

No-core shell model with JISP16 NN interaction: spectroscopy of light nuclei and neutron-nucleus scattering

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Highlights

JISP = J -matrix inverse scattering potential

PETs = phase-equivalent transformations

No-core shell model: *ab initio* \Leftrightarrow *ab exitu*

approach \Rightarrow JISP16 NN interaction

No three-nucleon forces

Spectroscopy of nuclei with $A \leq 16$

N -nucleus scattering

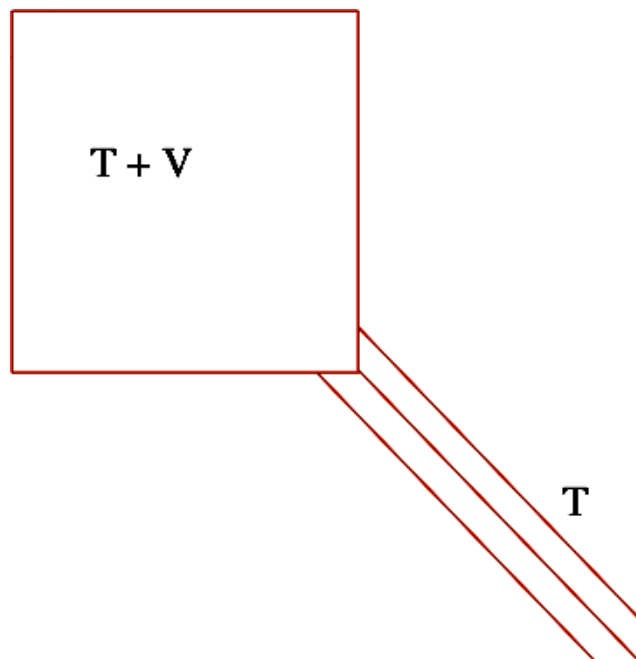
JISP type NN interaction:

- Description of NN data
- Small matrix of the NN interaction in the oscillator basis:

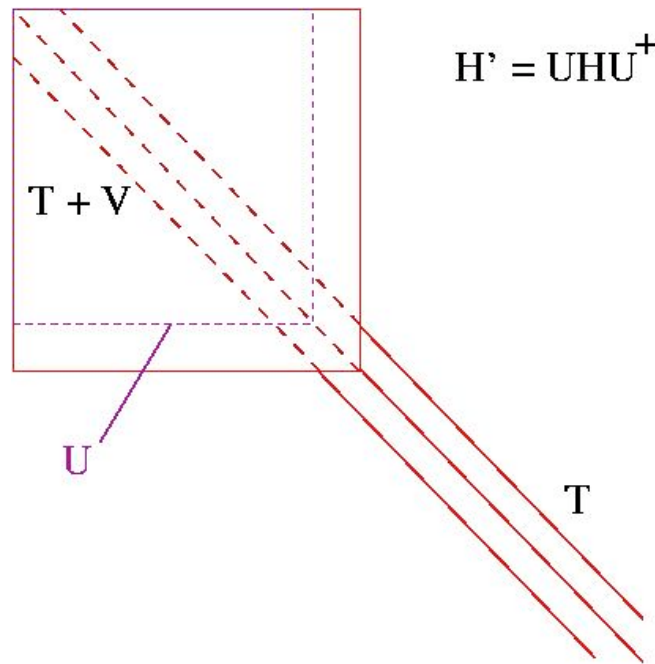
$$V = \sum_{\Gamma, \Gamma'} \sum_{n=0}^{N_{\Gamma}} \sum_{n'=0}^{N_{\Gamma'}} |n\Gamma\rangle V_{nn'}^{\Gamma\Gamma'} \langle n'\Gamma'|$$

- Description of (light) nuclei

J -matrix formalism: scattering in the oscillator basis



PETs



ab initio \Leftrightarrow *ab exitu*

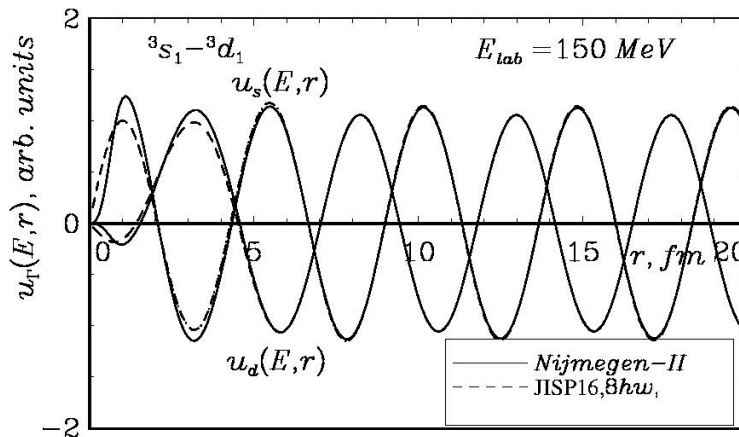
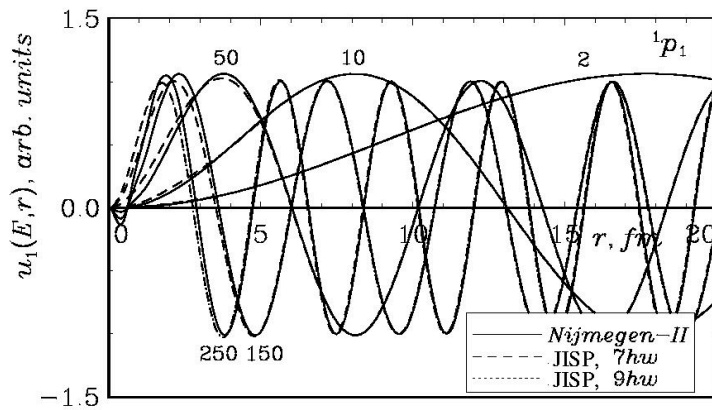
- JISP16: J -matrix inverse scattering $9h\Omega$ NN potential with $h\Omega = 40$ MeV fitted to nuclei up through ^{16}O
- Only simplest PETs generated by 2×2 unitary matrix U are used
- *Ab exitu* approach:
- PETs: sd wave - fitting deuteron properties (rms radius and quadrupole moment)
various p and one of d waves - fitting few levels of ^6Li and binding energy of ^{16}O in relatively small model spaces
- All the rest NCSM results (other nuclei, larger model spaces) are *ab initio*

JISP16 properties

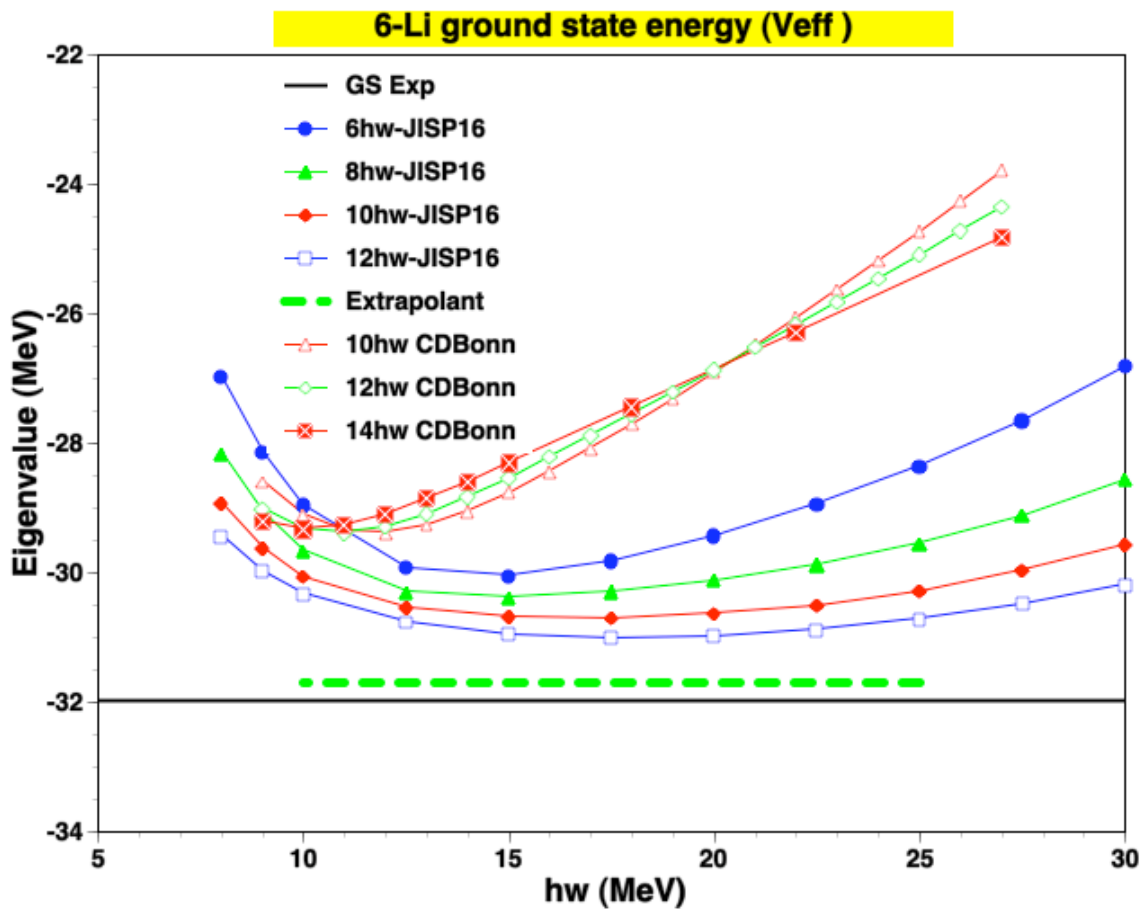
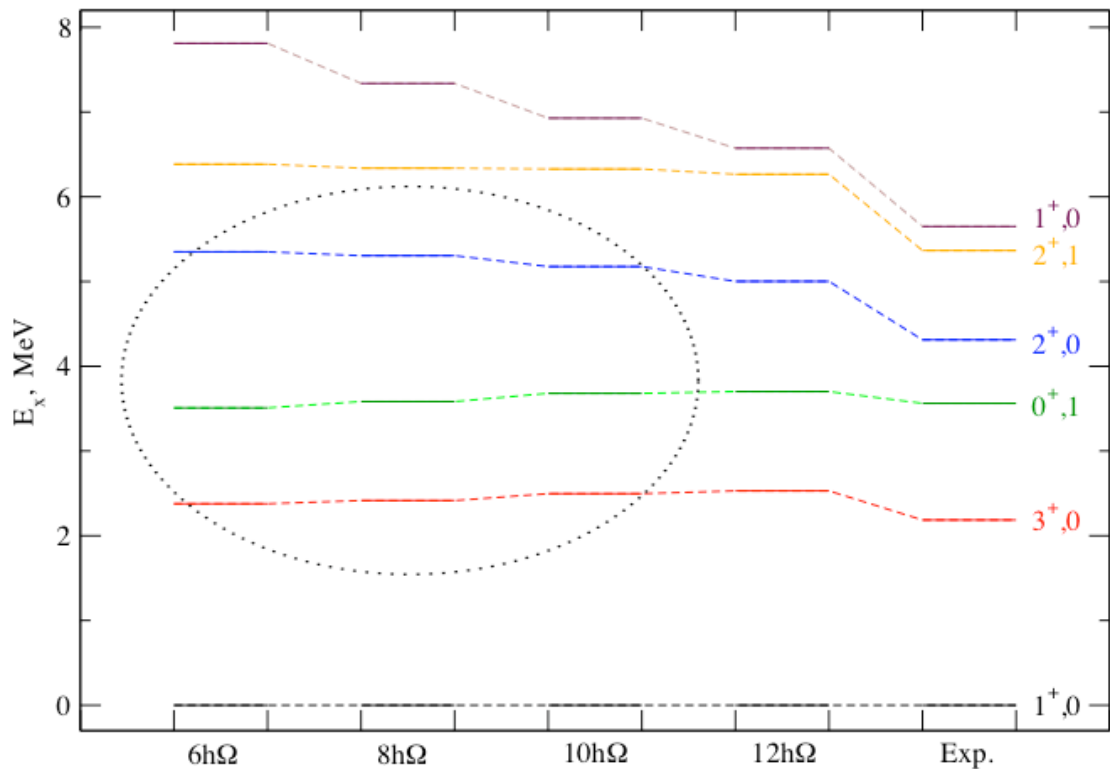
- 1992 np data base (2514 data): $\chi^2/\text{datum} = 1.03$
- 1999 np data base (3058 data): $\chi^2/\text{datum} = 1.05$

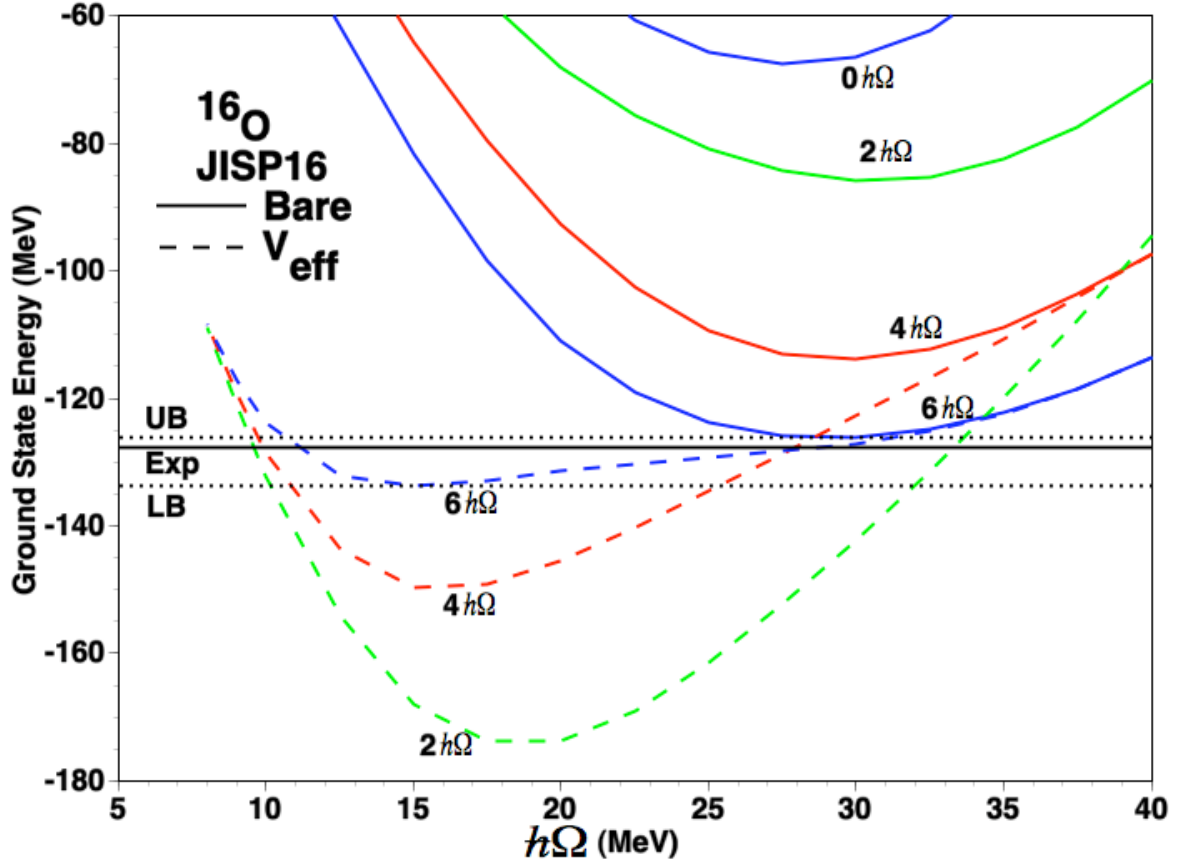
Table I: Deuteron properties.

Potential	E_d , MeV	d state	rms radius,	Q , fm ²	As. norm. const.	$\eta = \frac{\mathcal{A}_d}{\mathcal{A}_s}$
		probability, %	fm		\mathcal{A}_s , fm ^{-1/2}	
JISP16	-2.224575	4.1360	1.9643	0.2886	0.8629	0.0252
Nijmegen-II	-2.224575	5.635	1.968	0.2707	0.8845	0.0252
AV18	-2.224575	5.76	1.967	0.270	0.8850	0.0250
CD-Bonn	-2.224575	4.85	1.966	0.270	0.8846	0.0256
Nature	-2.224575(9)	—	1.971(6)	0.2859(3)	0.8846(9)	0.0256(4)



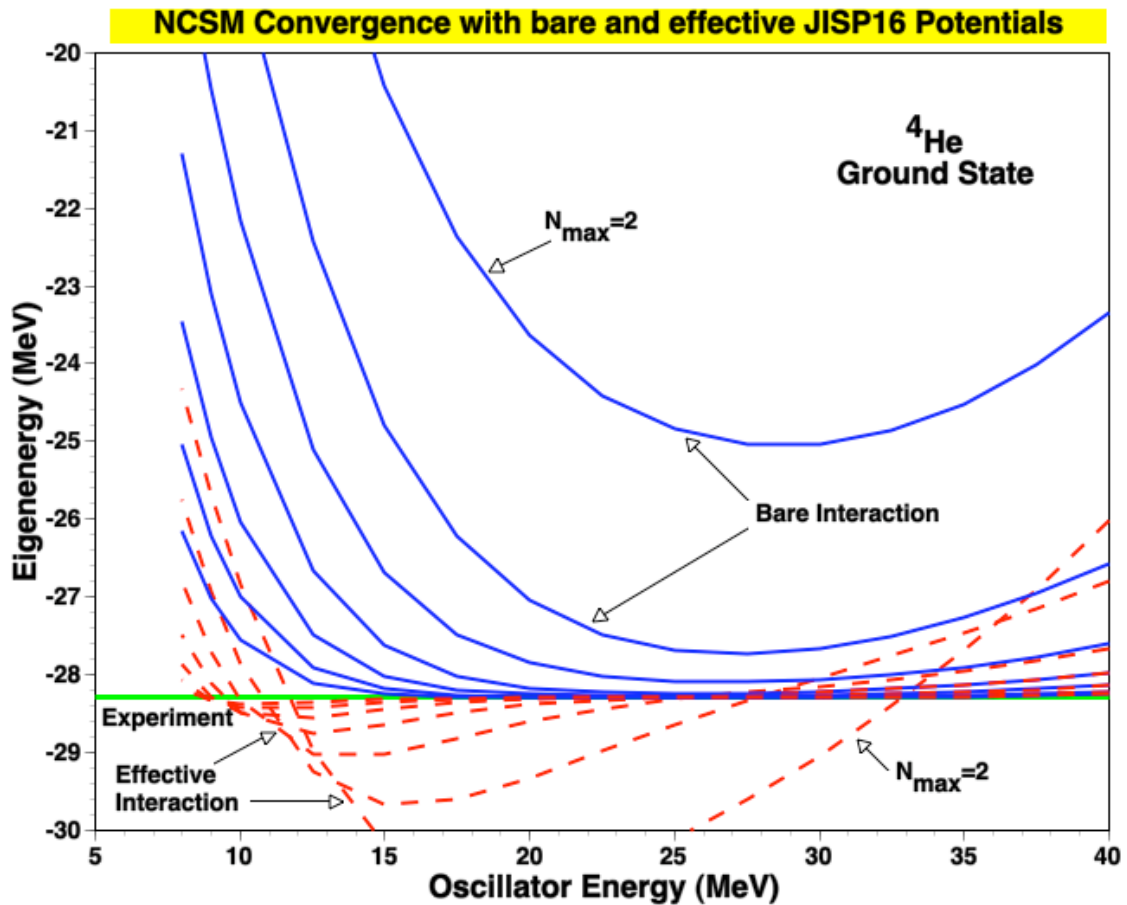
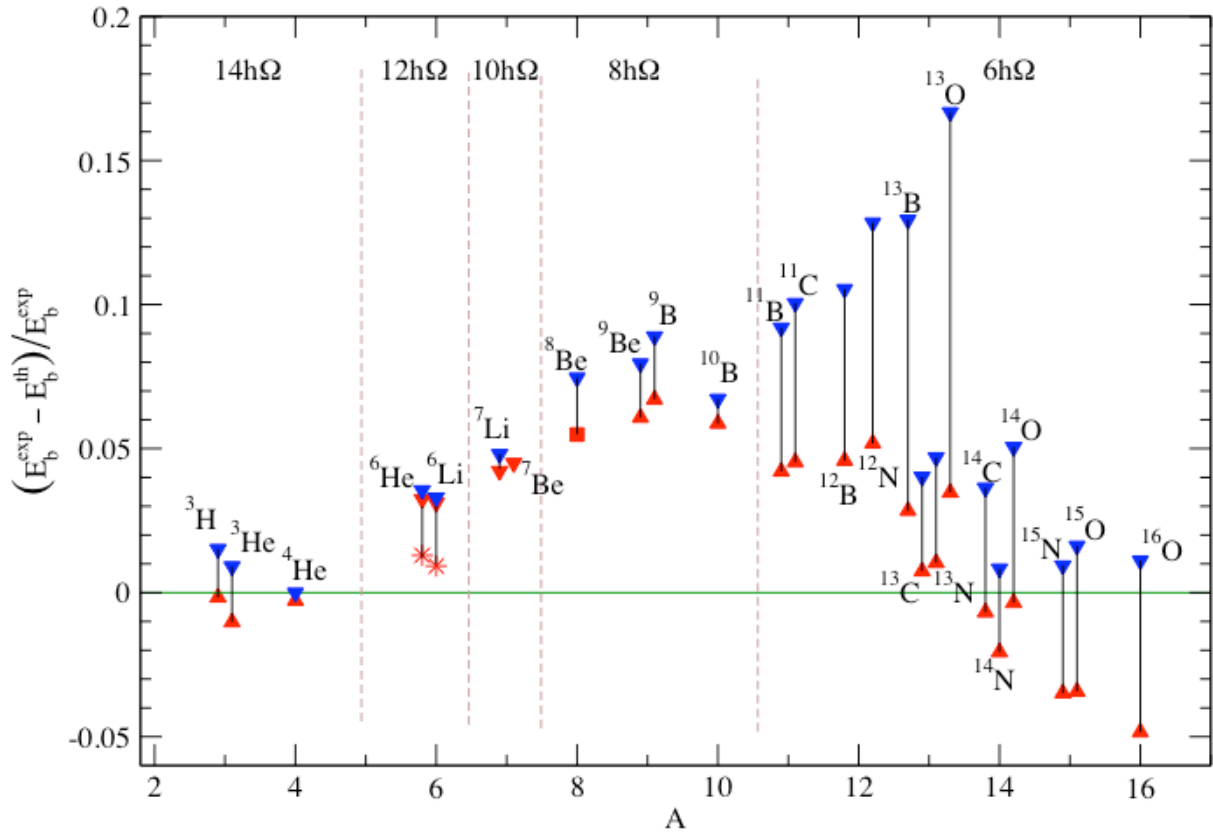
${}^6\text{Li}$ spectrum with JISP16 NN interaction, $\hbar\Omega=17.5$ MeV

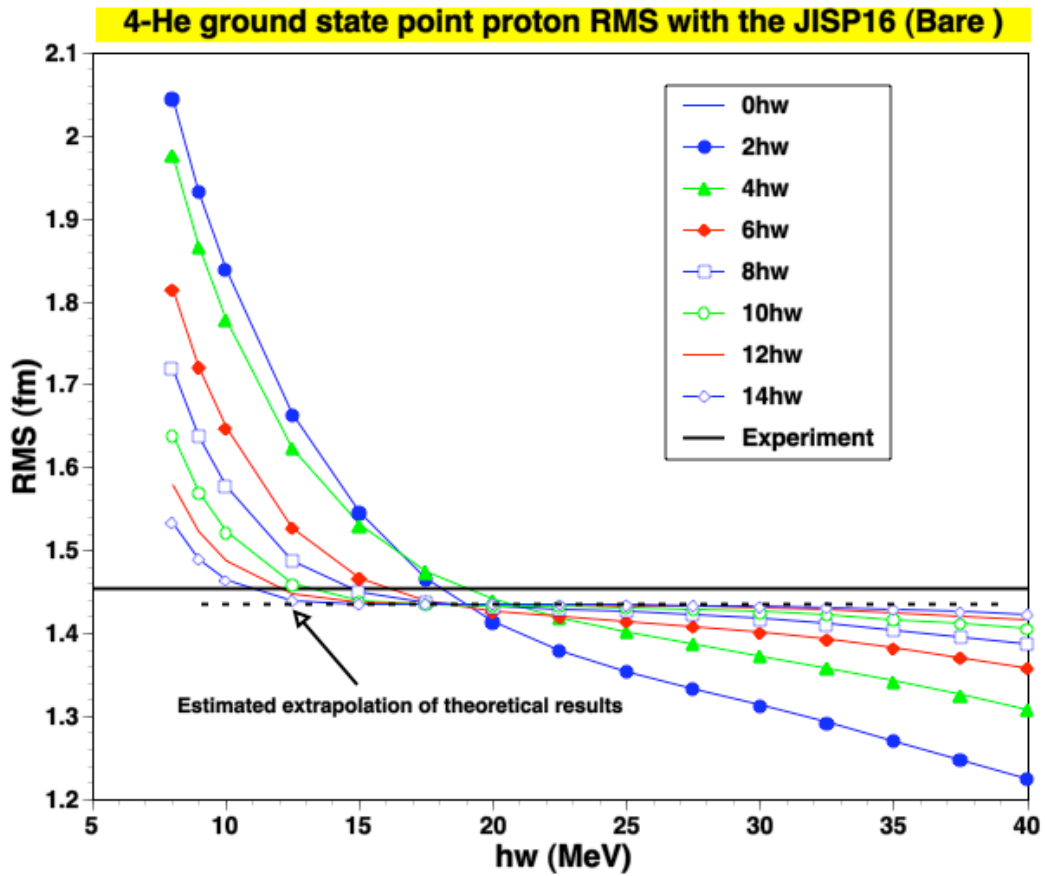
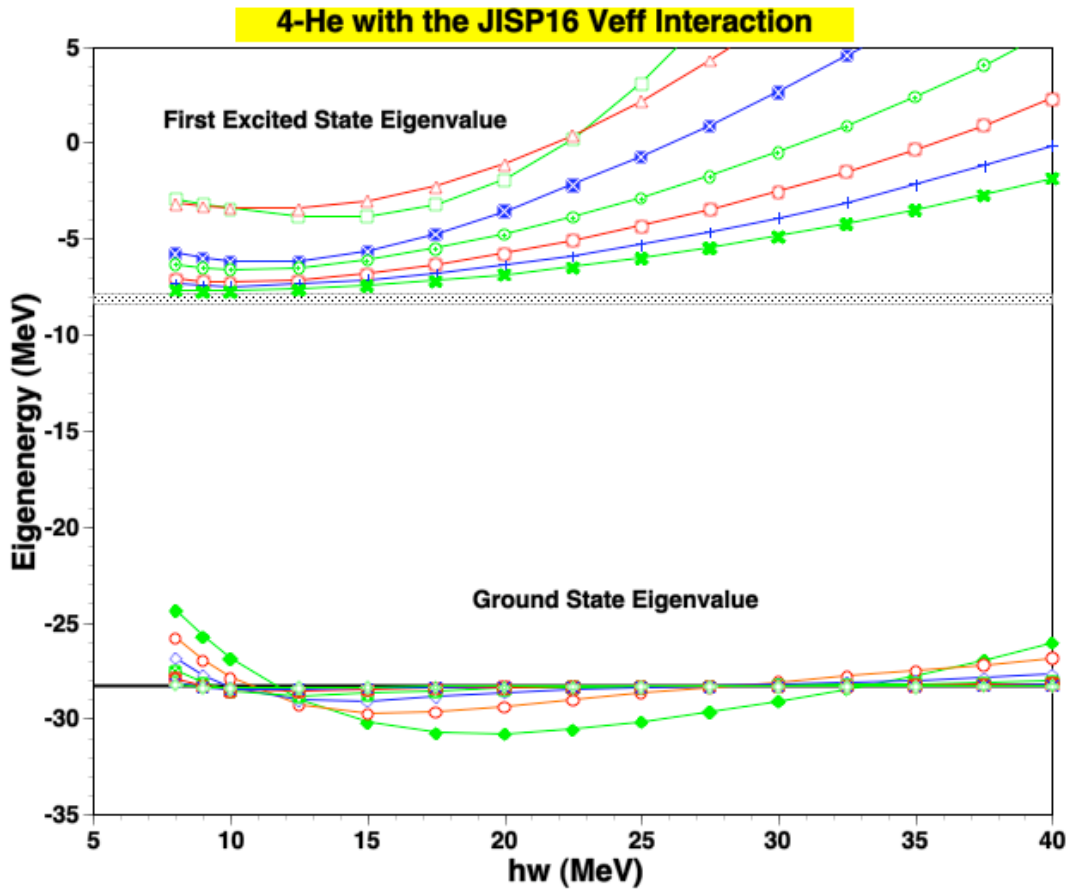


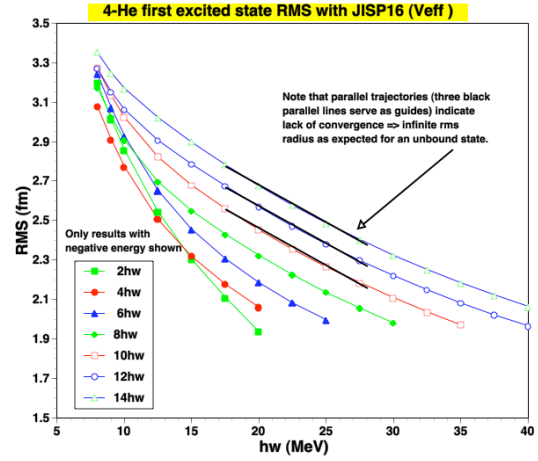
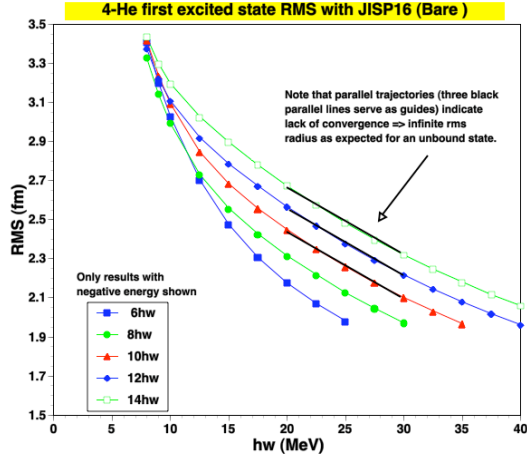
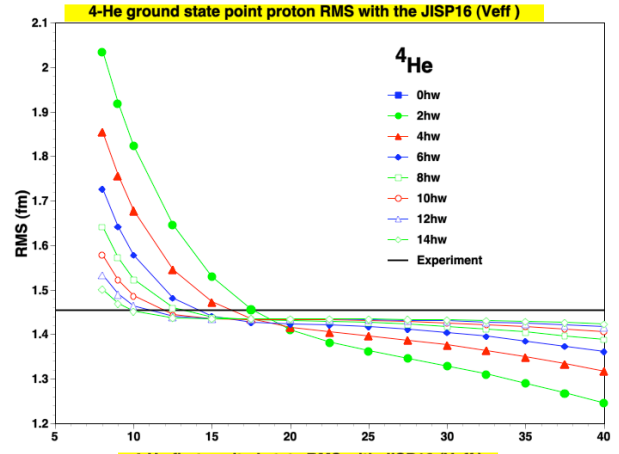
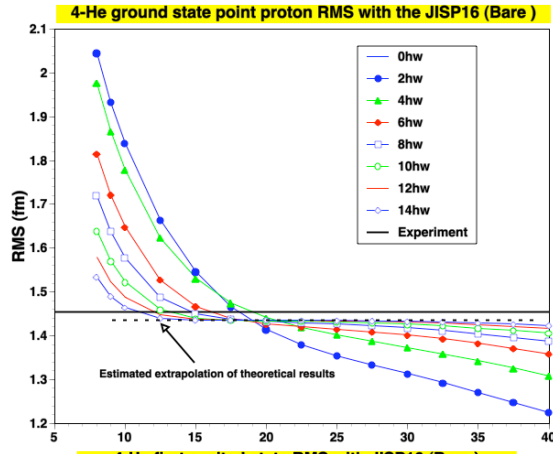


Nucleus	Nature	Bare	Effective	$\hbar\omega$ (MeV)	Model space	Nucleus	Nature	Bare	Effective	$\hbar\omega$ (MeV)	Model space
³ H	8.482	8.354	8.496(20)	7	14 $\hbar\omega$	¹¹ C	73.440	66.1	70.1(32)	17	6 $\hbar\omega$
³ He	7.718	7.648	7.797(17)	7	14 $\hbar\omega$	¹² B	79.575	71.2	75.9(48)	15	6 $\hbar\omega$
⁴ He	28.296	28.297	28.374(57)	10	14 $\hbar\omega$	¹² C	92.162	87.4	91.0(49)	17.5	6 $\hbar\omega$
⁶ He	29.269		28.32(28)	17.5	12 $\hbar\omega$	¹² N	74.041	64.5	70.2(48)	15	6 $\hbar\omega$
⁶ Li	31.995		31.00(31)	17.5	12 $\hbar\omega$	¹³ B	84.453	73.5	82.1(67)	15	6 $\hbar\omega$
⁷ Li	39.245		37.59(30)	17.5	10 $\hbar\omega$	¹³ C	97.108	93.2	96.4(59)	19	6 $\hbar\omega$
⁷ Be	37.600		35.91(29)	17	10 $\hbar\omega$	¹³ N	94.105	89.7	93.1(62)	18	6 $\hbar\omega$
⁸ Be	56.500		53.40(10)	15	8 $\hbar\omega$	¹³ O	75.558	63.0	72.9(62)	14	6 $\hbar\omega$
⁹ Be	58.165	53.54	54.63(26)	16	8 $\hbar\omega$	¹⁴ C	105.285	101.5	106.0(93)	17.5	6 $\hbar\omega$
⁹ B	56.314	51.31	52.53(20)	16	8 $\hbar\omega$	¹⁴ N	104.659	103.8	106.8(77)	20	6 $\hbar\omega$
¹⁰ Be	64.977	60.55	61.39(20)	19	8 $\hbar\omega$	¹⁴ O	98.733	93.7	99.1(92)	16	6 $\hbar\omega$
¹⁰ B	64.751	60.39	60.95(20)	20	8 $\hbar\omega$	¹⁵ N	115.492	114.4	119.5(126)	16	6 $\hbar\omega$
¹⁰ C	60.321	55.26	56.36(67)	17	8 $\hbar\omega$	¹⁵ O	111.956	110.1	115.8(126)	16	6 $\hbar\omega$
¹¹ B	76.205	69.2	73.0(31)	17	6 $\hbar\omega$	¹⁶ O	127.619	126.2	133.8(158)	15	6 $\hbar\omega$

Binding energies



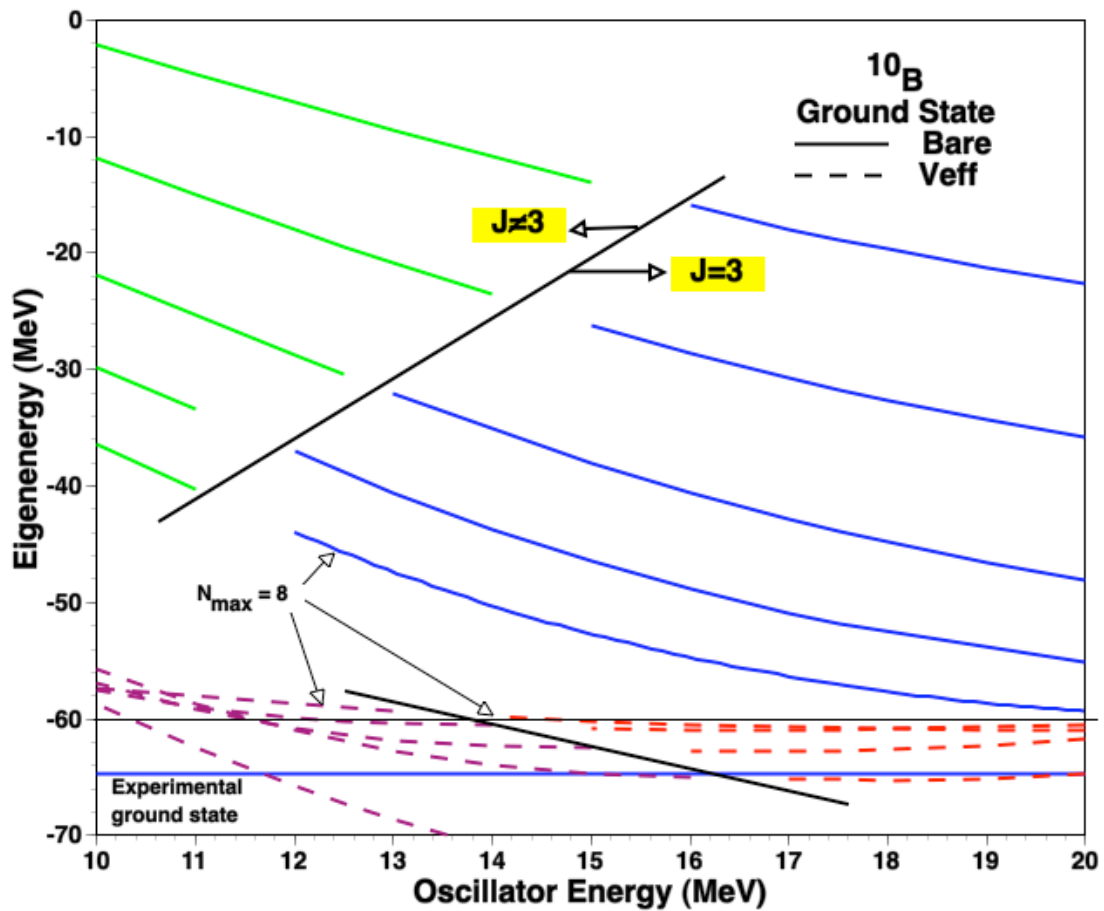
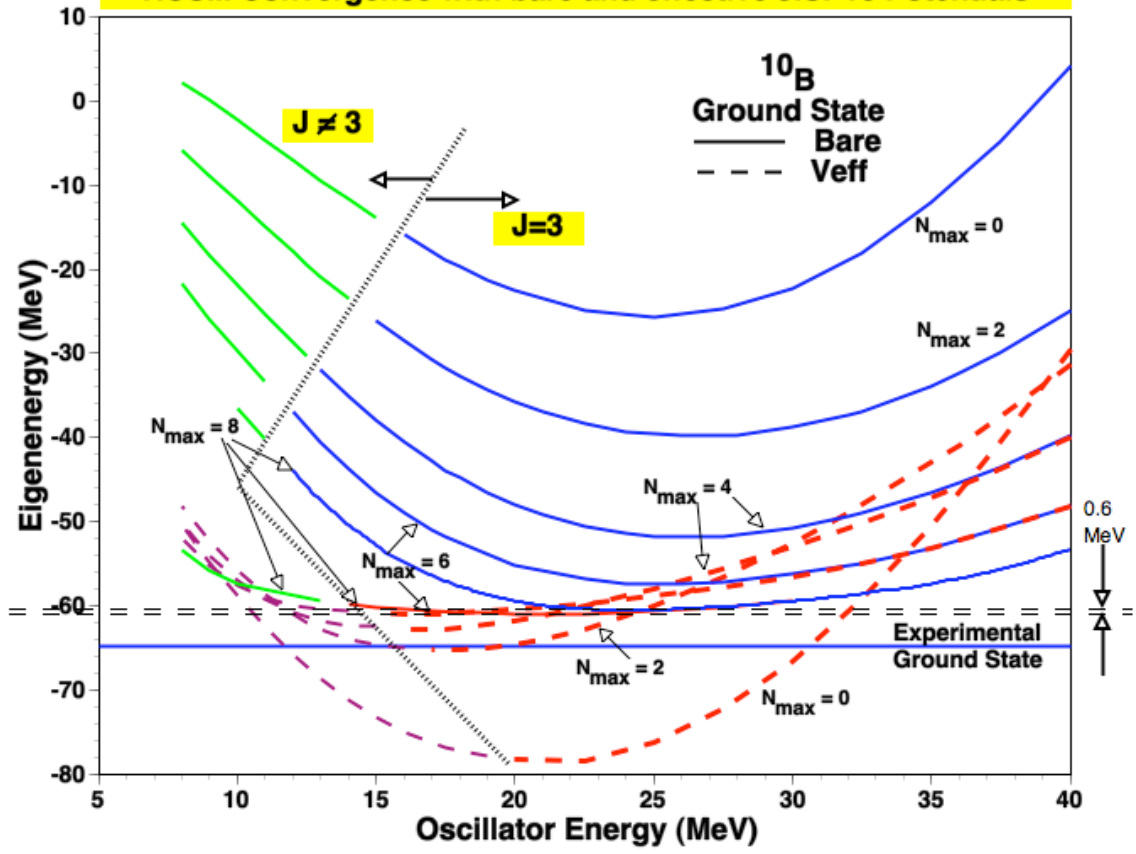




Ground state energy E_{gs} and excitation energies E_x (in MeV), ground state point-proton rms radius r_p (in fm) and quadrupole moment Q (in $e \cdot \text{fm}^2$) of the ${}^6\text{Li}$ nucleus; $\hbar\omega = 17.5$ MeV.

Interaction	Nature	JISP6	JISP16	AV8'+TM'	AV18+UIX	AV18+IL2
Method		NCSM, $10\hbar\omega$ [6]	NCSM, $12\hbar\omega$	NCSM, $6\hbar\omega$ [2]	GFMC [8,15]	GFMC [10,15]
$E_{gs}(1_1^+, 0)$	-31.995	-31.48	-31.00	-31.04	-31.25(8)	-32.0(1)
r_p	2.32(3)	2.083	2.151	2.054	2.46(2)	2.39(1)
Q	-0.082(2)	-0.194	-0.0646	-0.025	-0.33(18)	-0.32(6)
$E_x(3^+, 0)$	2.186	2.102	2.529	2.471	2.8(1)	2.2
$E_x(0^+, 1)$	3.563	3.348	3.701	3.886	3.94(23)	3.4
$E_x(2^+, 0)$	4.312	4.642	5.001	5.010	4.0(1)	4.2
$E_x(2^+, 1)$	5.366	5.820	6.266	6.482		5.5
$E_x(1_2^+, 0)$	5.65	6.86	6.573	7.621	5.1(1)	5.6

NCSM Convergence with bare and effective JISP16 Potentials



¹⁰B

Potential Approach	Nature	JISP16 NCSM, $8\hbar\omega^a$	AV8'+TM' NCSM, $4\hbar\omega^b$	AV18+IL2 GFMC ^c	ChPT NCSM, $6\hbar\omega^d$
$E_{gs}(3_1^+, 0)$	-64.751	-60.14	-60.57	-65.6(5)	-64.78
r_p	2.30(12)	2.168	2.168	2.33(1)	2.197
Q	+8.472(56)	6.484	+5.682	+9.5(2)	+6.327
$E_x(1_1^+, 0)$	0.718	0.555	0.340	0.9	0.523
$E_x(0^+, 1)$	1.740	1.202	1.259		1.279
$E_x(1_2^+, 0)$	2.154	2.379	1.216		1.432
$E_x(2_1^+, 0)$	3.587	3.721	2.775	3.9	3.178
$E_x(3_2^+, 0)$	4.774	6.162	5.971		6.729
$E_x(2_1^+, 1)$	5.164	5.049	5.182		5.315
$E_x(2_2^+, 0)$	5.92	5.548	3.987		4.835
$E_x(4^+, 0)$	6.025	5.775	5.229	5.6	5.960
$E_x(2_2^+, 1)$	7.478	7.776	7.491		7.823
$B(E2; 1_1^+0 \rightarrow 3_1^+0)$	4.13(6)	3.317	1.959		3.05
$B(E2; 1_2^+0 \rightarrow 3_1^+0)$	1.71(26)	0.627	1.010		0.50
$B(GT; 3_1^+0 \rightarrow 2_1^+1)$	0.083(3)	0.042	0.066		0.07
$B(GT; 3_1^+0 \rightarrow 2_2^+1)$	0.95(13)	1.652	1.291		1.22

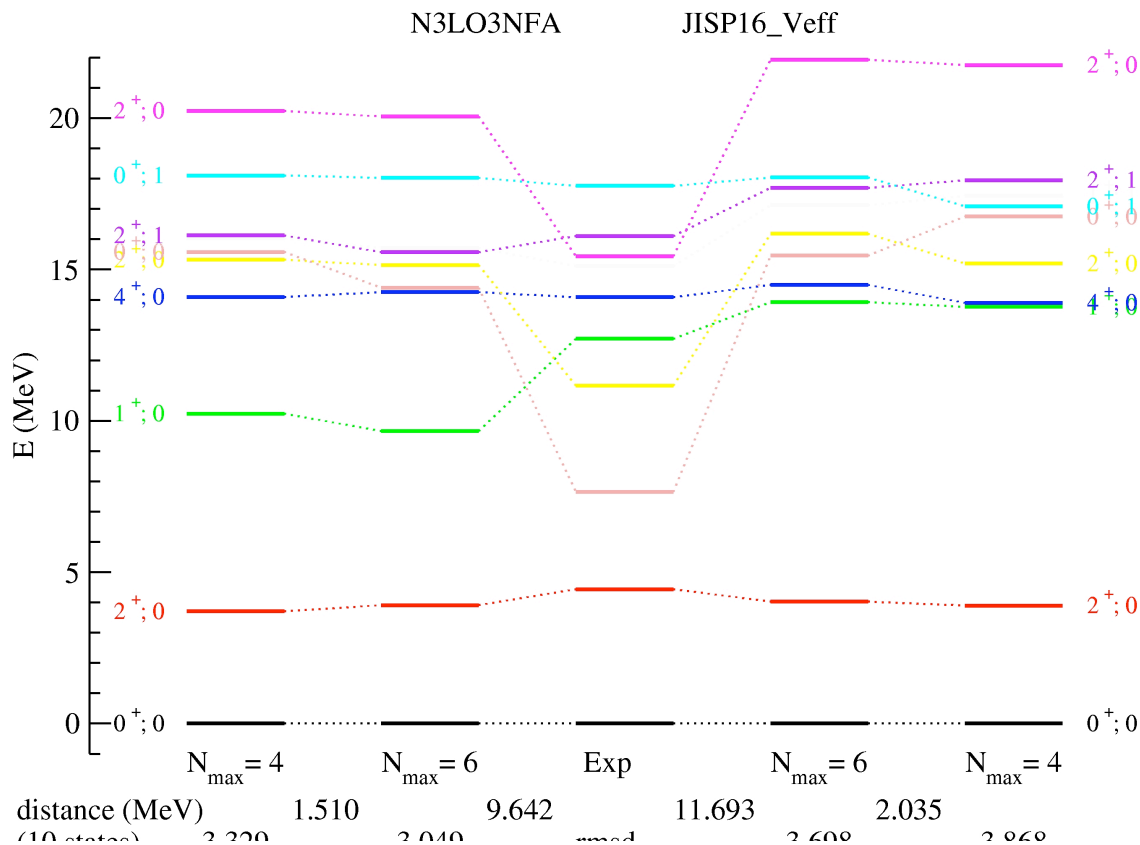
^aA.M.Shirokov, J.P.Vary, A.I.Mazur, T.A.Weber, Phys. Lett. **B644**, 33 (2007).

^bP. Navrátil, W. E. Ormand, Phys. Rev. C **68**, 034305 (2003).

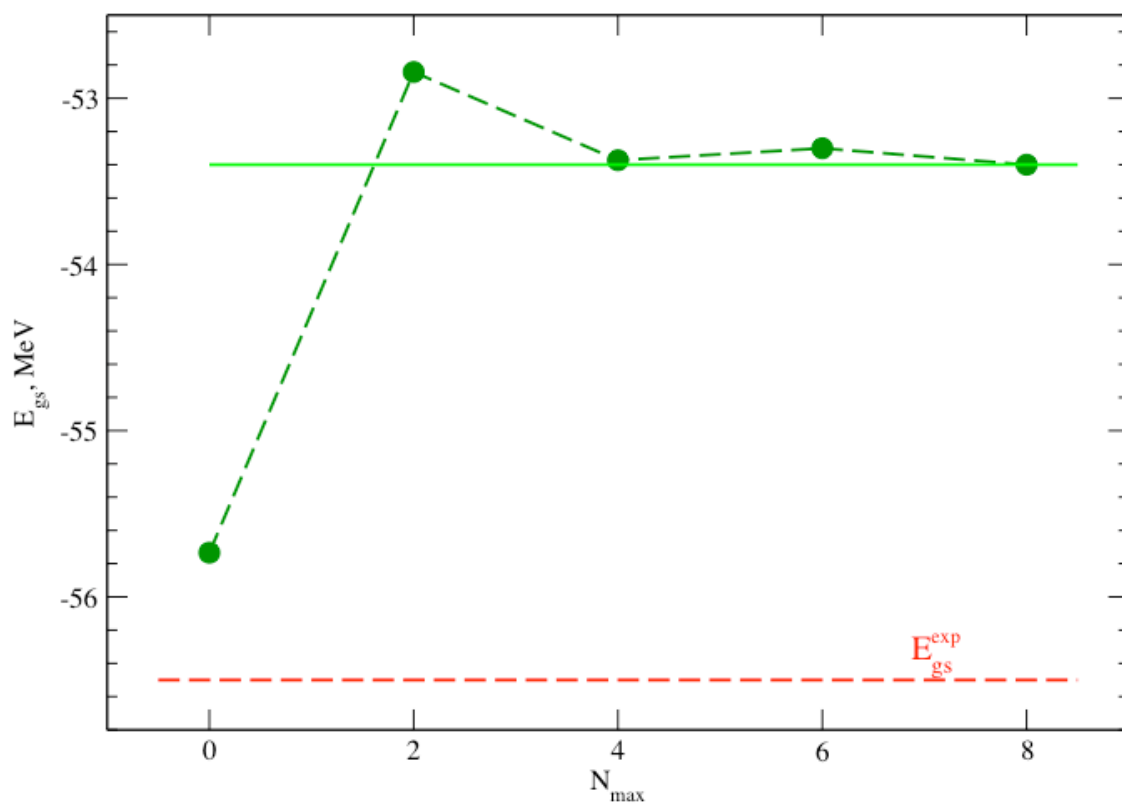
^cS. C. Pieper, K. Varga, R. B. Wiringa, Phys. Rev. C **66**, 044310 (2002).

^dP. Navrátil, V. G. Gueorguiev, J. P. Vary, W. E. Ormand, A. Nogga, Phys. Rev. Lett. **99**, 042501 (2007).

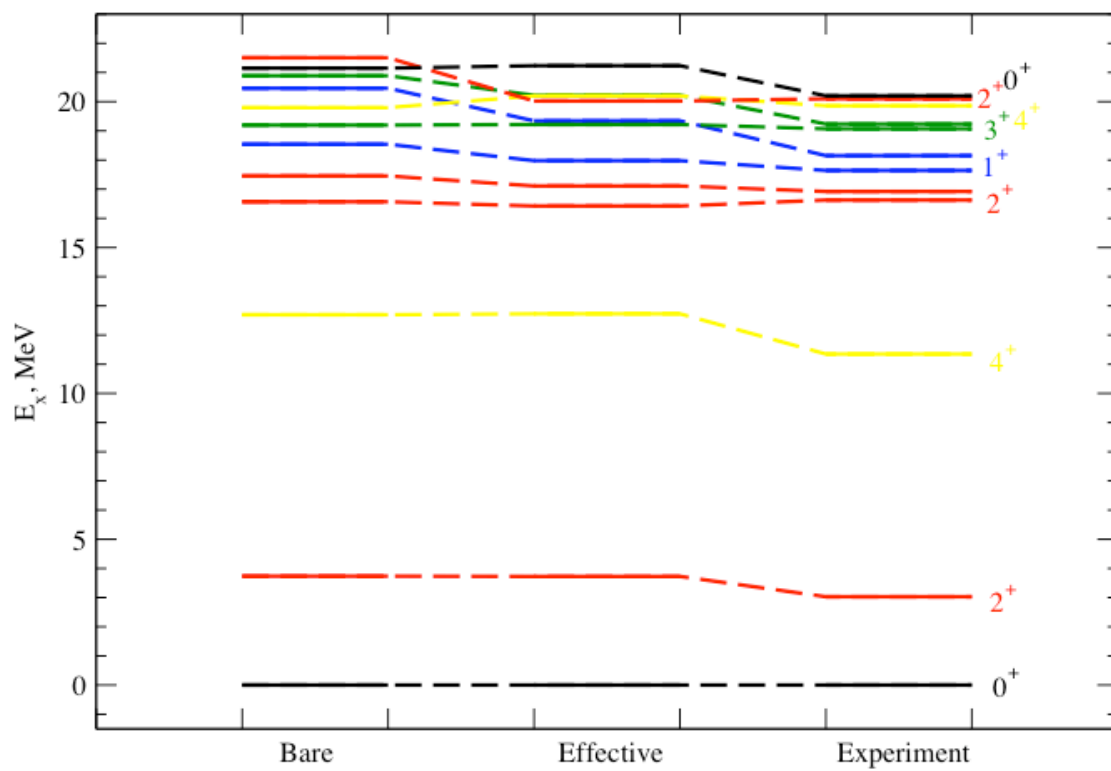
¹²C Spectral Convergence for $\hbar\Omega = 15$ MeV



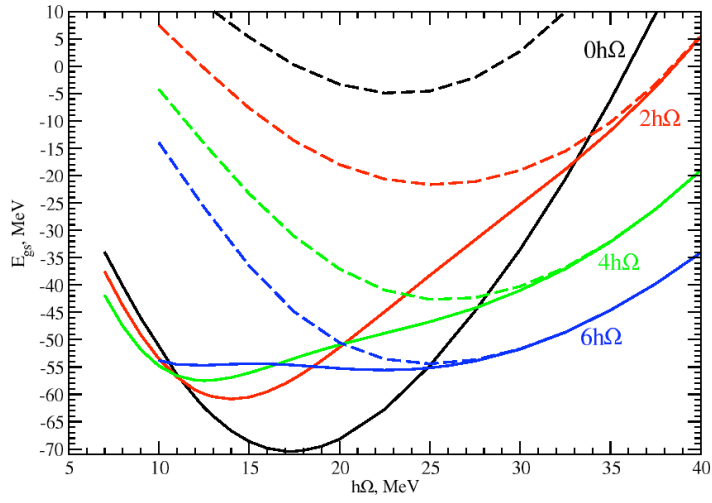
^8Be g.s. convergence with N_{max} $h\Omega$
 $h\Omega = 15 \text{ MeV}$



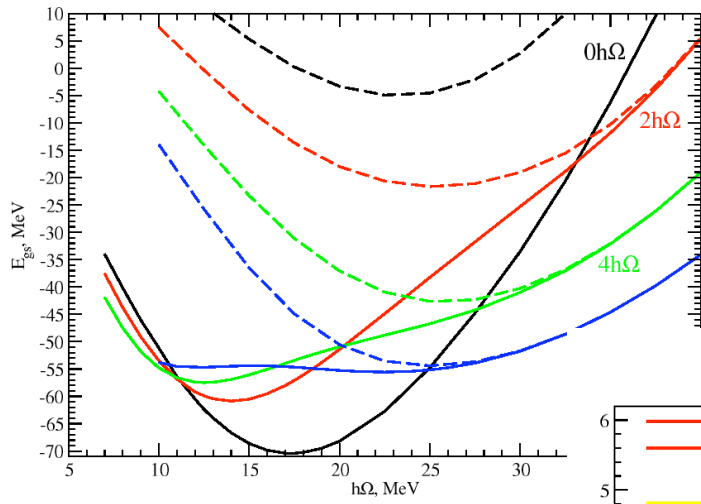
^8Be spectrum
 NCSM, $8h\Omega$ model space



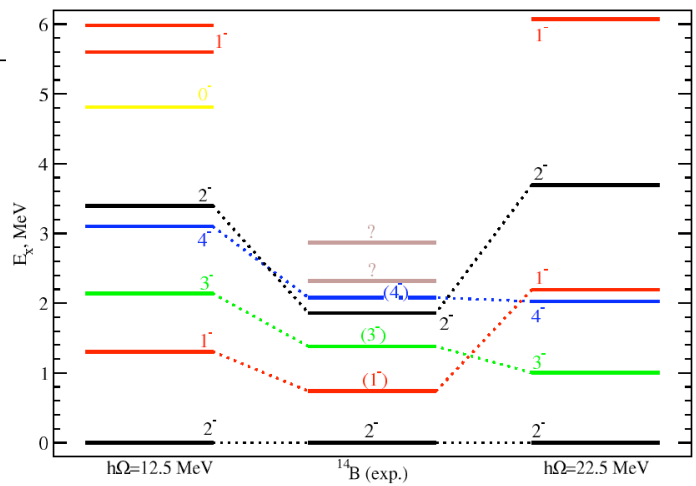
^{14}F ground state



^{14}F ground state



^{14}F spectrum



Role of *NNN* force?

- W. Polyzou and W. Glöckle theorem (Few-body Syst. **9**, 97 (1990)):

$$H=T+V_{ij} \Rightarrow H'=T+V'_{ij}+V_{ijk}$$

where V_{ij} and V'_{ij} are phase-equivalent, H and H' are isospectral.

Hope:

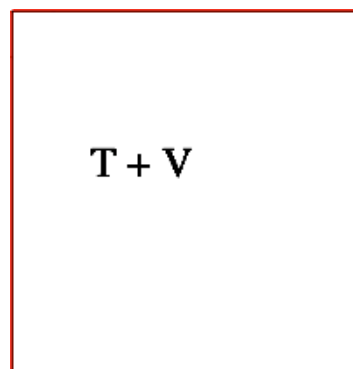
$$H'=T+V'_{ij}+V_{ijk} \Rightarrow H=T+V_{ij}$$

with (approximately) isospectral H and H' .

JISP16 seems to be *NN* interaction minimizing *NNN* force.

Without *NNN* force calculations are simpler, calculations are faster, larger model spaces become available.

J-matrix formalism: scattering in the oscillator basis



$T + V$

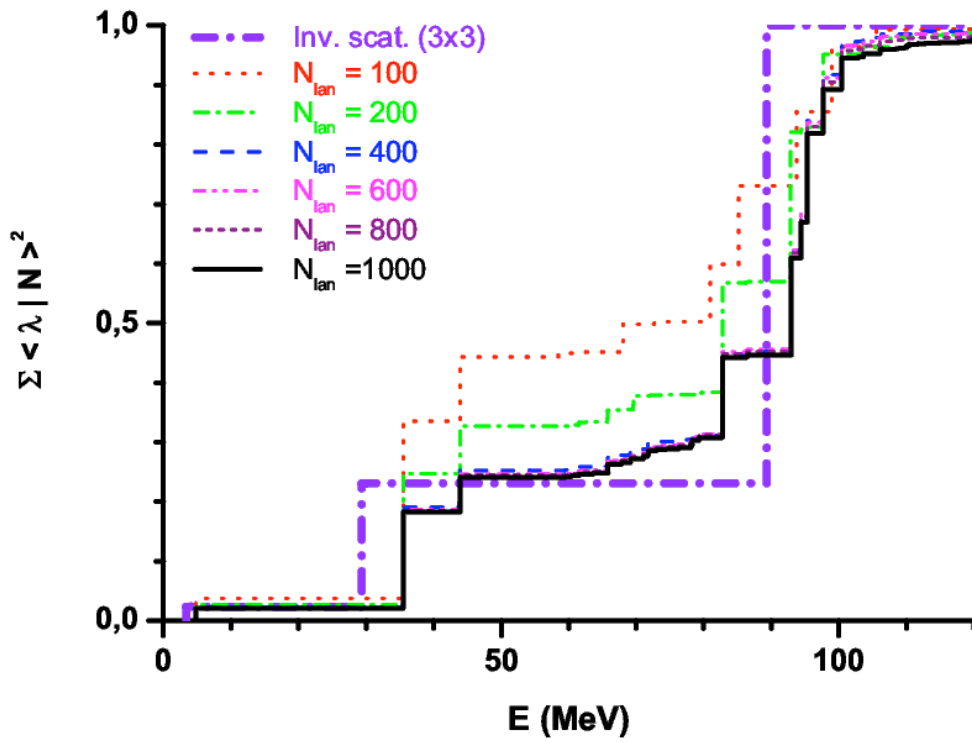
$$\sum_{n'=0}^N H_{nn'}^l \langle n' | \lambda \rangle = E_\lambda \langle n | \lambda \rangle, \quad n \leq N.$$

$$\mathcal{G}_{NN}(E) = - \sum_{\lambda=0}^N \frac{\langle N | \lambda \rangle^2}{E_\lambda - E}$$

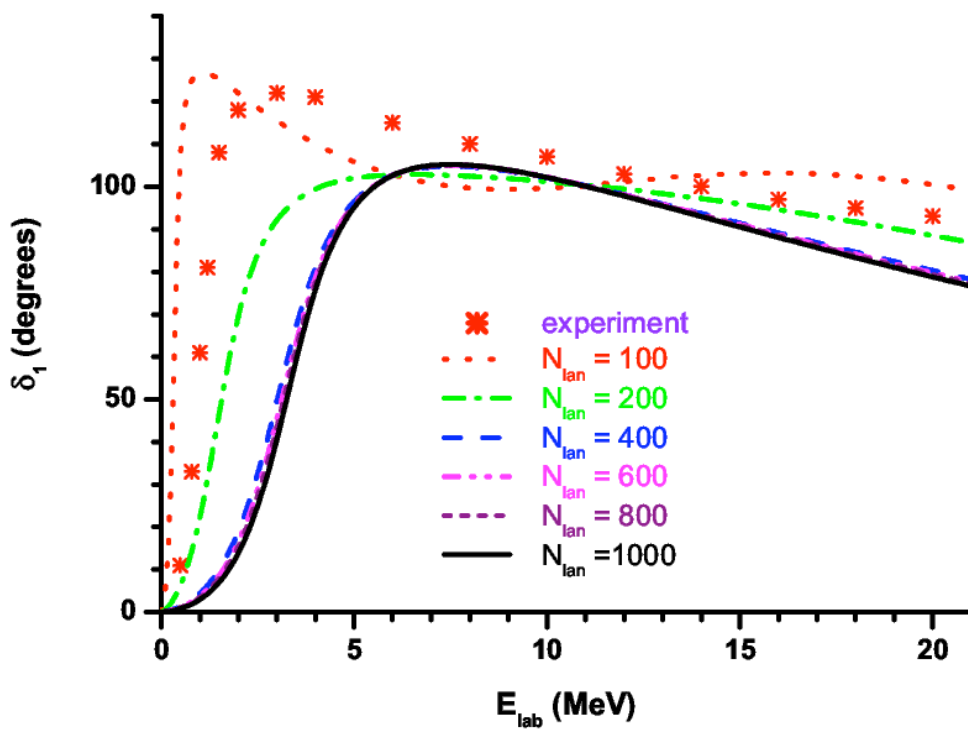
$$S = \frac{C_{Nl}^{(-)}(q) - \mathcal{G}_{NN}(E) T_{N,N+1}^l C_{N+1,l}^{(-)}(q)}{C_{Nl}^{(+)}(q) - \mathcal{G}_{NN}(E) T_{N,N+1}^l C_{N+1,l}^{(+)}(q)},$$

T

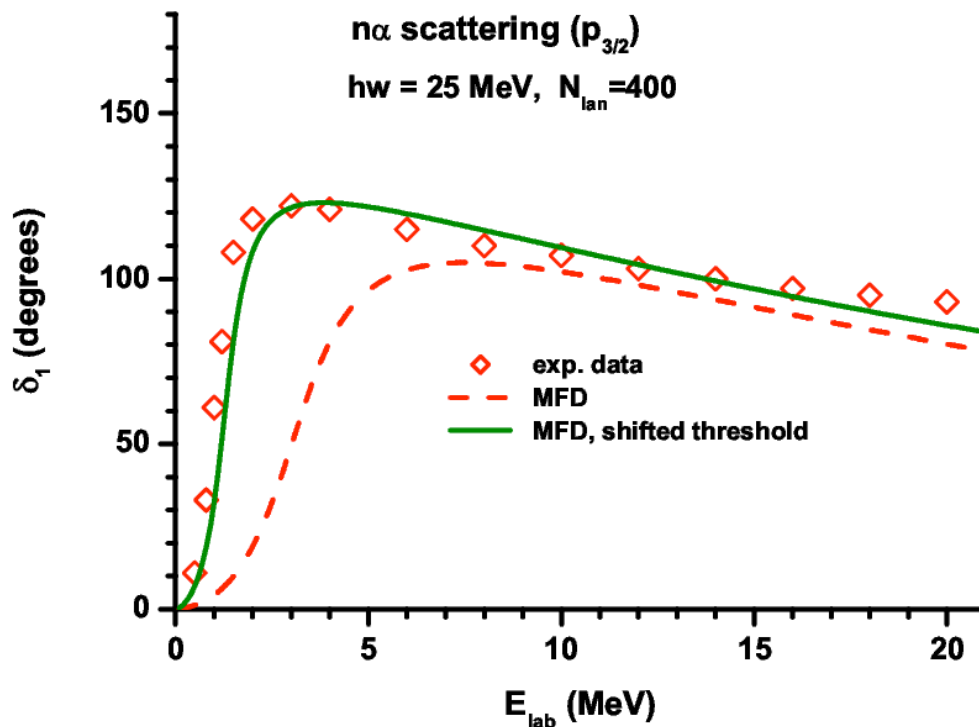
$n\alpha$ scattering



$n\alpha$ scattering

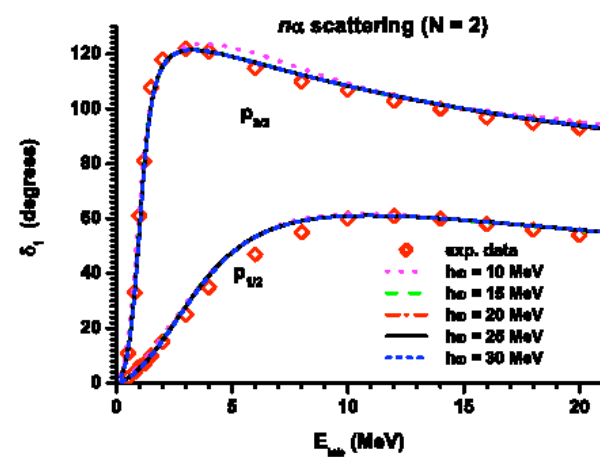
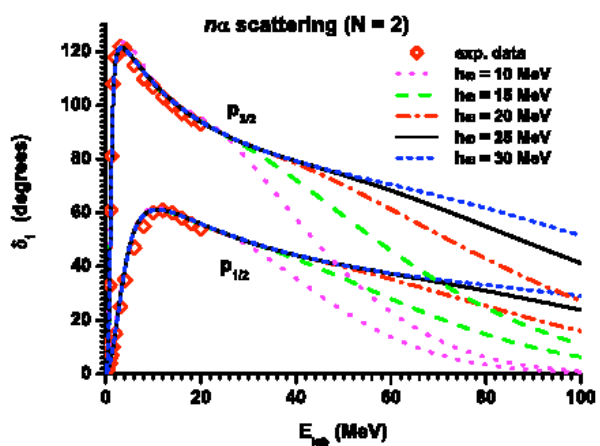
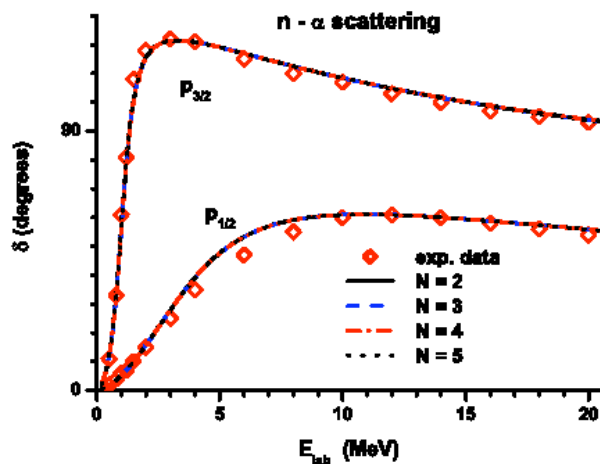
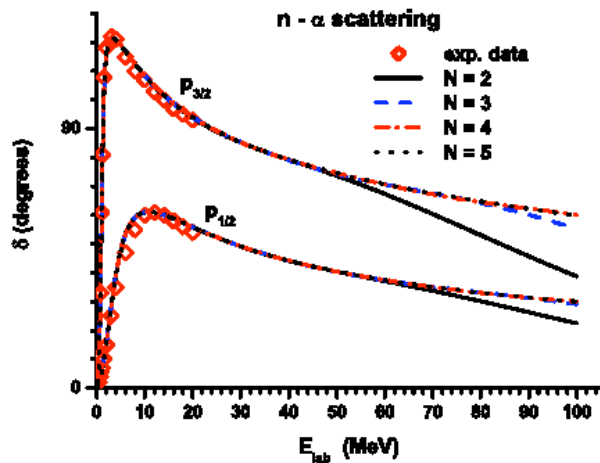


$n\alpha$ scattering



$n\alpha$ inverse scattering

- J -matrix inverse N -nucleus scattering analysis can be used to derive resonance parameters (position, width)
- J -matrix inverse N -nucleus scattering analysis suggests values for resonant and non-resonant states that should be compared with obtained in NCSM



$n\alpha$ inverse scattering

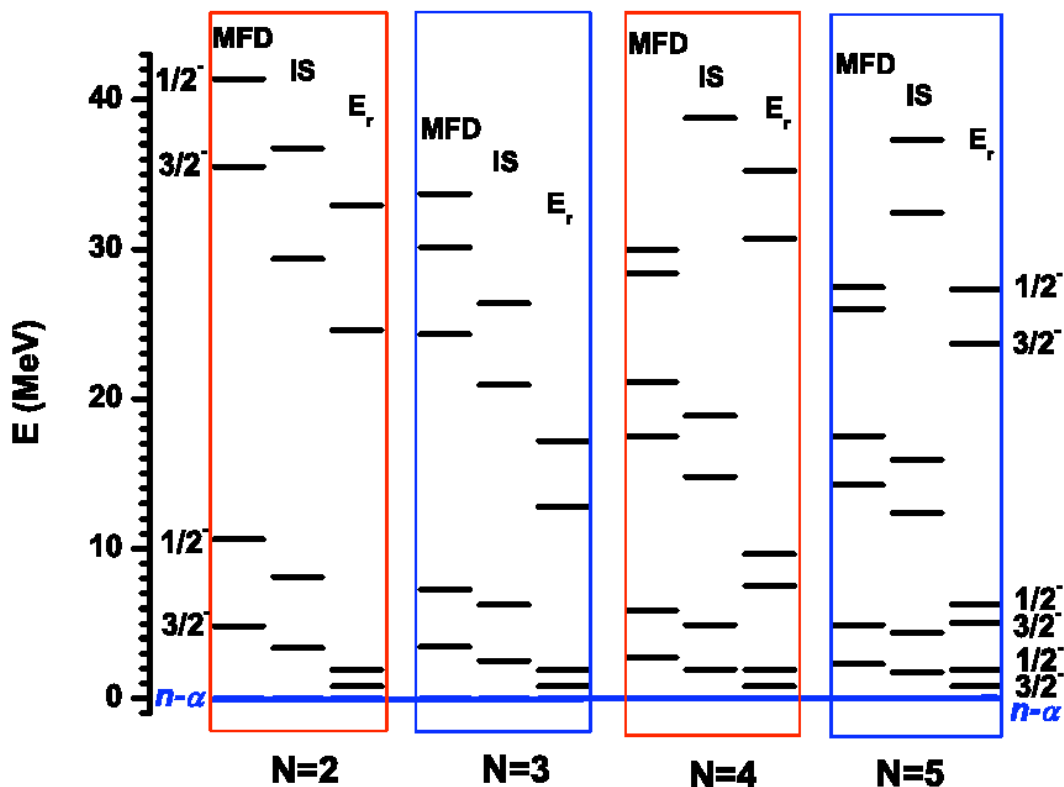
Table 1: S -matrix poles for $3/2^-$ states (energies and widths are in MeV).

	$N = 2$		$N = 3$		$N = 4$		$N = 5$	
	E_r	Γ	E_r	Γ	E_r	Γ	E_r	Γ
1	0.797	0.746	0.800	0.700	0.800	0.700	0.801	0.700
2	24.564	60.728	12.805	53.480	7.520	44.682	5.041	42.316
3	—	—	49.584	72.358	30.687	62.910	23.693	59.476
R -matrix (A. Csóttó, G.M. Hale. Phys. Rev. C55 (1997) 536):								
$E_r = 0.80, \Gamma = 0.65$								

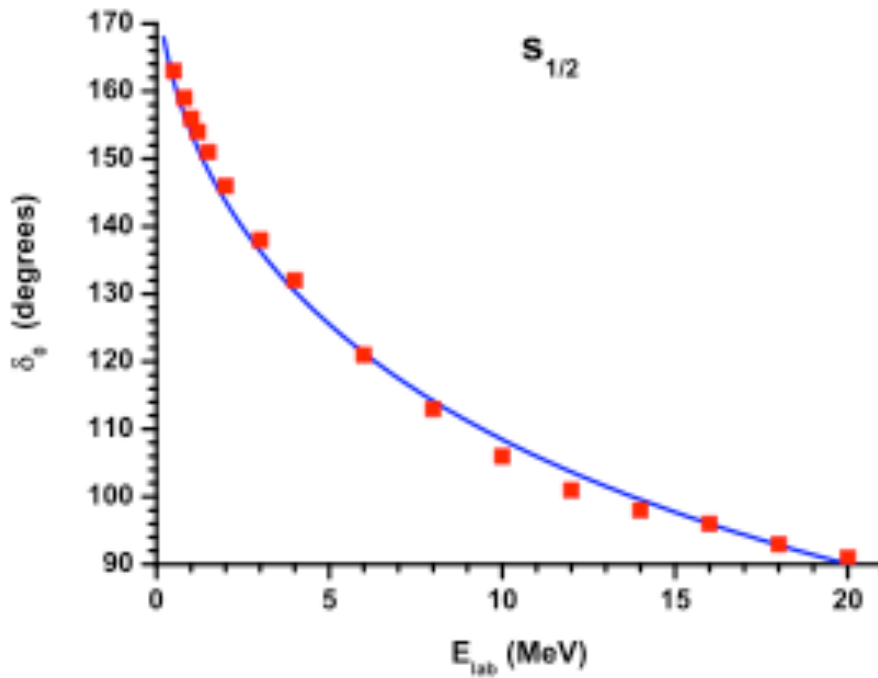
Table 2: S -matrix poles for $1/2^-$ states (energies and widths are in MeV).

	$N = 2$		$N = 3$		$N = 4$		$N = 5$	
	E_r	Γ	E_r	Γ	E_r	Γ	E_r	Γ
1	1.894	5.254	1.898	5.264	1.901	5.266	1.902	5.266
2	32.950	80.408	17.185	67.992	9.651	55.774	6.299	52.252
3	—	—	—	—	35.231	75.304	27.297	70.006
R -matrix (A. Csóttó, G.M. Hale. Phys. Rev. C55 (1997) 536):								
$E_r = 2.07, \Gamma = 5.57$								

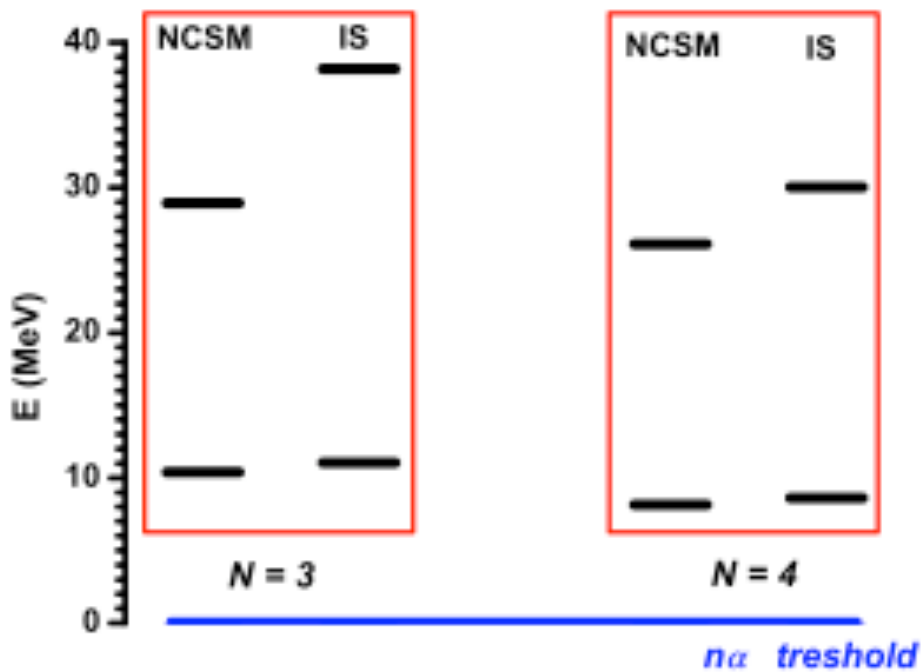
$n\alpha$ inverse scattering and NCSM



$n\alpha$ inverse scattering



$n\alpha$ inverse scattering and NCSM



Conclusions

- JISP16 provides a realistic description of two-body and many-body properties, comparable with modern realistic $NN + NNN$ forces
- Convergence of NCSM calculations with JISP16 is faster, even the bare JISP16 calculation convergence is reasonable, i.e. the results are more reliable. A confidence region of the binding energy predictions can be obtained for many nuclei by comparing the bare and effective interaction results
- Combining J -matrix formalism and NCSM one can perform scattering calculations with bare NN interaction; using JISP16 we describe well $n\alpha$ scattering. J -matrix inverse N -nucleus scattering analysis suggests values for resonant and non-resonant states that should be compared with obtained in NCSM

Plans

- JISP16 improvement by the fit to the same nuclei
- Charge-dependent JISP16
- Extending the calculations to the sd shell
- NCSM + J -matrix: Scattering calculations