

Conformal or Walking?
Monte Carlo Renormalization Group studies in
technicolor-inspired models

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Anna Hasenfratz
University of Colorado

Beyond Standard Model physics on the lattice

Many of interesting candidates are fermion-gauge systems

In Euclidean space they can be discretized

→ they turn into 4d statistical systems

→ identifying phases, critical points, measuring critical exponents
can be done just like in a statistical physics

Dictionary:

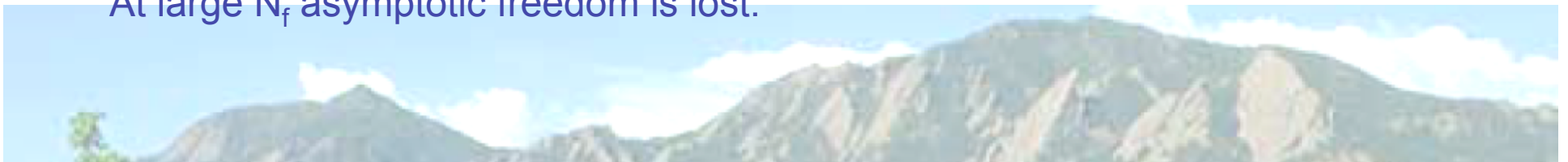
m or m_q	:	fermion mass of N_f degenerate fermions
g	:	gauge coupling of $SU(N_c)$ gauge system
β	:	either $\beta=2N/g^2$ or $\beta = \mu \, dg/d\mu = -b_0 g^3 - b_1 g^5 + \dots$ RG β function
γ_m	:	mass anomalous dimension, $\gamma_m = y_m - 1$

If the system is asymptotically free it can be

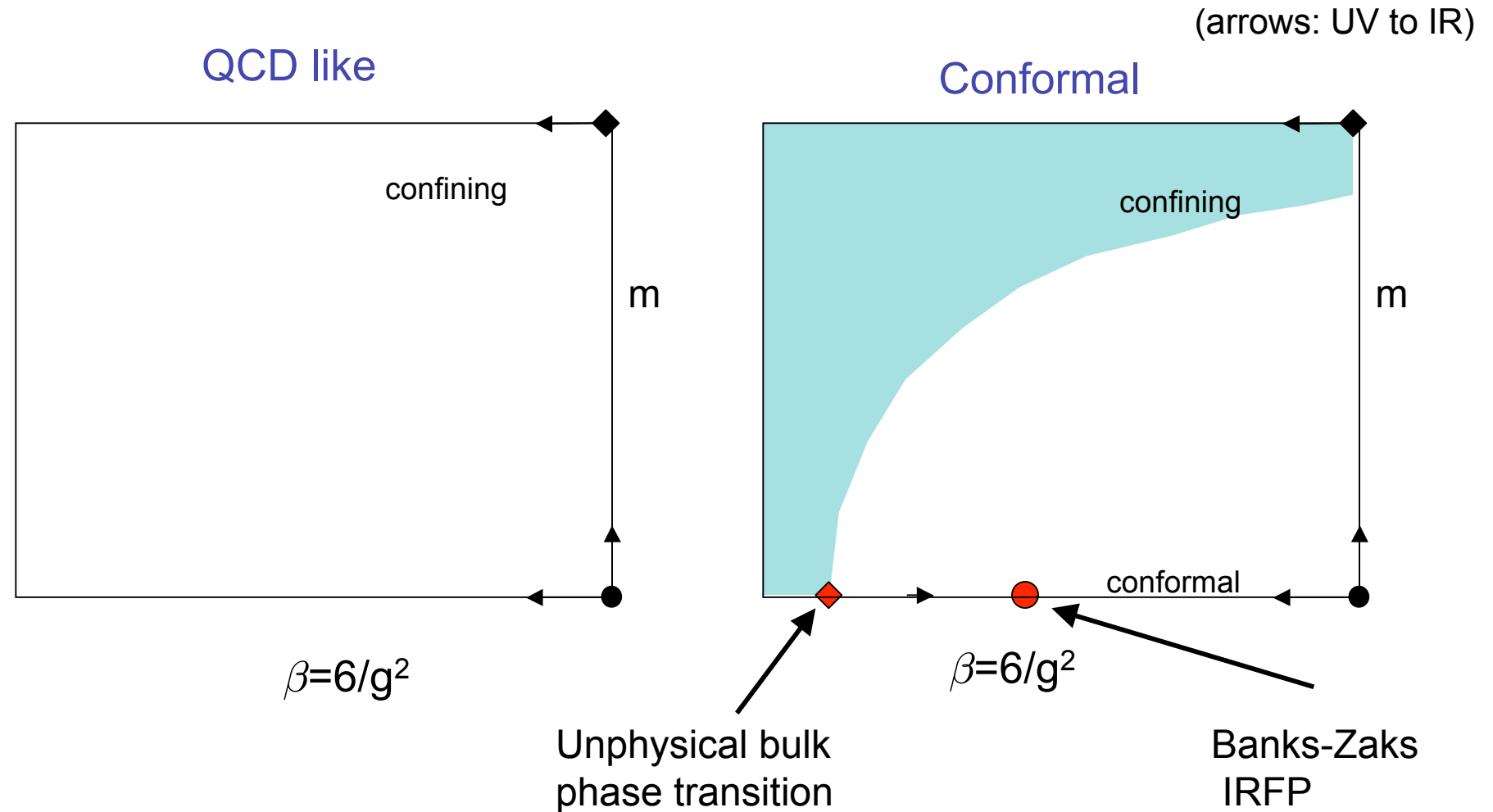
–confining and chirally broken (small N_f) or

–conformal in the $m_q=0$ limit (increasing N_f)

At large N_f asymptotic freedom is lost.

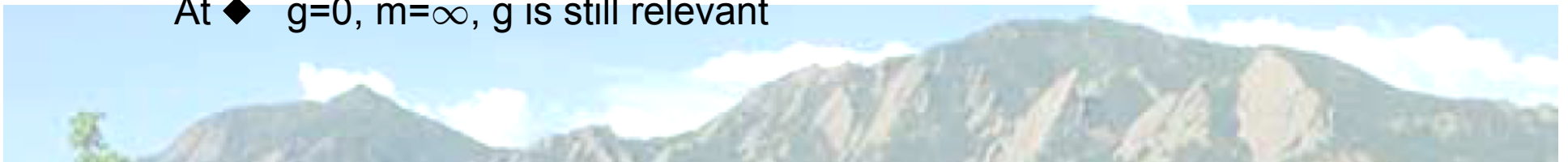


The lattice phase diagram



At ● $g=0, m=0$, both couplings are relevant

At ◆ $g=0, m=\infty$, g is still relevant

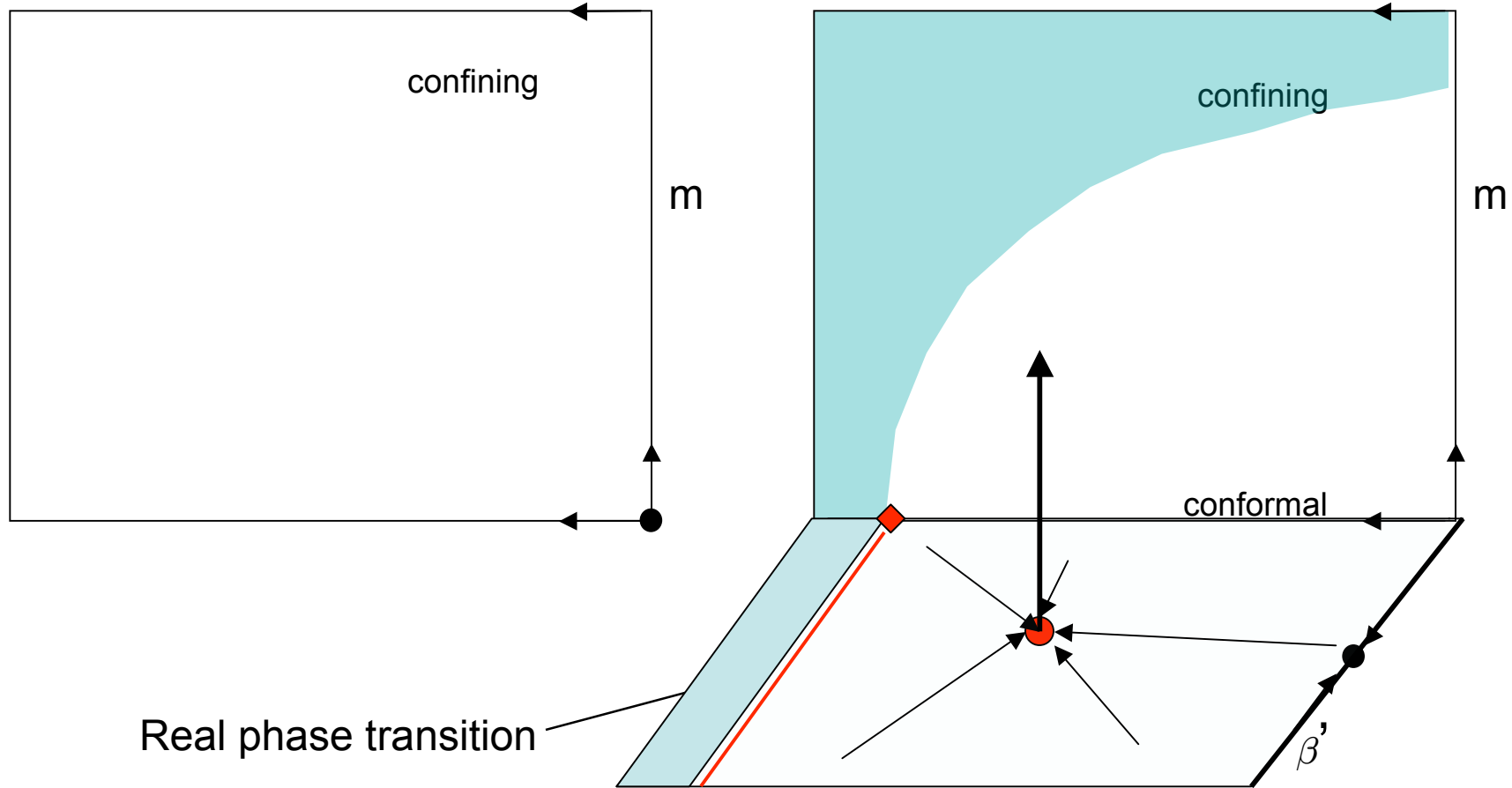


The lattice phase diagram

(arrows: UV to IR)

QCD like

Conformal



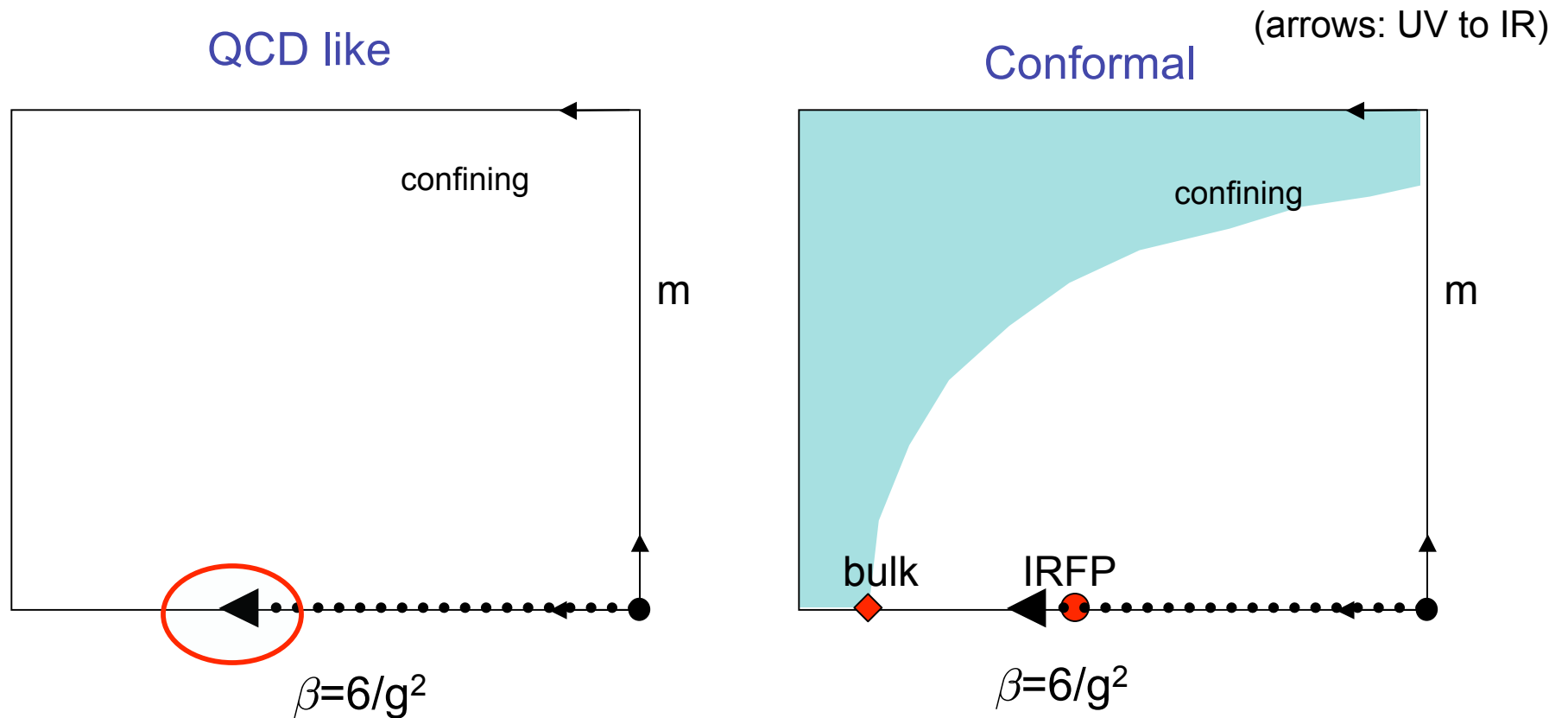
Real phase transition

Existence of IRFP is universal;
location is RG dependent

At \bullet IRFP only m is relevant



We want to distinguish QCD-like and conformal systems:



Lattice simulations can connect the perturbative FP and strong coupling

- Found IRFP ? Done ✓
- No IRFP? Show that it is confining before a bulk transition is reached
- Strong lattice artifacts can interfere

Outline

- The bare step scaling function as an alternative to the RG β function
- Monte Carlo Renormalization Group method
 - The 2-lattice matching approach
 - Optimization
 - The role of different RG transformations
- Some results:
 - Pure gauge SU(3) – testing case for MCRG
 - $N_f=8, 16, 12$ fundamental fermions, SU(3) gauge
 - The steps scaling function
 - The anomalous dimension of the mass



The step scaling function around a UVFP

I measure the bare differential step scaling function s_b
(instead of the RG β function)

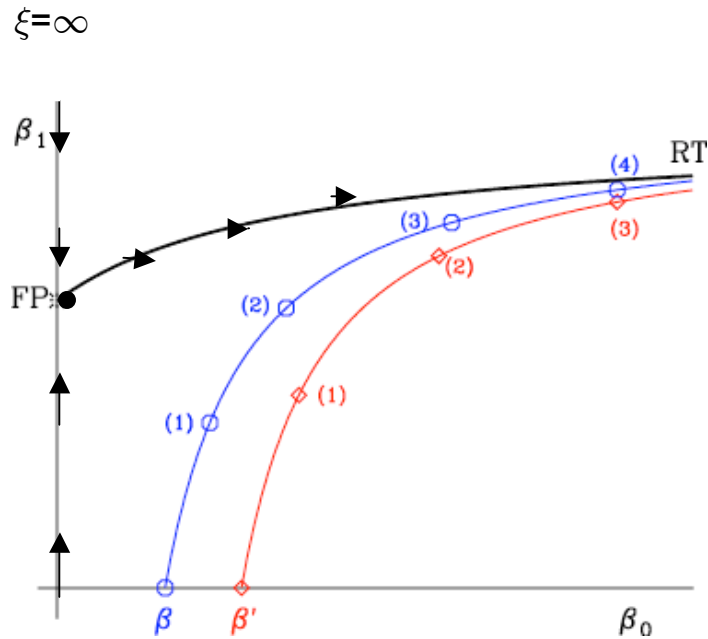
$$s_b(\beta) = \beta - \beta' \quad \text{where} \quad \xi(\beta) = \xi(\beta')/2 \quad (\beta = 2N_c/g_0^2)$$

ξ is the correlation length defined by some physical mass

- Can be measured directly or
- Use RG flow



RG flow lines along a relevant direction (UVFP):

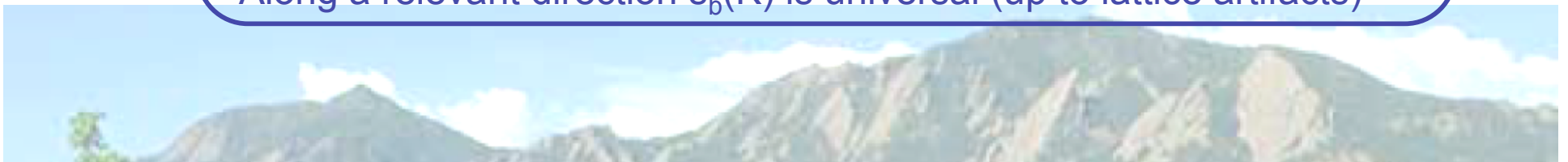


- do simulations at β and β' ($m=0$)
- do RG blocking and compare the blocked actions
- if $S(\beta^{(n)}) = S(\beta'^{(n-1)}) \rightarrow a(\beta) = a(\beta')/2$
- the step scaling function is

$$s_b(\beta) = \lim_{n_b \rightarrow \infty} (\beta - \beta')$$

- The location of the FP on the critical surface depends on the RG transformation
- Tuning free parameters in the RG transformation can pull the FP and its RT close, reducing systematical errors

Along a relevant direction $s_b(K)$ is universal (up to lattice artifacts)



The step scaling function in a conformal system

In the chiral limit $\xi = \infty$ everywhere !

$s_b(\beta)$ can be defined through the RG flow or the running coupling

$$s_b(\beta) = \beta - \beta' \quad \text{where} \quad S^{(n)}(\beta) = S^{(n-1)}(\beta')$$

or

$$g^2(\beta;L) = g^2(\beta';L/2) \quad (\beta=2N_c/g_0^2)$$

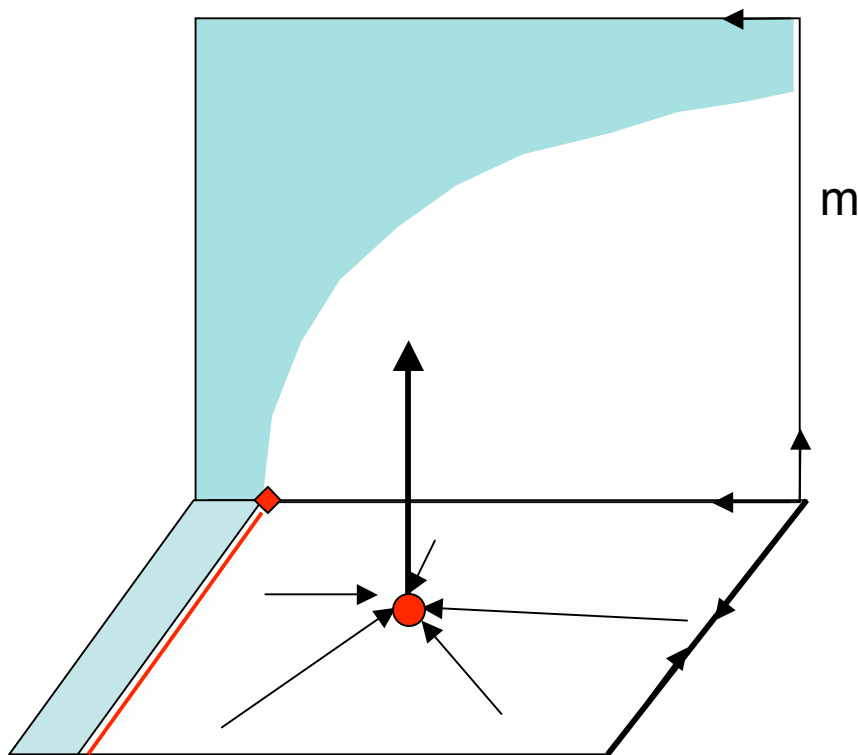
$(g^2(\beta;L))$ can be defined via Schroedinger functional or other RG transformation)



RG flow lines around an IRFP

On the critical surface ($m=0$) around an IRFP the flows converge to the FP when $n_b \rightarrow \infty$

With finite n_b the flow picks up the slowest flowing operator



The location of the IRFP depends on the RG transformation

$s_b(\beta)$ along an irrelevant direction depends on the blocking (scheme dependence)

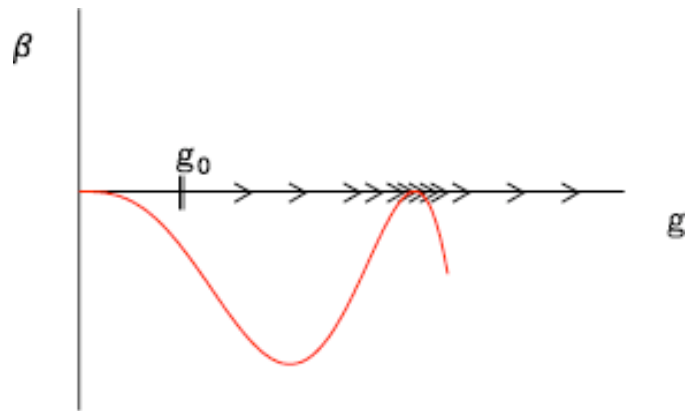
This is a signal for non-QCD-like behavior



The step scaling function of a walking theory

What happens with a walking theory?

- It is QCD-like, ξ is defined, in the scaling region of the UVFP $s_b(\beta)$ is universal (up to $O(a^2)$ corrections).
- The near-zero of the RG β function depends on the RG transformation. At that region $s_b(\beta)$ is RG dependent.



Can confinement and χ SB be established before lattice artifacts overwhelm the system?



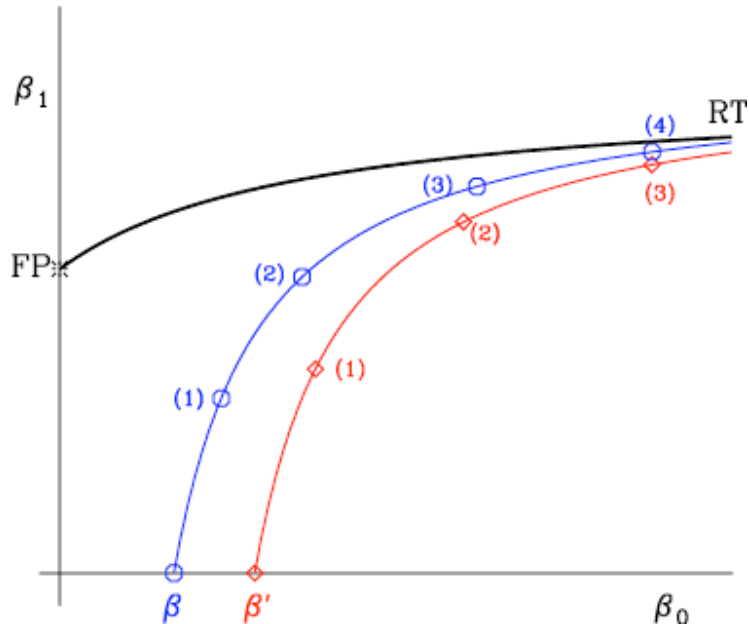
Properties of the step scaling function

This is the **bare** differential step scaling function

- at a fixed point $s_b(\beta^*)=0$
- the value of s_b is related to the scaling dimension of the coupling
 - for AF models $s_b = 3 \ln(2)/(4\pi^2) b_0 + O(g^2)$
 - $s_b > 0$ where the RG β -function is $\beta(g) < 0$ (sorry)
 - s_b in the mass predicts the anomalous dimension of the mass
$$m = m' 2^{1/y}$$



Calculating $s_b(\beta)$ with MCRG



Two actions are identical if all operator expectations values agree



Match operators after several blocking steps

Tests: SU(3) pure gauge (test system)

blocking: "Original"

$$\begin{array}{c}
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α : optimization parameter

$(1-\alpha)$

$\alpha/6$

$\alpha/6$



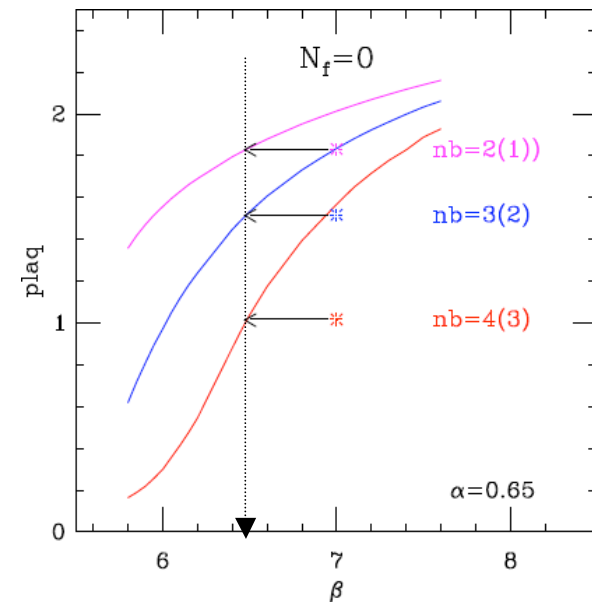
The plaquette* after 1-4 levels of blocking

$32^4 \rightarrow 16^4 \rightarrow 8^4 \rightarrow 4^4 \rightarrow 2^4$ (symbols) compare to

$16^4 \rightarrow 8^4 \rightarrow 4^4 \rightarrow 2^4$ (lines)

Repeat with many different operators. If they all give the same result, we found matching

$$s_b(\beta=7.0) = \Delta\beta = \beta - \beta' = 7.0 - 6.49 = 0.51$$

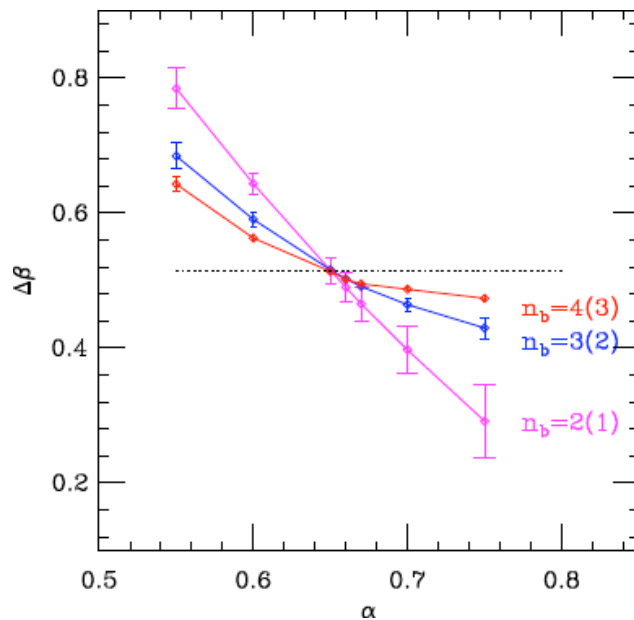


*plaquette: $\text{Tr}(U_{\square})$



Optimization of the RG transformations

$\Delta\beta = \beta - \beta'$ at $\beta = 7.0$ as the function of the RG parameter



- Optimizing the RG transformation is essential
 $\alpha_{opt} = 0.65$
- Optimized RG gives the same matching value at each level, for each operator

$$a(\beta=7.0) = a(\beta'=6.49)/2$$

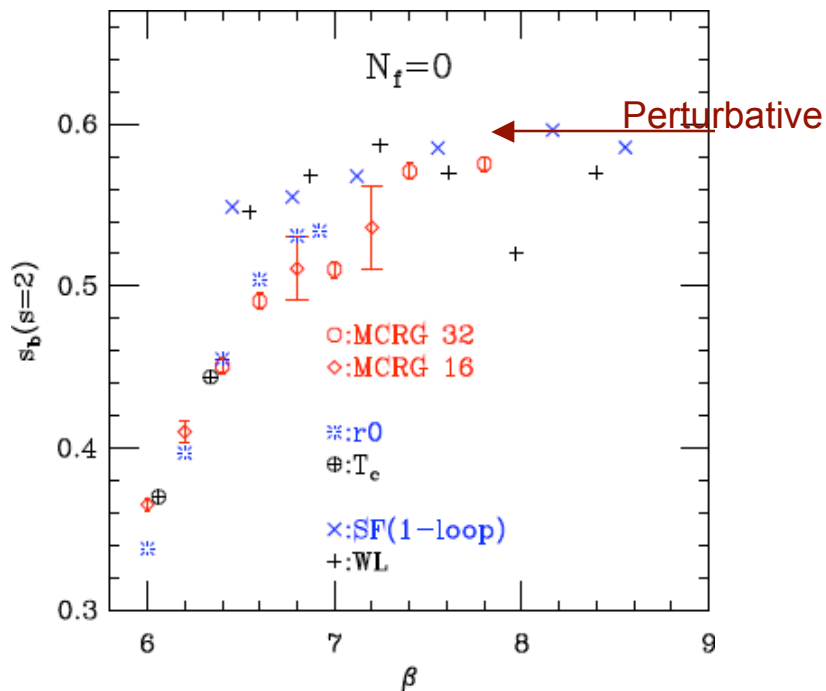
$$s_b = \beta - \beta' = 0.51$$



SU(3) pure gauge

The bare step scaling function can be calculated in many ways

- Schrodinger fn; Wilson loop ratios,
- physical observables r_0, T_c
- RG matching: $32^4 \rightarrow 16^4$ and $16^4 \rightarrow 8^4$



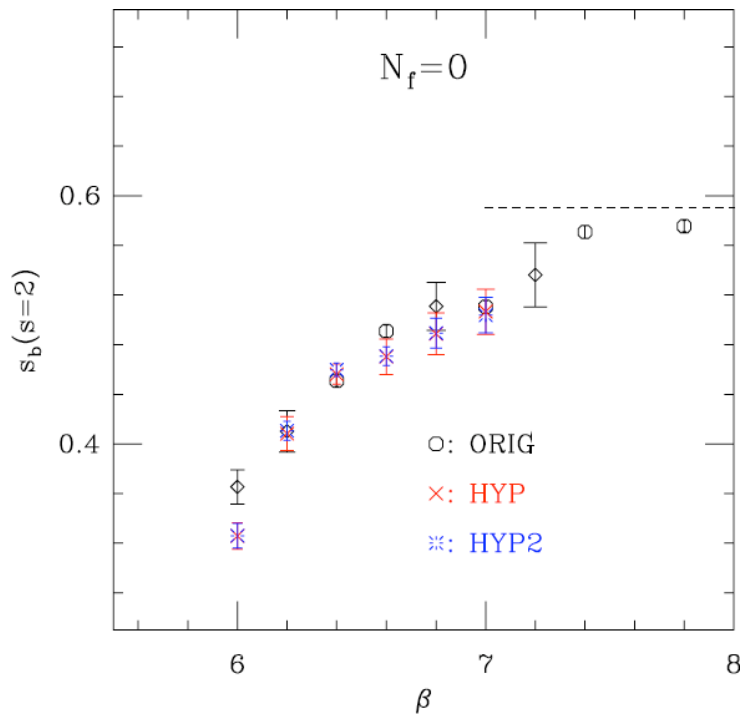
- Excellent agreement between r_0, T_c and MCRG
- Both SF and MCRG approach the perturbative value
- Since at $\beta=6$ we can test confinement, we know there is no physical IRFP



Compare different RG transformations:

When the flow is governed by a UVFP, $s_b(\beta)$ is universal (up to lattice corrections).

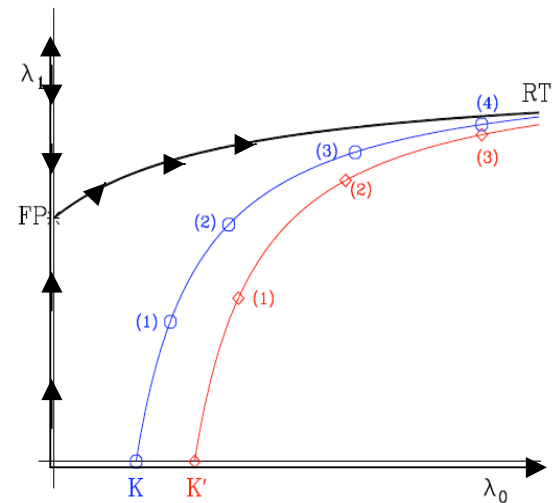
Compare 3 different RG transformations:



Excellent agreement between the
3 RG blockings
→ attractive region of a UVFP



Why do we need different RG's ?



For matching the RG flow has to approach the renormalized trajectory

- The RT describes “perfect actions”, i.e. no lattice artifacts

A “good” RG should remove most of the UV modes

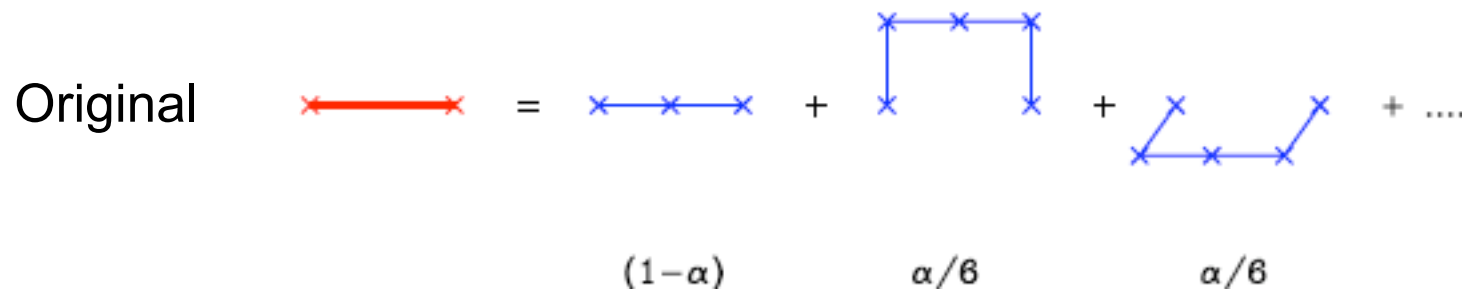
- A “good” RG will approach the RT faster

This is especially important towards strong coupling where UV fluctuations are large

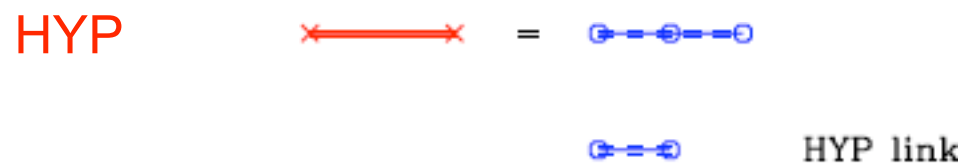


The 3 Renormalization Group transformations

A real space block transformation averages out the short distance modes
 Many possibilities - I tried 3 types:



optimize with α



optimize with α_1
 (play with α_2, α_3)

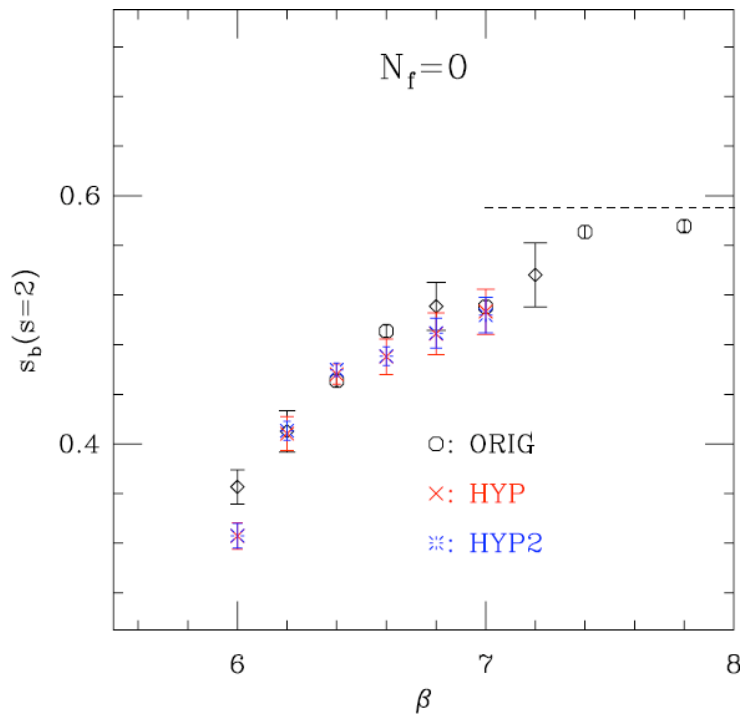
HYP2 like HYP, but with twice blocked links



Compare different RG transformations:

When the flow is governed by a UVFP, $s_b(\beta)$ is universal (up to lattice corrections).

Compare 3 different RG transformations:



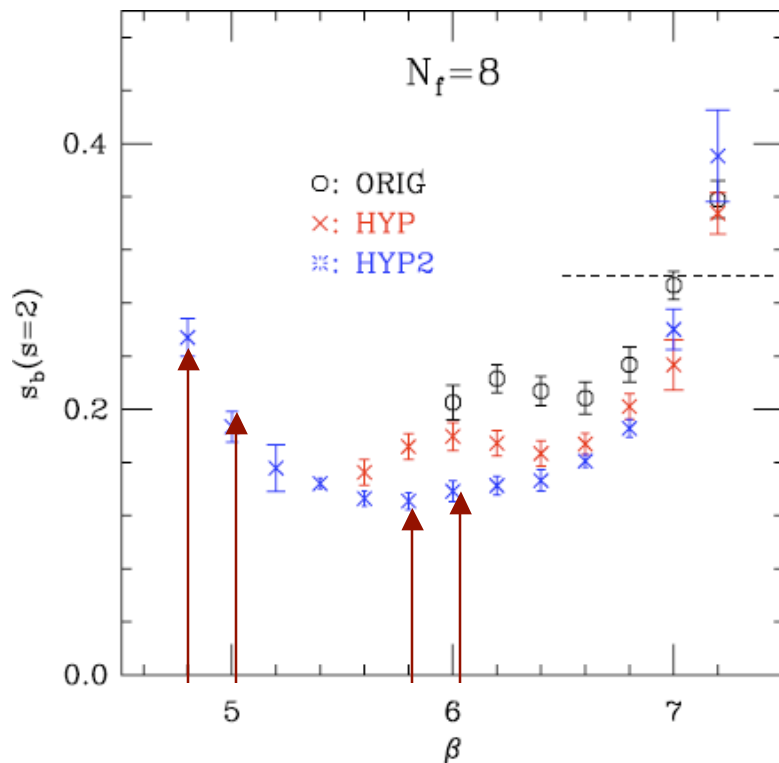
Excellent agreement between the 3 RG blockings
→ attractive region of a UVFP



$N_f=8$ flavors

Expected to be QCD-like: analytical & numerical results

Compare the different RG transformations:



$s_b > 0$ everywhere - no IRFP

At $\beta \sim 5.0$ RG matching gets difficult, but by then confinement develops (string tension is $a\sqrt{\sigma} \sim 0.2$)

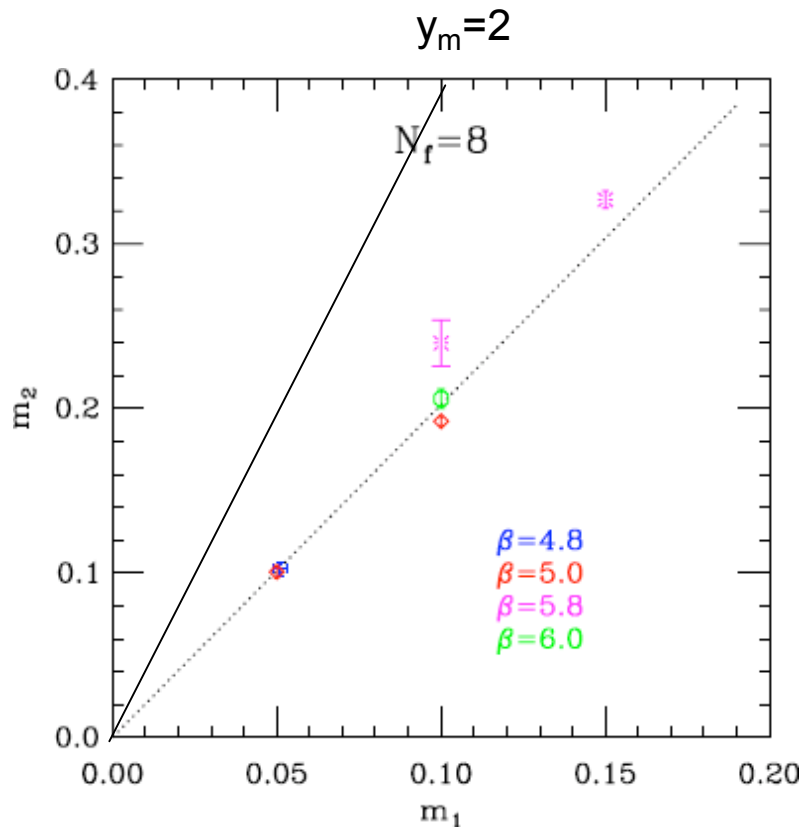
Considerably larger difference between the 3 RG blockings
non-QCD like behavior?

Look at the anomalous mass dimension



$N_f=8$ flavors, anomalous mass

4 different couplings ($\beta=4.8, 5.0, 5.8, 6.0$), optimal RG from $m=0$ data



$$m_2 = m_1 2^{-1/y_m}$$

$$\gamma_m = y_m - 1$$

All 3 couplings predict the same value

$$y_m = 1.02(5)$$

close to free field exponent



$N_f=16$ flavors

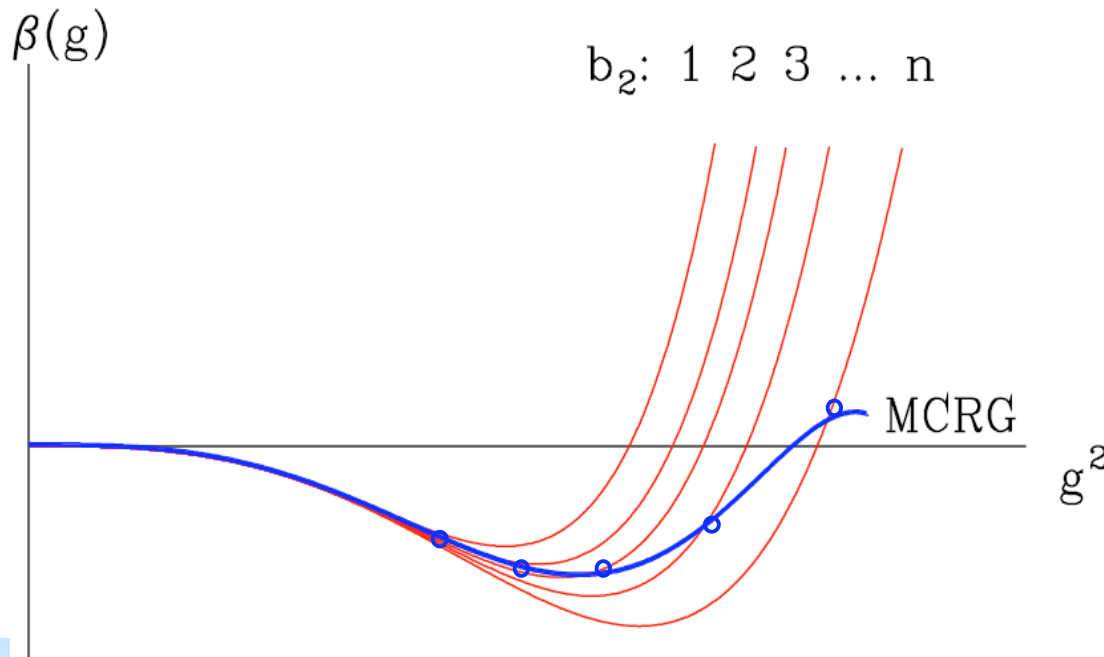
MCRG optimizes the RG as the function of the bare coupling . Along an irrelevant direction

-it can pick a different RG at each coupling

-the β function it determines does not correspond to any “real” β function

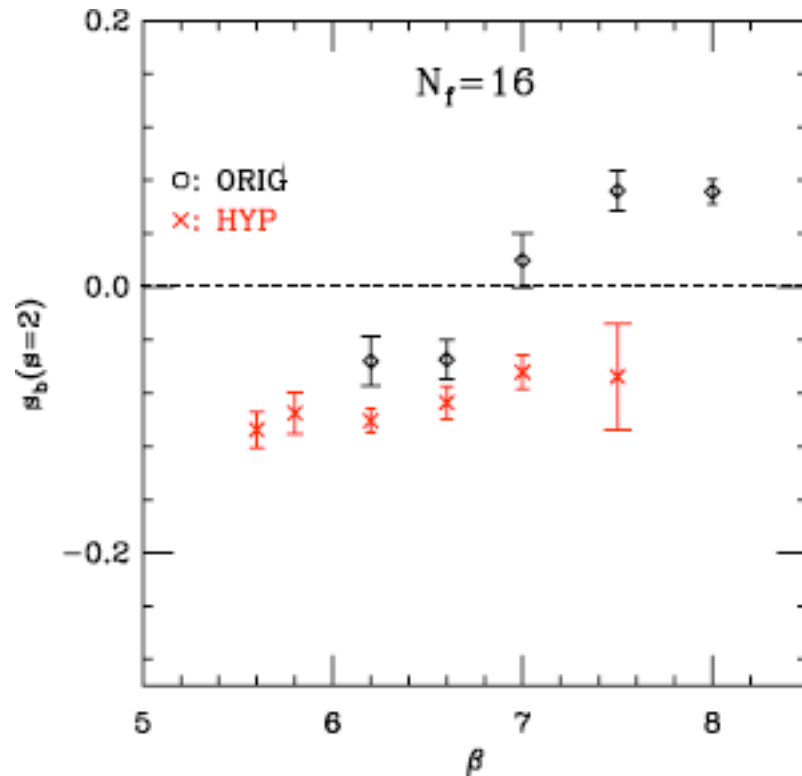
Warning!

But: a zero is a zero



$N_f=16$ flavors

$16^4 \rightarrow 8^4$ MCRG



ORIG blocking shows $s_b(\beta)=0$ around $\beta=7.0$

HYP blocking has an IRFP around $\beta=8.0$

Different block transformations predict different $s_b(\beta)=0$ but they both show a positive RG β function



$N_f=12$ flavors

Some history:

- The analytic work of Appelquist et al predicts $N_f=12$ is just above the conformal window
- Yale group found an IRFP at fairly strong coupling, using Schrodinger functional method
- Groningen group identified a bulk phase transition characteristic to a conformal system and claim chiral symmetry at weaker coupling
- Two groups (San Diego and Columbia) have studied the spectrum of the model with improved and unimproved actions. Both see QCD-like behavior, though at strong gauge couplings.

If $N_f=12$ is conformal,

the spectral measurements are in the strong coupling lattice artifact phase

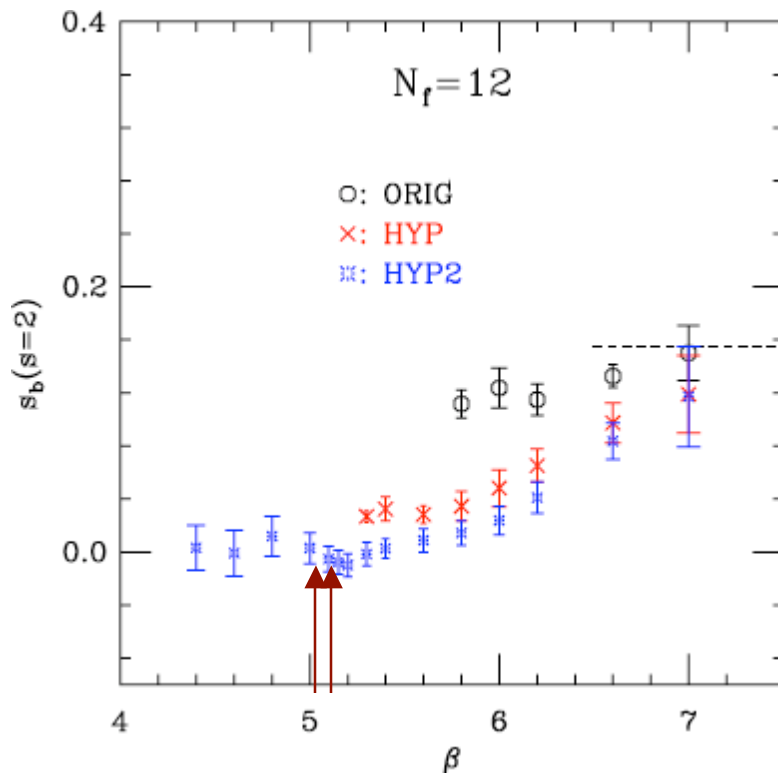
If $N_f=12$ is QCD-like,

the unimproved actions used with Schrodinger functional are unreliable



$N_f=12$ flavors with MCRG

Use the same techniques as before; $16^4 \rightarrow 8^4$



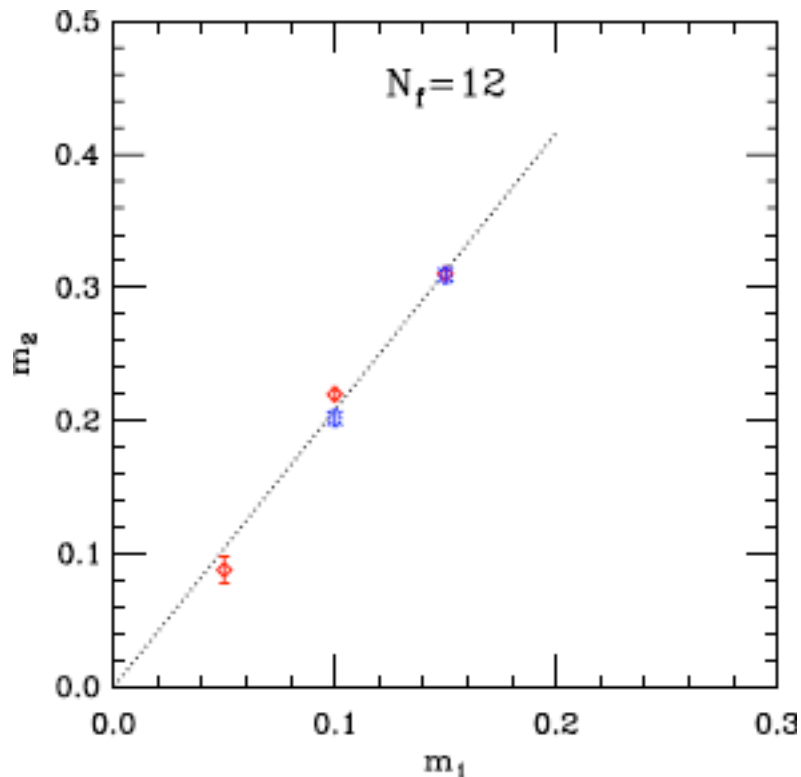
- **Orig/HYP** blockings predicts different $s_b(\beta)$ functions
- **HYP2** hovers around 0 -- likely IRFP
- String tension remains zero at $\beta=4.4$ on 16^4 volumes, but lattice artifacts are large

$N_f=12$ could be walking between $\beta=4$ and 6 - but that would be strange, the least



$N_f=12$ - anomalous mass

- $\beta=5.0, 5.1$: both where $s_b \sim 0$;
using optimal RG from $m=0$ data



Again, consistently

$$y_m = 1.06(3)$$

for both couplings, masses

At a strongly coupled IRFP one expects a large y_m

Maybe it is not strongly coupled

Note:

SU(2) with adjoint fermions (DeLDebbio et al) looks the same ; SU(3) with sextet (T. DeGrand) has $y_m \sim 1.5$



Summary: $N_f=12$ flavors

- It would be nice to firmly establish a back flow or confinement
 - Larger volume simulations
 - Different action
 - Different RG transformation
- For now, $N_f=12$ looks conformal, but with a trivial exponent



Summary: 2- lattice matching MCRG

- Can be optimized by tuning the free parameter(s) of the RG transformation
- Finite volume effects are largely controlled
- Requires relatively small statistics
- Has a lot of built-in consistency checks
 - compare several blocking levels
 - compare several operators
 - compare different RG transformations



Conclusion

MCRG is an effective alternative method to study the phase structure and scaling properties of lattice QFT's

- The method is very universal, straightforward to implement for any other system

$N_f=0-8,16$ as expected. $N_f=12$ is difficult:

- $y_m \sim 1.0$ for 12 flavors; far above the opening of the conformal window?

What is next?

- Could the different groups come up with a consistent picture for $N_f=12$?
- $N_f=10,(9,11?)$ would be very interesting
- SU(2) gauge, other fermion representations can be studied the same way



EXTRA SLIDES

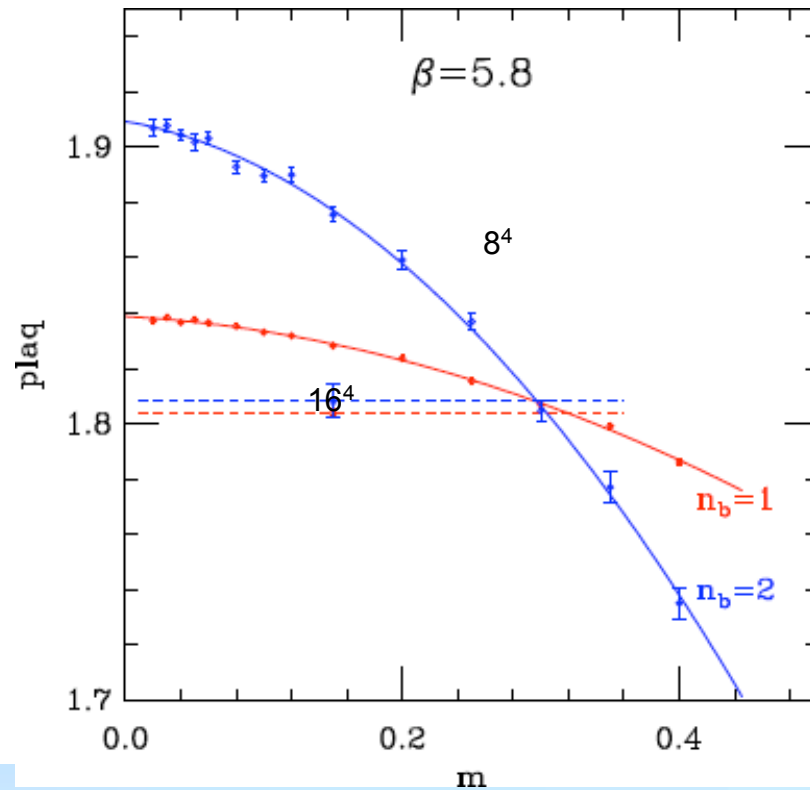


MCRG to find the mass anomalous dimension

$N_f=16$ flavor SU(3) model

Matching in the mass at fixed $\beta = 5.8$

$$m_2 = m_1 2^{1/\nu}$$



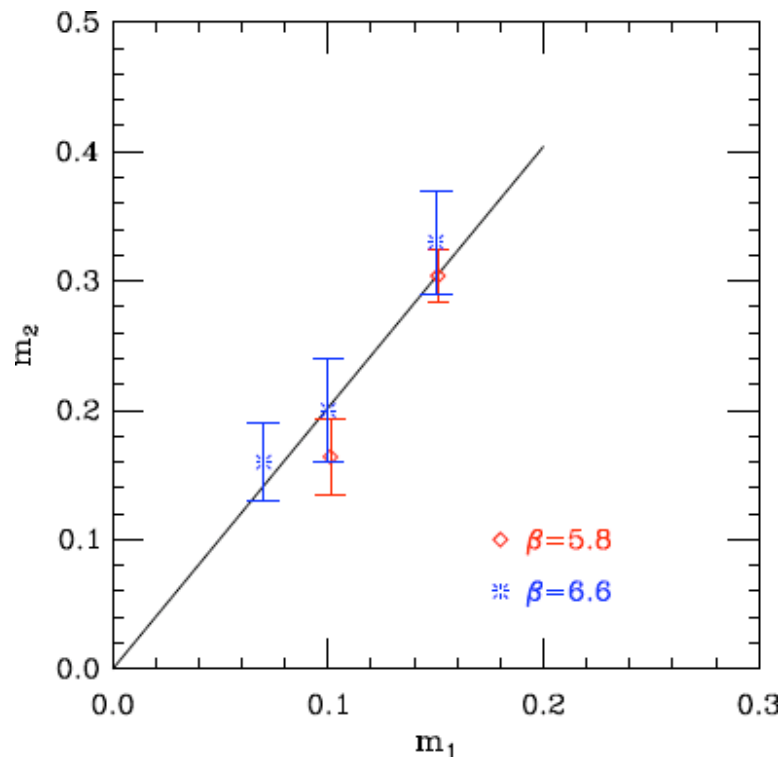
- use the same gauge observables (probably not the best choice)

-at α_{opt} both $n_b=2(1)$ and $3(2)$ predicts the same matching pair



The critical exponent for the mass

At several couplings, mass values



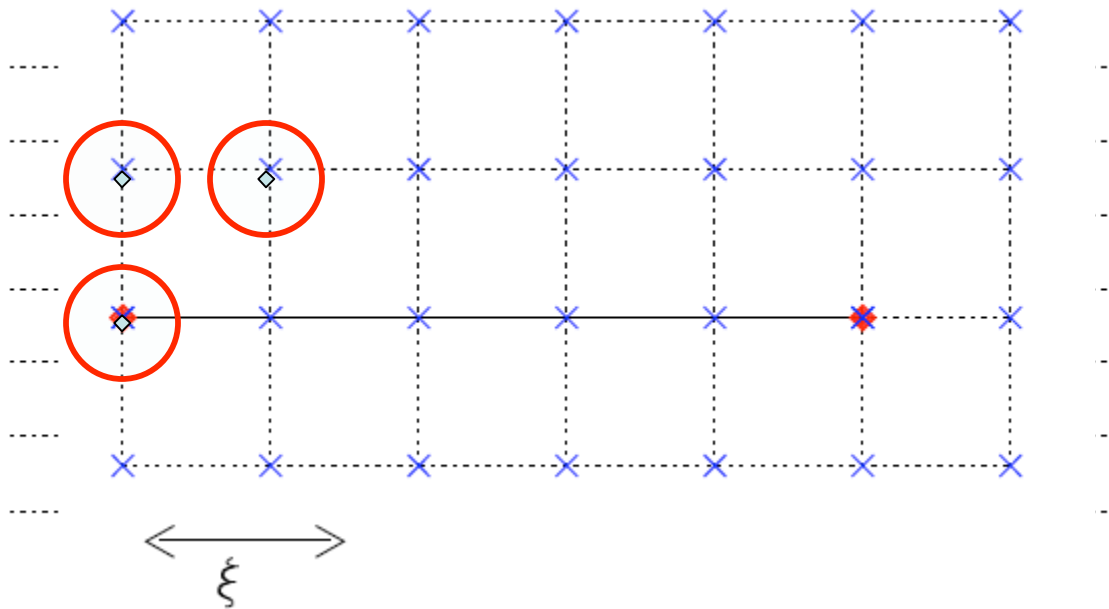
$$m_2 = m_1 2^{1/\nu}$$

$$\nu=1.0(1)$$

Free field exponent (close to GFP)



Real space block transformation



$$\langle s(0)s(x) \rangle \sim e^{-x/\xi}$$

if $x \gg \xi$

Original :

variables :

$s(x)$

lattice spacing:

a

correlation length:

$\xi = \xi_{lat} a$

Action:

$\{K\}$

Blocked in $(b=2)^d$

$s'(x) =$ (average of $s(x)$ in block)

$a' = 2a$

$\xi'_{lat} = \xi_{lat}/2$, (but $\xi' = \xi$!)

$\{K'\}$

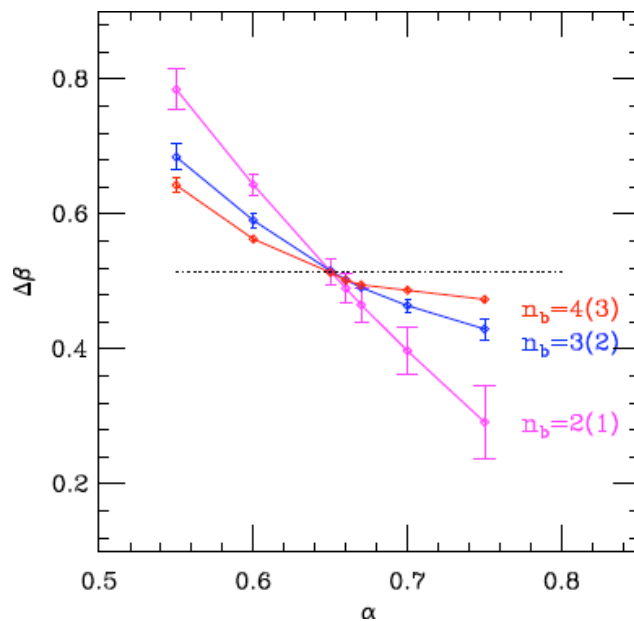
The two systems have the same IR (as long as $\xi_l > 1$)

$K'_i(K_j)$ describes the evolution of the action (flow lines)



Optimization of the RG transformations

$\Delta\beta = \beta - \beta'$ at $\beta = 7.0$ with different RG parameter, blocking levels, operators



- Optimizing the RG transformation is essential
 $\alpha_{opt} = 0.65$
- Optimized RG gives the same matching value at each level, for each operator

$$a(\beta=7.0) = a(\beta'=6.49)/2$$

$$s_b = \beta - \beta' = 0.51$$

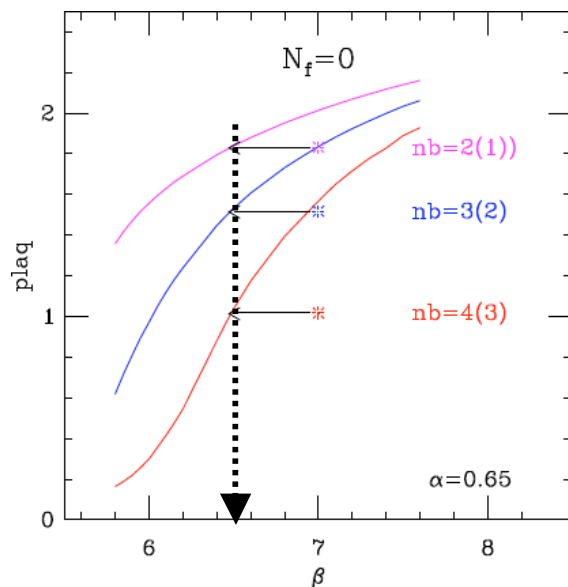


2- lattice matching MCRG - in practice:

Two actions are identical if every expectation value measured with the 2 actions are identical

MCRG identifies matched couplings (β, β') by comparing expectation values after n_b ($n_b - 1$) RG blocking steps.

Example: pure gauge SU(3)



The plaquette after 1-4 levels of blocking

$32^4 \rightarrow 16^4 \rightarrow 8^4 \rightarrow 4^4 \rightarrow 2^4$ (symbols)

$(n_b \quad 2 \quad 3 \quad 4)$

compared to

$16^4 \rightarrow 8^4 \rightarrow 4^4 \rightarrow 2^4$ (lines)

$(n_b \quad 1 \quad 2 \quad 3)$

$$\Delta \beta = \beta - \beta' = 0.51$$



Walking : not quite that simple....

The RG β function is scheme dependent:

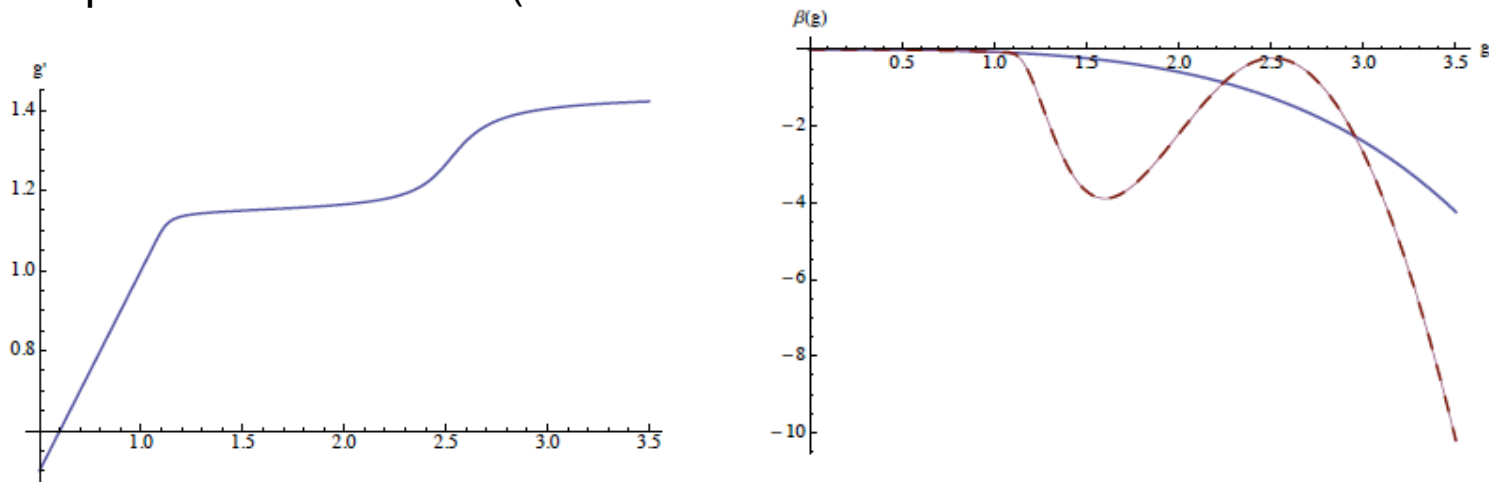
$$\beta(g) = \mu (d g / d \mu) = -\beta_0 g^3 - \beta_1 g^5 + O(g^7)$$

only β_0, β_1 are universal

Change $g \rightarrow g' = \Phi(g)$

$$\beta(g) \rightarrow \beta'(g') = (d \Phi(g) / d g) \beta(g)$$

Example from DeLDebbio (Leiden workshon)



(this is $N_f=0$ SU(3) !)

One really needs to look at physical quantities



Why walking?

2 energy scales:

- weak scale $v=250\text{GeV}$: techni-pion decay constant $F_\pi \sim 250 \text{ GeV}$
- cut-off scale (extended technicolor symm. breaking) : Λ_{ETC}

Light fermion masses :
$$m_q(\Lambda_{ETC}) \sim \frac{\langle \bar{\psi}\psi \rangle_{ETC}}{\Lambda_{ETC}^2}$$

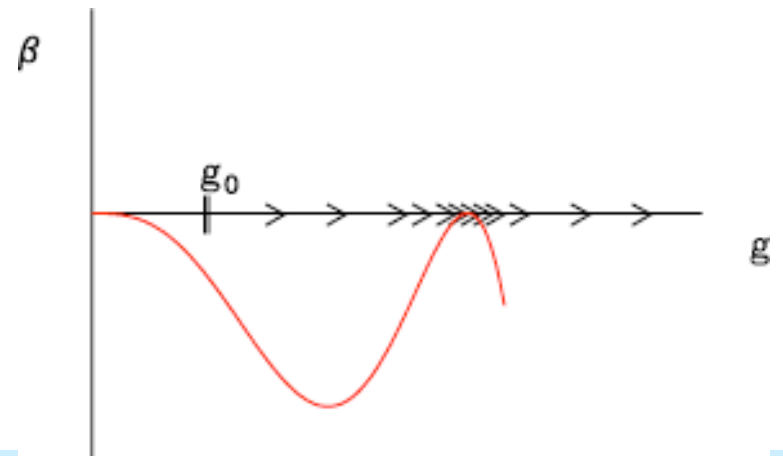
Flavor changing neutral current limits require $\Lambda_{ETC} > 10^3 \text{ TeV}$

→ large $\langle \bar{\psi}\psi \rangle_{ETC}$

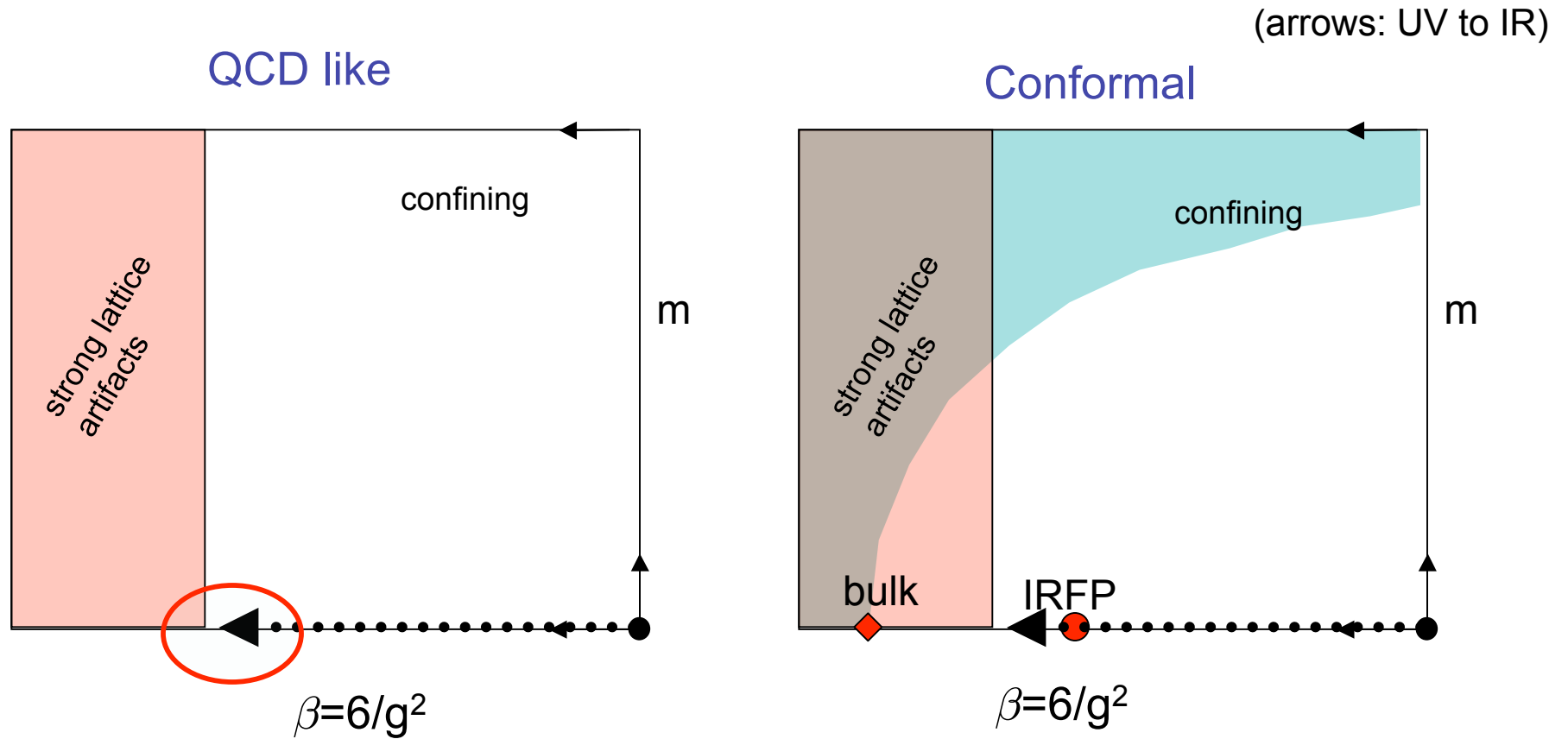
while QCD like theories have $\langle \bar{\psi}\psi \rangle_{TC} \approx 4\pi F_\pi^3$

→ small $\langle \bar{\psi}\psi \rangle_{TC}$

The problem can be solved if the running coupling and the anomalous mass remain strong across a large scale : the coupling walks



The lattice phase diagram



Lattice simulations can connect the perturbative FP and strong coupling

- Found IRFP ? Done ✓
- No IRFP? Show that it is confining before a bulk transition is reached
- Strong lattice artifacts can interfere

Calculating the step scaling function

- Schroedinger functional method:
 - generalization of the approach used to calculate the renormalized coupling in QCD
- 2-lattice matching Monte Carlo Renormalization Group method:
 - based on Wilson RG description to study critical behavior
 - has been used in QCD and statistical systems extensively
 - Has a lot of built-in checks and controls
 - works with bare quantities
 - could be used to calculate the renormalized coupling in QCD

