B_d$

- RG fixed points of $g_2(\Lambda)$: $a = 0$ and $a = \infty$

- Higher order corrections: perturbation theory



- Three-body equation :

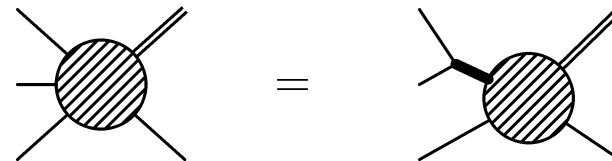


$$\mathcal{T}_3(k, p) = M(k, p) + \frac{4}{\pi} \int_0^\Lambda dq q^2 M(q, p) D_d(q) \mathcal{T}_3(k, q)$$

with $M(k, p) = \underbrace{F(k, p)}_{\text{1-atom exchange}} \underbrace{-\frac{g_3}{9g_2^2}}_{H(\Lambda)/\Lambda^2}$

($g_3 = 0, \Lambda \rightarrow \infty \rightarrow$ Skorniakov, Ter-Martirosian '57)

- Three-body recombination:



- Observables are independent of regulator/cutoff Λ

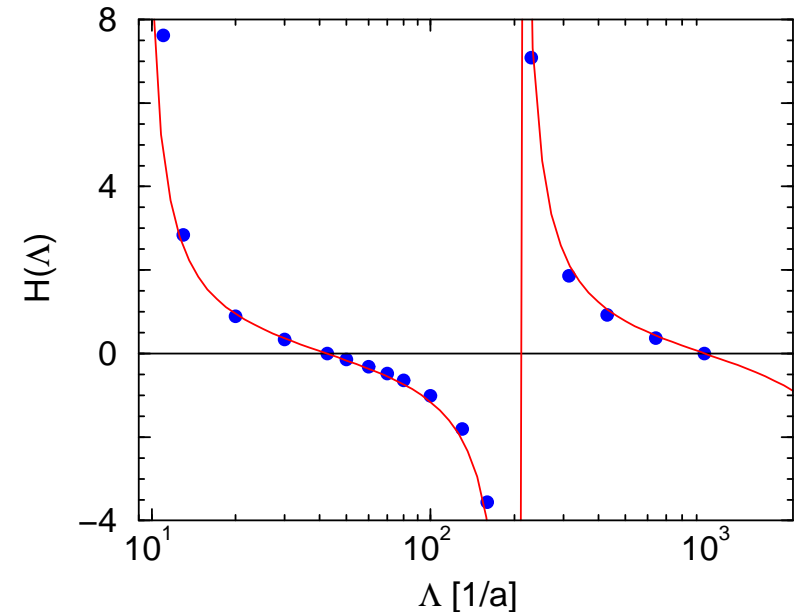
⇒ Running coupling $H(\Lambda)$

- $H(\Lambda)$ periodic: **limit cycle**

$$\Lambda \rightarrow \Lambda e^{n\pi/s_0} \approx \Lambda (22.7)^n$$

(cf. Wilson, 1971)

- Full scale invariance broken to discrete subgroup

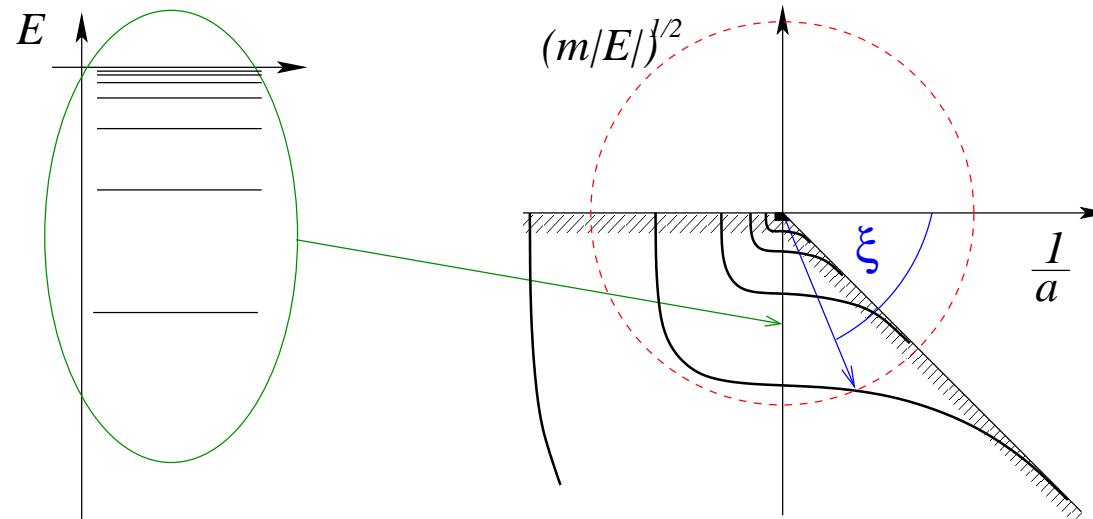


$$H(\Lambda) = \frac{\cos(s_0 \ln(\Lambda/\Lambda_*) + \arctan(s_0))}{\cos(s_0 \ln(\Lambda/\Lambda_*) - \arctan(s_0))}, \quad s_0 \approx 1.00624$$

- **Limit cycle** \iff **Discrete scale invariance**
- **Matching:** $\Lambda_* \longleftarrow B_t, K_3, \dots \longrightarrow \kappa_*, a_*, a'_*$

- Universal spectrum of three-body states

(V. Efimov, Phys. Lett. **33B** (1970) 563)

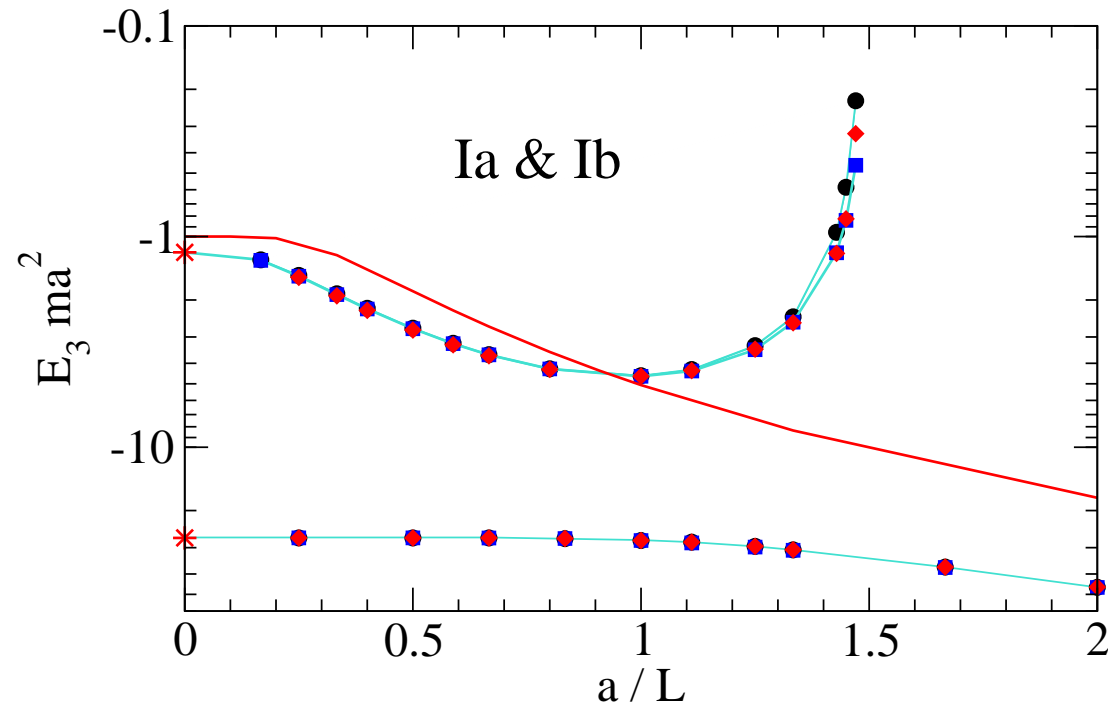


- Discrete scale invariance for fixed angle ξ
- **Geometrical spectrum** für $1/a \rightarrow 0$

$$B_3^{(n)} / B_3^{(n+1)} \xrightarrow{1/a \rightarrow 0} 515.035\dots$$

- Ultracold atoms \implies variable scattering length

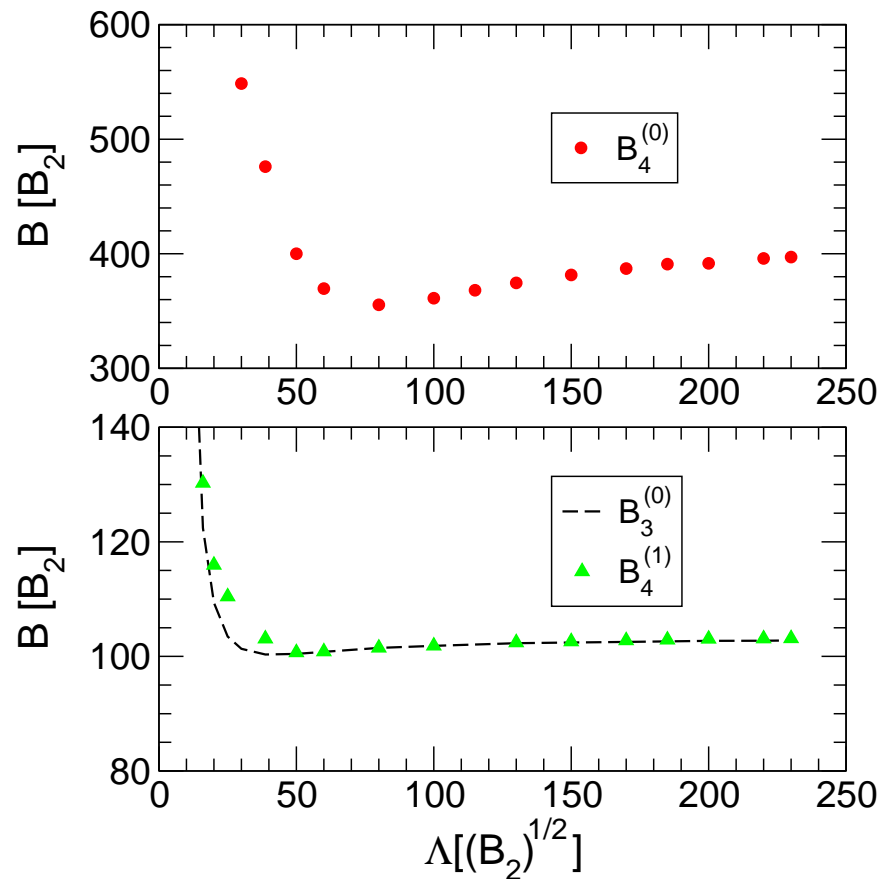
- Modification of spectrum by cubic box ($V = L^3$)
 - Box provides infrared cutoff $1/L \Rightarrow$ calculable in EFT
 - Box breaks rotational invariance \Rightarrow partial wave mixing
 - 3-momenta quantized $\vec{p} = \vec{n} (2\pi/L) \Rightarrow$ 3d sum equation



Kreuzer, HWH, Eur. Phys. J. A **43**, 229 (2010)

- Higher partial waves:
 - few percent effect from $l = 4$ for $L \gtrsim 2R$
 - very small for threshold state
- Indications for universal scaling of finite volume corrections
 - $\implies L_{10\%}$ scales with $1/\sqrt{B_3} \sim R_3$
- First results for 3 nucleons \implies Lattice QCD
 - $\implies \Delta B_t/B_t \approx 300\%$ for $L \approx 3$ fm
- Limit cycle in finite volume
(Kreuzer, HWH, in progress)

- Extension to 4-body system in effective QM approach
- No four-body parameter at LO (Platter, HWH, Meißner, 2004)



- 3- and 4-body observables are correlated (Tjon line, ...)

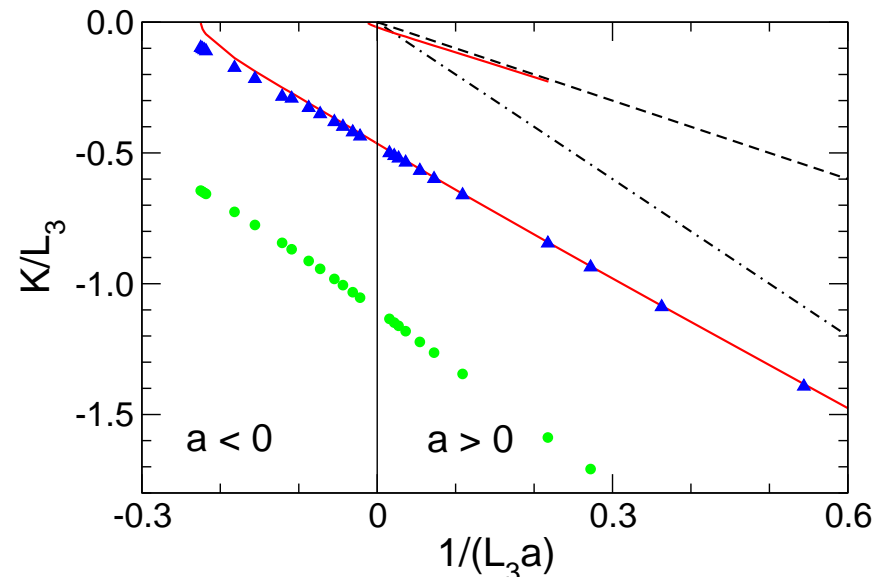
- Universal properties of 4-body system with large a
 - Bound state spectrum, scattering observables, ...
- “Efimov-plot”: 4-body bound state spectrum as function of $1/a$

$$K = \text{sign}(E) \sqrt{m|E|}$$

$$B_4^{(0)} = 5B_3^{(0)} \quad (1/a \equiv 0)$$

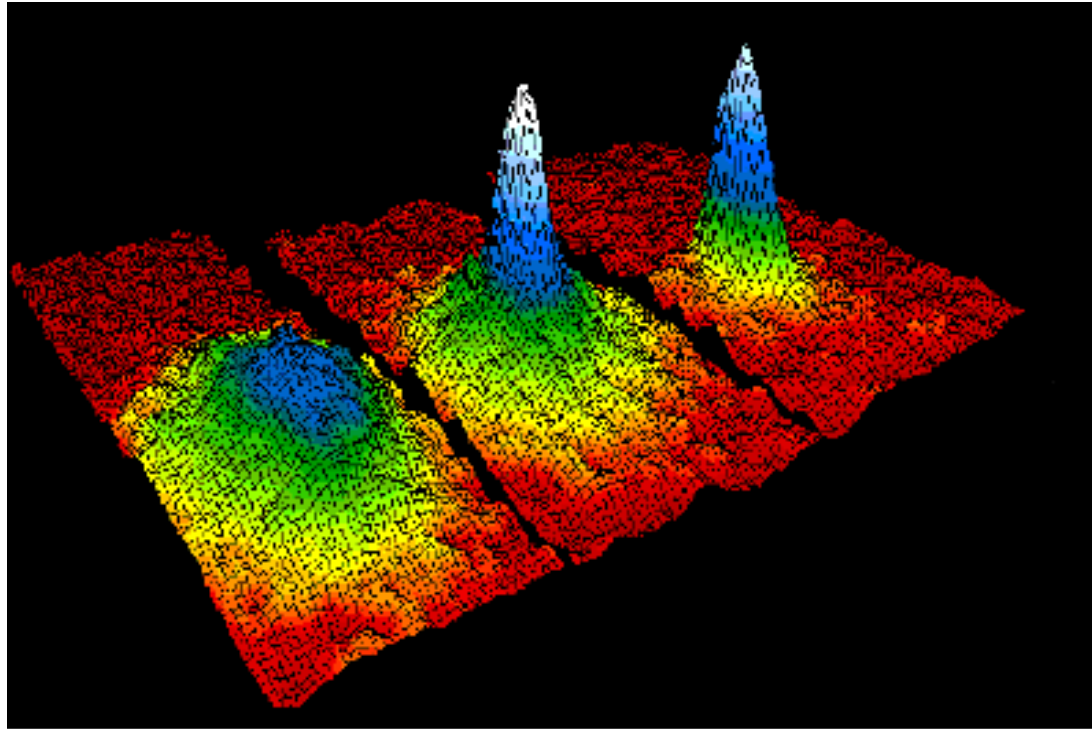
$$B_4^{(1)} = 1.01B_3^{(0)}$$

(Platter, HWH, EPJA **32** (2007) 113)



- Improved theoretical description and signature in Cs loss data
von Stecher, D’Incao, Greene, Nature Physics **5** (2009) 417
Ferlandino, Knoop, Berninger, Harm, D’Incao, Nägerl, Grimm, PRL **102** (2009) 140401
- Four-body recombination (Wang, Esry; Mehta et al.)

- Velocity distribution ($T = 400$ nK, 200 nK, 50 nK)

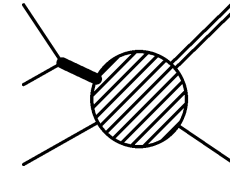


(Source: <http://jilawww.colorado.edu/bec/>)

- Few-body loss rates provide window on Efimov physics
- Variable scattering length via Feshbach resonances

- Recombination into weakly-bound dimer:

3 atoms \rightarrow dimer + atom \Rightarrow **loss of atoms**



- Recombination constant: $\dot{n}_A = -K_3 n_A^3$

- Scattering length dependence for $a > 0$:

(Nielsen, Macek, 1999; Esry, Greene, Burke, 1999; Bedaque, Braaten, HWH, 2000)

$$K_3 \approx 201.3 \sin^2 [s_0 \ln(a\kappa_*) + 1.16] \frac{\hbar a^4}{m}, \quad s_0 \approx 1.00624..$$

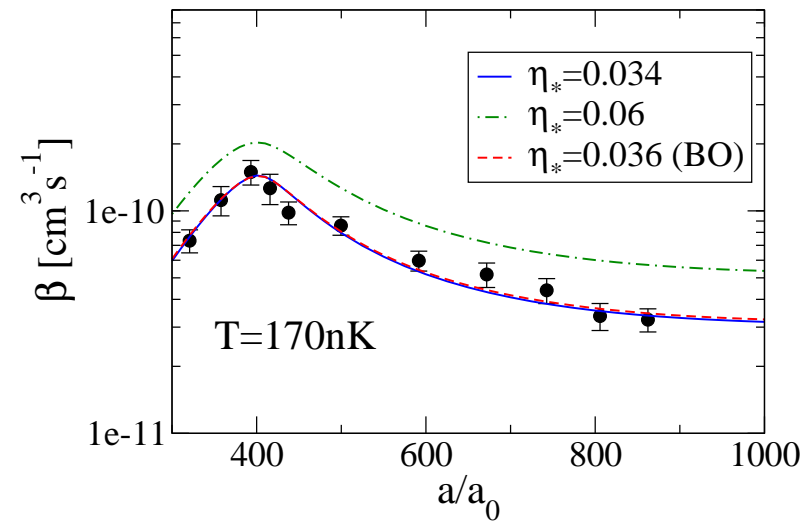
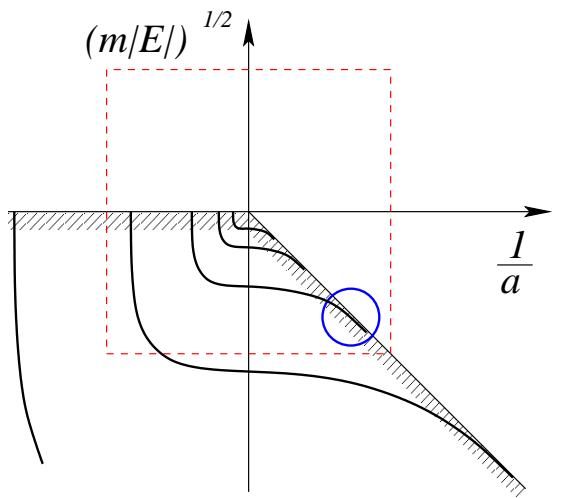
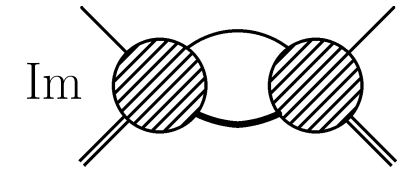
- Modification from deep dimers: Efimov states acquire width

$$\implies \kappa_* \rightarrow \kappa_* \exp(i\eta_*/s_0)$$

- Recombination into deep dimers \implies Efimov resonances

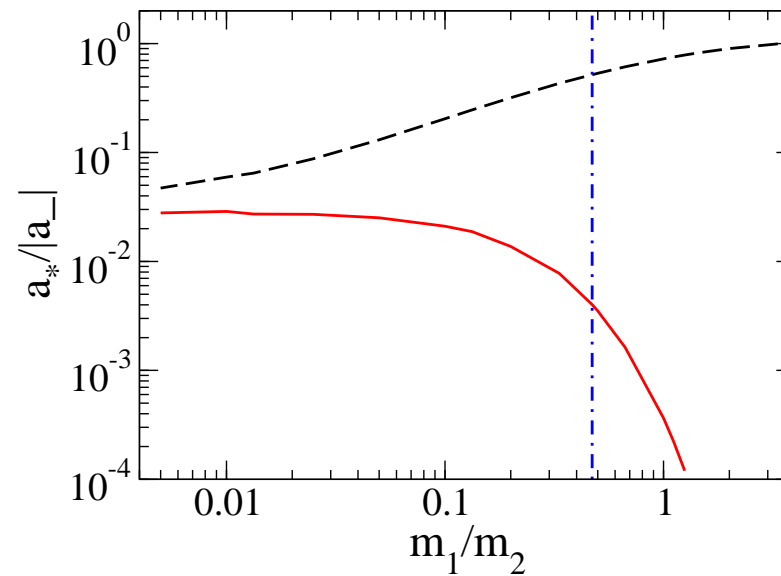
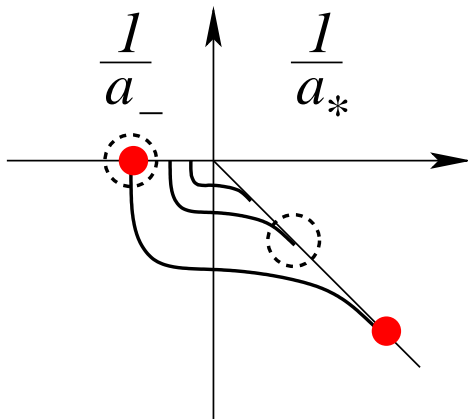
- Evidence for Efimov trimers in ^{133}Cs , ^6Li , ^7Li , ^{39}K

- Dimer Relaxation: $a + d \rightarrow a + D$ (energetic)
- Relaxation constant: $\dot{n}_A = \dot{n}_D = -\beta n_A n_D$
- Recent experiment: Knoop et al. (Innsbruck), Nature Physics **5** (2009) 227
- Finite temperature $T \sim T_c$: Bose-Einstein average ?



Helfrich, HWH, EPL **86** (2009) 53003

- Recent experiment with heteronuclear mixture of ^{41}K and ^{87}Rb atoms (Barontini et al. (Florence), Phys. Rev. Lett. **103** (2009) 043201)
 \Rightarrow Connected K-Rb-Rb resonances for $a > 0$ and $a < 0$
- Ratio of resonance positions: $a_*/|a_-|$

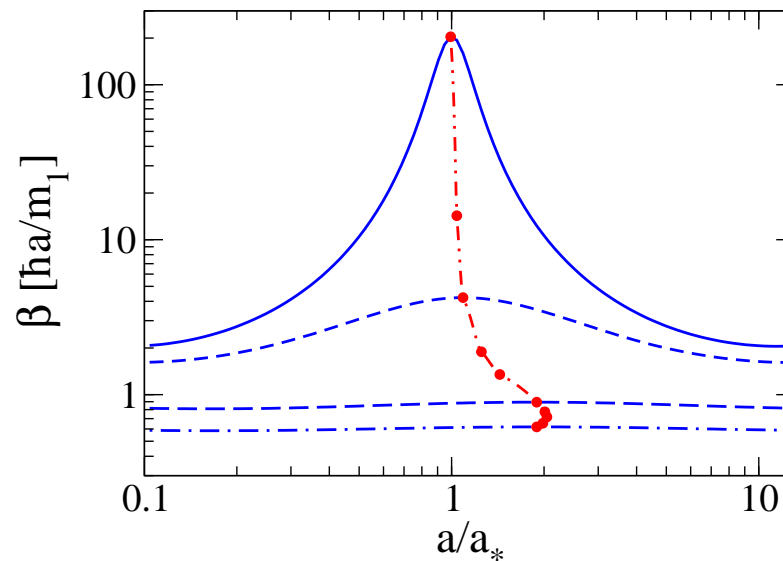


Helfrich, HWH, Petrov, arXiv:1001.4371

- K-Rb-Rb:** $m_1/m_2 = 41/87 \approx 0.47 \Rightarrow \exp(\pi/s_0) \approx 128$
 $a_*/|a_-|: 2.7 \text{ (Exp)} \Leftrightarrow 0.52 \text{ (Th)}$

- Identification of atom-dimer resonance in purely atomic sample through rescattering processes? (cf. Barontini et al., 2009)
- Calculate loss rate above threshold: $E/B_d = -1, -0.95, -0.5, 0$

$$\eta_* = 0.1$$



Helfrich, HWH, Petrov, arXiv:1001.4371

- Should not exclude other interpretations of data
- Analytical and numerical results for mass dependence of recombination and relaxation rates (cf. D’Incao, Esry, 2006)

- Universal properties for resonant P -wave interactions?
 \implies not in the usual sense
- Interactions corresponding to a_1 and r_1 required for consistent renormalization (Bertulani, HWH, van Kolck, 2002)
- Two relevant directions near RG fixed point (Barford, Birse, 2003)
- Causal wave propagation requires $(2L + d \geq 5)$
(HWH, Lee, Phys. Lett. B **681**, 500 (2009))

$$r_{L,d} \leq -2\Gamma(L + \frac{d}{2} - 2)\Gamma(L + \frac{d}{2} - 1)/\pi \times (2\Lambda)^{2L+d-4}$$

- $d = 3, L = 0 \implies r_{0,3} \leq 2/\Lambda$

(Phillips, Cohen, Phys. Lett. B **390**, 7 (1997))

- $d = 3, L \geq 1 \implies r_{0,3} \leq -|\text{const.}| \Lambda^{2L-1}$

\implies effective range cannot be tuned to zero

- Many new $c\bar{c}$ -mesons at B-factories: X, Y, Z
 - Challenge for understanding of QCD
 - Large scattering length physics important
- Example: $X(3872)$ (Belle, CDF, BaBar, D0)

$$m_X = (3871.55 \pm 0.20) \text{ MeV} \quad \Gamma < 2.3 \text{ MeV} \quad J^{PC} = 1^{++}$$

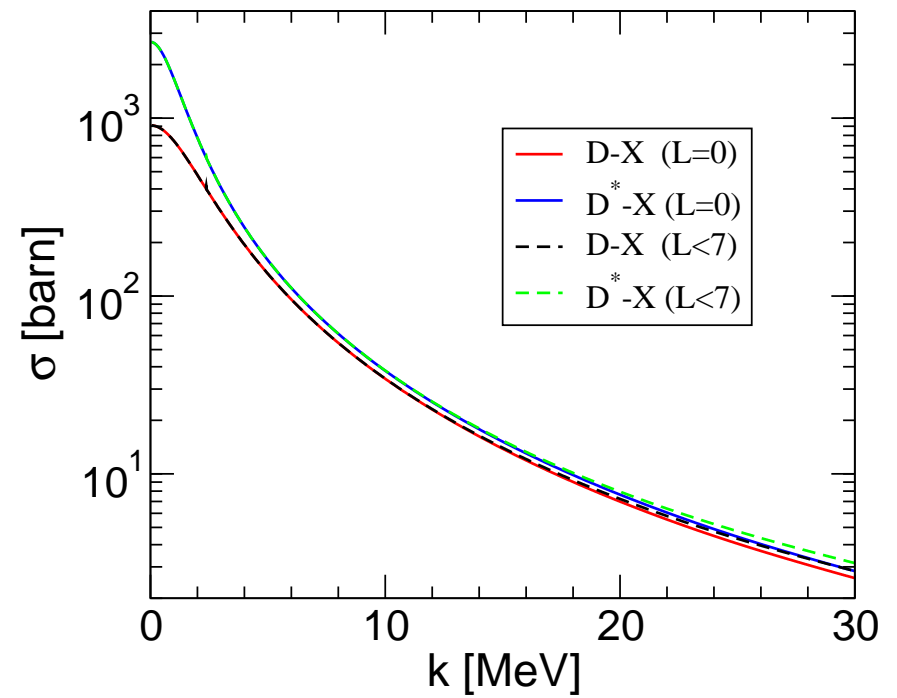
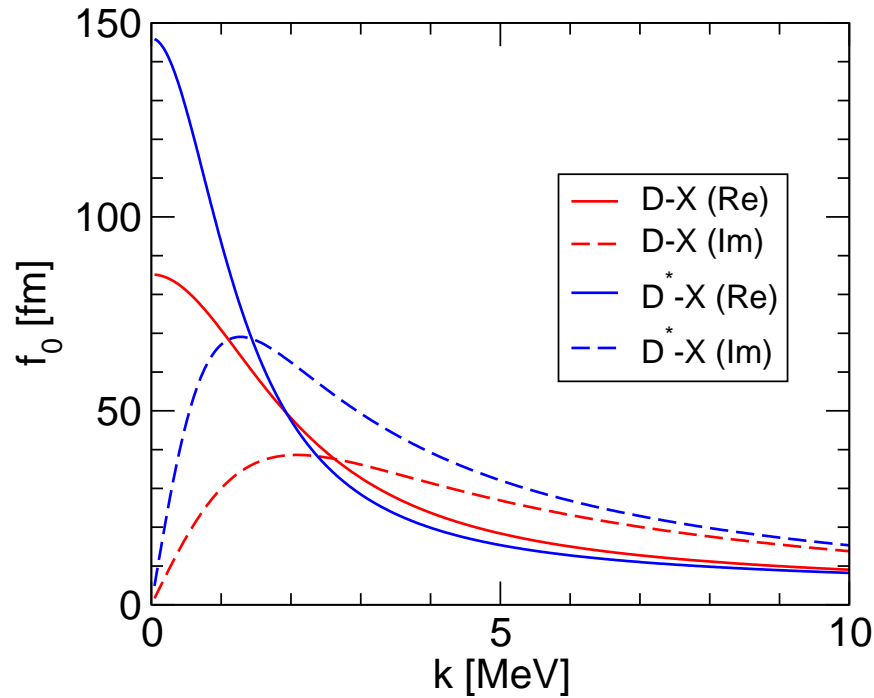
- No ordinary $c\bar{c}$ -state
 - Decays violate isospin
 - Measured mass depends on decay channel
- Nature of $X(3872)$?
 - $D^0 D^{0*}$ -molecule? (cf. Tornquist, 1991)
 - Tetraquark
 - Charmonium Hybrid
 - ...

- Nature of $X(3872)$ not finally resolved
- Assumption: $X(3872)$ is weakly-bound D^0 - \bar{D}^{0*} -molecule
 - $\implies |X\rangle = (|D^0\bar{D}^{0*}\rangle + |\bar{D}^0D^{0*}\rangle)/\sqrt{2}$, $B_X = (0.26 \pm 0.41)$ MeV
 - \implies **universal properties** (cf. Braaten et al., 2003-2008, ...)
 - Explains isospin violation in decays of $X(3872) \Rightarrow$ superposition of $I = 1$ and $I = 0$
 - Different masses due to different line shapes in decay channels
- EFT with explicit pions: short distance contributions dominate (Fleming, Kusunoki, Mehen, van Kolck, 2007)
 - \implies EFT for large scattering length is applicable
- Large scattering length determines interaction of $X(3872)$ with D^0 and D^{0*}

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- Efimov effect?
 - ⇒ occurs if 2 out of 3 pairs have resonant interactions
- $X(3872)$: only 3 out of 6 pairs have resonant interactions
 - ⇒ **no Efimov effect** (Braaten, Kusunoki, 2003)
 - ⇒ no X - D^0 - and X - D^{0*} -molecules
 - ⇒ no three-body interaction at leading order

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- **But:** parameter-free prediction of X - D^0 -, X - D^{0*} -scattering
- **Low-energy parameters:** $B_X = (0.26 \pm 0.41)$ MeV
 - ⇒ Scattering length in the X channel: $a = (8.8_{-3.3}^{+\infty})$ fm

- Predictions for scattering amplitude/cross section

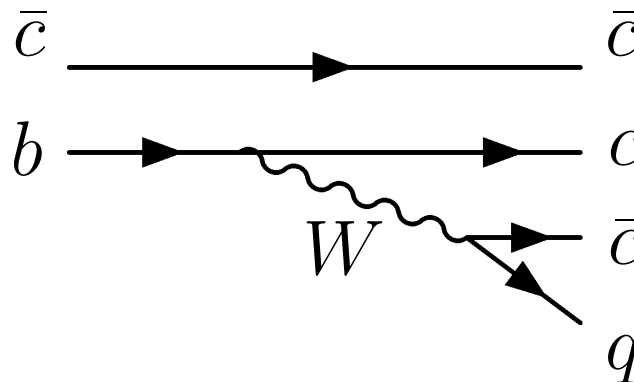


Canham, HWH, Springer, Phys. Rev. D **80**, 014009 (2009)

- Three-body scattering lengths

$$a_{D^0 X} = a_{\bar{D}^0 X} = -9.7a, \quad \text{and} \quad a_{D^{*0} X} = a_{\bar{D}^{*0} X} = -16.6a$$

- Behavior of $X(3872)$ produced in isolation should be distinguishable from its behavior when in the presence of $D^0, D^{*0}, \bar{D}^0, \bar{D}^{*0}$
- Rare events in $B\bar{B}$ production ($B \rightarrow X, \bar{B} \rightarrow D, D^*$)
- Final state interaction of D, D^* mesons in B_c -decays
- Example: quark-level B_c decay yielding three charmed/anticharmed quarks in final state



- Process may be accessible at the LHC

- Effective field theory for large scattering length
 - Limit cycle in three-body system \Leftrightarrow Efimov physics
 - Universal correlations (Phillips, Tjon line,...)
- Applications in atomic, nuclear, and particle physics
 - Cold atoms close to Feshbach resonance
 - Halo nuclei
 - Scattering properties of the $X(3872)$
- Future directions:
 - **Cold atoms:** heteronuclear systems, $N \geq 4$, 2d-systems, ...
 - **Halo nuclei:** reactions, external currents, ...
 - **Hadronic molecules:** universal properties, three-body molecules?
 - **Three-nucleon system on the lattice:** finite volume corrections, limit cycle in “deformed” QCD?