

# Non-Empirical Pairing Functional

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# Outline

- 1 Introduction
  - EDF with pairing
- 2 Non-empirical pairing functional: Formalism
  - Separable operator representation of  $V_{\text{low } k}$
  - EDF calculations in spherical nuclei
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  - Results: gaps, Coulomb
- 4 Effective-mass approximation
  - $\Lambda$ -dependence
  - Choice of effective mass
- 5 Summary

# Introduction : Non-empirical Energy Density Functional

- EDF: only general method applicable to the heavier nuclei
- ✗ Empirical in nature: limited to observables used for fitting
- ➔ Long-term goal: build a nuclear EDF (partly) from first principles
  
- Pairing part
  - ✗ Parameters hard to constrain (esp. isovector dependence)
  - ✓ Dominated by  $^1S_0$  channel at low density
  - ➔ **Starting point:** use perturbative  $V_{\text{low } k}$  at first order
  
- Origin of superfluidity
  - Relative importance of bare  $V_{\text{low } k, \text{ NN}}$  ?
  
- Assess first-order  $V_{\text{low } k, \text{ NN}}$  contribution to pairing gaps by using
  - $V_{\text{low } k, \text{ NN}}$  (pp channel)
  - A Skyrme EDF with consistent  $m^* = m_{\text{HF, INM}}^*$  (ph channel)

## Introduction : EDF with pairing

- Starting from effective zero-range, momentum-dependent vertices :

$$\mathcal{E}[\rho, \kappa^*, \kappa] \equiv \sum_{ij} t_{ij} \rho_{ij} + \frac{1}{2} \sum_{ijkl} \bar{V}_{ijkl}^{\rho\rho} \rho_{ik} \rho_{jl} + \frac{1}{4} \sum_{ijkl} \bar{V}_{ijkl}^{\kappa\kappa} \kappa_{ij}^* \kappa_{kl}$$

$$\rho_{ij} = \langle \Phi | c_j^\dagger c_i | \Phi \rangle, \quad \kappa_{ij} = \langle \Phi | c_j c_i | \Phi \rangle$$

- Bogolyubov-de Gennes equations

$$\begin{pmatrix} h - \lambda & \Delta \\ -\Delta^* & -h^* + \lambda \end{pmatrix} \begin{pmatrix} U_i \\ V_i \end{pmatrix} = E_i \begin{pmatrix} U_i \\ V_i \end{pmatrix},$$

$$h_{ij} = \frac{\delta \mathcal{E}}{\delta \rho_{ji}}, \quad \Delta_{ij} = \frac{\delta \mathcal{E}}{\delta \kappa_{ij}^*}$$

- Quasiparticle energy  $E_i = \sqrt{(h_{ii} - \lambda)^2 + \Delta_{ii}^2}$  ( $i$ : q.p. or canonical state)
- Odd-mass nucleus  $\sim$  1-q.p. state: linked to odd-even mass difference

$$\Delta_n^{(3)} = \frac{(-1)^N}{2} [\mathcal{E}(N-1) - 2\mathcal{E}(N) + \mathcal{E}(N+1)] \simeq \Delta_{\text{LCS}}^n(N) \quad (N \text{ odd})$$

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# Introduction : “Skyrme” EDF

- Kohn-Sham-like approach : Minimize  $\mathcal{E}[\rho, \kappa]$ , with  $\rho$  and  $\kappa$  characterizing a Bogolyubov quasiparticle vacuum (BCS state)  $|\Phi\rangle$ .
- Form of the functional (particle-hole)
  - Integral of an energy density bilinear in  $\rho$  + “density-dependent” terms
  - Local + up to 2<sup>nd</sup>-order derivatives

$$\mathcal{H}_{\text{Skyrme}} = \sum_{t=0,1} \left\{ C_t^\rho [\rho_0] \rho_t^2 + C_t^{\Delta\rho} \rho_t \Delta \rho_t + C_t^\tau \rho_t \tau_t + \frac{1}{2} C_t^J \vec{j}_t^2 + C_t^{\nabla \cdot J} \rho_t \nabla \cdot \vec{j}_t \right\}$$

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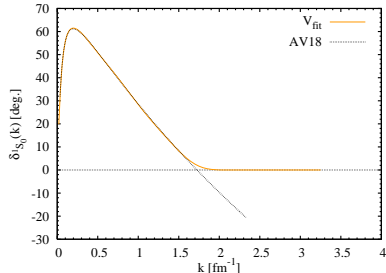
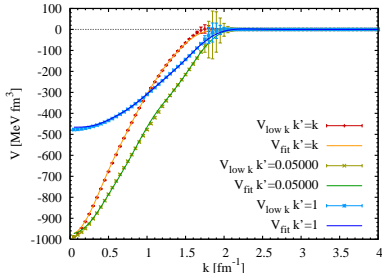


# Separable representation of $V_{\text{low } k}(\Lambda) + V_{\text{Coulomb}}$

- High precision separable representation of rank  $n$

$$V_n^1 S_0(k, k'; \Lambda) = \sum_{\alpha=1}^n g_{\alpha}(k) \lambda_{\alpha} g_{\alpha}(k')$$

- Fit of  $g_{\alpha}(k)$  and  $\lambda_{\alpha}$  to  $V_{\text{low } k}^1 S_0(k, k'; \Lambda)$  and  $T_{\text{low } k}^1 S_0(k, k'; k^2; \Lambda)$
- For  $\Lambda = 1.8/4.0/\infty \text{ fm}^{-1}$  (rank 2/4/9) and smooth cutoff



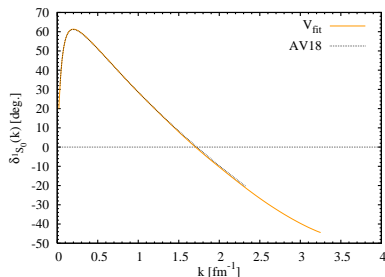
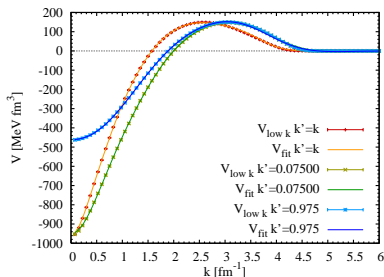
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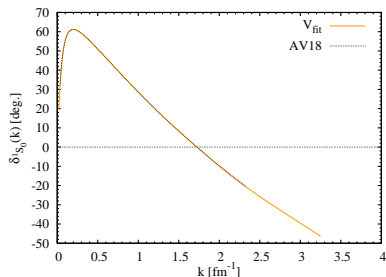
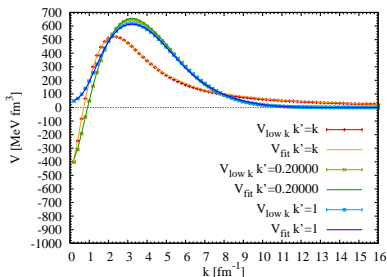
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## EDF calculations in spherical nuclei

- Separable force in coordinate-space

$$\langle \mathbf{r}'_1 \mathbf{r}'_2 | V^{1S_0} | \mathbf{r}_1 \mathbf{r}_2 \rangle = \sum_{\alpha} G_{\alpha}(s') \lambda_{\alpha} G_{\alpha}(s) \delta(\mathbf{R}' - \mathbf{R}),$$

$$\mathbf{s} = \mathbf{r}_1 - \mathbf{r}_2, \quad \mathbf{R} = (\mathbf{r}_1 + \mathbf{r}_2)/2$$

- Coordinate-space form factor  $G_{\alpha}(s)$ : Fourier transform  $g_{\alpha}(k)$
- “Pseudo-local” pairing functional

$$E_{\text{pair}}[\kappa^*, \kappa] = \frac{1}{2} \int d^3 \mathbf{R} \sum_{\alpha=1}^n \check{\chi}_{\alpha}^*(\mathbf{R}) \lambda_{\alpha} \check{\chi}_{\alpha}(\mathbf{R})$$

- The effective pair densities  $\check{\chi}_{\alpha}(\mathbf{R})$  contain the finite range/non-locality of the pairing vertex

$$\check{\chi}_{\alpha}(\mathbf{R}) = \int d^3 \mathbf{s} \sum_{\sigma} (-)^{\frac{1}{2} - \sigma} G_{\alpha}(s) \kappa \left( \mathbf{R} + \frac{\mathbf{s}}{2}, \sigma; \mathbf{R} - \frac{\mathbf{s}}{2}, -\sigma \right)$$

- ➔ Non-local potential and pairing field

## EDF calculations in spherical nuclei (2)

- Define a reduced two-body wave function

$$\check{\Psi}_{ij}^{\alpha}(\mathbf{R}) = \int d^3\mathbf{s} G_{\alpha}(s) \Psi_{ij}(\mathbf{R} + \mathbf{s}/2, \mathbf{R} - \mathbf{s}/2)$$

$$\Psi_{ij}(\mathbf{r}, \mathbf{r}') = \sum_{\sigma} (-)^{s-\sigma} \phi_i(\mathbf{r}, \sigma) \phi_j(\mathbf{r}', -\sigma).$$

- The  $\phi_i$  are basis functions : the  $\check{\Psi}_{ij}^{\alpha}(\mathbf{R})$  are computed **once**
- Build densities and pairing field matrix elements

$$\check{\chi}_{\alpha}(\mathbf{R}) = \sum_{ij; i>0, j<0} \check{\Psi}_{ij}^{\alpha}(\mathbf{R}) \kappa_{ij}$$

$$\Delta_{\alpha}(\mathbf{R}) = \frac{1}{2} \sum_{\alpha} \lambda_{\alpha} \check{\chi}_{\alpha}(\mathbf{R})$$

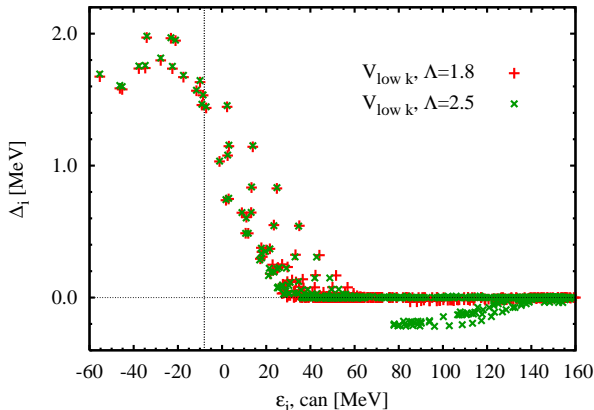
$$\Delta_{ij} = \sum_{\alpha} \int d^3\mathbf{R} \check{\Psi}_{ij}^{\alpha}(\mathbf{R}) \Delta_{\alpha}(\mathbf{R})$$

- Basic tool: CoM/relative coordinate separation in spherical symmetry
  - Basis of spherical Bessel functions  $j_{\ell}(kr)$ ,  $k < k_{\max} = \Lambda + 1.0 \text{ fm}^{-1}$
  - ✓ Natural representation for continuum states ( $\rightarrow$  drip line...)

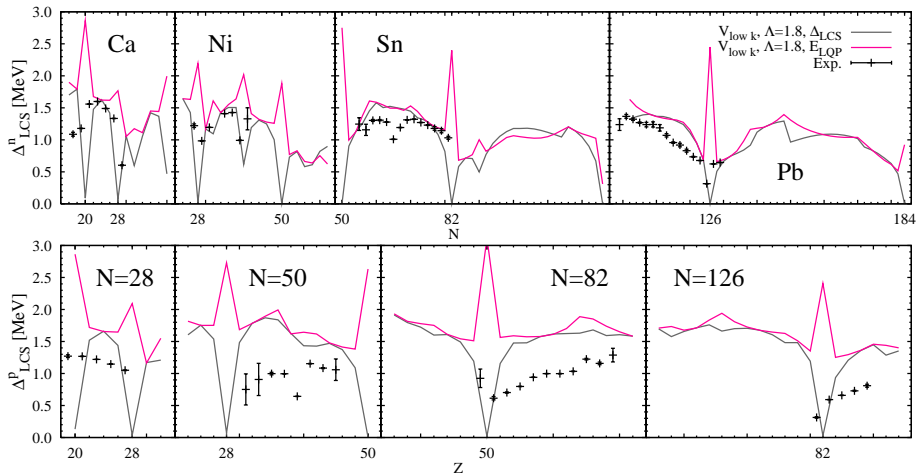
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# Results: State-dependent gaps ( $^{120}\text{Sn}$ )



- SLy4,  $m^*/m = 0.7$
- Gaps almost identical below  $\varepsilon \simeq 60$  MeV
- $\Lambda = 2.5 \text{ fm}^{-1}$ : hint of a repulsive “core” at high energy

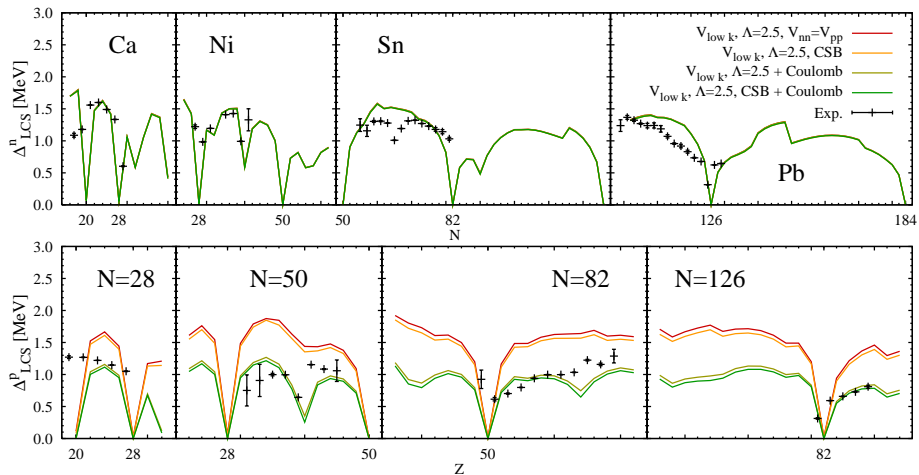
Results: LQP energies and LCS Gaps, no Coulomb,  $V_{nn} = V_{pp}$ 

- $E_{LQP}$ : lowest quasiparticle energy,  
 $\Delta_{LCS}$ : corresponding gap

- Theoretical  $\Delta_{LCS}^n$  close (n) or above (p) exp.  $\Delta_{odd}^{(3)}$  values

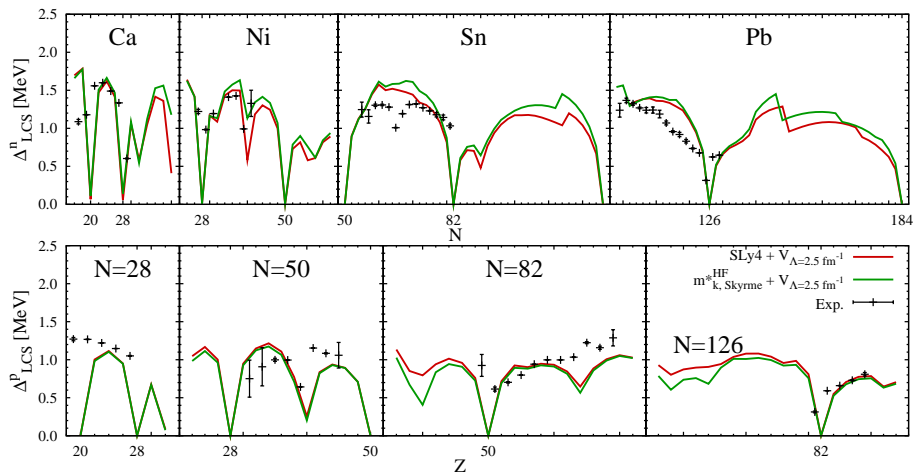


## Results: LCS Proton gaps, CSB, Coulomb



- Coulomb decreases gaps by  $\sim 30\%$
- $V_{NN}$  (incl. Coulomb) at first order matches the magnitude of the gaps
- Coulomb effect consistent with Madrid group (Gogny)

## HF effective mass

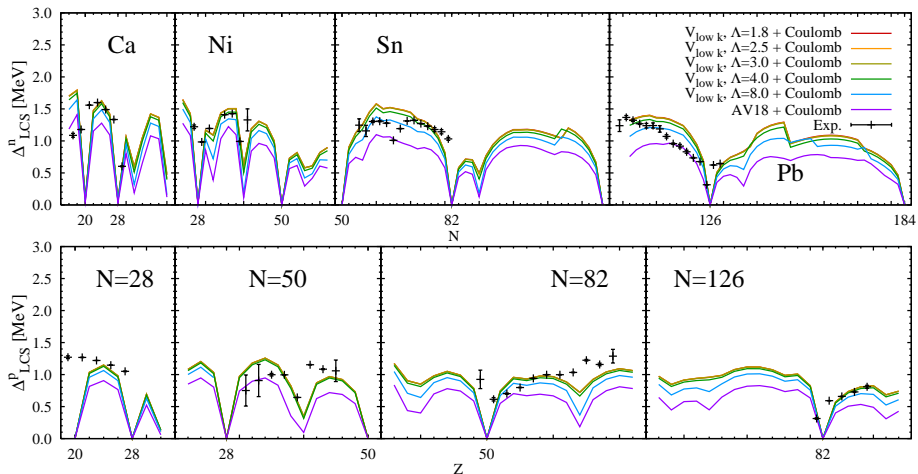


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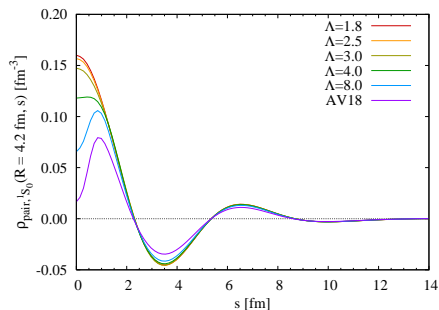
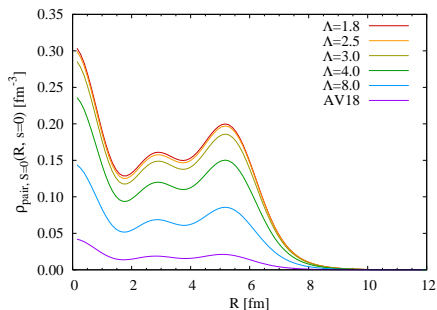
# $\Lambda$ -dependence: SLy4

- Milan group: SLy4 + Argonne  $v_{14}$  yields ca. 1 MeV gap in  $^{120}\text{Sn}$   
(A. Pastore et al. PRC 78 (2008) 024315)



# $\Lambda$ -dependence: SLy4

- Non-local pair density in  $^{120}\text{Sn}$ 
  - Short-range part  $\Lambda$ -dependent
  - Long-range part  $\Lambda$ -independent up to  $\Lambda \simeq 4.0 \text{ fm}^{-1}$



# Some comments

- Missing in this calculation:
  - ▼ Higher partial waves
  - ▼ Three-body force
  - ▼ Two-body CoM correction (H. Hergert, R. Roth)
  - ▲ Phonon coupling
  - ▲ Coulomb screening ?
  - ▼ Blocking
- These effects seem to cancel out “near” the valley of stability
- ✗ Still valid in highly asymmetric nuclei ?
  - ➔ High isospin asymmetry  $\Rightarrow$  effect of NNN on proton gaps ?
  - ➔ Continuum effects beyond first order ?
- $\Lambda$ -dependence: soft  $\neq$  hard interactions
- Genuine disagreement likely
- Issue: effective-mass approximation (cf. Kai Hebeler yesterday)
  - Skyrme-EDF has  $m^*(\rho)$ , not  $m^*(\rho, k)$
  - Reliable approximation with low-momentum  $V_{NN}$
  - Uncertainty due to  $m^*$  averaging with hard  $V_{NN}$

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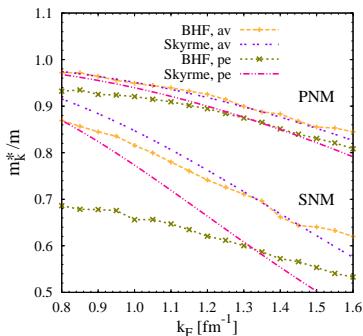
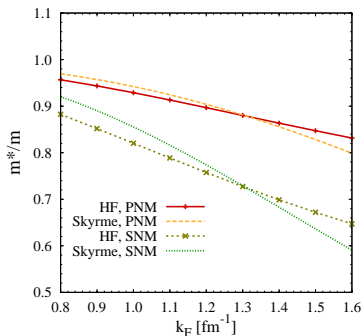
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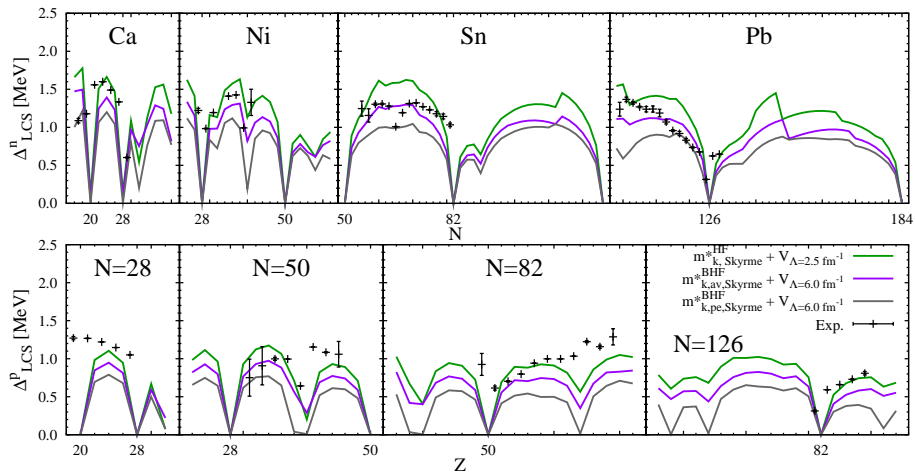


# Effective-mass approximation

- Soft  $V_{NN}$ : HF  $m^*$
- Hard  $V_{NN}$ : BHF  $k$ -mass, averaged over  $k$  or point-evaluated at  $k_F$

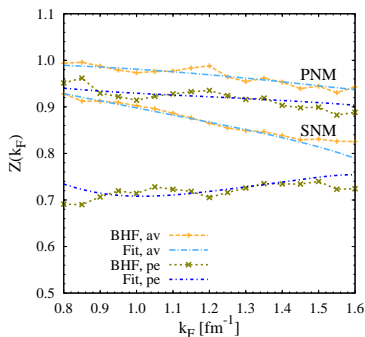
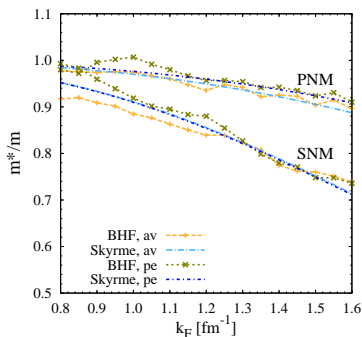


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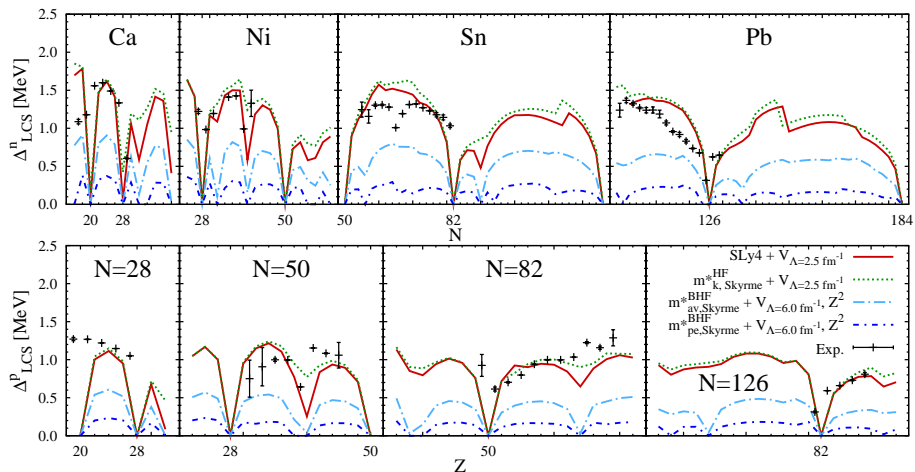
# Effective-mass approximation – full self-energy

- Consistent treatment of self-energy effects:  $\omega$ -dependence of  $\Sigma(k, \omega)$ 
  - Pole approximation:  $Z(k) = [1 + \partial_\omega \Sigma(k, \omega)|_{\omega=\varepsilon_k}]^{-1}$
- Soft interaction: HF  $m^*$
- Hard interaction: BHF total  $m^*$  and  $Z(k)$ , averaged over  $k$  or point-evaluated



$$\hat{\Delta}(k) = \int \frac{p^2 dp}{2\pi^2} \frac{Z(k) V(k, p) Z(p) \hat{\Delta}(p)}{2\sqrt{\varepsilon_p^2 + \hat{\Delta}(p)^2}}$$

## Effective-mass approximation – full self-energy



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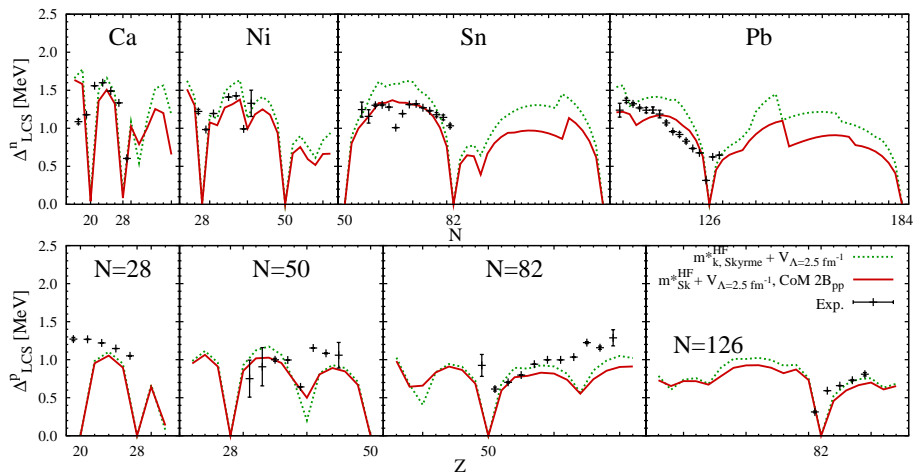
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# Summary and outlook

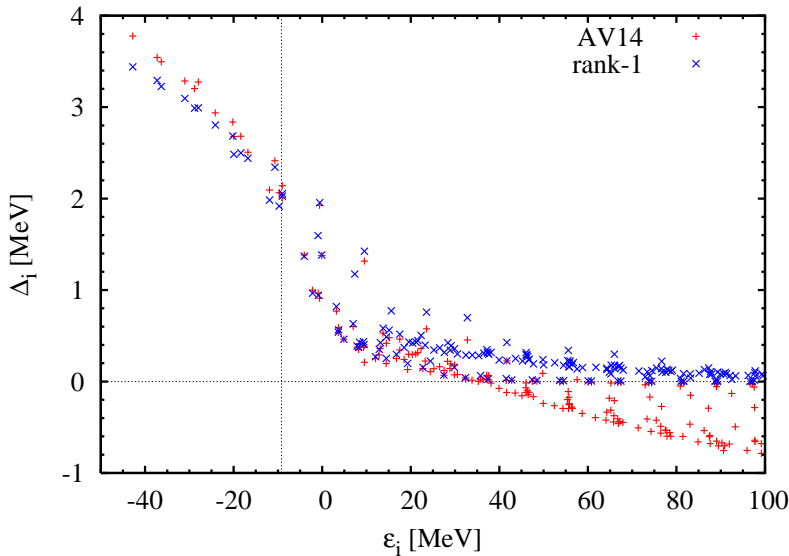
- Systematic calculation of pairing gaps in spherical nuclei
- Low-momentum NN interaction (including Coulomb potential) yields gaps of the order of exp. mass differences
  - Strong reduction of proton gaps due to Coulomb
  - Effects beyond present calculation cancel each other ?
    - Corrections at first order (NNN,CoM...) should leave room for higher orders
- Future work with separable forces
  - 3-body force, through an LDA or DME scheme (TL, K. Hebeler, A. Schwenk, T. Duguet)
  - DME: non-empirical local pairing functional (B. Gebremariam, S. Bogner, T.Duguet)
  - Other channels of the interaction ( $^3P, ^1D$ ) to be added (S. Baroni)

# Thank you !

## Two-body CoM correction





$V_{\text{low } k}$  vs. Argonne

## Coulomb / INM gaps

