

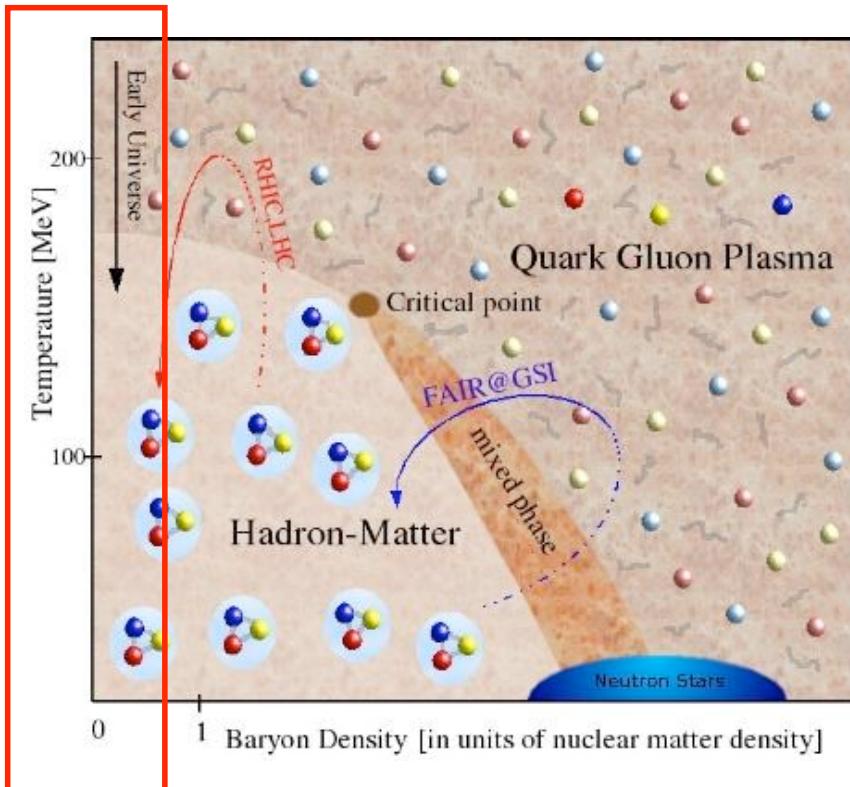
PNJL Model and QCD Phase Transitions

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Phase Transitions in Quantum Chromodynamics

This Talk



- Low Temperature

Lattice and Phenomenology say it is confined.

- High Temperature

Perturbation theory is reliable and shows deconfinement.

Want to make a phenomenological model which describe both low and high temperature QCD behavior in a single picture.

→ PNJL

Outline

- Confinement and Chiral Symmetry Breaking
- Nambu-Jona Lasinio Model (NJL)
- Polyakov-Nambu-Jona Lasinio Model (PNJL)
 - Fundamental Fermions (Finite-Temperature QCD)
 - Adjoint Fermions
 - Antiperiodic Boundary Conditions (ABC)
 - Periodic Boundary Conditions (PBC)

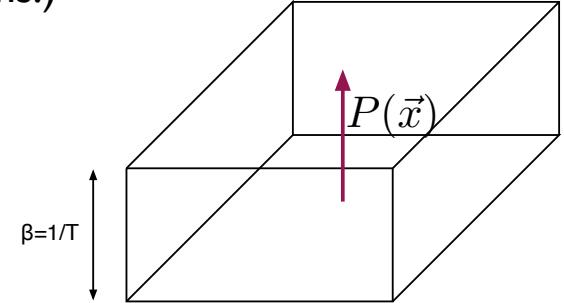
Confinement

Consider pure gauge theory (i.e. only gauge bosons and no fermions.)

Order Parameter: Polyakov Loop

$P(\vec{x}) = \mathcal{P} e^{i \int_0^\beta dt A_0(x)} \rightarrow$ The world line of massive quark at \vec{x} .

$$\langle \text{Tr} P(\vec{x}) \rangle = e^{-\beta F_q}$$



$$g \in SU(N_c) \quad z \in Z(N_c): P(\vec{x}) \rightarrow zP(\vec{x})$$

$$\langle \text{Tr} P(\vec{x}) \rangle = 0 \longrightarrow F_q = \infty \quad \text{Confinement}$$

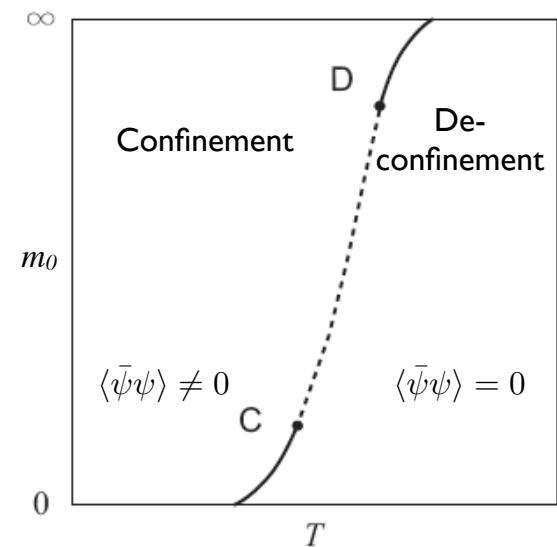
$$\langle \text{Tr} P(\vec{x}) \rangle \neq 0 \longrightarrow F_q = \text{Finite} \quad \text{Deconfinement}$$

χSB

Condensate of bound quark-antiquark pairs,
almost like Cooper pairs in BCS.

Order Parameter: $\langle \bar{\psi} \psi \rangle$

Dynamically generated quark mass (constituent mass)



<K. Fukushima, PRD69, 2004>

NJL Model (1/2): Effective Lagrangian

How does low energy QCD lagrangian look like?

$$\mathcal{L}_{QCD} = \bar{\psi}(i\gamma^\mu \partial_\mu - m_0)\psi - \frac{1}{4}F_{\mu\nu}^a F_a^{\mu\nu} + g\bar{\psi}\gamma^\mu A_\mu\psi$$

- ↓
- Guess
1. Effective Quark Theory.
 2. “Integrate out” high energy degrees of freedom.
 3. Chiral Symmetry.

<Y. Nambu, and G. Jona-Lasinio, PR122 & PR124, 1961>

$$\mathcal{L}_{NJL} = \bar{\psi}(i\gamma \cdot \partial - m_0)\psi + \frac{g_S}{2} \left[(\bar{\psi}\lambda^a\psi)^2 + (\bar{\psi}i\gamma_5\lambda^a\psi)^2 \right] + g_D [\det \bar{\psi}(1 - \gamma_5)\psi + h.c.]$$

4 Fermion Interaction

U_A(1) Breaking

- Similar to the microscopic theory of superconductivity.
- Non-renormalization theory in three dimension.

2 Flavors: Set $g_D = 0$.

$$\mathcal{L}_{NJL} \rightarrow \bar{\psi}(i\gamma \cdot \partial - m_0)\psi + \frac{g_S}{2} \left[(\bar{\psi}\lambda^a\psi)^2 + (\bar{\psi}i\gamma_5\lambda^a\psi)^2 \right]$$

NJL Model (2/2): Gap Equation

$$\langle \bar{\psi} \psi \rangle = \sigma$$

$$S_{eff} = -i \sum_{j=1}^{N_f} \ln [\det (i\gamma \cdot \partial - m_{0j} + 2g_S \sigma_j)] + \int d^4x \left[- \sum_{j=1}^{N_f} g_S \sigma_j^2 \right]$$

Effective mass: $m = m_0 - 2g_s \sigma$

$$V_{eff} = 2g_S \sigma^2 - 4N_c \int^\Lambda \frac{d^3p}{(2\pi)^3} \sqrt{p^2 + m^2} \quad \xrightarrow{\frac{dV_{eff}}{dm} = 0} \quad m = m_0 + 4g_S N_c \int^\Lambda \frac{d^3p}{(2\pi)^3} \frac{m}{\sqrt{p^2 + m^2}}$$

“Gap Equation”

If $g_S \Lambda^2 > \frac{\pi^2}{8}$, then $\sigma \neq 0$ and chiral symmetry is broken.

- There is a great deal of work on the model.
- A downside is NJL model lacks confinement.

PNJL Model (1/3): Setup

$$S_{eff} = -i \sum_{j=1}^{N_f} \ln [\det (i\gamma \cdot (\partial - \underline{iA}) - m_0 j + 2g_S \sigma_j)] + \int d^4x \left[- \sum_{j=1}^{N_f} g_S \sigma_j^2 \right]$$

- The covariant derivative couples fermions to a background Polyakov loop.
- A_0 is constant and diagonal. (“Polyakov Gauge”)
- Interplay between Polyakov loop and σ is now easy to see.

$$V_{PNJL} = V_{cond} + V_{zero} + V_{quark,T} + V_{gluon,T}$$

$$V_{cond} + V_{zero} = 2g_S \sigma^2 - 4N_c \int^\Lambda \frac{d^3 p}{(2\pi)^3} \sqrt{p^2 + m^2}$$

$$V_{quark,T} = -2N_f Tr_R [T \int \frac{d^3 p}{(2\pi)^3} \ln(1 \pm Pe^{-\omega_p/T}) + h.c.] = \frac{4m^2 T^2}{\pi^2} \sum_{n=1}^{\infty} \frac{(\mp 1)^n Tr_R P^n}{n^2} K_2(nm/T)$$

↑
 + Antiperiodic B.C.
 - Periodic B.C.

<P. Meisinger, and M. Ogilvie, PRD65, 2002>

$V_{gluon,T} \cdots$ Need a phenomenological model.

PNJL Model (2/3): Gluonic Part

$$V_{pert} = -2Tr_A \left[\int \frac{d^3 p}{(2\pi)^3} \ln(1 - Pe^{-\omega_p/T}) \right] \propto T^4$$

$$\omega_p = |p|$$

Valid at high T

Add subleading term (T^2) by hand to get correct behavior at low T. Introduce “a mass term”, M:

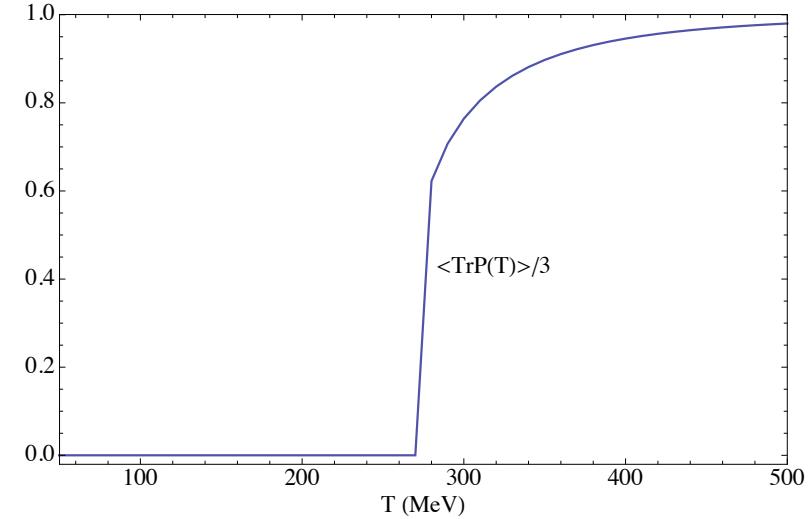
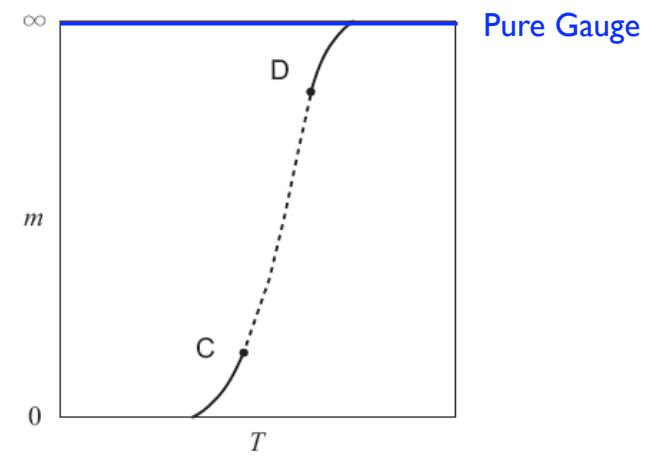
$$\omega_p = \sqrt{p^2 + M^2}$$

M is just a parameter, not mass of gluon.

<P. Meisinger, T. Miller, and M. Ogilvie, PRD65, 2002>

$$V_{gluon,T} = f(P)T^4 + g(P)M^2T^2$$

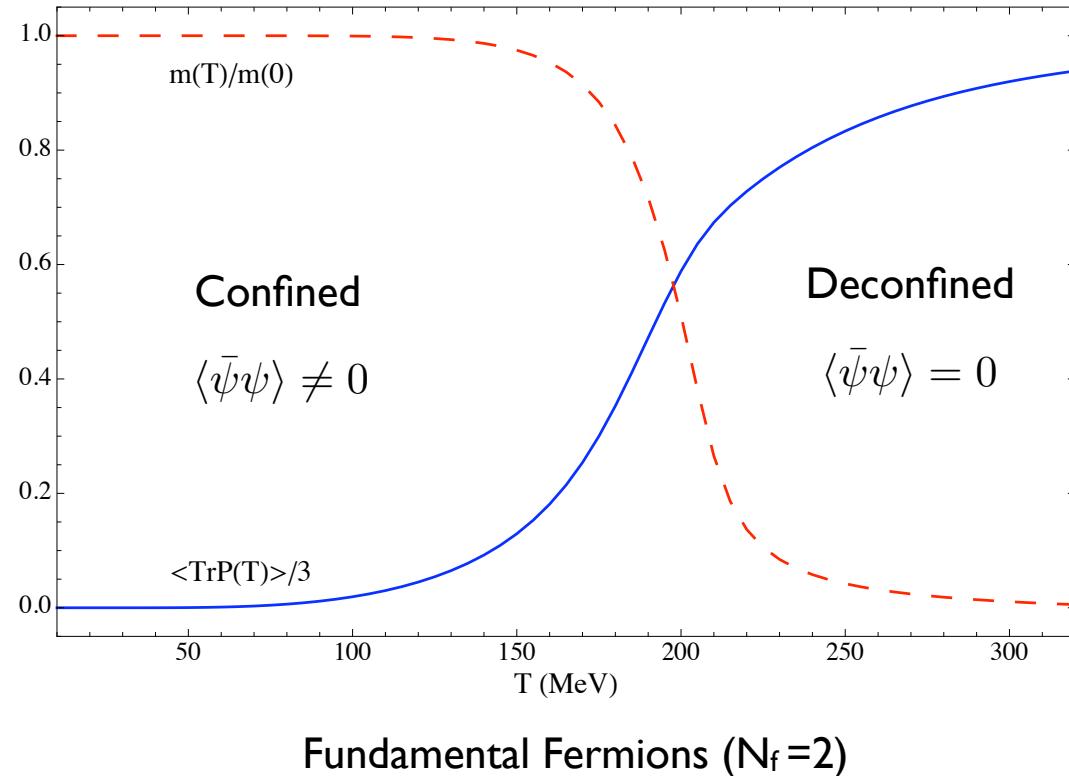
$$M = 596 \text{ MeV} \rightarrow T_d = 270 \text{ MeV}$$



Pure Gauge

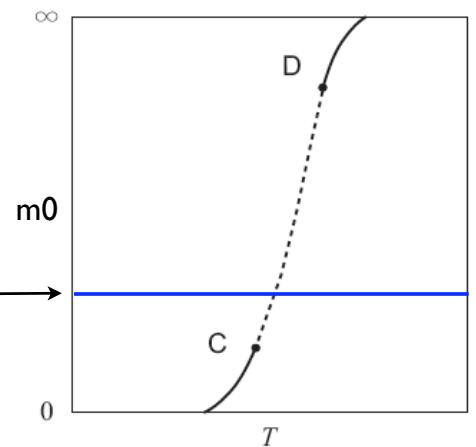
PNJL Model (3/3): Results

Now we have a whole potential: $V_{PNJL} = V_{cond} + V_{zero} + V_{quark,T} + V_{gluon,T}$



PNJL model works!

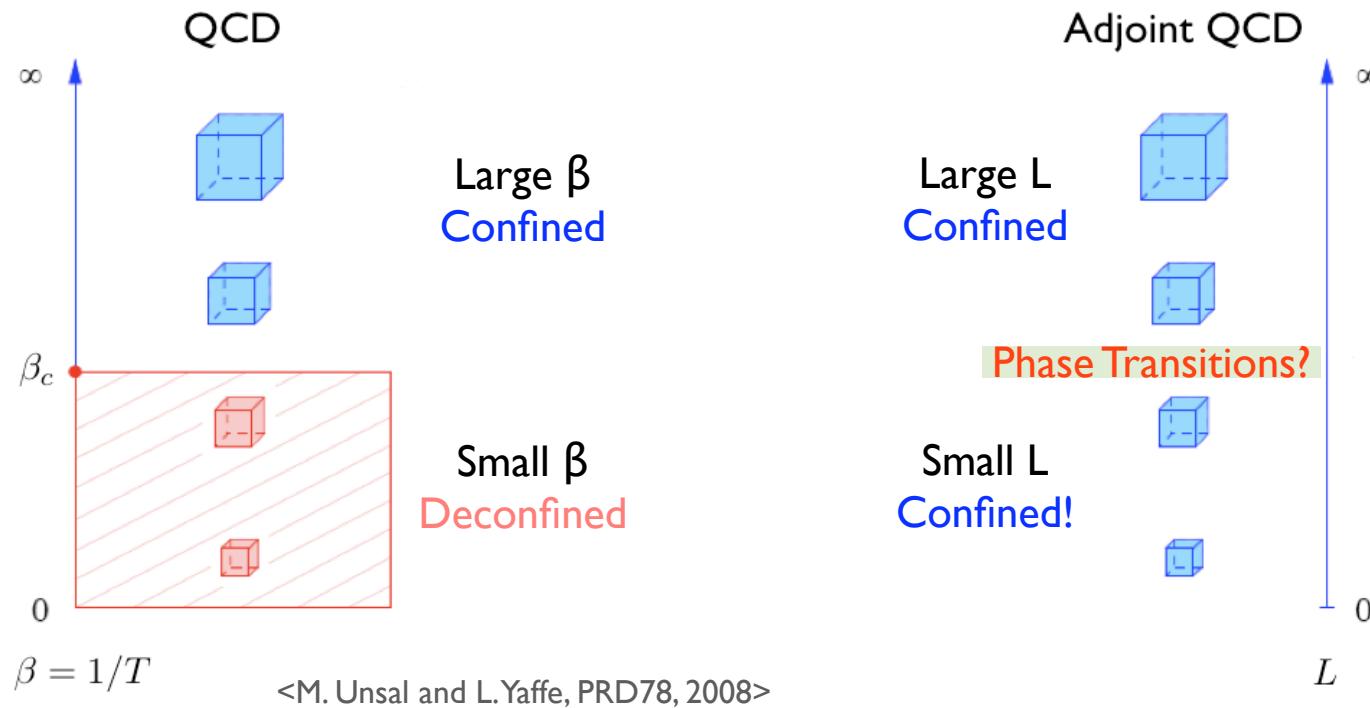
$$m_0 = 5.5 \text{ MeV}$$



Why Adjoint QCD?

- SUSY
 - $N_f = 1/2$
- Conformal Theory <E. Poppitz and M. Unsal, JHEP0909, 2009>
 - $2 < N_f < N^*$ is in the conformal window.
- Eguchi Kawai Model
 - Volume independence of SU(N) gauge theories in the large N limit.
- Confinement
 - Confinement at “high temperature” or small L.

Confinement with Small Radius



$$\mathcal{Z}_E = \int \mathcal{D}A \mathcal{D}\Psi \mathcal{D}\bar{\Psi} e^{-\int_0^\beta d\tau \int d^3x \mathcal{L}_E} = \text{Tr } e^{-\beta H}$$

ABC:Temporal Component $\rightarrow \beta = 1/T$
 PBC:Spatial Component $\rightarrow \beta = L$

Adjoint Fermions with PBC

- Confinement at “high T” or small radius.
 <P. Kovtun, M. Unsal, and L.Yaffe, JHEP, 2007>
- New confinement is analytically tractable.
 <M. Unsal, PRD, 2009>

Adjoint fermions with ABC (finite T) in PNJL

Lattice shows that scale of deconfinement and chiral symmetry restoration are different!

$$T_\chi/T_d \approx 7.8$$

<J. Engels, S. Holtmann, and T. Schulzel, NPB724, 2005>

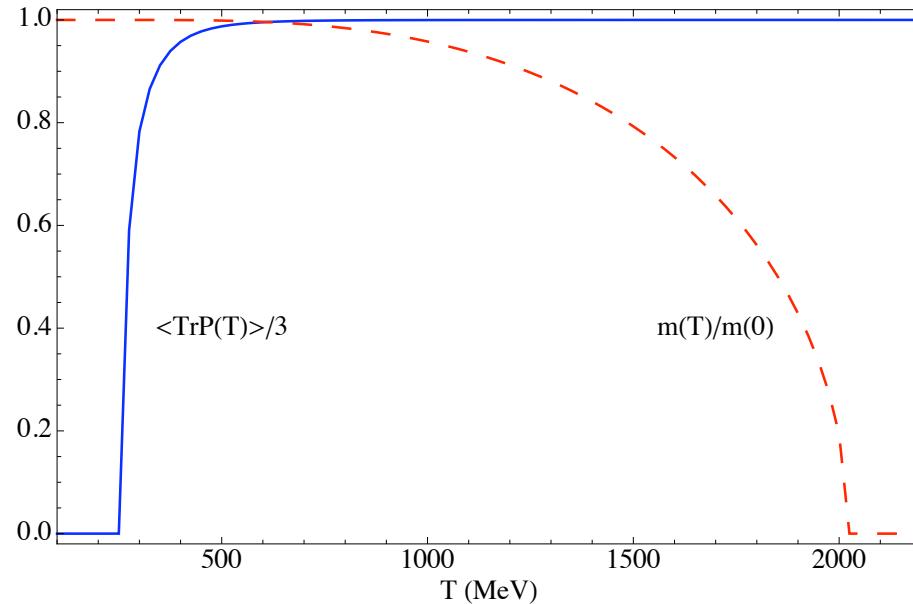
Choose the parameters;

$$m_0 = 0$$

$$\Lambda = 23.22 \text{ GeV}$$

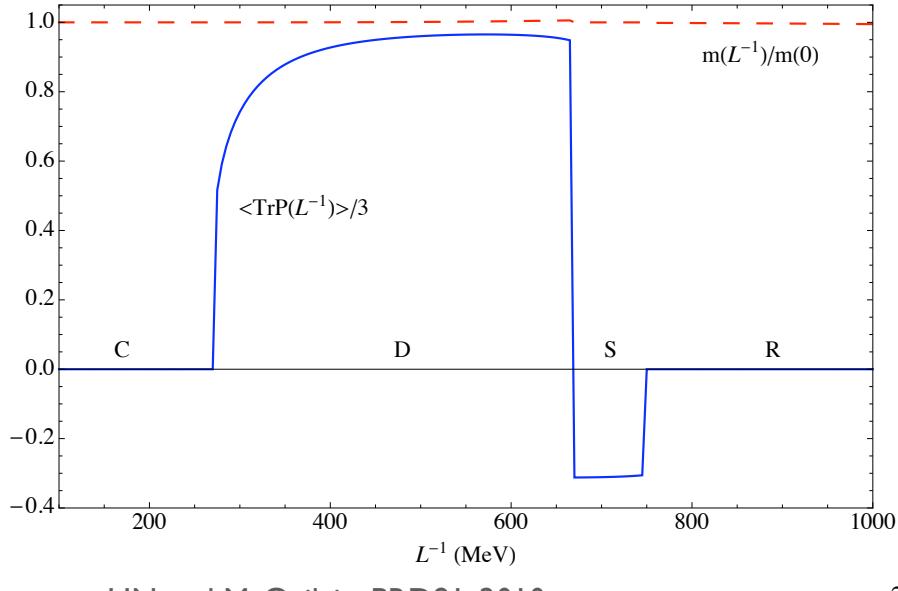
$$g_S \Lambda^2 = 1.2653$$

such that $T_\chi/T_d \approx 7.8$.



<HN and M. Ogilvie, PRD81, 2010>

Adjoint fermions with PBC in PNJL (1/2)



<HN and M. Ogilvie, PRD81, 2010>

$$m_0 = 0$$

$$\Lambda = 23.22 \text{ GeV}$$

$$\kappa = g_S \Lambda^2 = 1.2653$$

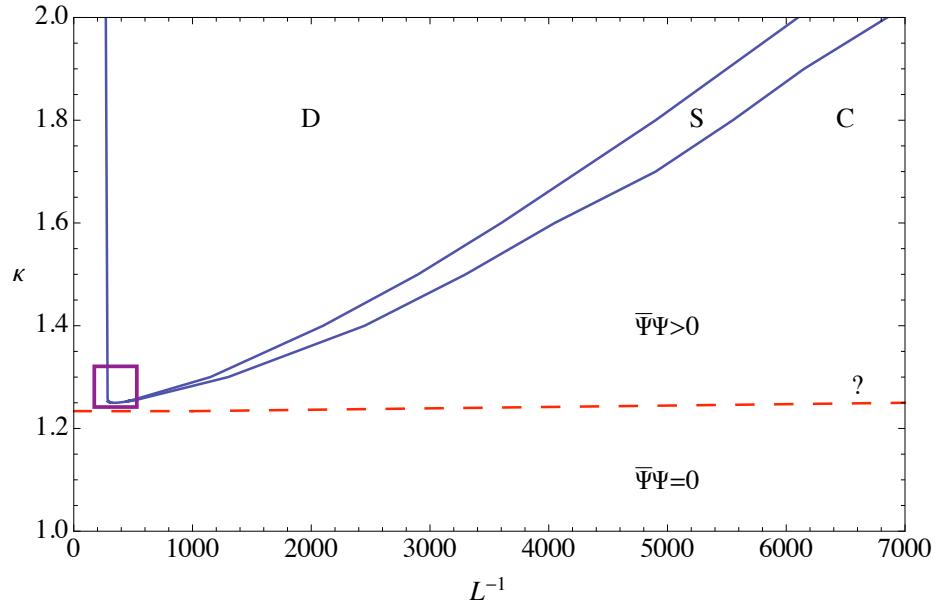
C = Confined

D = Deconfined

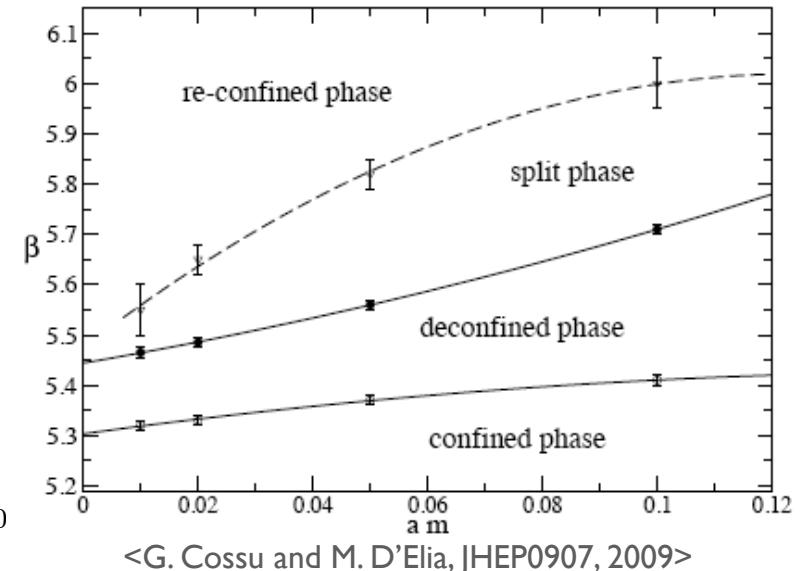
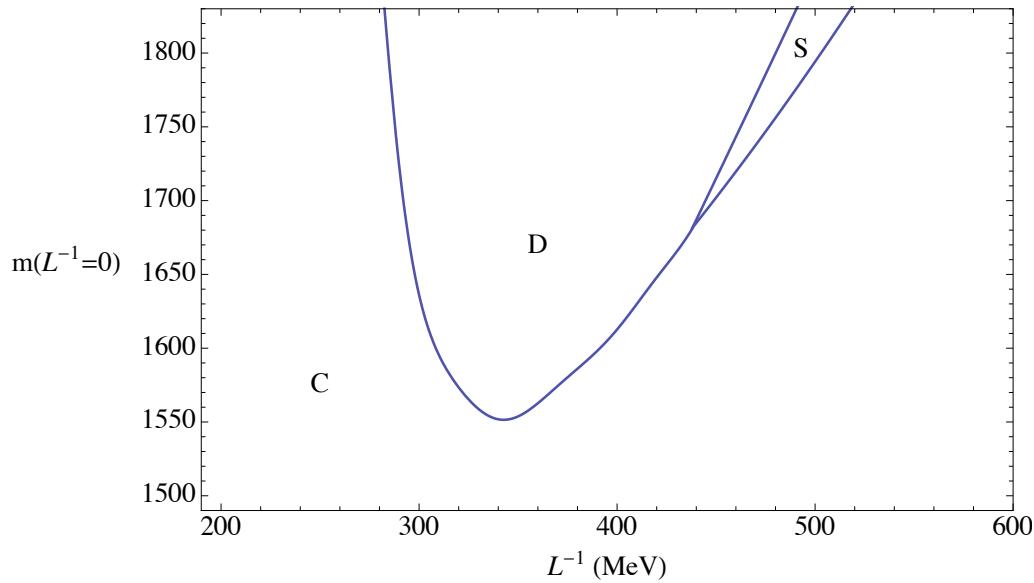
S = Skewed Phase

R = Reconfined

- No chiral symmetry restoration up to 10 GeV.
- 3 distinct 1st order phase transitions for TrP .



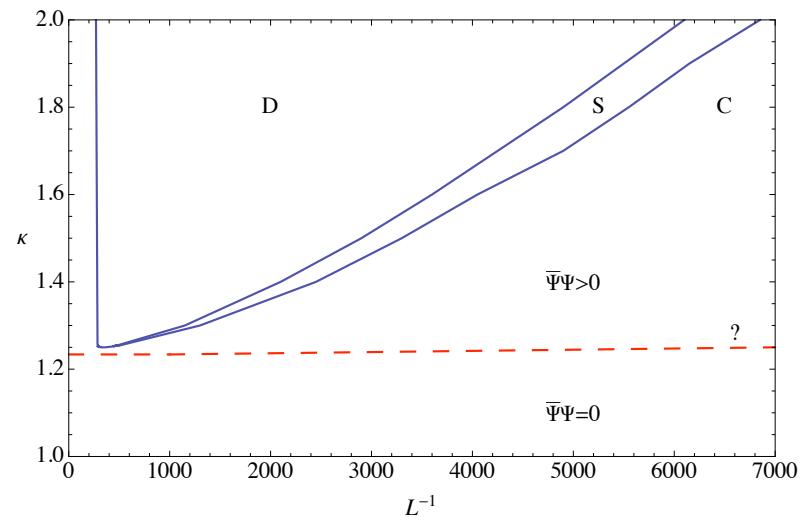
Adjoint fermions with PBC in PNJL (2/2)



Remaining Puzzle: Casher's Argument

- Unsal's scenario predicts confinement without chiral symmetry breaking at small L .

- We need more lattice results!



Conclusions

- Two fundamental fermions with ABC.

The PNJL model can reproduce two crossover transitions at $T \approx 170$ MeV.

- Two adjoint fermions with ABC.

The PNJL model can reproduce $T_\chi/T_d \approx 7.8$.

PNJL works for finite-temperature cases very well.

- Two adjoint fermions with PBC.

The PNJL model shows that two confining regions are connected.

We need more lattice results.