#### **SELF-CONSISTENT APPROACH TO THE**

#### **GAMOW-TELLER BETA DECAY OF PROTON-RICH KR ISOTOPES**

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 $\cdot$  <sup>74</sup>Kr  $\rightarrow$  <sup>74</sup>Br  $0^{+}_{ground-state}$   $\rightarrow 1^{+}$ 

- $\cdot$  <sup>72</sup>Kr  $\rightarrow$  <sup>72</sup>Br
	- $0^+_{ground-state} \rightarrow 1^+$ <br>  $0^+_{first-excited} \rightarrow 1^+$ <br>  $2^+_{yrast} \rightarrow 1^+$

## within

the complex EXCITED VAMPIR variational approach

## **VAMPIR** - **Variational approaches to the nuclear many-body problem**

### **Framework**

- $\bullet$  the model space is defined by a finite dimensional set of spherical single particle states
- $\bullet$  the effective many-body Hamiltonian is represented as a sum of one- and two-body terms
- the basic buiding blocks are Hartree-Fock-Bogoliubov (HFB) vacua
- the HFB transformations are essentially *complex* and allow for proton-neutron, parity and angular momentum mixing being only restricted by time-reversal and axial symmetry
- the broken symmetries (s=N, Z, I,  $\pi$ ) are restored before variation by projection techniques

Variational approaches to the nuclear many-body problem with symmetry projection before variation

Model space

 $\{|i\rangle \equiv |\tau n l j m\rangle\}$  $\{c_i^\dagger, c_k^\dagger, ...\}_M$  ${c_i, c_k, ...\}_M$ 

Effective many-body Hamiltonian

$$
\hat{H} = \sum_{i=1}^{M} \varepsilon(i) c_i^{\dagger} c_i + \frac{1}{4} \sum_{i,k,r,s=1}^{M} v(ikrs) c_i^{\dagger} c_k^{\dagger} c_s c_r
$$

Hartree-Fock-Bogoliubov transformation

$$
\begin{pmatrix} a^{\dagger} \\ a \end{pmatrix} = F \begin{pmatrix} c^{\dagger} \\ c \end{pmatrix} = \begin{pmatrix} A^T & B^T \\ B^{\dagger} & A^{\dagger} \end{pmatrix} \begin{pmatrix} c^{\dagger} \\ c \end{pmatrix}
$$

 $\it Quasi\mbox{-}particle$ vacuum

$$
|F\rangle = \prod_{\alpha=1}^{M'} a_{\alpha}|0\rangle \quad \text{with} \quad \left\{ \begin{array}{ll} a_{\alpha}|0\rangle \neq 0 & \text{for } \alpha = 1, ..., M' \leq M \\ a_{\alpha}|0\rangle = 0 & \text{else} \end{array} \right\}
$$

 $\hat{\Theta}_{MK}^s \equiv \hat{P}(I; MK)\hat{Q}(N)\hat{Q}(Z)\hat{p}(\pi)$  $\hat{p}(\pi) \equiv \frac{1}{2} \left( 1 + \pi \hat{\Pi} \right)$ 

 $\hat{Q}(N_{\tau}) \, \equiv \, \frac{1}{2\pi} \int_0^{2\pi} d\phi_{\tau} \exp\{i\phi_{\tau}(N_{\tau} \, - \, \hat{N}_{\tau})\})$ 

 $\hat{P}(I;MK) \equiv \frac{2I+1}{8\pi^2} \int d\Omega D_{MK}^{I*}(\Omega) \hat{R}(\Omega)$ 

$$
|\psi(F^s); sM\rangle = \sum_{K=-I}^{+I} \hat{\Theta}^s_{MK} |F^s\rangle f^s_K
$$

$$
|\psi(F^s);sM\rangle\,=\,\frac{\hat{\Theta}^s_{M0}|F^s\rangle}{\sqrt{\langle F^s|\hat{\Theta}^s_{00}|F^s\rangle}}
$$

$$
|F\rangle = \left\{\prod_{m=1/2}^{m_{max}} \left(\prod_{\alpha}^{(m)} [u_{\alpha} + v_{\alpha}b_{\alpha}^{\dagger}b_{\bar{\alpha}}^{\dagger}]]\right)\right\}|0\rangle
$$

$$
b^{\dagger}_{\alpha}\,=\,\sum\limits_{\tau_i,n_i,l_i,j_i}^{(m_{\alpha}>0)}D_{i\alpha}^*\,c^{\dagger}_i
$$

$$
b^\dagger_\alpha b^\dagger_{\bar{\alpha}} = \mathop{\sum}\limits_{\tau = p,n}\limits^{\left(m_\alpha\tau\right)}\mathop{\sum}\limits_{\substack{\mathbf{i}<\mathbf{k}}} \left[1+\delta(\mathbf{i},\mathbf{k})\right]^{-1}\mathop{\sum}\limits_{I} (-)^{j_{k}+l_{k}-m_\alpha}(j_{i}j_{k}I|m_\alpha-m_\alpha 0)
$$

 $\times\{[Re(D_{i_{\tau}\alpha}^{*}D_{k_{\tau}\alpha})[1+(-)^{l_{i}+l_{k}+I}]+iIm(D_{i_{\tau}\alpha}^{*}D_{k_{\tau}\alpha})[1-(-)^{l_{i}+l_{k}+I}]] [c_{1}^{\dagger}c_{k}^{\dagger}]_{12\tau}^{I0}\}$ 

$$
+\sum_{\underline{i}}^{(m_{\alpha}p)(m_{\alpha}n)} \sum_{\underline{k}} [1/21/2T] - 1/21/20) (-)^{j_{k}+l_{k}-m_{\alpha}} (j_{i}j_{k}I|m_{\alpha}-m_{\alpha}0)
$$

 $\times\{[Re(D_{i_{p}\alpha}^{*}D_{k_{n}\alpha})[1+(-)^{l_{i}+l_{k}+I}]+iIm(D_{i_{p}\alpha}^{*}D_{k_{n}\alpha})[1-(-)^{l_{i}+l_{k}+I}]][c_{1}^{\dagger}c_{k}^{\dagger}]_{T_{0}}^{I_{0}}\}$ 

$$
\left[c_{\underline{i}}^{\dagger}c_{\underline{k}}^{\dagger}\right]_{TT_{z}}^{IM} \equiv \sum_{m_{i}m_{k}\tau_{i}\tau_{k}}(j_{i}j_{k}I|m_{i}m_{k}M)(\frac{1}{2}\frac{1}{2}T|\tau_{i}\tau_{k}T_{z})c_{i}^{\dagger}c_{k}^{\dagger}
$$

**Variational procedures**

*complex* Vampir approach

$$
E^s[F_1^s] = \frac{\langle F_1^s | \hat{H} \hat{\Theta}_{00}^s | F_1^s \rangle}{\langle F_1^s | \hat{\Theta}_{00}^s | F_1^s \rangle}
$$

$$
|\psi(F_1^s); sM\rangle = \frac{\hat{\Theta}_M^s e |F_1^s \rangle}{\sqrt{\langle F_1^s | \hat{\Theta}_0^s | F_1^s \rangle}}
$$

## *complex* **Excited Vampir approach**

$$
\begin{aligned}\n|\psi(F_2^s); sM\rangle &= \hat{\Theta}_{M0}^s \left\{ |F_1^s \rangle \alpha_1^2 + |F_2^s \rangle \alpha_2^2 \right\} \\
|\psi(F_i^s); sM\rangle &= \hat{\Theta}_{M0}^s \sum_{j=1}^i |F_j^s \rangle \alpha_j^i \quad \text{for} \quad i = 1, ..., n \\
|\Psi_{\alpha}^{(n)}; sM \rangle &= \sum_{i=1}^n |\psi_i; sM \rangle f_{i\alpha}^{(n)}, \quad \alpha = 1, ..., n\n\end{aligned}
$$

$$
(H - E^{(n)}N) f^n \, = \, 0
$$

$$
(f^{(n)})^+ N f^{(n)} = 1
$$

 $A = 70 - 90$  mass region

 $40Ca - core$ 

model space  $(\pi,\nu)$ :  $1p_{1/2}$   $1p_{3/2}$   $0f_{5/2}$   $0f_{7/2}$   $1d_{5/2}$   $0g_{9/2}$ renormalized G-matrix (OBEP, Bonn A)

- $\bullet$  short range Gaussians in the  $nn$ , pp, np channels
- monopole shifts:

 $\langle 0g_{9/2}0f; T=0|\hat{G}|0g_{9/2}0f; T=0 \rangle$  $\langle 1p1d_{5/2}; T=0|\hat{G}|1p1d_{5/2}; T=0 \rangle$ 

 $f_{5/2}$   $f_{7/2}$  $(ms1): -0.590 \, MeV / -0.060 \, MeV$  $(ms2): -0.500 \, MeV / -0.150 \, MeV$  $(ms3): -0.400 \; MeV / -0.250 \; MeV$ 

# *Gamow-Teller β Decay of 74Kr*

*CERN/ISOLDE E. Poirier et al., Phys.Rev. C69(2004)034307*

 ${}^{74}\text{Kr} \rightarrow {}^{74}\text{Br}$  $0^{+}_{ground-state}$   $\rightarrow 1^{+}$ 

 $Q_{EC} = 3.140 \pm 0.060 \text{ MeV}$ 

The amount of mixing for the ground-state of  $74$ Kr.



The amount of mixing for the lowest calculated  $1^+$  states of  $7^4Br$  (msl).

 $o$ -mixing /p-mixing

 $94(3)(3)\%$  $61(35)(2)(1)\%$  $89(3)(2)(2)(1)(1)(1)\%$  $44(28)(19)(4)(1)(1)(1)\%$ 97%  $69(19)(5)(2)(2)\%$  $70(7)(3)(2)(1)(1)(1)\%4(3)(2)(2)\%$  $9(3)\% 25(24)(11)(10)(3)(2)(2)(1)(1)(1)(1)(1)(1)(1)\%$  $7(1)(1)\% 71(8)(5)(1)(1)(1)(1)(1)\%$  $57(3)(2)(1)(1)(1)(1)\%13(5)(4)(2)(2)(1)(1)(1)(1)\%$  $26(1)(1)\%36(20)(4)(3)(2)(1)(1)(1)\%$  $21(21)(14)(14)(6)(5)(4)(2)(2)(1)(1)(1)(1)(1)(1)\%$  $2(1)(1)(1)\%$  36(14)(12)(6)(6)(5)(4)(2)(2)(1)(1)(1)(1)(1)(1)(3)%  $10(2)(2)(1)(1)\% 27(13)(11)(9)(3)(3)(2)(2)(2)(2)(1)(1)(1)(1)(1)(1)(1)(1)\%$  $50(16)(9)(5)(3)(3)(2)(2)(2)(1)(1)(1)\% 2(2)\%$  $33(21)(12)(8)(5)(5)(3)(2)(2)(1)(1)\% 1(1)(1)\%$  $1(1)\%34(18)(13)(9)(4)(3)(2)(2)(2)(2)(1)(1)(1)(1)(1)(1)\%$ 



The amount of mixing for the lowest calculated  $1^+$  states of  $74Br$  (msl).

 $\overline{1}$ 

The amount of mixing for the lowest calculated  $1^+$  states of  $74$ Br (ms2).

 $25(19)(11)(10)(8)(8)(3)(3)(2)(2)(1)(1)(1)(1)\% 2(1)\%$ 



The amount of mixing for the lowest calculated  $1^+$  states of  $74\text{Br (ms1)}$ .

The amount of mixing for the lowest calculated  $1^+$  states of  $7^4Br$  (ms3).

Spectroscopic quadrupole moments  $Q_2^{sp}$  (in  $efm^2$  ) for the lowest calculated  $1^+$ states of the  $^{74}Br$  nucleus (ms1).

 $1_{II}^{+}$   $1_{III}^{+}$  ........  $1^+_I$ 





Spectroscopic quadrupole moments  $Q_2^{sp}$  (in  $efm^2$  ) for the lowest calculated  $1^+$  states of the  $^{74}\rm{Br}$  nucleus (ms1).

Spectroscopic quadrupole moments  $Q_2^{sp}$  (in  $efm^2$ ) for the lowest calculated 1<sup>+</sup> states of the <sup>74</sup>Br nucleus (ms<sup>2</sup>).



 $1^+$  $1^+_{III}$  ........  $1^+_I$ 



Spectroscopic quadrupole moments  $Q_2^{sp}$  (in  $efm^2$ ) for the lowest calculated 1<sup>+</sup> states of the  $^{74}Br$  nucleus (ms3).

 $1<sup>+</sup>$ 

 $1^+$ 

 $1^+_{III}$  – ………

 $-49.5$   $47.8$   $-48.0$   $40.9$   $-51.3$   $45.5$   $-49.8$  $-49.3$   $48.4$   $-48.7$ 47.1  $-51.3$   $45.0$   $-50.1$   $-51.4$   $-44.9$   $38.1$   $-53.7$   $-49.6$   $-50.9$   $44.7$ 6.6  $-3.4$   $-47.3$   $-52.1$   $-46.2$   $-52.0$   $4.9$   $-11.9$   $-43.4$   $-45.4$   $38.9$   $-27.5$ 39.2 -44.9 30.9  $-42.0$   $-42.5$   $-45.3$   $-52.9$   $-44.2$   $-27.3$   $28.1$ 16.7 39.4  $-0.1$  $-4.7 -45.8$  $-37.3$   $-45.4$   $32.0$ 39.6 -40.3 -41.4 30.4  $13.0 \quad -16.0 \quad 38.4 \quad -34.4 \quad 25.8 \quad -55.7 \quad -15.5$  $-43.5 -24.2$  20.6 41.9 14.5 -52.2 40.9 35.2  $29.1$  $2.7\,$  $-15.9$  4.7



Excitation energy (MeV)



## *Gamow-Teller β Decay of 72Kr*

*CERN/ISOLDE I. Piqueras, Eur. Phys. J. A16(2003)313*

 ${}^{72}\text{Kr} \rightarrow {}^{72}\text{Br}$ 

 $Q_{EC}$  = 5.040  $\pm$  0.375 MeV

 $0^{+}_{ground-state}$   $\rightarrow 1^{+}$ 

 $0^{+}_{first-excited}$   $\rightarrow 1^{+}$ *E*  $\theta_2^+$  = 0.671 MeV

 $2^+_{yrast} \rightarrow 1^+$ *E*  $_{2_{I}^{+}}$  = 0.710 MeV The amount of mixing for the calculated states of the  ${}^{72}\text{Kr}$  nucleus (ms3).



The amount of mixing for the lowest calculated  $1^+$  states of <sup>72</sup>Br (ms3) with significant B(GT).

 $o$ -mixing / $p$ -mixing

 $85(12)\%$  $81(11)(4)\%$  $87(2)(2)(2)(2)(1)(1)\%$  $81(4)(4)(2)(2)(1)(1)(1)\%$  $78(16)(2)(1)\%$  $78(4)(3)(3)(2)(2)(1)(1)(1)(1)(2)$  $49(24)(8)(6)(5)(2)(1)(1)(1)\%$  $32(31)(15)(9)(3)(2)(1)(1)(1)(1)$ %  $79(15)(1)\%$  $31(2)(2)(1)\%20(16)(13)(2)(1)(1)(1)(1)(1)(1)(1)(1)\%$  $50(5)(1)(1)\% 12(10)(8)(2)(1)(1)(1)(1)(1)\%$  $2\%$  68(10)(5)(3)(3)(2)(1)(1)(1)\%  $36(24)(7)(6)(5)(4)(3)(3)(2)(1)(1)(1)\%$  $72(12)(4)(2)(1)(1)(1)(1)(1)\%1\%$  $62(17)(8)(4)(2)\% 1\%$  $56(15)(11)(2)(2)(1)(1)(1)(1)(1)\% 1(1)\%$ 

	$1^+_I$ $1^+_{II}$ $1^+_{III}$				
	48.5 48.7 -49.9 -49.4 46.5 45.5 -51.6 -50.1 -49.5 46.8 -11.5				
	8.7 -46.5 -48.7 45.4 44.0 -53.5 -39.1 27.0 41.0 -48.9 -46.5				
	$-49.2$ $42.5$ $-39.8$ $35.8$ $-46.3$ $41.8$ $-45.0$ $-43.5$ $42.4$ $-46.9$ $-46.6$				
	$-26.3$ 10.7 $-37.3$ 37.4 $-36.5$ 35.5 $-46.6$ 47.6 $-48.8$ $-40.0$ $-1.2$				
	$-24.0$ $-35.8$ $37.1$ $-47.7$ $-53.2$ $-42.8$ $27.0$ $-7.2$ $10.2$ $-45.8$ $-32.8$				
	30.8 40.7 -24.2 21.8 -23.9 -41.8 15.0 -13.5 -38.3 39.6 11.8				
	36.4 -47.6 -24.7 21.8 41.7 37.4 29.5 12.1 -20.2 -23.6 -39.3				
	$-33.2$ $37.6$ $27.7$ $-50.9$ $43.6$ $24.1$ $-10.7$ $15.6$ $-32.7$ $44.3$ $-46.4$				
	$-33.9$ $32.5$ $-42.2$ $-23.1$ $43.3$ $20.9$ $38.6$ $-44.1$ $-52.3$ $-45.8$ $21.0$				
	$-45.0$ $1.5$ $-1.8$ $-37.6$ $39.6$ $45.1$ $-48.9$ $-43.6$ $-23.9$ $31.5$ $36.1$				
	16.1 34.9 -53.6 43.2 -41.8 -45.9 -43.5				

Spectroscopic quadrupole moments  $Q_2^{sp}$  (in  $efm^2$  ) for the calculated  $1^+$  states of the  $^{72}\rm{Br}$  nucleus (ms3).







$$
\frac{1}{T_{1/2}} = \frac{g_A^2}{D} \sum_i f(Z, E_i) |\langle 1_i^+ || \beta^+ || 0^+ \rangle|^2
$$

 $D = 6146 s$   $g_A = 1.26$ 

 $T_{1/2}^{exp} = 17.1(2)$  s

 $T_{1/2}$ <sup>(gs)</sup> = 20.8 s

 $T_{1/2}$ (first-excited 0<sup>+</sup>) = 17.3 s

## **Summary and outlook**

• **the Gamow-Teller** β **decay of 74Kr was investigated for the first time within shape-coexistence and –mixing in both parent and daughter nucleus the complex Excited Vampir variational approach, describing consistently the** 

• **the first results concerning the Gamow-Teller strength distribution as well as the accumulated strength for the ground state, the first-excited 0+ and the yrast 2+ of 72Kr are obtained in a self-consistent approach. A good agreement with available data is revealed**

• **the uncertainties in the effective interaction require systematic investigations**

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