

Three-nucleon forces and neutron-rich nuclei

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TRIUMF

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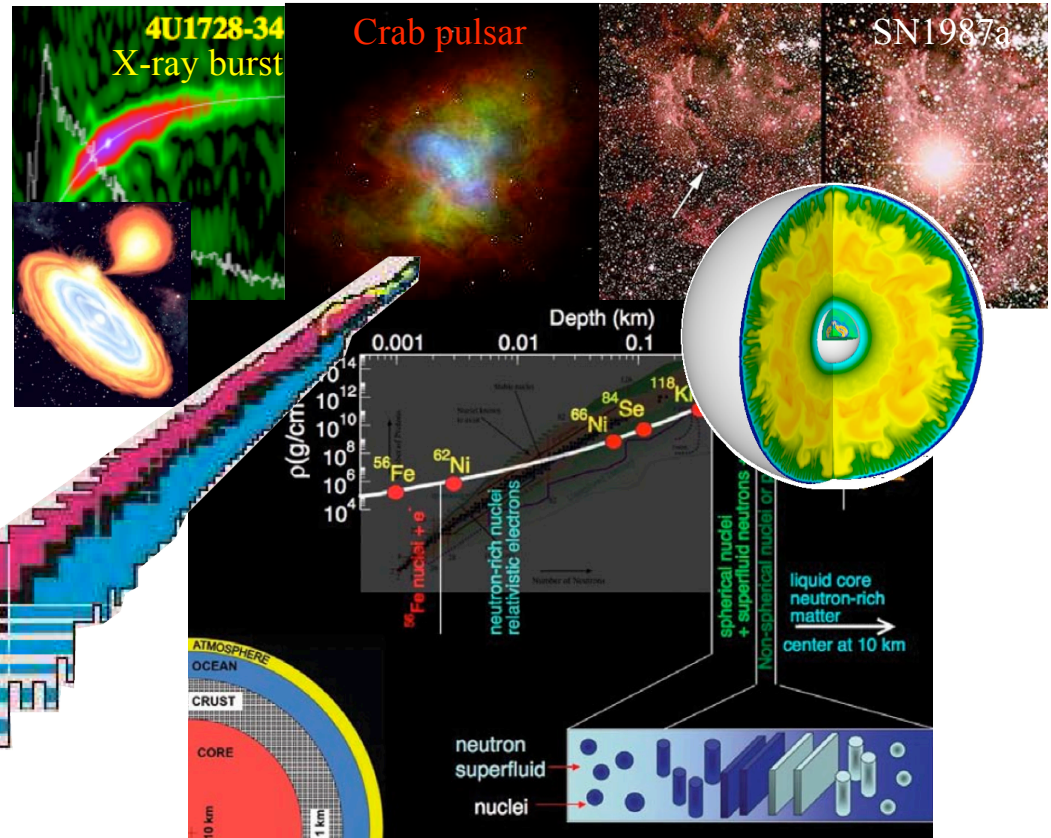
Outline

1. EFT and RG for nuclear forces
2. Three-nucleon forces in light nuclei
3. Three-nucleon forces towards heavier and neutron-rich nuclei
4. Three-nucleon forces in nuclear and neutron matter
5. Summary

Towards a unified description from light to heavy nuclei and strongly-interacting matter at the extremes

development of EFT and RG for nuclear forces

advances in ab-initio methods and in nuclear matter theory



three-nucleon forces play a central role with new qualitative effects



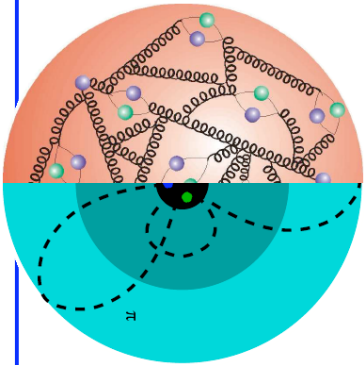
Outline

1. EFT and RG for nuclear forces

Λ / Resolution dependence of nuclear forces

with high-energy probes:
quarks+gluons

cf. scale/scheme dependence
of parton distribution functions



Lattice QCD

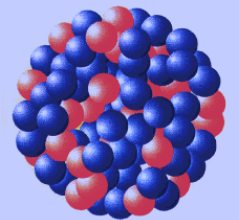
Effective theory for NN, 3N, many-N interactions and electroweak operators: resolution scale/ Λ -dependent

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$

Λ_{chiral}

momenta $Q \sim \lambda^{-1} \sim m_{\pi}$: chiral EFT - typical momenta in nuclei

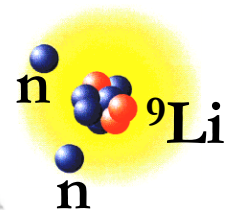
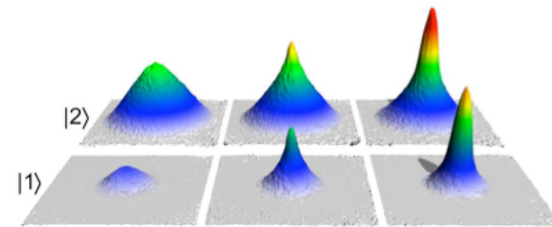
nucleons interacting via pion exchanges and shorter-range contact interactions



$\Lambda_{\text{pionless}}$

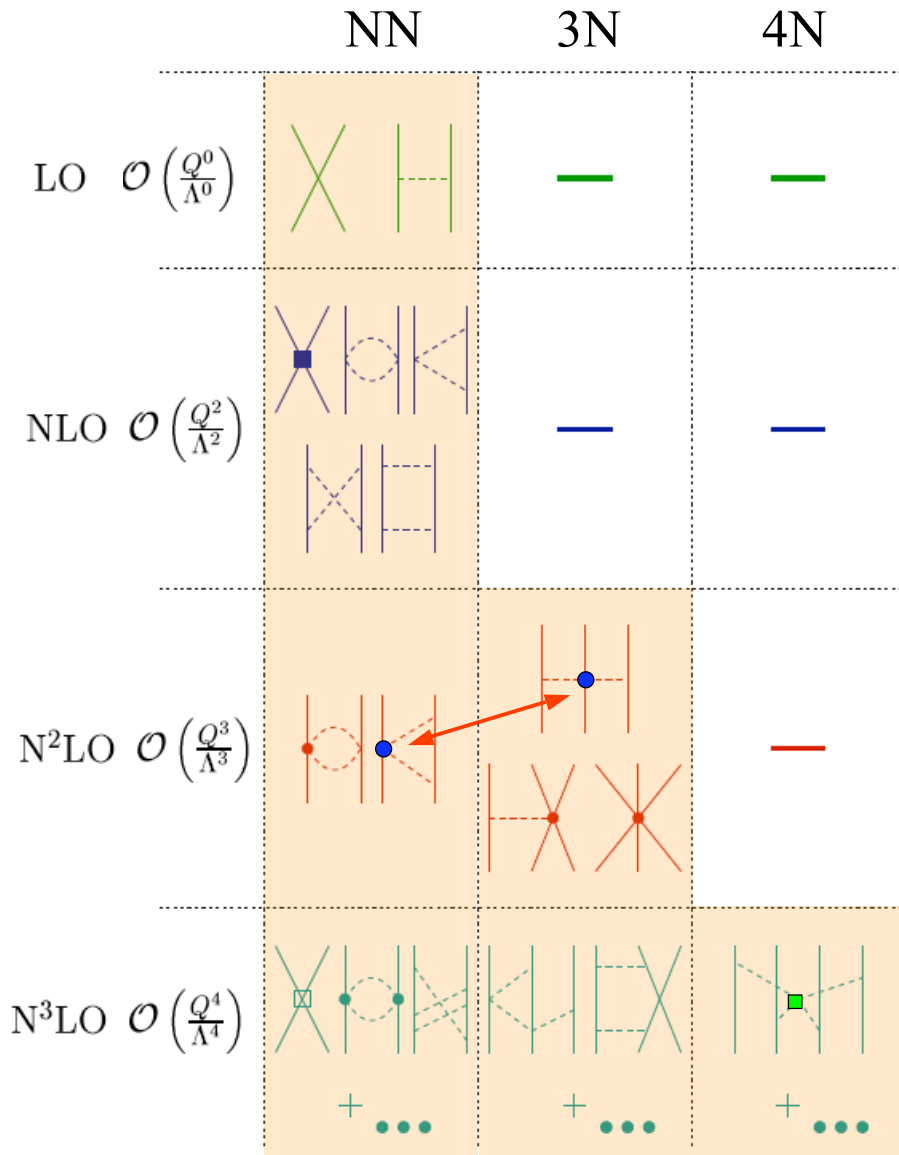
$Q \ll m_{\pi} = 140 \text{ MeV}$: pionless EFT

large scattering length physics and non-universal corrections



Chiral EFT for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



explains pheno hierarchy:

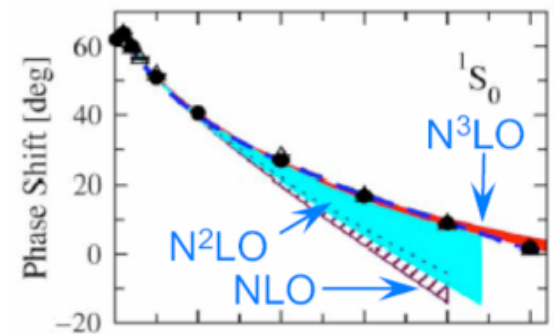
NN > 3N > 4N > ...

consistent NN-3N, π N, $\pi\pi$,
electroweak operators

3N,4N: 2 new couplings to N³LO

resolution/ Λ -dep. contact interactions

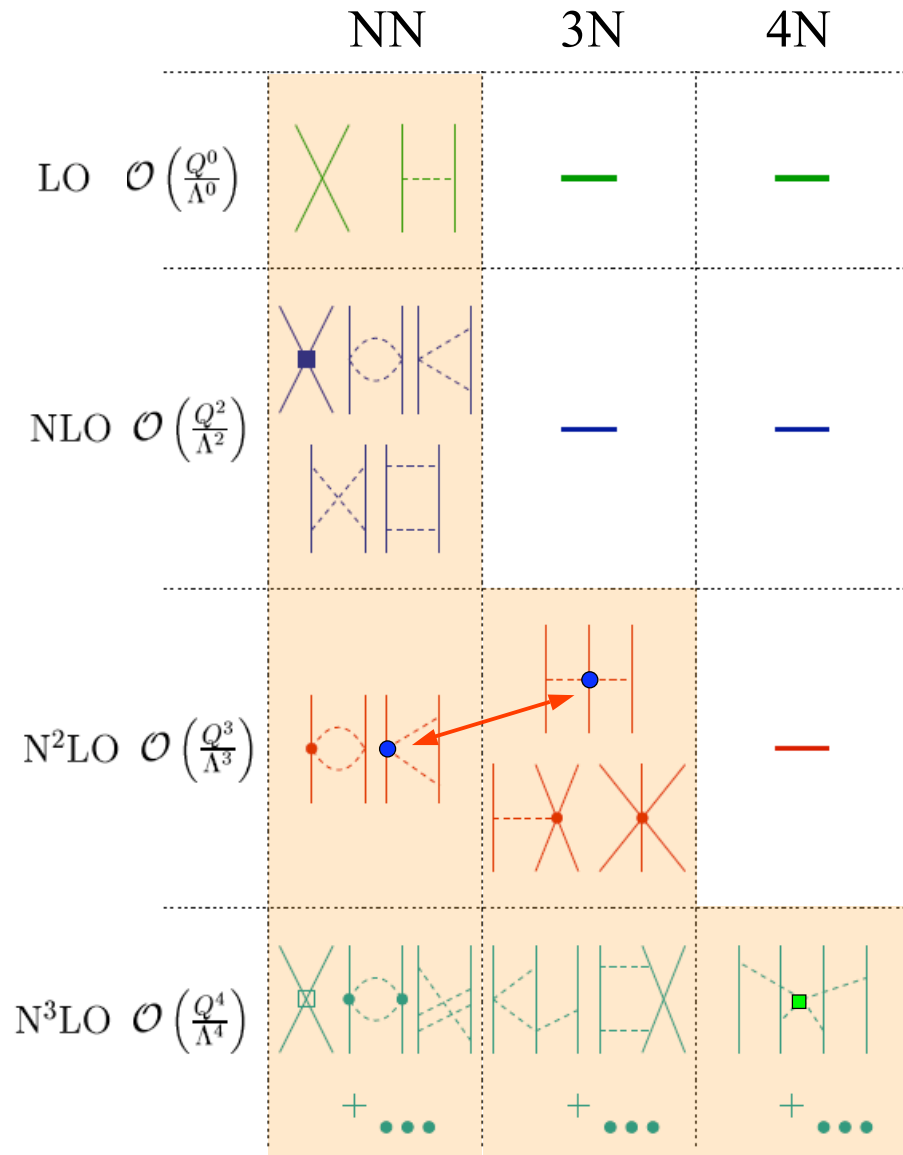
error estimates from truncation order
and resolution/ Λ -variation



Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner, Nogga, Machleidt, ...

Chiral EFT for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV



Open questions:

Power counting with singular pion exchanges, tensor parts $\sim 1/r^n$

promotion of contact interactions

Nogga, Timmermans, van Kolck (2005)

Delta-full $m_\Delta - m_N \sim m_\pi \rightarrow$ 3N at NLO vs. Delta-less EFT

Counting of $1/m$ corrections,

$m \sim \Lambda^2$ or Λ , could be important for A_y

This talk: 3N forces from N²LO

Delta-less chiral EFT

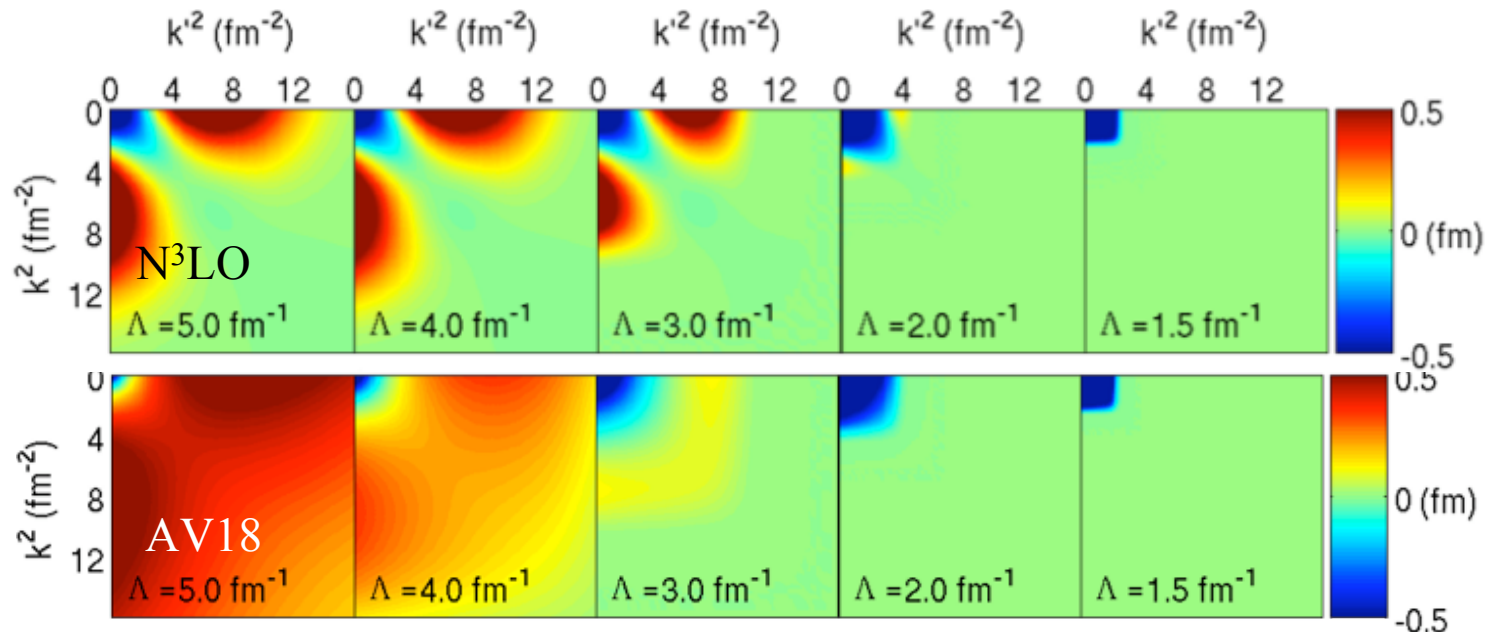
Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner, Nogga, Machleidt,...

Low-momentum interactions from the Renormalization Group

RG evolution to lower resolution/cutoffs

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$

exact RG for NN interactions Bogner, Kuo, AS, Brown, Furnstahl,... (2003...)



\approx universal low-momentum interaction $V_{\text{low } k}(\Lambda)$, sharp/smooth regulators

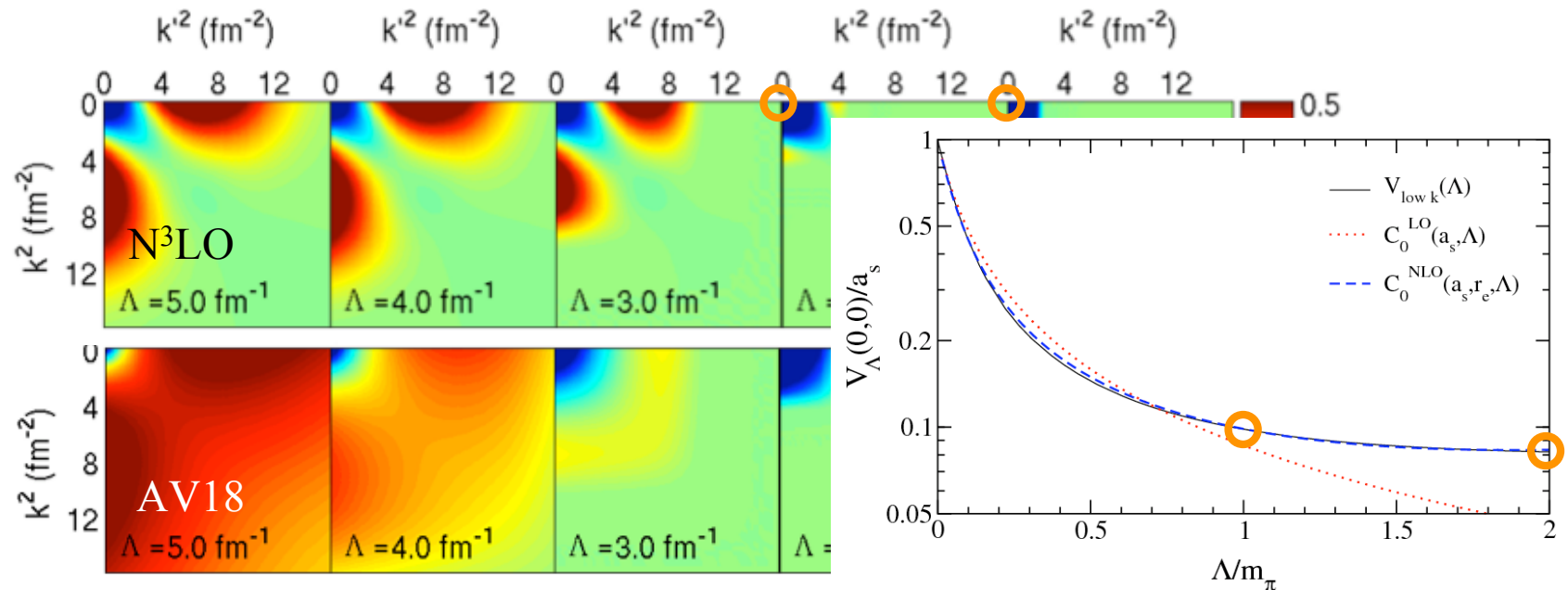
RG evolution decouples high momenta (short-range repulsion and tensor parts)

Low-momentum interactions from the Renormalization Group

RG evolution to lower resolution/cutoffs

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$

exact RG for NN interactions **Bogner, Kuo, AS, Brown, Furnstahl, ... (2003...)**



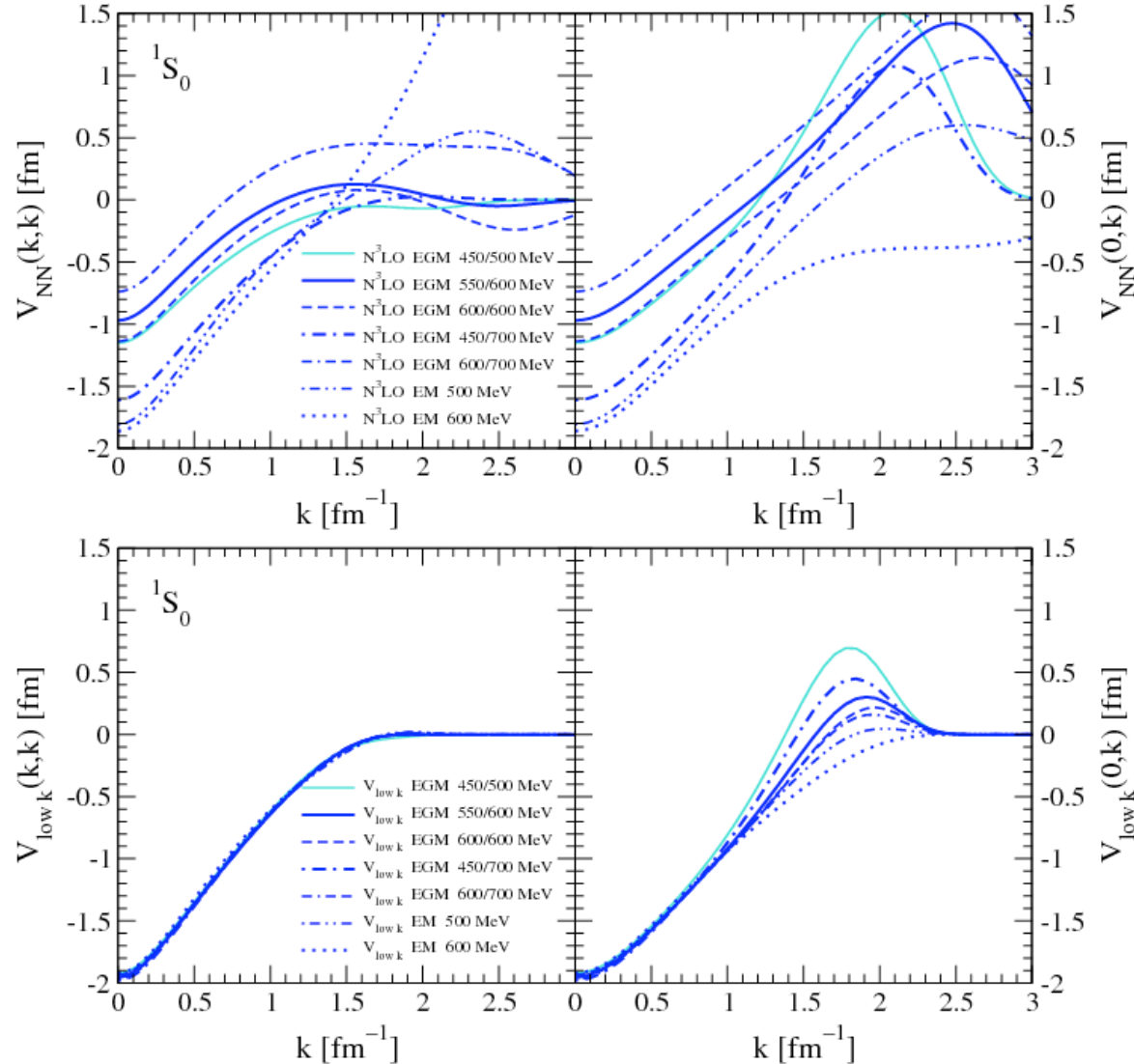
\approx **universal** low-momentum interaction $V_{\text{low } k}(\Lambda)$, sharp/smooth regulators

RG evolution decouples high momenta (short-range repulsion and tensor parts)

same RG evolution as in EFT, but without expansion on operators

$V_{\text{low } k}(0,0;\Lambda)$ follows contact interaction $c_0(\Lambda)$ at NLO in pionless EFT regime

Chiral EFT and RG



\approx **universality** from different $N^3\text{LO}$ potentials

RG preserves long range and generates higher-order contact interactions

Weinberg eigenvalue diagnostic

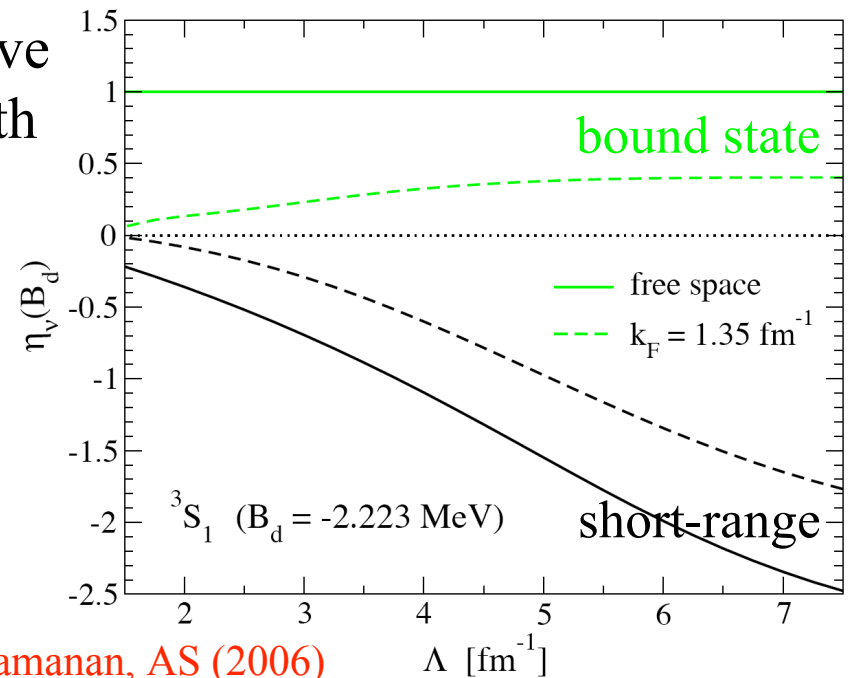
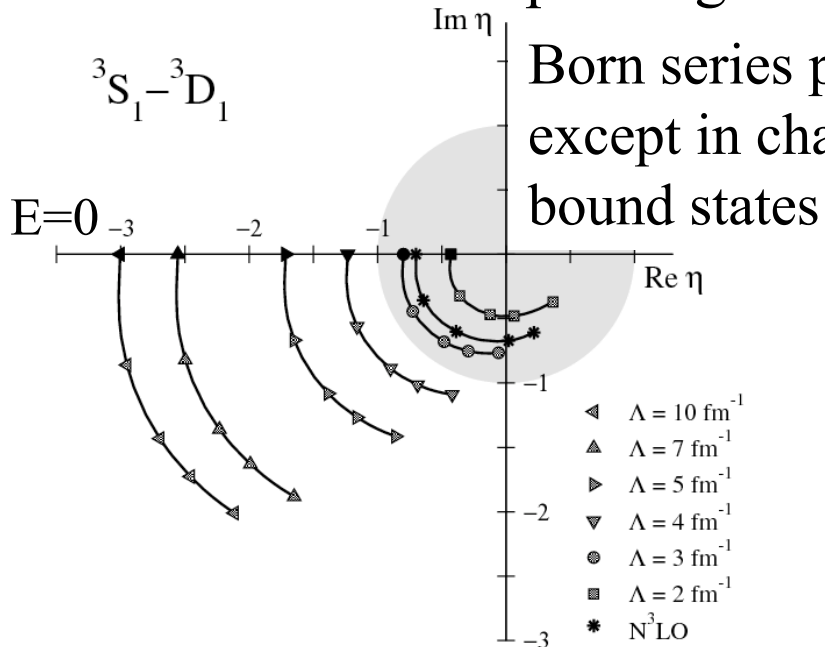
study spectrum of $G_0(z)V |\Psi_\nu(z)\rangle = \eta_\nu(z) |\Psi_\nu(z)\rangle$ at fixed energy z

governs convergence $T(z) |\Psi_\nu(z)\rangle = (1 + \eta_\nu(z) + \eta_\nu(z)^2 + \dots) V |\Psi_\nu(z)\rangle$

can write as Schrödinger equation $(H_0 + \frac{1}{\eta_\nu(z)} V) |\Psi_\nu(z)\rangle = z |\Psi_\nu(z)\rangle$

high momenta/large cutoffs lead to flipped-potential bound states of $-\lambda V$ for small λ /large $\eta \rightarrow$ strong coupling to high momenta/short range and Born series always nonperturbative

RG evolution decouples high momenta (short-range repulsion and tensor parts)



Bogner, AS, Furnstahl, Nogga (2005), Bogner, Furnstahl, Ramanan, AS (2006)

Advantages of low-momentum interactions for nuclei

lower cutoffs need smaller basis

also via similarity RG (SRG)

evolution towards band diagonal

Bogner, Furnstahl, Perry, Roth, Hergert,... (2007...)

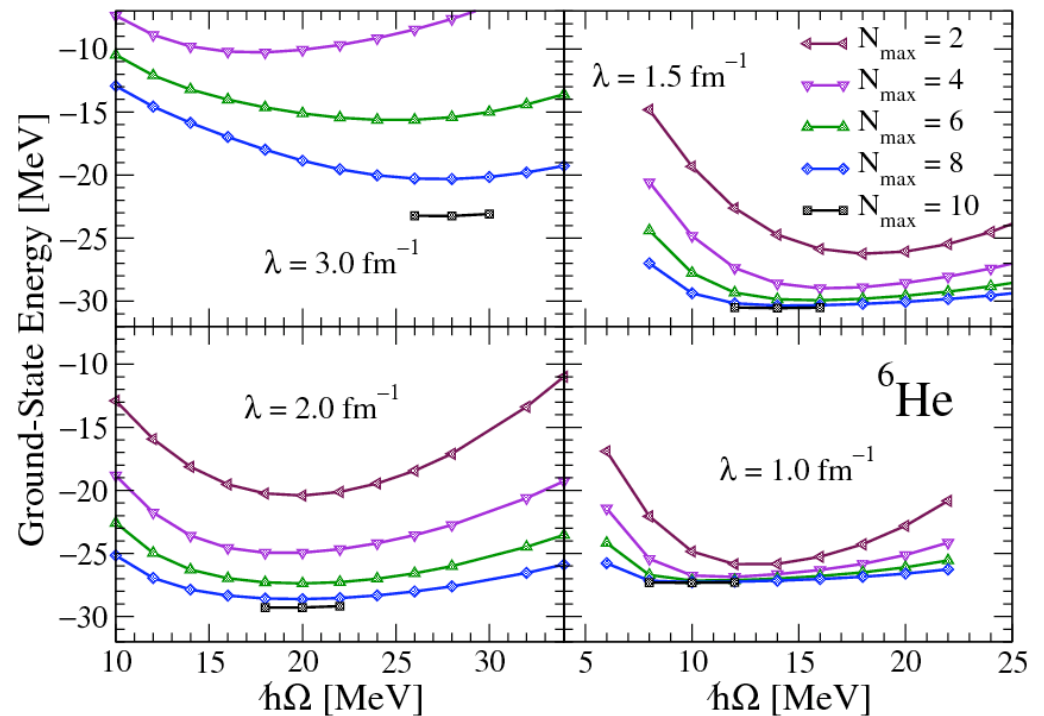
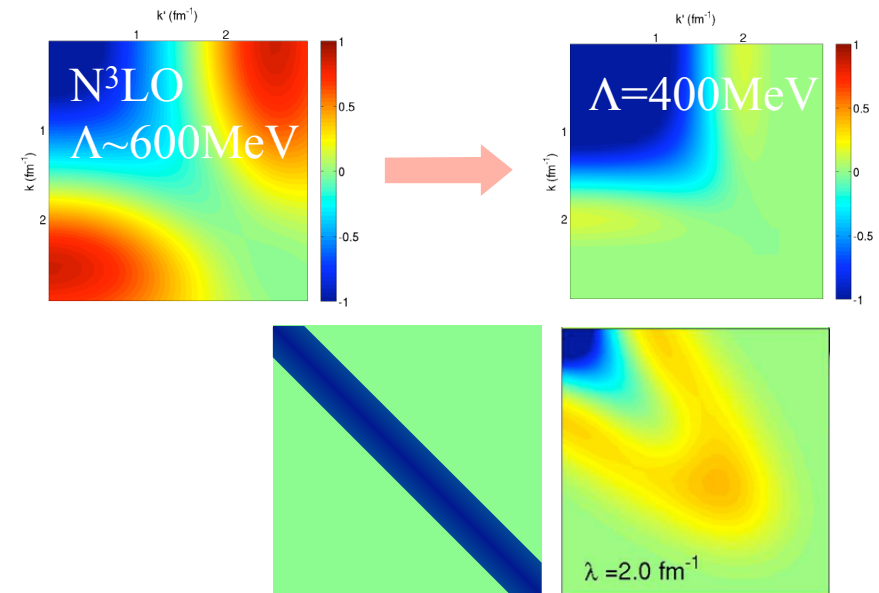
connection to EFT?

improved convergence for nuclei

Bogner Furnstahl, Maris, Perry, AS, Vary (2008)

10^3 states for $N_{\max}=2$ vs.

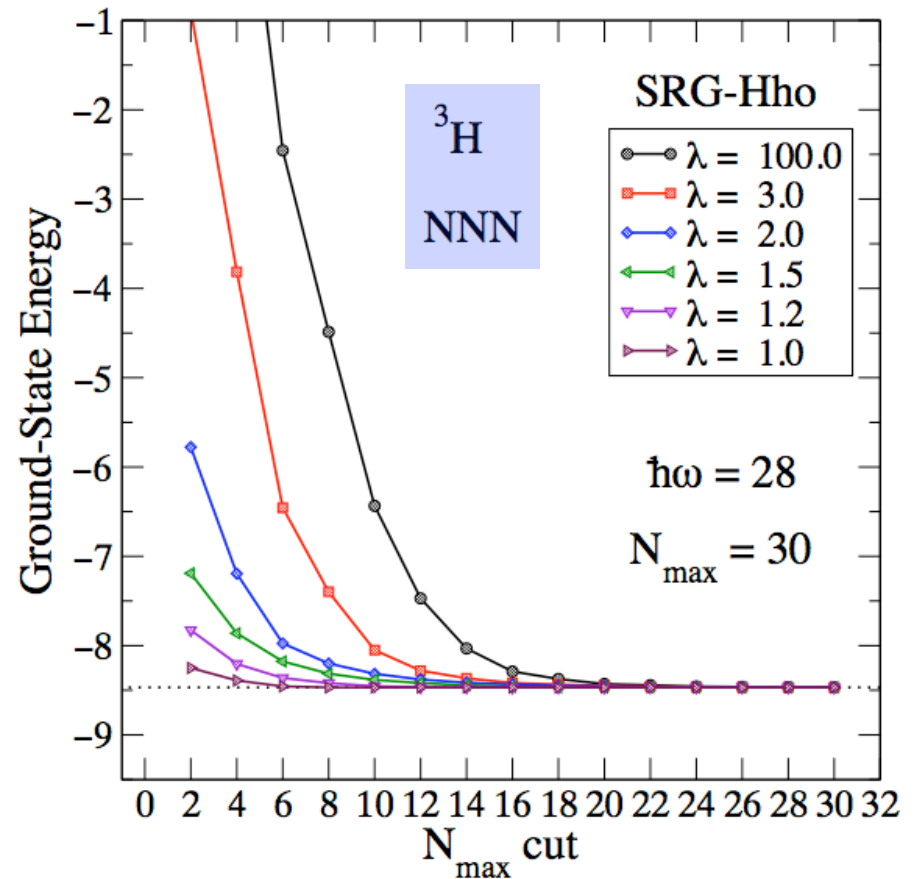
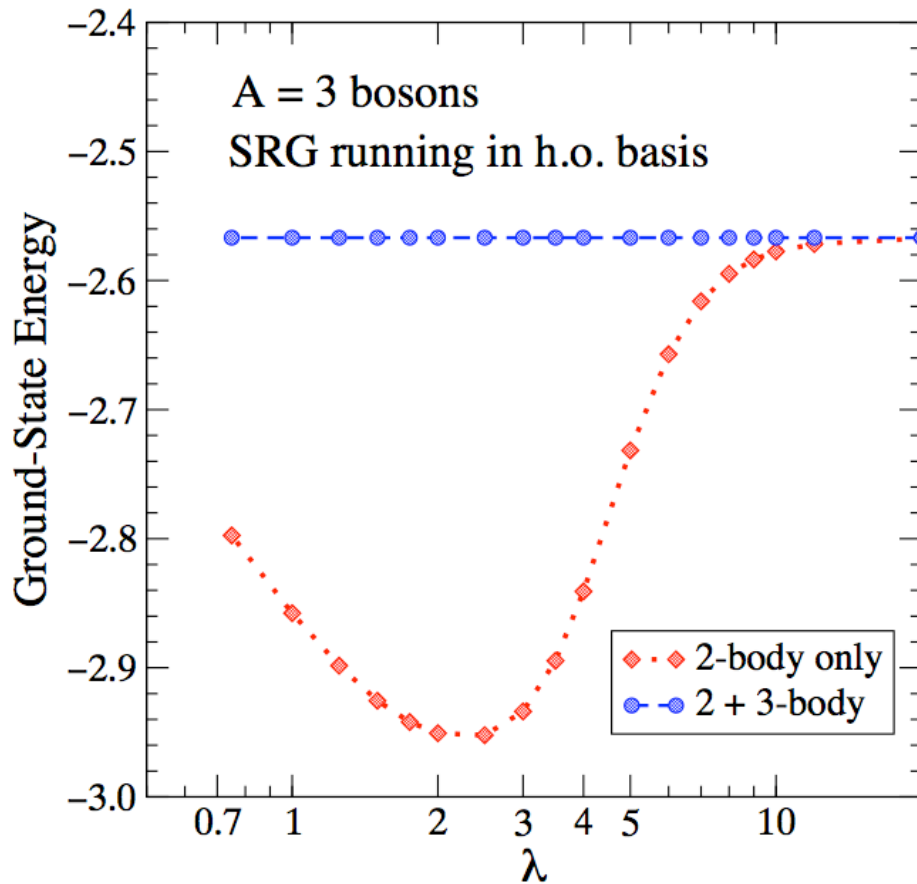
10^7 states for $N_{\max}=10$



Towards evolving 3N interactions

SRG evolution with harmonic oscillator Hamiltonian as generator for bosons in 1d and for chiral 3N interactions

Jurgenson, Furnstahl (2008) and from Eric Jurgenson

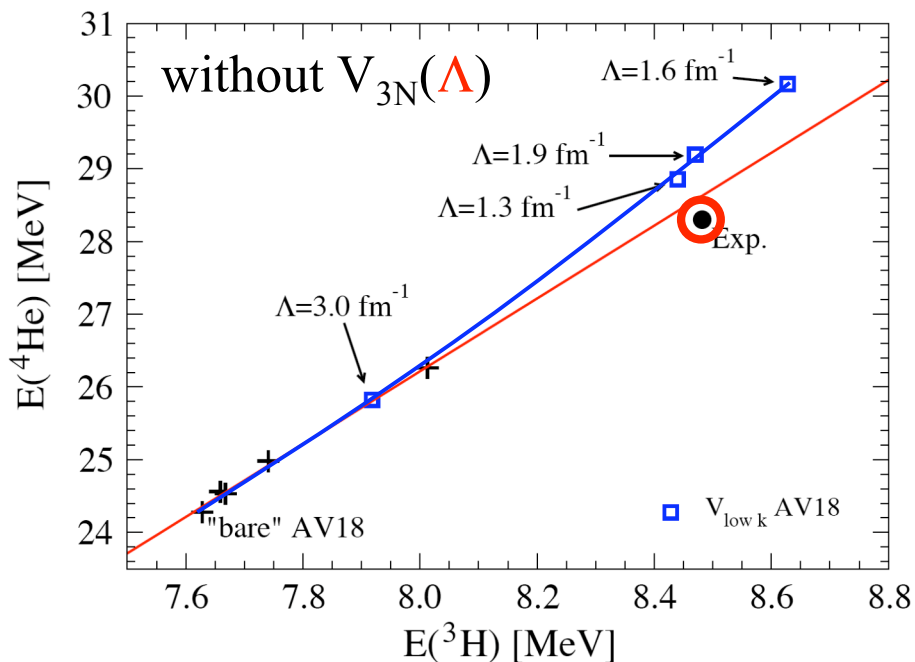


This talk: chiral EFT is complete basis, $V_{3N}(\Lambda)$ fits for lower cutoffs

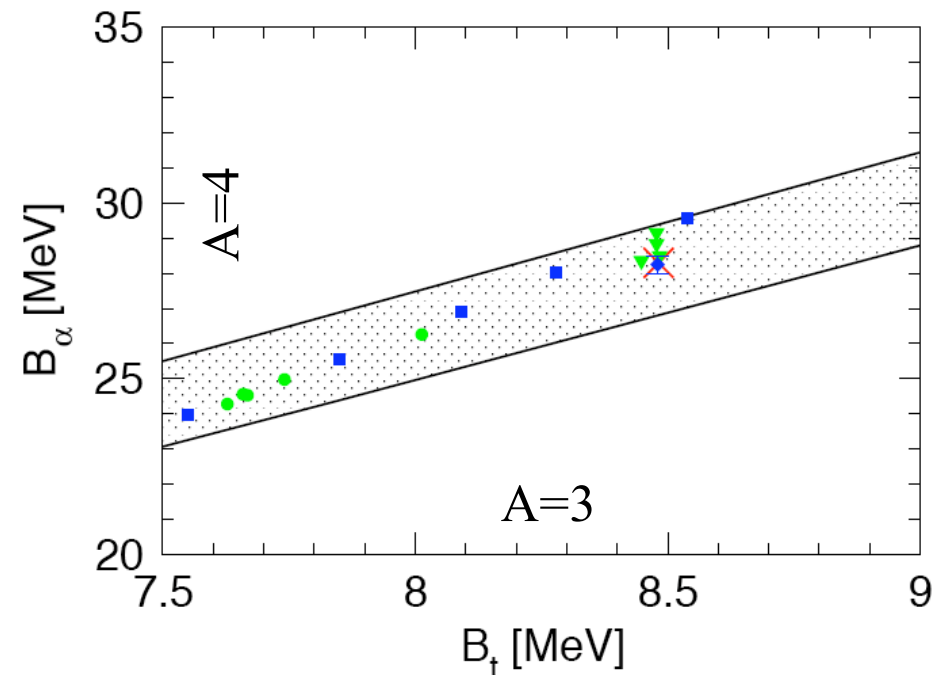
3N interactions required for renormalization

$V_{\text{NN}}(\Lambda)$ defines NN interactions with cutoff-independent NN observables
 cutoff variation estimates errors due to neglected parts in $H(\Lambda)$

cutoff dependence explains “Tjon line” \rightarrow 3N for renormalization



Nogga, Bogner, AS (2004)



pionless EFT: Platter, Hammer, Meissner (2005)

large scattering lengths drive correlation

Tjon lines in $^{16}\text{O}_{\pm 1}$ Hagen et al., in prep.

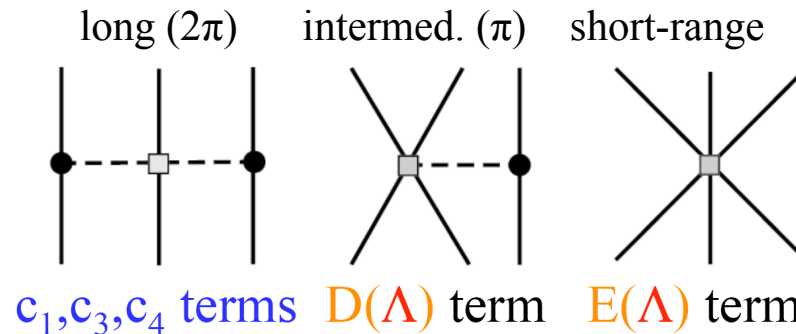
A scenic winter landscape featuring a dense forest of snow-laden evergreen trees in the foreground. The trees are heavily coated in white snow, creating a textured, white canopy. In the background, a range of rugged mountains is visible, with some peaks partially obscured by a light mist or low clouds. The sky is a pale, overcast blue. The overall scene is peaceful and serene, typical of a high-altitude winter environment.

Outline

2. Three-nucleon forces in light nuclei

Chiral EFT 3N interactions

leading N²LO $\sim (Q/\Lambda)^3$ van Kolck (1994), Epelbaum et al. (2002)



c_i from πN , consistent with NN $c_1 = -0.9^{+0.2}_{-0.5}$, $c_3 = -4.7^{+1.2}_{-1.0}$, $c_4 = 3.5^{+0.5}_{-0.2}$ Meissner et al.

c_3, c_4 important for structure, large uncertainties at present

$V_{3N}(\Lambda)$ based on fits of D, E couplings to $A=3,4$ for range of cutoffs

chiral EFT is a complete low-mom basis \rightarrow 3N up to truncation errors

D term can be fixed by tritium beta decay Gardestig, Phillips (2006), Gazit et al. (2008)

generally improves 3N scattering Gloeckle, Kamada, Witala,.... (2002...)

Subleading chiral EFT 3N interactions

parameter-free N³LO from Epelbaum; Bernard et al. (2007), Ishikawa, Robilotta (2007)

- 1/m-corrections to 1 insertion from $\mathcal{L}_{1/m}^{(2)} = \text{---} \blacksquare \text{---} + \text{---} \blacksquare \text{---} + \text{---} \blacksquare \text{---} + \mathcal{O}(\pi^3)$

— rich operator structure (includes spin-orbit interactions)

- 1-loop diagrams with all vertices from $\mathcal{L}_{\text{eff}}^{(0)}$

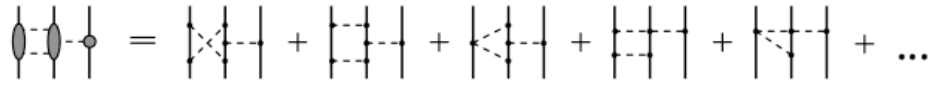
2π - exchange



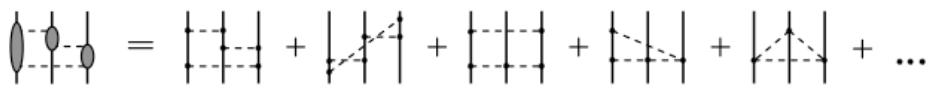
The calculated corrections simply shift the LECs c_i as follows:

$$\delta c_1 = \frac{g_A^2 M_\pi}{64\pi F_\pi^2} \sim 0.13 \text{ GeV}^{-1} \quad \delta c_3 = \frac{3g_A^4 M_\pi}{16\pi F_\pi^2} \sim 2.5 \text{ GeV}^{-1} \quad \delta c_4 = -\frac{g_A^4 M_\pi}{16\pi F_\pi^2} \sim -0.85 \text{ GeV}^{-1}$$

2π-1π - exchange



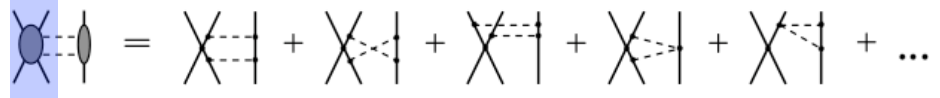
ring diagrams



contact-1π - exchange



contact-2π - exchange



3N interactions involving NN contacts in progress

Epelbaum et al.

Lattice QCD and 3N forces

Long-range couplings:

pion-NN coupling g_A from full QCD

Edwards et al. (2006)

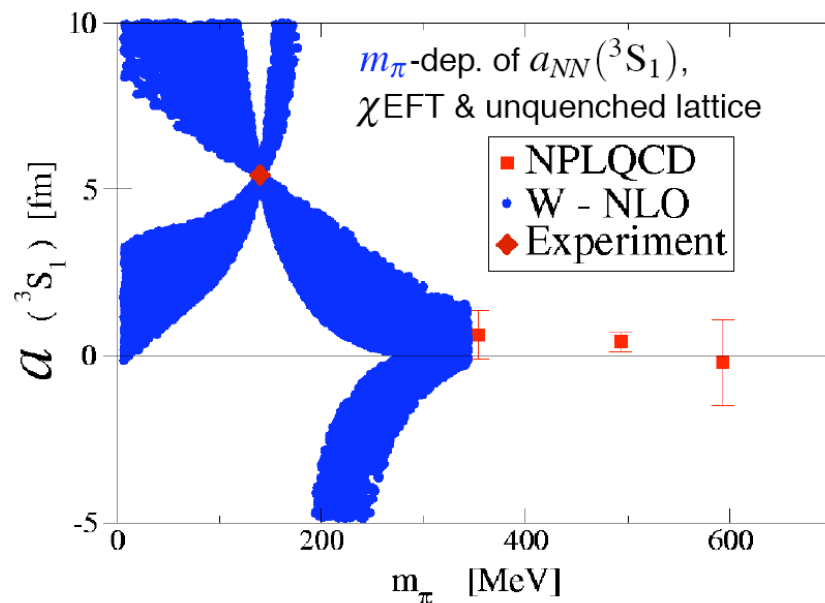
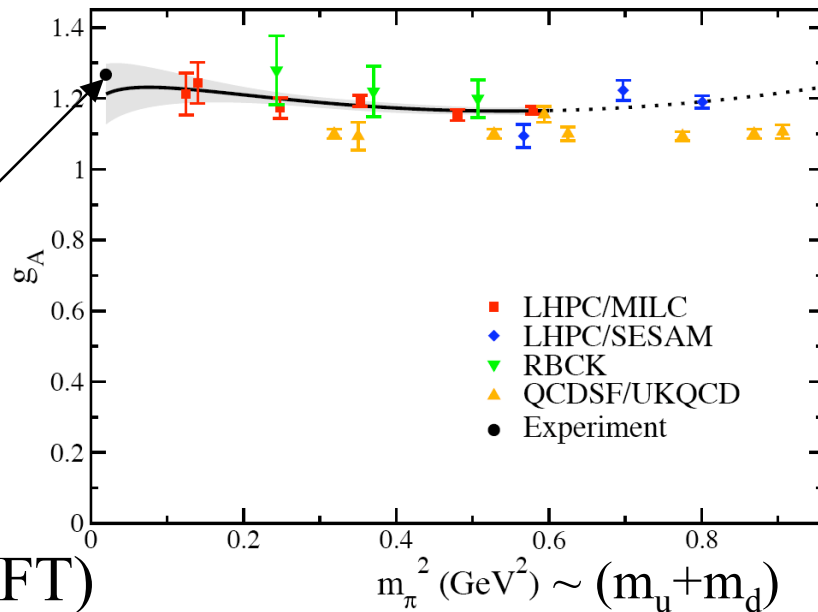
chiral EFT extrapolation to physical pion mass agrees with experiment

constrain higher-order long-range pion-, N-,... couplings (scheme: chiral EFT)
 c_i from lattice QCD?

Few-nucleon observables:

NN scattering lengths from full QCD, dependence on quark masses Beane et al. (2006)

Constrain experimentally difficult observables: 3-neutron properties
first steps: 3 pions on a lattice Detmold et al. (2008)



Low-momentum 3N fits

fit D, E couplings to ${}^3\text{H}$, ${}^4\text{He}$ binding energies
for range of cutoffs

linear dependences in fits to triton binding

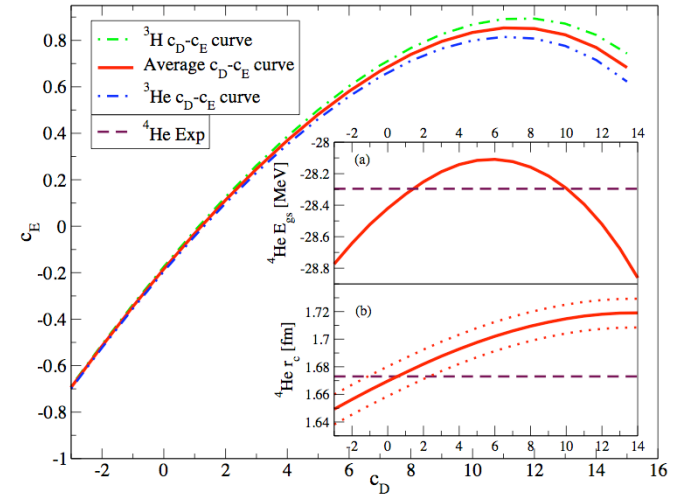
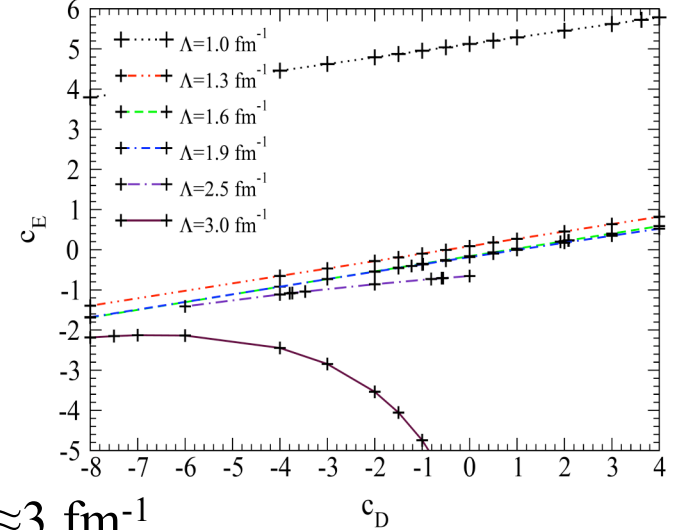
3N interactions perturbative for $\Lambda \lesssim 2 \text{ fm}^{-1}$

Nogga, Bogner, AS (2004)

nonperturbative at larger cutoffs, cf. chiral EFT $\Lambda \approx 3 \text{ fm}^{-1}$

new 3N fits to ${}^3\text{H}$ binding energy, ${}^4\text{He}$ radius

Bogner, Furnstahl, Nogga, AS, 0903.3366.



Navratil et al. (2007)

Λ or $\lambda/\Lambda_{3\text{NF}}$ [fm^{-1}]	$V_{\text{low } k}$		SRG	
	c_D	c_E	c_D	c_E
1.8/2.0	-0.0112	-0.2212		
2.0/2.0	-0.3000	-0.2761	-1.023	-0.3397
2.0/2.5	-2.000	-0.7564	-2.991	-0.8797
2.2/2.0	-0.9000	-0.3673		
2.8/2.0	-1.552	-0.4058		

TABLE I: Results for the c_D and c_E couplings fit to $E_{3\text{H}} = -8.482 \text{ MeV}$ and $r_{4\text{He}} = 1.46 - 1.47 \text{ fm}$ for the NN/3N cutoffs used here. For $V_{\text{low } k}$ (SRG) interactions, the 3NF fits lead to $E_{4\text{He}} = -28.22 \dots 28.45 \text{ MeV}$ ($-28.53 \dots 28.71 \text{ MeV}$).

perturbative in $A=3,4$ for these cutoffs and generally $\Lambda_{3\text{NF}} \lesssim \Lambda$

Low-momentum 3N interactions in light nuclei

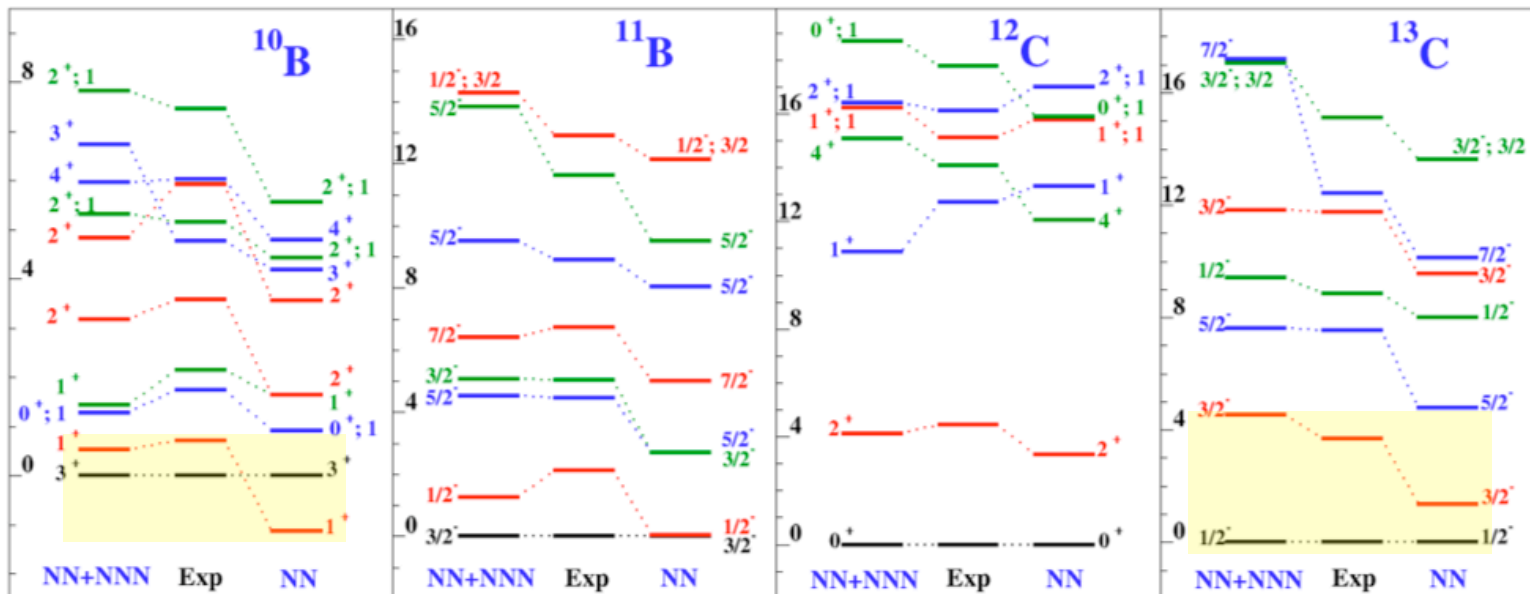
Nogga, Bogner, AS (2004)

Λ	${}^3\text{H}$					${}^4\text{He}$					max $ V_{3N}/V_{\text{low } k} $	${}^4\text{He}$ k_{rms}
	T	$V_{\text{low } k}$	c -terms	D -term	E -term	T	$V_{\text{low } k}$	c -terms	D -term	E -term		
1.0	21.06	-28.62	0.02	0.11	-1.06	38.11	-62.18	0.10	0.54	-4.87	0.08	0.55
1.3	25.71	-34.14	0.01	1.39	-1.46	50.14	-78.86	0.19	8.08	-7.83	0.10	0.63
1.6	28.45	-37.04	-0.11	0.55	-0.32	57.01	-86.82	-0.14	3.61	-1.94	0.04	0.67
1.9	30.25	-38.66	-0.48	-0.50	0.90	60.84	-89.50	-1.83	-3.48	5.68	0.06	0.70
2.5(a)	33.30	-40.94	-2.22	-0.11	1.49	67.56	-90.97	-11.06	-0.41	6.62	0.12	0.74
2.5(b)	33.51	-41.29	-2.26	-1.42	2.97	68.03	-92.86	-11.22	-8.67	16.45	0.18	0.74
3.0(*)	36.98	-43.91	-4.49	-0.73	3.67	78.77	-99.03	-22.82	-2.63	16.95	0.23	0.80

natural size of individual 3N expectation values ~ 0.1 of NN,
consistent with chiral power-counting estimates

long-range c_i -terms attractive in $A=3,4$ for most cutoffs,
but repulsive in nuclear/neutron matter and drive saturation

3N interactions and nuclear structure



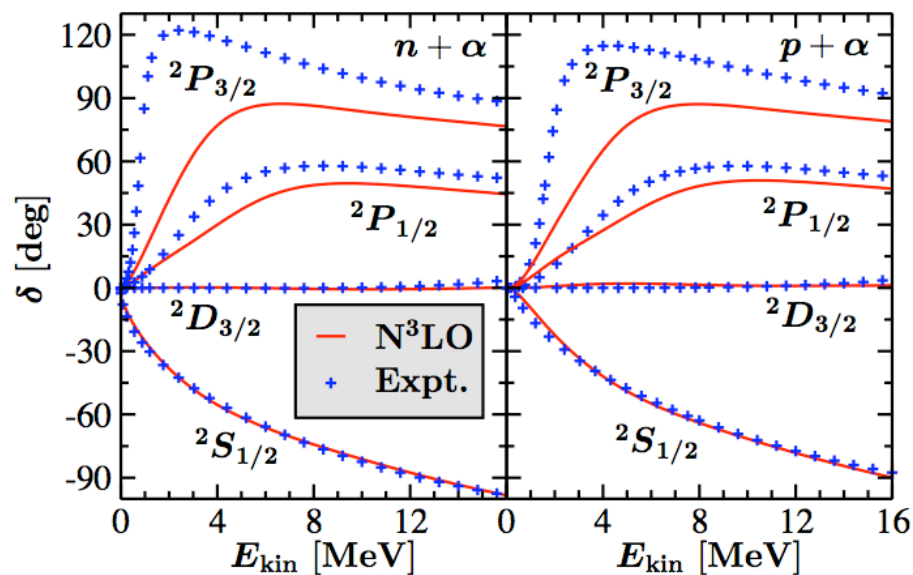
Navratil et al. (2007)

highlights importance of 3N forces

crucial for spin-orbit splittings

spin-orbit shell closures ($^{48}\text{Ca}, \dots$)

see, e.g., Zuker, AS (2006)



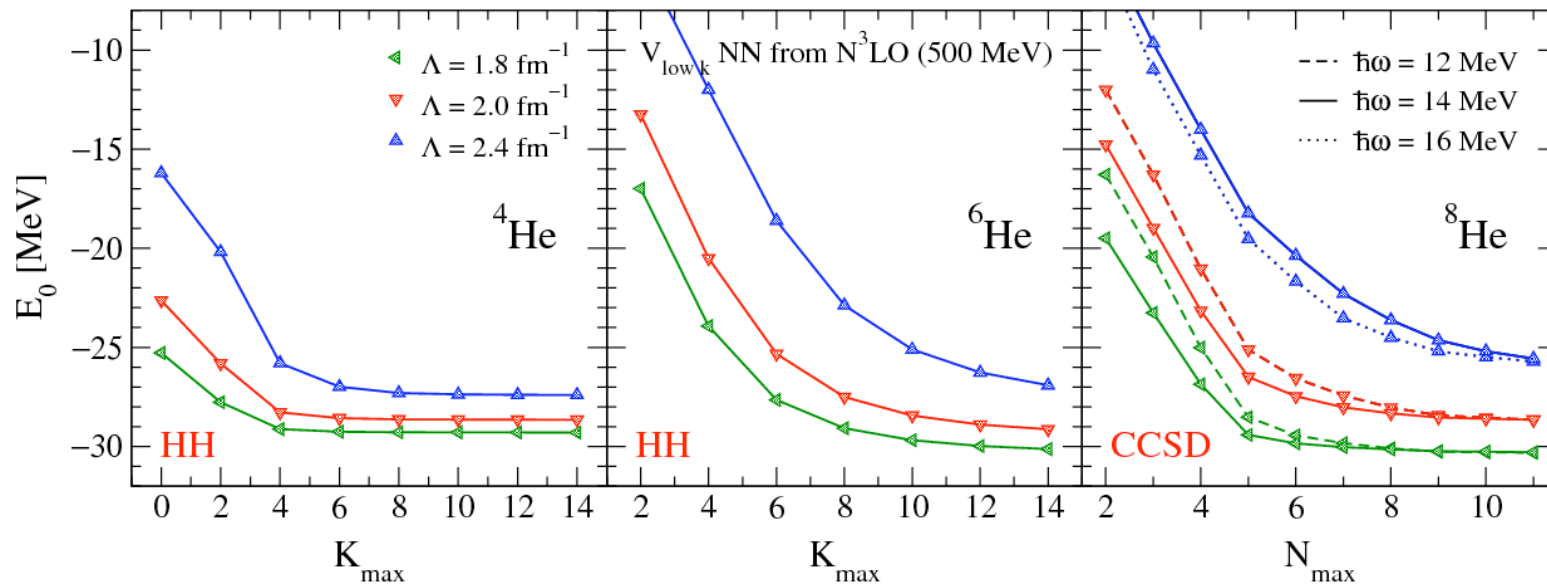
Quaglioni, Navratil (2008)

Helium halos

Hyperspherical Harmonics for ${}^6\text{He}$ and Coupled-Cluster theory for ${}^8\text{He}$

Bacca et al., 0902.1696; see talk by Sonia Bacca.

based on $N^3\text{LO}$ NN potential and RG,
describe weakly-bound nuclei with correct asymptotics



Λ [fm^{-1}]	$E_0({}^4\text{He})$ [MeV]	$E_0({}^6\text{He})$ extrapolated [MeV]	$E_0({}^8\text{He})$ Λ -CCSD(T) [MeV]
1.8	-29.30	-30.23(7)	-31.20
2.0	-28.65	-29.28(2)	-29.84
2.4	-27.40	-27.99(40)	-25.71 (CCSD)
experiment	-28.296	-29.268 [43]	-31.395 [4]

need attractive part of $3N$ interaction in neutron-rich Helium halos

Long-term goal

EFT/RG: more accurate with higher orders and cutoff variation estimates theoretical uncertainties

apply to pivotal matrix elements needed to constrain beyond Standard Model physics

isospin-symmetry breaking corrections for superallowed beta decay

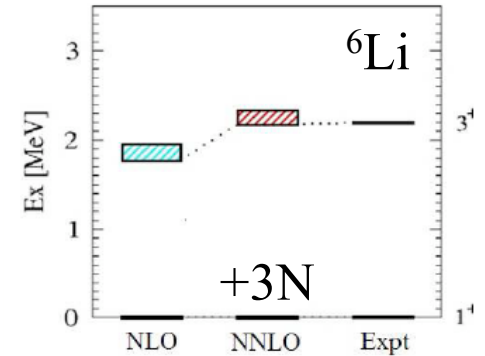
new 2008 Towner-Hardy calculation of ISB corrections (~enlarged model space) 3.1 σ shift in world-average Ft, 1.5 σ in V_{ud} largest changes in 20 years

^{62}Ga : 0.04% precise ft value ~lighter A

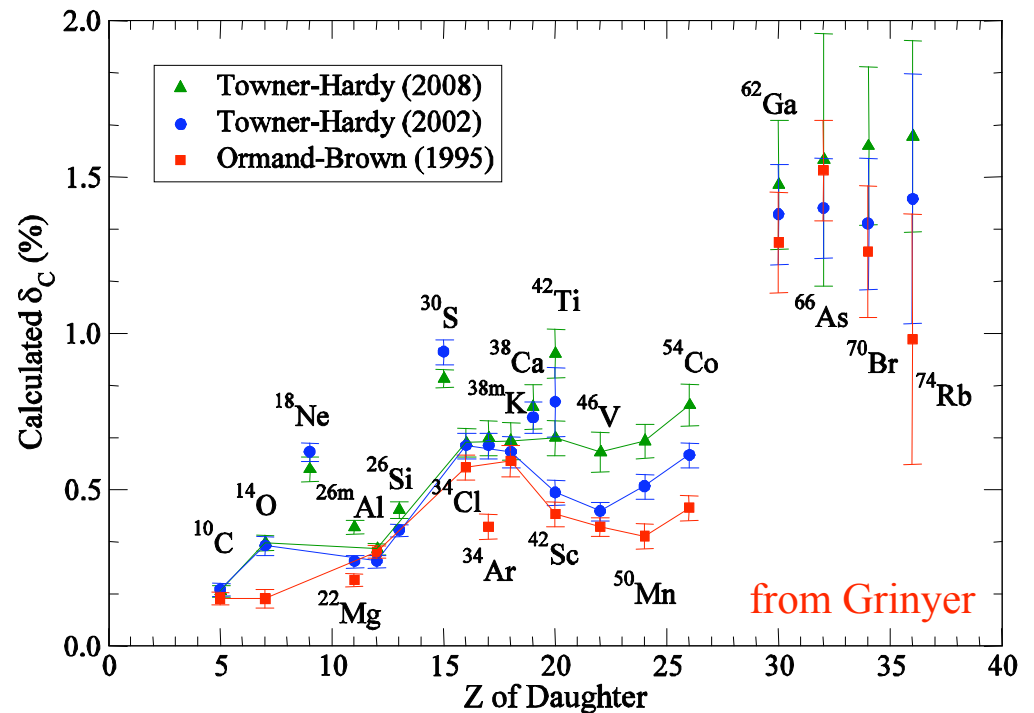
Miller, AS (2008)

Towner-Hardy formalism violates isospin commutation relations, cannot separate ISB corrections into radial wf + SM space parts

nuclear matrix elements for $0\nu\beta\beta$ decay, octupole EDM enhancement,... limited by nuclear theory uncertainties



from Nogga



from Grinyer

A scenic winter landscape featuring a dense forest of snow-laden evergreen trees in the foreground. In the background, a range of rugged mountains is visible under a cloudy sky, with a body of water nestled between the peaks. The overall atmosphere is serene and cold.

Outline

3. Three-nucleon forces towards heavier and neutron-rich nuclei

Pushing the limits

First ab-initio calculations for heavier systems:
Coupled-cluster theory based on $V_{\text{low } k}(\Lambda)$

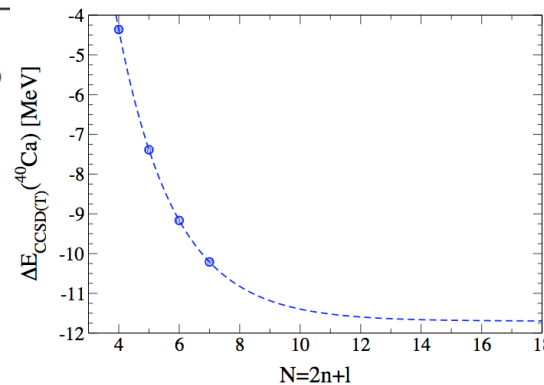
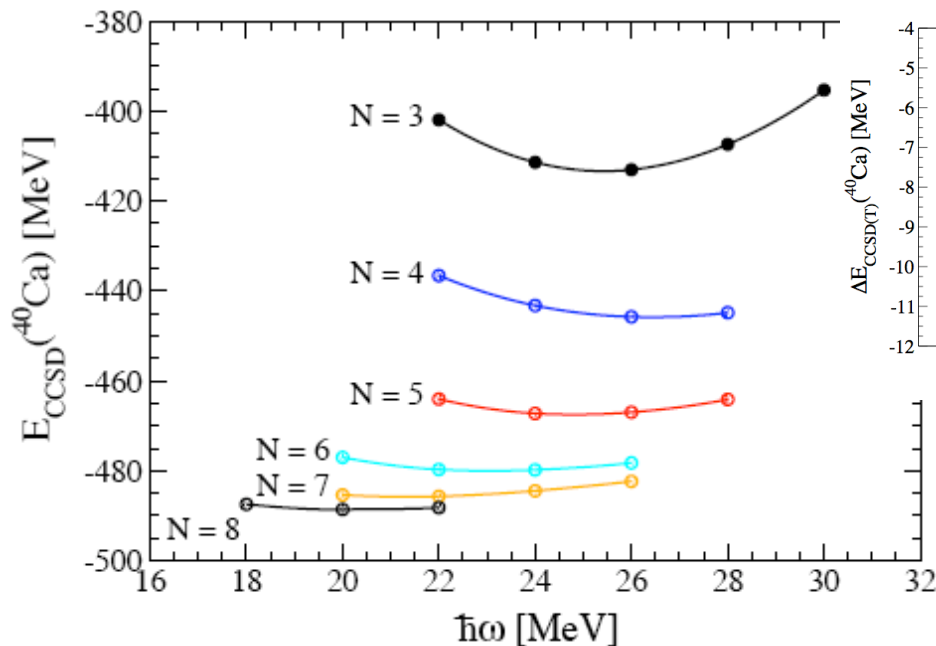
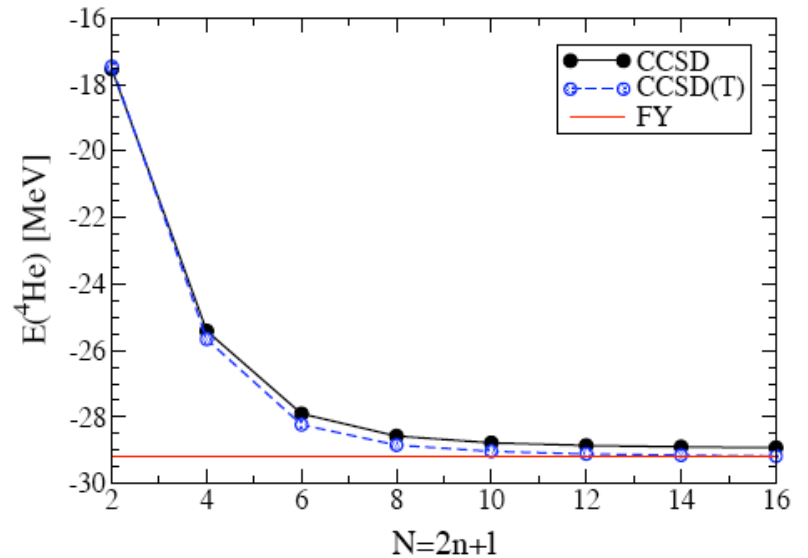
Hagen, Dean, Hjorth-Jensen, Papenbrock, AS (2007)

meets and sets benchmarks:

within 10 keV of exact FY for ${}^4\text{He}$

accurate for ${}^{16}\text{O}$ and ${}^{40}\text{Ca}$

$N=8$: basis dimension $\sim 10^6$



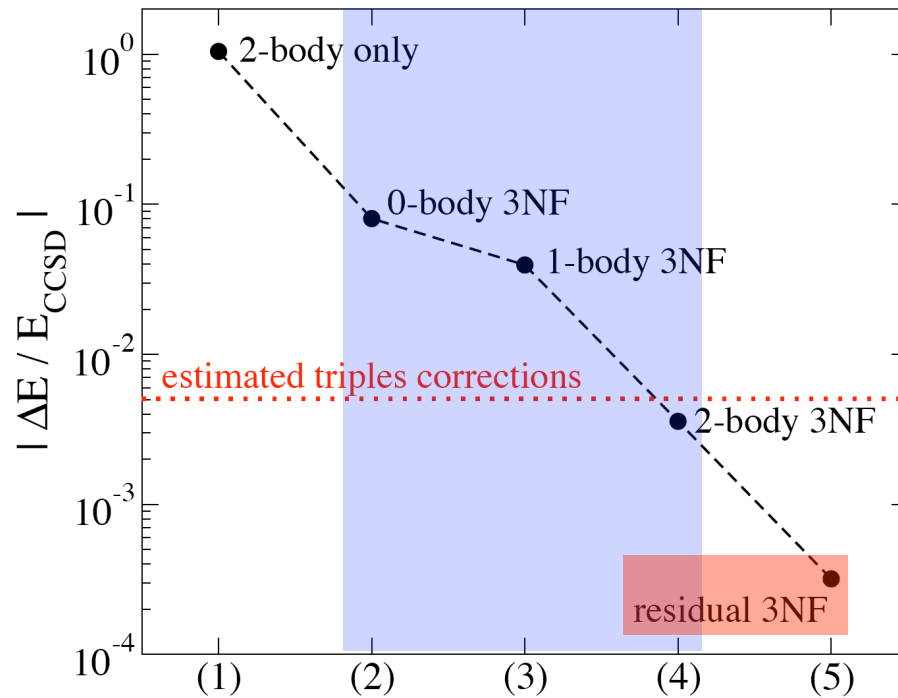
	${}^4\text{He}$	${}^{16}\text{O}$	${}^{40}\text{Ca}$
E_0	-11.8	-60.2	-347.5
ΔE_{CCSD}	-17.1	-82.6	-143.7
$\Delta E_{\text{CCSD(T)}}$	-0.3	-5.4	-11.7
$E_{\text{CCSD(T)}}$	-29.2	-148.2	-502.9
exact (FY)	-29.19(5)		

< 2 MeV triples corrections
in ${}^{40}\text{Ca}$ extrapolated

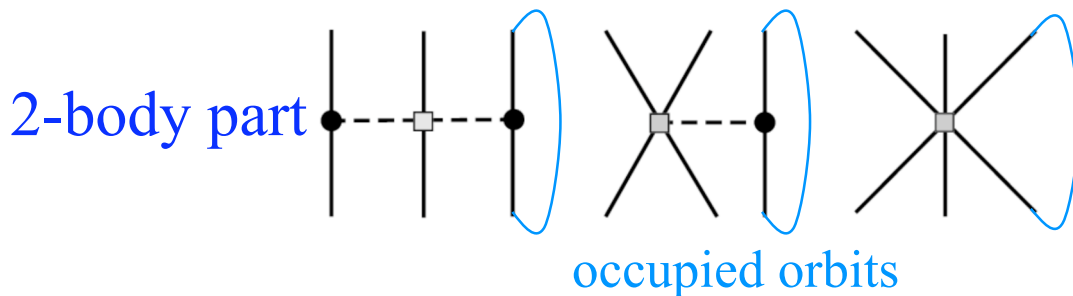
for CC developments, see talk by Gaute Hagen

Towards 3N interactions in medium-mass nuclei

developed coupled-cluster theory with 3N interactions [Hagen et al. \(2007\)](#)
 first benchmark for ^4He based on low-momentum interactions



show that 0-, 1- and 2-body parts of 3N interaction dominate



residual 3N interaction can be neglected: very promising

supports that monopole shifts are due to 3N interactions [cf. Zuker \(2003\)](#)

3N contribution to two-body monopole interaction

take into account 2-body part of 3N interactions to two-body monopoles
determines interaction of s with t orbit

$$V_{st}^T = \frac{\sum_J V_{stst}^{JT} (2J+1) [1 - (-)^{J+T} \delta_{st}]}{\sum_J (2J+1) [1 - (-)^{J+T} \delta_{st}]}$$

small changes in monopoles enhanced by number operators

monopoles from $V_{\text{low } k}(\Lambda) + 2\text{nd order (6hw)}$, cutoff independent in $T=1$

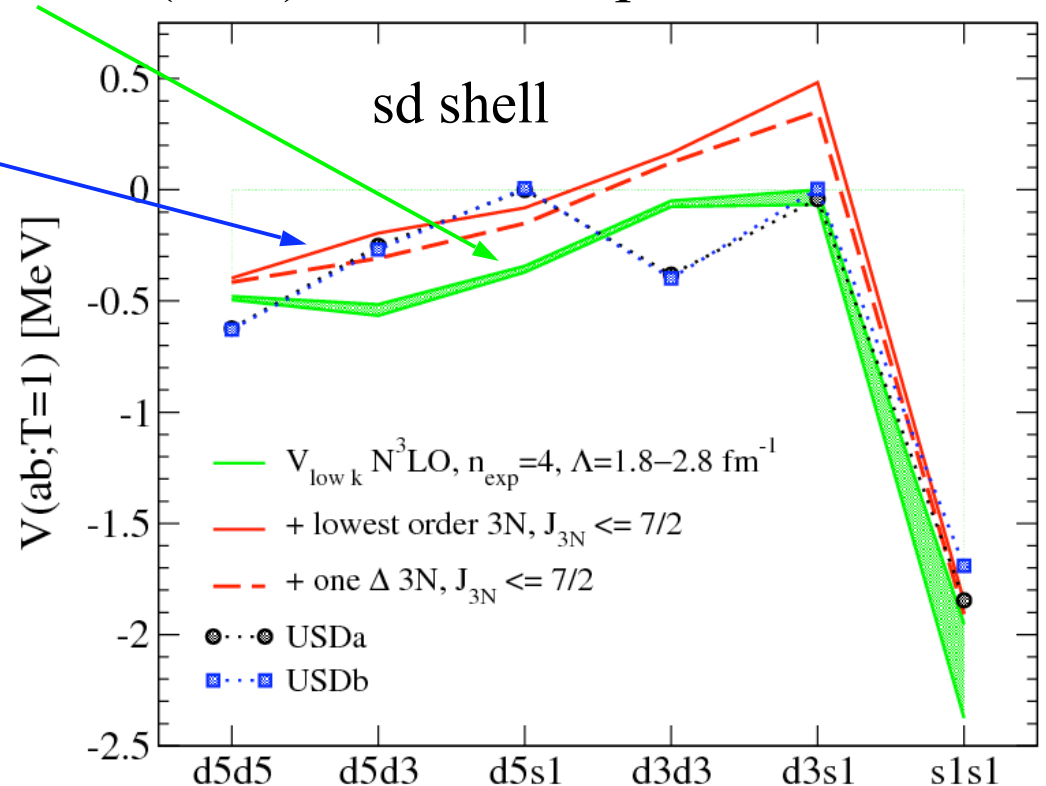
repulsive contributions from
3N interactions fit to $A=3,4$

reproduces hierarchy for
low orbits

$(d_{3/2})^2$ with ^{40}Ca core?

dominated by c_i -terms
→ Δ -hole contributions

shell model matrix elements
for different cores probe 3N interactions

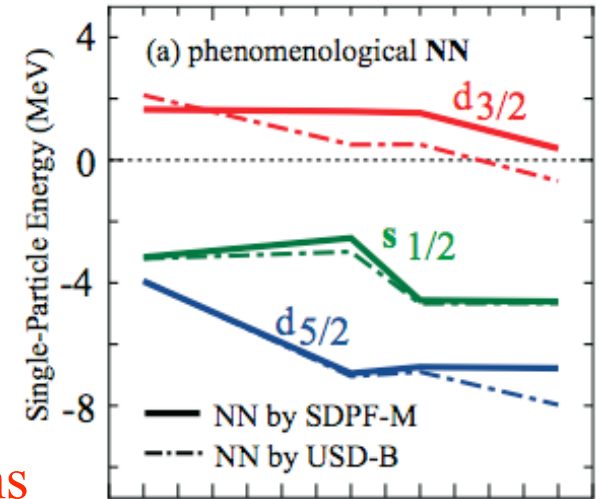
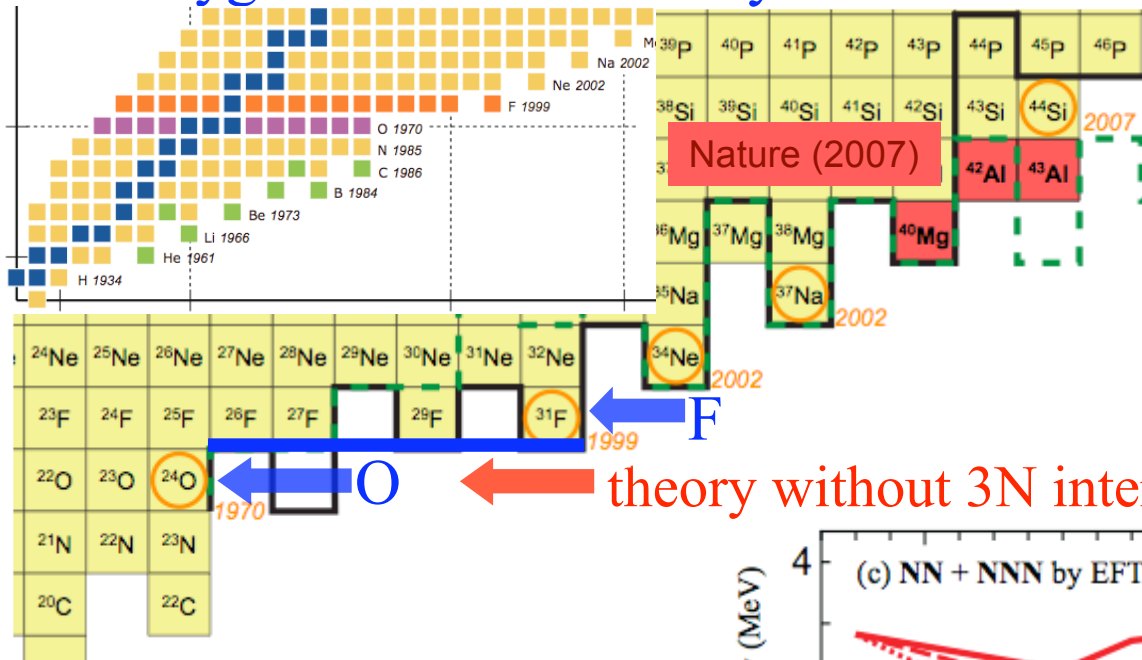


Otsuka, Suzuki, Holt, AS, Akaishi, in prep.

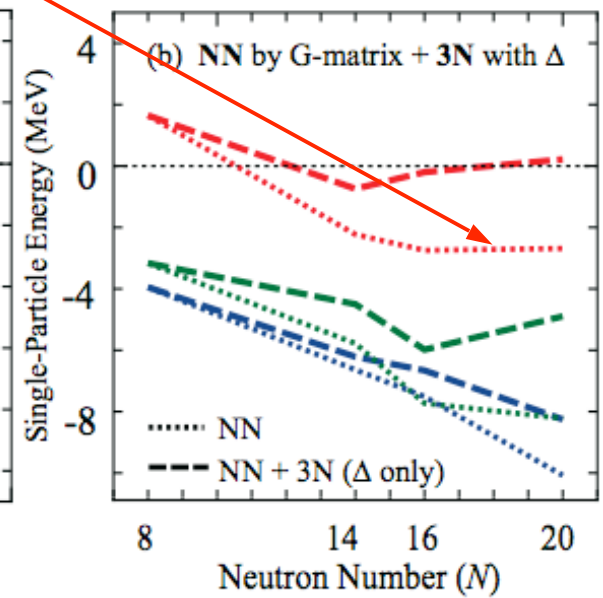
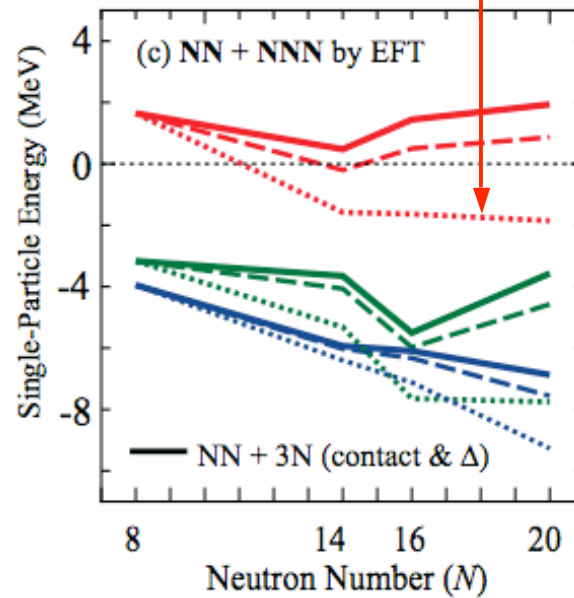
3N forces and neutron-rich nuclei

evolution to the neutron drip-line: limits of existence and shell structure

The oxygen-fluorine anomaly



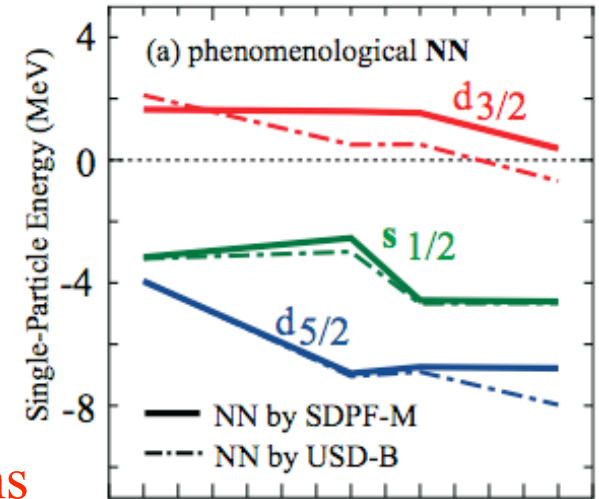
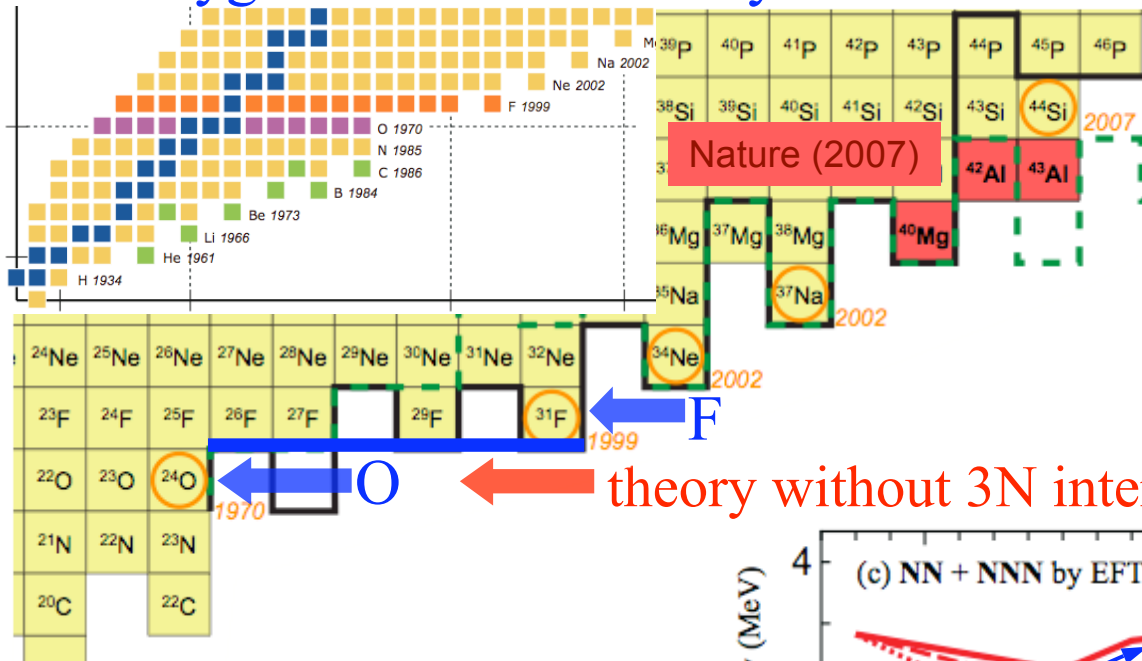
theory without 3N interactions



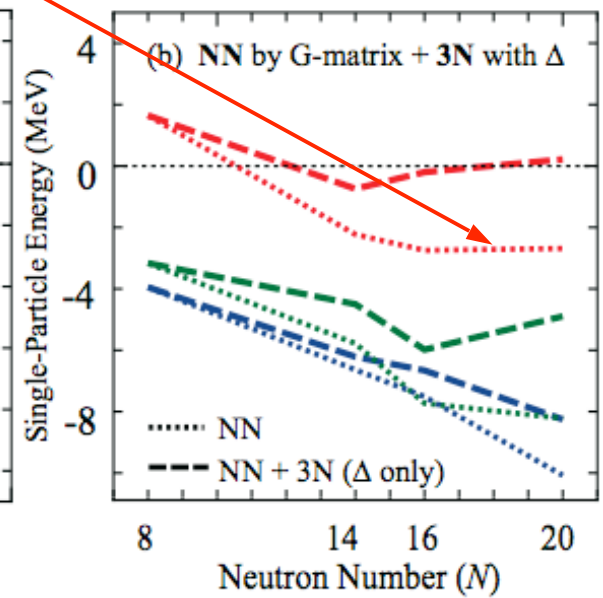
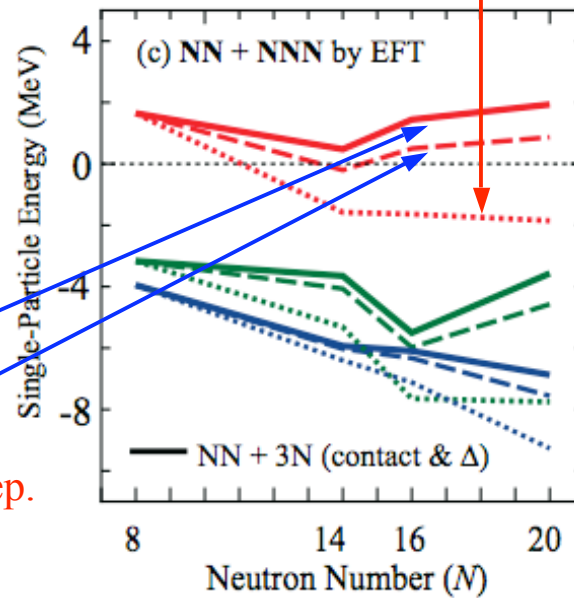
3N forces and neutron-rich nuclei

evolution to the neutron drip-line: limits of existence and shell structure

The oxygen-fluorine anomaly



theory without 3N interactions



first results with
3N interactions:

3N fit to $E(^3\text{H})$, ^4He radius,
3N from single- Δ excitation

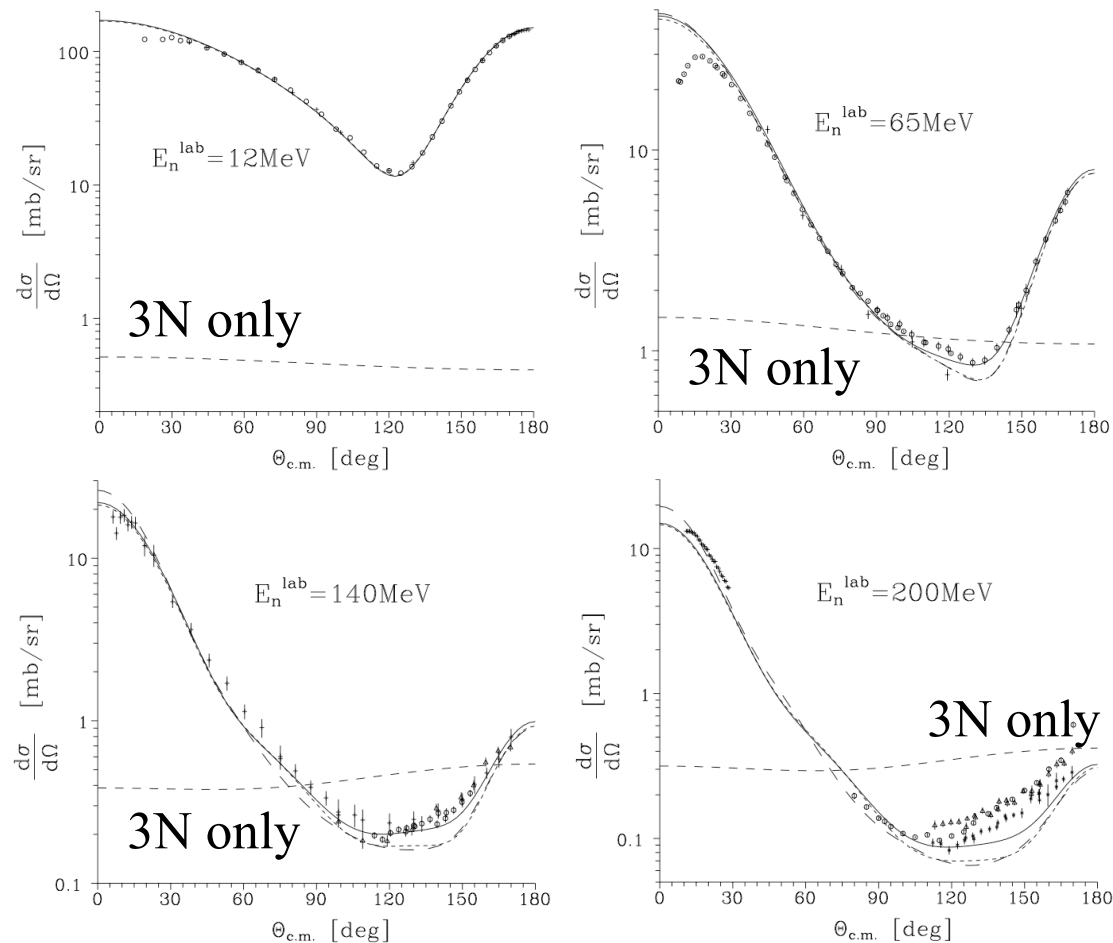
Otsuka, Suzuki, Holt, AS, Akaishi, in prep.

reproduce $N=16$ dripline

neutron-proton interaction stronger, binds $d_{3/2}$ neutrons in fluorine

Is there an EFT argument/counting for the hierarchy of n-body parts of many-body forces? for the SM multipole expansion?

maybe interesting to understand connection to
weak dependence on angle in N-d scattering



Witala et al. (1998)

A scenic winter landscape featuring a dense forest of snow-laden evergreen trees in the foreground. In the background, a range of rugged mountains is visible under a cloudy sky, with a body of water nestled between the peaks. The overall scene is serene and cold.

Outline

4. Three-nucleon forces in nuclear and neutron matter

Weinberg eigenvalue diagnostic

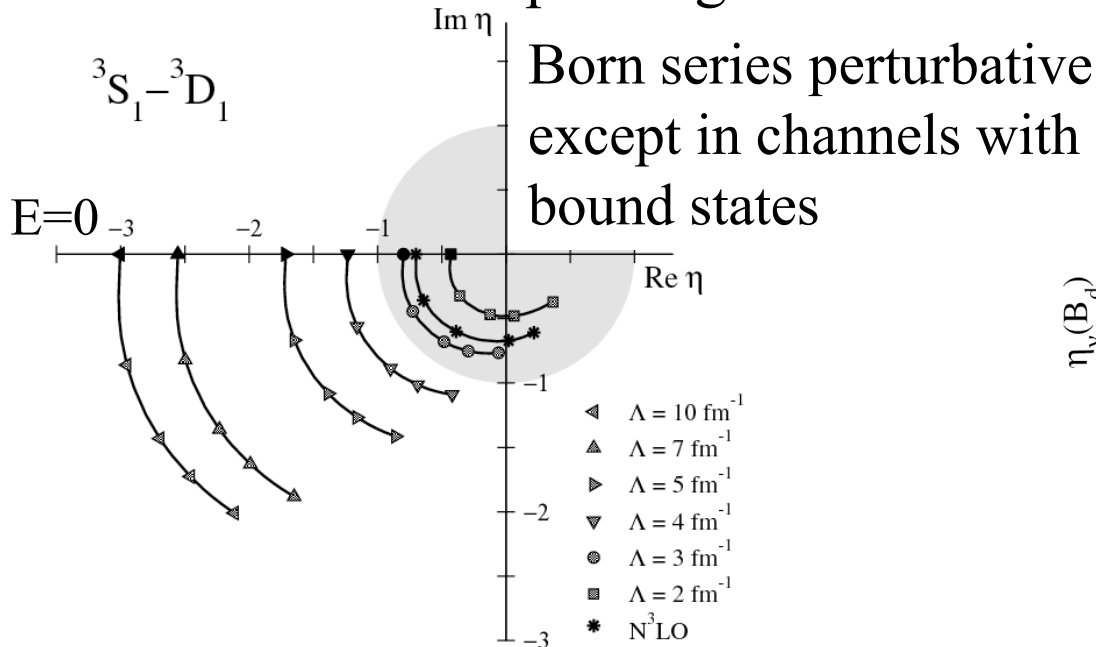
study spectrum of $G_0(z)V |\Psi_\nu(z)\rangle = \eta_\nu(z) |\Psi_\nu(z)\rangle$ at fixed energy z

governs convergence $T(z) |\Psi_\nu(z)\rangle = (1 + \eta_\nu(z) + \eta_\nu(z)^2 + \dots) V |\Psi_\nu(z)\rangle$

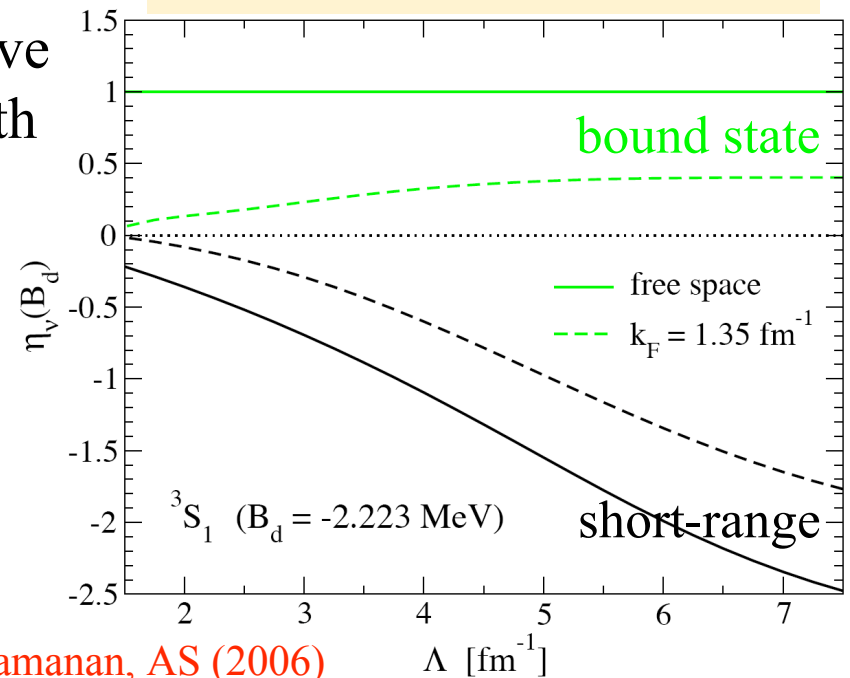
can write as Schrödinger equation $(H_0 + \frac{1}{\eta_\nu(z)} V) |\Psi_\nu(z)\rangle = z |\Psi_\nu(z)\rangle$

high momenta/large cutoffs lead to flipped-potential bound states of $-\lambda V$ for small λ /large $\eta \rightarrow$ strong coupling to high momenta/short range and Born series always nonperturbative

RG evolution decouples high momenta



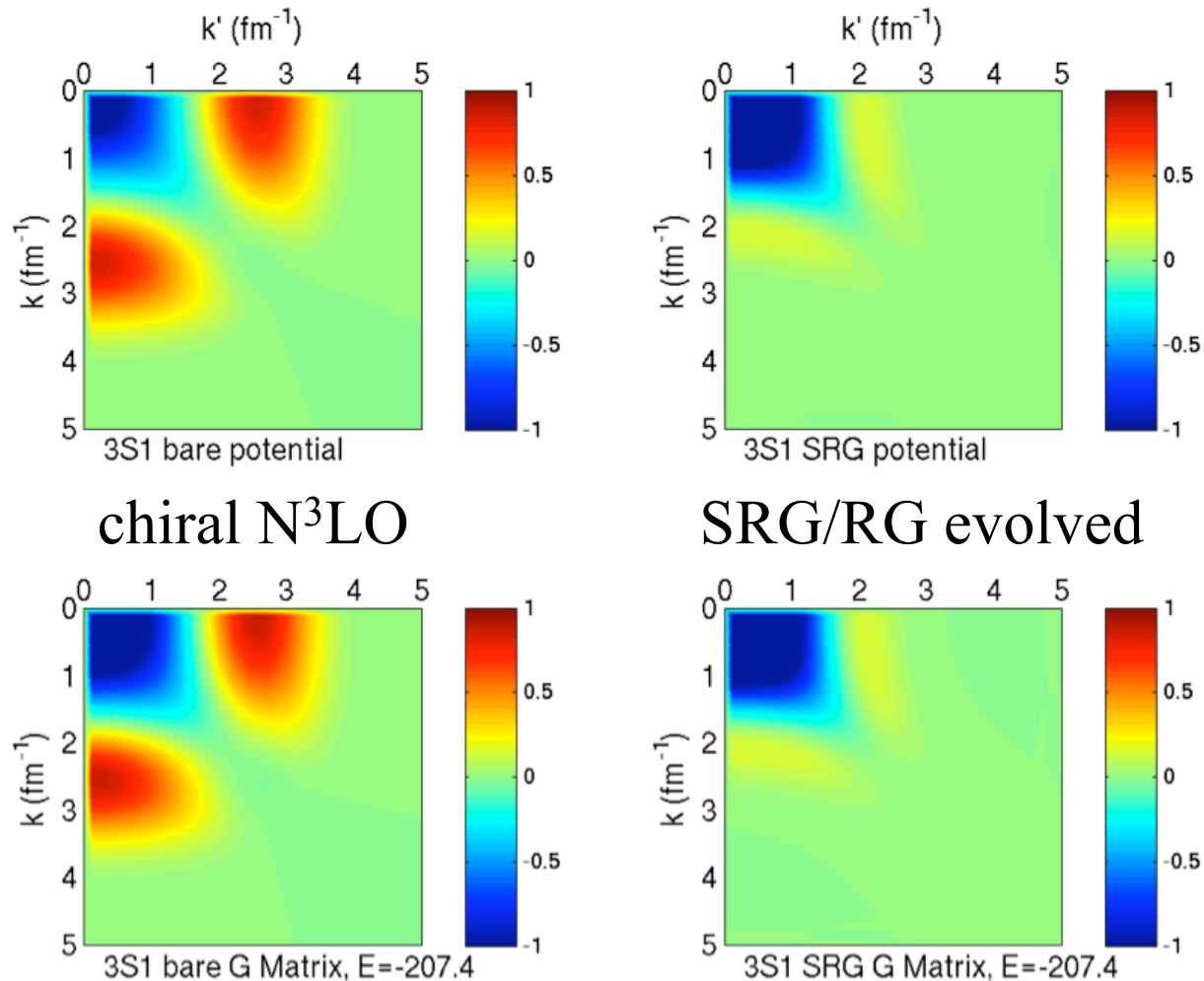
loosely-bound state is dissolved in nuclear matter



Is nuclear matter perturbative with chiral EFT and RG?

conventional Bethe-Brueckner-Goldstone expansion:

no, due to nonpert. cores (flipped-V bound states) and off-diag coupling



G matrix

conventional G-matrix approach does not solve off-diagonal coupling

Is nuclear matter perturbative with chiral EFT and RG?

conventional Bethe-Brueckner-Goldstone expansion:

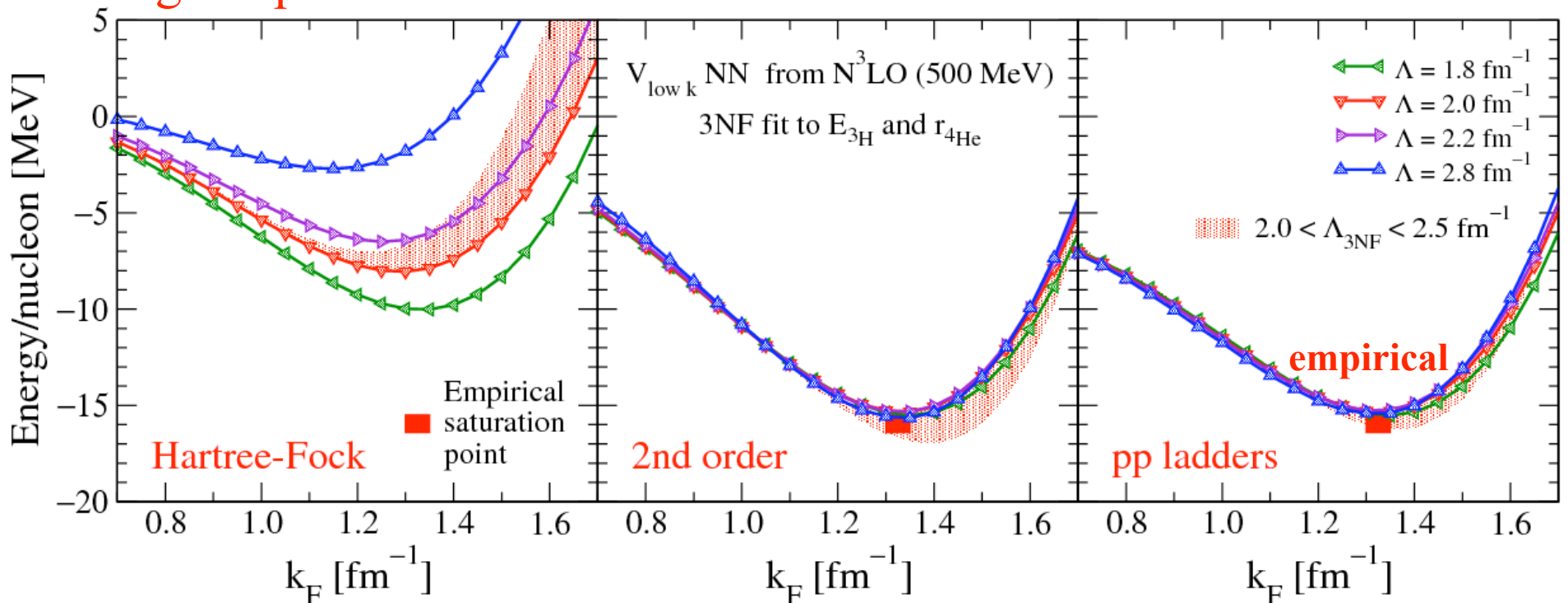
no, due to nonpert. cores (flipped-V bound states) and off-diag coupling

start from chiral EFT and RG evolution:

nuclear matter converged at \approx 2nd order, 3N drives saturation

weak cutoff dependence, improved by 3N fits to ^4He radius

exciting: empirical saturation within theoretical uncertainties

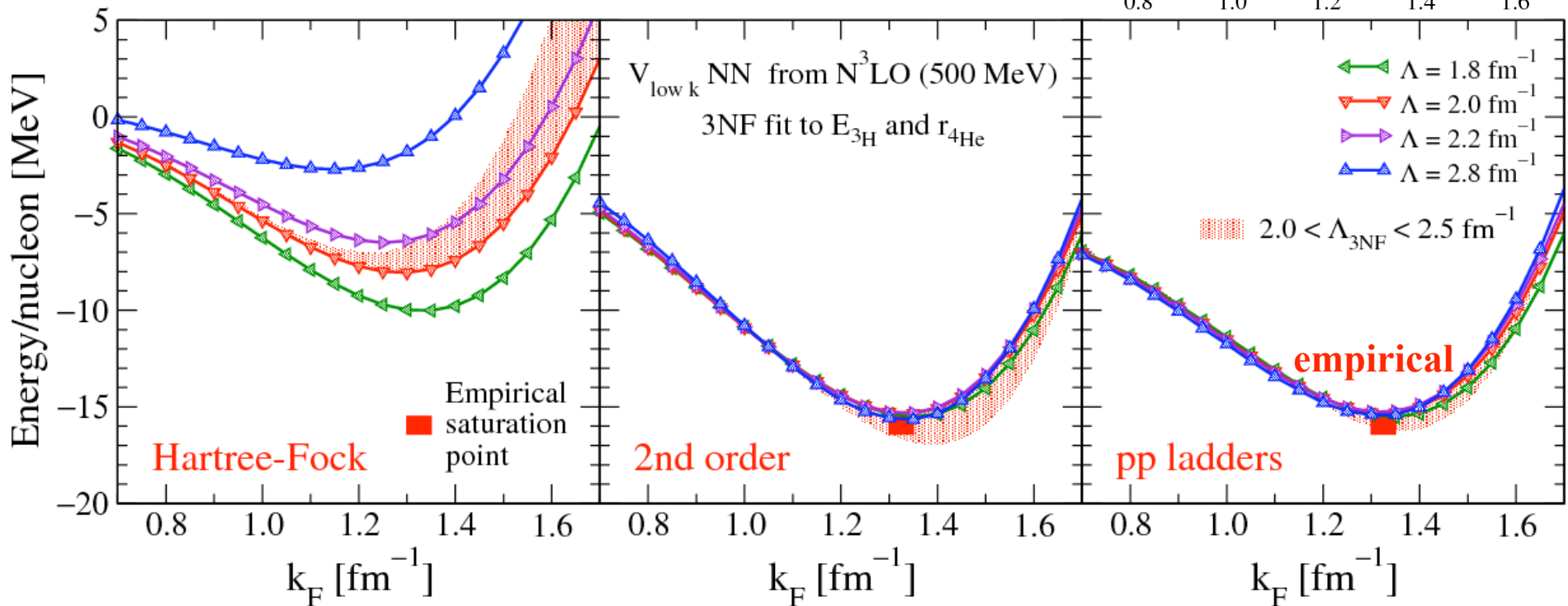
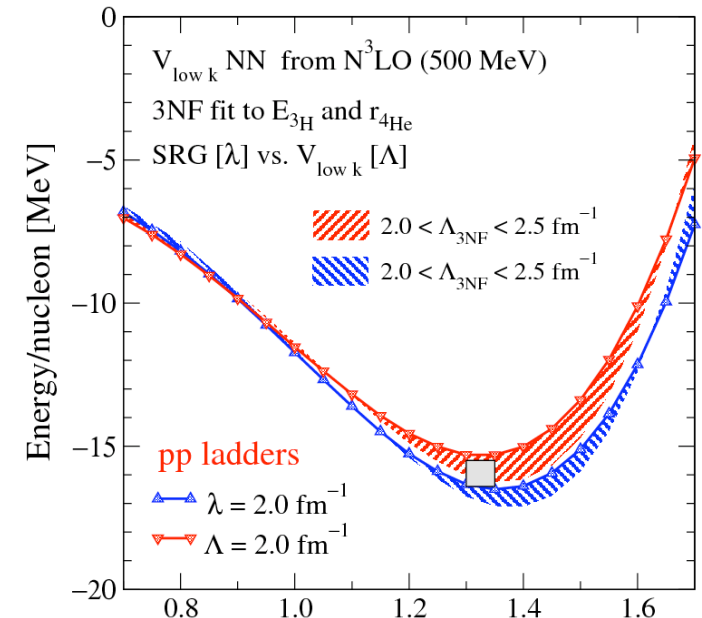


Bogner, AS, Furnstahl, Nogga (2005) and 0903.3366.

Is nuclear matter perturbative with chiral EFT and RG?

SRG nuclear matter results similar,
 supports general nature of 3N fit and results
 theoretical uncertainty $\sim 1\text{-}2$ MeV per particle

compressibility $K = 190 - 240$ MeV
 mainly from $\Lambda_{3\text{NF}} = 2.0 - 2.5 \text{ fm}^{-1}$

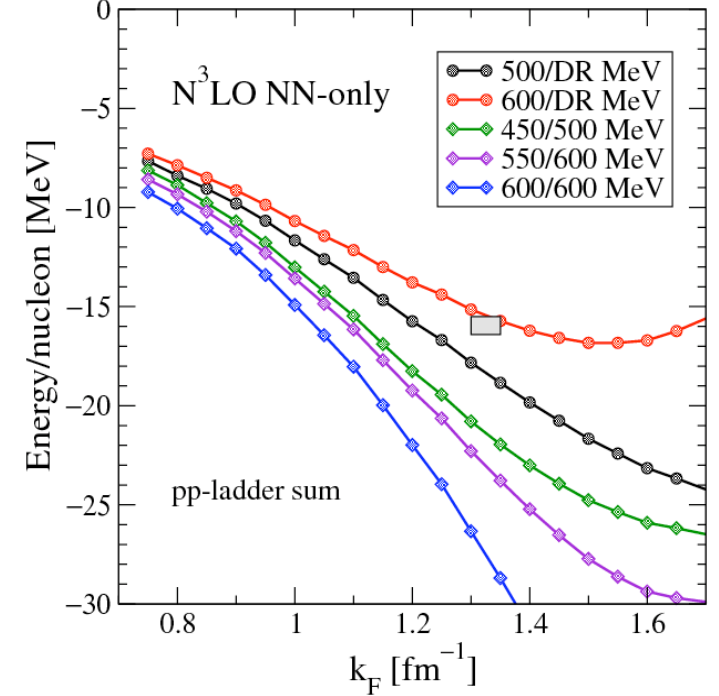
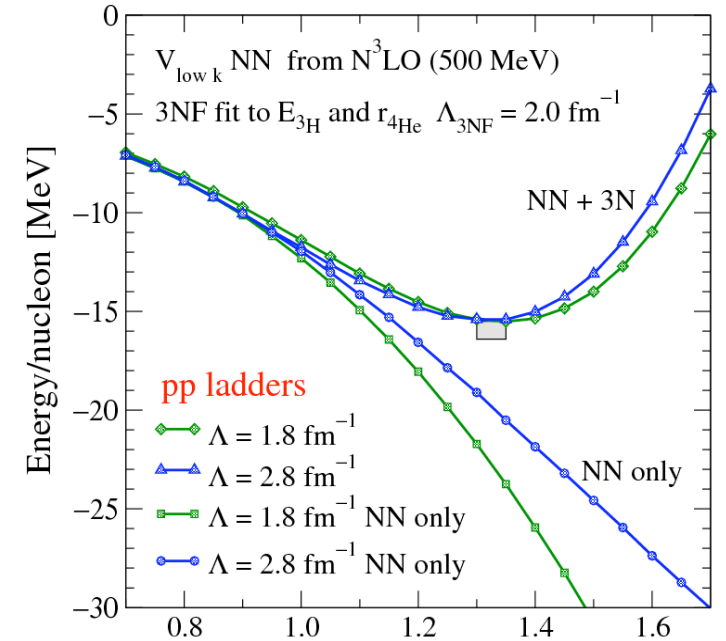


Impact of 3N interactions on nuclear matter

3N contributions natural but not small for RG-evolved and “bare” $N^3\text{LO}$ interactions

long-range c_i -terms repulsive in nuclear matter and drive saturation

k_F	Λ	Hartree-Fock					Hartree-Fock + dominant second order				
		T	$V_{\text{low } k}$	V_c	V_D	V_E	T	$V_{\text{low } k}$	V_c	V_D	V_E
1.0	1.6	12.44	-19.62	1.65	0.42	-0.22	15.50	-26.58	1.49	0.34	-0.29
	1.9	12.44	-18.18	1.67	-0.25	0.40	16.29	-26.81	0.85	-0.09	0.55
	2.1	12.44	-17.35	1.67	-0.42	0.62	16.92	-27.04	0.11	0.05	0.79
	2.3	12.44	-16.56	1.67	-0.56	0.81	17.60	-27.27	-0.89	0.43	0.85
1.2	1.6	17.92	-31.47	5.37	1.31	-0.64	20.86	-37.66	4.59	1.03	-0.65
	1.9	17.92	-28.95	5.61	-0.81	1.18	21.80	-37.38	3.99	-0.50	1.28
	2.1	17.92	-27.51	5.67	-1.37	1.84	22.87	-37.53	2.27	-0.37	1.82
	2.3	17.92	-26.13	5.70	-1.86	2.42	24.32	-37.95	-0.38	0.51	1.78
1.35	1.6	22.67	-42.47	10.75	2.59	-1.21	26.09	-47.85	8.73	1.96	-1.12
	1.9	22.67	-38.82	11.95	-1.69	2.34	26.75	-46.72	9.14	-1.16	2.24
	2.1	22.67	-36.74	12.19	-2.91	3.68	28.05	-46.47	6.99	-1.33	3.22
	2.3	22.67	-34.77	12.30	-3.97	4.89	30.06	-46.45	3.10	-0.35	3.26



Impact of 3N interactions on neutron matter

Tolos, Friman, AS (2007); Hebeler, AS, in prep., see talk by Kai Hebeler

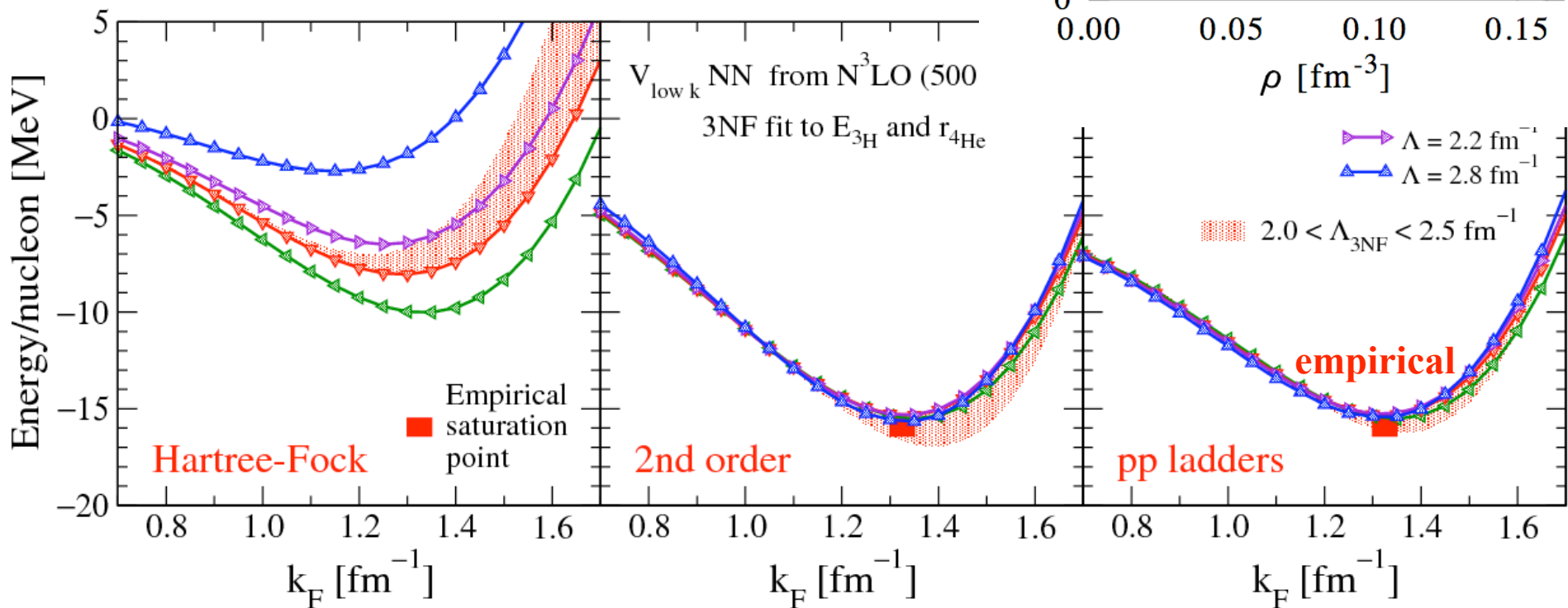
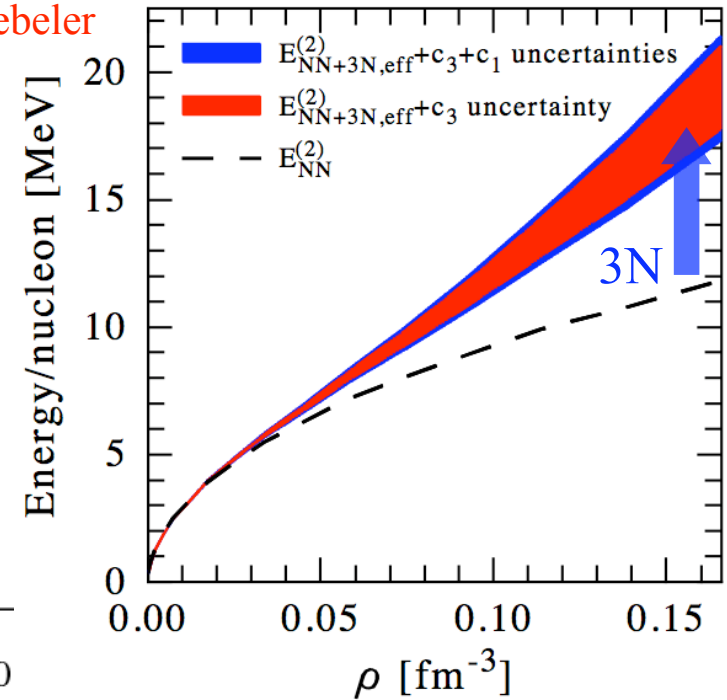
symmetry energy and density dependences

c_i uncertainties overwhelm cutoff variation

D, E terms vanish in spin-sat. neutron matter

developing density-dep. NN interactions

to include 3N effects



Towards ab-initio density functional theory

low-momentum interactions key to develop universal density functional

first DFT results with microscopic pairing functional

Lesinski, Duguet, Bennaceur, Meyer (2008), see talk by Thomas Lesinski

density matrix expansion

Bogner, Furnstahl, Platter (2008), see also Finelli, Kaiser, Weise (2003)

$$\mathcal{E} = \frac{\tau}{2M} + A[\rho] + B[\rho]\tau + C[\rho]|\nabla\rho|^2 + \dots$$

improved DME: use phase space averaging for nuclei

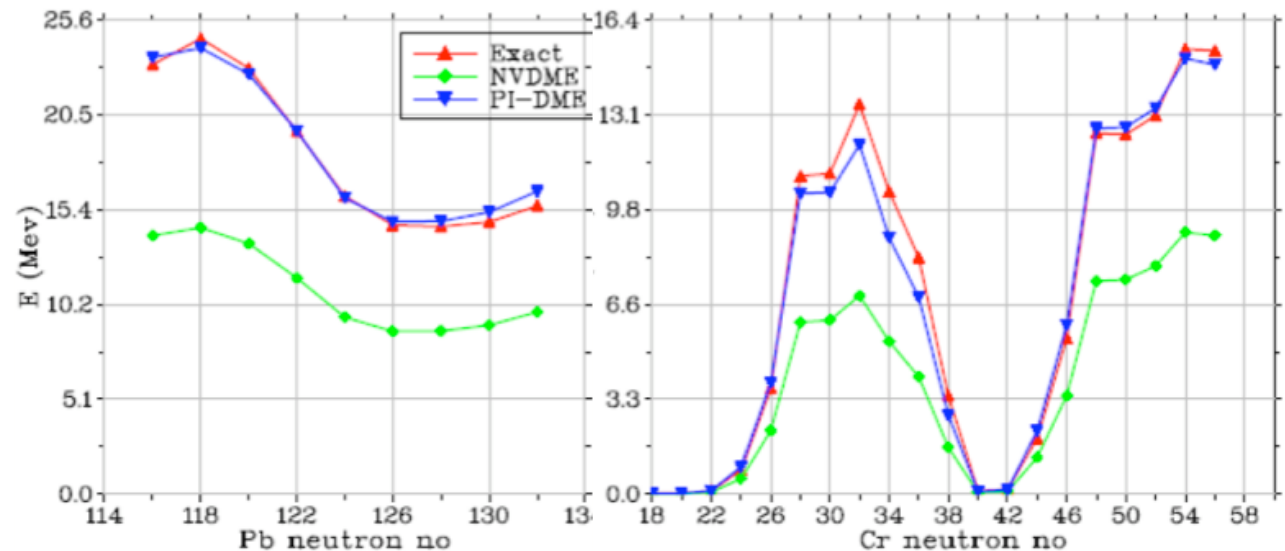
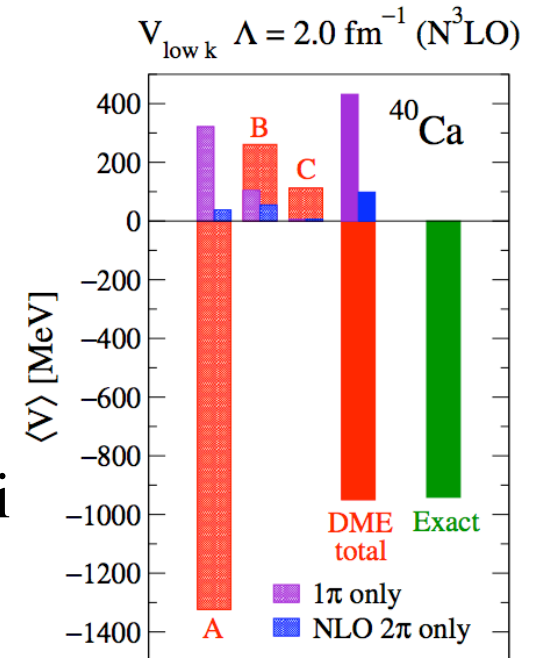
Gebremarian, Bogner, Duguet, in prep.

use EFT/RG interactions:

to identify new terms

to quantify errors

to benchmark with
ab-initio methods





Outline

5. Summary: impact of three-nucleon forces



Thanks to collaborators

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東京大学
THE UNIVERSITY OF TOKYO

T. Otsuka



T. Suzuki

Summary

development of EFT and RG for nuclear forces

advances in ab-initio methods and in nuclear matter theory

EFT and RG interactions enable a unified description
from light to heavy nuclei and matter in astrophysics

three-nucleon forces play a central role

spin-orbit splittings and spin-orbit shell closures

helium halo nuclei, location of the neutron drip line

saturation, symmetry energy, density dependences

individual 3N parts can have interesting dependences
 c_i -terms attractive in light nuclei, repulsive in bulk matter
(scheme dependent)