

Momentum-space techniques for reactions involving weakly bound nuclei

A. Deltuva

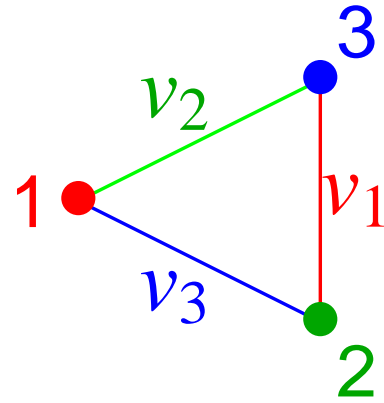
Centro de Física Nuclear da Universidade de Lisboa

Few-body scattering

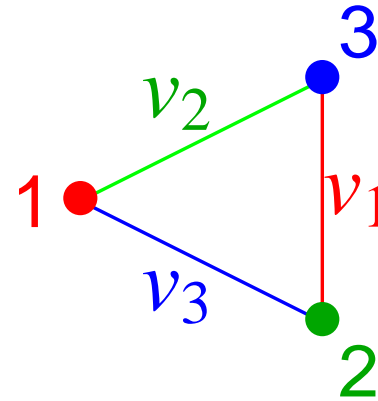
- Three-body scattering equations
- Inclusion of the Coulomb interaction
- Three-nucleon system
- Three-body direct nuclear reactions
- Four-nucleon scattering
- Extension to atomic physics

Three-body system

Hamiltonian $H_0 + \sum_{\alpha} v_{\alpha}$



Three-body system



Hamiltonian $H_0 + \sum_{\alpha} v_{\alpha}$

- Faddeev equations

$$(E - H_0 - v_{\alpha}) |\Psi_{\alpha}\rangle = v_{\alpha} \sum_{\sigma} \bar{\delta}_{\alpha\sigma} |\Psi_{\sigma}\rangle$$

$$|\Psi\rangle = \sum_{\alpha} |\Psi_{\alpha}\rangle$$

Alt-Grassberger-Sandhas equations

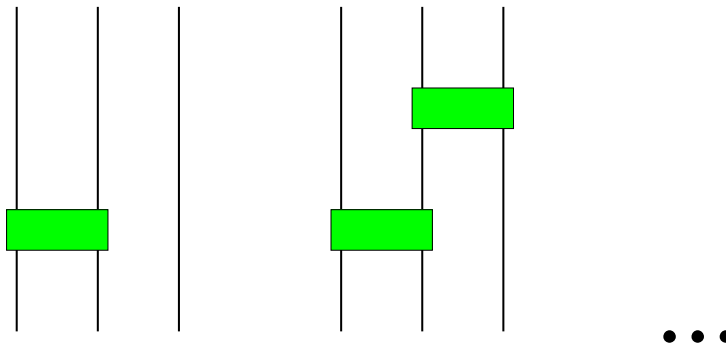
$$U_{\beta\alpha} = \bar{\delta}_{\beta\alpha} G_0^{-1} + \sum_{\sigma} \bar{\delta}_{\beta\sigma} T_{\sigma} G_0 U_{\sigma\alpha}$$

$$U_{0\alpha} = G_0^{-1} + \sum_{\sigma} T_{\sigma} G_0 U_{\sigma\alpha}$$

$$T_{\sigma} = v_{\sigma} + v_{\sigma} G_0 T_{\sigma}$$

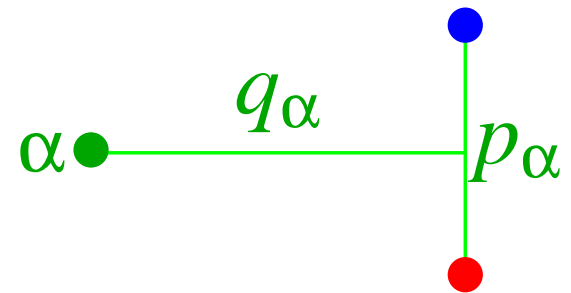
$$G_0 = (E + i0 - H_0)^{-1}$$

channel states $(E - H_0 - v_{\alpha})|\phi_{\alpha}\rangle = 0$



AGS equations: numerical solution

$$U_{\beta\alpha} = \bar{\delta}_{\beta\alpha} G_0^{-1} + \sum_{\sigma} \bar{\delta}_{\beta\sigma} T_{\sigma} G_0 U_{\sigma\alpha}$$



- 3 sets of Jacobi momenta
- momentum-space partial wave basis
- set of coupled 2-variable integral equations
- integrable singularities in kernel
- Gaussian integration, spline interpolation, Padé summation

Inclusion of Coulomb: screening

$$w_R(r) = \frac{\alpha_e}{r} e^{-\left(\frac{r}{R}\right)^n}$$

- standard scattering theory

$$v_\sigma \rightarrow v_\sigma + w_{\sigma R} : \quad T_\sigma, U_{\beta\alpha} \rightarrow T_\sigma^{(R)}, U_{\beta\alpha}^{(R)}$$

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- nature: Coulomb is screened at large distances

- large R :

physical observables insensitive to screening,
screened and full Coulomb physically indistinguishable

Screening and renormalization

J. R. Taylor, Nuovo Cimento **B23**, 313 (1974),
V. G. Gorshkov, Sov. Phys.-JETP **13**, 1037 (1961):

Screening and renormalization

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Renormalization of the on-shell screened Coulomb transition matrix $T_R = w_R + w_R G_0 T_R$ and wave function

$$T_R z_R^{-1} \xrightarrow{R \rightarrow \infty} T_C \quad \text{as distribution}$$

$$(1 + G_0 T_R) |\mathbf{p}\rangle z_R^{-\frac{1}{2}} \xrightarrow{R \rightarrow \infty} |\Psi_C^{(+)}(\mathbf{p})\rangle$$

in the limit $R \rightarrow \infty$ yields **Coulomb amplitude**
and **Coulomb wave function**

$$z_R \xrightarrow{R \rightarrow \infty} \exp(-2i(\sigma_L - \eta_{LR})) \xrightarrow{R \rightarrow \infty} \exp(-2i\alpha_e M/p [\ln(2pR) - C/n])$$

Two-particle scattering

transition matrix

$$T^{(R)} = v + w_R + (v + w_R)G_0T^{(R)}$$

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with long-range and Coulomb-distorted short-range parts

$$T^{(R)} = T_R + (1 + T_R G_0) \tilde{T}^{(R)} (1 + G_0 T_R)$$
$$\tilde{T}^{(R)} = v + v G_R \tilde{T}^{(R)}$$

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Renormalized amplitude:

$$T^{(R)} z_R^{-1} \xrightarrow{R \rightarrow \infty} T = T_C + \langle \psi_C^{(-)} | \tilde{T}^{(C)} | \psi_C^{(+)} \rangle$$

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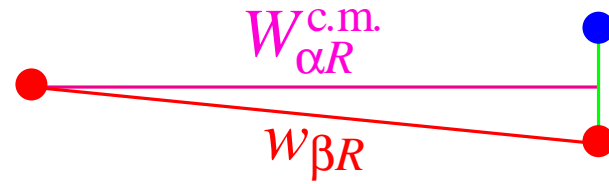
$$T^{(R)} z_R^{-1} \xrightarrow{R \rightarrow \infty} T = T_C + \langle \Psi_C^{(-)} | \tilde{T}^{(C)} | \Psi_C^{(+)} \rangle$$
$$= T_C + \lim_{R \rightarrow \infty} z_R^{-\frac{1}{2}} [T^{(R)} - T_R] z_R^{-\frac{1}{2}}$$

short-range part: fast convergence with R

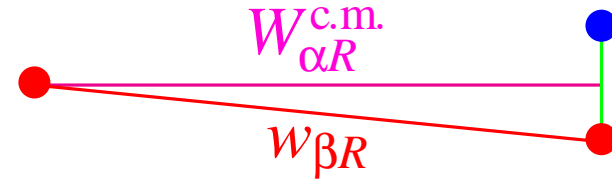
Three-particle scattering

long-range part

$$T_{\alpha R}^{\text{c.m.}} = W_{\alpha R}^{\text{c.m.}} + W_{\alpha R}^{\text{c.m.}} G_{\alpha}^{(R)} T_{\alpha R}^{\text{c.m.}}$$



Three-particle scattering



Split into **long-range** part

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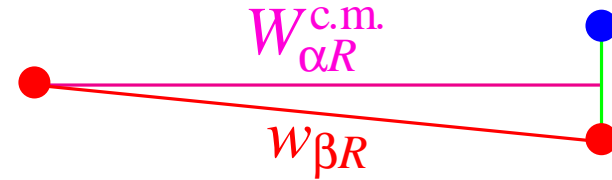
and **Coulomb-distorted short-range** part

$$U_{\beta\alpha}^{(R)} = \delta_{\beta\alpha} T_{\alpha R}^{\text{c.m.}} + [U_{\beta\alpha}^{(R)} - \delta_{\beta\alpha} T_{\alpha R}^{\text{c.m.}}]$$

$$[U_{\beta\alpha}^{(R)} - \delta_{\beta\alpha} T_{\alpha R}^{\text{c.m.}}] = [1 + T_{\beta R}^{\text{c.m.}} G_{\beta}^{(R)}] \tilde{U}_{\beta\alpha}^{(R)} [1 + G_{\alpha}^{(R)} T_{\alpha R}^{\text{c.m.}}]$$

$$U_{0\alpha}^{(R)} = [1 + T_{\rho R} G_0] \tilde{U}_{0\alpha}^{(R)} [1 + G_{\alpha}^{(R)} T_{\alpha R}^{\text{c.m.}}] \quad [\rho \text{ is neutral}]$$

Three-particle scattering



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Renormalized amplitudes:

$$U_{\beta\alpha} = \delta_{\beta\alpha} T_{\alpha C}^{\text{c.m.}} + \lim_{R \rightarrow \infty} Z_{Rf}^{-\frac{1}{2}} [U_{\beta\alpha}^{(R)} - \delta_{\beta\alpha} T_{\alpha R}^{\text{c.m.}}] Z_{Ri}^{-\frac{1}{2}}$$

$$U_{0\alpha} = \lim_{R \rightarrow \infty} z_R^{-\frac{1}{2}} U_{0\alpha}^{(R)} z_{Ri}^{-\frac{1}{2}} \quad \text{fast convergence with } R$$

Practical realization

- Calculation of short-range part using standard scattering theory (Faddeev/AGS) for nuclear + screened Coulomb interaction

$$U_{\beta\alpha}^{(R)} = \bar{\delta}_{\beta\alpha} G_0^{-1} + \sum_{\sigma} \bar{\delta}_{\beta\sigma} T_{\sigma}^{(R)} G_0 U_{\sigma\alpha}^{(R)}$$
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+ renormalization

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- **Additional difficulties:**
quasi-singular nature of screened Coulomb potential
slow partial-wave convergence

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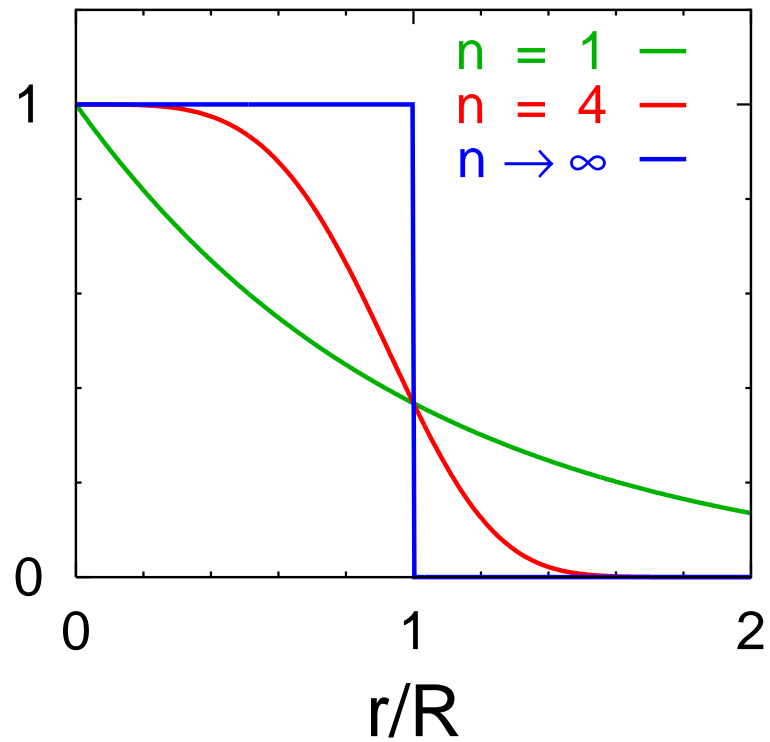
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+ renormalization

- **Additional difficulties:**
quasi-singular nature of screened Coulomb potential
slow partial-wave convergence
- Success of the method depends strongly on the choice of screening function

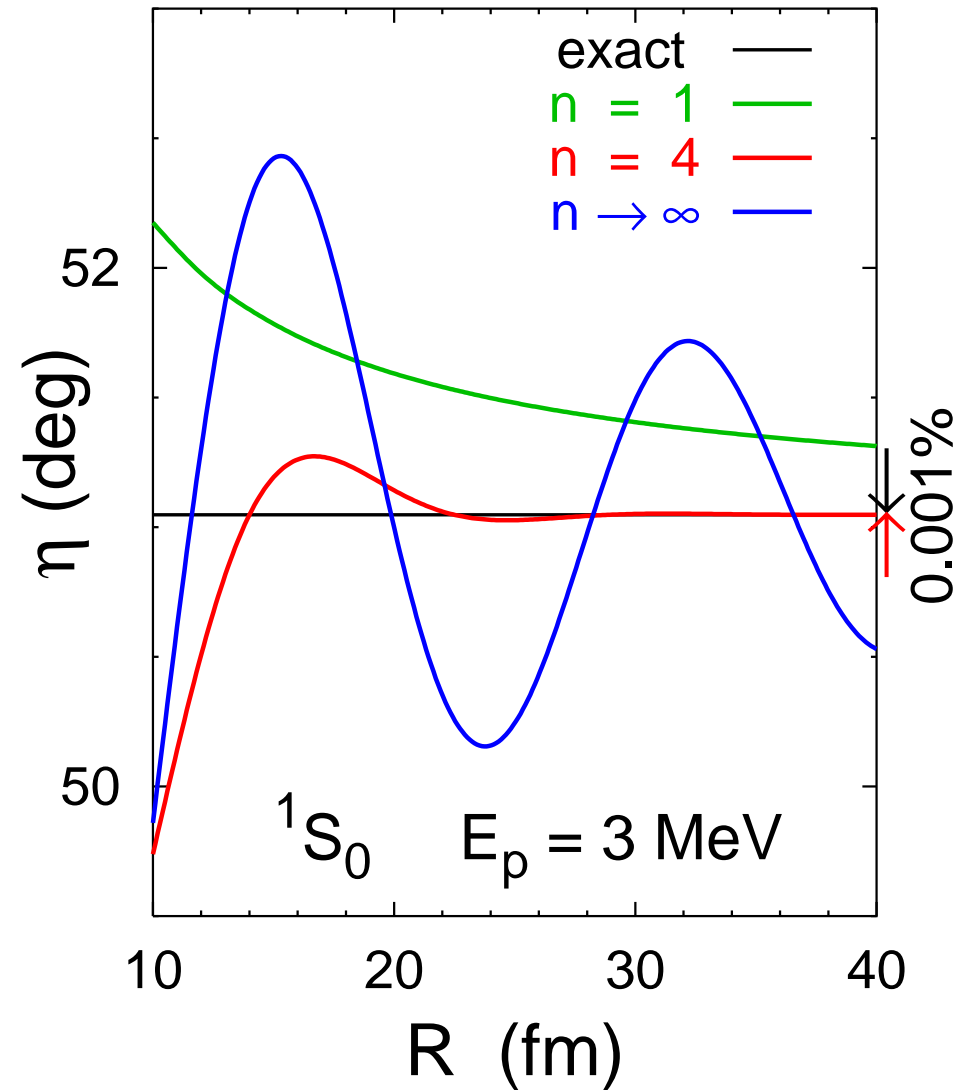
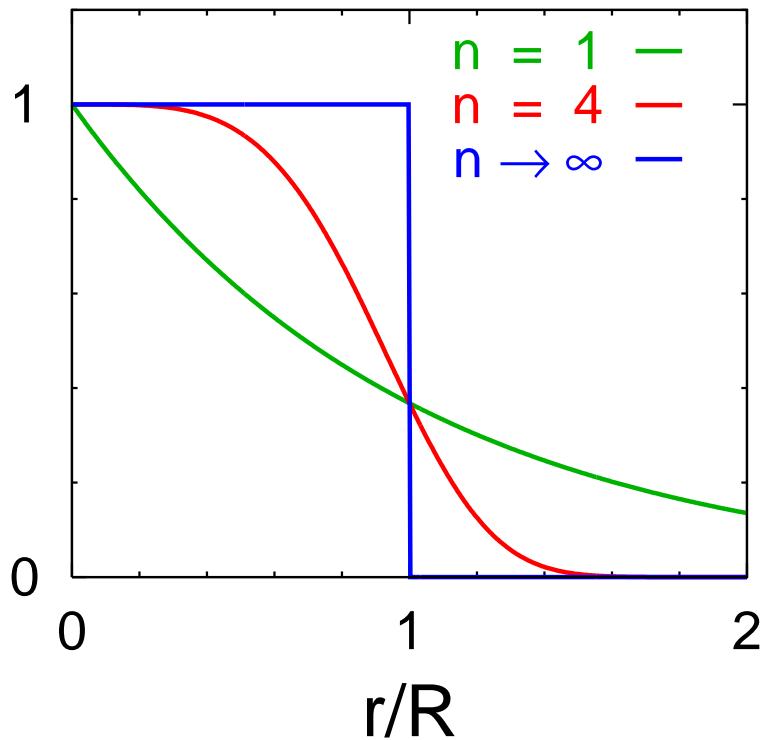
Screened Coulomb potential

$$\frac{w_R(r)}{w(r)} = e^{-\left(\frac{r}{R}\right)^n}$$



Screened Coulomb potential

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optimal choice: $3 \leq n \leq 8$

Limits of practical applicability

$p \rightarrow 0$:

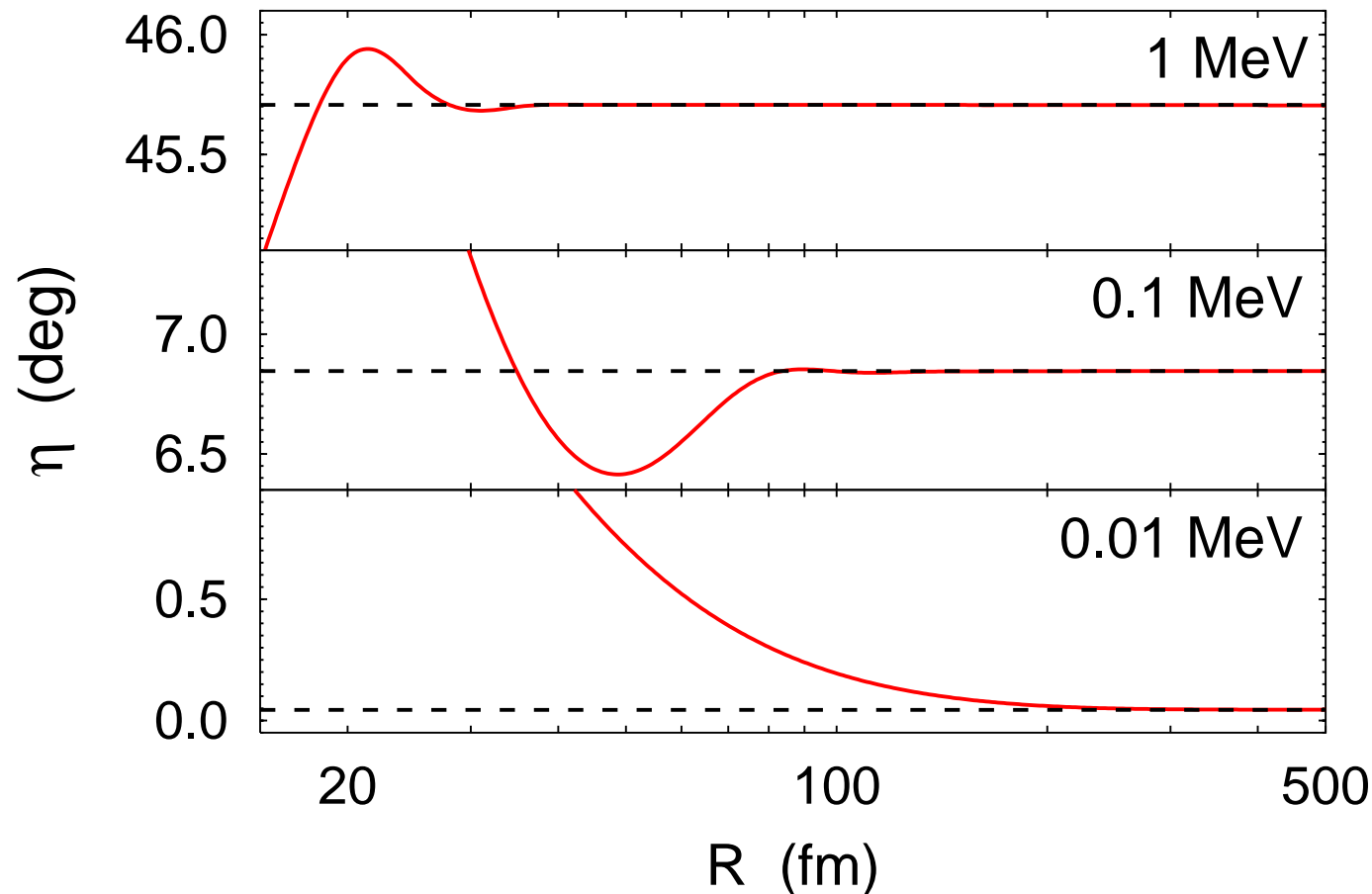
$\kappa = \alpha M/p$, $\sigma_L = \arg \Gamma(1 + L + i\kappa)$, and z_R diverge,
renormalization procedure ill-defined

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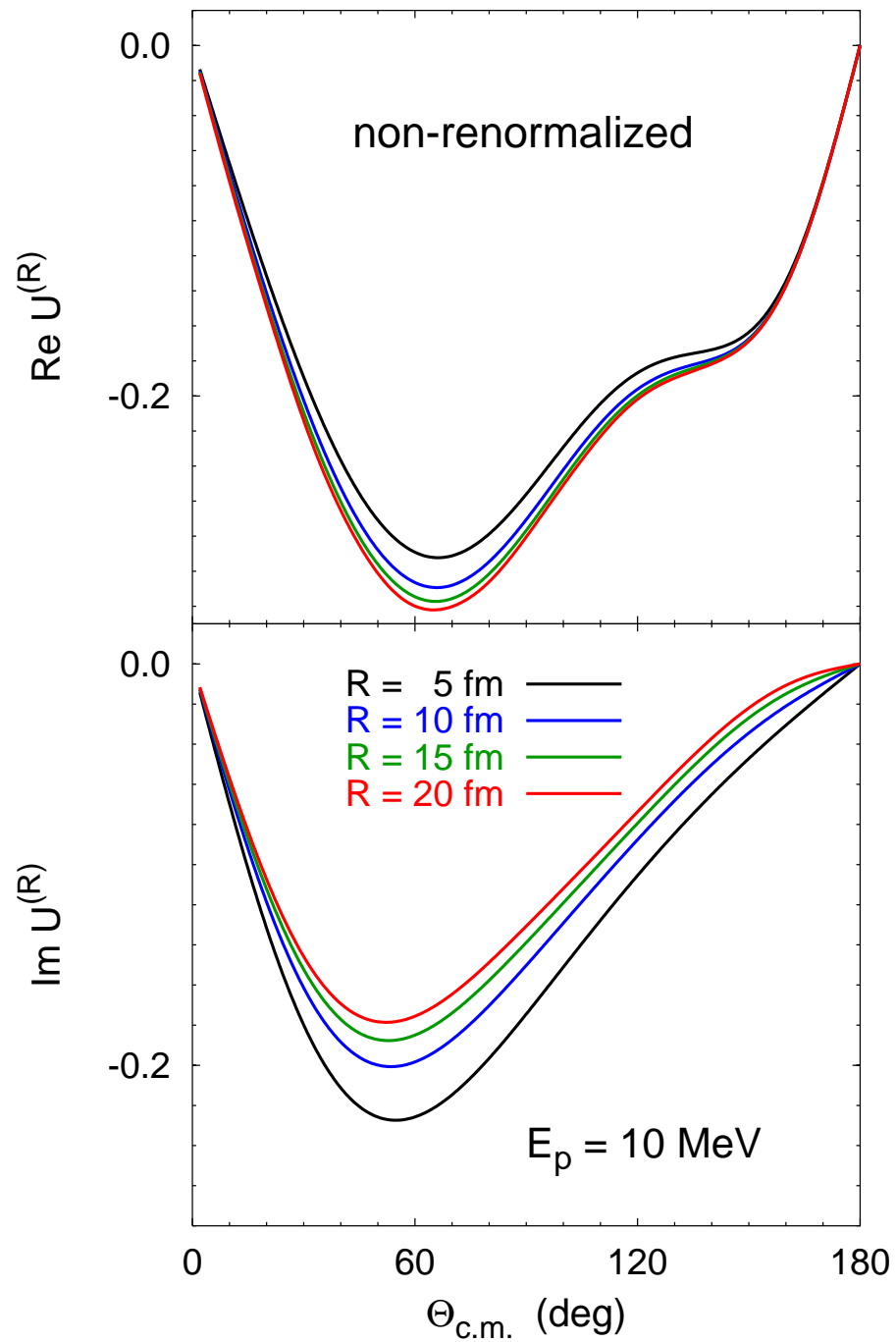
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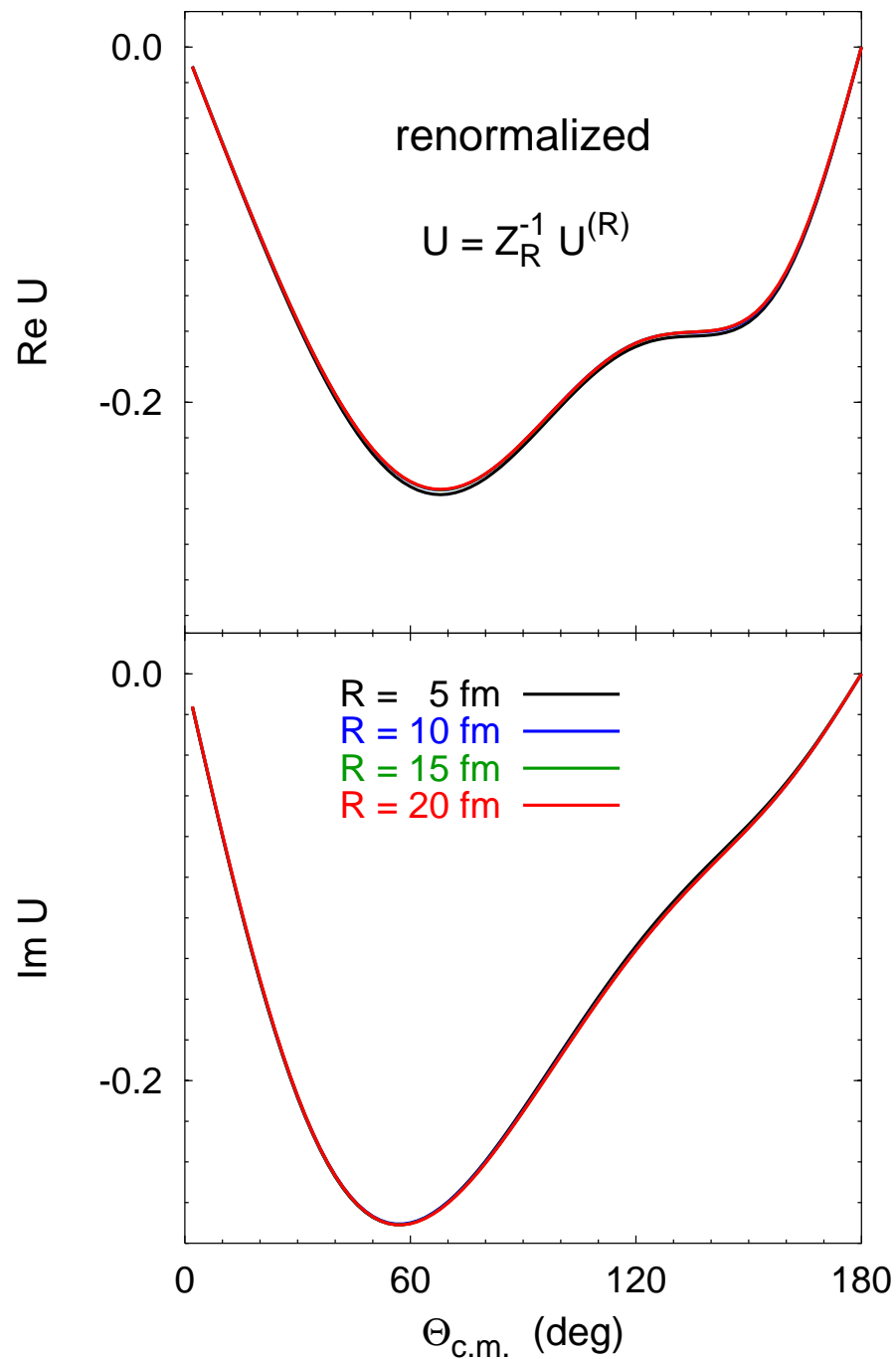
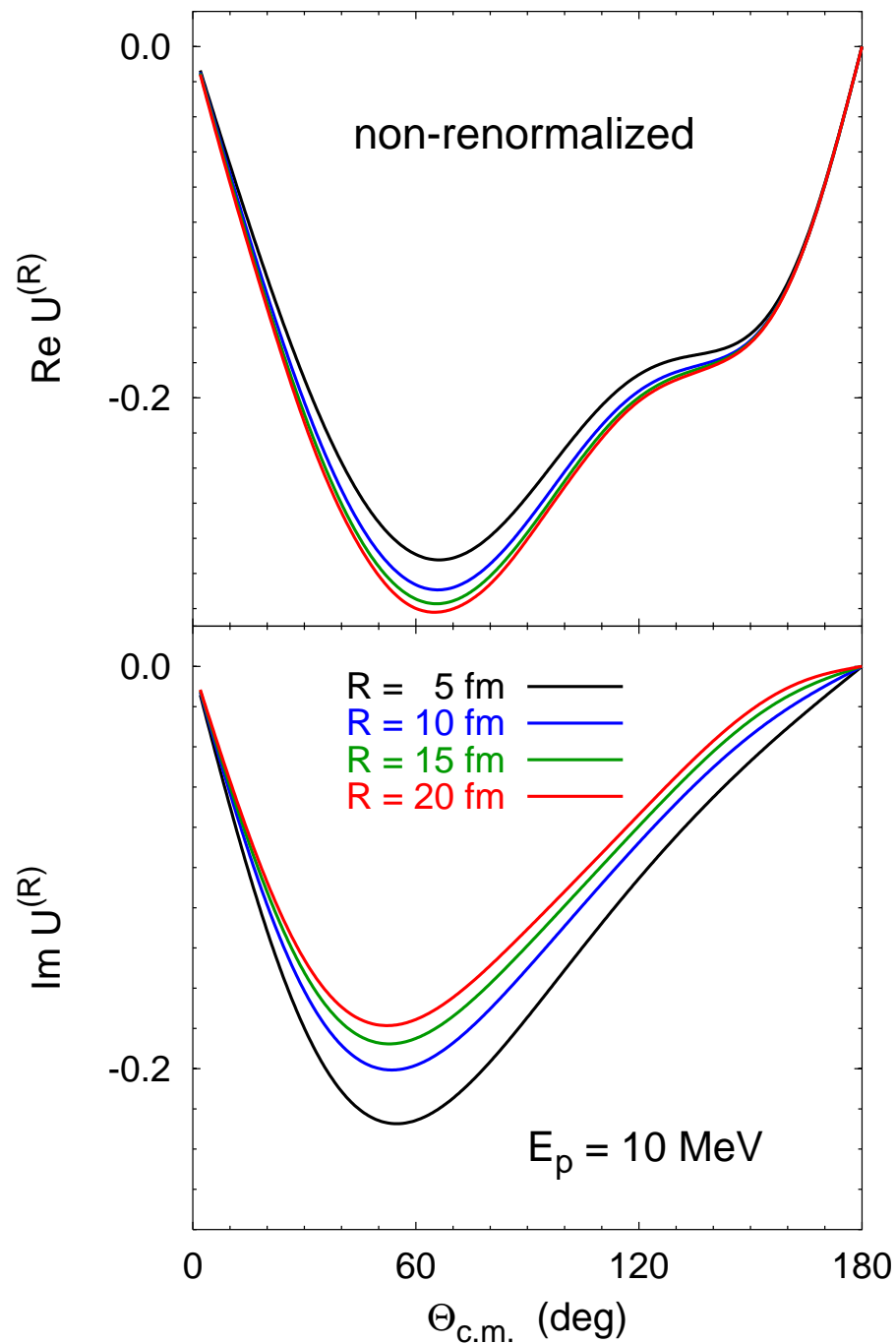
\Rightarrow slow convergence with R at low relative energies



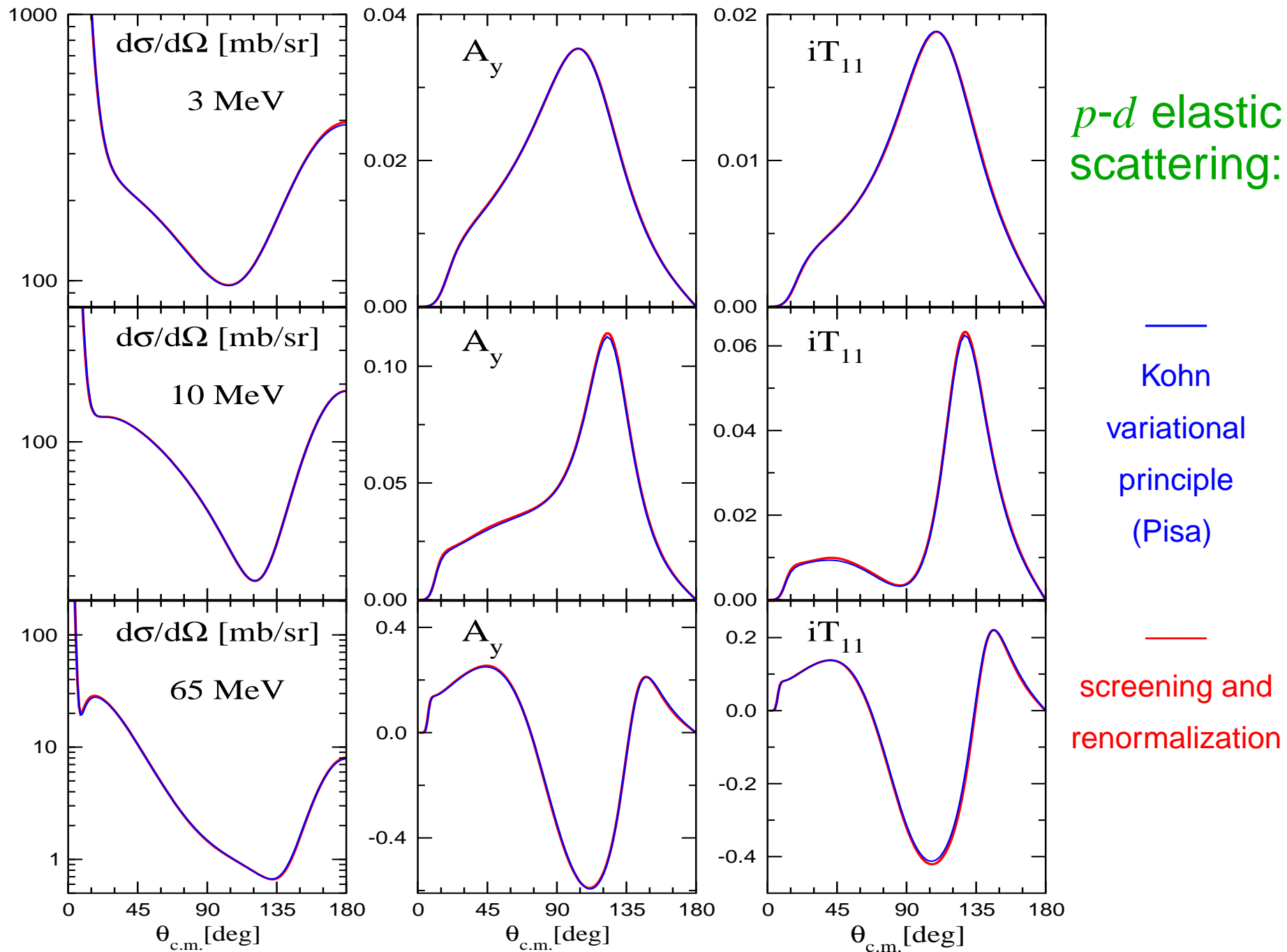
pd elastic amplitude: convergence with R



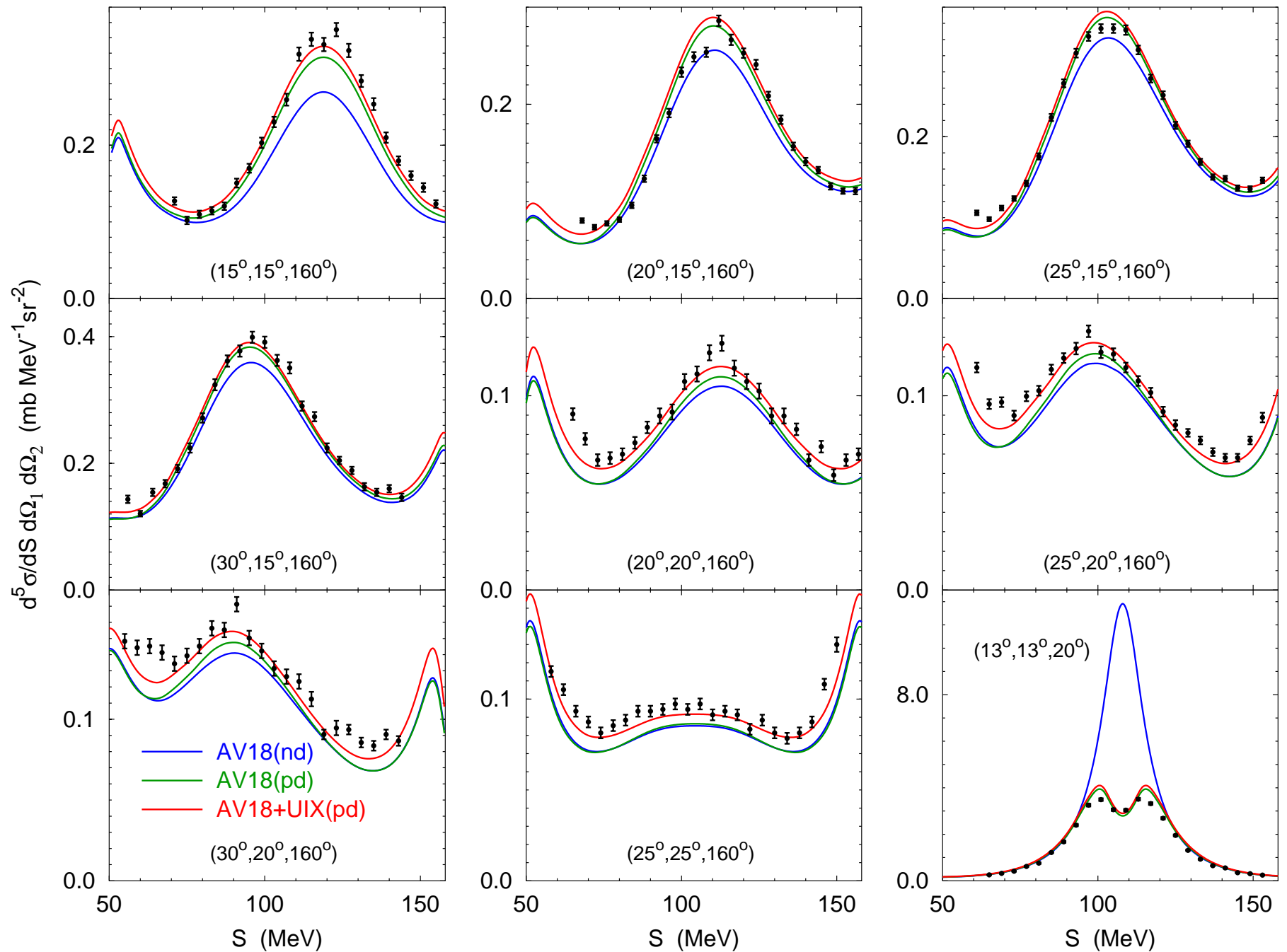
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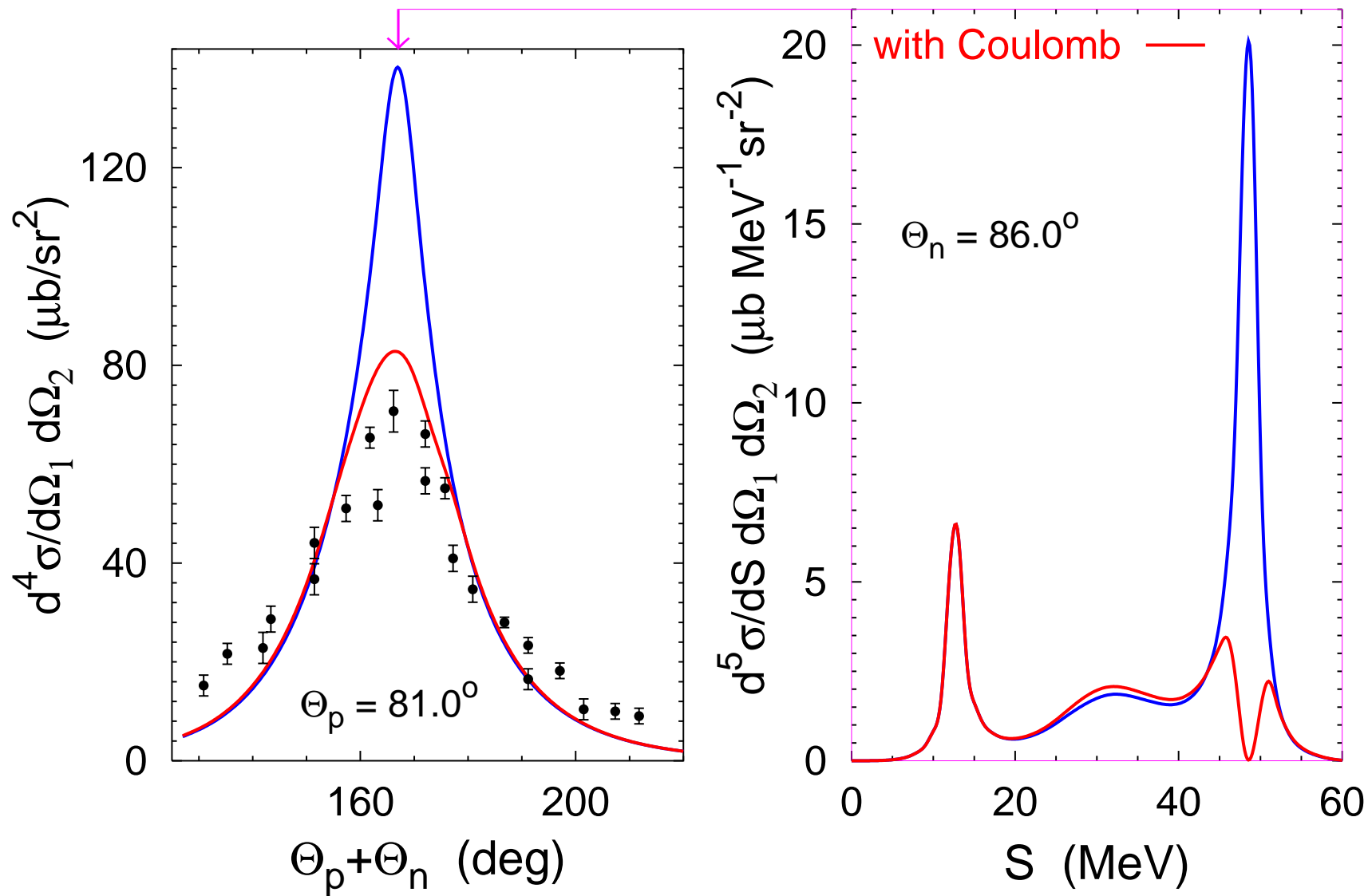
Comparison with configuration-space results



Coulomb vs 3NF: $^1\text{H}(d,pp)n$ at $E_d = 130$ MeV



${}^3\text{He}(\gamma, pn)p$ at $E_\gamma = 55$ MeV



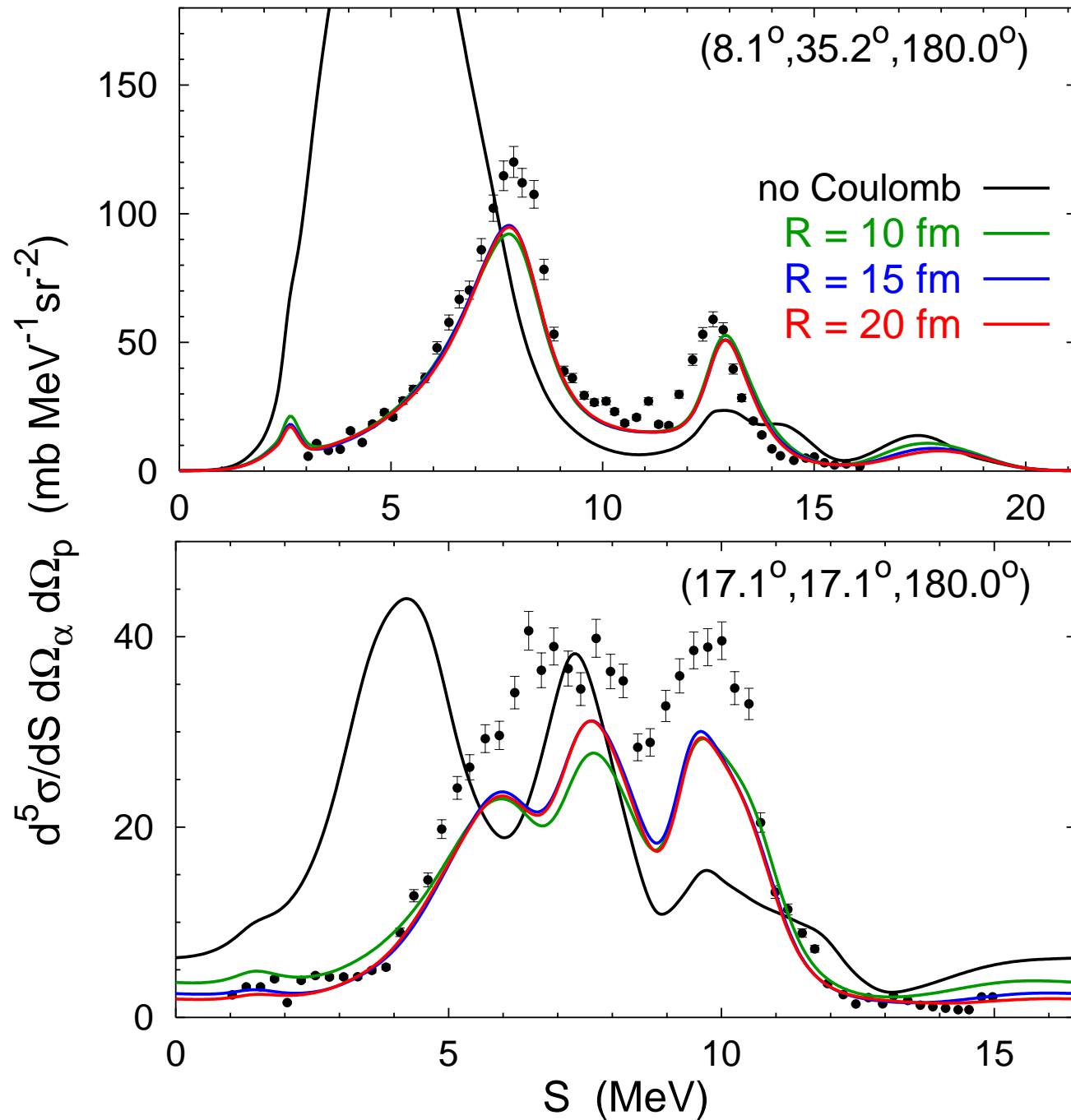
Application to 3-body nuclear reactions

$$\left. \begin{array}{l} p + (nA) \\ d + A \end{array} \right\} \rightarrow \left\{ \begin{array}{l} n + (pA) \\ p + (nA) \\ d + A \\ p + n + A \end{array} \right.$$

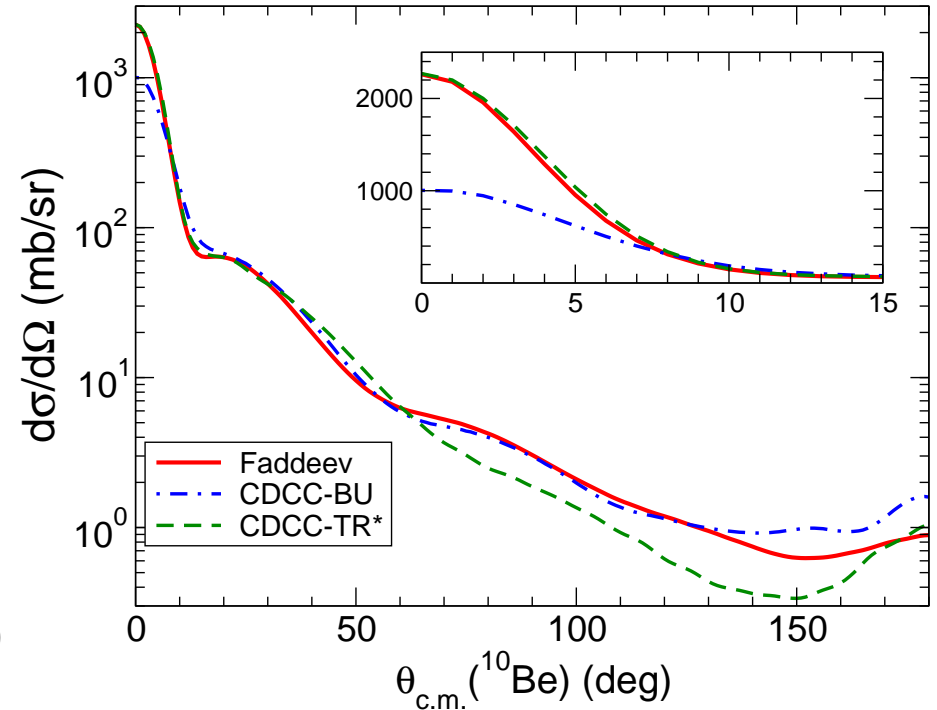
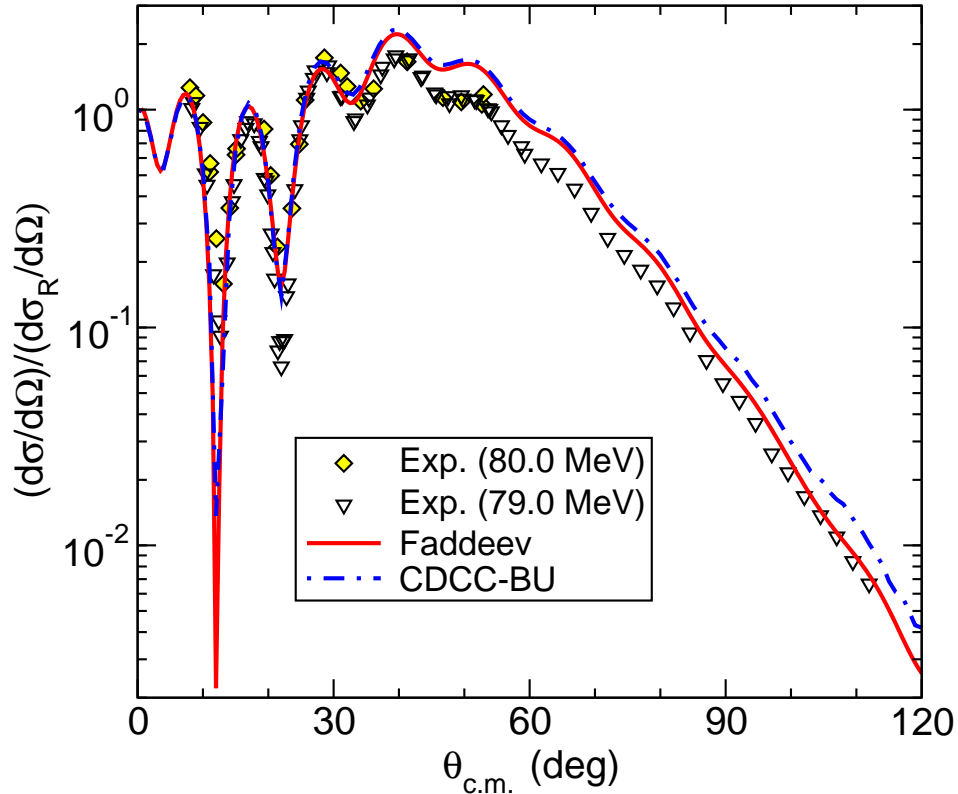
with $A = {}^4\text{He}, {}^{10}\text{Be}, {}^{12}\text{C}, {}^{14}\text{C}, {}^{16}\text{O}, {}^{28}\text{Si}, {}^{40}\text{Ca}, {}^{58}\text{Ni}, \dots$

- Validity test of approximate nuclear reaction methods: CDCC, DWBA, Glauber, ...
- Novel dynamic input: nonlocal potentials, ...

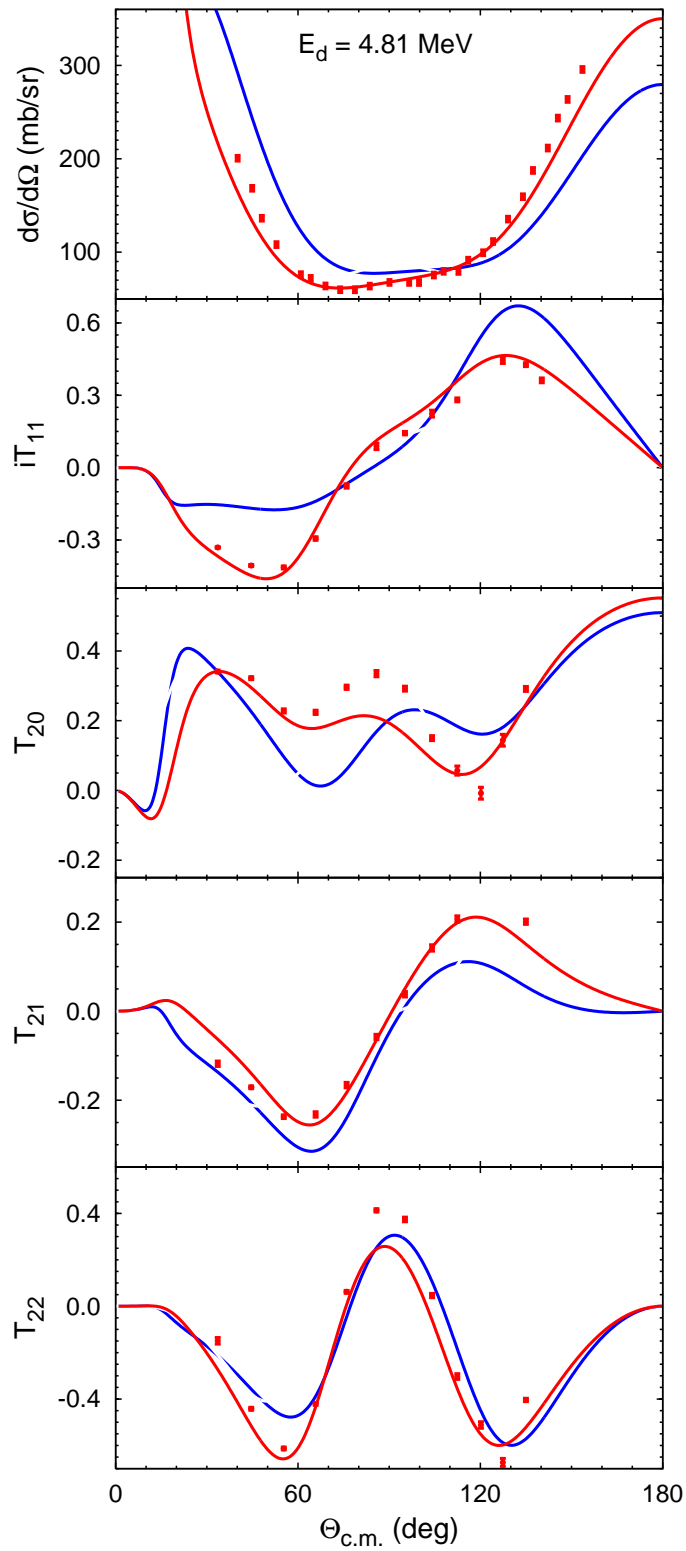
α - d breakup at $E_\alpha = 15$ MeV: convergence with R



CDCC test: $^{58}\text{Ni}(d,d)^{58}\text{Ni}$ and $^1\text{H}(^{11}\text{Be},^{10}\text{Be})np$



CDCC: A. M. Moro and F. M. Nunes



$\vec{d} + \alpha$ elastic scattering

N- α S-wave:

—

$$V_{N\alpha} \rightarrow V_{N\alpha} + |b\rangle\Gamma\langle b|$$

Pauli forbidden bound state
projected out

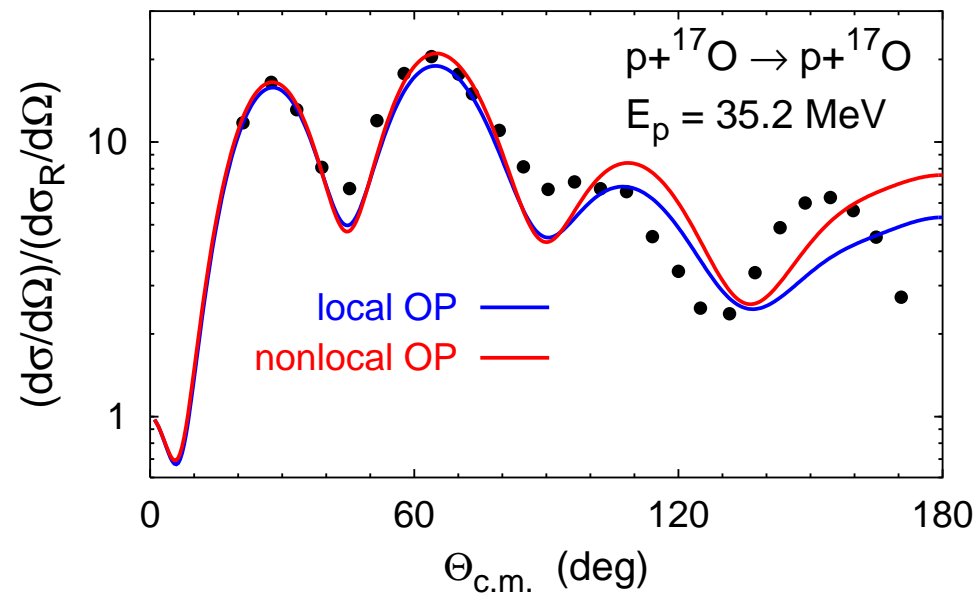
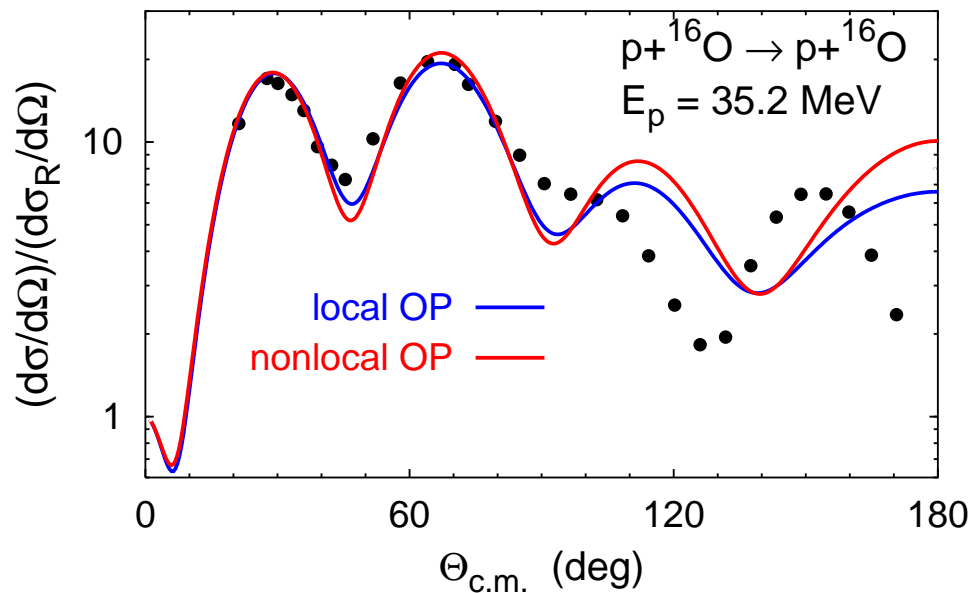
—

$V_{N\alpha}$ local repulsive

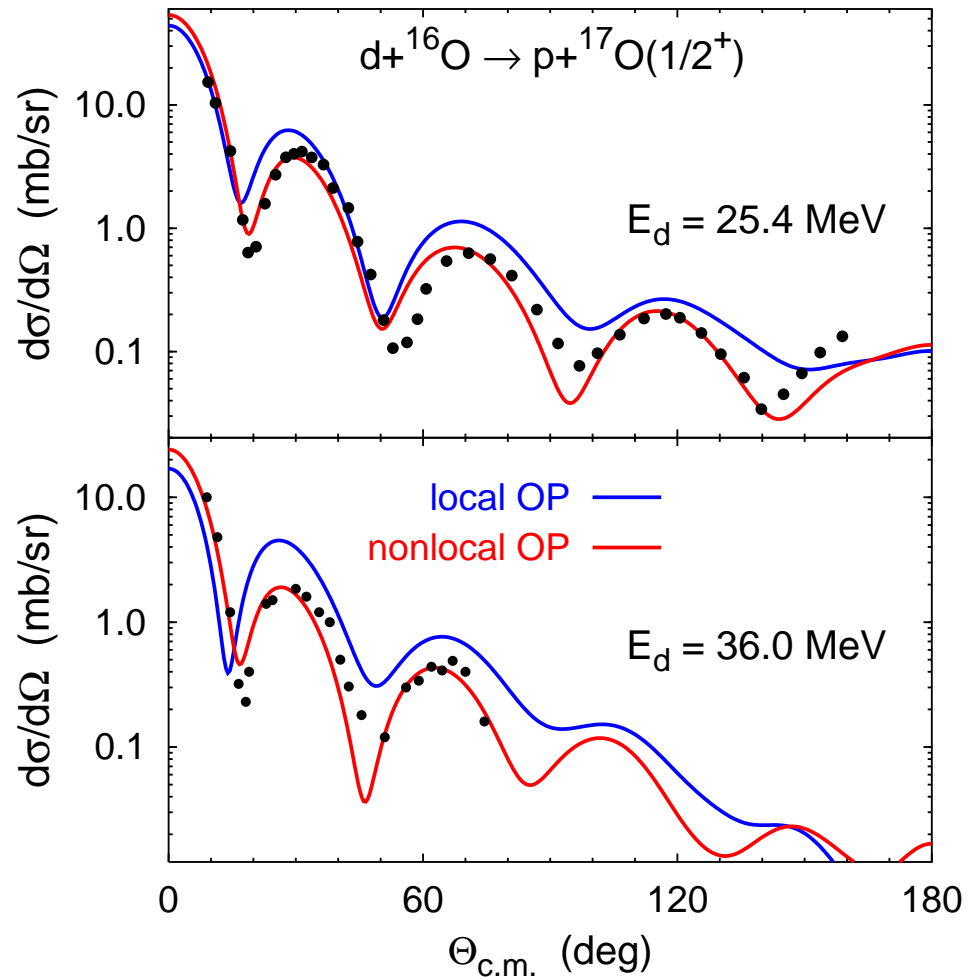
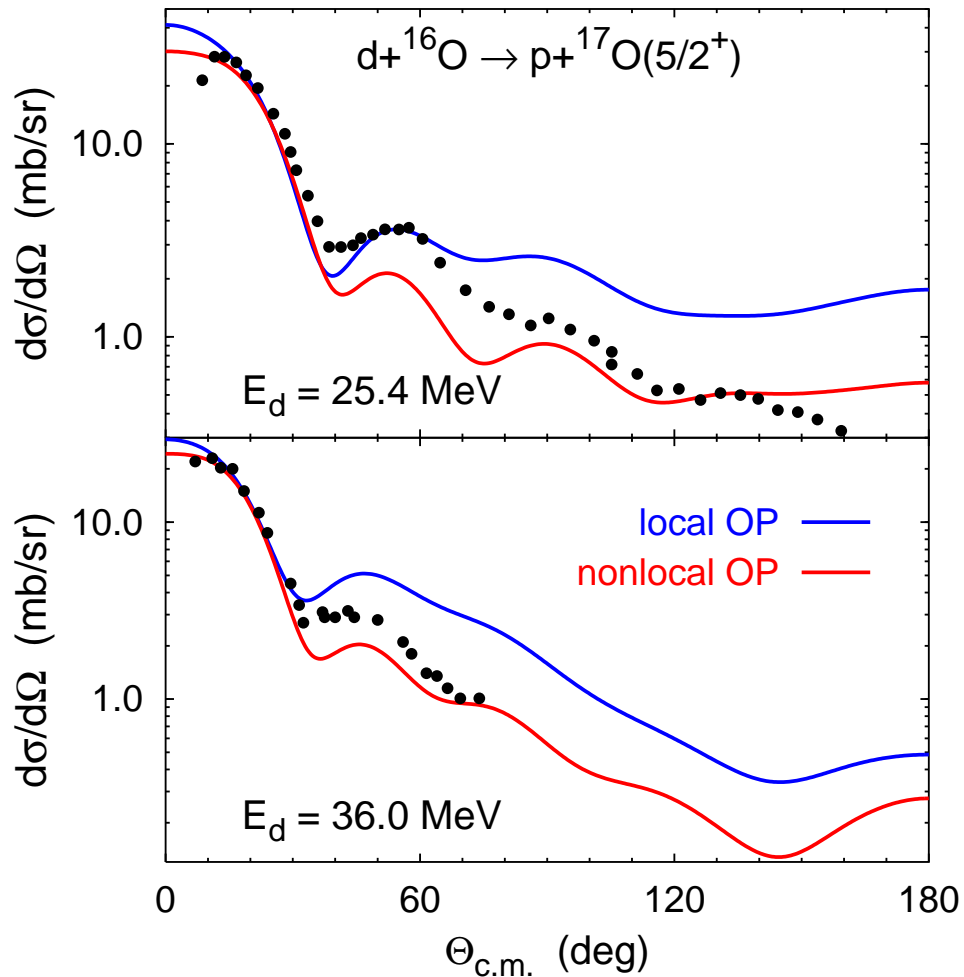
Nonlocal optical potential: proton elastic scattering

$$V_N(\mathbf{r}', \mathbf{r}) \sim e^{-(\mathbf{r}' - \mathbf{r})^2 / \beta^2} V((\mathbf{r}' + \mathbf{r}) / 2)$$

[M. M. Giannini *et al.*,
Ann. Phys. (NY) 102, 458 (1976) & 124, 208 (1980)]



Nonlocal OP: transfer reactions $^{16}\text{O}(d, p)^{17}\text{O}$



4N scattering: symmetrized AGS equations

two-cluster **1+3** and **2+2** transition operators

$$\mathcal{U}_{11} = - (G_0 T G_0)^{-1} P_{34} - P_{34} U_1 G_0 T G_0 \mathcal{U}_{11} + U_2 G_0 T G_0 \mathcal{U}_{21}$$

$$\mathcal{U}_{21} = (G_0 T G_0)^{-1} (1 - P_{34}) + (1 - P_{34}) U_1 G_0 T G_0 \mathcal{U}_{11}$$

$$\mathcal{U}_{12} = (G_0 T G_0)^{-1} - P_{34} U_1 G_0 T G_0 \mathcal{U}_{12} + U_2 G_0 T G_0 \mathcal{U}_{22}$$

$$\mathcal{U}_{22} = (1 - P_{34}) U_1 G_0 T G_0 \mathcal{U}_{12}$$

$$U_j = P_j G_0^{-1} + P_j T G_0 U_j$$

$$P_1 = P = P_{12} P_{23} + P_{13} P_{23}$$

$$P_2 = \tilde{P} = P_{13} P_{24}$$

$$T = v + v G_0 T$$

scattering amplitude $\mathcal{T}_{fi} = S_{fi} \langle \mathbf{p}_f \phi_f | \mathcal{U}_{fi} | \mathbf{p}_i \phi_i \rangle$

$$|\phi_j\rangle = G_0 T P_j |\phi_j\rangle$$

Screening and renormalization in 4N scattering

$$v \rightarrow v + w_R$$

$$T, U_j, \mathcal{U}_{fi}, \mathcal{T}_{fi} \rightarrow T^{(R)}, U_j^{(R)}, \mathcal{U}_{fi}^{(R)}, \mathcal{T}_{fi}^{(R)}$$

isolate **long-range** interaction
and Coulomb distortion between c.m. of two clusters



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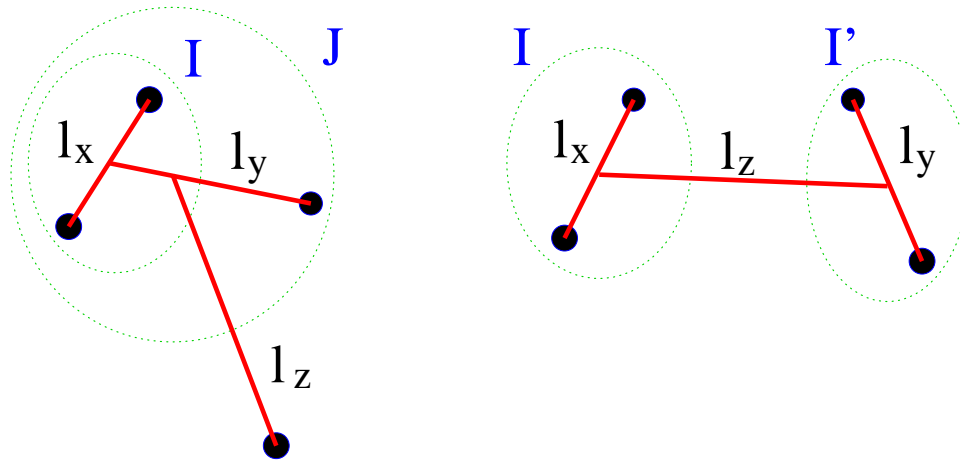
Renormalization:

$$\begin{aligned} \mathcal{T}_{fi} &= \lim_{R \rightarrow \infty} Z_{Rf}^{-\frac{1}{2}} \mathcal{T}_{fi}^{(R)} Z_{Ri}^{-\frac{1}{2}} \\ &= \delta_{fi} T_{Ci}^{\text{c.m.}} + \lim_{R \rightarrow \infty} Z_{Rf}^{-\frac{1}{2}} [\mathcal{T}_{fi}^{(R)} - \delta_{fi} T_{Ri}^{\text{c.m.}}] Z_{Ri}^{-\frac{1}{2}} \end{aligned}$$

Coulomb-distorted short-range part: fast convergence with R

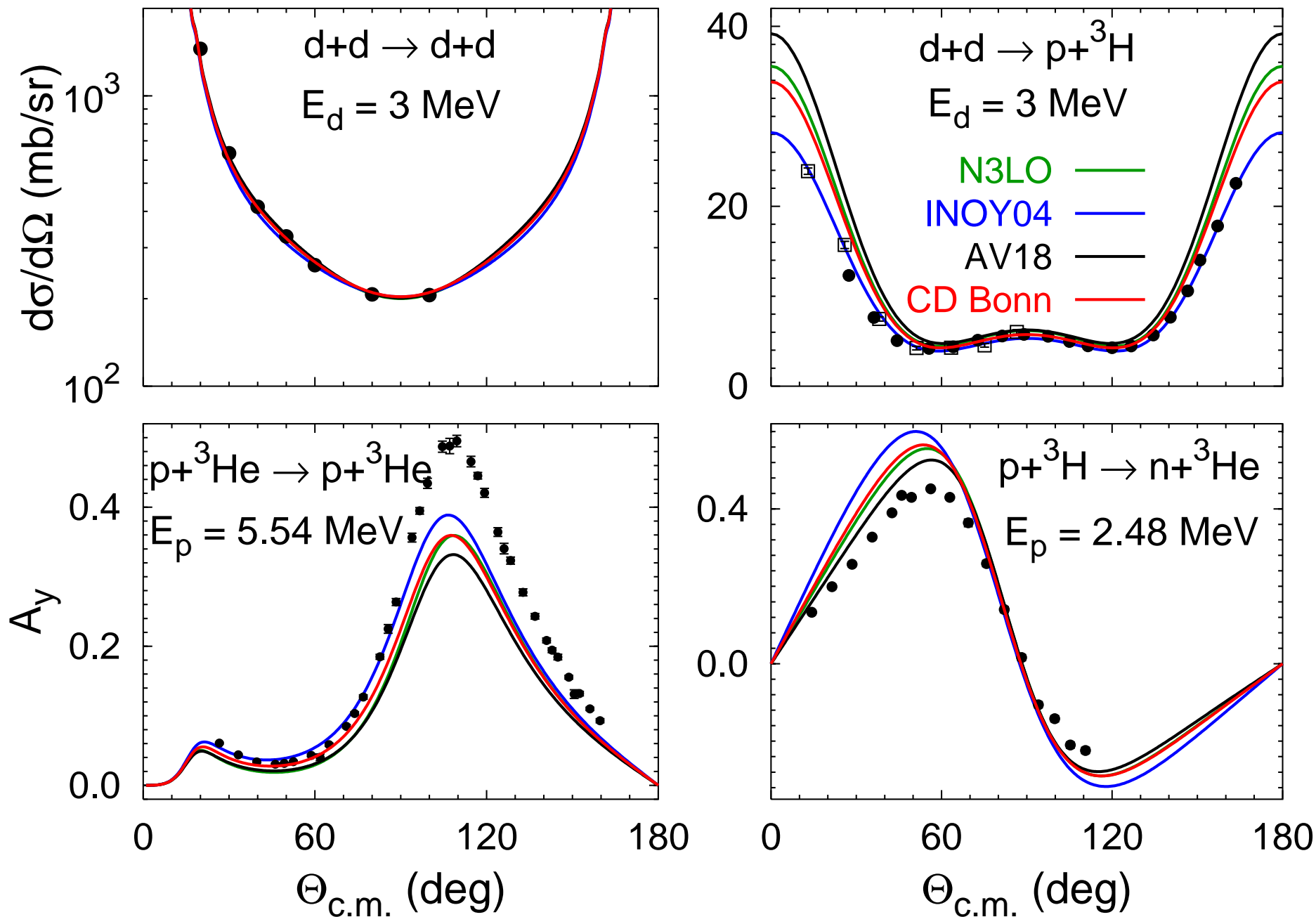
Practical realization

$$u_{12} = (G_0 T G_0)^{-1} - P_{34} U_1 G_0 T G_0 u_{12} + U_2 G_0 T G_0 u_{22}$$



- momentum-space partial-wave basis (up to 30000 partial waves)
- set of coupled integral equations in 3 variables
- integrable singularities in the kernel
- Gaussian integration, spline interpolation, double Padé summation

4N elastic, transfer, and charge exchange reactions



Atomic systems

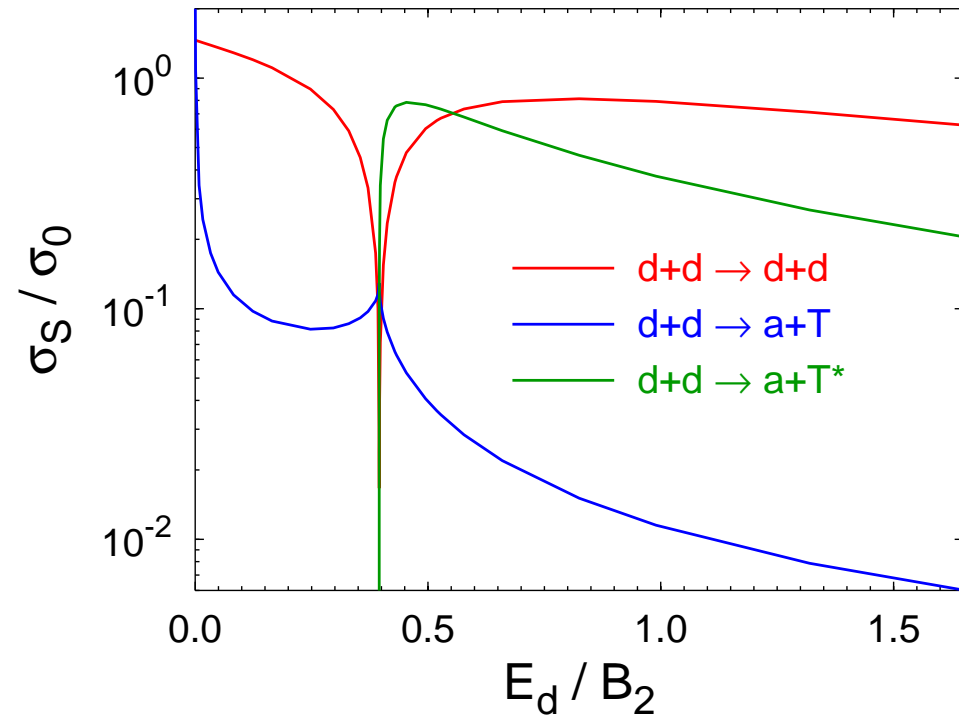
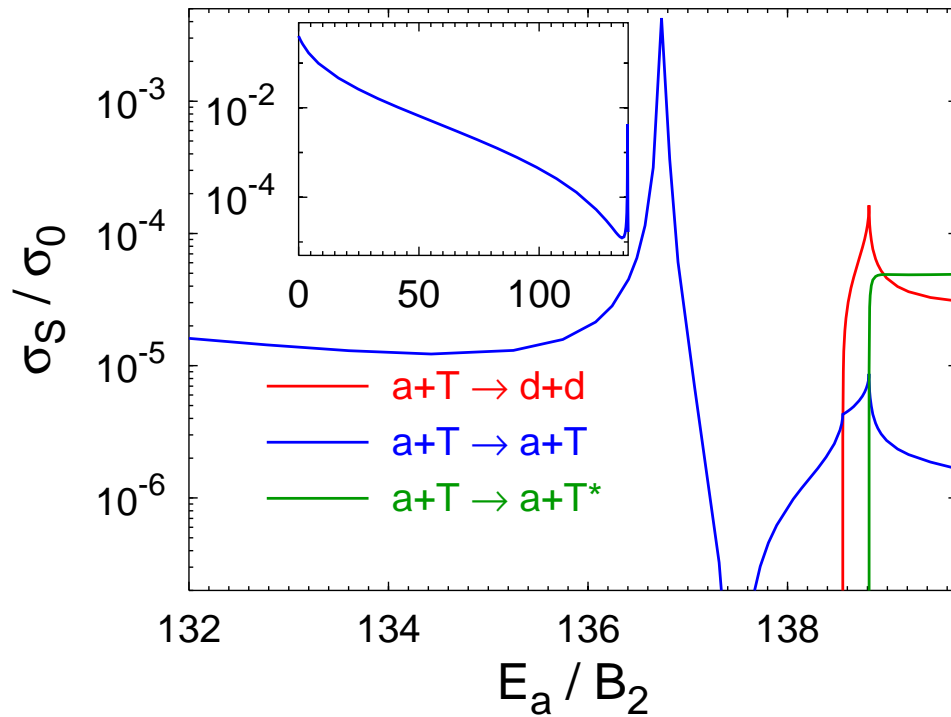
- Realistic interactions very repulsive at short distances
- Cutoff: ^4He trimer BE and atom-dimer scattering in agreement with r -space results by R. Lazauskas

Four-boson scattering

Effective separable potentials between ^4He atoms:

$$B_2 = 1.213 \text{ mK}, B_3 = 128.43 \text{ mK}, B_{3*} = 2.186 \text{ mK},$$

$$a_0 = 104 \text{ \AA}, \sigma_0 = 4\pi a_0^2$$



Summary

- Faddeev/AGS equations in momentum space
- Coulomb interaction: screening and renormalization

Summary

- Faddeev/AGS equations in momentum space
- Coulomb interaction: screening and renormalization
- hadronic and electromagnetic 3N reactions
- 3-body nuclear reactions
- low energy 4N scattering
- atomic systems