



Gesellschaft für Schwerionenforschung, Darmstadt,
Theory Division

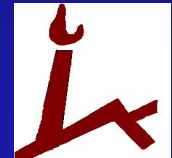
Application of the LIT Method to Electromagnetic Reactions with Nuclei: Recent Results

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Application of the LIT Method to Electromagnetic Reactions with Nuclei: Recent Results

Outline

- Photodisintegration of Light Nuclei
 - Six- and Seven-body Reactions with Semirealistic Potentials
 - Realistic Description of ^4He Photoabsorption
- Electron Scattering off ^4He :
 - Transverse Response Function with a Semirealistic Interaction
 - Longitudinal Response Function with AV18+UIX
- Conclusions and Outlook

LIT + EHH

Solve the Schrödinger-like equation

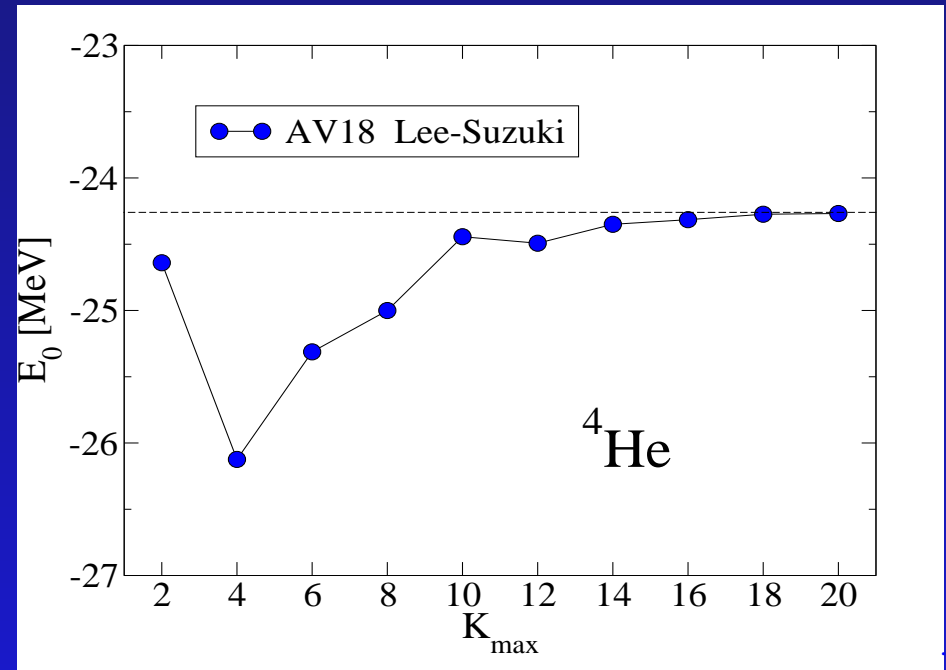
$$(H - E_0 + \sigma) |\tilde{\Psi}\rangle = \hat{O} |\Psi_0\rangle \quad \text{LIT equation}$$

using an **Hyperspherical Harmonics** expansion of the w.f.

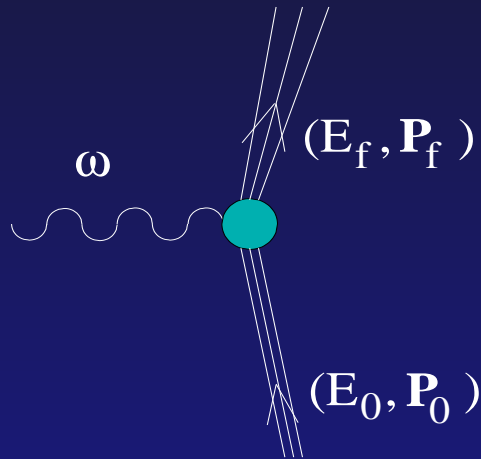
$$\Psi = \sum_{n [\mathbf{K}]}^{n_{\max} \mathbf{K}_{\max}} c_n^{[\mathbf{K}]} e^{-\frac{\rho}{2}} \rho^{\frac{\nu}{2}} L_n^{\nu}(\rho) [\mathcal{Y}_{[\mathbf{K}]}^{\mu}(\Omega) \chi_{ST}^{\bar{\mu}}]_{JT}^a$$

We define an **Effective Interaction** according to the Lee-Suzuki method

- Similar to the NCSM, but uses HH basis instead of HO.
- HO parameter dependence
- fast convergence
- can be used for $A > 4$



Photoabsorption of Nuclei



Real Photon

$$|\mathbf{q}| = \omega$$

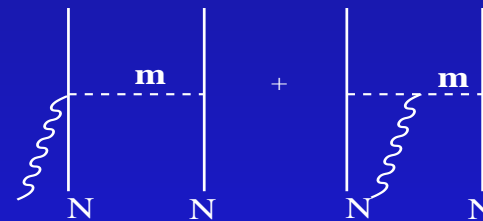
$$\sigma(\omega) = 4\pi^2 \alpha \omega R(\omega)$$

$$R(\omega) = \sum_f |\langle \Psi_f | \mathbf{E1} | \Psi_0 \rangle|^2 \delta(E_f - E_0 - \omega)$$

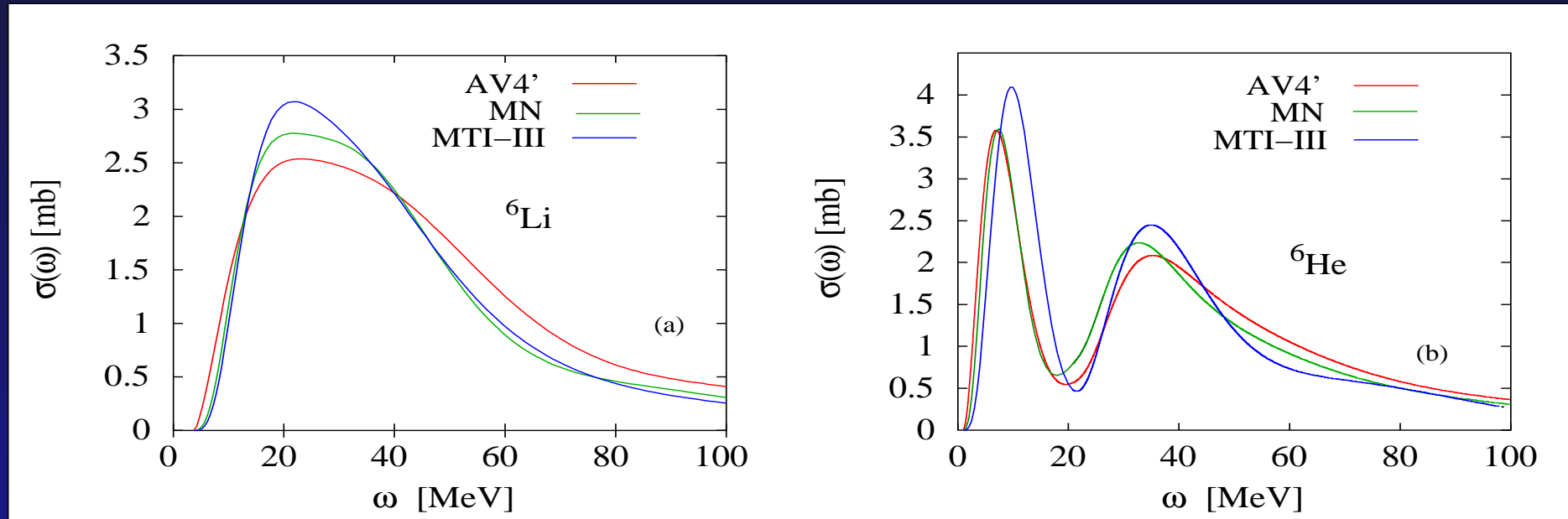
$$\mathbf{E1} = \sum_i^A (\mathbf{r}_i - \mathbf{R}_{\text{CM}}) \frac{1 + \tau_i^3}{2}$$

The **MEC** are implicitly included in the **E1**-response via the Siegert theorem

Alternatively....

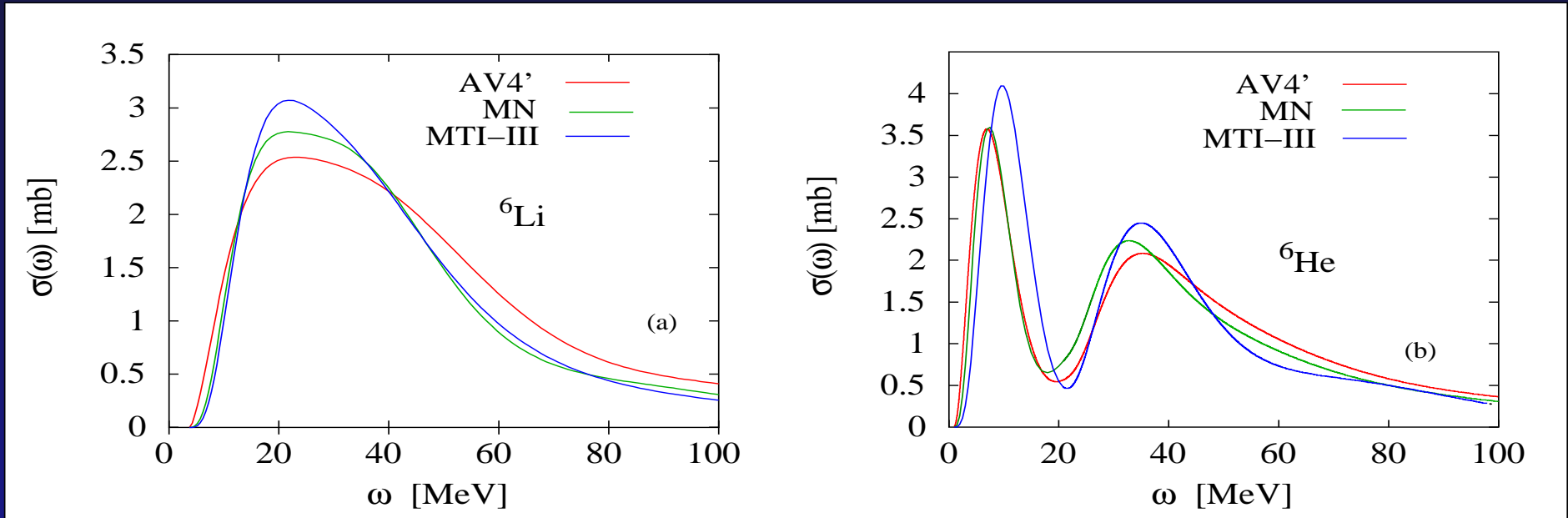


Six-body Photoabsorption



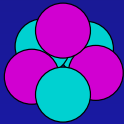
S. B., M. A. Marchisio, N. Barnea, W. Leidemann, G. Orlandini, Phys. Rev. Lett. **89** (2002)

Six-body Photoabsorption



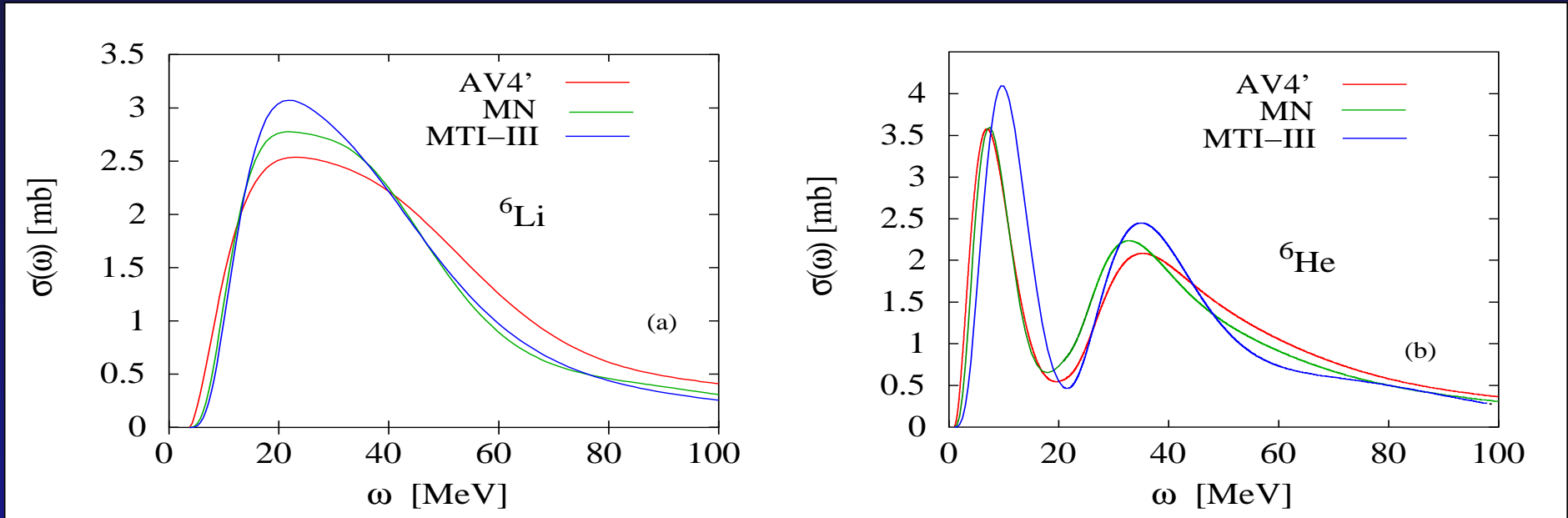
S. B., M. A. Marchisio, N. Barnea, W. Leidemann, G. Orlandini, Phys. Rev. Lett. **89** (2002)

Giant Dipole Mode



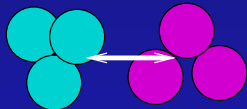
${}^6\text{Li}$

Six-body Photoabsorption



S. B., M. A. Marchisio, N. Barnea, W. Leidemann, G. Orlandini, Phys. Rev. Lett. **89** (2002)

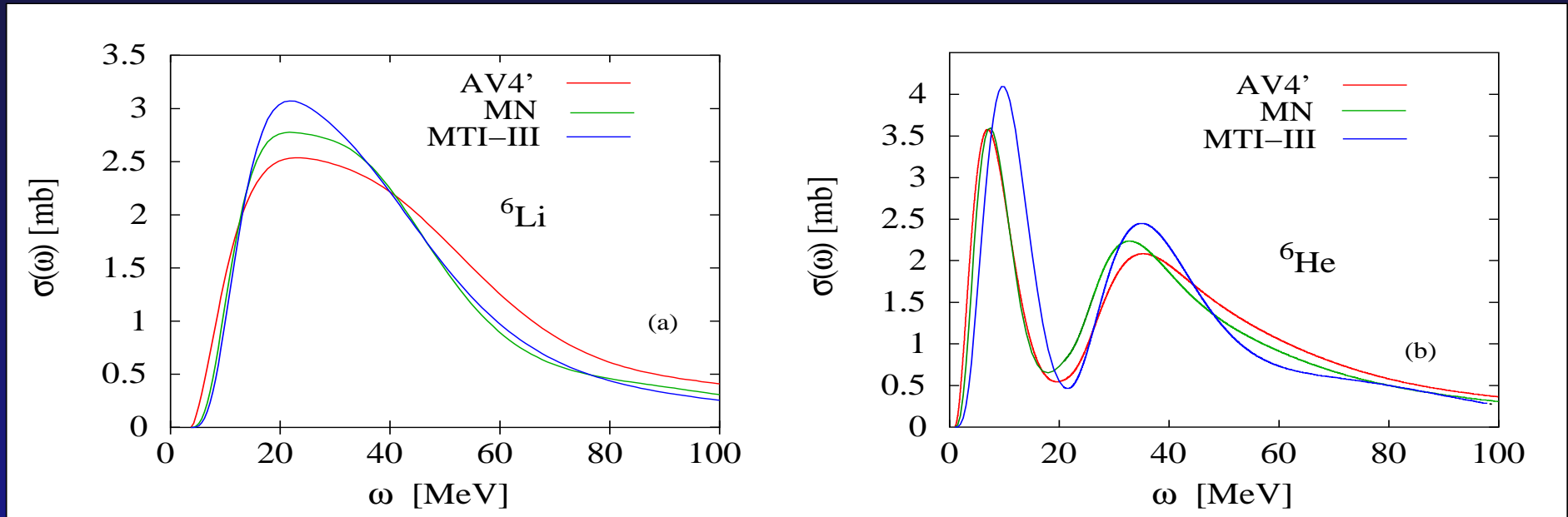
Giant Dipole Mode



neutrons protons

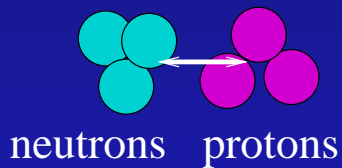
${}^6\text{Li}$

Six-body Photoabsorption



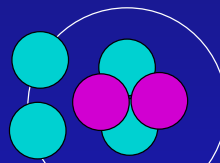
S. B., M. A. Marchisio, N. Barnea, W. Leidemann, G. Orlandini, Phys. Rev. Lett. **89** (2002)

Giant Dipole Mode



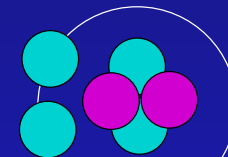
${}^6\text{Li}$

Soft Dipole Mode



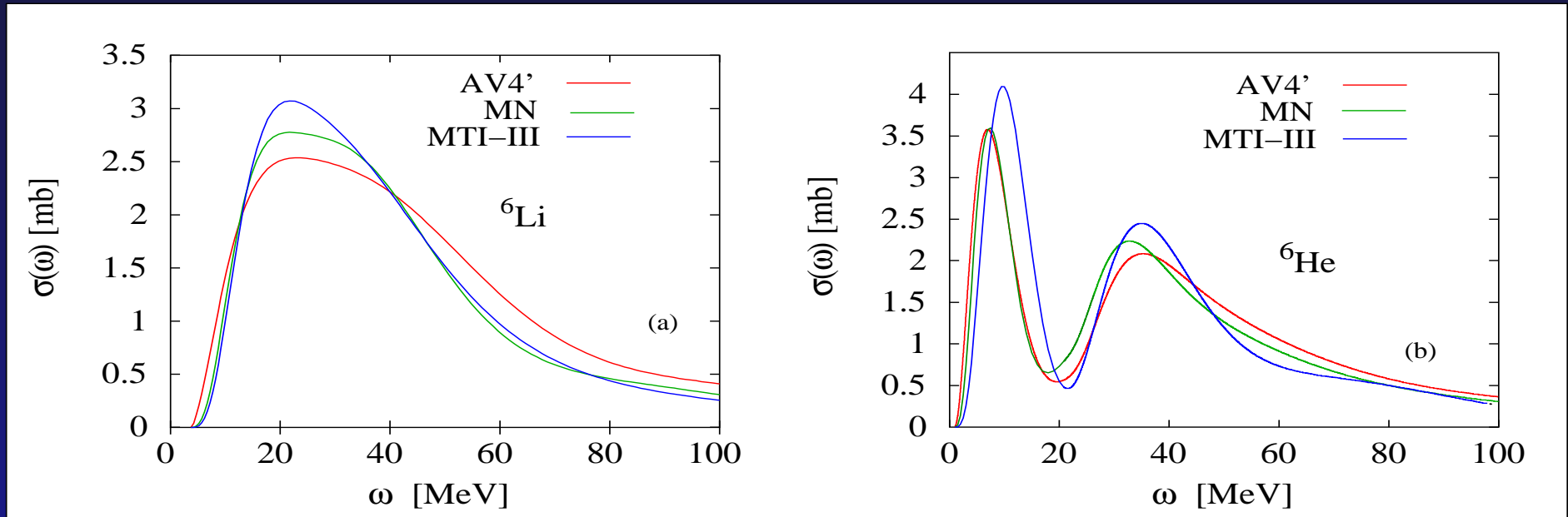
${}^6\text{He}$

Giant Dipole Mode



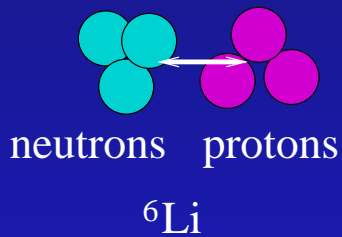
${}^6\text{He}$

Six-body Photoabsorption

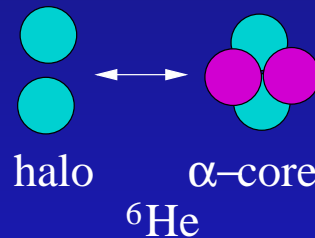


S. B., M. A. Marchisio, N. Barnea, W. Leidemann, G. Orlandini, Phys. Rev. Lett. **89** (2002)

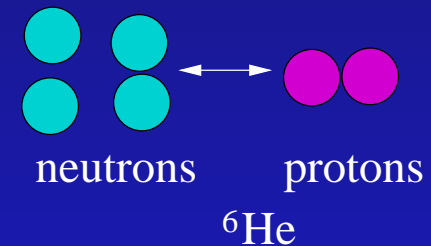
Giant Dipole Mode



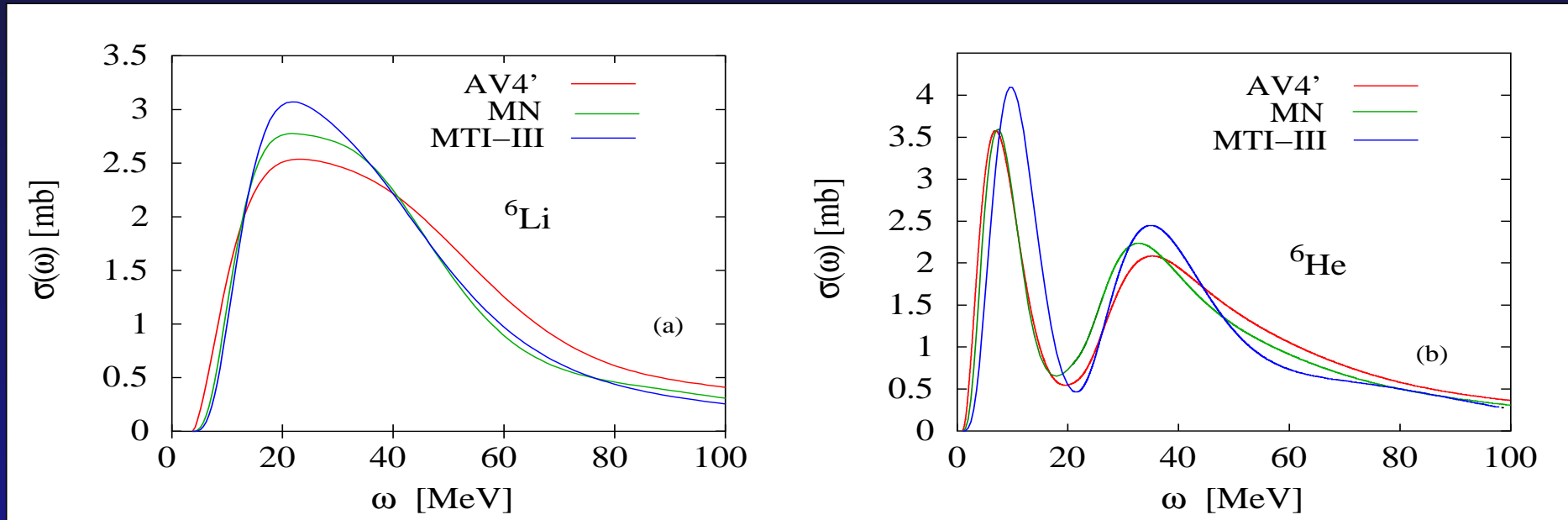
Soft Dipole Mode



Giant Dipole Mode

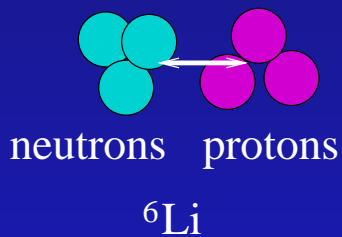


Six-body Photoabsorption

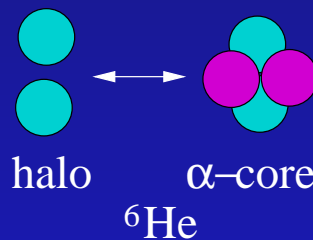


S. B., M. A. Marchisio, N. Barnea, W. Leidemann, G. Orlandini, Phys. Rev. Lett. **89** (2002)

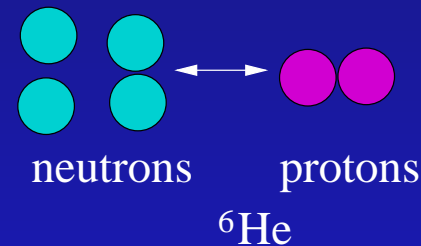
Giant Dipole Mode



Soft Dipole Mode



Giant Dipole Mode



AV4'

$$\sqrt{\langle r^2 \rangle} \text{ fm}$$

$$\sqrt{\langle r_p^2 \rangle} \text{ fm}$$

2.41(5)

1.80(4)

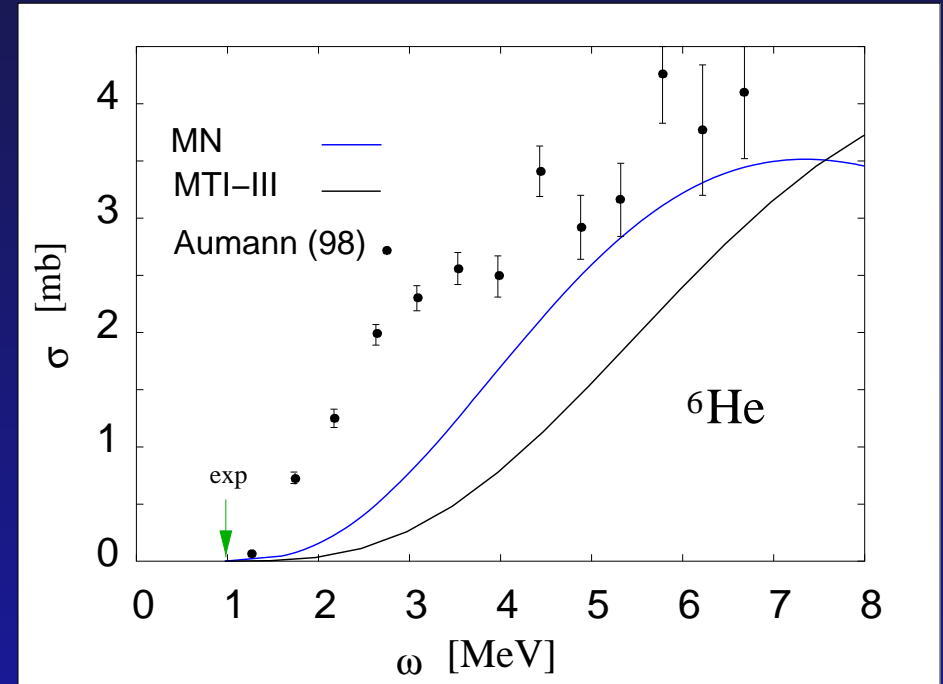
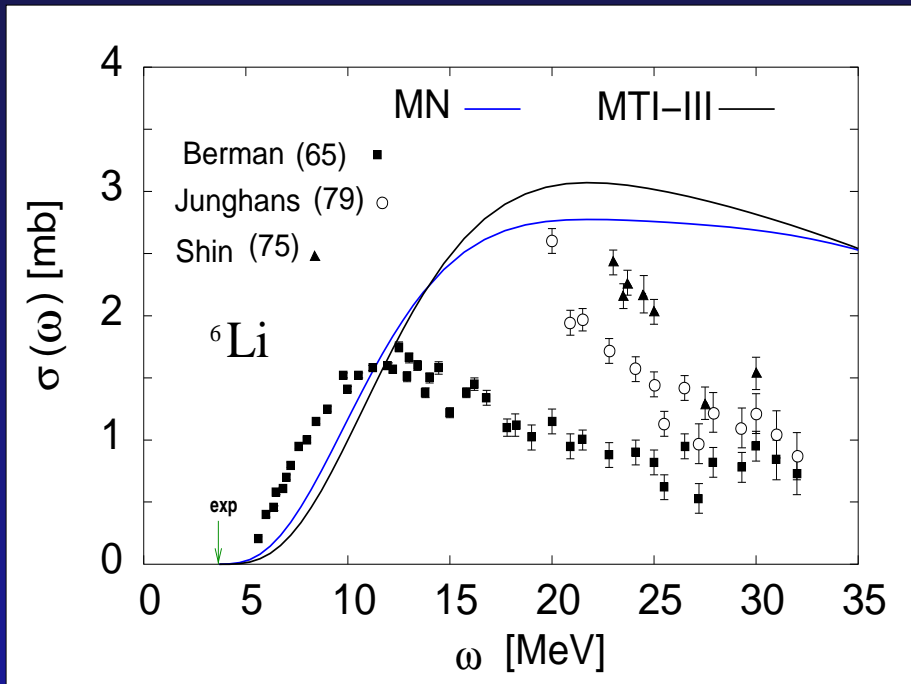
Exp

PRL **93** (2004) 142501

1.912 ± 0.018

Six-body Photoabsorption

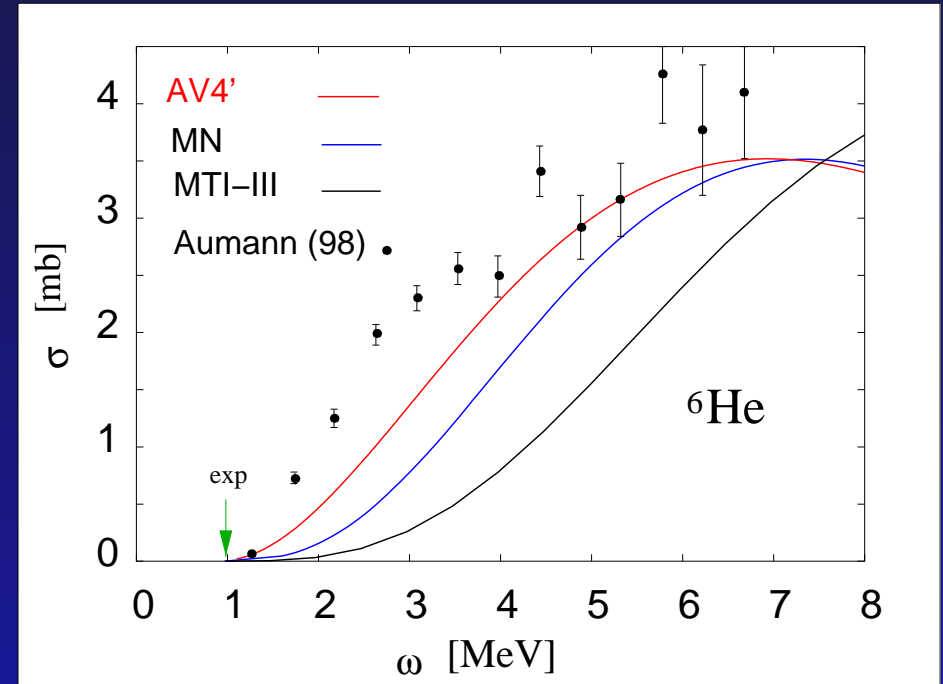
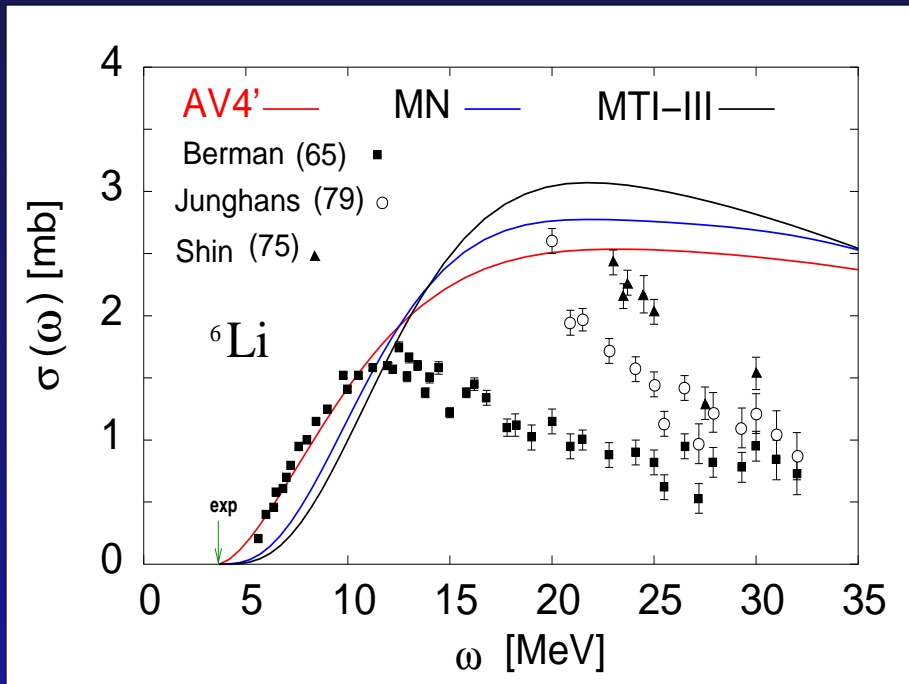
Comparison with experiment



- MTI-III and MN are dominant S -wave potentials (L -even) interactions

Six-body Photoabsorption

Comparison with experiment

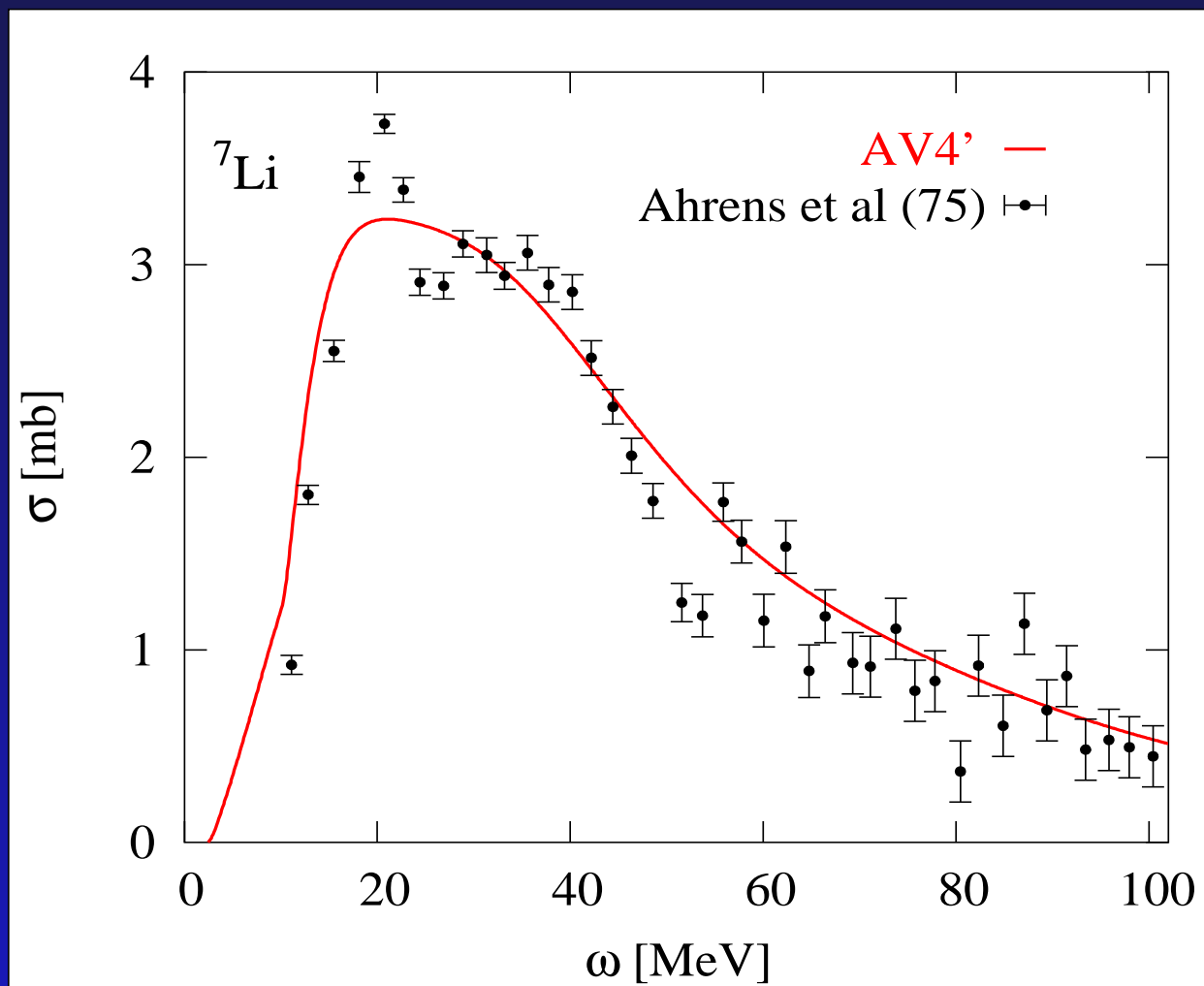


- MTI-III and MN are dominant S -wave potentials (L -even) interactions
- AV4' contains also (L -odd) interactions
Dominant S - and P -wave Potential

S.B., N. Barnea, W. Leidemann, and G. Orlandini, Phys. Rev. C **69**, 057001 (2004)

Seven-body Photoabsorption

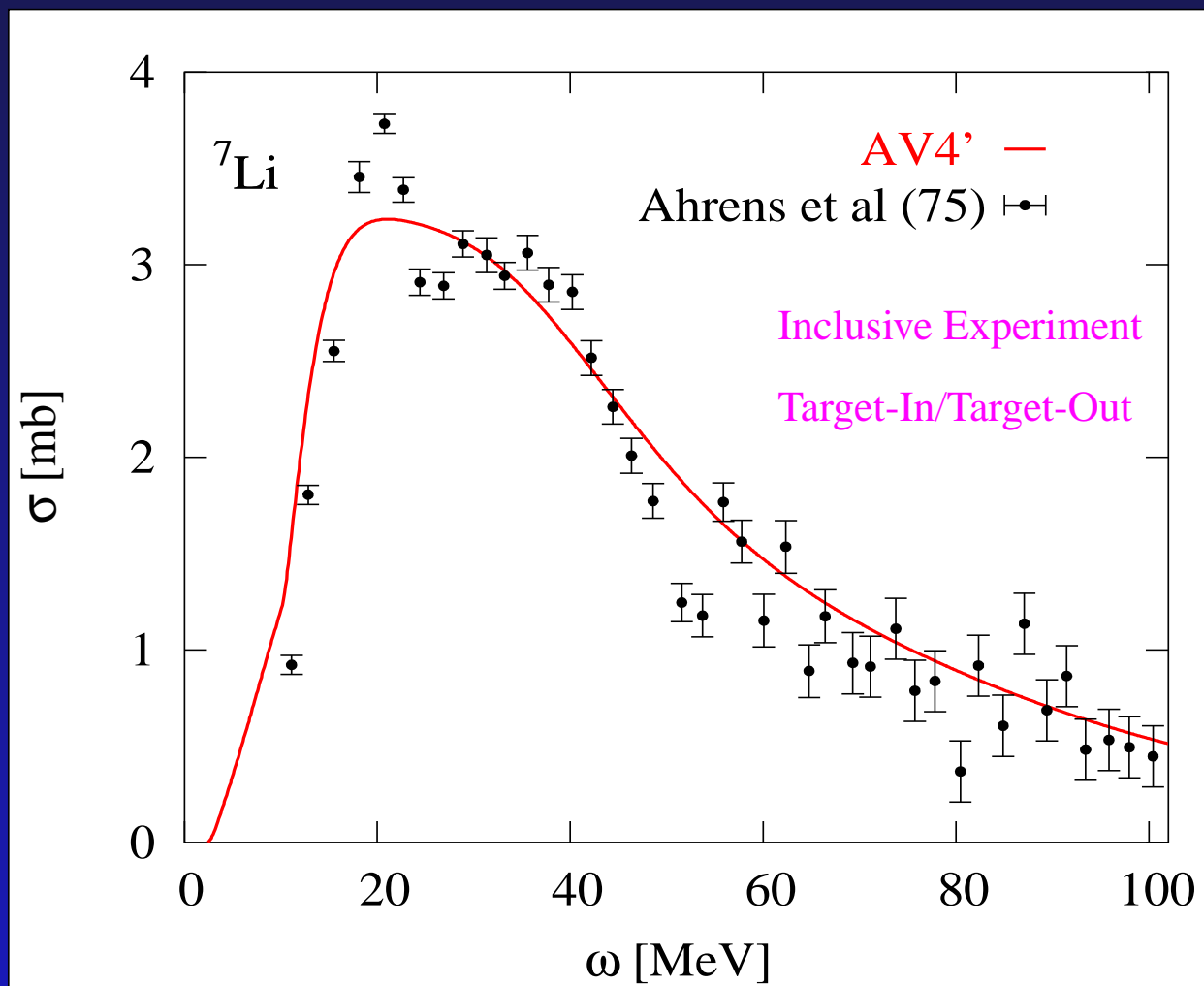
Comparison with experiment



S.B., H. Arenhövel, N. Barnea, W. Leidemann, and G. Orlandini, Phys. Lett. B **603** (2004)

Seven-body Photoabsorption

Comparison with experiment



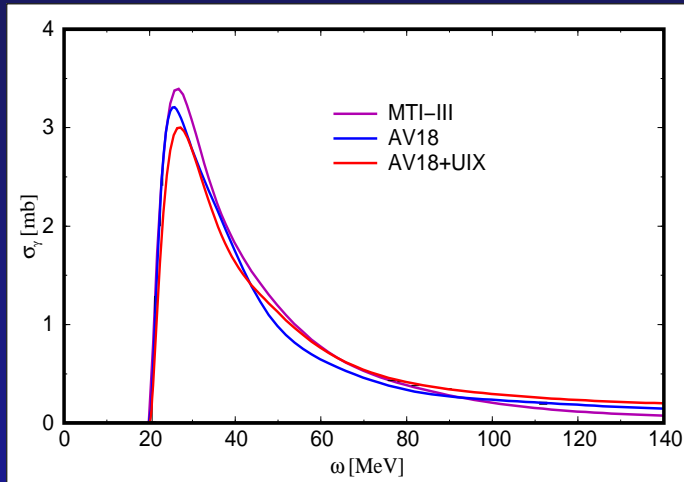
S.B., H. Arenhövel, N. Barnea, W. Leidemann, and G. Orlandini, Phys. Lett. B **603** (2004)

Photodisintegration of the α -particle

Calculation with realistic forces

AV18+UIX

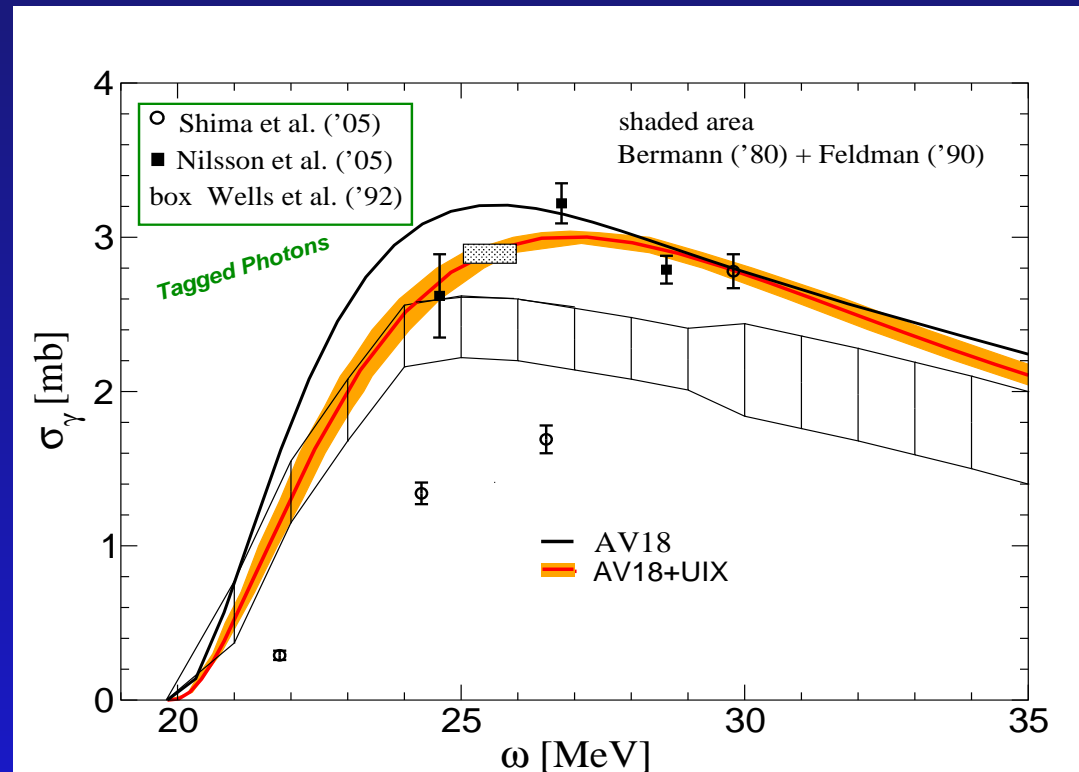
D.Gazit, S.B., N.Barnea, W.Leidemann, G.Orlandini, PRL **96**, 112301 (2006)



- MTI-III overestimates of 10 – 15% the peak and strongly underestimate the tail
- 3NF lead to a 6% peak damping and 35% tail enhancement

Comparison with experiment improves when adding 3NF.

Experimental situation should be improved!

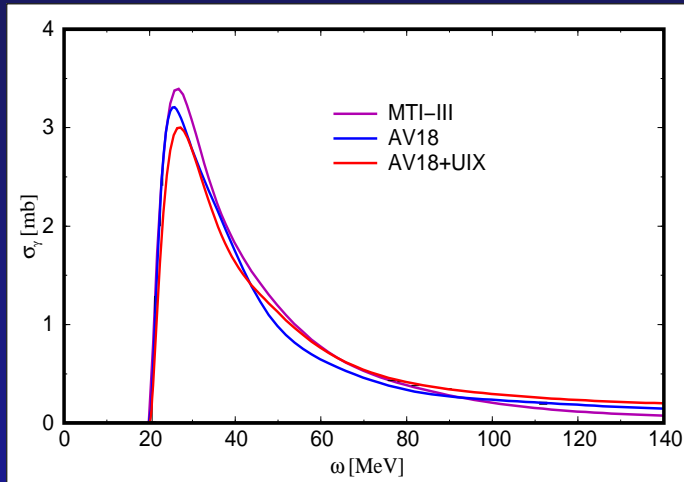


Photodisintegration of the α -particle

Calculation with realistic forces

AV18+UIX

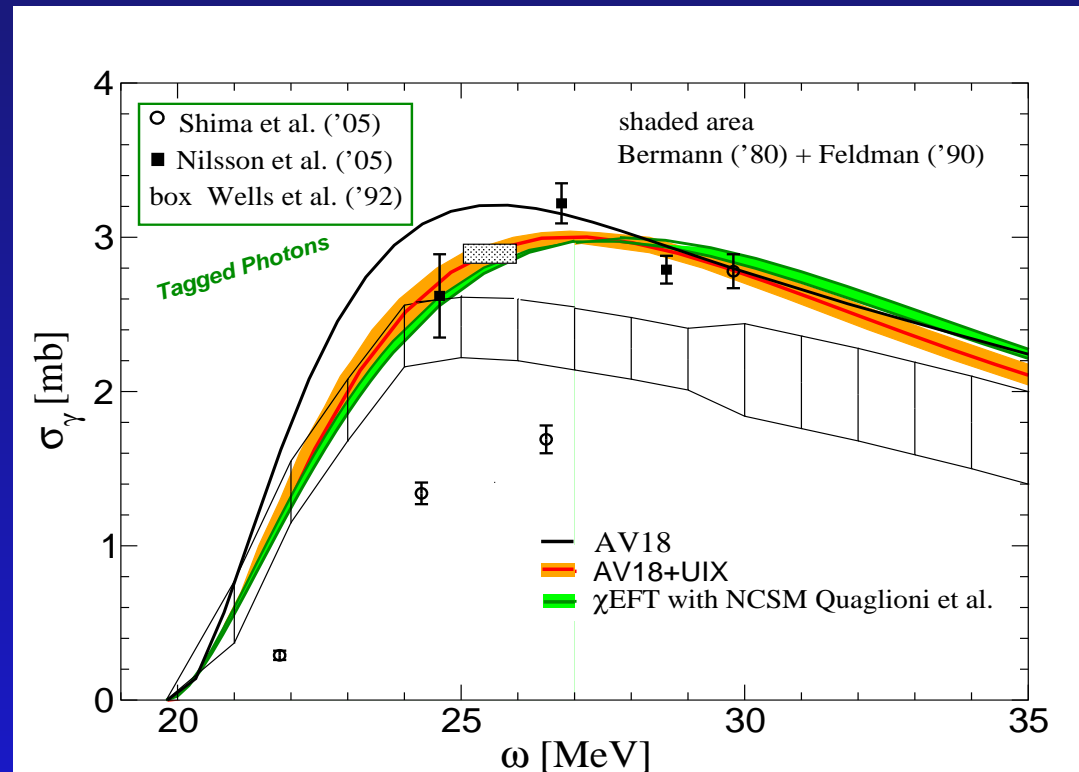
D.Gazit, S.B., N.Barnea, W.Leidemann, G.Orlandini, PRL **96**, 112301 (2006)



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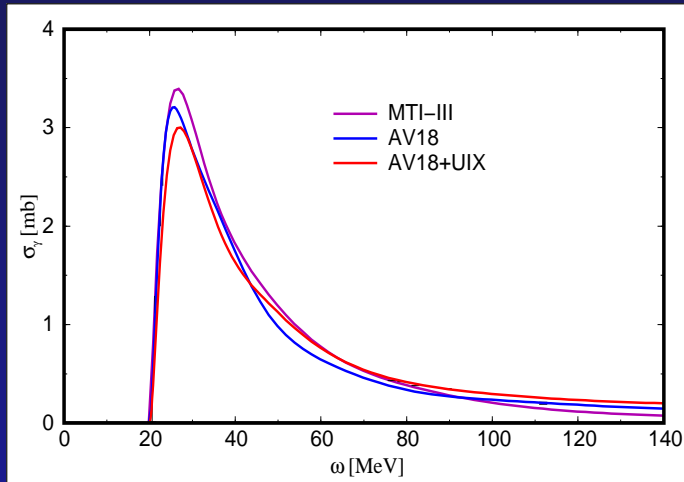
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Photodisintegration of the α -particle

Calculation with realistic forces

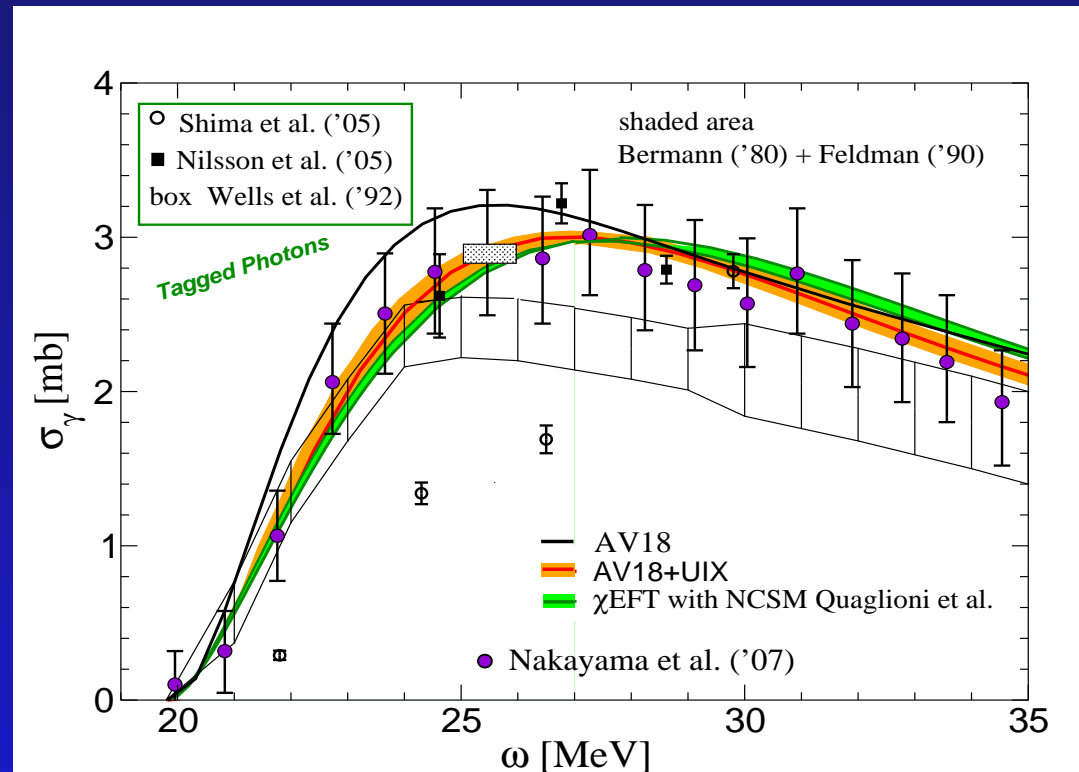
AV18+UIX D.Gazit, S.B., N.Barnea, W.Leidemann, G.Orlandini, PRL **96**, 112301 (2006)



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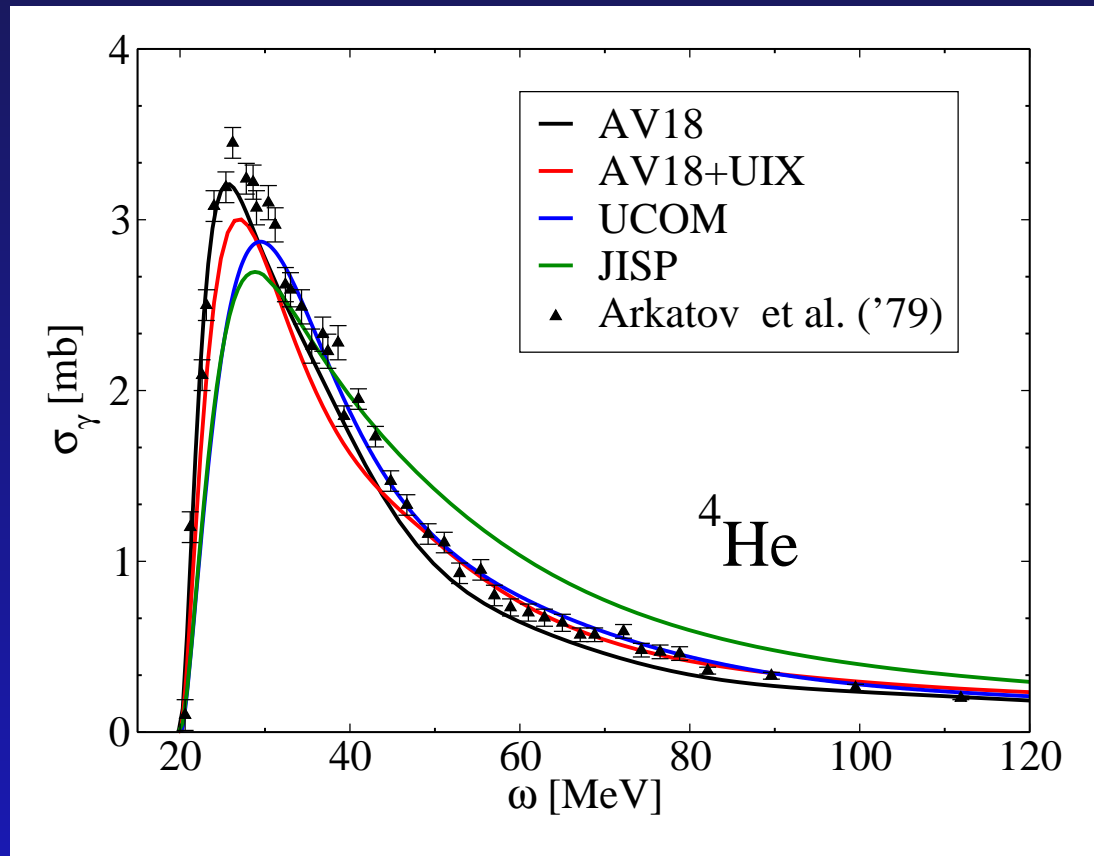
Experimental situation should be improved!



Photodisintegration of the α -particle

Calculation with realistic forces:

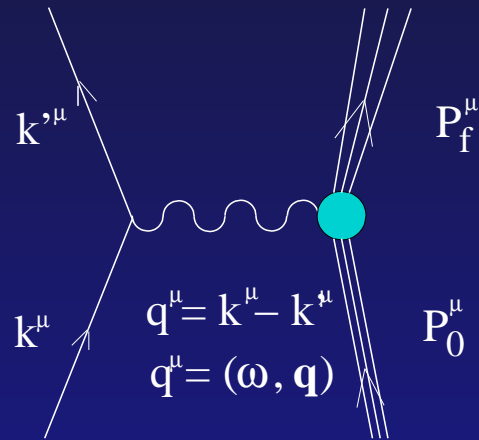
non-local interaction vs. 3NF



JISP N.Barnea, W.Leidemann, G.Orlandini, PRC **74**, 034003 (2006)

UCOM S.B., PRC **75**, 044001 (2007)

Electron Scattering off a nucleus



Virtual Photon
 (ω, \mathbf{q})
 vary independently

Inclusive Cross Section $A(e, e')X$

$$\frac{d^2\sigma}{d\Omega d\omega} = \sigma_M \left[\frac{Q^4}{\mathbf{q}^4} R_L(\omega, \mathbf{q}) + \left(\frac{Q^2}{2\mathbf{q}^2} + \tan^2 \frac{\theta}{2} \right) R_T(\omega, \mathbf{q}) \right]$$

with $Q^2 = -q_\mu^2 = \mathbf{q}^2 - \omega^2$ and θ scattering angle
 and σ_M Mott cross section

Electron Scattering off a nucleus

Longitudinal and transverse response functions can be disentangled via Rosenbluth separation

$$R_L(\omega, \mathbf{q}) = \sum_f |\langle \Psi_f | \rho(\mathbf{q}) | \Psi_0 \rangle|^2 \delta \left(E_f - E_0 - \omega + \frac{\mathbf{q}^2}{2M} \right)$$
$$R_T(\omega, \mathbf{q}) = \sum_f |\langle \Psi_f | \mathbf{J}_T(\mathbf{q}) | \Psi_0 \rangle|^2 \delta \left(E_f - E_0 - \omega + \frac{\mathbf{q}^2}{2M} \right)$$

In a **non relativistic** framework:

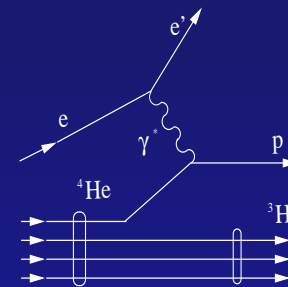
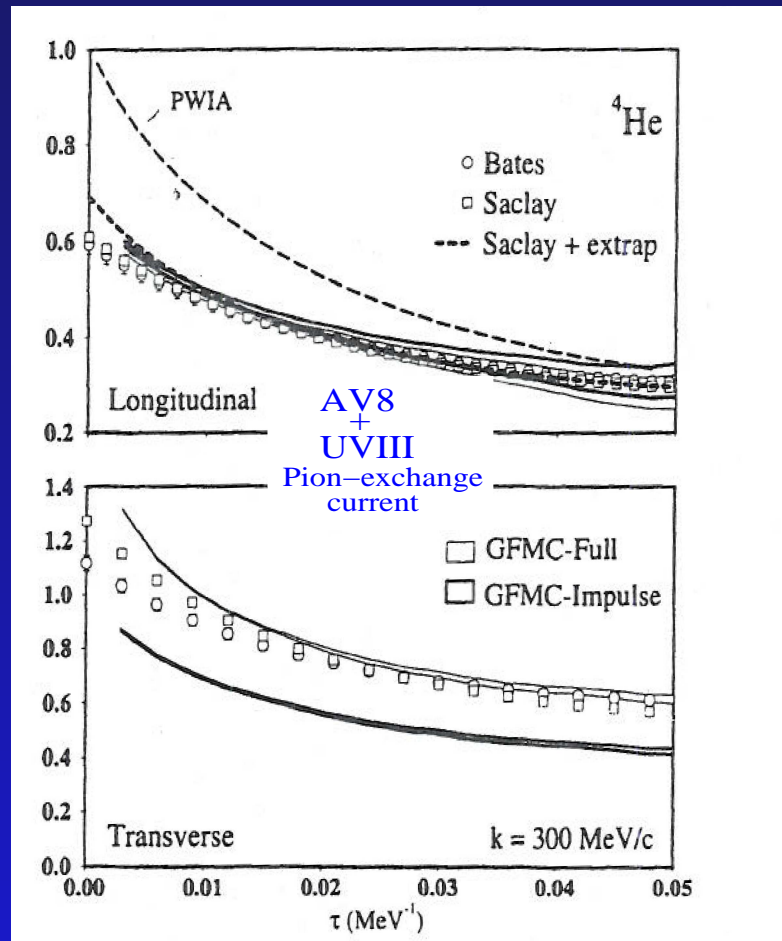
- The charge $\rho(\mathbf{q})$ is a one-body operator $\rho_1(\mathbf{q})$
- The current $\mathbf{J}_T(\mathbf{q})$ includes a one-body operator $\mathbf{J}_1(\mathbf{q})$, spin and convection current, and a two-body operator $\mathbf{J}_2(\mathbf{q})$, the so-called **Meson Exchange Current** required by gauge invariance

$$\nabla \cdot \mathbf{J}_2(\mathbf{x}) = -i [V, \rho_1(\mathbf{x})]$$

Previous Work

For ${}^4\text{He}$ the only realistic calculation with **FSI** and a consistent **MEC** is performed with the Laplace Transform using GFMC

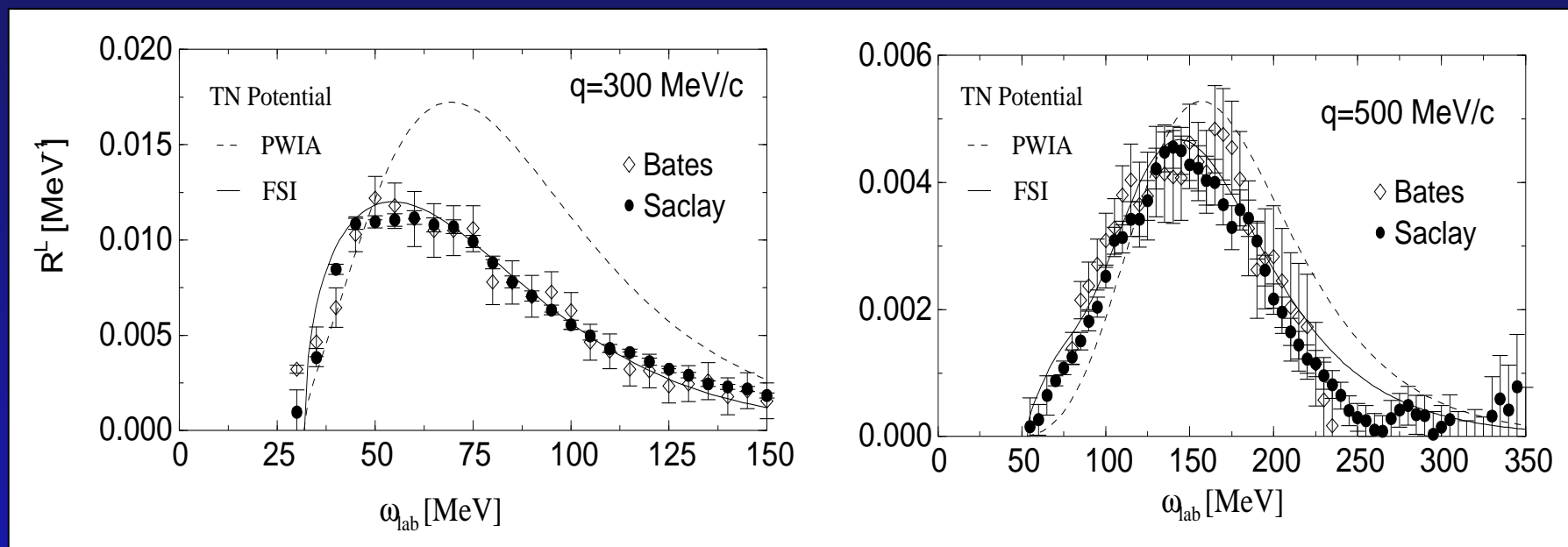
J. Carlson and R. Schiavilla, PRC 49, R2880 (1994)



Previous Work

Calculation for ${}^4\text{He}$ with the LIT method and CHH

Semirealistic TN potential

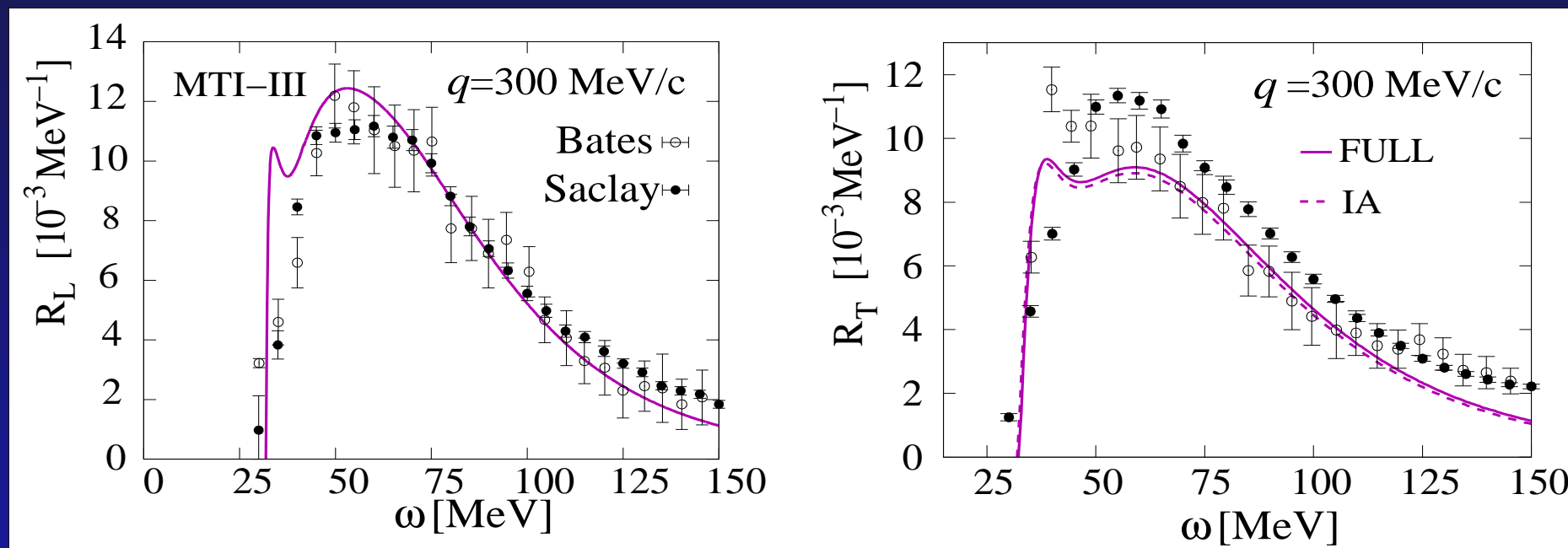


V. E. Efros, W. Leidemann and G. Orlandini, PRL **78**, 432 (1997).

Confirms the importance of **FSI** to correctly describe the reaction mechanism!

Electron Scattering with the LIT

Electron scattering off ^4He with a semirealistic forces



S.B., H.Arenhövel, N. Barnea, W. Leidemann, G. Orlandini, PRC 76, 014003 (2007)

MTI-III leads to a reasonable description of R_L

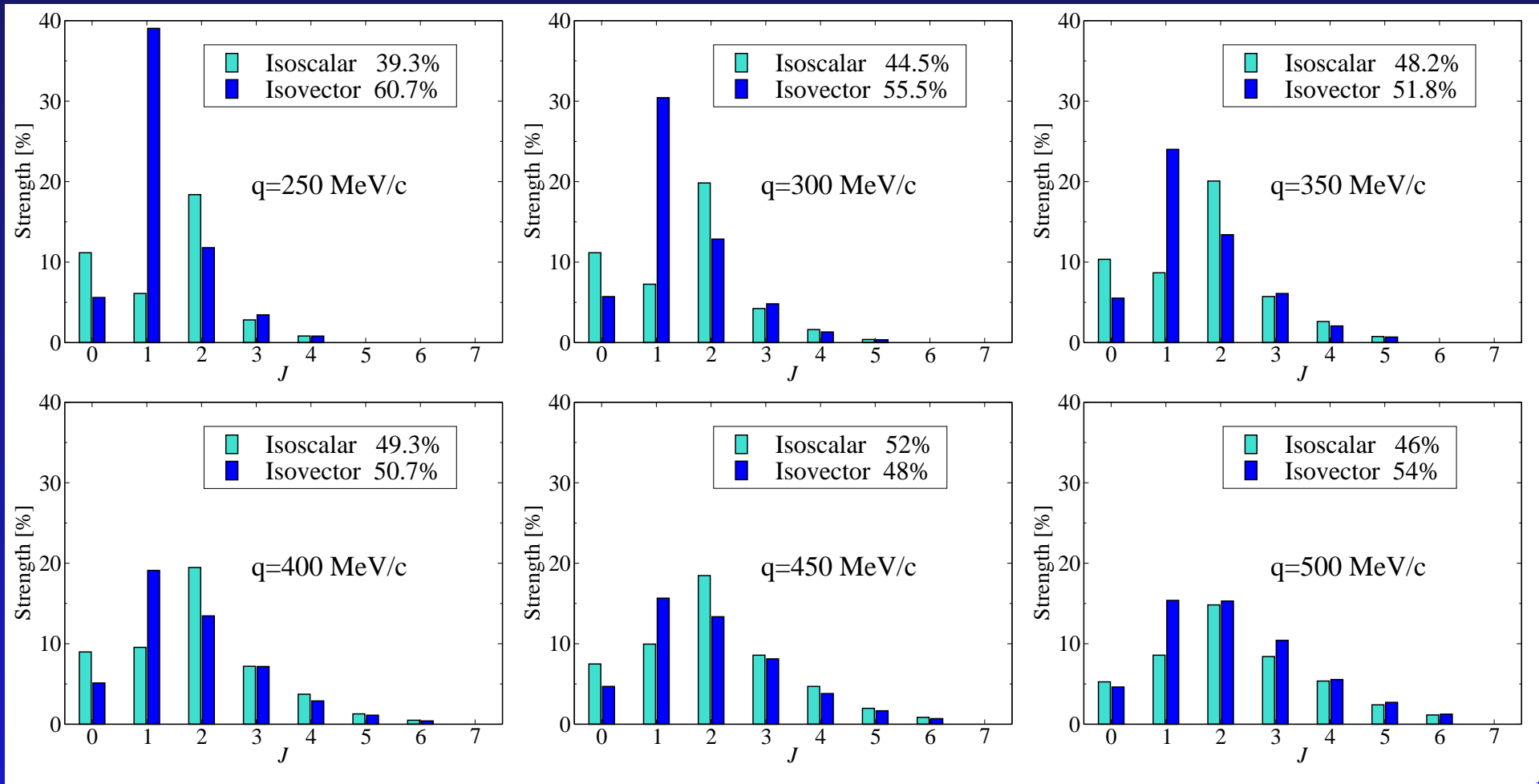
Small effect of MEC in R_T probably due to missing π -exchange

Longitudinal Response Function: AV18

$$\rho_{(1)}(\mathbf{q}, \omega) = \sum_k G_E^S(Q^2) e^{i\mathbf{q}\cdot\mathbf{r}'_k} \frac{1}{2} + G_E^V(Q^2) e^{i\mathbf{q}\cdot\mathbf{r}'_k} \frac{\tau_k^3}{2}, \text{ with } G_E^{S/V} = G_E^p \pm G_E^n$$

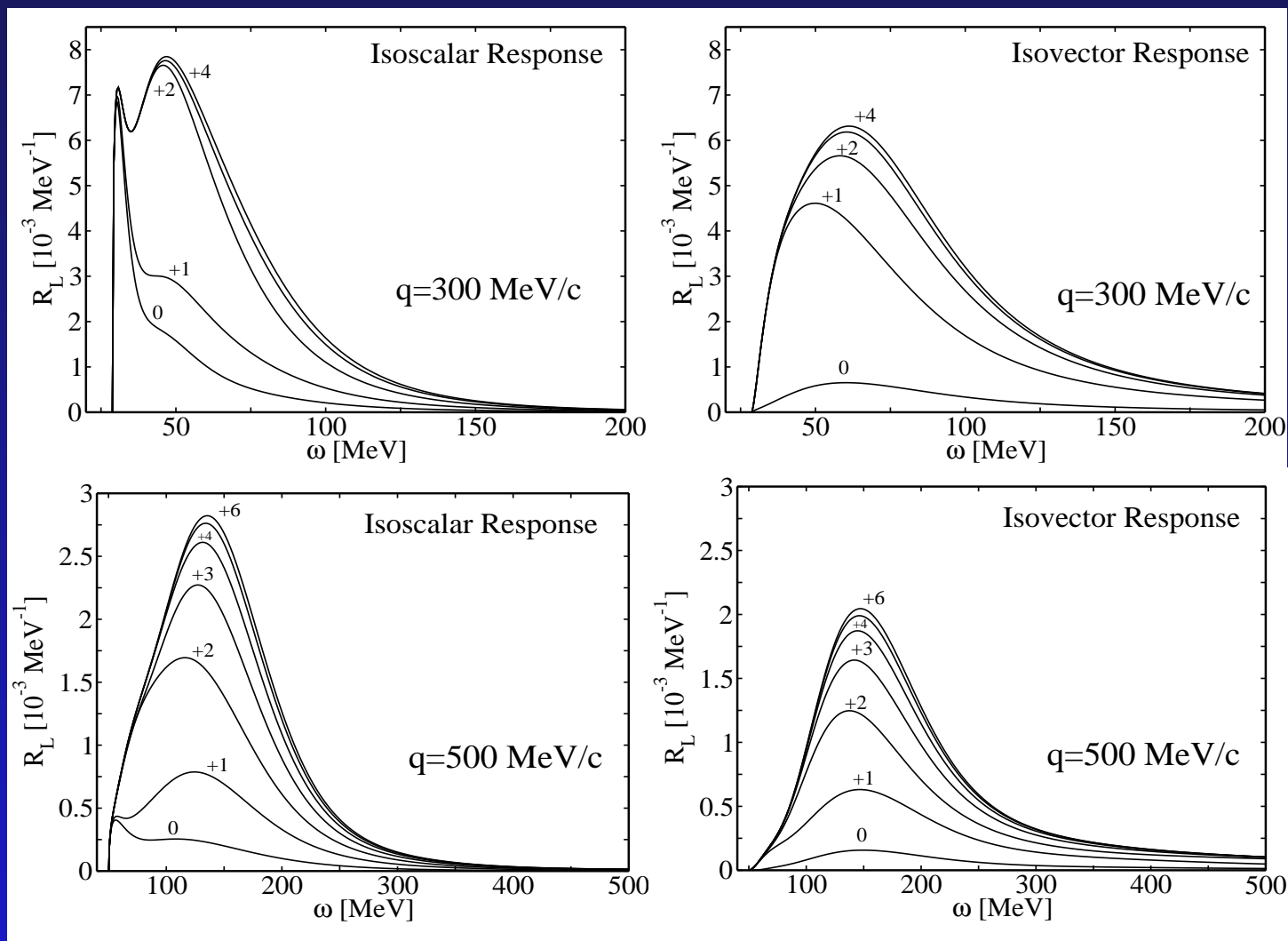
$$\rho_{(1)}(\mathbf{q}) = \sum_J C_J^S(\mathbf{q}) + C_J^V(\mathbf{q})$$

Distribution of the total strength among the multipoles



Longitudinal Response Function: AV18

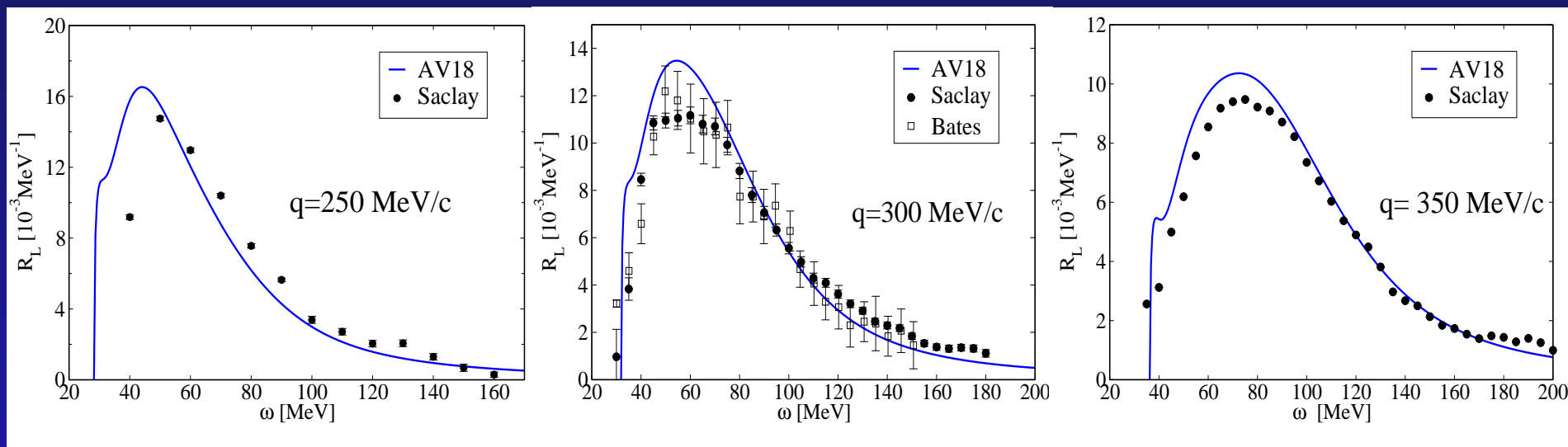
Distribution of the multipole strengths as function of the energy



Longitudinal Response Function: AV18

Comparison with experiment

Medium-low momentum transfer

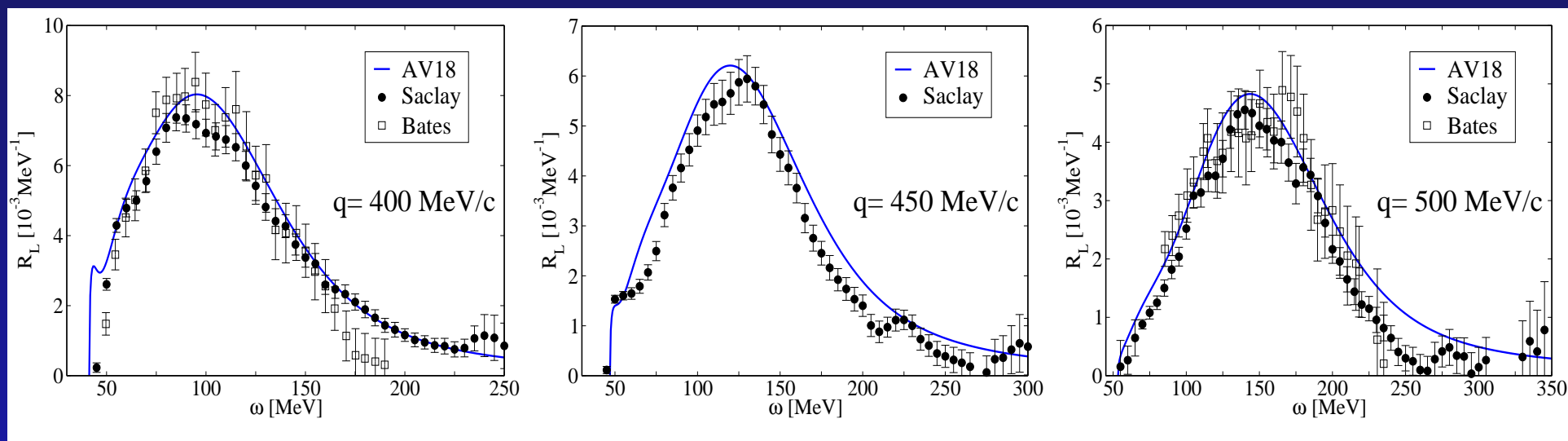


- With AV18 we find larger strength with respect to the experiment in the q.e. peak region.
- The shoulder at threshold is due to isoscalar monopole transition.

Longitudinal Response Function: AV18

Comparison with experiment

Medium-high momentum transfer

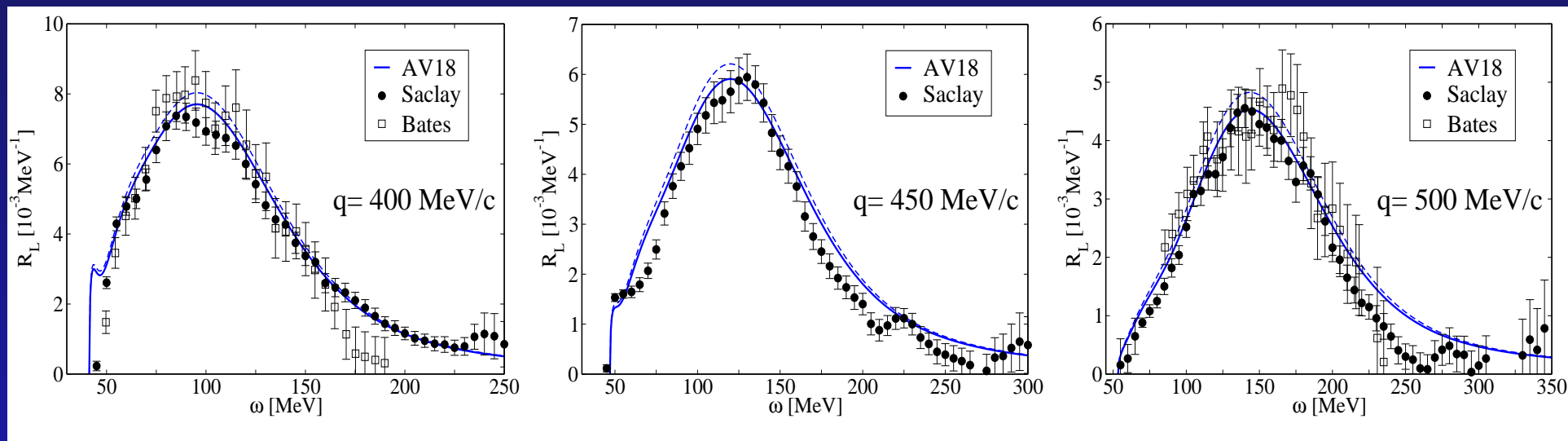


- With AV18 we find larger strength with respect to the experiment in the q.e. peak region.

Longitudinal Response Function: AV18

Comparison with experiment

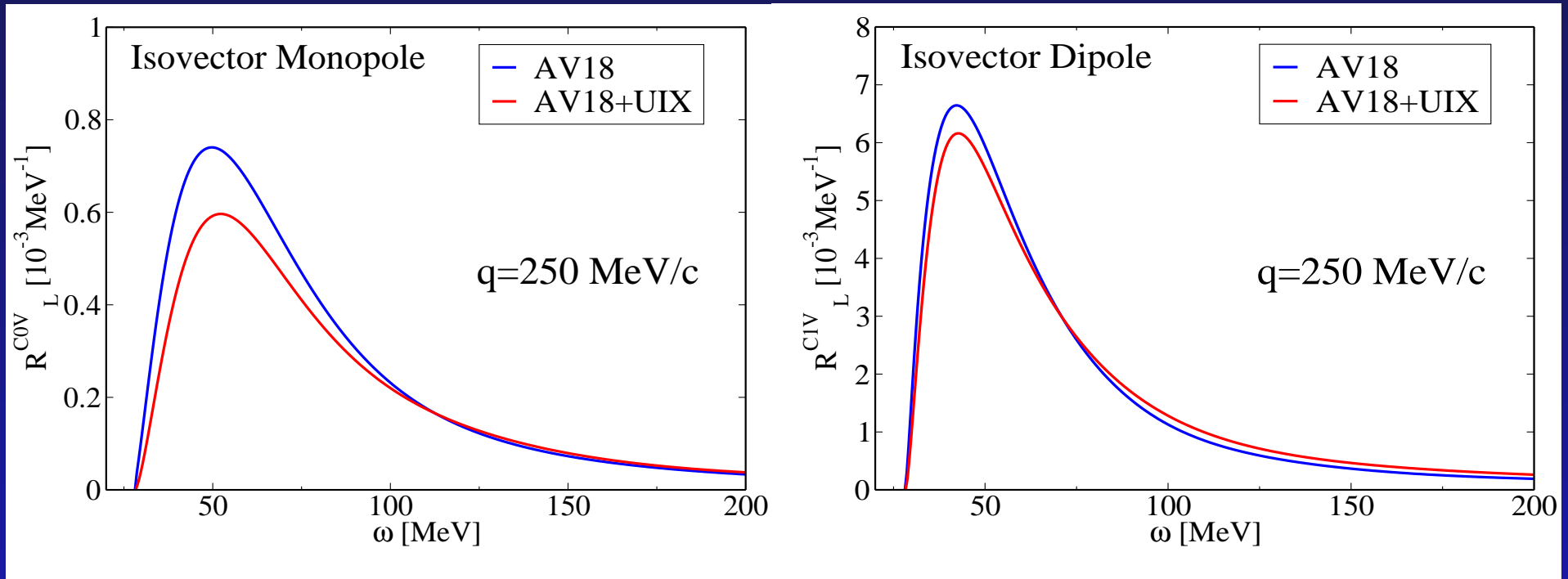
Medium-high momentum transfer: with relativistic form factors



- With **AV18** we find larger strength with respect to the experiment in the q.e. peak region.
- Considering the effect of relativity in the nucleon form factor the agreement with experiment improves

Addition of 3NF

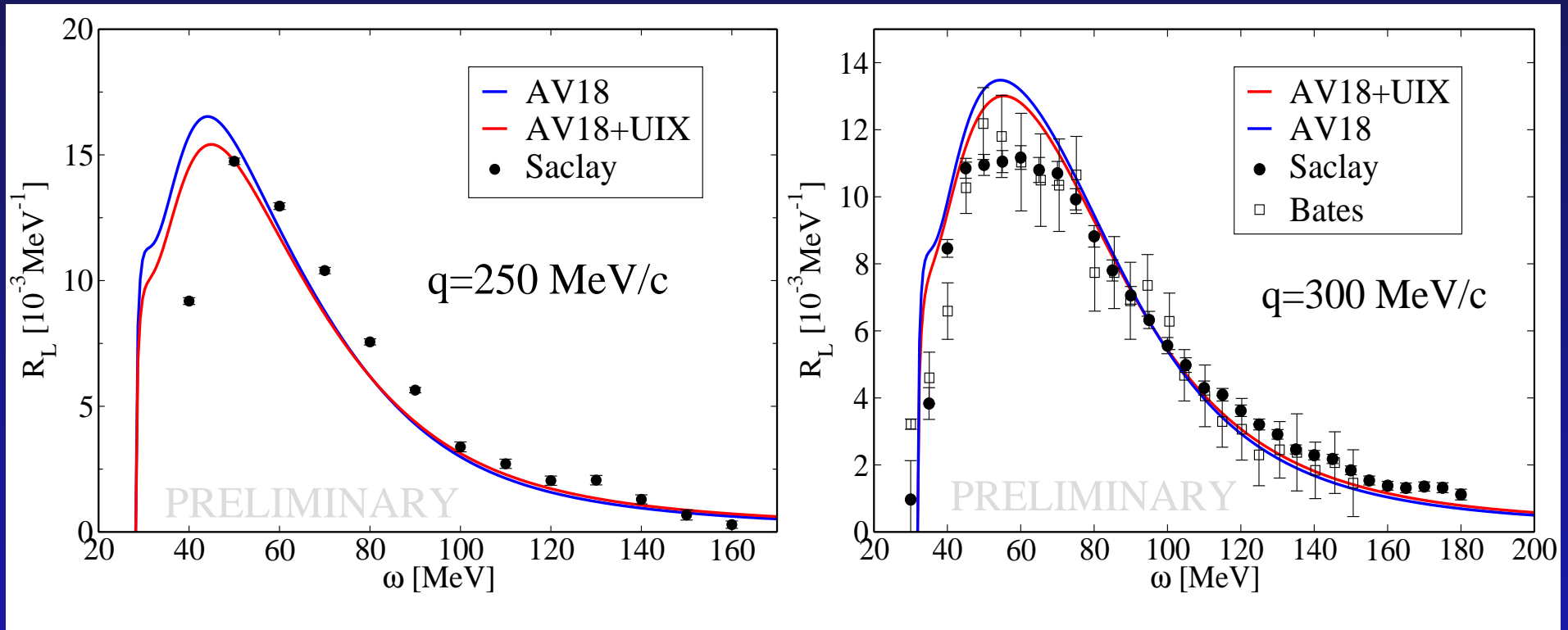
Isovector Monopole C_0^V and Dipole C_1^V responses



- The addition of UIX reduces the isovector monopole peak of 20%
- The addition of UIX reduces the isovector dipole peak of 7% and enhances the tail of 20-30% at pion-threshold (as for the photon!)

Longitudinal Response Function: AV18+UIX

Three-body force only in $C_0^{V/S}$ and C_1^V



- The preliminary addition of the UIX 3-body force improves the comparison with the experiment
- Larger effect of UIX is expected when including it in all multipoles.

Conclusion and Outlook

- We have presented a summary of the recent results obtained with the **LIT and EIH** methods, which demonstrate the power of this approach.
- Exact calculations of photoabsorption reactions have renewed the experimental interest in the field!
New measurement planned at MaxLab@Lund and TUNL.
- Investigation of electron scattering with realistic forces and consistent **MEC** can shed more light on our knowledge of 3NF and mesonic degrees of freedom
- **Future**: extend the “realistic” investigation of e.m. reactions to heavier system to investigate astrophysical reactions, like radiative capture
 ${}^2\text{H}(\alpha, \gamma){}^6\text{He}$, ${}^3\text{H}(\alpha, \gamma){}^7\text{Be}$, ${}^7\text{Be}(p, \gamma){}^8\text{B}$

Explicit Inclusion of MEC

- The MTI-III has the form

$$V = V_1(r) + V_2(r)\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 + V_3(r)\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2 + V_4(r)(\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2)(\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2)$$

where $V(r) \propto J_m(r)$ with J_m scalar meson propagator and

$$r = |\mathbf{r}_1 - \mathbf{r}_2|$$

Explicit Inclusion of MEC

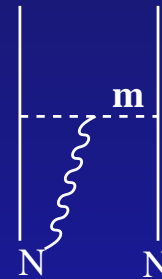
- The MTI-III has the form

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where $V(r) \propto J_m(r)$ with J_m **scalar meson propagator** and
 $r = |\mathbf{r}_1 - \mathbf{r}_2|$

- The consistent **MEC** has only **meson in flight** terms

$$\mathbf{J}_2(\mathbf{x}, \mathbf{r}_1, \mathbf{r}_2) \sim (\boldsymbol{\tau}_1 \times \boldsymbol{\tau}_2)_3 J_m(\mathbf{r}_1 - \mathbf{x}) \overleftrightarrow{\nabla}_{\mathbf{x}} J_m(\mathbf{x} - \mathbf{r}_2)$$



Explicit Inclusion of MEC

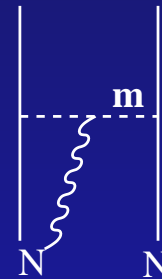
- The MTI-III has the form

$$V = V_1(r) + V_2(r)\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 + V_3(r)\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2 + V_4(r)(\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2)(\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2)$$

where $V(r) \propto J_m(r)$ with J_m **scalar meson propagator** and $r = |\mathbf{r}_1 - \mathbf{r}_2|$

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$$\mathbf{J}_2(\mathbf{x}, \mathbf{r}_1, \mathbf{r}_2) \sim (\boldsymbol{\tau}_1 \times \boldsymbol{\tau}_2)_3 J_m(\mathbf{r}_1 - \mathbf{x}) \overleftrightarrow{\nabla}_{\mathbf{x}} J_m(\mathbf{x} - \mathbf{r}_2)$$



- In momentum space

$$\mathbf{J}_2(\mathbf{q}, \mathbf{r}_1, \mathbf{r}_2) \longrightarrow \mathbf{J}_2(\mathbf{q}, \mathbf{r}, \mathbf{R}) = \frac{1}{4\pi^3} e^{i\mathbf{R} \cdot \mathbf{q}} (\nabla_{\mathbf{r}} I_m(\mathbf{q}, \mathbf{r}))$$

$$\mathbf{R} = \frac{\mathbf{r}_1 + \mathbf{r}_2}{2}, \quad \mathbf{r} = \mathbf{r}_1 - \mathbf{r}_2.$$

Explicit Inclusion of MEC

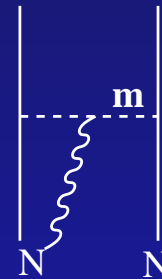
- The MTI-III has the form

$$V = V_1(r) + V_2(r)\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 + V_3(r)\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2 + V_4(r)(\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2)(\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2)$$

where $V(r) \propto J_m(r)$ with J_m **scalar meson propagator** and $r = |\mathbf{r}_1 - \mathbf{r}_2|$

- The consistent **MEC** has only **meson in flight** terms

$$\mathbf{J}_2(\mathbf{x}, \mathbf{r}_1, \mathbf{r}_2) \sim (\boldsymbol{\tau}_1 \times \boldsymbol{\tau}_2)_3 J_m(\mathbf{r}_1 - \mathbf{x}) \overleftrightarrow{\nabla}_{\mathbf{x}} J_m(\mathbf{x} - \mathbf{r}_2)$$



- In momentum space

$$\mathbf{J}_2(\mathbf{q}, \mathbf{r}_1, \mathbf{r}_2) \longrightarrow \mathbf{J}_2(\mathbf{q}, \mathbf{r}, \mathbf{R}) = \frac{1}{4\pi^3} e^{i\mathbf{R}\cdot\mathbf{q}} (\nabla_{\mathbf{r}} I_m(\mathbf{q}, \mathbf{r}))$$

$$\mathbf{R} = \frac{\mathbf{r}_1 + \mathbf{r}_2}{2}, \quad \mathbf{r} = \mathbf{r}_1 - \mathbf{r}_2.$$

- Approximation**

$$\mathbf{J}_2(\mathbf{q}, \mathbf{r}, \mathbf{R}) \simeq \mathbf{J}_2^A(\mathbf{q}, \mathbf{r}) = \frac{1}{4\pi^3} (\nabla_{\mathbf{r}} I_m(\mathbf{q}, \mathbf{r}))$$

Previous Work

- The quasi-elastic region $\omega \simeq \mathbf{q}^2/2m$ is seen as an interesting playground to study **nuclear correlations**.

Previous Work

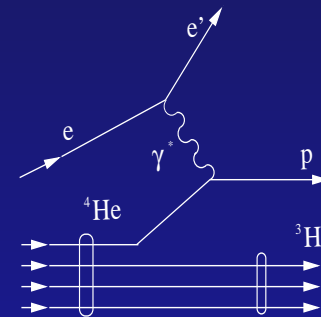
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They lead to:

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 - (2) Slight underestimation of the transverse response function
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- (2) probably due to the missing **Meson Exchange Currents**

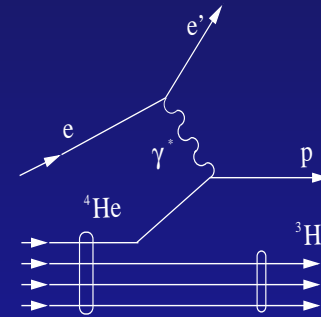


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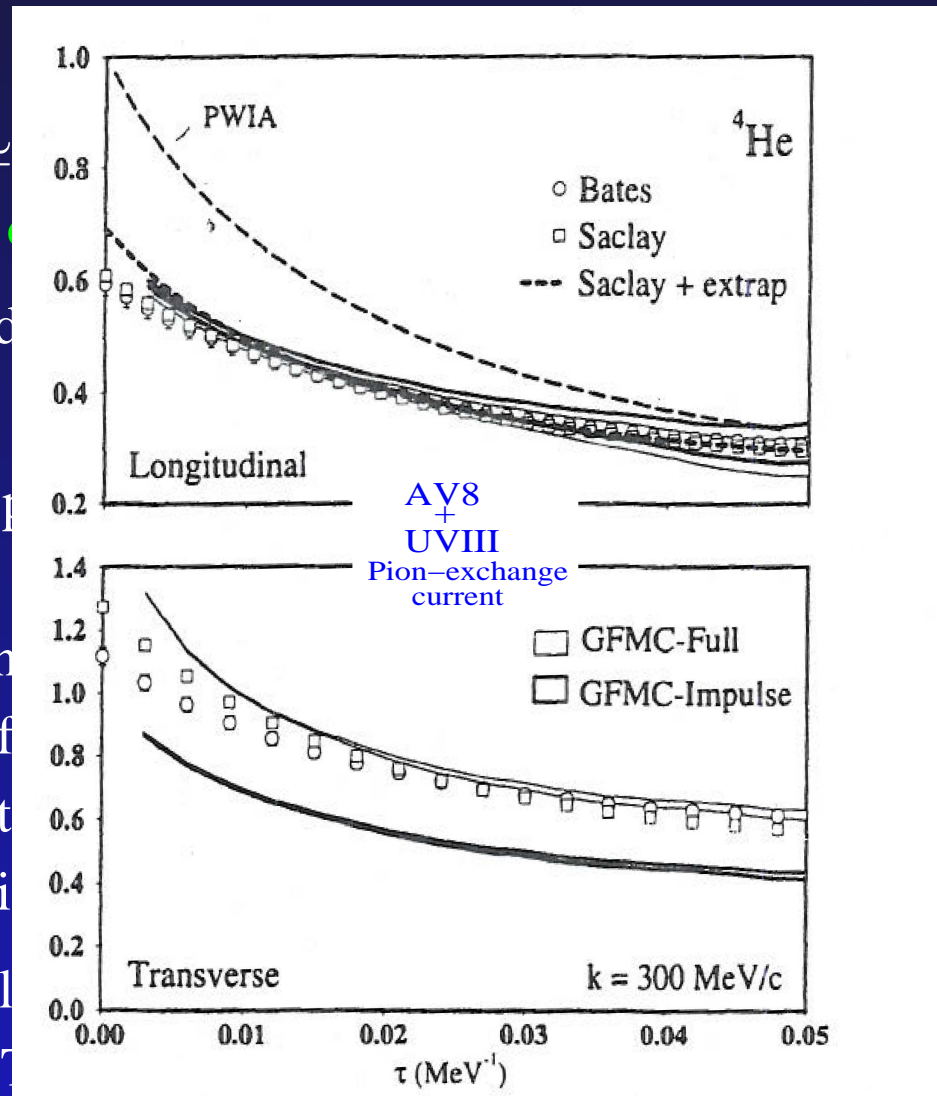


- For ${}^4\text{He}$ the only realistic calculation with **FSI** and a consistent **MEC** is performed with the Laplace Transform using GFMC

J. Carlson and R. Schiavilla, PRC **49**, R2880 (1994)

Previous Work

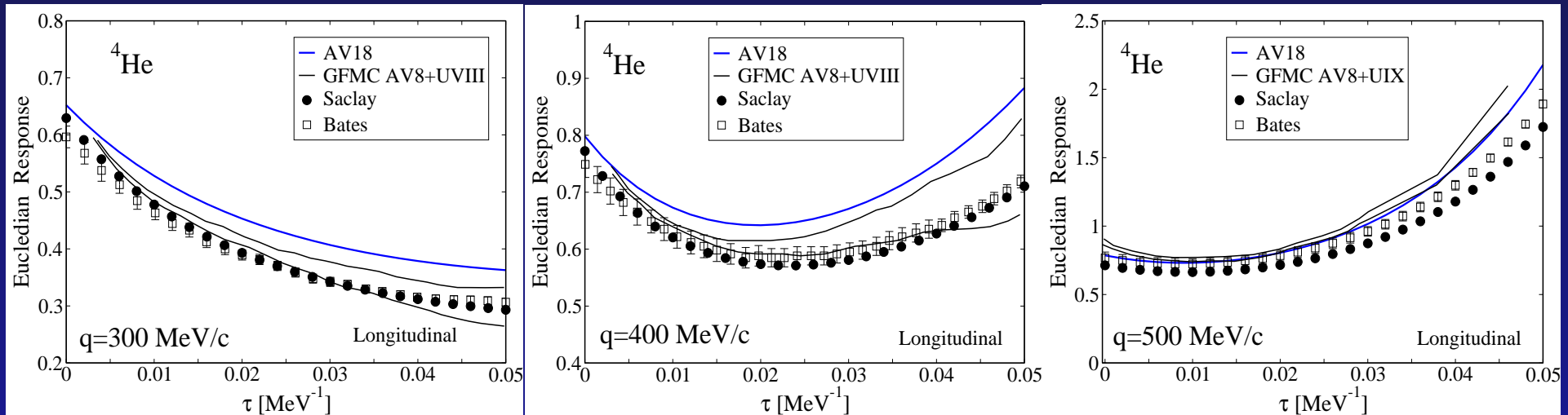
- The quasi-elastic region $\omega \simeq \tau$ is a playground to study nuclear structure
- Technical difficulty of providing realistic *ground state* wave functions
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J. Carlson and R. Schiavilla, PRC **49**, R2880 (1994)

Longitudinal Response Function

Comparison with GFMC



GFMC from J. Carlson and R. Schiavilla, PRL **68**, 3682 (1992)

J. Carlson, J. Jourdan, R. Schiavilla and I. Sick, PRC **65**, 024002 (2002).

Estimate effect of 3NF \Rightarrow If the AV8 and AV18 lead to similar result, we expect 3NF to produce:

- “global” decrease of the total strength, with larger effect at low energy for $q = 300$ and 400 MeV/c
- only at tiny effect $q = 500$ MeV/c.

Transverse Response Function

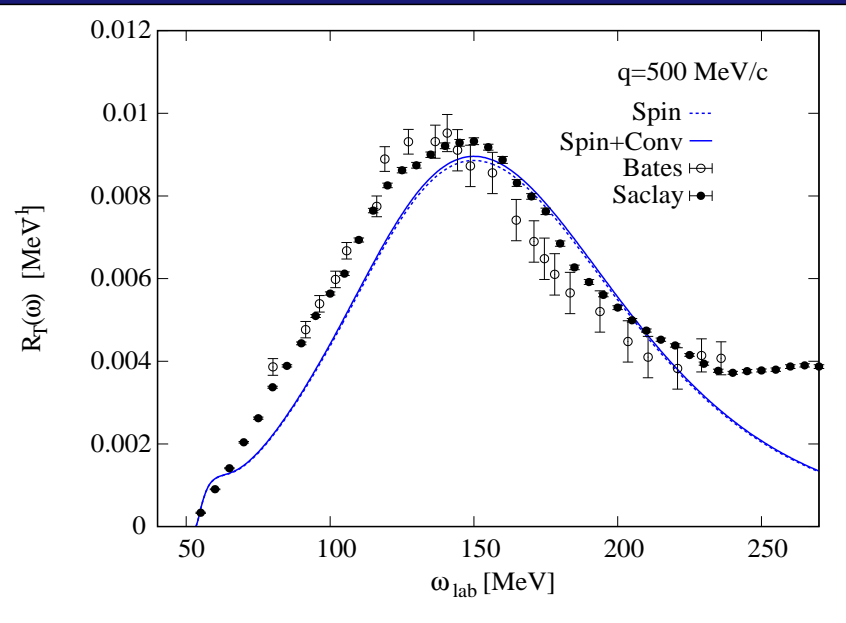
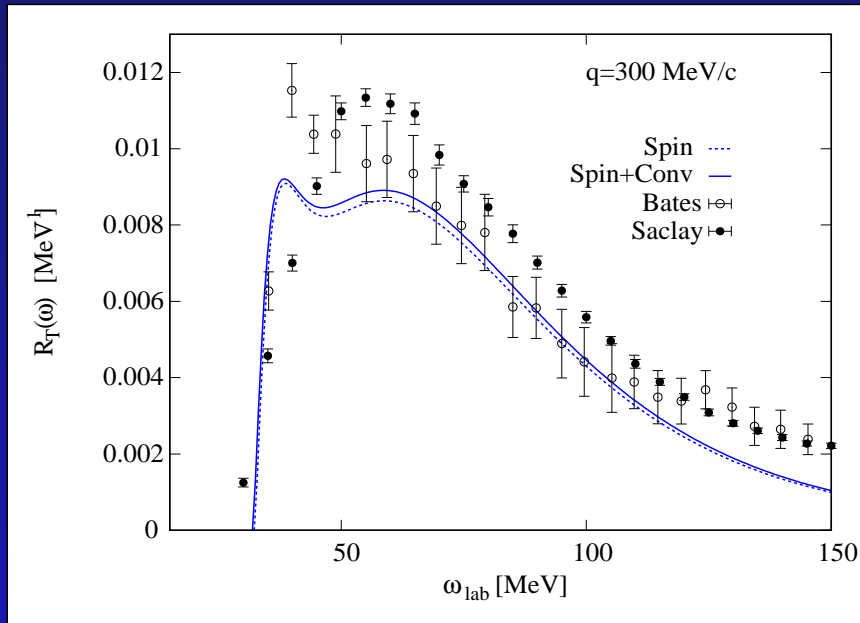
One-body current

Spin Current

$$\mathbf{J}_1^s \sim \frac{i}{2m} \sum_k (\boldsymbol{\sigma}_k \times \mathbf{q}) e^{i\mathbf{q} \cdot \mathbf{r}_k}$$

Convection Current

$$\mathbf{J}_1^c \sim \frac{1}{2m} \sum_k \{\mathbf{p}_k, e^{i\mathbf{q} \cdot \mathbf{r}_k}\}$$



Transverse Response Function

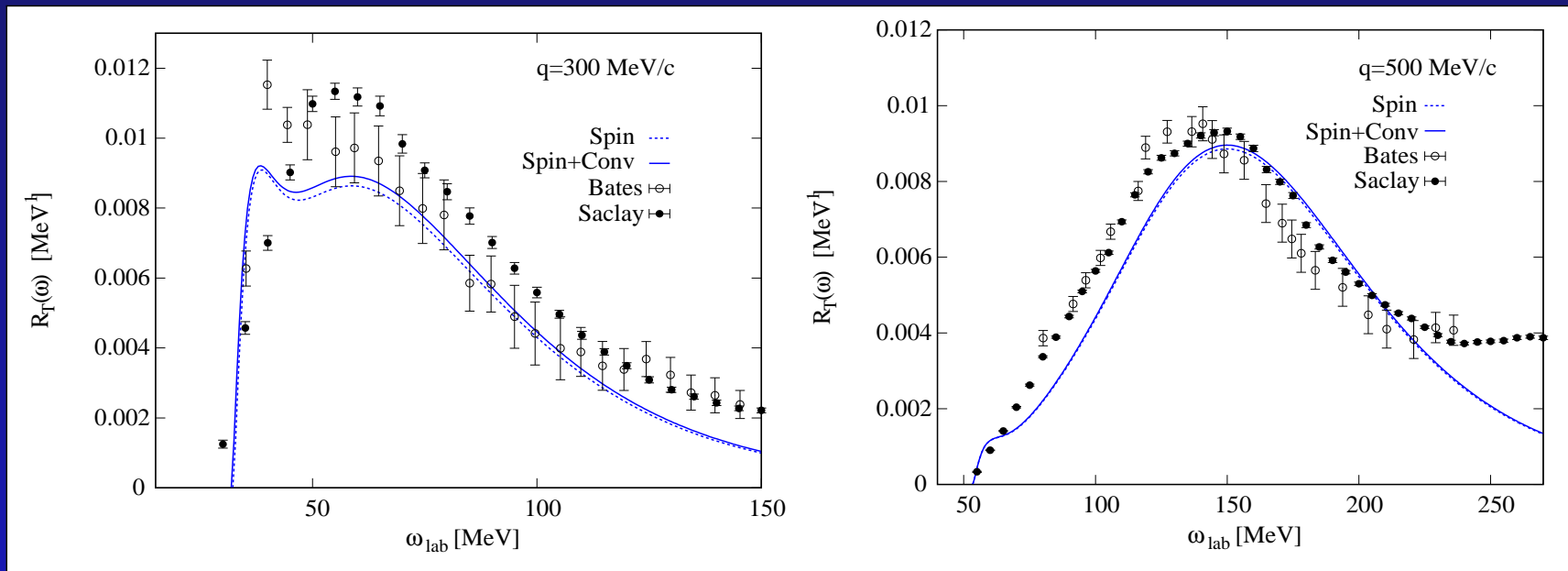
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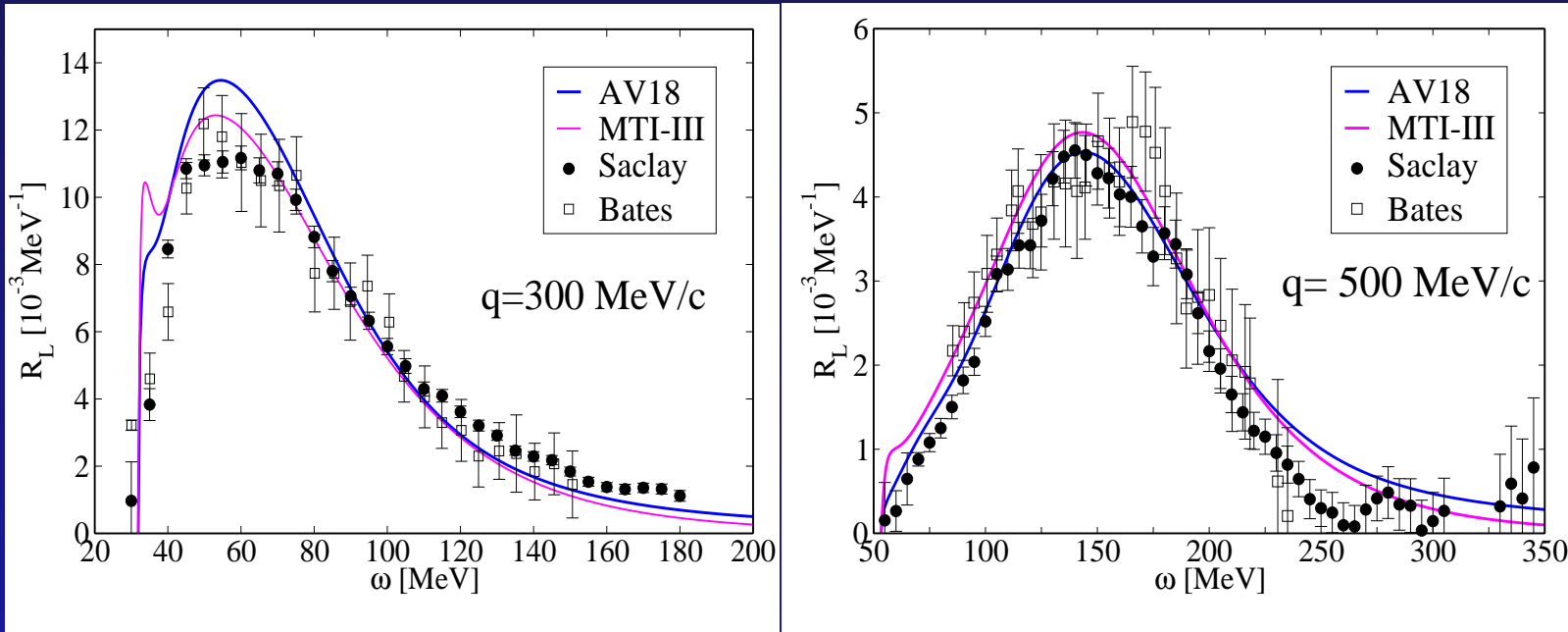
$$\mathbf{J}_1^C \sim \frac{1}{2m} \sum_k \{\mathbf{p}_k, e^{i\mathbf{q} \cdot \mathbf{r}_k}\}$$



Is the missing strength due to the **two-body current**?

Longitudinal Response Function

Comparison of Realistic/Semirealistic Interaction



MTI-III from S.B., H.Arenhövel, N. Barnea, W. Leidemann, G. Orlandini, PRC **76**, 014003 (2007)

- In the q.e. peak with the AV18 we find larger strength for $q = 300 \text{ MeV}/c$ and less for $q = 500 \text{ MeV}/c$ with respect to the MTI-III.
- C_1^V behaves like for photoabsorption, but in ee' other multipoles are involved
- R_L at $q = 500 \text{ MeV}/c$ may not yet be fully converged in J .