

Universal Correlations in Pion-less EFT with the Resonating Group Model

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Ultimate Goals

- ▶ analysis of **universal correlations** between bound & scattering observables in $A > 4$ systems
 - ⇒ objects which **differ on a microscopic scale** exhibit the **same low-energy behavior**

- ▶ How far in A (and density) can one **push the pionless EFT**?
 - ⇒ development of the most **simple** theory, **rooted in QCD**, appropriate for the description of low-energy nuclear properties

The Effective Field Theory “without pions”

see e.g. P.F. Bedaque, U. van Kolck,
Ann. Rev. Nucl. Part. Sci. **52**, 2002

- ▶ **appropriate** theory for nuclear systems with $p_{\text{typ}} \ll m_{\pi}$
- ▶ **systematic** expansion of observables in $\frac{p_{\text{typ}}}{\Lambda_{\not{f}}}$
- ▶ provides theoretical **uncertainties**
- ▶ systematic analysis of **universal** aspects of few-nucleon systems feasible
- ▶ **model independent** statements about the consistency between theory & experiment

(Refined) Resonating Group Model

see e.g. H. M. Hofmann, proceedings of Models and Methods
in Few-Body Physics, 1986.

- ▶ **versatile** method:
 - ▶ suited for calculations of **bound & scattering** observables
 - ▶ applicable to other **degrees of freedom**, e.g., (α, n) & $({}^9\text{Li}, t, n)$
 - ▶ Coulomb interaction included
- ▶ with EFT $_{\not{f}}^{\text{NLO}}$ relatively modest **CPU time** requirements

The RRGM model space

Ritz variation: $\delta \left(\langle \psi | \hat{H} - E | \psi \rangle \right) = 0$

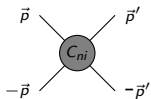
$$\psi_{\text{BS}}^{J\pi}(\vec{\rho}_m, \vec{s}_m) = \mathcal{A} \left\{ \sum_{d,i,j} C_{dij} \left[\left[\prod_{k=1}^{N-1} e^{-\gamma_{dk} \vec{\rho}_k^2} \mathcal{Y}_{l_{ki}}(\vec{\rho}_k) \right]^{L_i} \otimes \Xi_{S_j} \right]^J \cdot \Upsilon \right\}$$

width parameters
antisymmetrizer Jacobi coordinates

$$|3\text{H}\rangle^{J\pi=\frac{1}{2}^+} = \left| \begin{array}{c} S=\frac{1}{2} \\ S_{12}=1, 0 \\ L=0 \end{array} \right\rangle^{J\pi=\frac{1}{2}^+} + \left| \begin{array}{c} S=\frac{3}{2} \\ S_{12}=1 \\ L=2 \end{array} \right\rangle^{J\pi=\frac{1}{2}^+} + \dots$$

NN potential & LEC determination

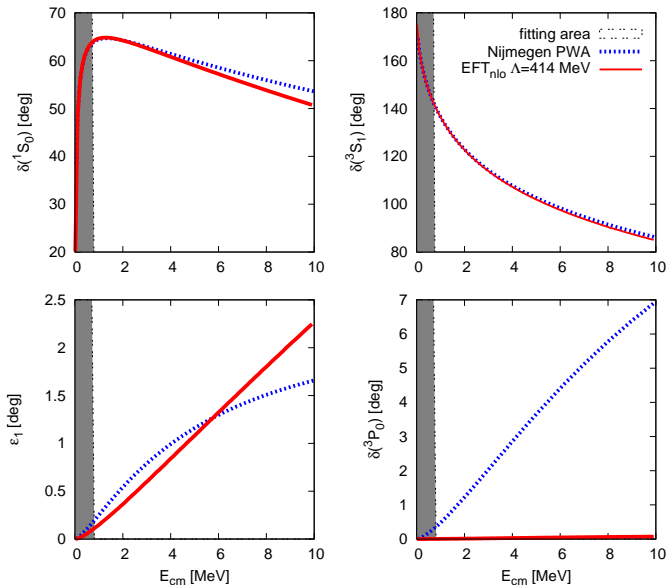
$$\begin{aligned}
 V_{\neq}^{\text{NLO}}(\vec{r}) = & \quad l_0(r) (A_1 + A_2 \vec{\sigma}_1 \cdot \vec{\sigma}_2) + (A_3 + A_4 \vec{\sigma}_1 \cdot \vec{\sigma}_2) \left\{ e^{-\frac{\Lambda^2}{4} r^2}, \nabla^2 \right\} + \\
 & \quad l_0(r) (A_5 + A_6 \vec{\sigma}_1 \cdot \vec{\sigma}_2) r^2 + l_0(r) A_7 \vec{L} \cdot \vec{S} + l_0(r) A_8 \left[\vec{\sigma}_1 \cdot \vec{r} \vec{\sigma}_2 \cdot \vec{r} - \frac{1}{3} r^2 \vec{\sigma}_1 \cdot \vec{\sigma}_2 \right] \\
 & \quad - A_9 \left\{ e^{-\frac{\Lambda^2}{4} r^2}, \left[[\partial^r \otimes \partial^s]^2 \otimes [\sigma_1^p \otimes \sigma_2^q]^2 \right]^{00} \right\}
 \end{aligned}$$



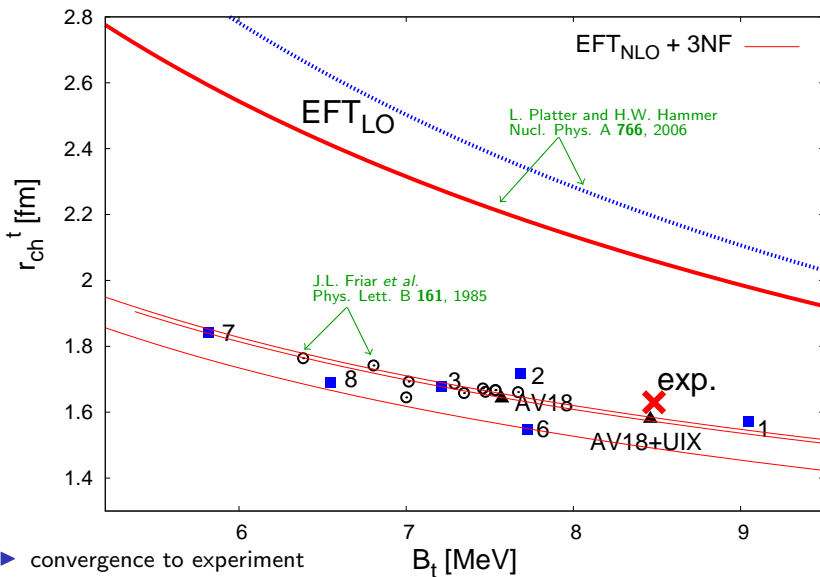
- ▶ Gaussian regulator functions $l_0(r, \Lambda) \propto e^{-\frac{\Lambda^2}{4} r^2}$
- ▶ low-energy-constants depend on cutoff, $A_i = A_i(\Lambda)$
- ▶ LEC of leading order 3NF **not** fitted

- ▶ low-energy input data: $B_d, \delta (^1S_0, ^3S_1, ^1,3P_{0,1,2}), \epsilon_1$ for $E_{\text{cm}} < 1$ MeV
- ▶ two methods to obtain different LEC sets \leftrightarrow different short-range physics:
 - ▶ variation of **cutoff** parameter Λ
 - ▶ variation of **low-energy input**

NN phase shifts at NLO



- ▶ deviation from $\delta^{Nij}(^{1,3}S_{0,1})$ is $\lesssim 10\%$
- ▶ deviation from ϵ_1^{Nij} is $\lesssim 30\%$
- ▶ P-wave phase shifts fitted to various values $< \delta^{Nij}(P)$



▶ convergence to experiment

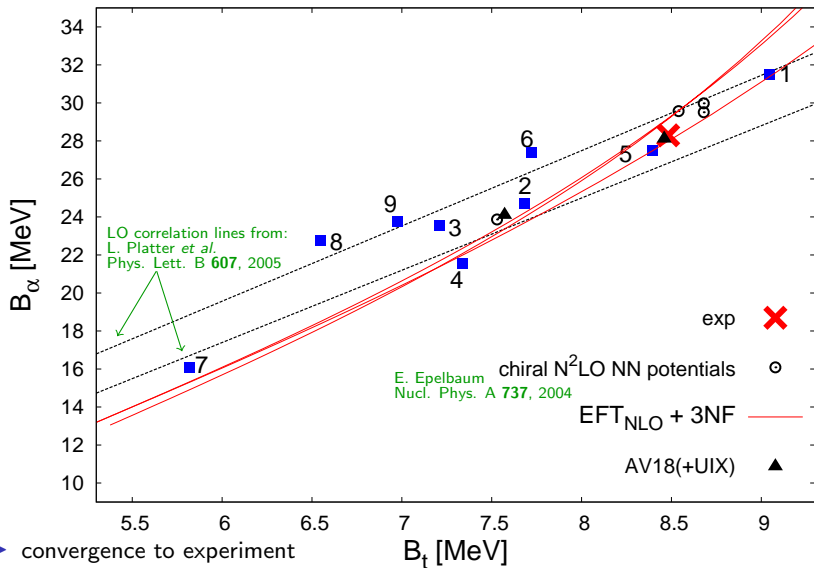
▶ $LO \rightarrow NLO \Rightarrow Q \approx \frac{1}{3}$

▶ NLO band width $\Rightarrow Q \approx \frac{1}{3}$

B_t [MeV]

▶ one three-body parameter needed for a result independent of short distance physics

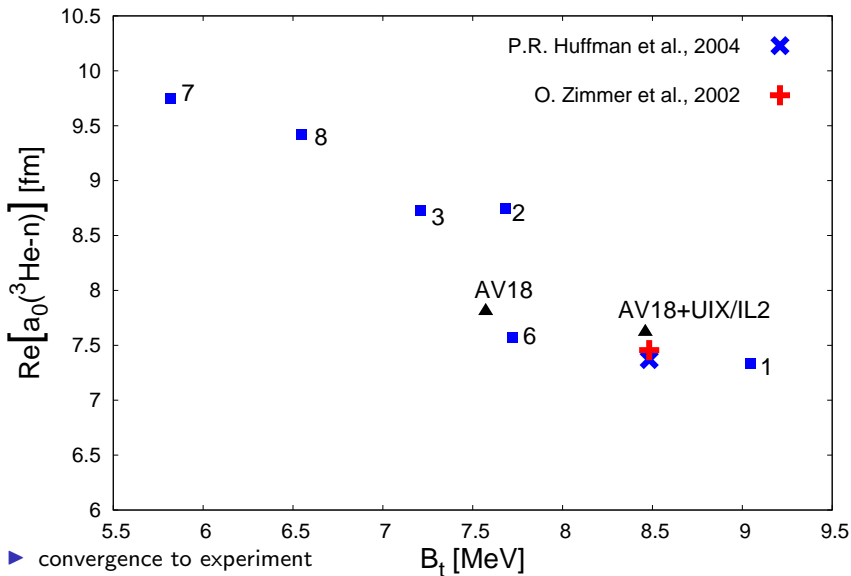
▶ $EFT_{\not{\tau}}$ works up to NLO



- ▶ convergence to experiment
- ▶ NLO band width $\Rightarrow Q \approx \frac{1}{3}$
- ▶ **no** four-body force needed at NLO

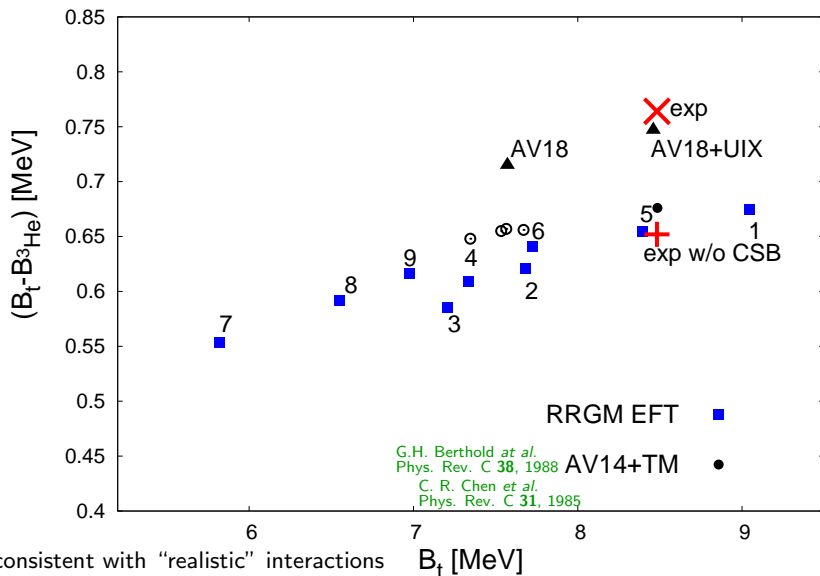
B_t [MeV]

- ▶ one three-body parameter needed for a result independent of short distance physics
- ▶ $EFT_{\not{\tau}}$ still works up to NLO



- ▶ convergence to experiment
- ▶ NLO band width $\Rightarrow Q \approx \frac{1}{3}$
- ▶ **no** four-body force needed at NLO

- ▶ calculation **not accurate enough** to discriminate between conflicting data
- ▶ EFT $\not{\neq}$ still² works up to NLO



- ▶ consistent with “realistic” interactions
- ▶ discrepancy to experiment due to higher order CSB interactions
- ▶ missing contribution consistent with χ PT calculations

Conclusions

- ▶ evidence for the **applicability** of EFT($\not\neq$) at NLO in the ^4He -channel
 - ▶ **convergence** from LO to NLO with $Q \approx \frac{1}{3}$
 - ▶ new correlation for a_0 ($^3\text{He} - n$)
 - ▶ consistency with high precision models & experiment
 - ▶ **four-nucleon** contact interaction **not** necessary at NLO
- ▶ model independent calculation of CIB/CSB effect of the Coulomb interaction in ^3H and ^3He
- ▶ RRGM and EFT($\not\neq$) *resonate* well



The “next order”

- ▶ universality in **halo nuclei**
- ▶ **electro-weak** processes in the four-nucleon system

RRGM for scattering states: $\delta \left(\langle \psi_\lambda | \hat{H} - E | \psi_\lambda \rangle - \frac{1}{2} a_{\lambda\lambda} \right) = 0$

fragment wave functions

$$\psi_{SS,\lambda}^{J\pi} = \mathcal{A} \sum_j^{n_k} \left[\frac{1}{R_j} Y_{L_j}(\hat{R}_j) \otimes [\psi_j^{J_1^{\pi_1}} \otimes \psi_j^{J_2^{\pi_2}}] S_{c_j} \right]^J \left(\delta_{\lambda j} F_{L_j}(R_j) + a_{\lambda j} \tilde{G}_{L_j}(R_j) + \underbrace{\sum_m b_{\lambda j m} R_j^{L_j+1} e^{-\omega_{jm} R_j^2}}_{\text{approximates state in interaction region}} \right)$$

open & distortion channels Coulomb functions

$$|{}^4\text{He}\rangle = \left[\left| \begin{array}{c} \gamma_{i1} \\ L=0,2 \\ S_{12}=1 \end{array} \right\rangle \right]_{L_{\text{rel}}} \left| \begin{array}{c} \tilde{\gamma}_{i1} \\ L=0,2 \\ S_{12}=1 \end{array} \right\rangle \Bigg|_{J=0^+} + \left[\left| \begin{array}{c} \gamma_{i1} \\ L=2 \\ S_{12}=1 \\ S=\frac{3}{2} \end{array} \right\rangle \right]_{L=0} \left| \begin{array}{c} \gamma_{i2} \\ L=0 \\ J_2=\frac{1}{2} \end{array} \right\rangle \Bigg|_{J=0^+} + \dots$$

open physical channels

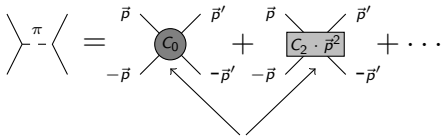
$$\left[\left| \begin{array}{c} \tilde{\gamma}_{i1} \\ L=0 \\ S_{12}=0 \end{array} \right\rangle \right]_{L_{\text{rel}}} \left| \begin{array}{c} \tilde{\gamma}_{i1} \\ L=0 \\ S_{12}=0 \end{array} \right\rangle \Bigg|_{J=0^+} + \left[\left| \begin{array}{c} \gamma_{i1} \\ L=2 \\ S_{12}=1 \\ S=\frac{3}{2} \end{array} \right\rangle \right]_{L=0} \left| \begin{array}{c} \tilde{\gamma}_{\text{rel}} \\ L=0 \\ J_2=\frac{1}{2} \end{array} \right\rangle \Bigg|_{J=0^+} + \dots$$

distortion & unphysical channels

narrow widths, i.e. no asymptotic tail

The Effective Field Theory:

- ▶ theory breaks down as momenta approach pion mass m_π
- ▶ (unnaturally) small deuteron binding energy relative to the triplet $n - p$ effective range
- ▶ non-relativistic isospin doublet of Pauli spinors: $N = \begin{pmatrix} p \\ n \end{pmatrix}$ LO: $B_d = \frac{1}{a_t^2 M} \ll \frac{m_\pi^2}{M}$
- ▶ Lorentz symmetry at small momentum



truncation based on two estimates:

coupling strength & relative contribution of graphs

naïve dimensional analysis & naturalness

- ▶ encode short distance physics
- ▶ match to underlying theory or low-energy data

$$\mathcal{L}_{\not{n}} = N^\dagger \left(iD_0 + \frac{\mathbf{D}^2}{2M} \right) N + C_0 (N^\dagger N)(N^\dagger N) + C_2 (N^\dagger N)(N^\dagger \partial_i^2 N) + \dots$$

$\frac{1}{a} = \mathcal{O}(p_{\text{typ}}) \Rightarrow$ certain operators require resummation
others can still be treated perturbatively

The two philosophies:

- ▶ power counting on the Lagrangean level **and** on the graph level:

$$\begin{aligned} \textcircled{T} &= \textcircled{C_0} + \textcircled{C_0} \textcircled{C_0} + \textcircled{C_0} \textcircled{C_0} \textcircled{C_0} + \dots && \text{LO} \\ &\textcircled{C_2} + \textcircled{C_2} \textcircled{C_0} + \textcircled{C_2} \textcircled{C_0} \textcircled{C_0} + \dots && \text{NLO} \end{aligned}$$

- ▶ iterate **effective potential** by solving Schrödinger/Lippmann-Schwinger equation:

$$\begin{aligned} \textcircled{T} &= \textcircled{C_0} + \textcircled{C_0} \textcircled{C_0} + \textcircled{C_0} \textcircled{C_0} \textcircled{C_0} + \dots && \text{LO} \\ &\textcircled{C_2} + \textcircled{C_4} \textcircled{C_2} + \textcircled{C_2} \textcircled{C_2} \textcircled{C_0} + \dots && \text{"NLO"} \end{aligned}$$

In **both** approaches higher order terms are expected to contribute $\mathcal{O}\left[\left(\frac{p_{\text{typ}}}{m_\pi}\right)^{N>n}\right]$ at order n