Hartree-Fock and Hartree-Fock-Bogoliubov with Modern Effective Interactions

Heiko Hergert

Institut für Kernphysik, TU Darmstadt

Overview

- Motivation
- Modern Effective Interactions
  - Unitary Correlation Operator Method
  - Similarity Renormalization Group
- Few-Body Systems
- Many-Body Systems
  - Hartree-Fock, Many-Body Perturbation Theory & Beyond
  - Hartree-Fock-Bogoliubov
- Conclusions

From QCD to Nuclear Structure

Nuclear Structure

Low-Energy QCD

From QCD to Nuclear Structure

Nuclear Structure

Realistic Nuclear Interactions

Low-Energy QCD

- chiral interactions: consistent NN & 3N interaction derived within χEFT
- traditional NN-interactions: Argonne V18, CD Bonn,...
- reproduce experimental NN phase-shifts with high precision
- induce strong short-range central & tensor correlations

From QCD to Nuclear Structure

Nuclear Structure

↑

Exact / Approx. Many-Body Methods

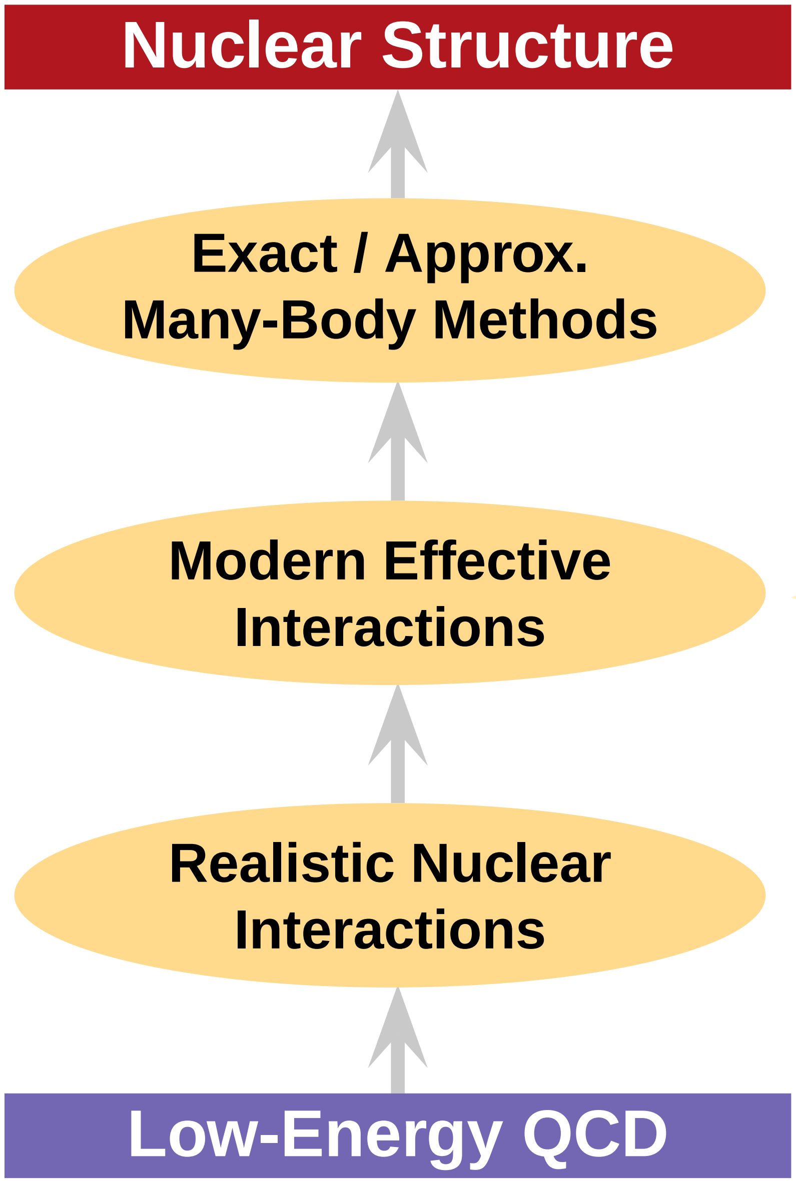
- ‘exact’ solution of the many-body problem for light and intermediate masses (GFMC, NCSM, CC,...)
- controlled approximations for heavier nuclei (HF & MBPT,...)
- rely on restricted model spaces of tractable size
- not suitable for the description of short-range correlations

Realistic Nuclear Interactions

↑

Low-Energy QCD

From QCD to Nuclear Structure



- adapt realistic potential to the available model space
  - tame short-range correlations
  - improve convergence behavior
- conserve experimentally constrained properties (phase shifts)
  - generate new realistic interaction
- provide consistent effective interaction & effective operators
- unitary transformations most convenient

Deuteron: Manifestation of Correlations

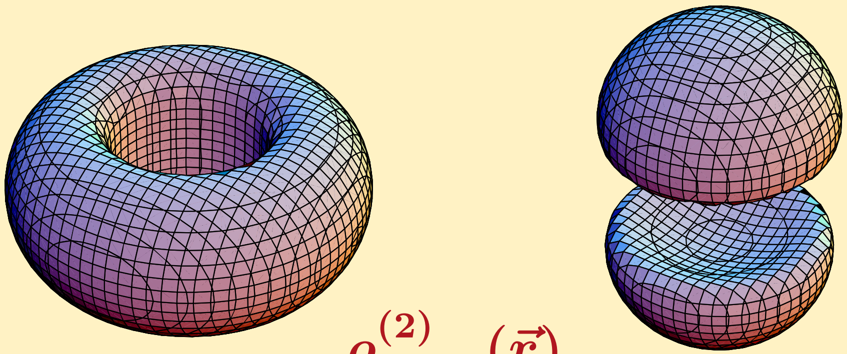
Argonne V18 Deuteron Solution

$M_S = 0$

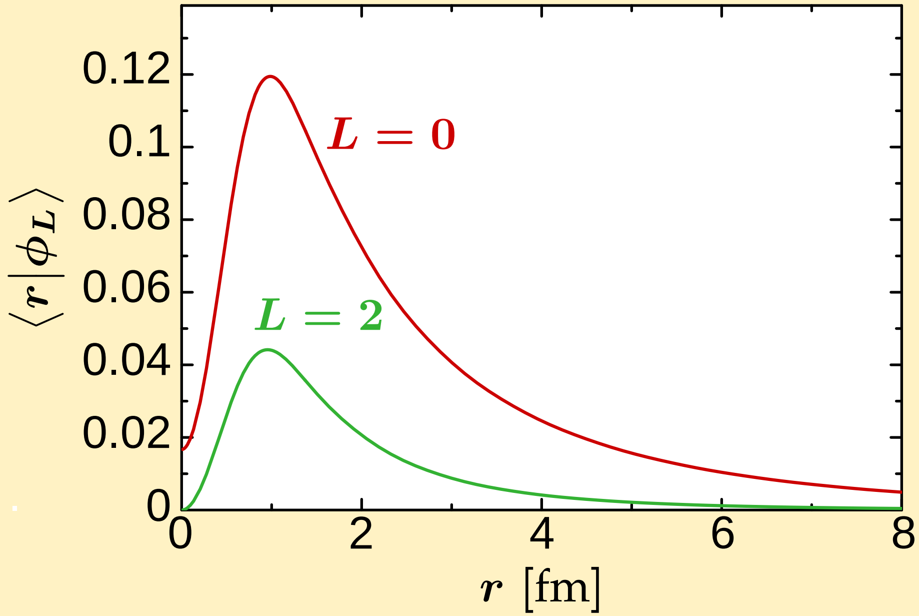
$\frac{1}{\sqrt{2}}\left(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle \right)$

$M_S = \pm 1$

$|\uparrow\uparrow\rangle, \; |\downarrow\downarrow\rangle$



$\rho^{(2)}_{1,M_S}(\vec{r})$



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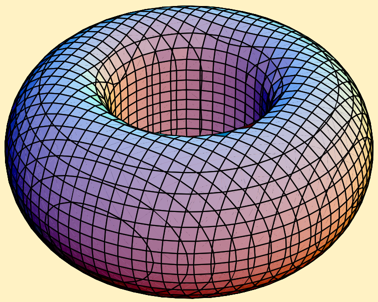

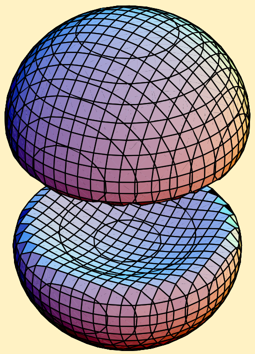

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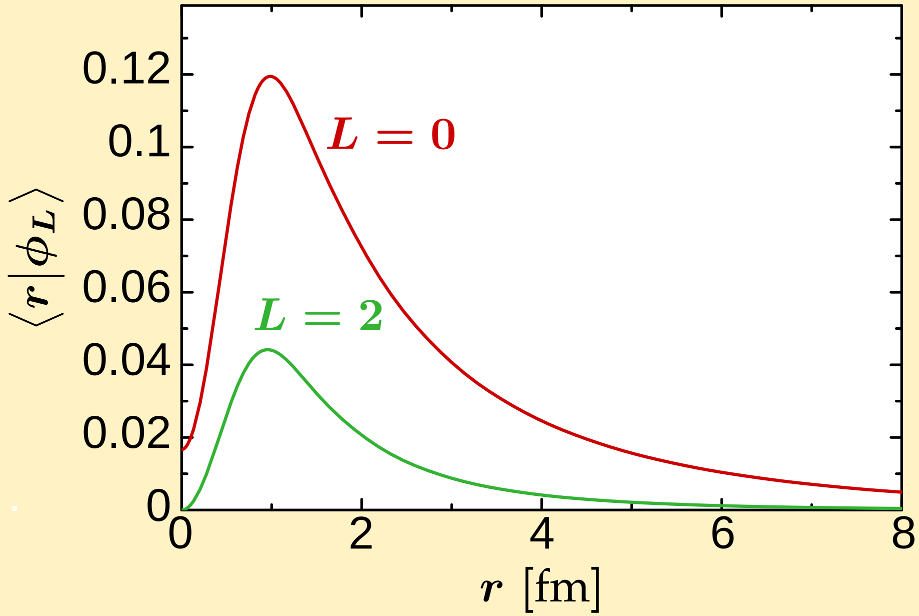

short-range repulsion
supresses wavefunction at
small distances r

central correlations

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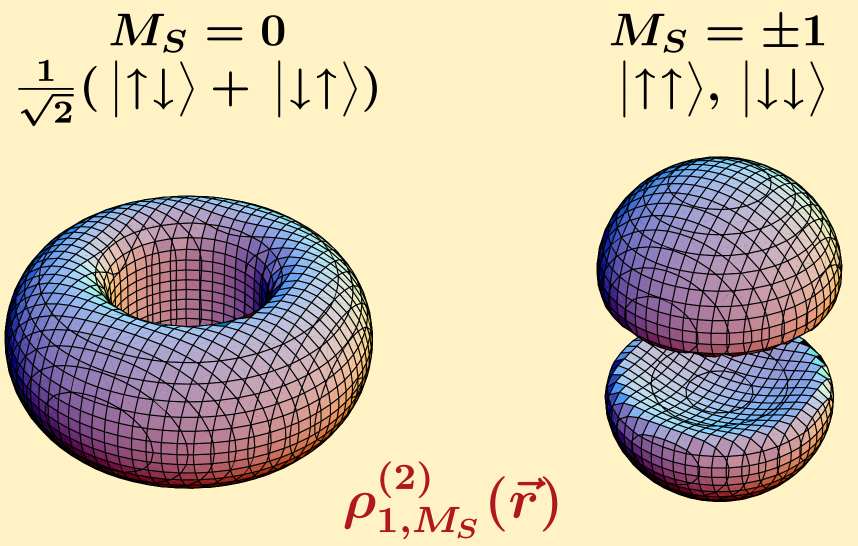

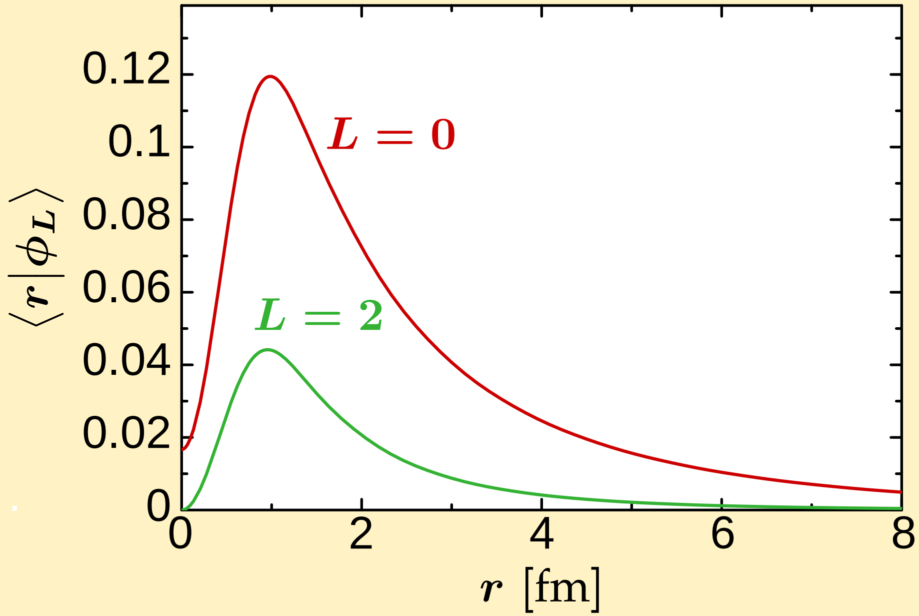

short-range repulsion supresses wavefunction at small distances r

central correlations

tensor interaction generates D-wave admixture in the ground state

tensor correlations

Modern Effective Interactions I

Unitary Correlation Operator Method (UCOM)

Unitary Correlation Operator Method

Correlation Operator

define an unitary operator $\mathbf{C}$ to describe the effect of short-range correlations

$$\mathbf{C} = \exp[-\mathrm{i}\,\mathbf{G}] = \exp\Big[-\,\mathrm{i}\sum_{i<j}\mathbf{g}_{ij}\Big]$$

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Correlated States

imprint short-range correlations onto uncorrelated many-body states

$$|\widetilde{\psi}\rangle = \mathrm{C}\;|\psi\rangle$$

Correlated Operators

adapt Hamiltonian and all other observables to uncorrelated many-body space

$$\widetilde{\mathrm{O}} = \mathrm{C}^{\dagger}\,\mathrm{O}\,\mathrm{C}$$

$$\langle\widetilde{\psi}|\,\mathrm{O}\,|\widetilde{\psi}'\rangle = \langle\psi|\,\mathrm{C}^{\dagger}\,\mathrm{O}\,\mathrm{C}\,|\psi'\rangle = \langle\psi|\,\widetilde{\mathrm{O}}\,|\psi'\rangle$$

Unitary Correlation Operator Method

explicit ansatz for the correlation operator motivated by the **physics of short-range central and tensor correlations**

Central Correlator C_r

- radial distance-dependent shift in the relative coordinate of a nucleon pair

$$\mathrm{g}_r = \frac{1}{2}\big[s(\mathrm{r})\ \mathrm{q}_r + \mathrm{q}_r\ s(\mathrm{r})\big]$$

$$\mathrm{q}_r = \tfrac{1}{2}\big[\tfrac{\vec{\mathrm{r}}}{\mathrm{r}} \cdot \vec{\mathrm{q}} + \vec{\mathrm{q}} \cdot \tfrac{\vec{\mathrm{r}}}{\mathrm{r}}\big]$$

Tensor Correlator C_Ω

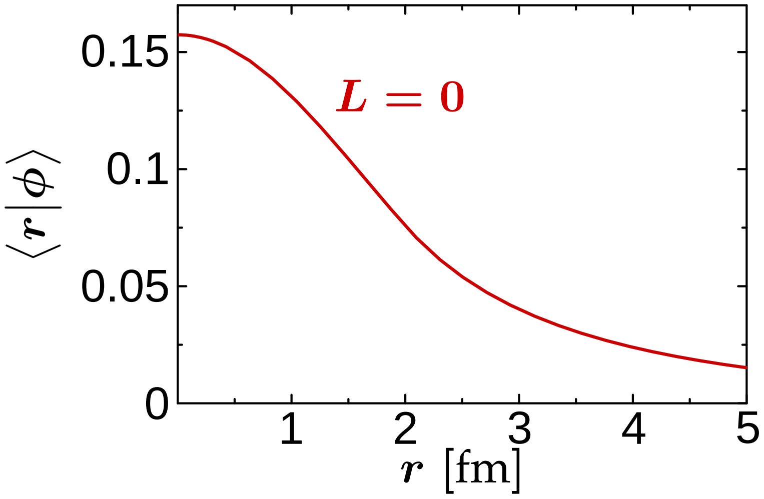
- angular shift depending on the orientation of spin and relative coordinate of a nucleon pair

$$\mathrm{g}_\Omega = \frac{3}{2}\vartheta(\mathrm{r})\big[(\vec{\sigma}_1 \cdot \vec{\mathrm{q}}_\Omega)(\vec{\sigma}_2 \cdot \vec{\mathrm{r}}) + (\vec{\mathrm{r}} \leftrightarrow \vec{\mathrm{q}}_\Omega)\big]$$

$$\vec{\mathrm{q}}_\Omega = \vec{\mathrm{q}} - \tfrac{\vec{\mathrm{r}}}{\mathrm{r}}\ \mathrm{q}_r$$

- $s(r)$ and $\vartheta(r)$ for given potential determined by energy minimization in the two-body system (for each S, T)

Correlated States: The Deuteron



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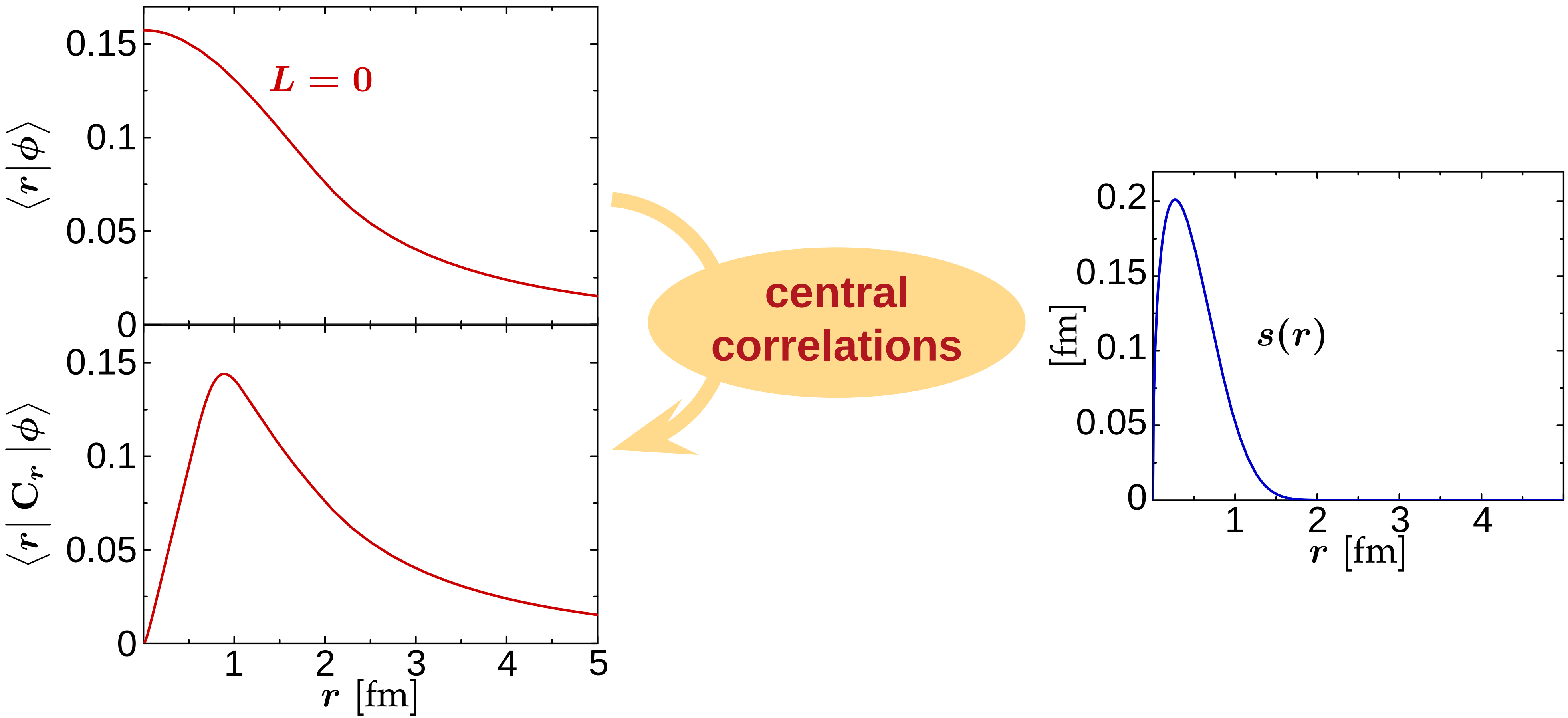

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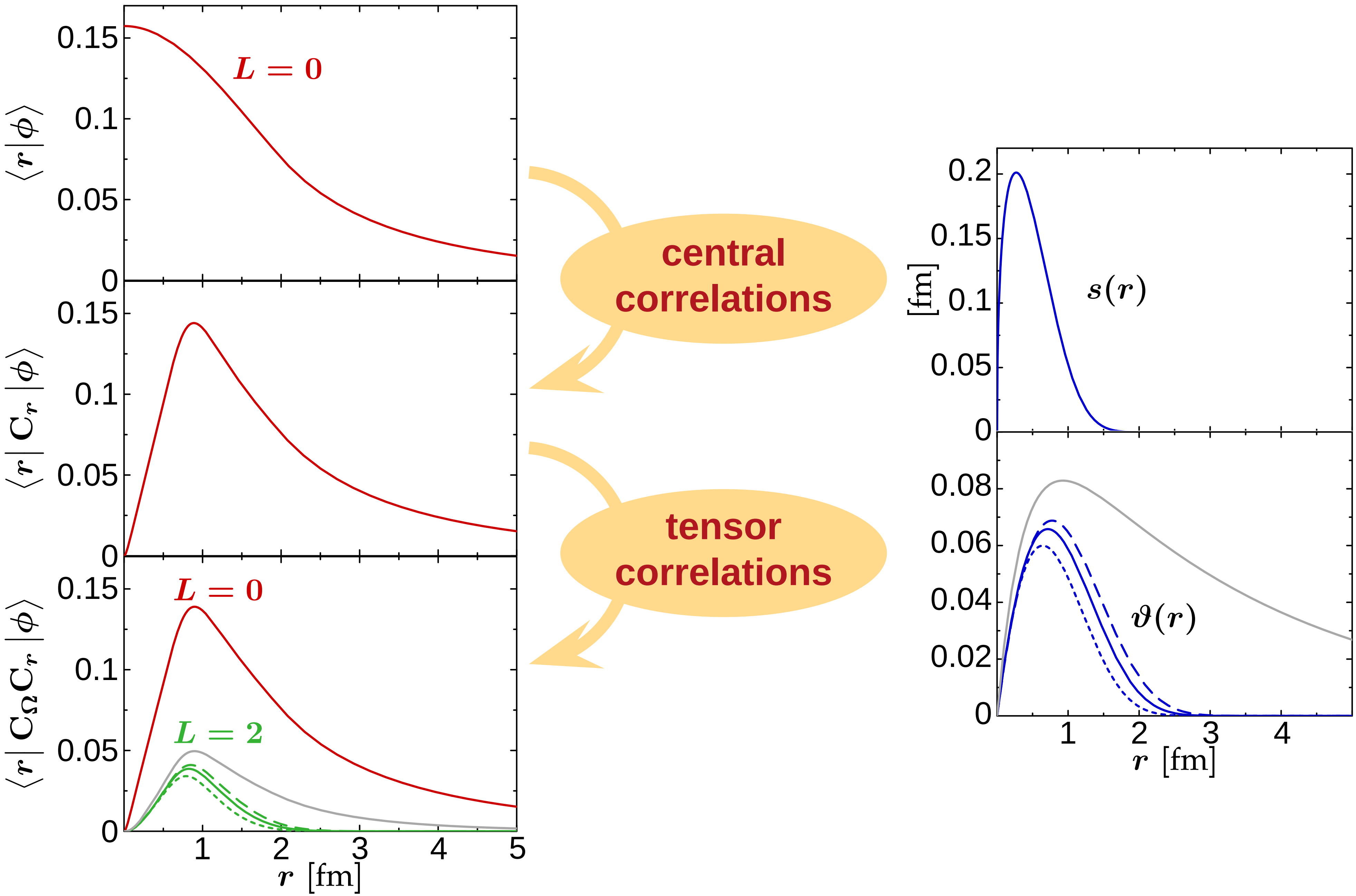

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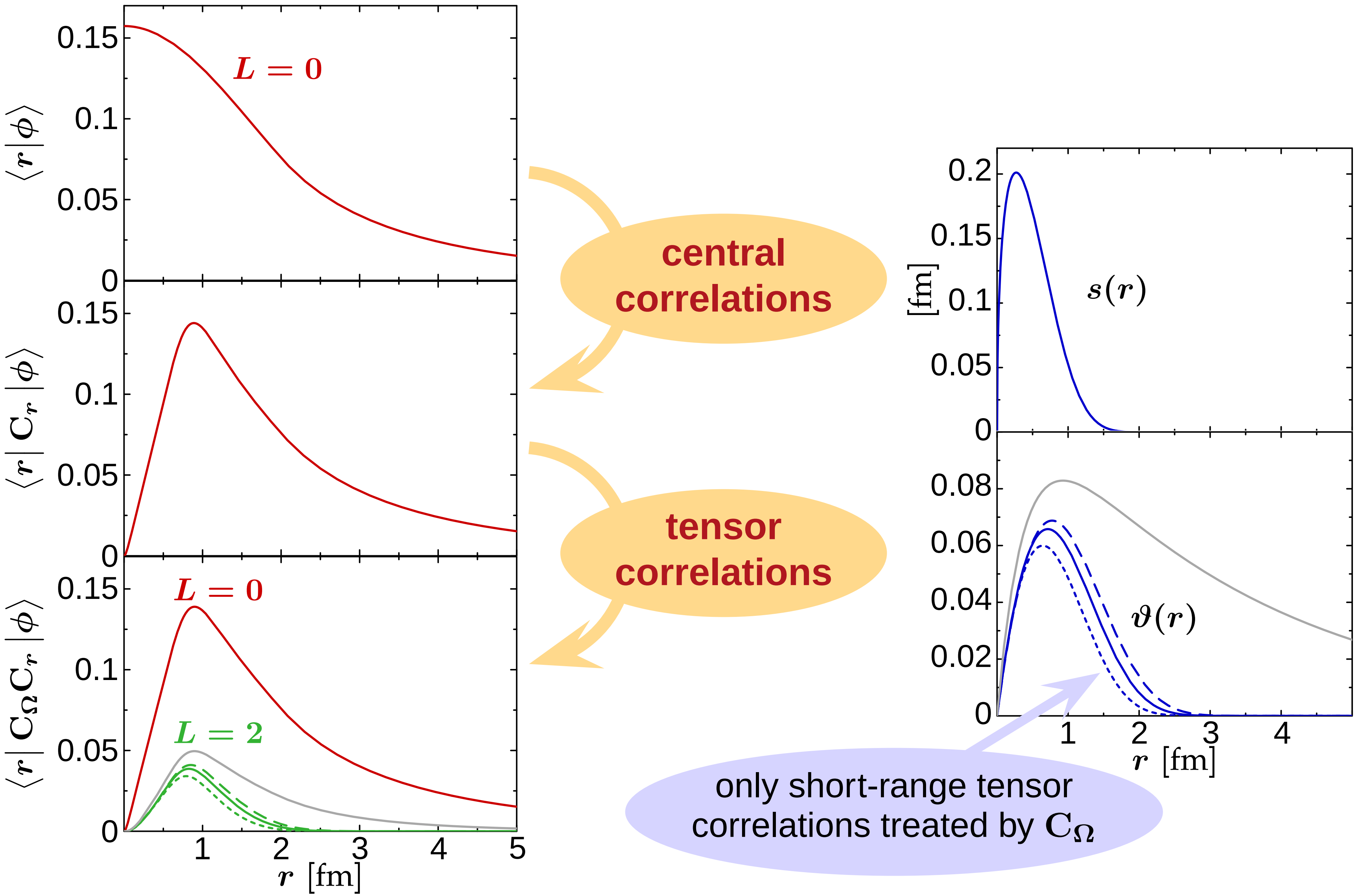
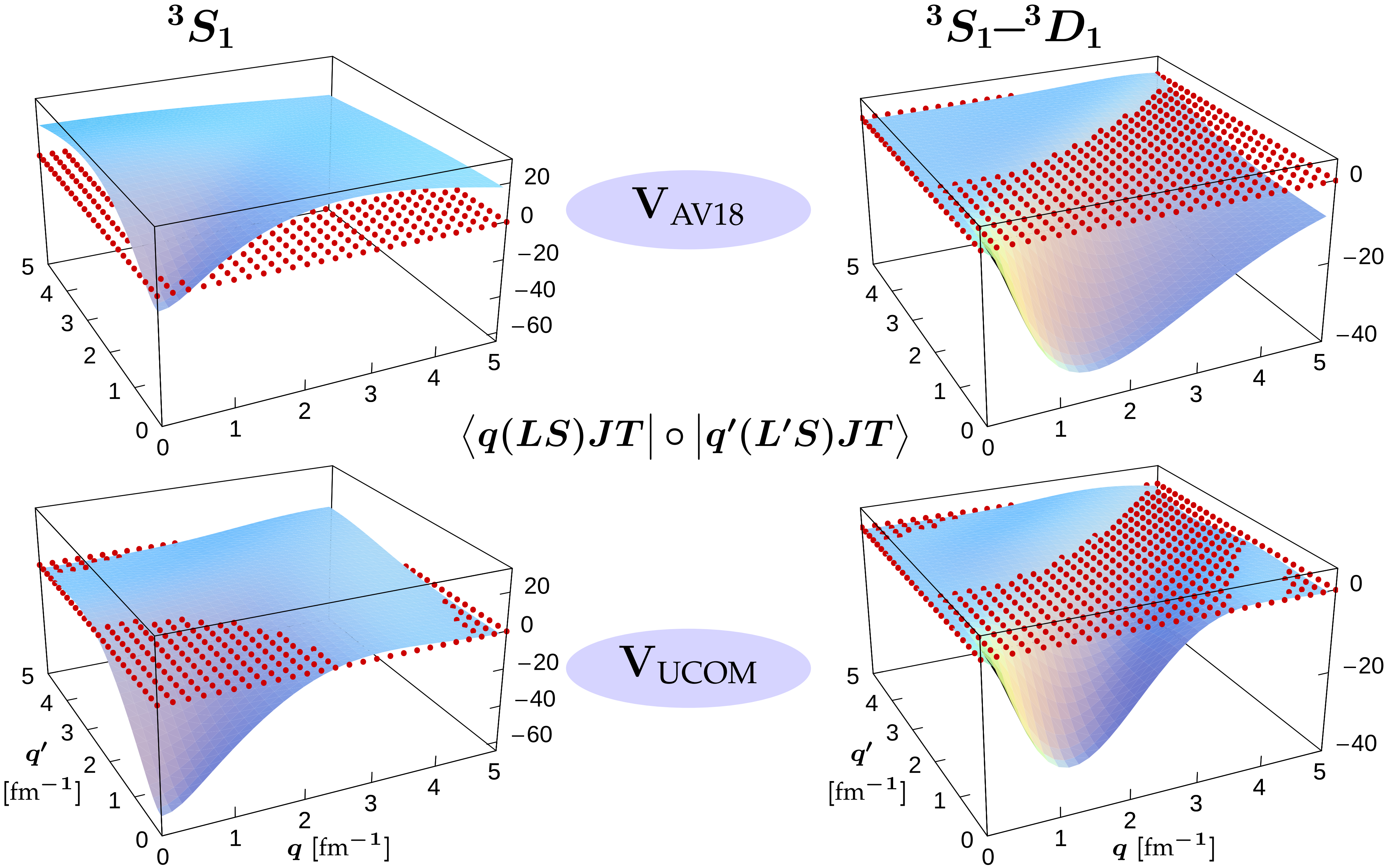
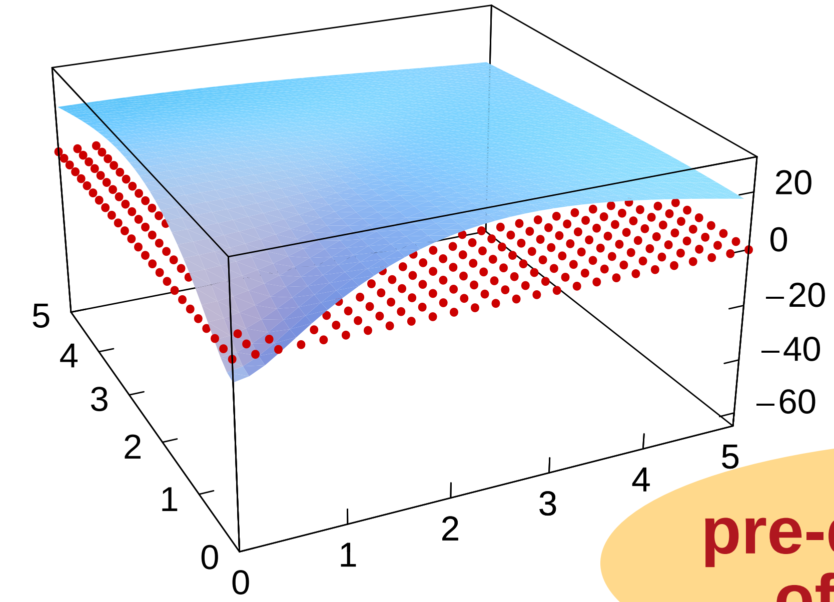


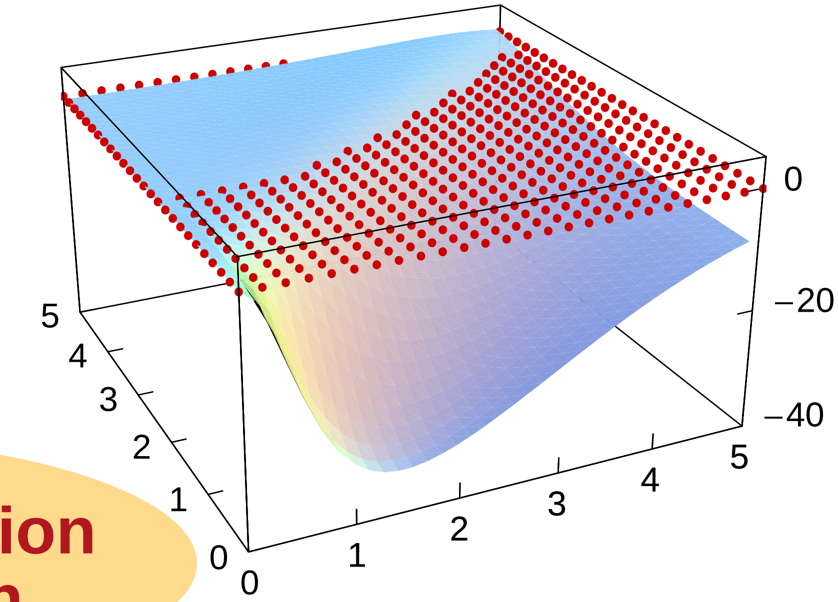
Correlated Interaction: V_{UCOM}



Correlated Interaction: V_{UCOM}

$^3S_1$ | $^3S_1$–$^3D_1$

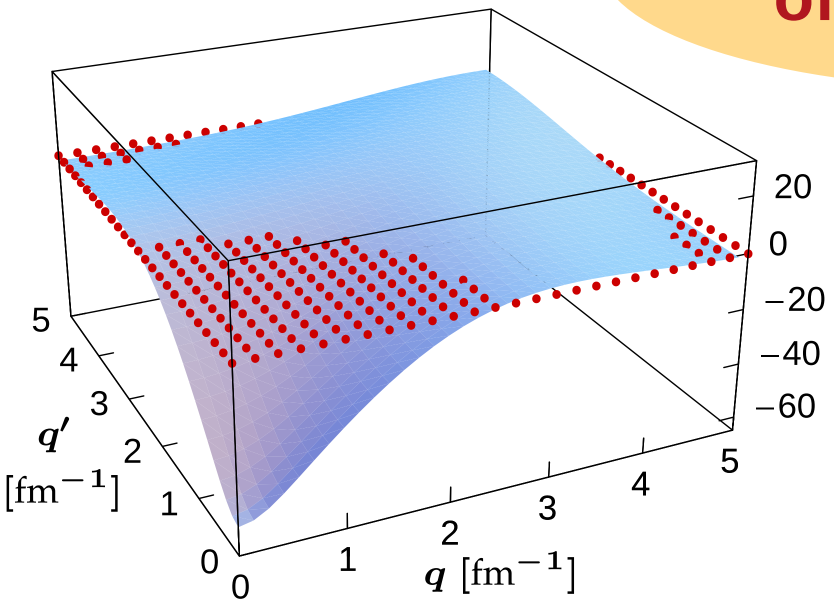


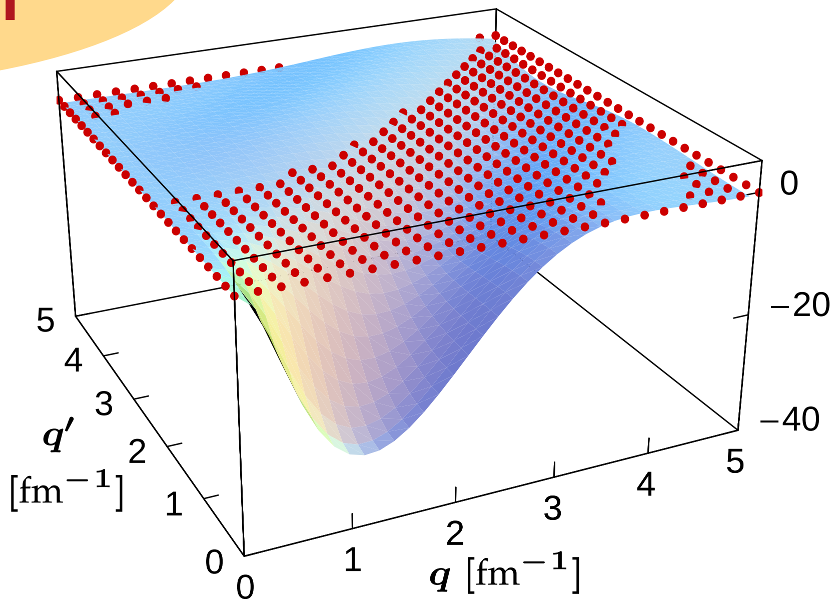


V_{AV18}

pre-diagonalization of Hamiltonian

V_{UCOM}





Modern Effective Interactions II

Similarity Renormalization Group (SRG)

Similarity Renormalization Group

unitary transformation of the **Hamiltonian to a band-diagonal form** with respect to a given uncorrelated many-body basis

Flow Equation for Hamiltonian

- evolution equation for Hamiltonian

$$\widetilde{\mathrm{H}}(\alpha) = \mathrm{C}^{\dagger}(\alpha)\,\mathrm{H}\,\mathrm{C}(\alpha) \qquad \longrightarrow \qquad \frac{\mathrm{d}}{\mathrm{d}\alpha}\widetilde{\mathrm{H}}(\alpha) = \left[\eta(\alpha), \widetilde{\mathrm{H}}(\alpha)\right]$$

- dynamical generator defined as commutator with the operator in whose eigenbasis H shall be diagonalized

$$\eta(\alpha) = \left[\mathrm{T}_{\mathrm{int}}, \widetilde{\mathrm{H}}(\alpha)\right] \overset{2B}{=} \frac{1}{2\mu}\left[\vec{\mathrm{q}}^{2}, \widetilde{\mathrm{H}}(\alpha)\right]$$

[Bogner et al., PRC**75** 061001(R) (2007); Hergert & Roth, PRC**75** 051001(R) (2007)]

The SRG Generator: A Closer Look

- typical NN interaction operators:

$$\mathrm{O}_p \in \{\mathbb{1}, \vec{\sigma}_1 \cdot \vec{\sigma}_2, \vec{\mathrm{l}}^{\,2}, \vec{\mathrm{l}} \cdot \vec{\mathrm{s}}, \mathrm{s}_{12}(\vec{\mathrm{r}}, \vec{\mathrm{r}}), \ldots\} \otimes \{\mathbb{1}, \vec{\tau}_1 \cdot \vec{\tau}_2, \ldots\}$$

Radial Kinetic Energy

$$\eta_r(0) \sim \left[\mathrm{q}_r^2, \mathrm{V}\right] = \sum_p \left[\mathrm{q}_r^2, v_p(\mathrm{r})\mathrm{O}_p\right] = \sum_p \left(\mathrm{q}_r v_p'(\mathrm{r})\mathrm{O}_p + \mathrm{O}_p v_p'(\mathrm{r})\mathrm{q}_r\right)$$

Angular Kinetic Energy

$$\eta_\Omega(0) \sim \left[\vec{\mathrm{l}}^{\,2}, \mathrm{V}\right] = \left[\vec{\mathrm{l}}^{\,2}, v_t(\mathrm{r})\mathrm{s}_{12}(\vec{\mathrm{r}}, \vec{\mathrm{r}})\right] = -4i v_t(\mathrm{r})\mathrm{s}_{12}(\vec{\mathrm{r}}, \vec{\mathrm{q}}_\Omega)$$

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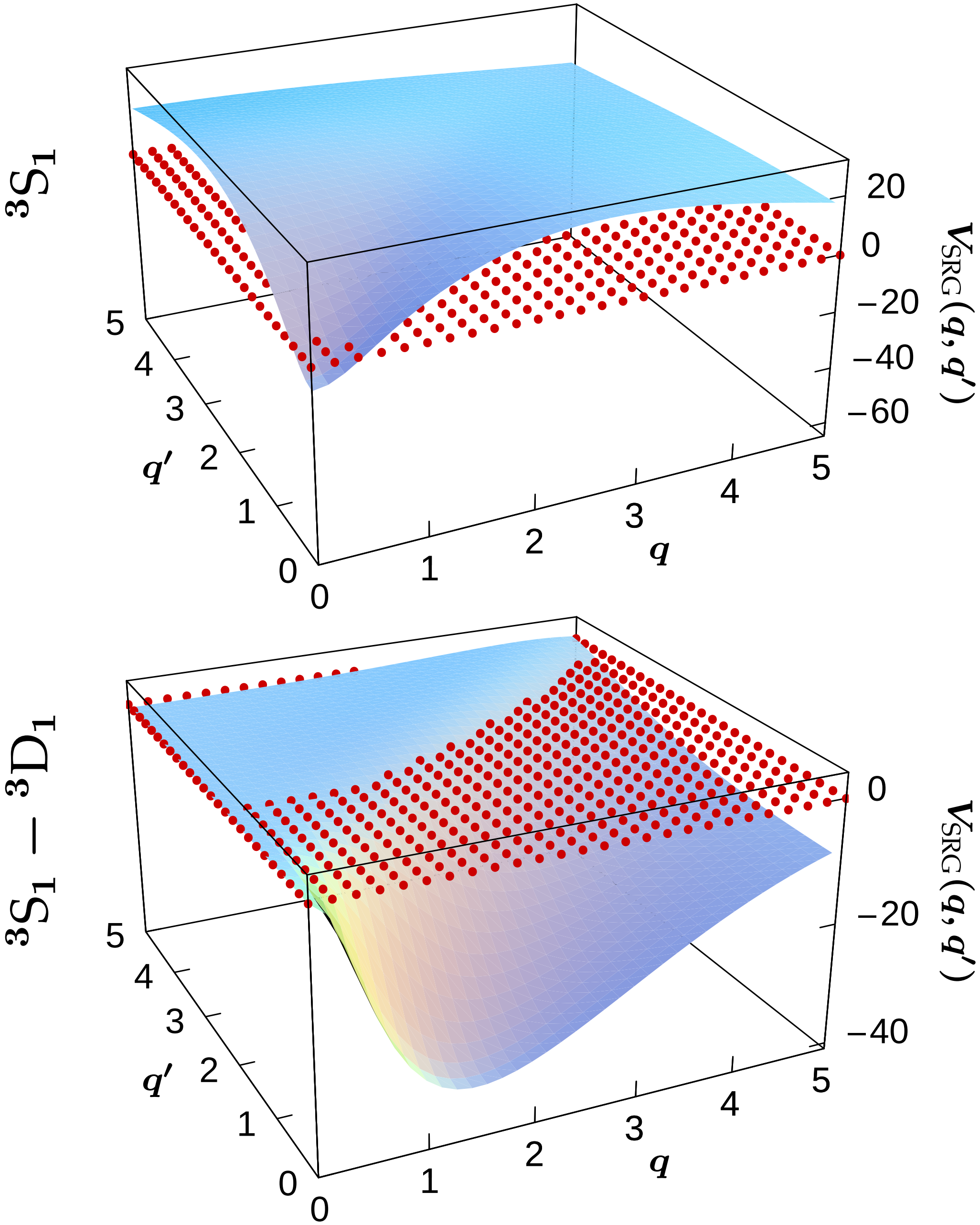
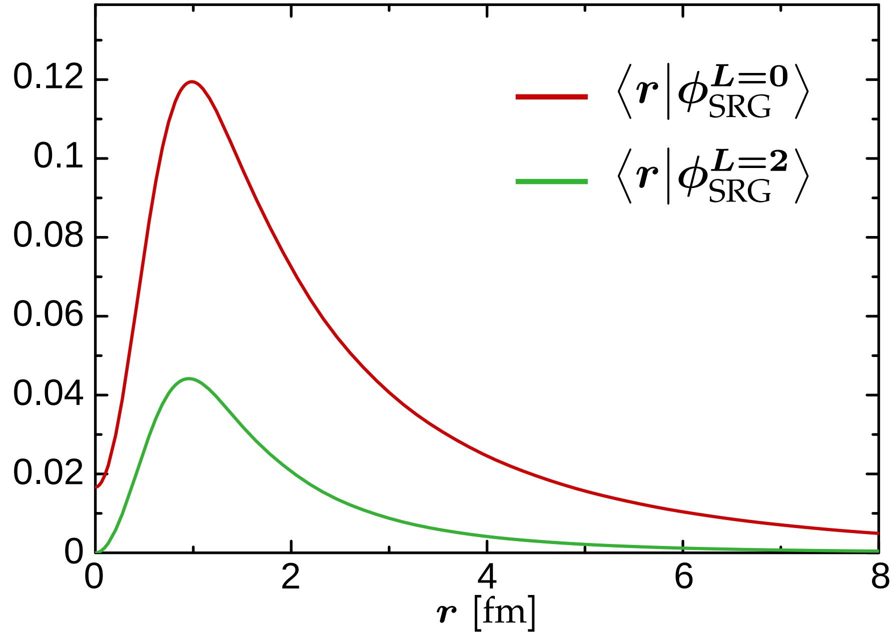
☞ $\eta(0)$ has the same structure as the UCOM generators g_r and g_Ω

SRG Evolution: The Deuteron



Argonne V18



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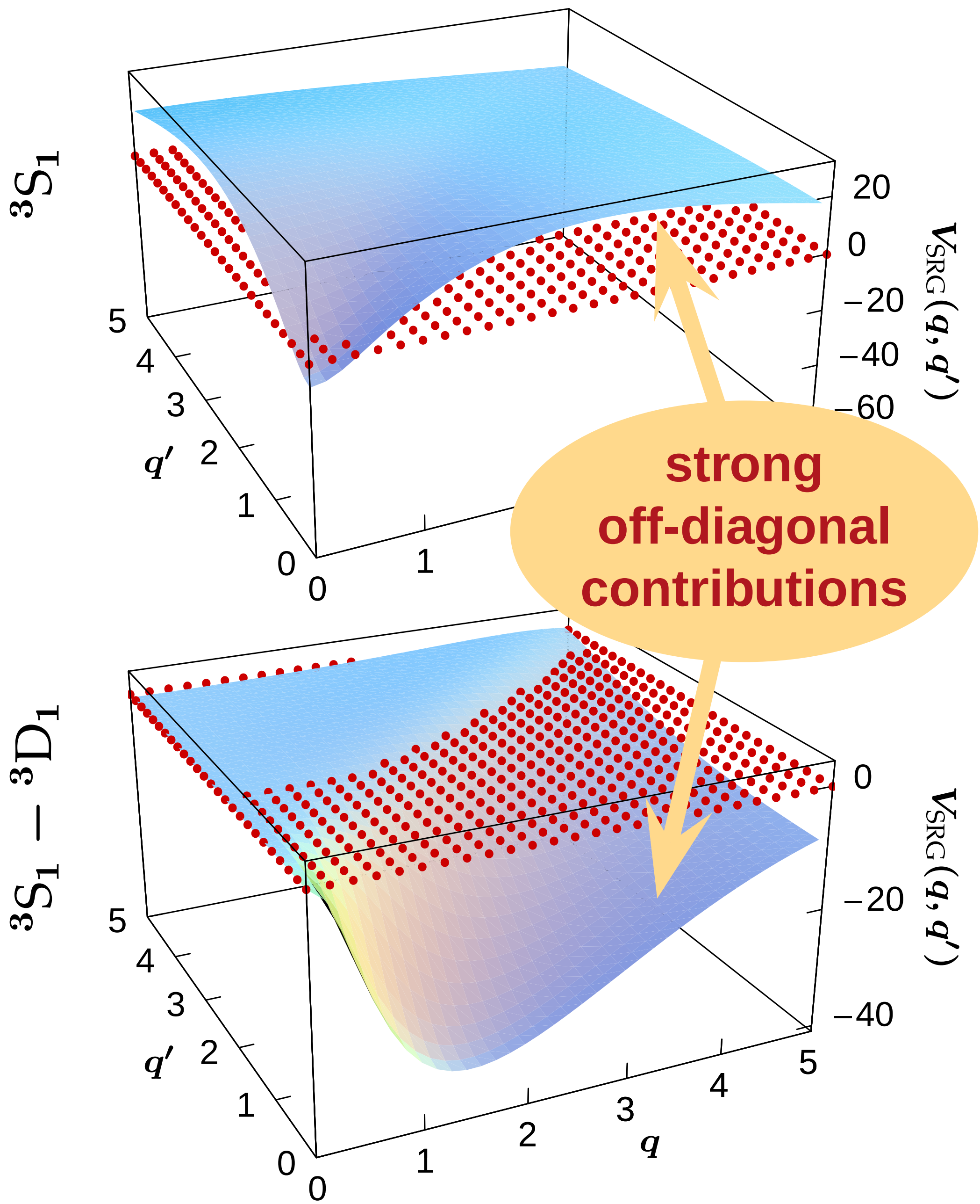

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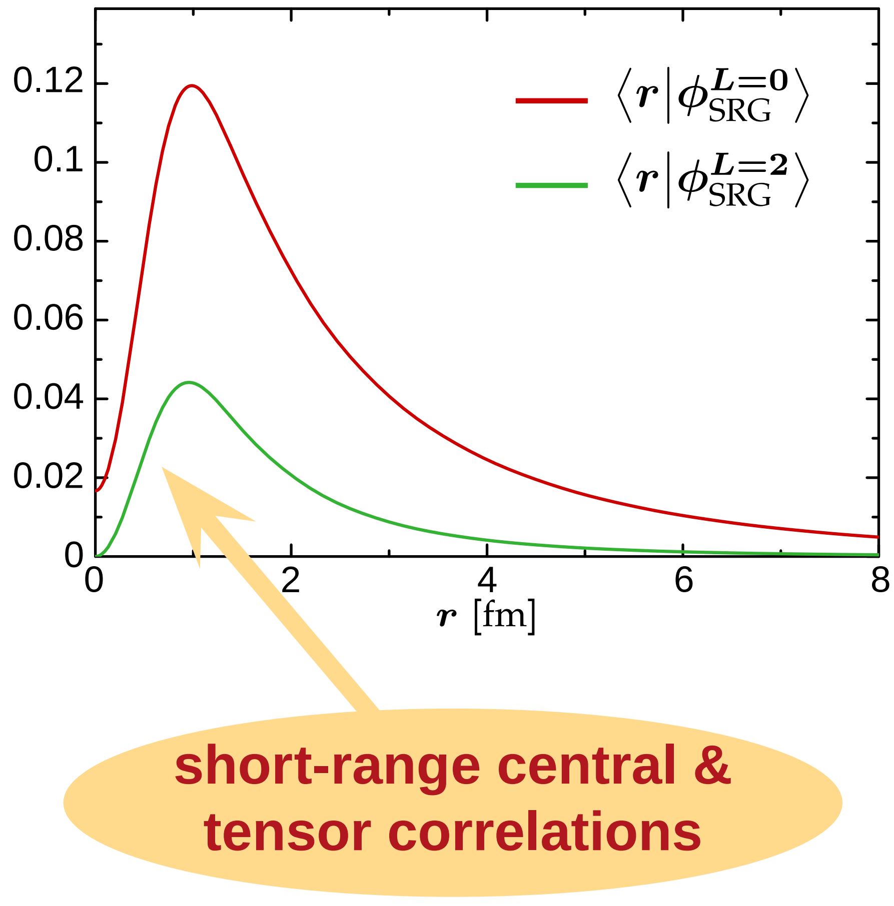

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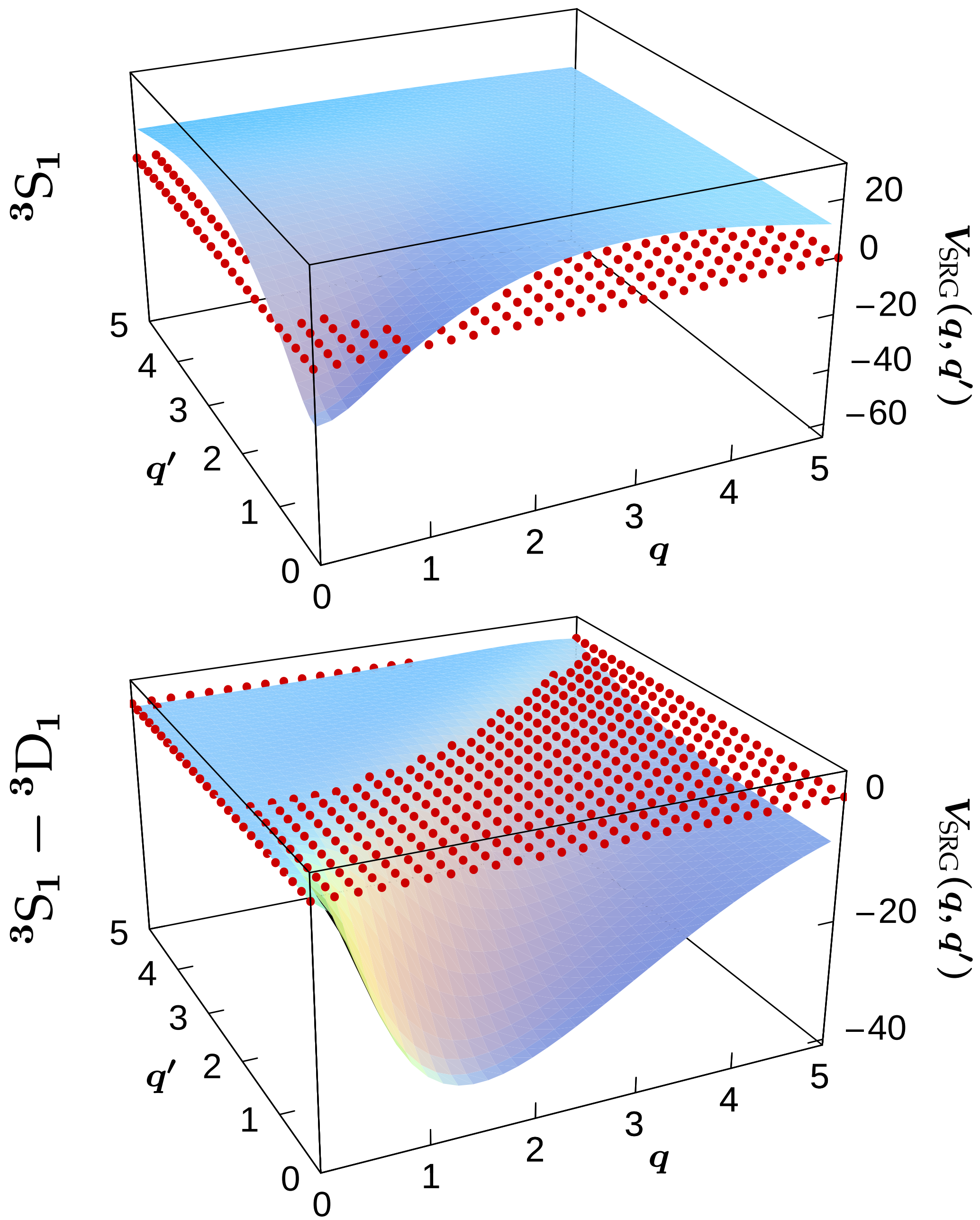
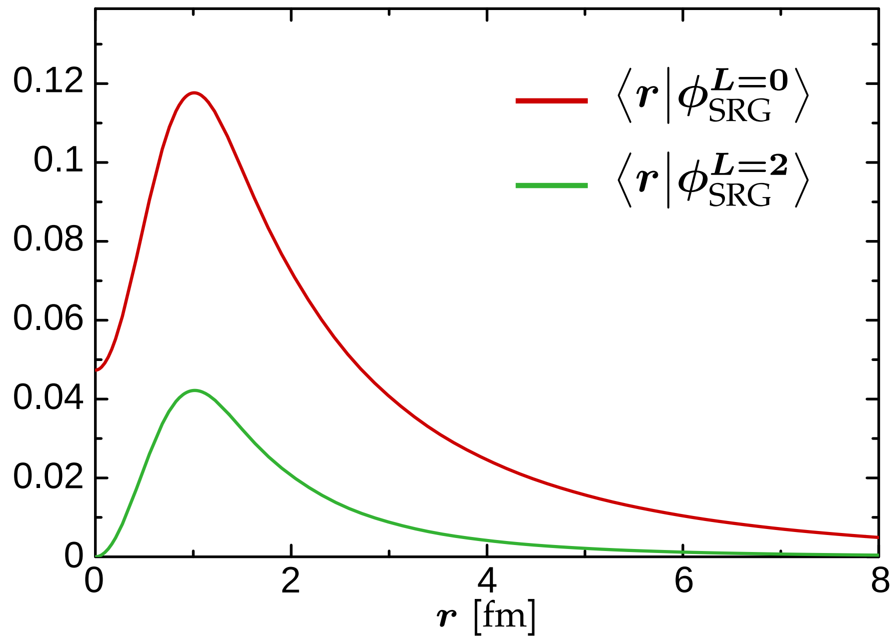


$$\alpha = 0.0004\,\text{fm}^4$$



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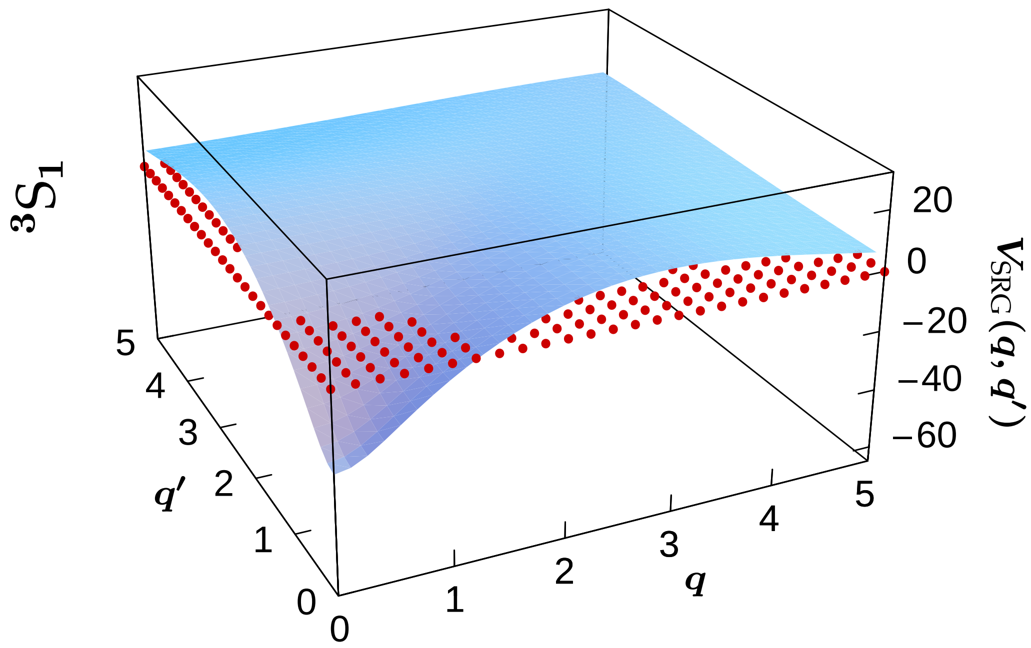

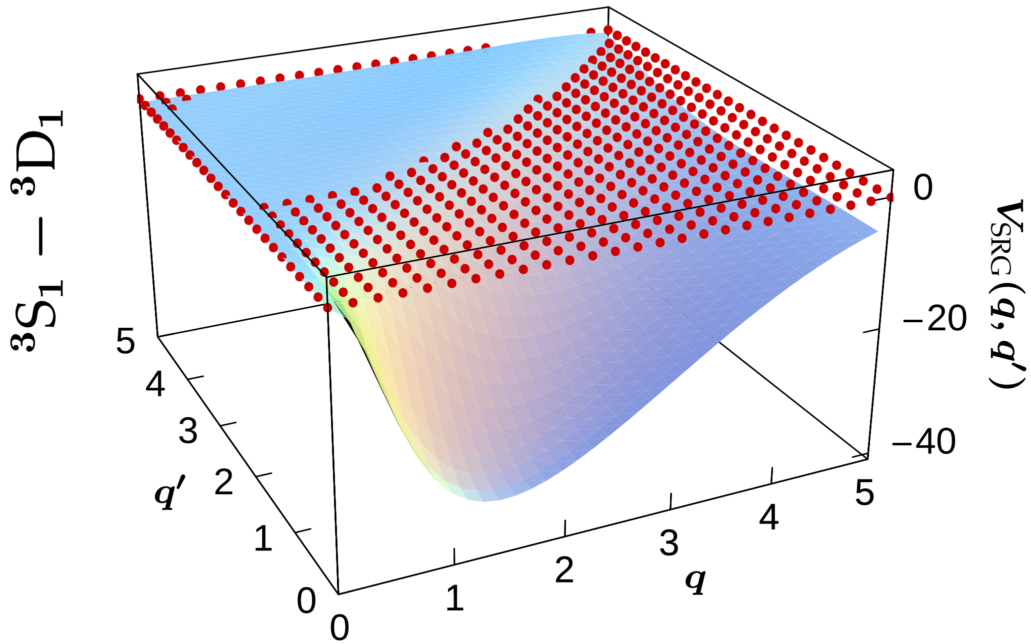

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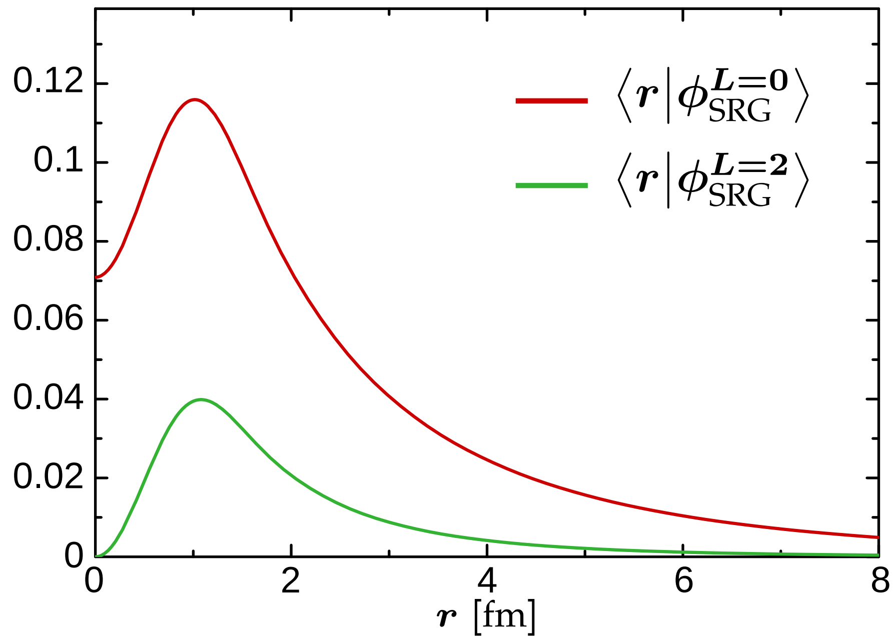
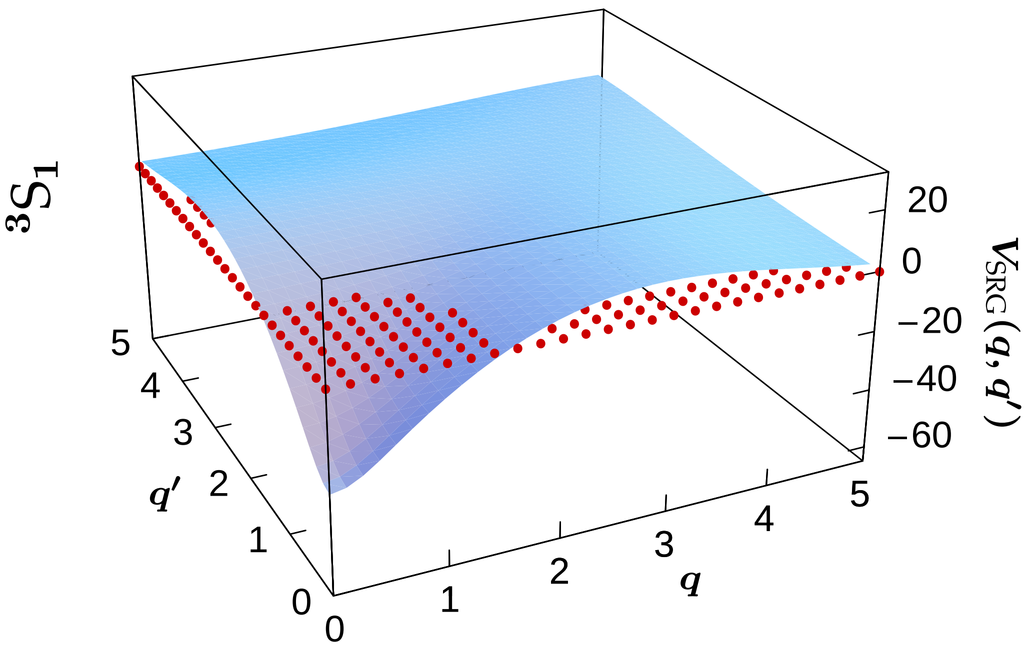
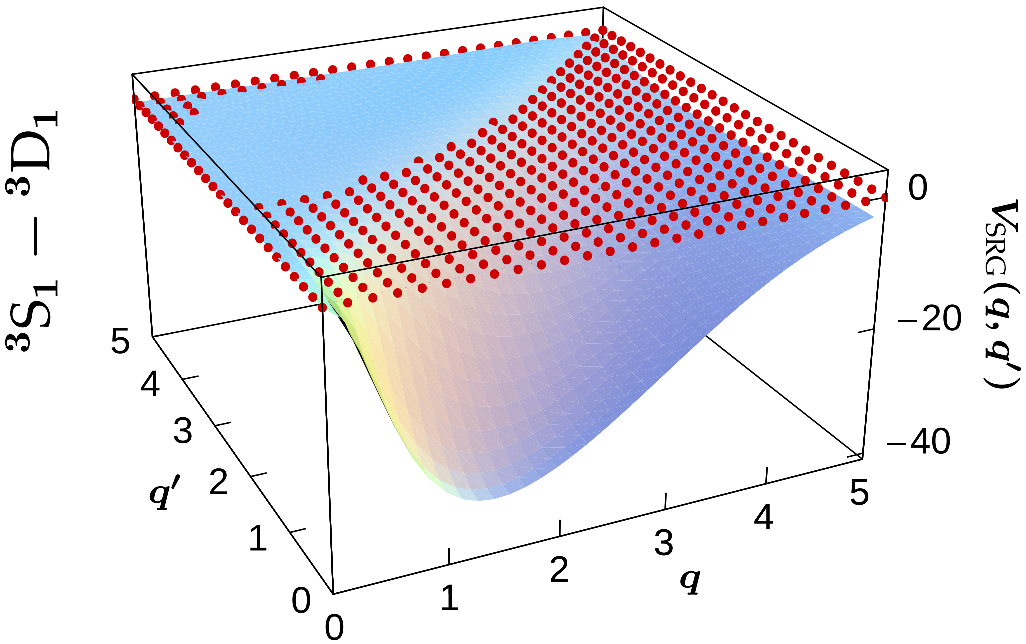


SRG Evolution: The Deuteron





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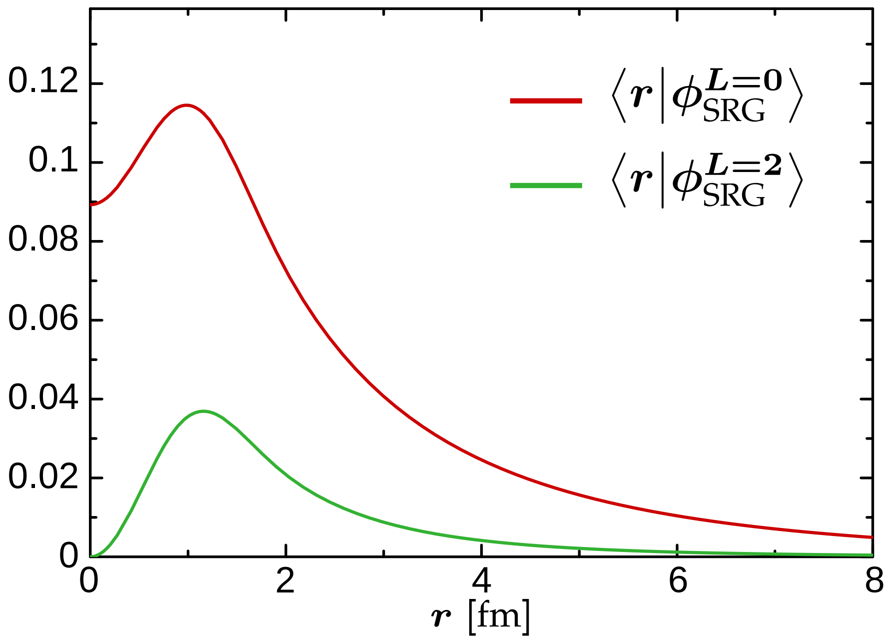
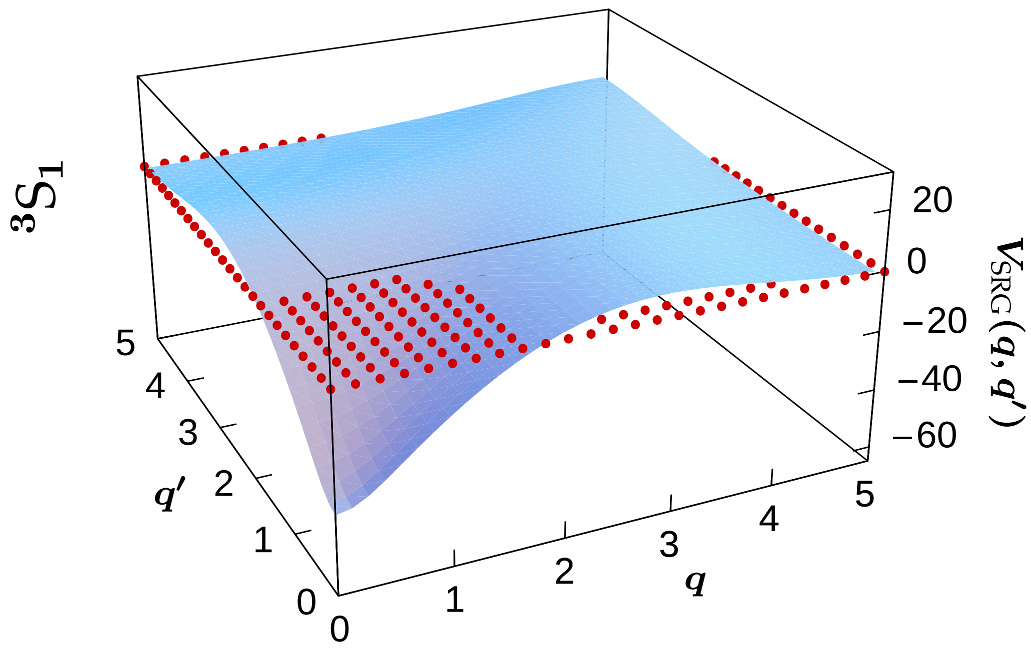
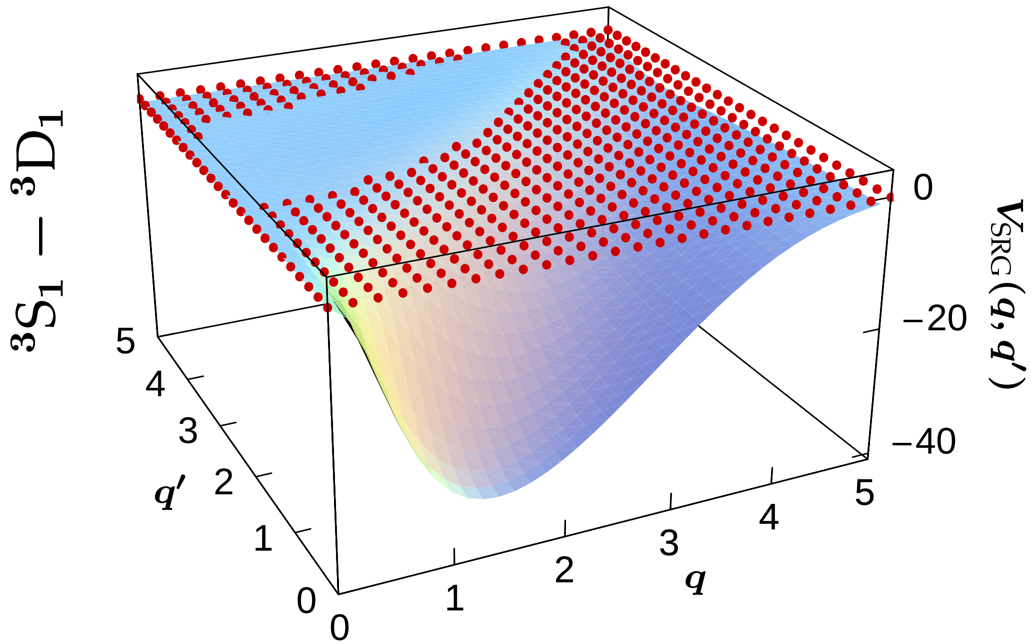
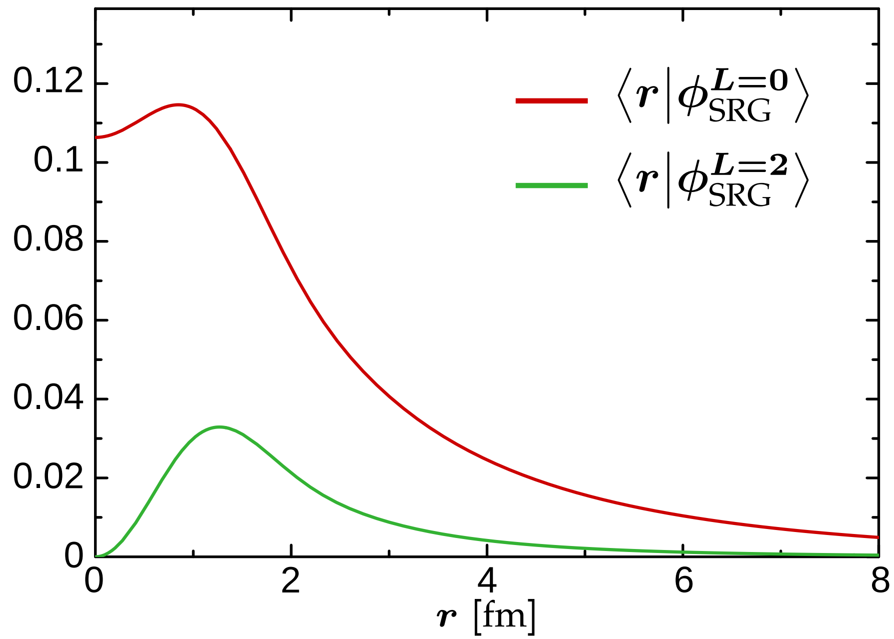


SRG Evolution: The Deuteron

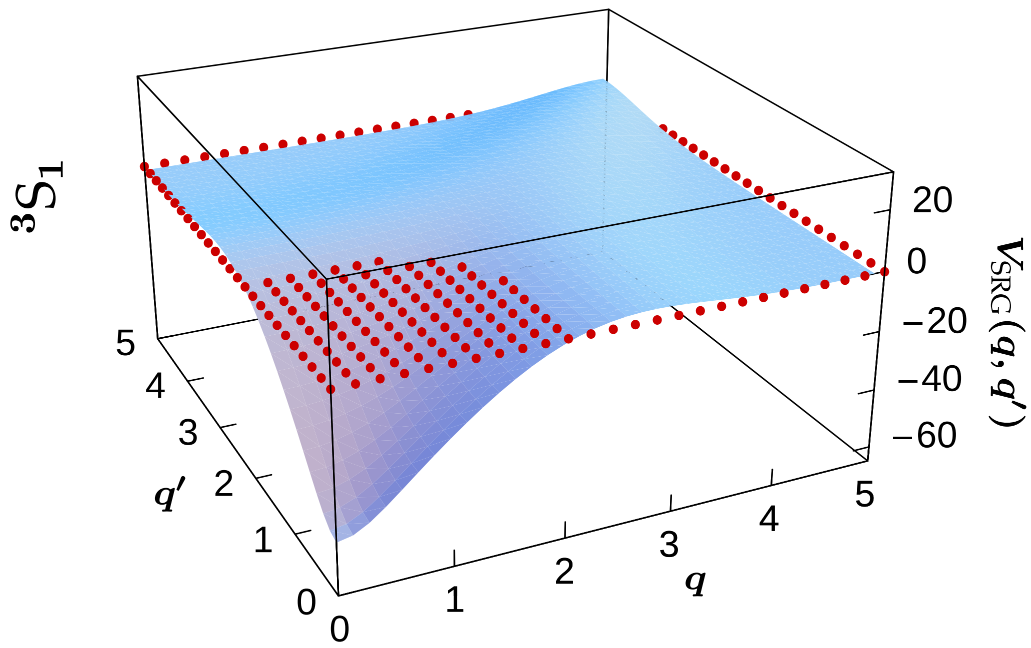


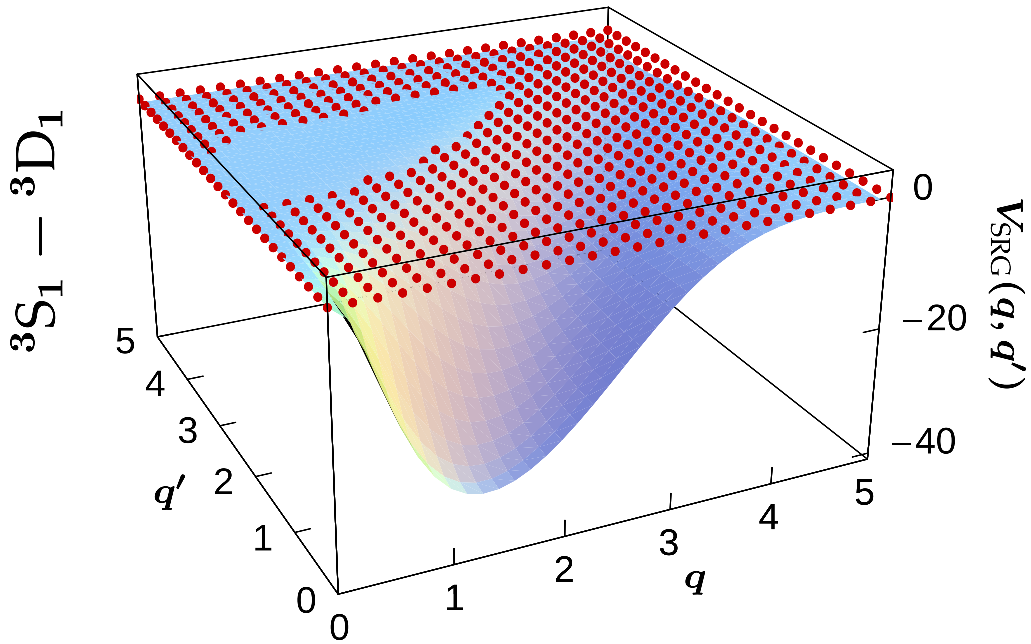


$\alpha = 0.0040\,\mathrm{fm}^4$

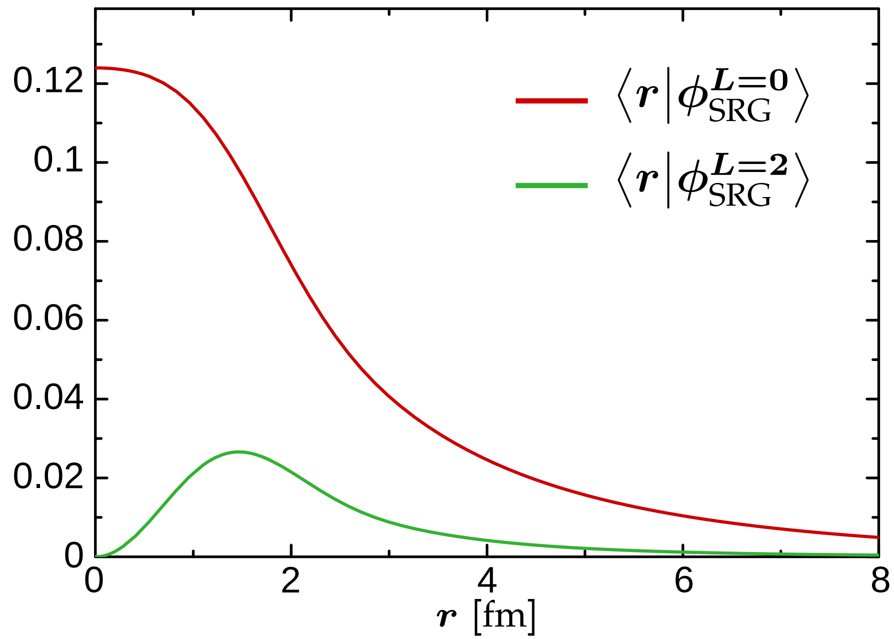


SRG Evolution: The Deuteron

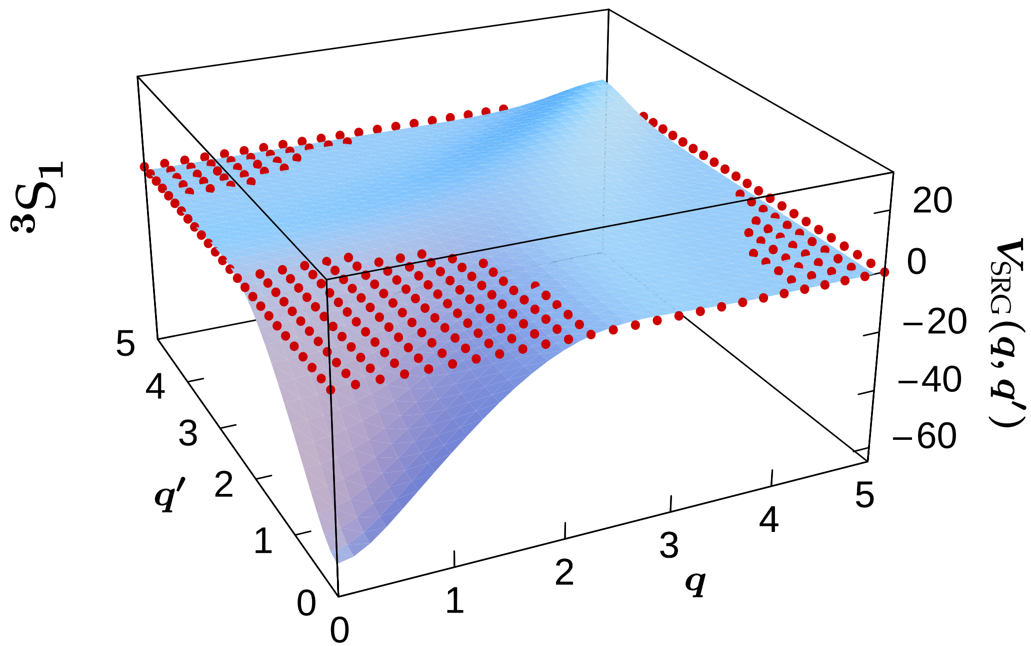


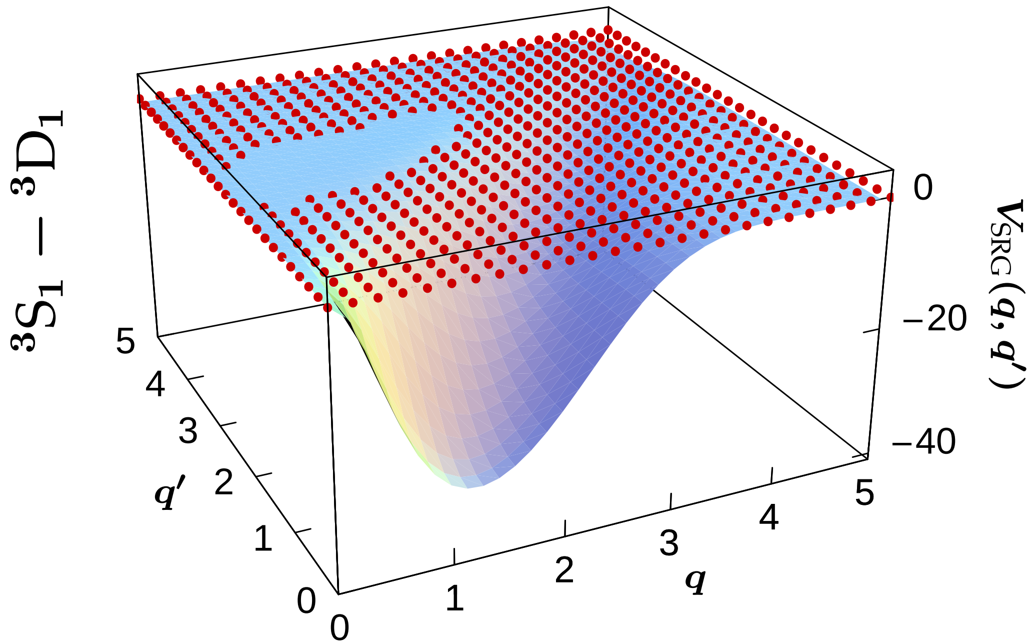


$$\alpha = 0.0100\,\mathrm{fm}^4$$



SRG Evolution: The Deuteron





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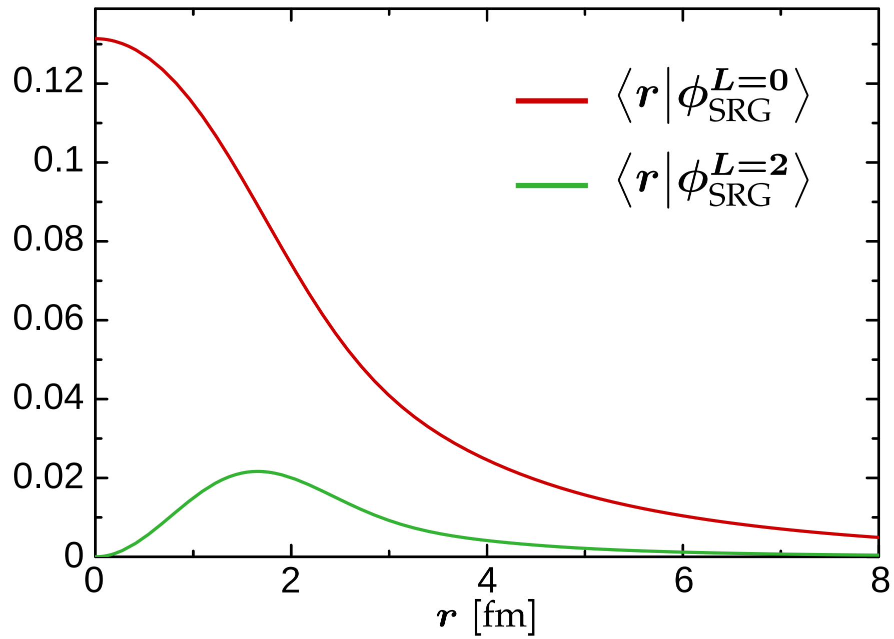
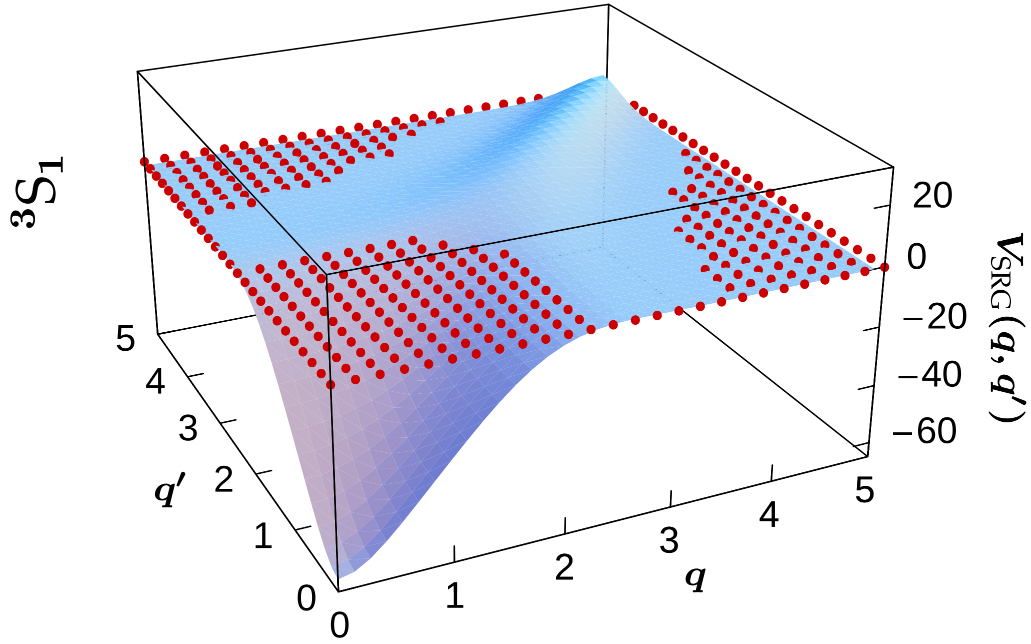
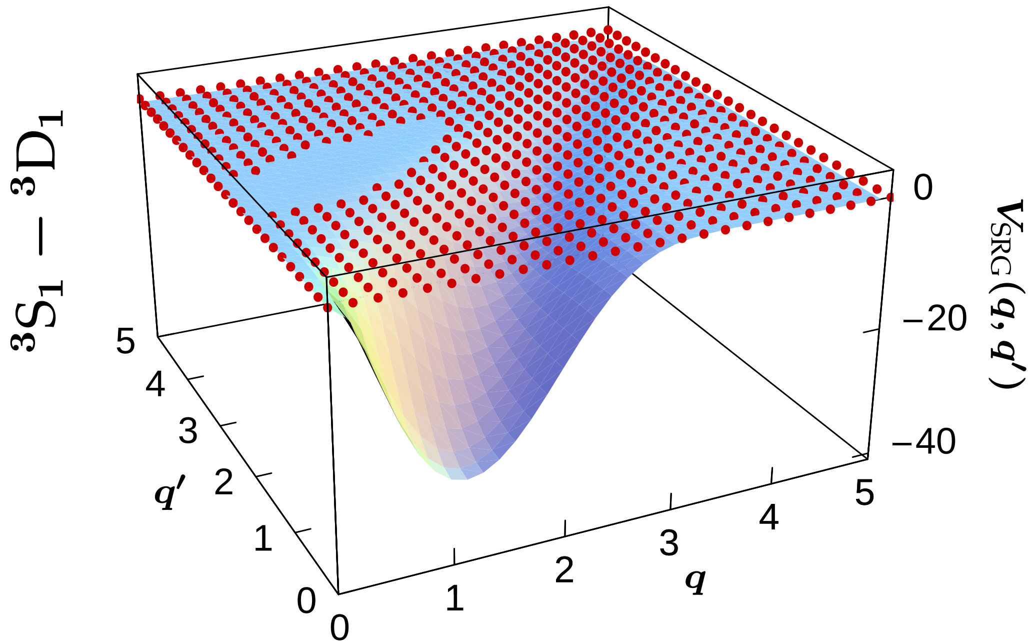
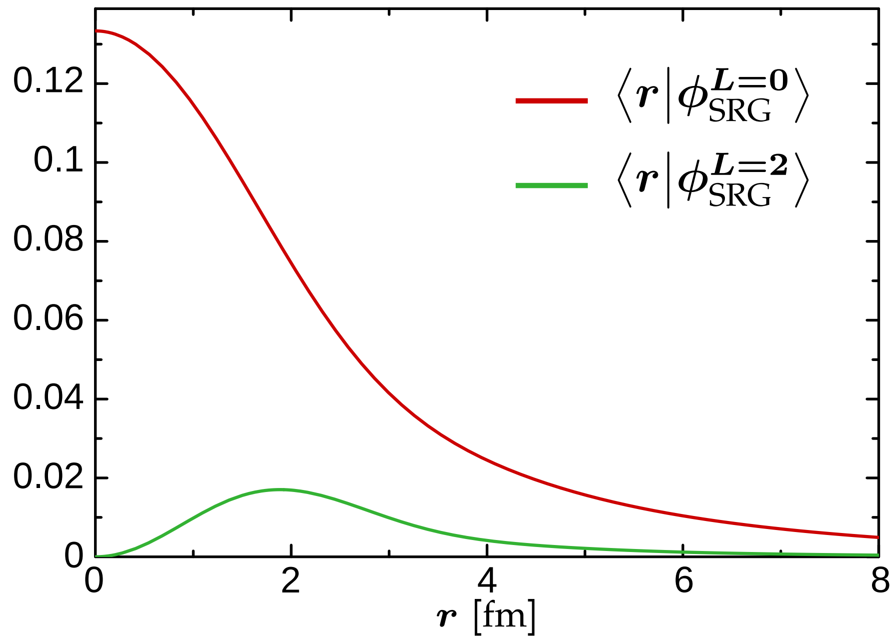


SRG Evolution: The Deuteron

$\alpha = 0.0400\,\text{fm}^4$







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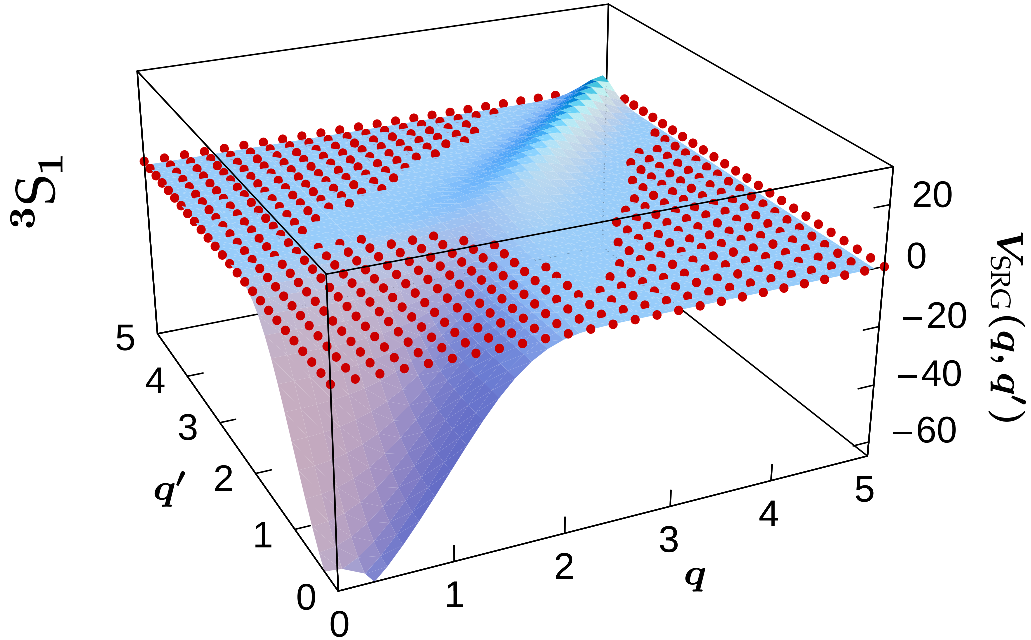
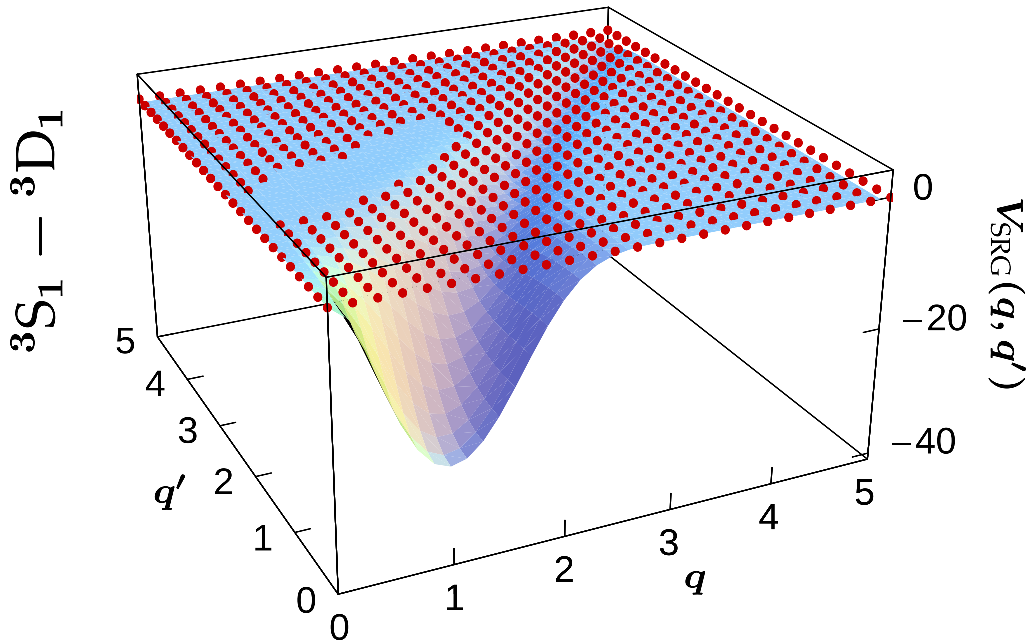
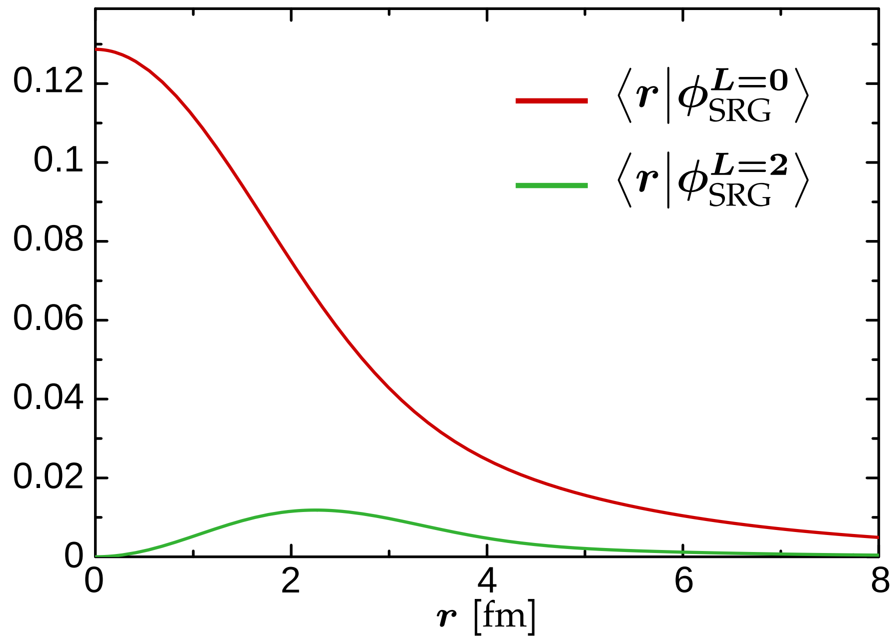
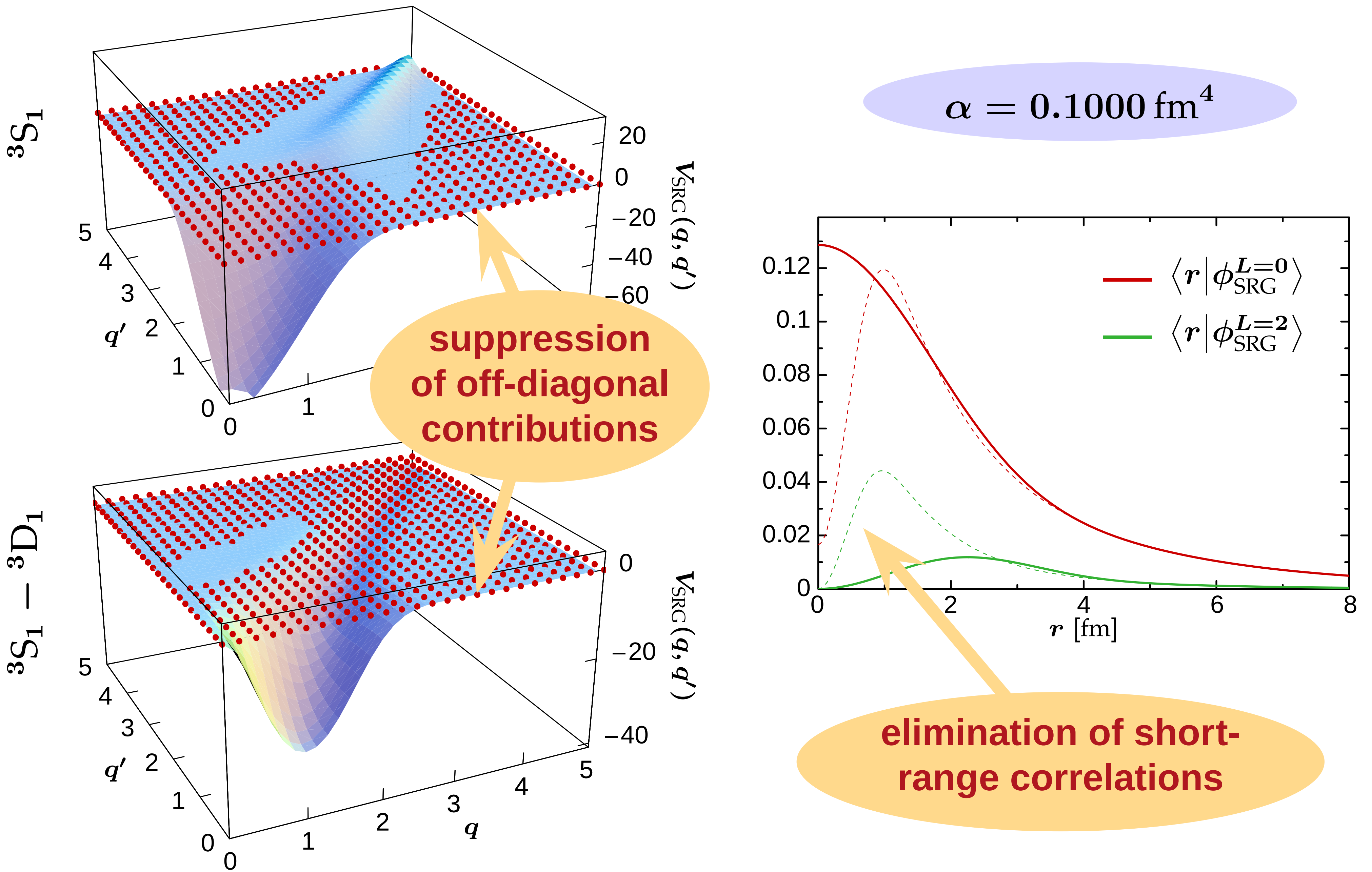




$\alpha = 0.1000\,\mathrm{fm}^4$



SRG Evolution: The Deuteron



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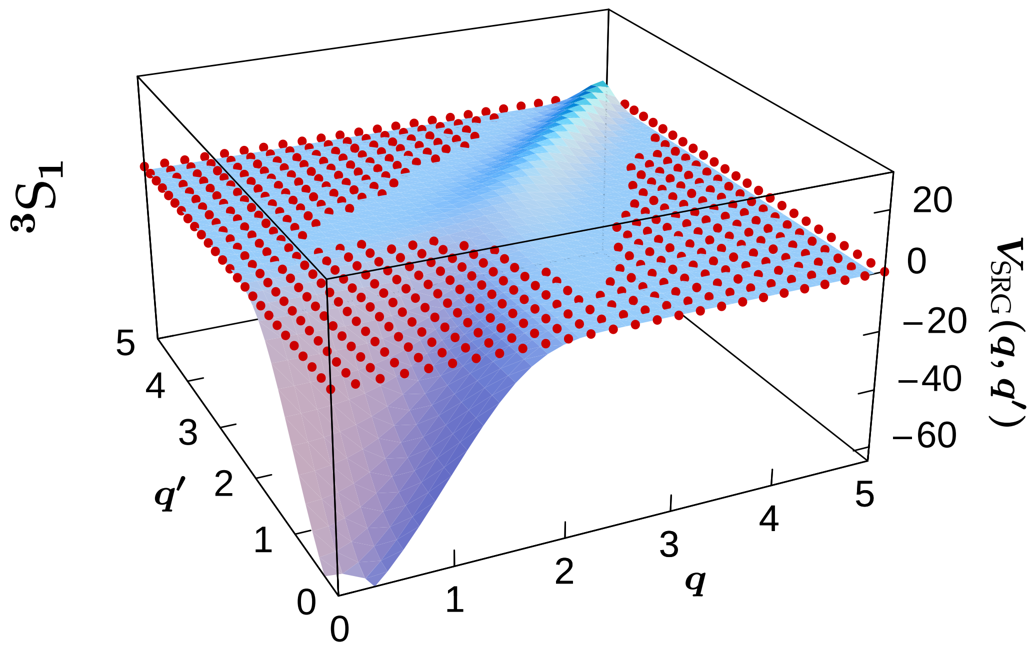

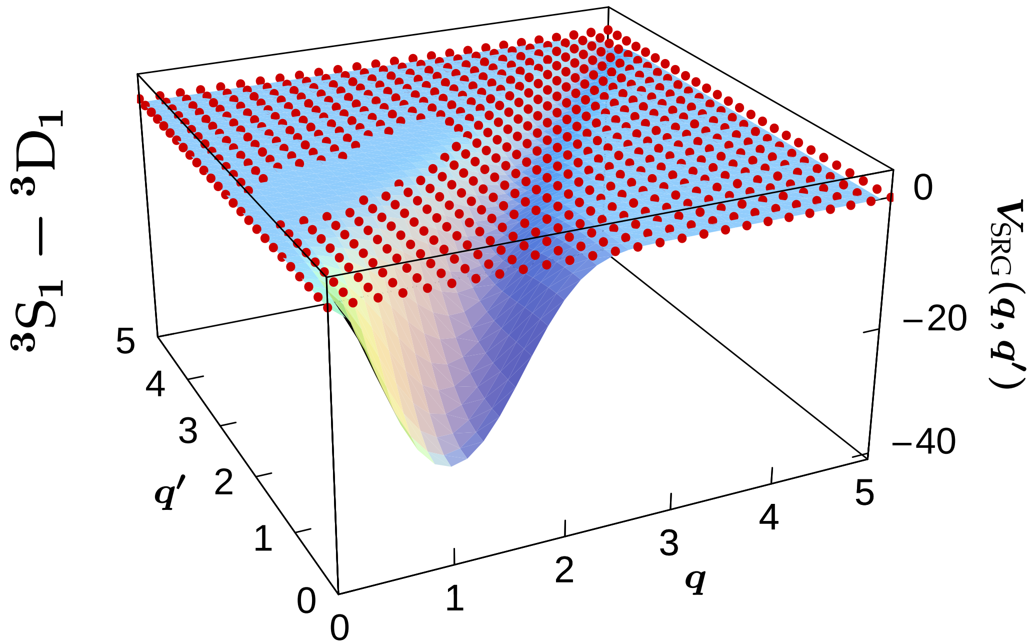

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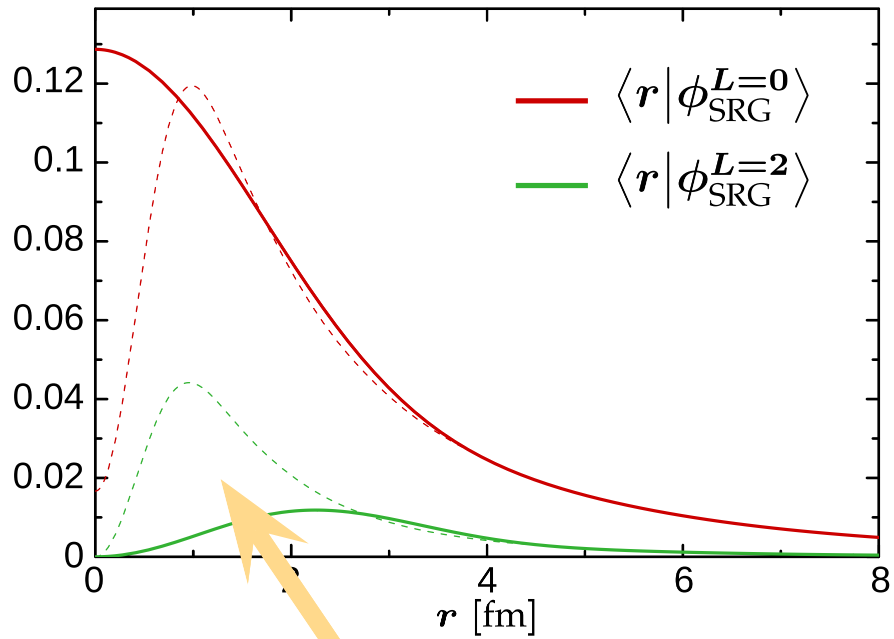
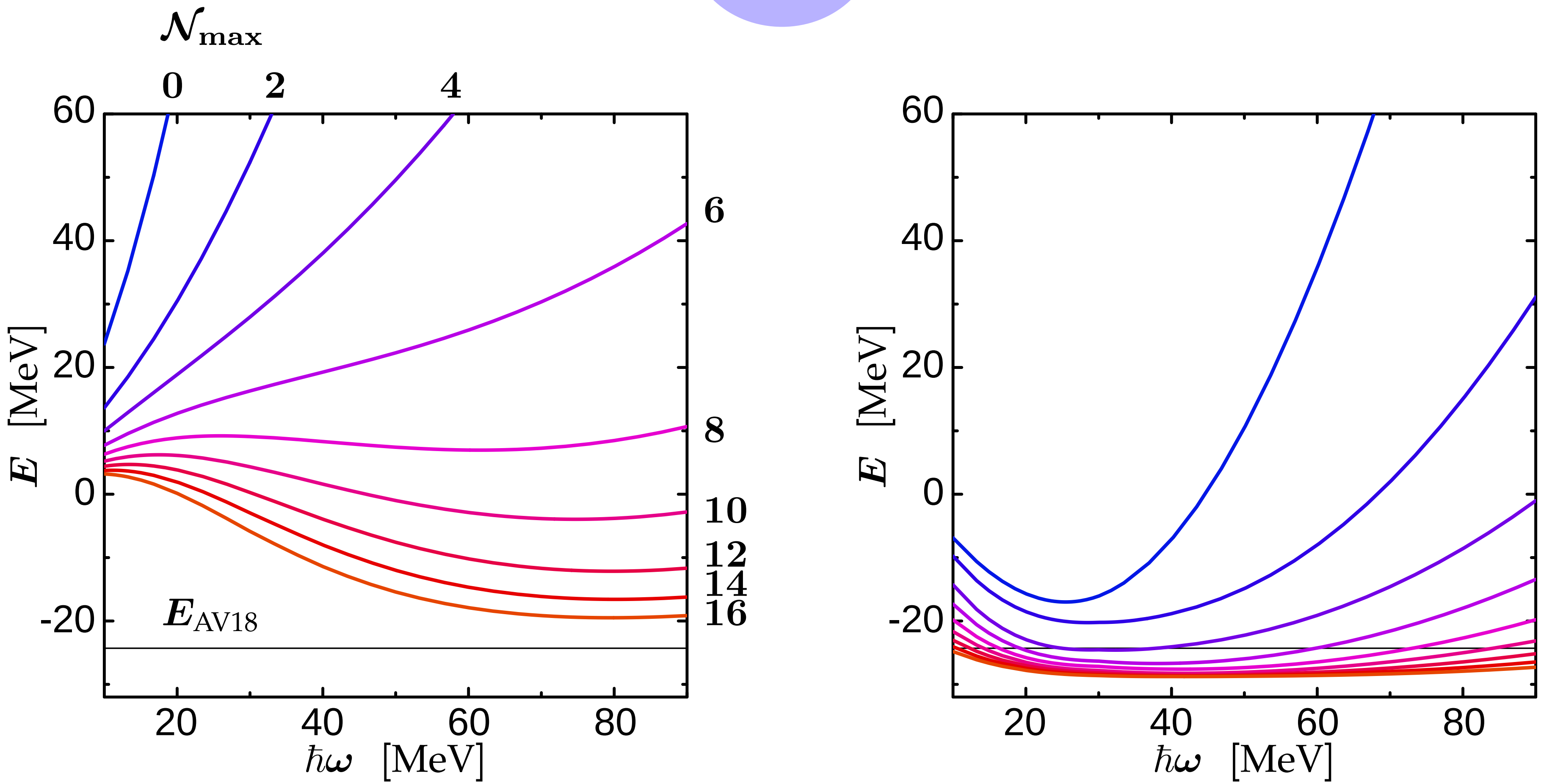


extract UCOM correlation functions $s(r)$ and $\vartheta(r)$

Few-Body Systems

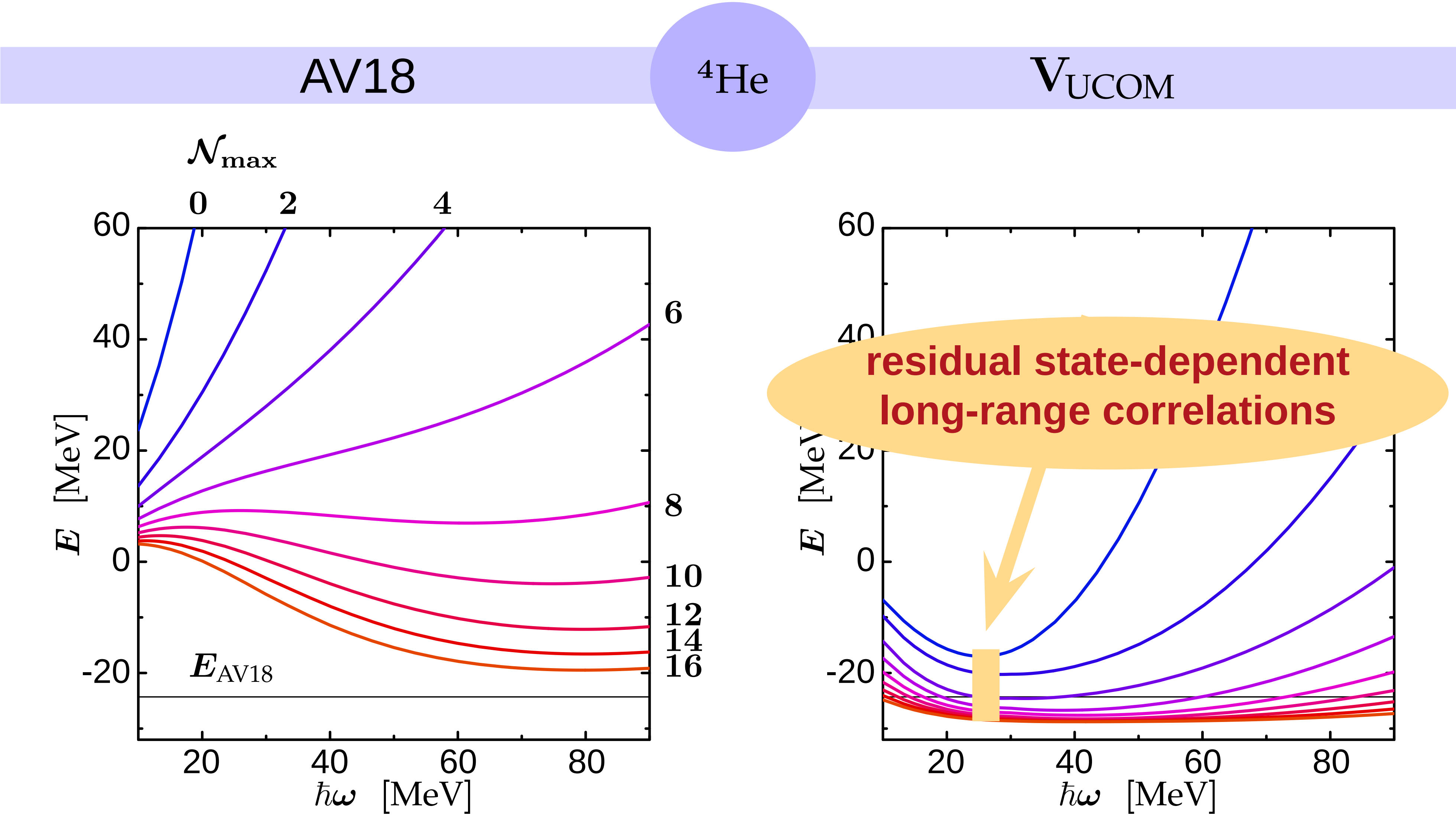
^{4}He: Convergence

AV18 | ^{4}He | V_{UCOM}



NCSM code by P. Navrátil [PRC 61, 044001 (2000)]

^{4}He: Convergence



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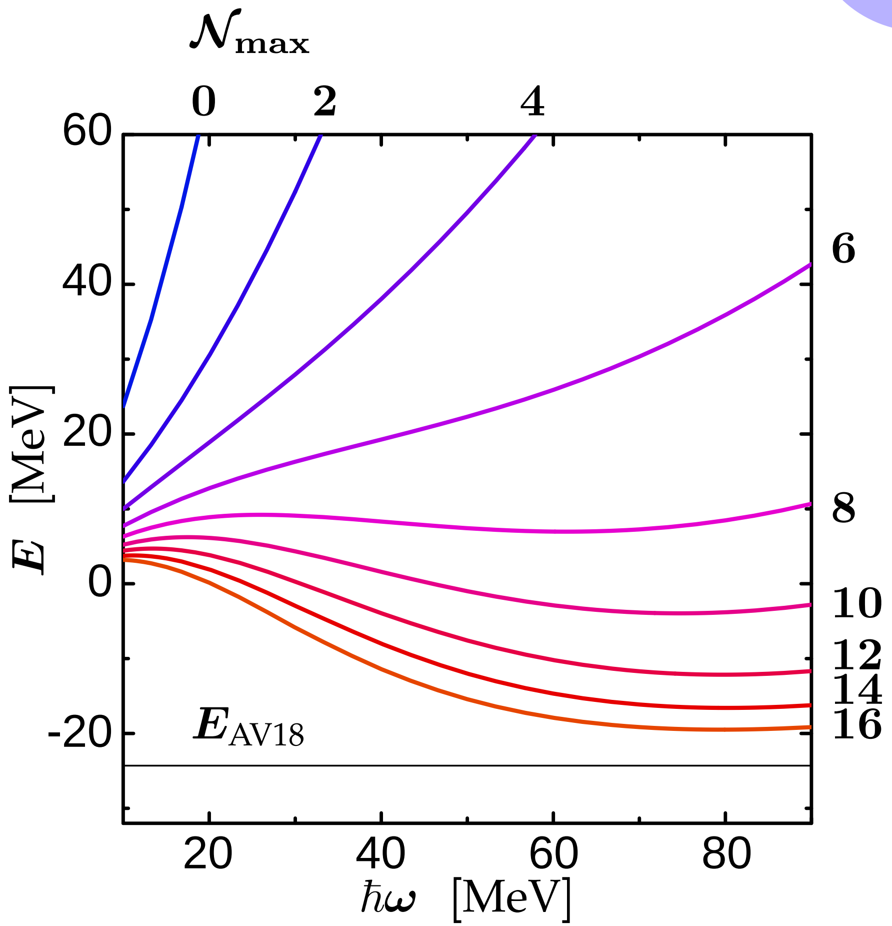
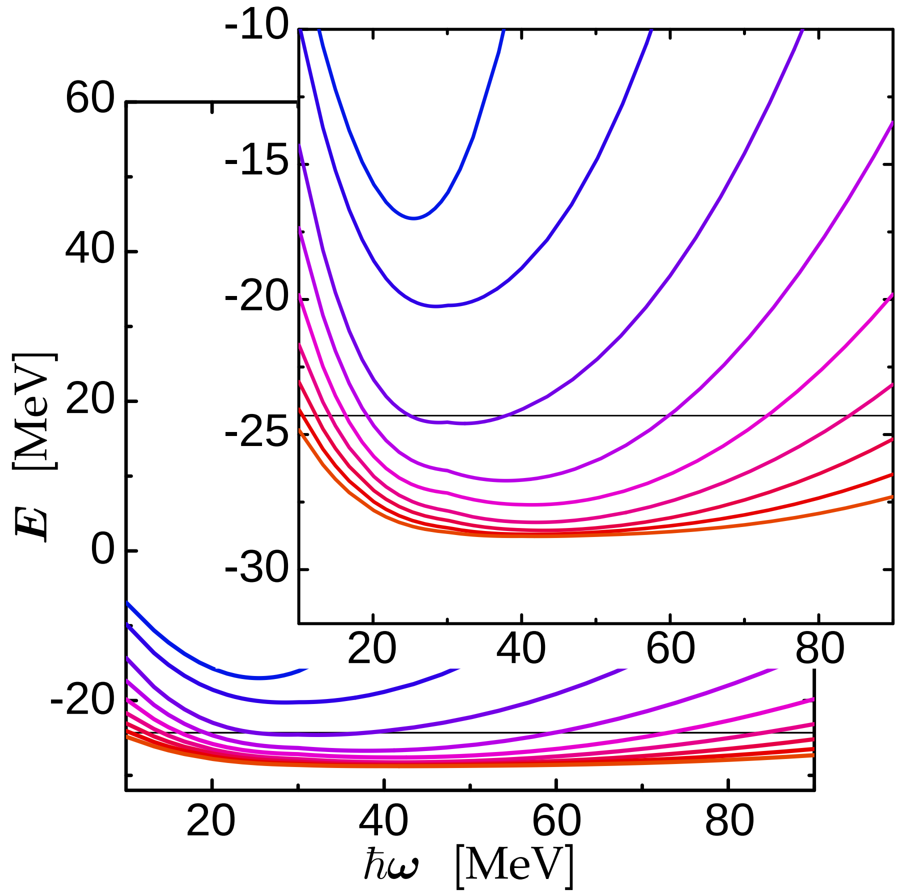
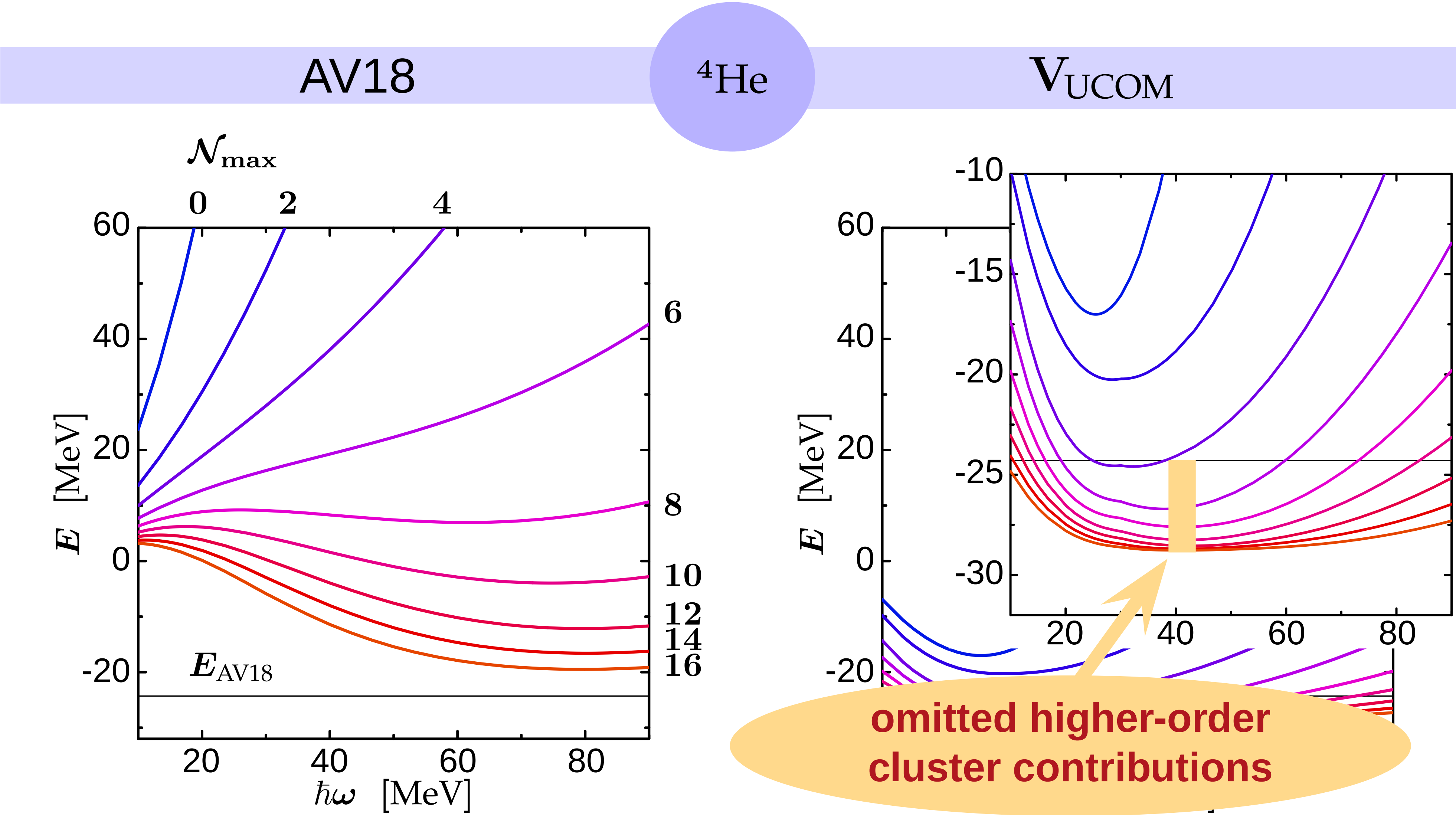
AV18 | ^{4}He | V_{UCOM}





NCSM code by P. Navrátil [PRC 61, 044001 (2000)]

^{4}He: Convergence



NCSM code by P. Navrátil [PRC 61, 044001 (2000)]

Three-Body Interactions — Strategies

Correlated Hamiltonian in Many-Body Space

$$\begin{aligned}
\widetilde{\mathrm{H}} &= \mathrm{C}^\dagger\,(\mathrm{T} + \mathrm{V}_{NN} + \mathrm{V}_{3N})\,\mathrm{C} \\
&= \widetilde{\mathrm{T}}^{[1]} + (\widetilde{\mathrm{T}}^{[2]} + \widetilde{\mathrm{V}}^{[2]}_{NN}) + (\widetilde{\mathrm{T}}^{[3]} + \widetilde{\mathrm{V}}^{[3]}_{NN} + \widetilde{\mathrm{V}}^{[3]}_{3N}) + \cdots \\
&= \mathrm{T} + \mathrm{V}_{\mathrm{UCOM}} + \mathrm{V}^{[3]}_{\mathrm{UCOM}} + \cdots
\end{aligned}$$

Three-Body Interactions — Strategies

Correlated Hamiltonian in Many-Body Space

$$\widetilde{\mathrm{H}} = \mathrm{C}^\dagger\,(\mathrm{T} + \mathrm{V}_{NN} + \mathrm{V}_{3N})\,\mathrm{C}$$

$$= \widetilde{\mathrm{T}}^{[1]} + (\widetilde{\mathrm{T}}^{[2]} + \widetilde{\mathrm{V}}^{[2]}_{NN}) + (\widetilde{\mathrm{T}}^{[3]} + \widetilde{\mathrm{V}}^{[3]}_{NN} + \widetilde{\mathrm{V}}^{[3]}_{3N}) + \cdots$$

$$= \mathrm{T} + \mathrm{V}_{\mathrm{UCOM}} + \mathrm{V}^{[3]}_{\mathrm{UCOM}} + \cdots$$

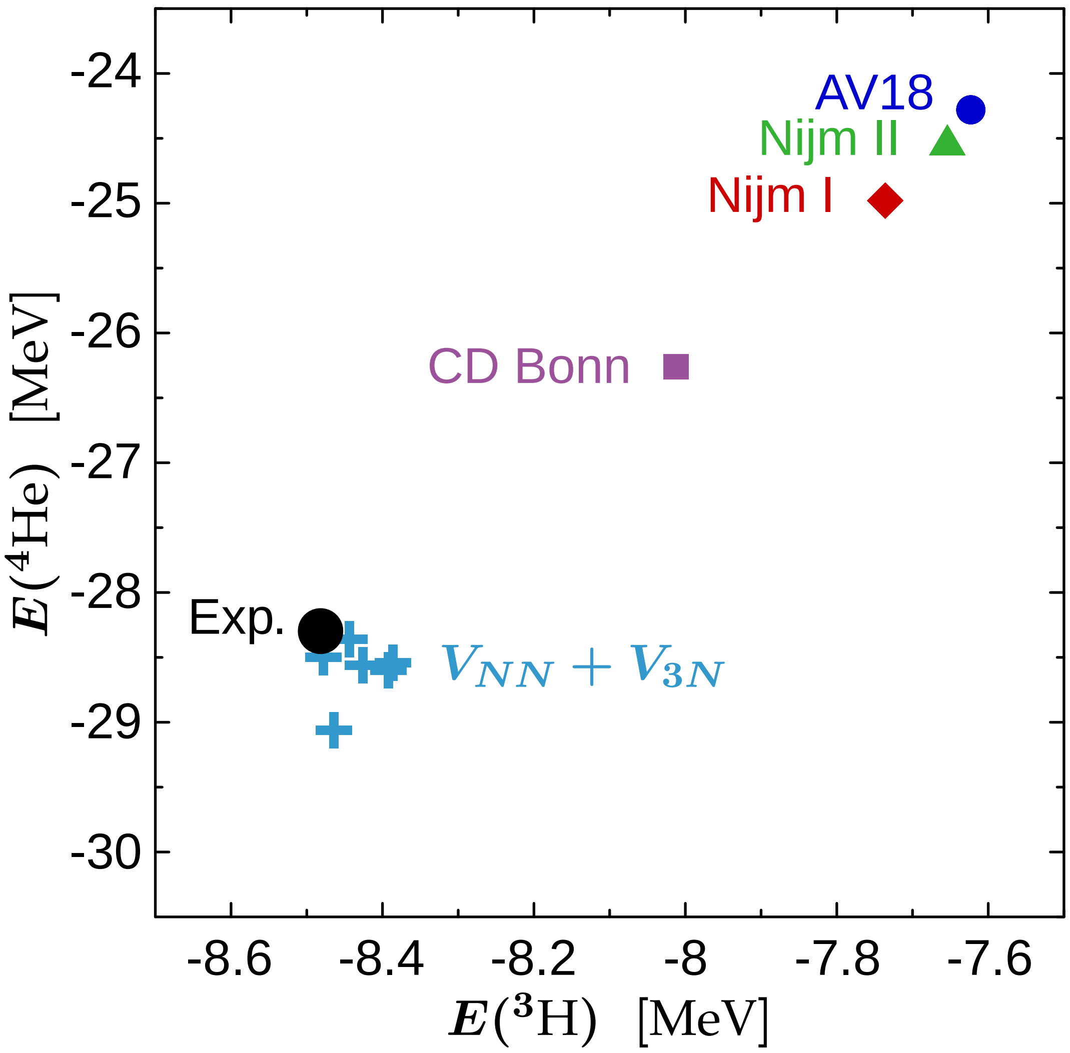
- strategies for treating the three-body contributions:

❶ **include full $\mathrm{V}^{[3]}_{\mathrm{UCOM}}$** consisting of genuine and induced 3N terms

❷ **replace $\mathrm{V}^{[3]}_{\mathrm{UCOM}}$** by “phenomenological” three-body force

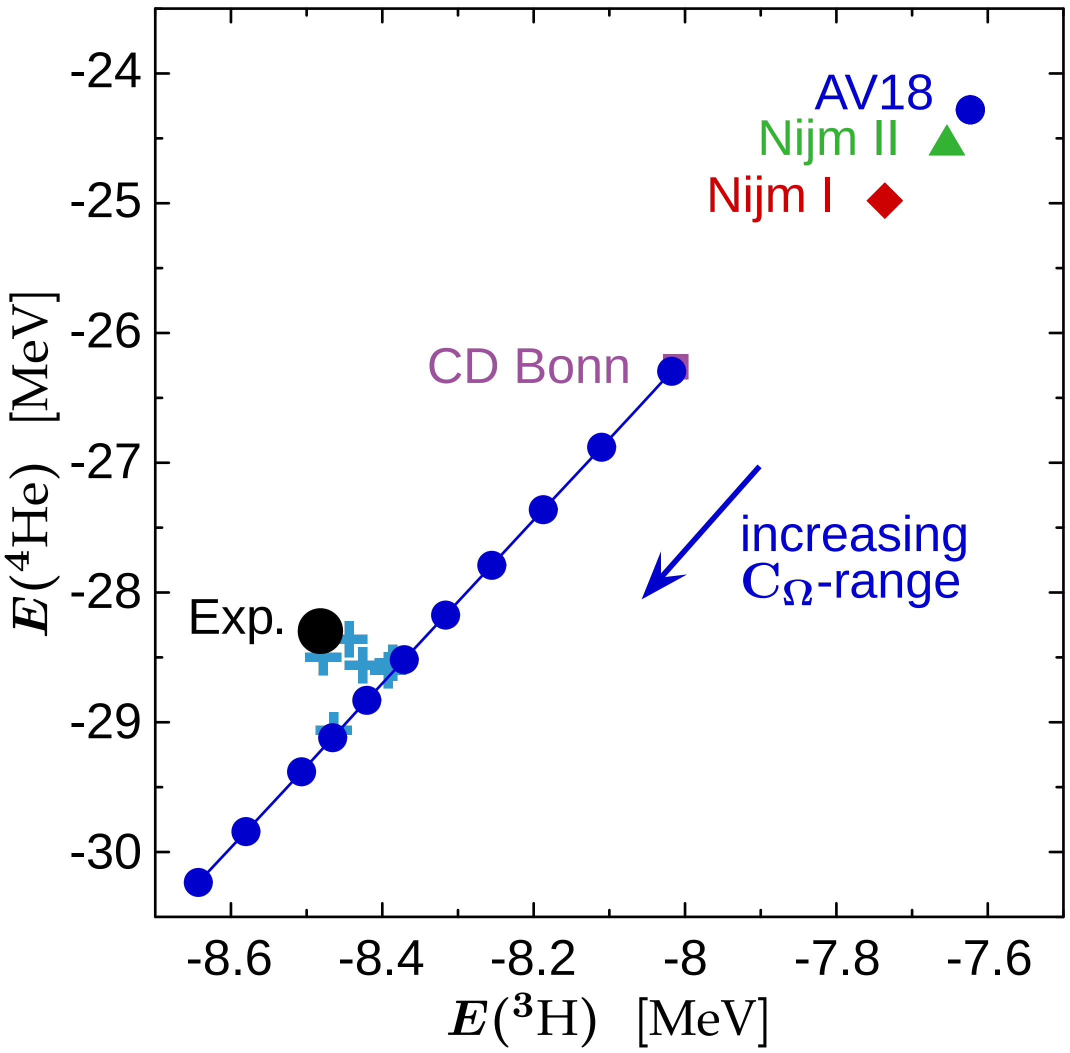
❸ **minimize $\mathrm{V}^{[3]}_{\mathrm{UCOM}}$** by proper choice of unitary transformation

Three-Body Interactions — Tjon Line



- **Tjon-line**: $E(^4\mathrm{He})$ vs. $E(^3\mathrm{H})$ for phase-shift equivalent NN-interactions

Three-Body Interactions — Tjon Line



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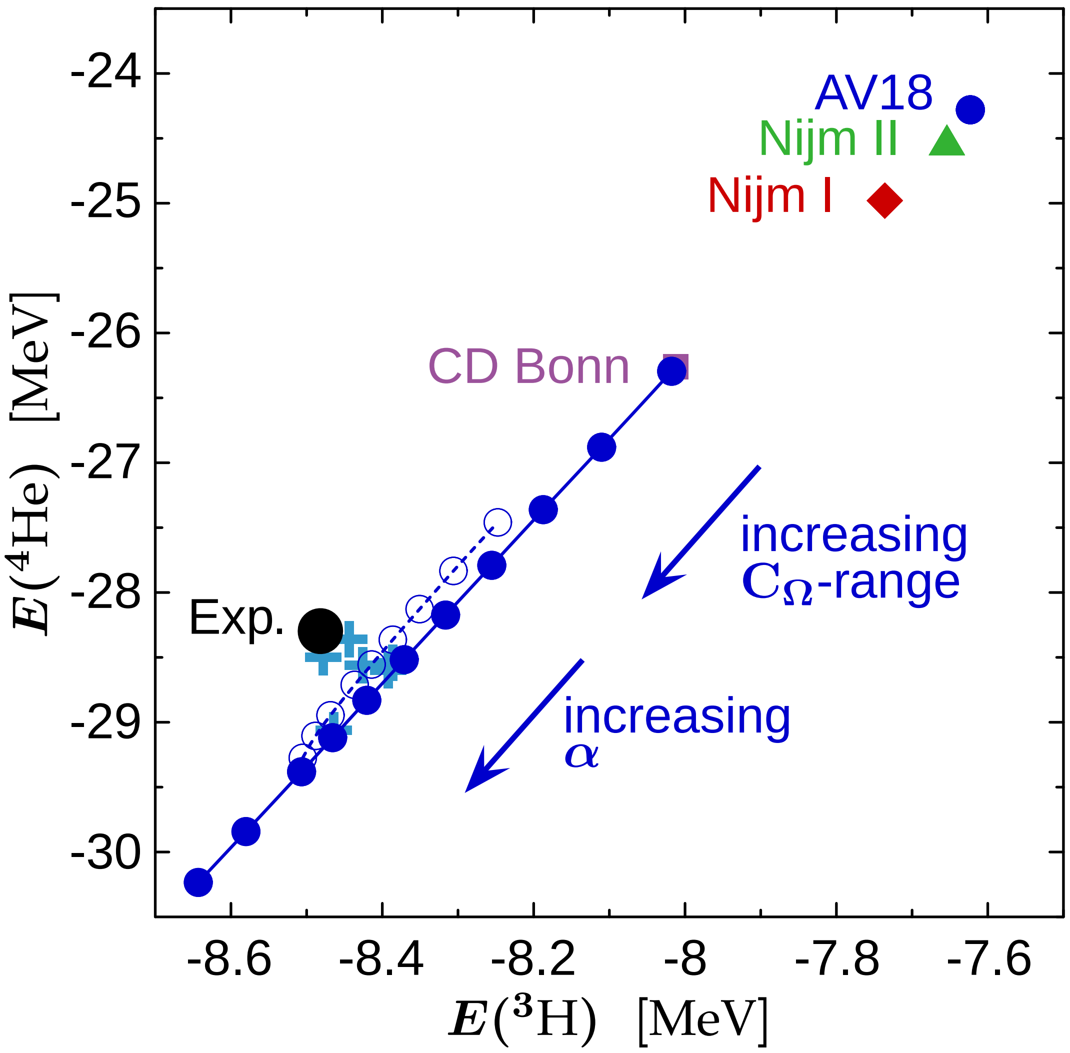

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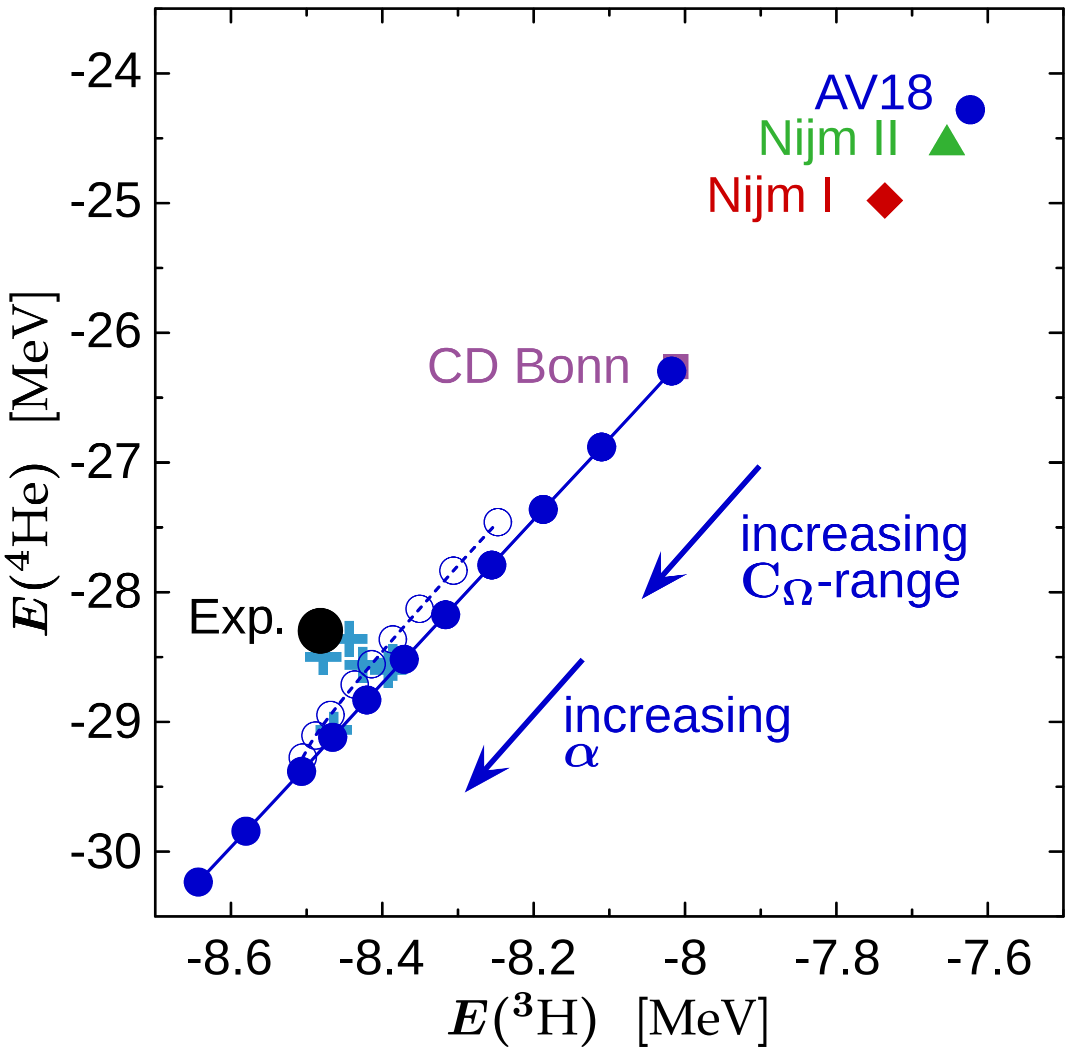
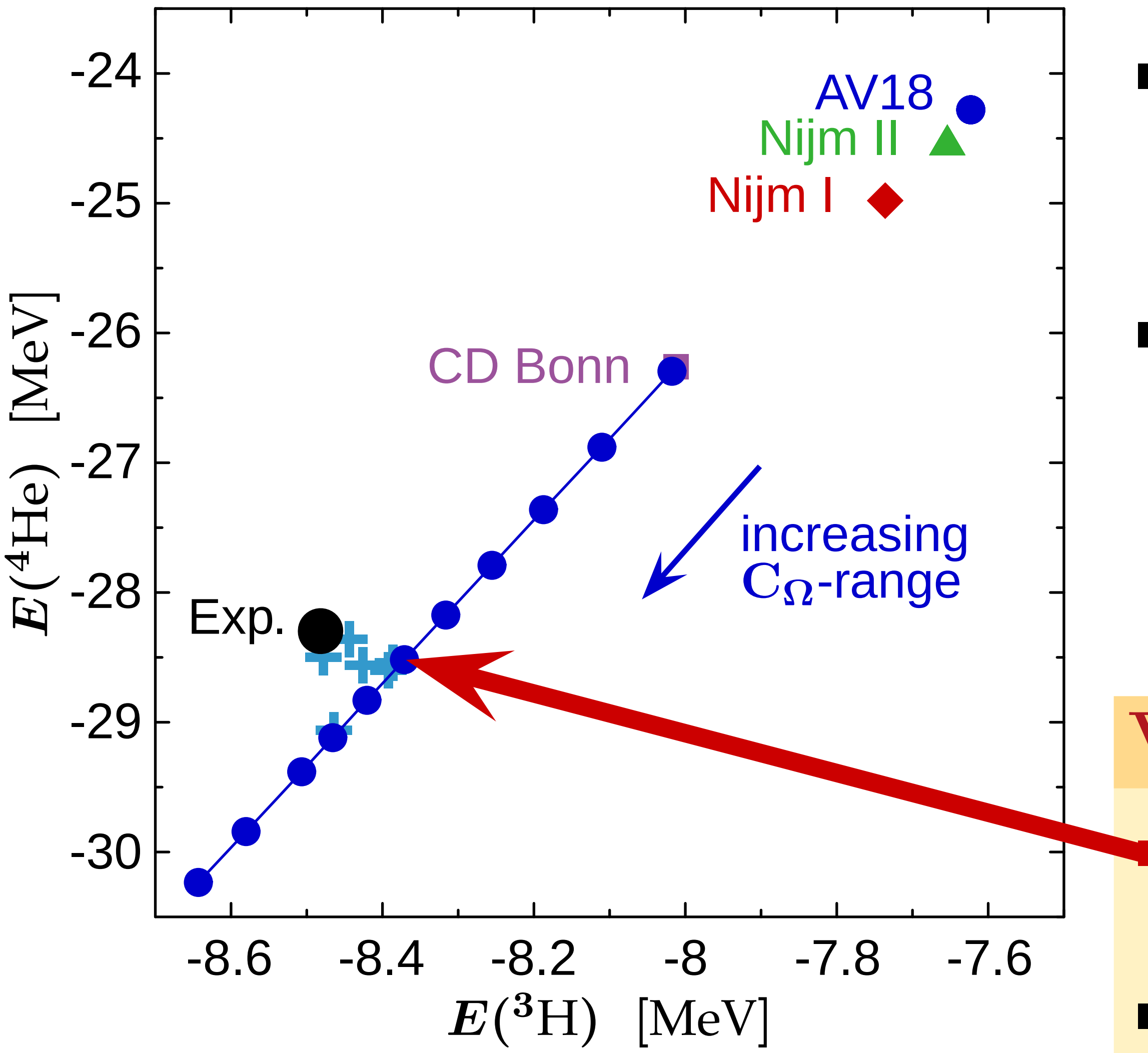


- **Tjon-line**: $E(^4\text{He})$ vs. $E(^3\text{H})$ for phase-shift equivalent NN-interactions

- use α / range of C_Ω to
  - test dependence of V_α or V_{UCOM}
  - tune contributions of **net 3N force**

Three-Body Interactions — Tjon Line



- **Tjon-line**: $E(^4\mathrm{He})$ vs. $E(^3\mathrm{H})$ for phase-shift equivalent NN-interactions

- use α / range of C_Ω to
  - test dependence of V_α or V_{UCOM}
  - tune contributions of **net 3N force**

V_{UCOM}

- **min. net $3N$ force**:

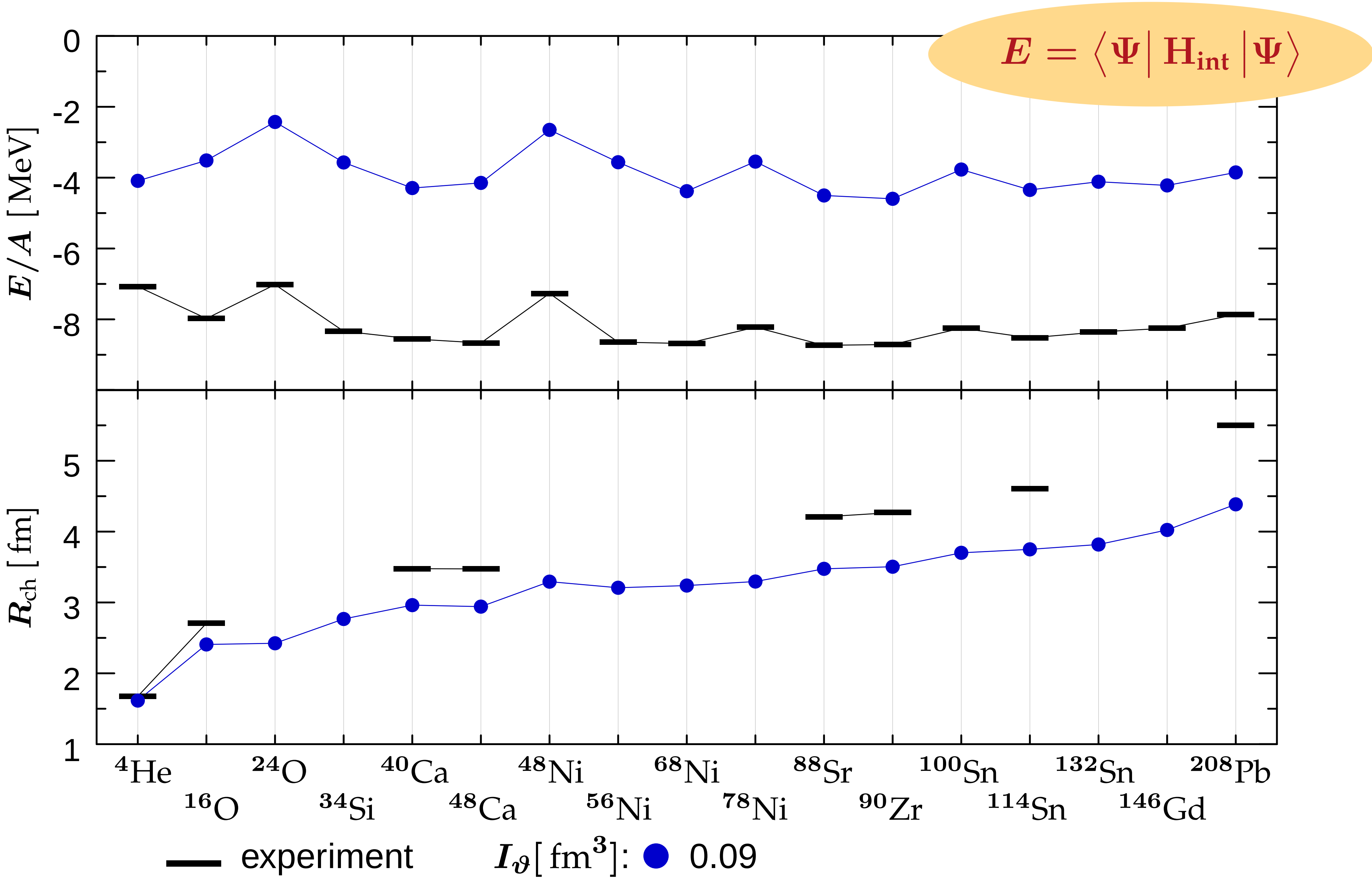
$$I_\vartheta = 0.09\,\mathrm{fm}^3$$

- with **phen. $3N$ force**:

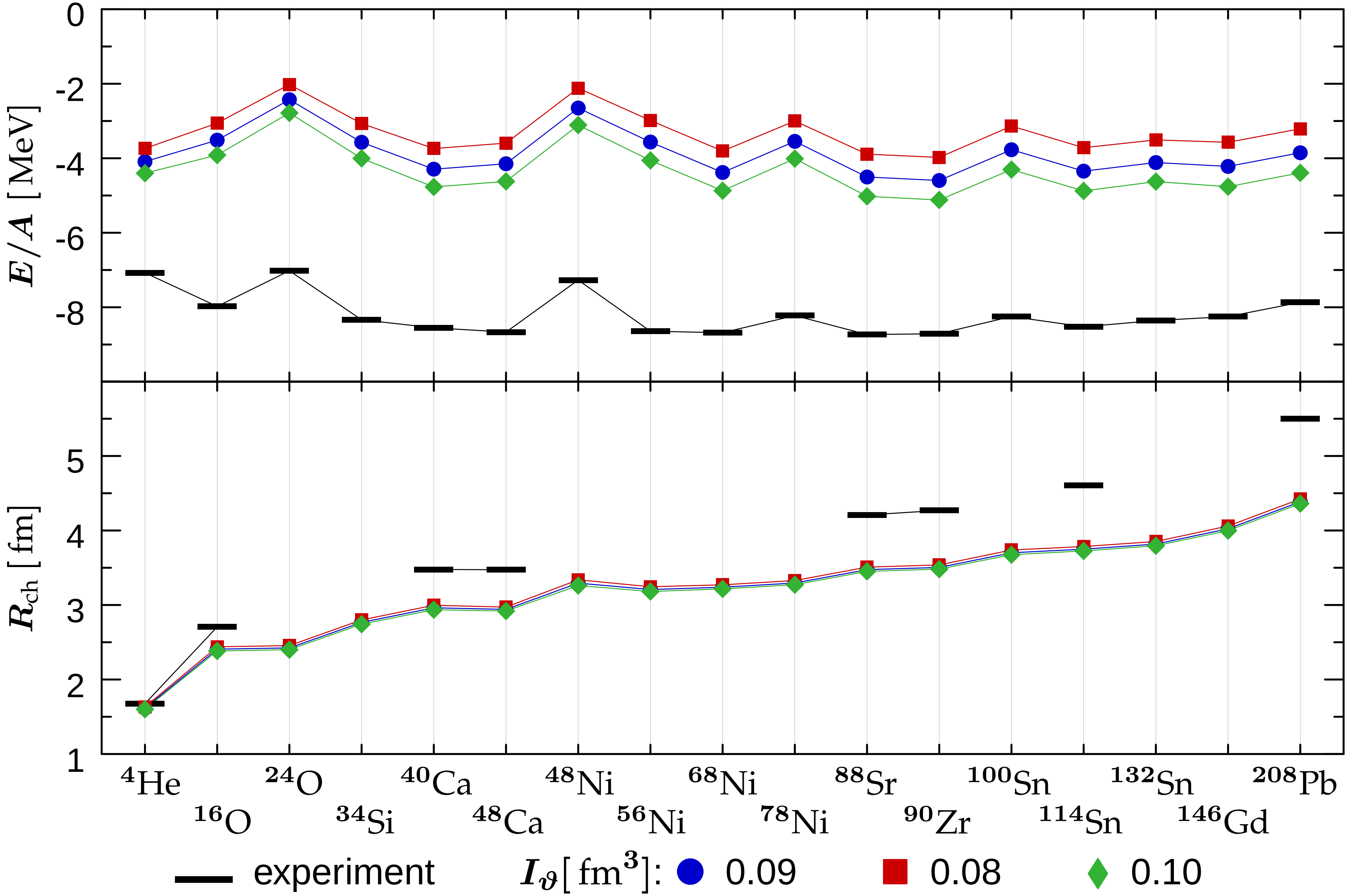
$$I_\vartheta = 0.20\,\mathrm{fm}^3$$

Many-Body Systems

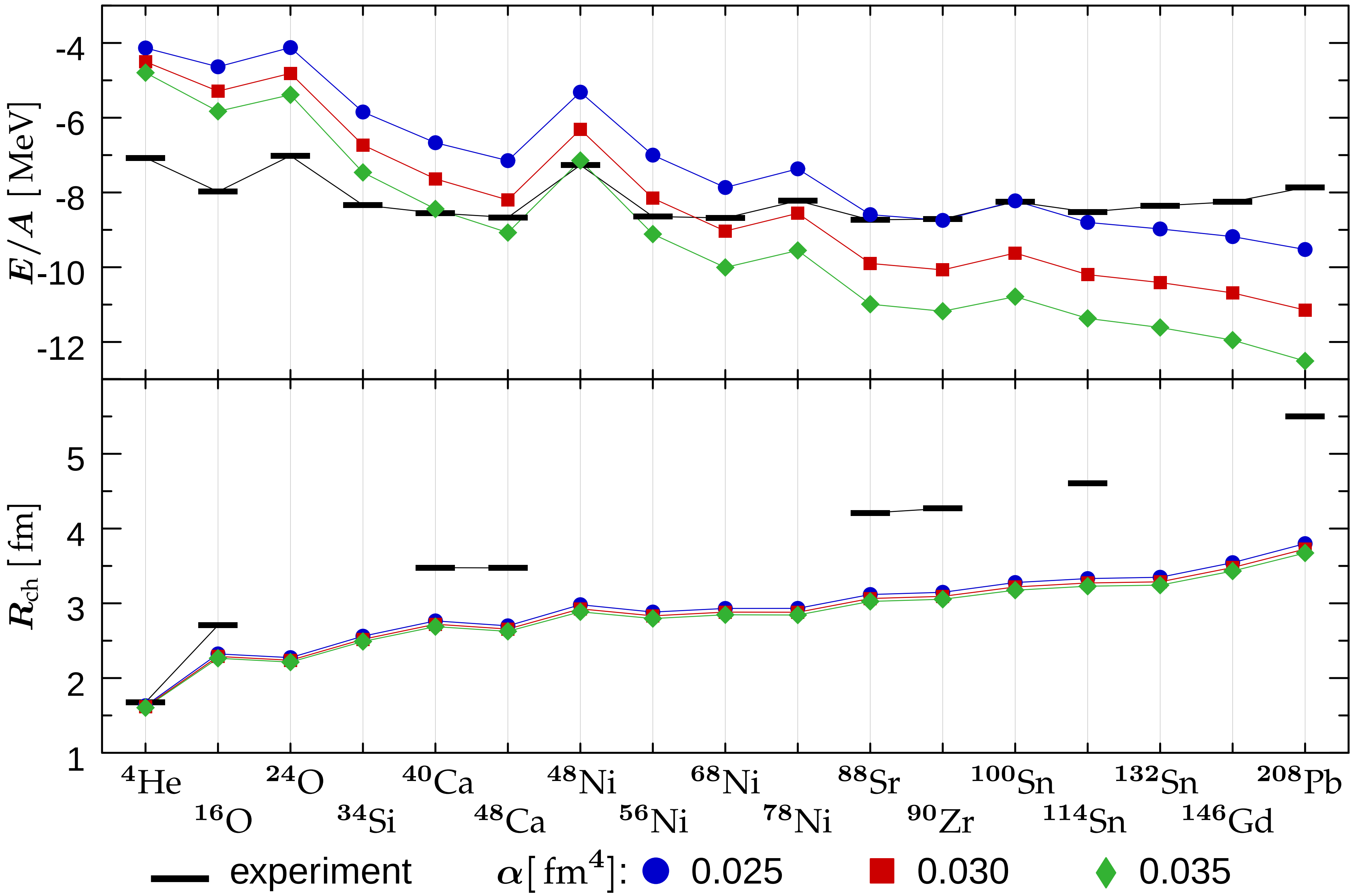
Hartree-Fock with V_{UCOM}



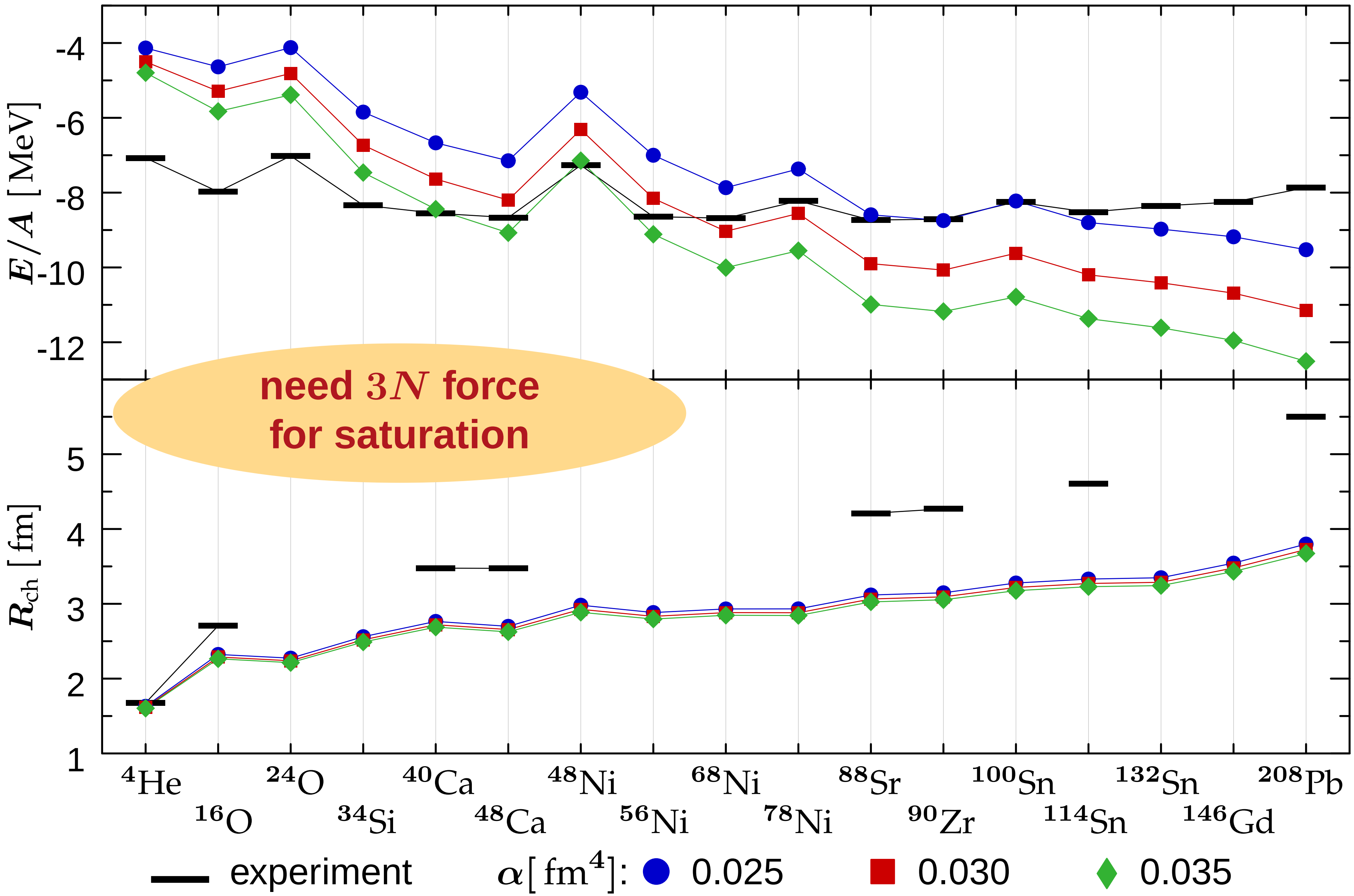
Hartree-Fock with V_{UCOM}



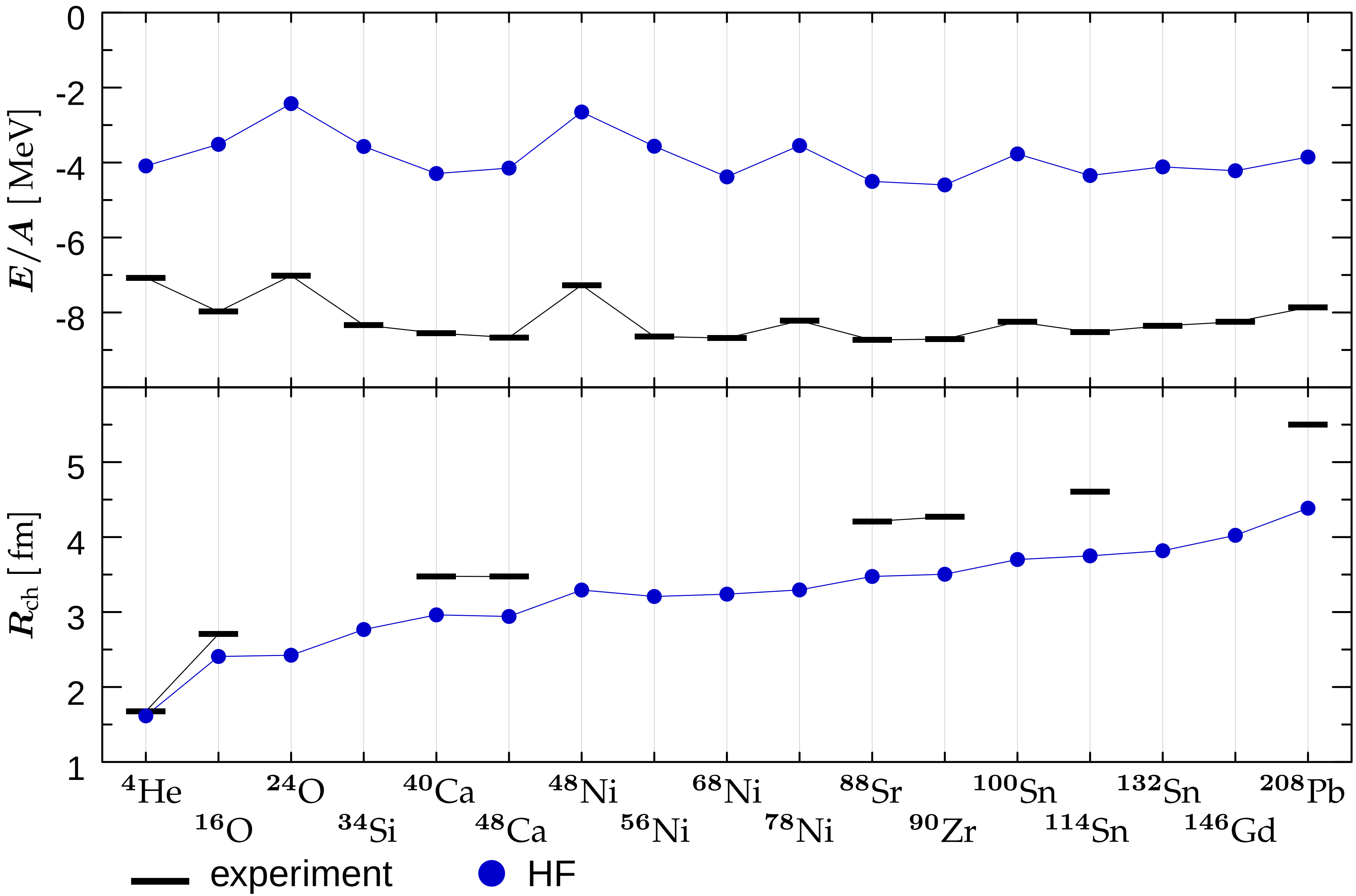
Hartree-Fock with SRG Potentials



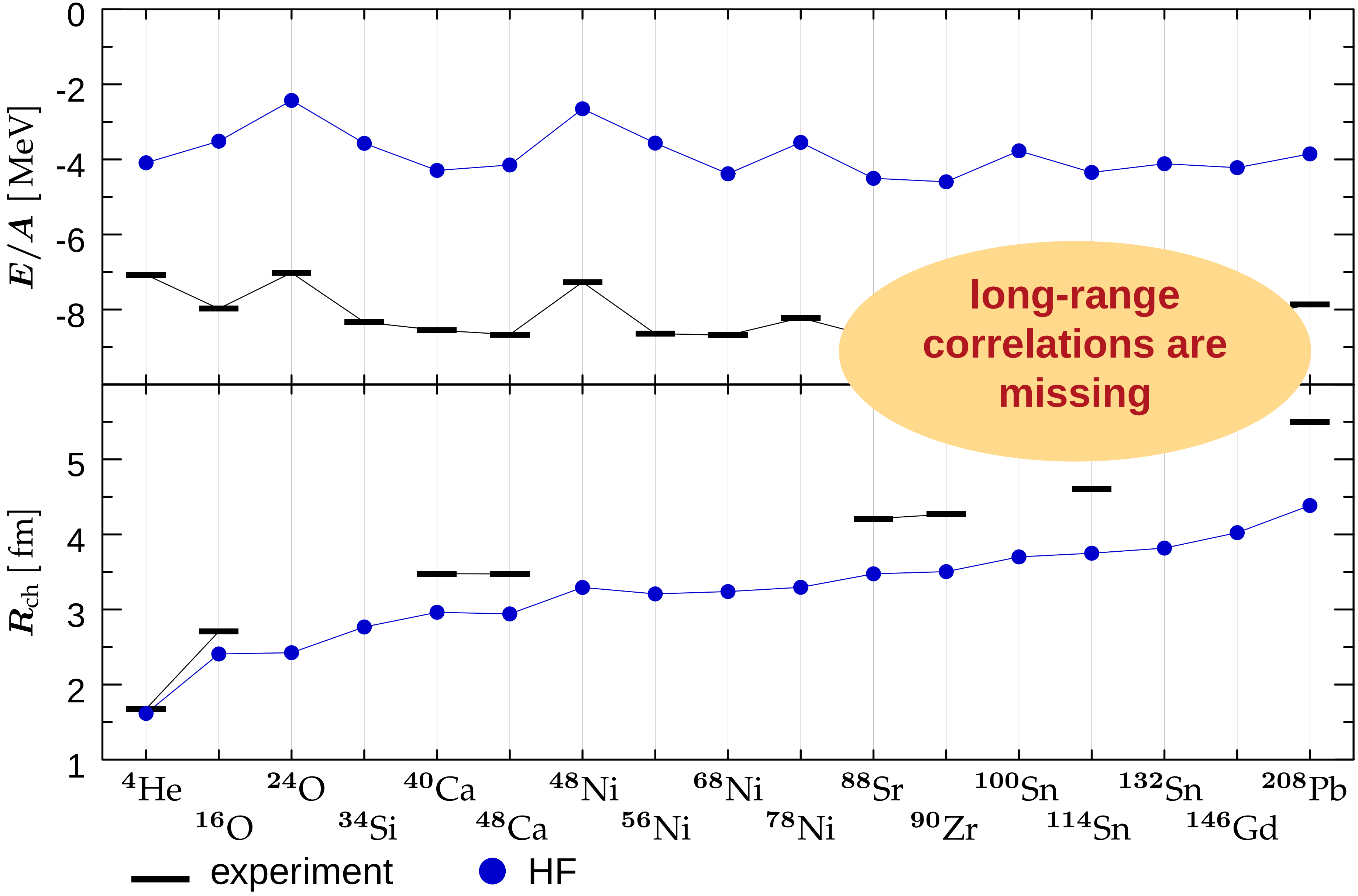
Hartree-Fock with SRG Potentials



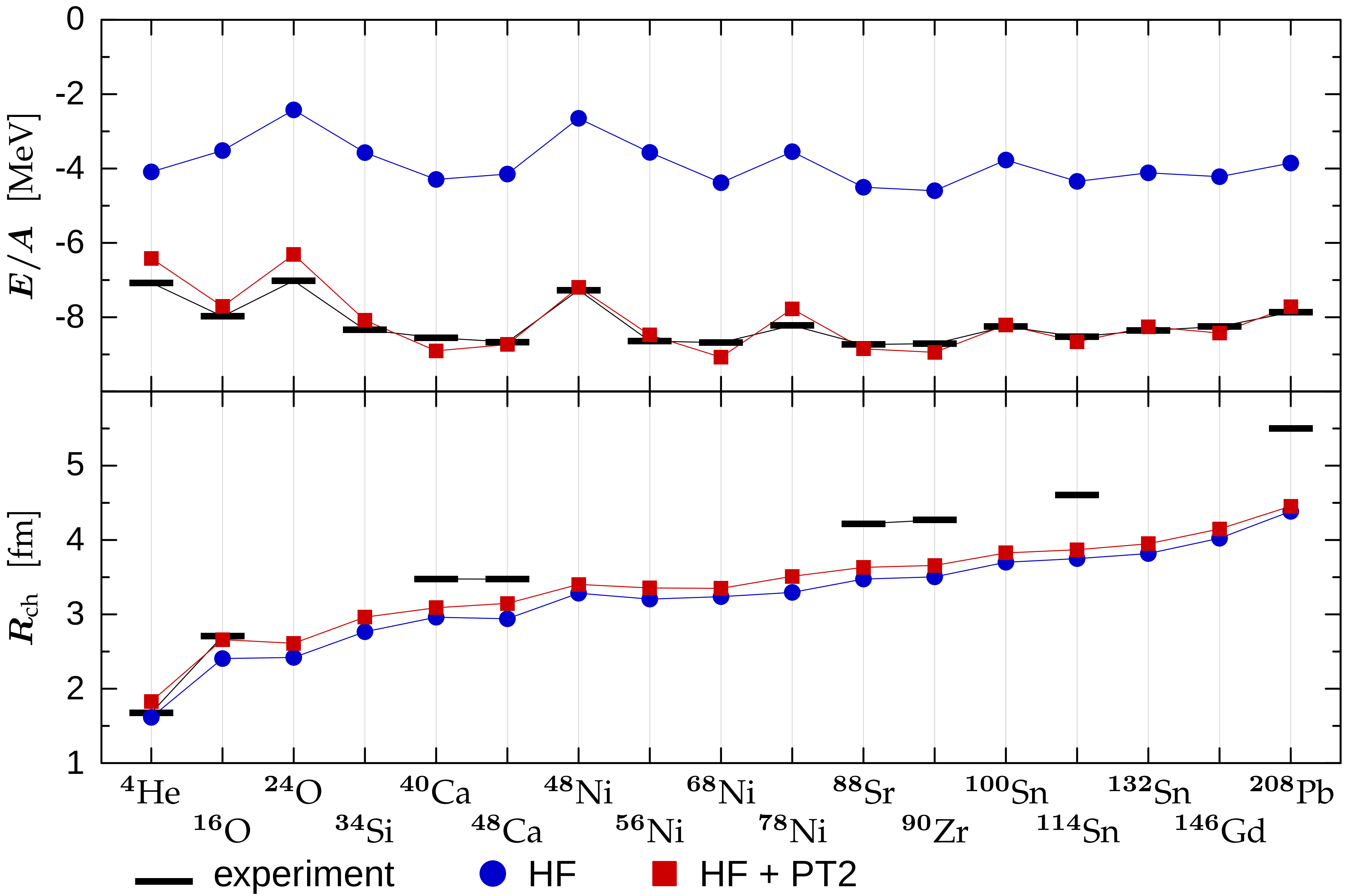
Hartree-Fock with V_{UCOM}



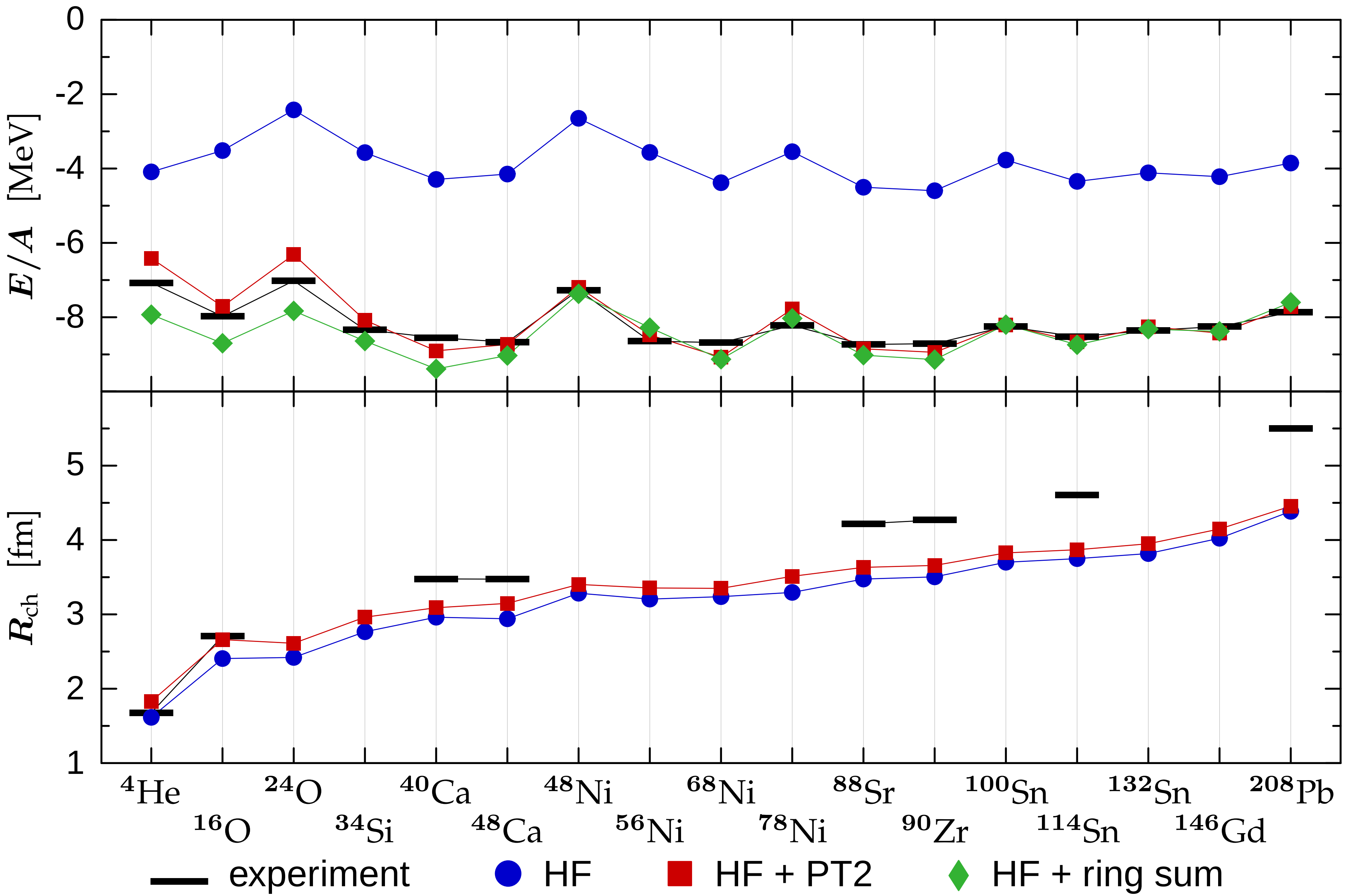
Hartree-Fock with V_{UCOM}



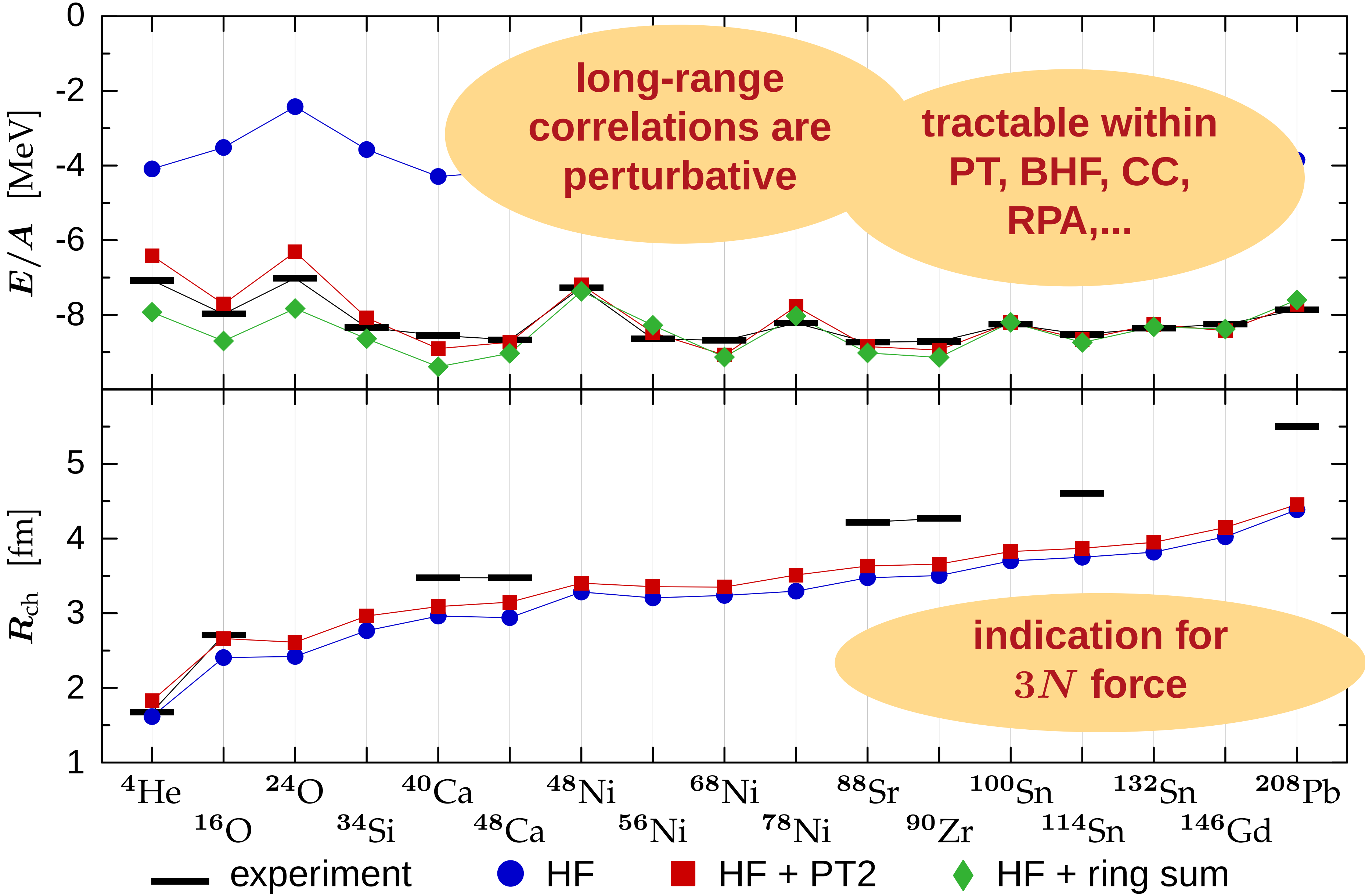
Perturbation Theory with V_{UCOM}



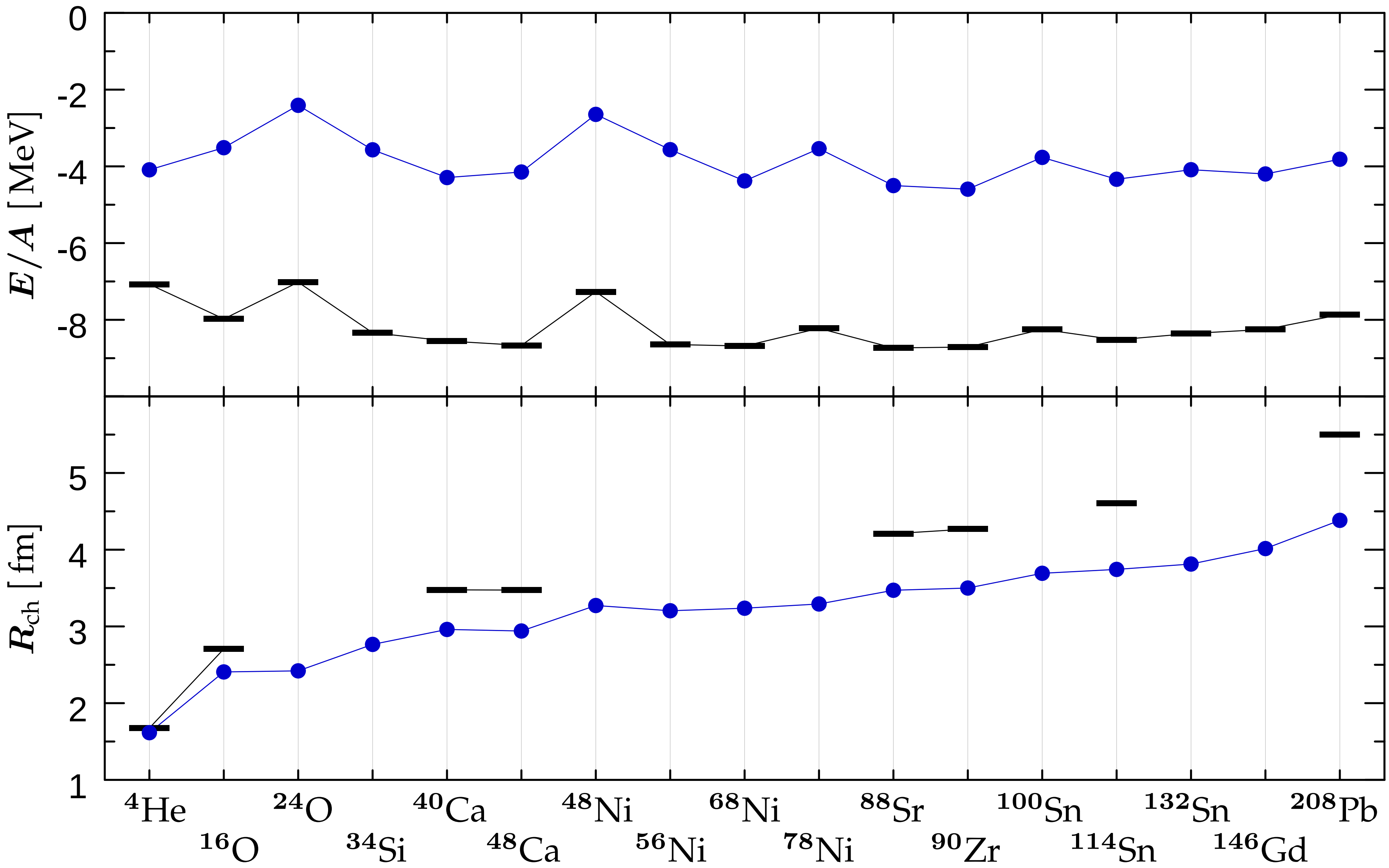
RPA Ring Summation with V_{UCOM}



RPA Ring Summation with V_{UCOM}

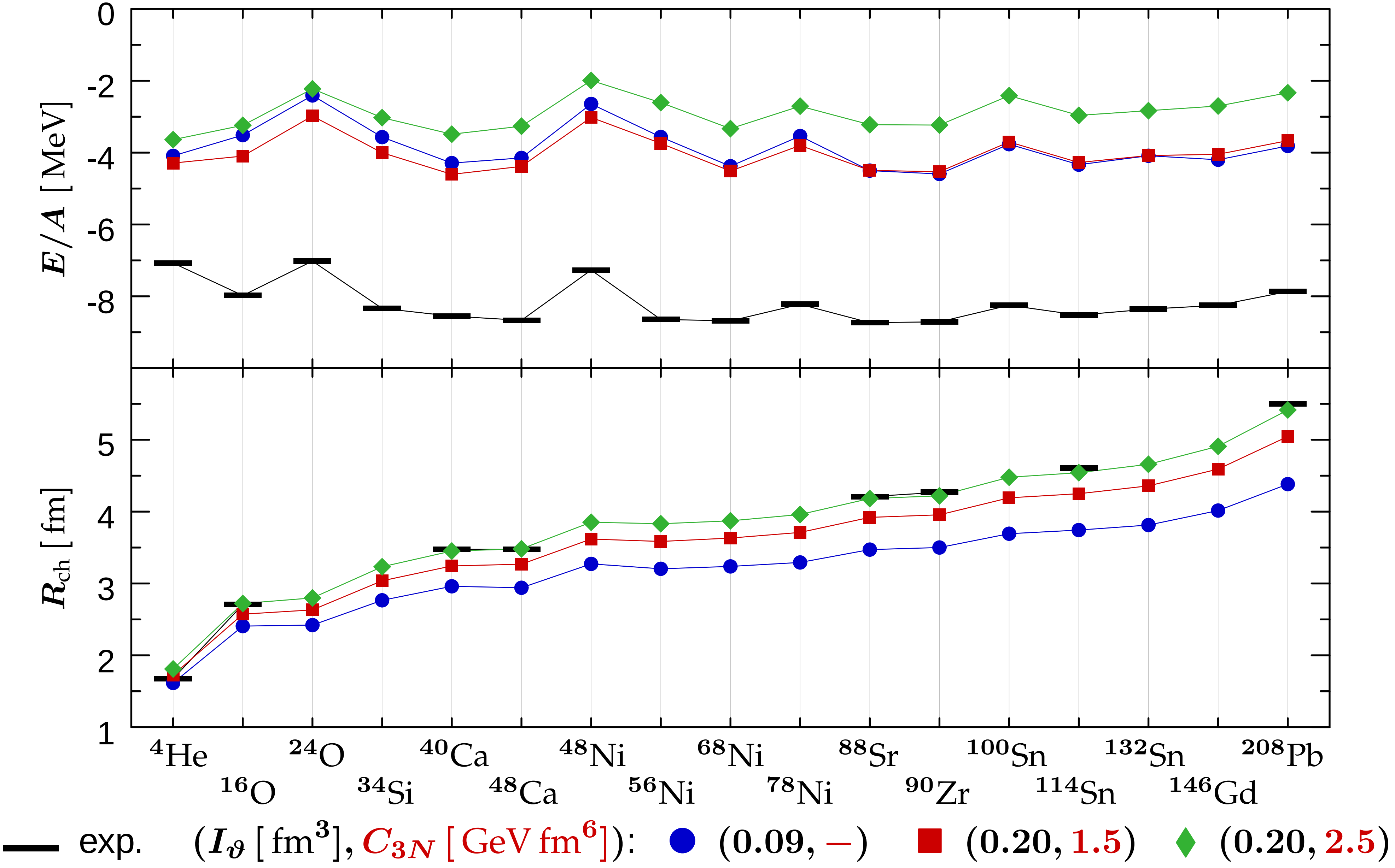


3N Forces: Energies & Radii

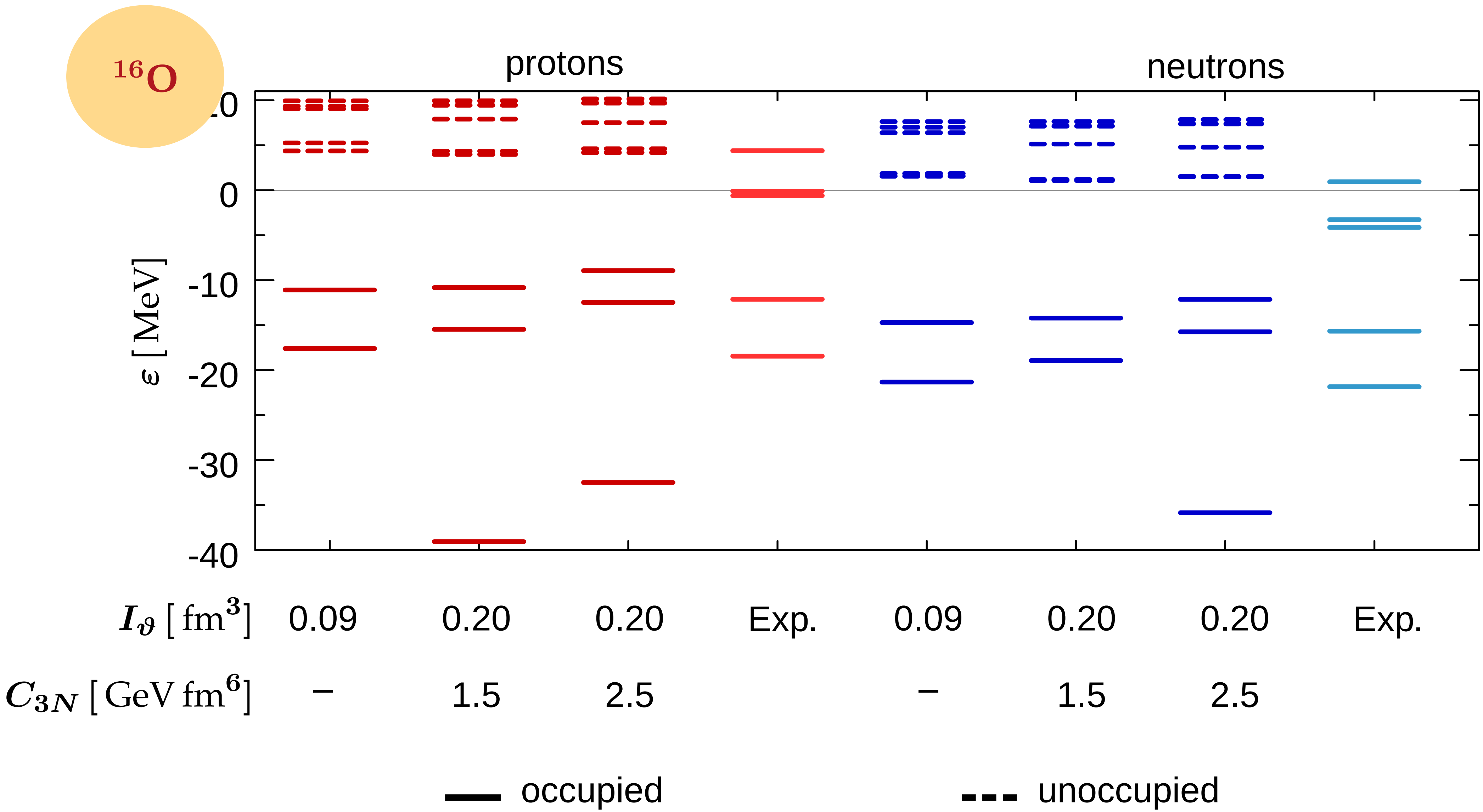


— exp. $(I_\vartheta\,[\mathrm{fm}^3], C_{3N}\,[\mathrm{GeV\,fm}^6])$: ● $(0.09, -)$

$3N$ Forces: Energies & Radii



$3N$ Forces: HF Single-Particle Energies



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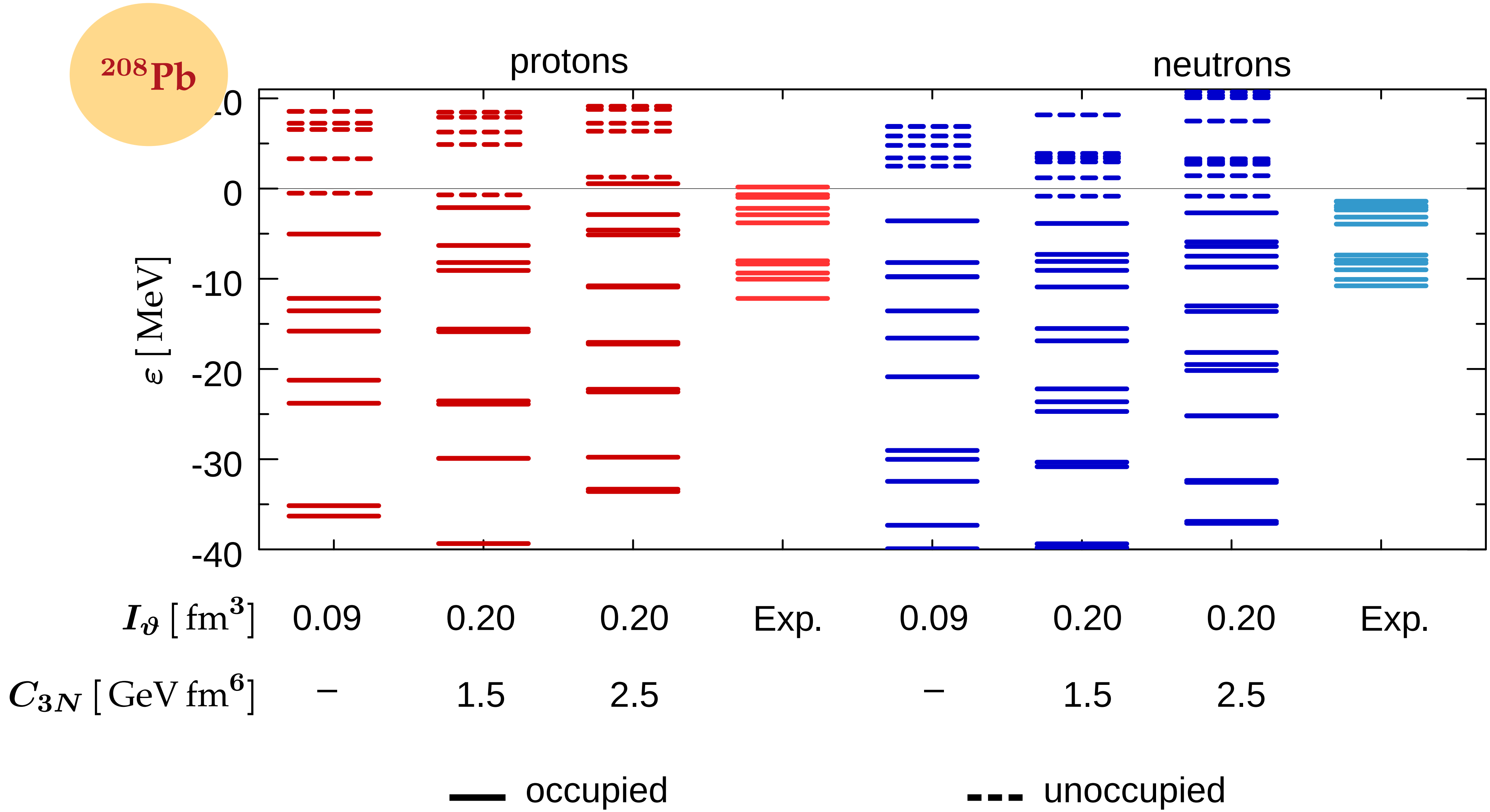

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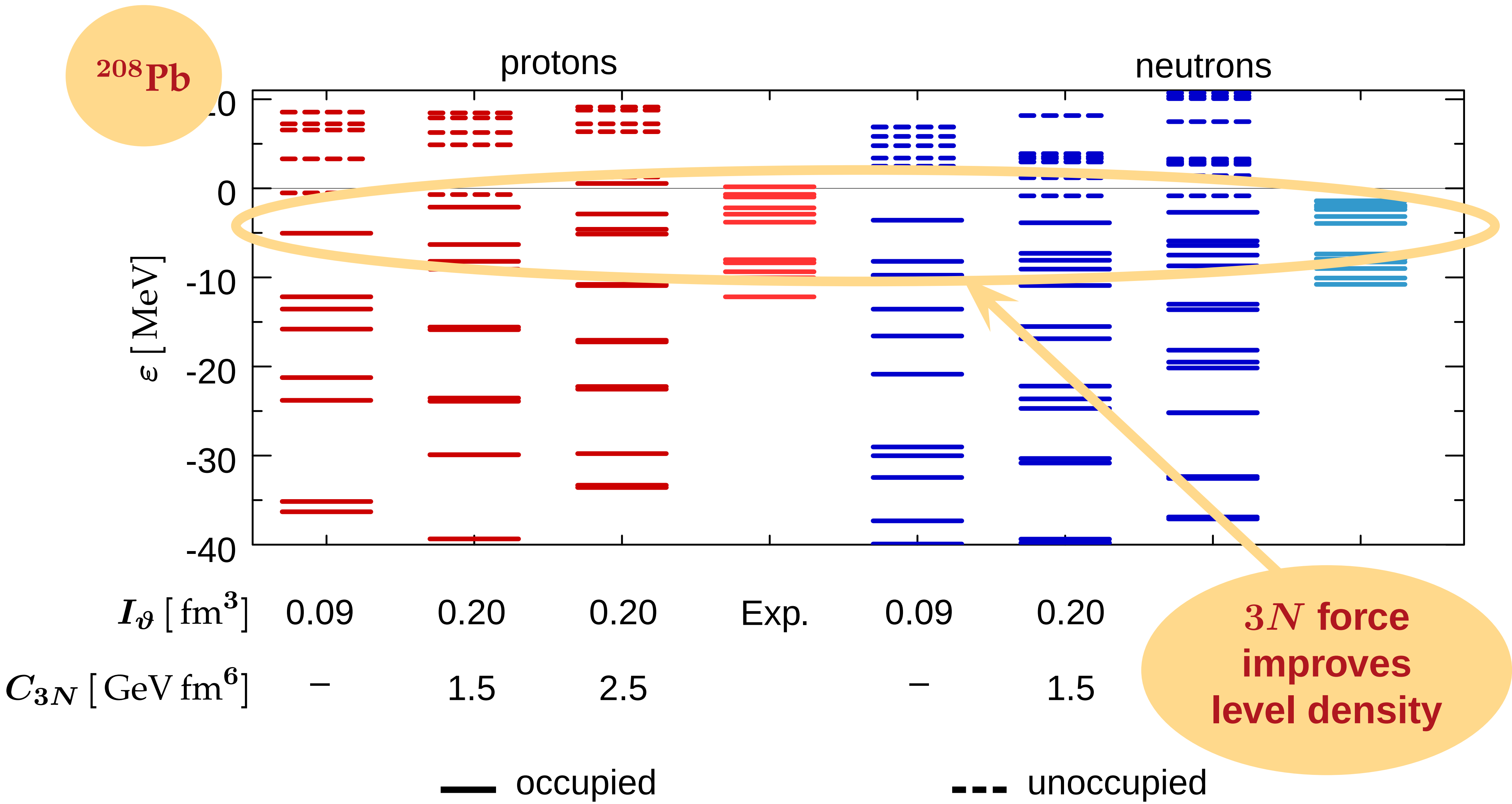
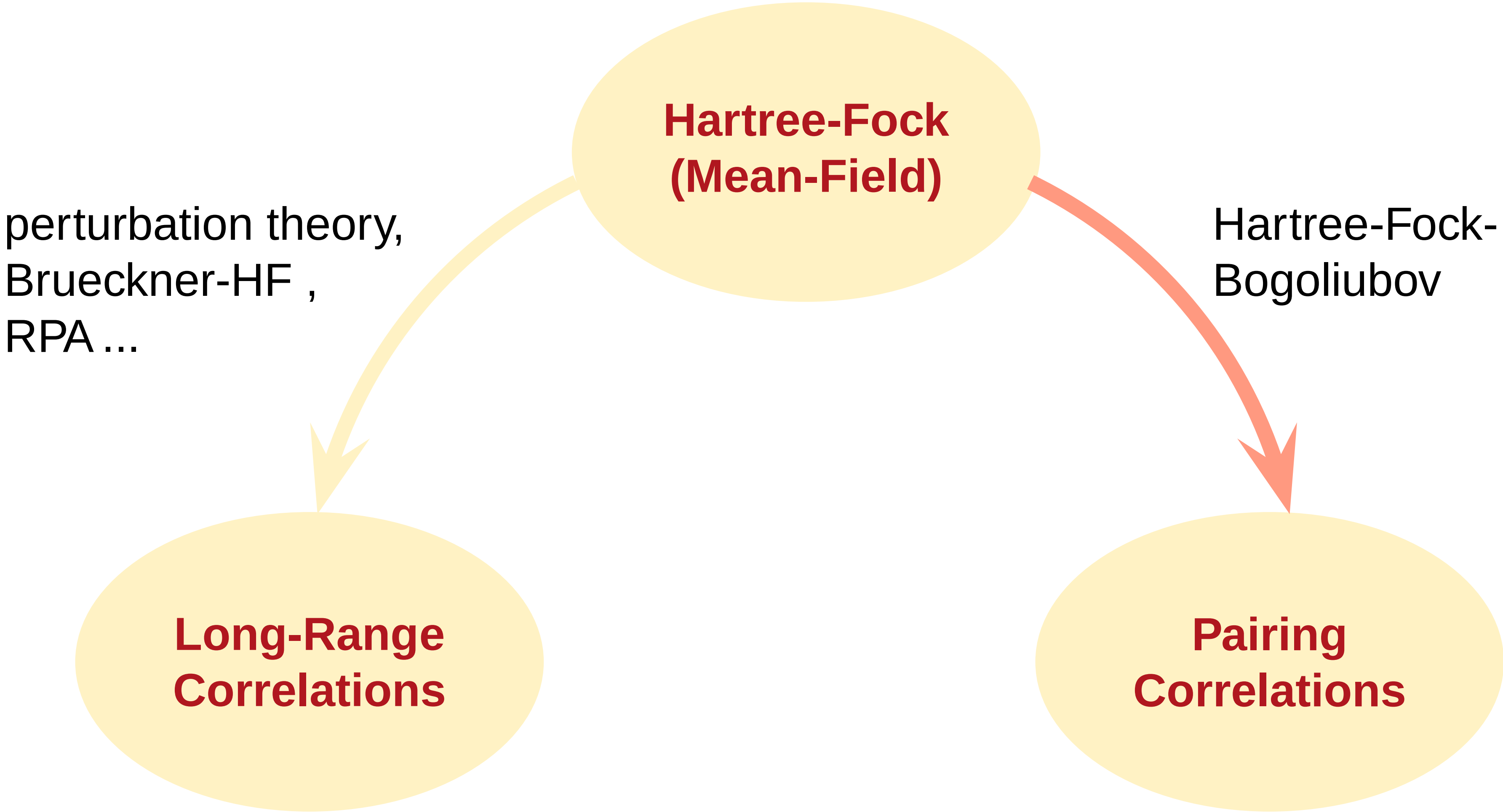


Beyond Hartree-Fock



HFB Theory Overview

Bogoliubov Transformation

$$\beta_k^\dagger = \sum_q U_{qk} \mathbf{c}_q^\dagger + V_{qk} \mathbf{c}_q$$

$$\beta_k = \sum_q U_{qk}^* \mathbf{c}_q + V_{qk}^* \mathbf{c}_q^\dagger$$

where

$$\{\beta_k, \beta_{k'}\} \stackrel{!}{=} \{\beta_k^\dagger, \beta_{k'}^\dagger\} \stackrel{!}{=} 0$$

$$\{\beta_k, \beta_{k'}^\dagger\} \stackrel{!}{=} \delta_{kk'}$$

HFB Densities & Fields

$$\rho_{kk'} \equiv \left\langle \Psi \middle| \mathbf{c}_{k'}^\dagger \mathbf{c}_k \middle| \Psi \right\rangle = (V^* V^T)_{kk'}$$

$$\kappa_{kk'} \equiv \left\langle \Psi \middle| \mathbf{c}_{k'} \mathbf{c}_k \middle| \Psi \right\rangle = (V^* U^T)_{kk'}$$

$$\Gamma_{kk'} = \sum_{qq'} \left(\frac{2}{A} \bar{\mathrm{t}}_{\mathrm{rel}} + \bar{\mathrm{v}} \right)_{kq',k'q} \rho_{qq'}$$

$$\Delta_{kk'} = \sum_{qq'} \left(\frac{2}{A} \bar{\mathrm{t}}_{\mathrm{rel}} + \bar{\mathrm{v}} \right)_{kk',qq'} \kappa_{qq'}$$

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Energy

$$E[\rho, \kappa, \kappa^*] = \frac{\langle \Psi | H | \Psi \rangle}{\langle \Psi | \Psi \rangle} \equiv \frac{1}{2} (\mathrm{tr}\ \Gamma\rho - \mathrm{tr}\ \Delta\kappa^*)$$

HFB Equations

$$(\mathcal{H} - \lambda\mathcal{N}) \begin{pmatrix} U \\ V \end{pmatrix} \equiv \begin{pmatrix} \Gamma - \lambda & \Delta \\ -\Delta^* & -\Gamma^* + \lambda \end{pmatrix} \begin{pmatrix} U \\ V \end{pmatrix} = E \begin{pmatrix} U \\ V \end{pmatrix}$$

Particle Number Projection

Projected Energy

$$E(N_0) = \frac{\langle \Psi | \mathrm{H} \mathrm{P}_{N_0} | \Psi \rangle}{\langle \Psi | \mathrm{P}_{N_0} | \Psi \rangle} = \frac{1}{2\pi \langle \mathrm{P}_{N_0} \rangle} \int_0^{2\pi} d\phi \, \langle \Psi | \mathrm{H} e^{i\phi(\mathrm{N} - N_0)} | \Psi \rangle$$

Flocard & Onishi, Ann. Phys. 254, **275** (1997, approx. PNP); Sheikh et al., Phys. Rev. **C66**, 044318 (2002, exact PNP)

Particle Number Projection

Variation of Projected Energy

$$\delta E(N_0) = \frac{1}{2\pi\langle \mathrm{P}_{N_0}\rangle} \int_0^{2\pi} d\phi \, \langle e^{i\phi(\mathrm{N}-N_0)}\rangle \left\{ \delta\langle \mathrm{H}\rangle_\phi - \left(E(N_0) - \langle \mathrm{H}\rangle_\phi \right) \delta \log \langle e^{i\phi \mathrm{N}}\rangle \right\}$$

$$\langle \mathrm{H}\rangle_\phi \equiv \langle \mathrm{H} e^{i\phi\mathrm{N}}\rangle / \langle e^{i\phi\mathrm{N}}\rangle$$

Flocard & Onishi, Ann. Phys. 254, **275** (1997, approx. PNP); Sheikh et al., Phys. Rev. **C66**, 044318 (2002, exact PNP)

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Lipkin-Nogami + PAV

- power series expansion
- expansion coefficients **not varied**
- indeterminate / numerically unstable at shell closures
- exact PNP after variation

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VAP

- higher (but managable) computational effort
- implement with care: **subtle cancellations between divergences of direct, exchange, and pairing terms**

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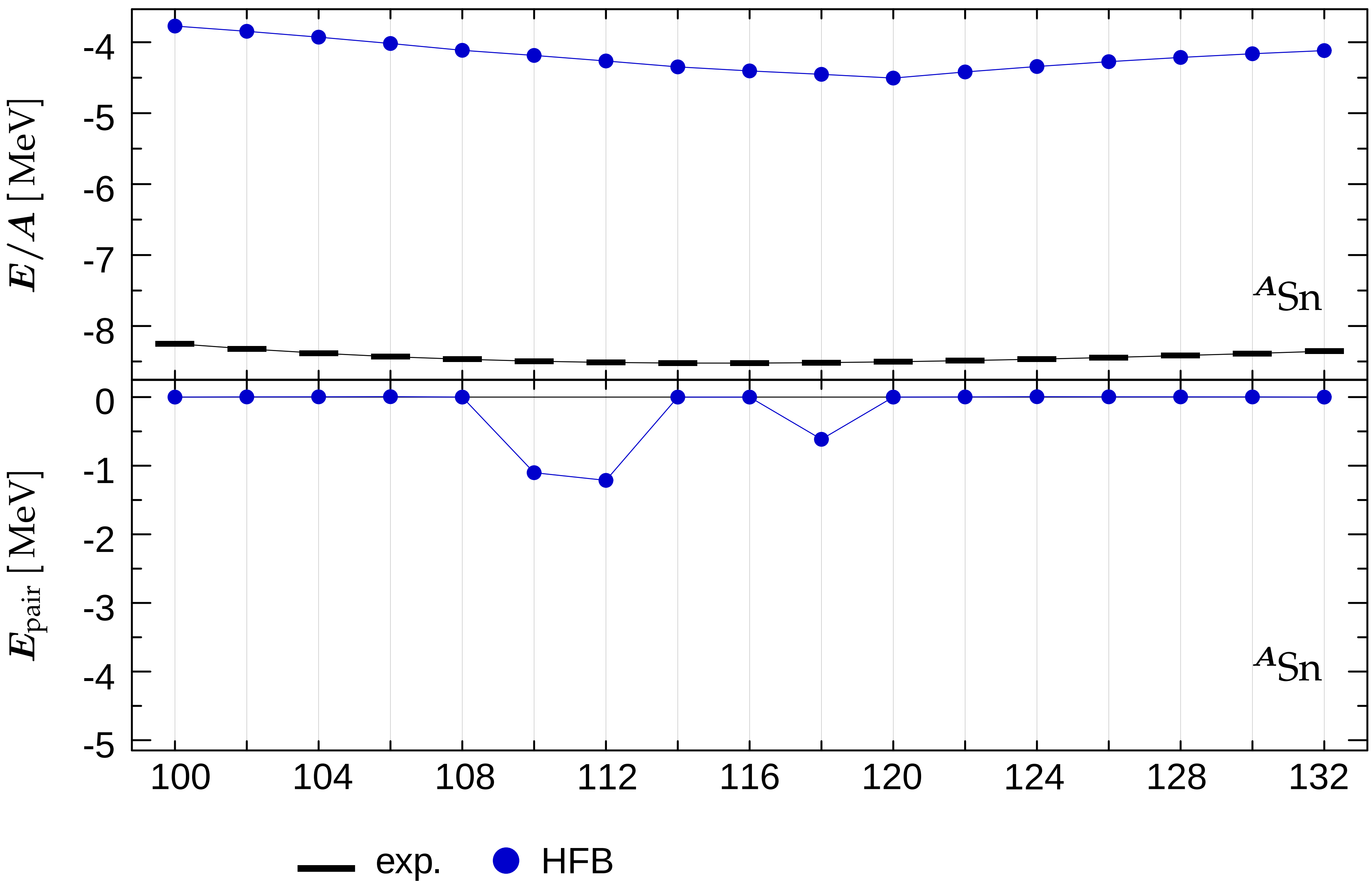
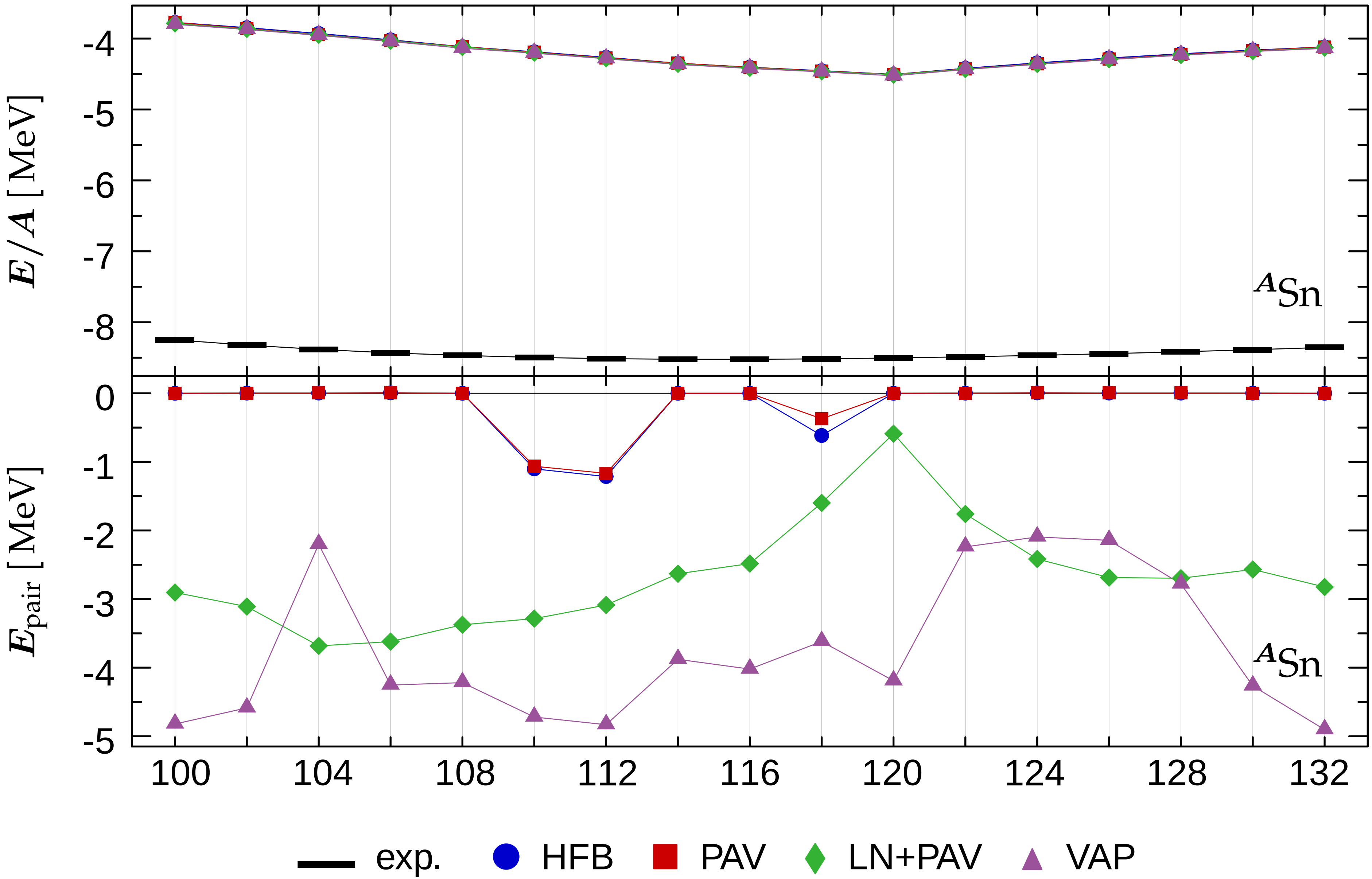
☞ Structure of **HFB equations is preserved** by both methods!

Flocard & Onishi, Ann. Phys. 254, **275** (1997, approx. PNP); Sheikh et al., Phys. Rev. **C66**, 044318 (2002, exact PNP)

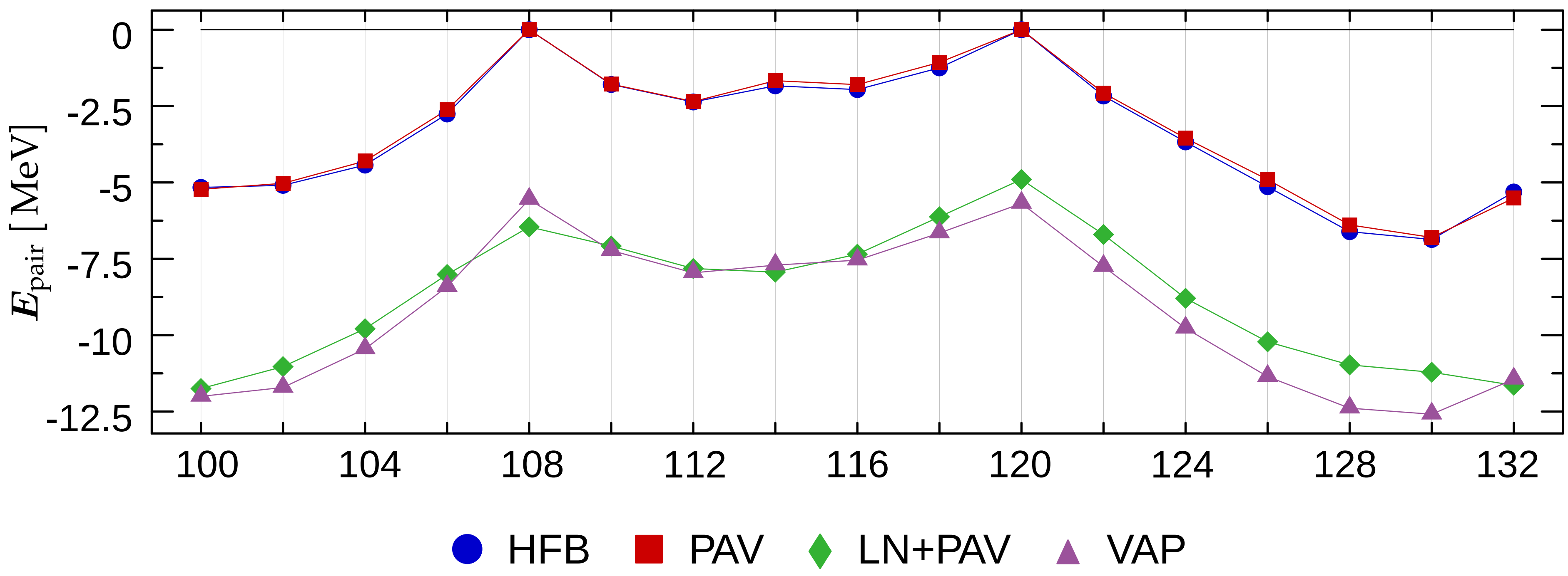
Sn Isotopes: Binding & Pairing Energies



Sn Isotopes: Binding & Pairing Energies



Density-Dependent Force: Pairing



- **linear density dependence**, $t_0 = \frac{1}{6} C_{3N}$:

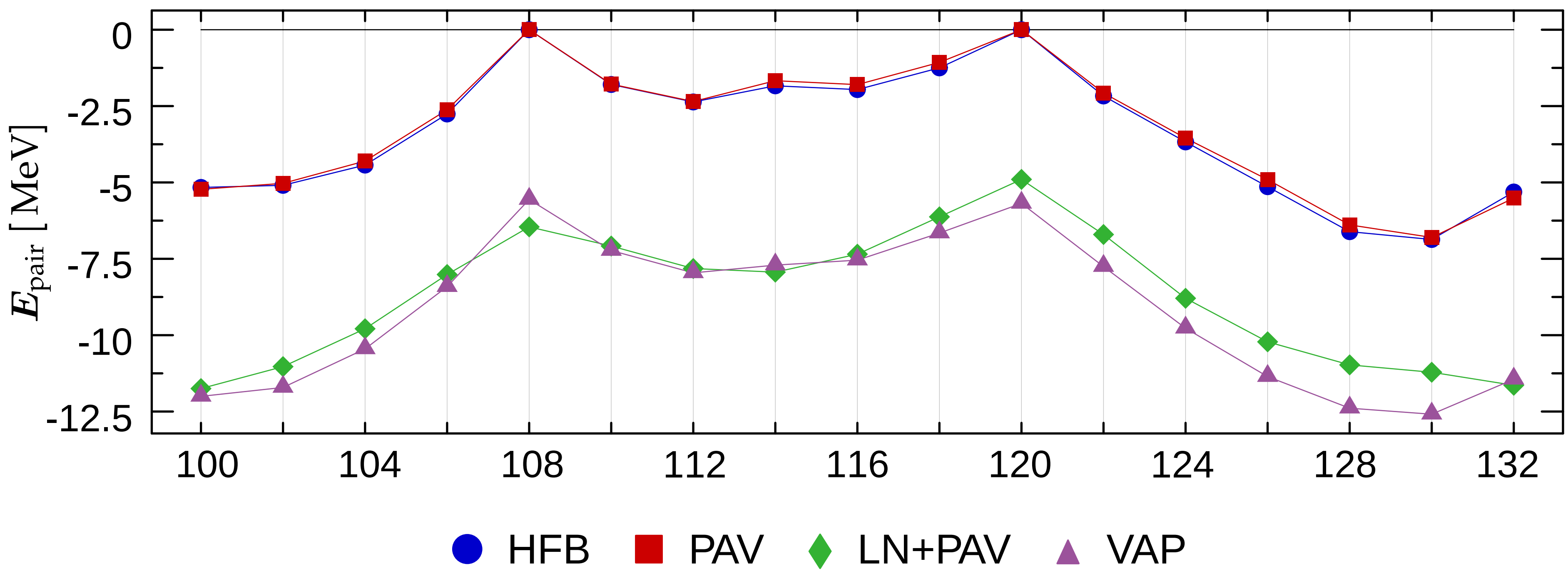
$$V_\rho = t_0 \left(1 + P_\sigma\right) \rho \left(\tfrac{1}{2}(\vec{r}_1 + \vec{r}_2)\right) \delta^3 \left(\vec{r}_1 - \vec{r}_2\right)$$

- **mixed density** for PAV/VAP: $\rho(\vec{R}) \longrightarrow \rho_{\phi_p,\phi_n}(\vec{R})$

- phenomenological VAP calculations: $E_{\rm pair} \simeq 10 - 20$ MeV

(Stoitsov et al., nucl-th/0610061; Anguiano et al., Phys. Lett. **B545** (2002), 62)

Density-Dependent Force: Pairing



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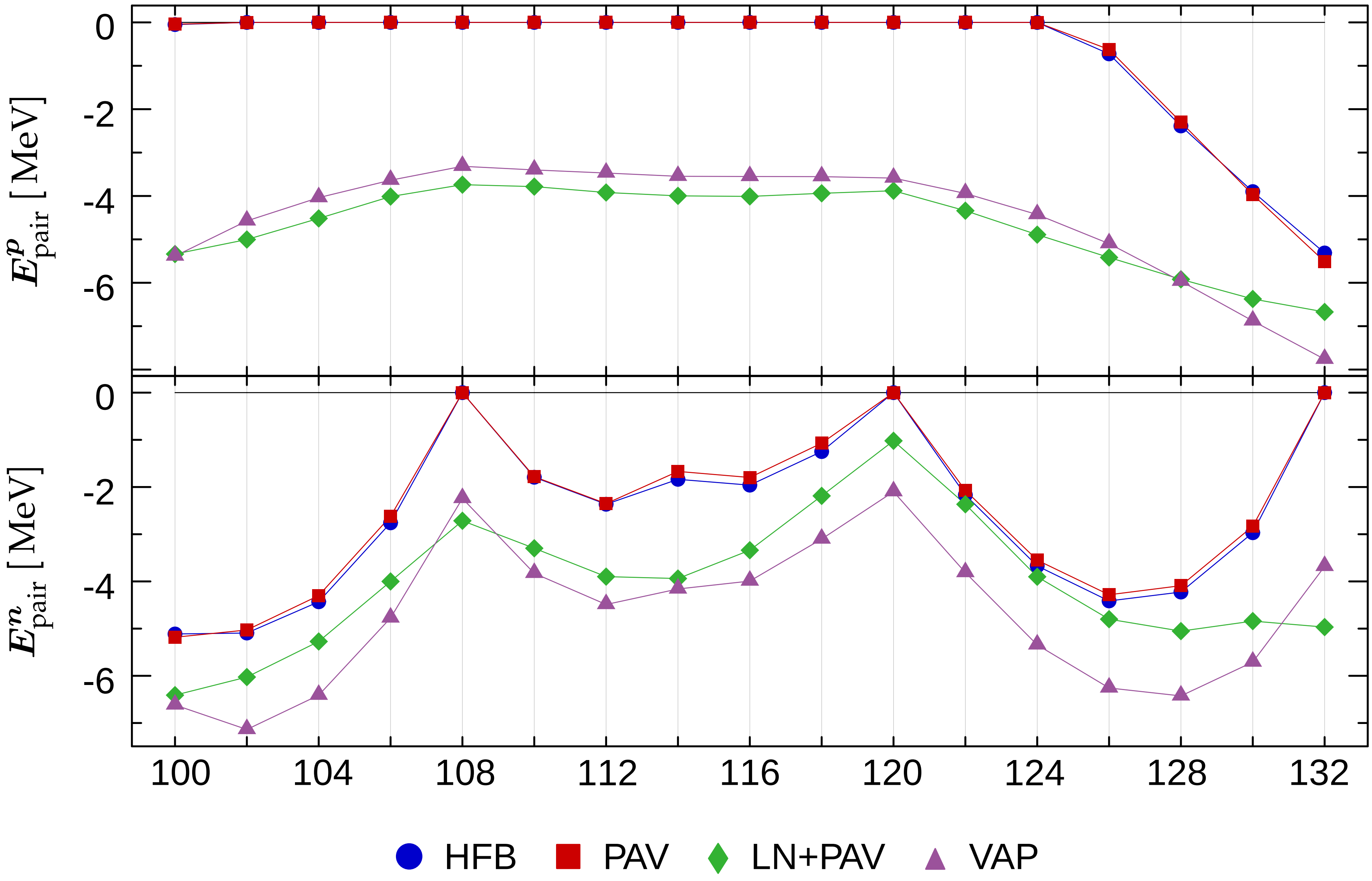
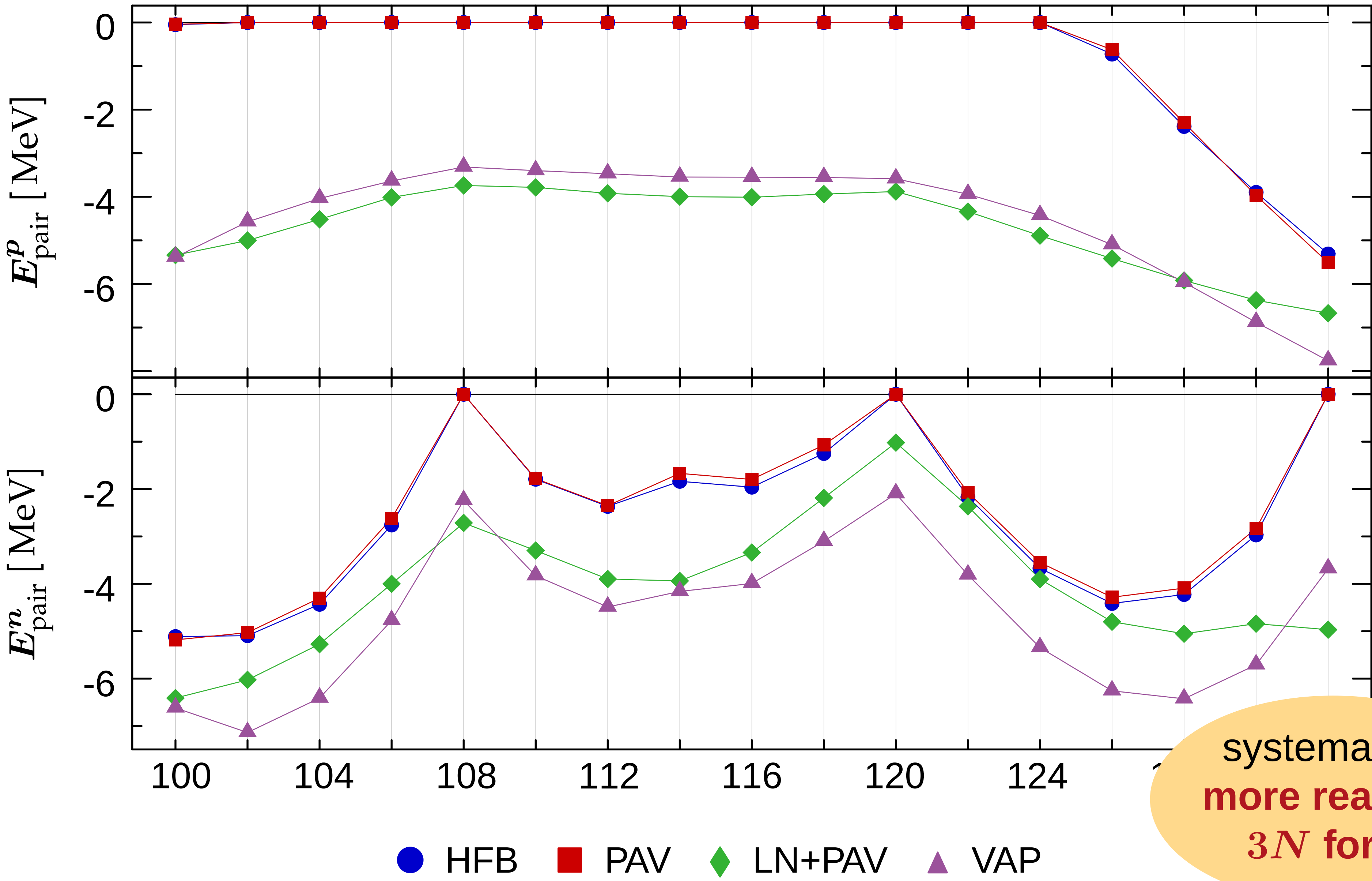
(Stoitsov et al., nucl-th/0610061; Anguiano et al., Phys. Lett. **B545** (2002), 62)

"correct" order of magnitude **with realistic NN int.**

Density-Dependent Force: Pairing



Density-Dependent Force: Pairing



systematics:
more realistic $3N$ force

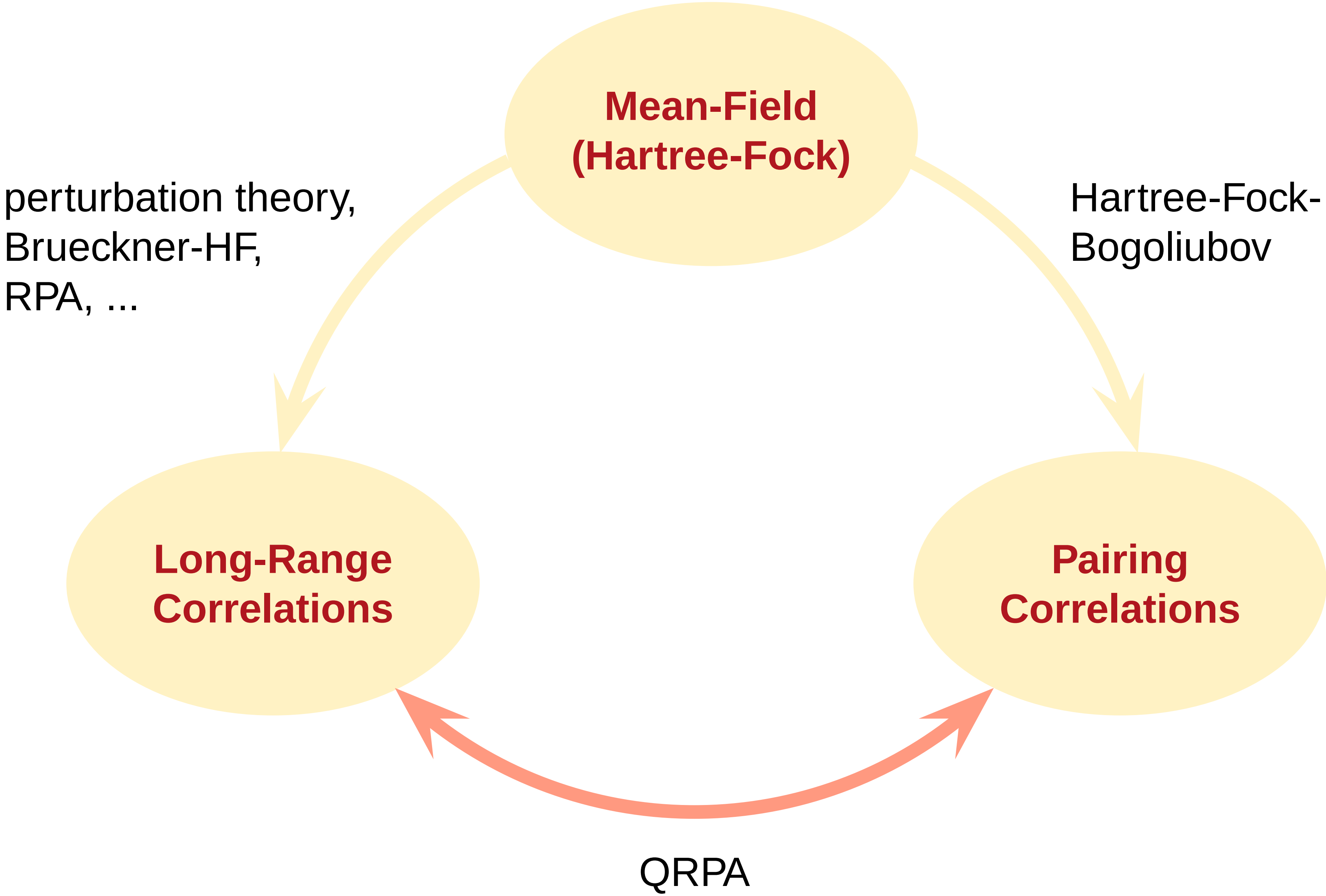
Conclusions

Modern Effective Interactions

- **Status**
  - treatment of **short-range central** and **tensor correlations** by unitary transformations:
    - Unitary Correlation Operator Method
    - Similarity Renormalization Group
  - **universal phase-shift equivalent** correlated interaction $\mathbf{V}_{\text{UCOM}}$

- **Outlook**
  - **connections** between UCOM and SRG
  - inclusion & treatment of $3N$ **Forces**, in particular...
  - **chiral interactions**

HF, HFB, and Beyond



HF, HFB, and Beyond

- **Status**
  - **fully consistent** HFB calculations with particle number projection
  - **inclusion of $3N$-forces**: **contact & finite range** matrix elements for HF, density-dependent force for HFB, RPA, ...
  - **Like-particle-** & pn**-QRPA** (benchmarked)
- **Outlook**
  - pn **pairing**, Isobaric Analog & Gamow-Teller Resonances
  - **deformation** and **symmetry restoration** by projection (isospin, parity, angular momentum)
  - **caveat:** analytic structure of density-dependent forces is very important in symmetry-projected HFB (J. Dobaczewski, arXiv: 0708.0441)

Perspectives

- **Modern Effective Interactions**
  - treatment of short-range central and tensor correlations by unitary transformations: UCOM, SRG, Lee-Suzuki,...
  - phase-shift equivalent correlated interaction $\mathbf{V}_{\mathrm{UCOM}}$
  - universal input for...
- **Innovative Many-Body Methods**
  - No-Core Shell Model,...
  - Importance Truncated NCSM, Coupled Cluster Method,...
  - Hartree-Fock plus MBPT, Padé Resummed MBPT, BHF, HFB, RPA,...
  - Fermionic Molecular Dynamics,...

Last Words...

My Collaborators

- R. Roth, P. Papakonstantinou, A. Zapp, P. Hedfeld, S. Reinhardt

  Institut für Kernphysik, TU Darmstadt

- T. Neff, H. Feldmeier

  Gesellschaft für Schwerionenforschung (GSI)

- N. Paar

  Department of Physics — Faculty of Science, University of Zagreb, Croatia

References

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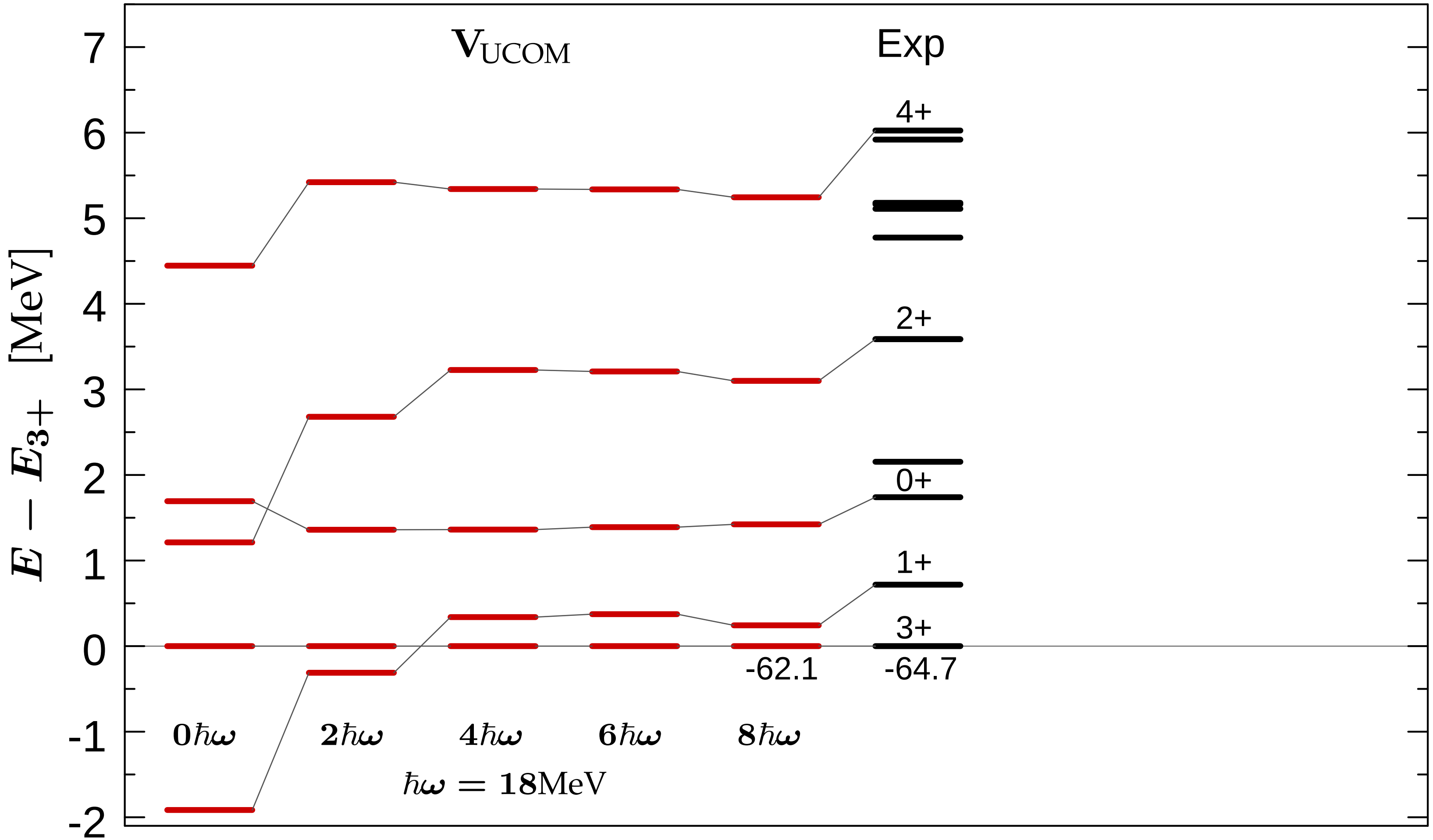
- N. Paar, P. Papakonstantinou, H. Hergert, and R. Roth, Phys. Rev. C**74**, 014318 (2006)
- R. Roth, P. Papakonstantinou, N.Paar, H. Hergert, T. Neff, and H. Feldmeier, Phys. Rev. C**73**, 044312 (2006)
- http://crunch.ikp.physik.tu-darmstadt.de/tnp/

Appendix

^{10}B: Hallmark of a $3N$ Interaction?



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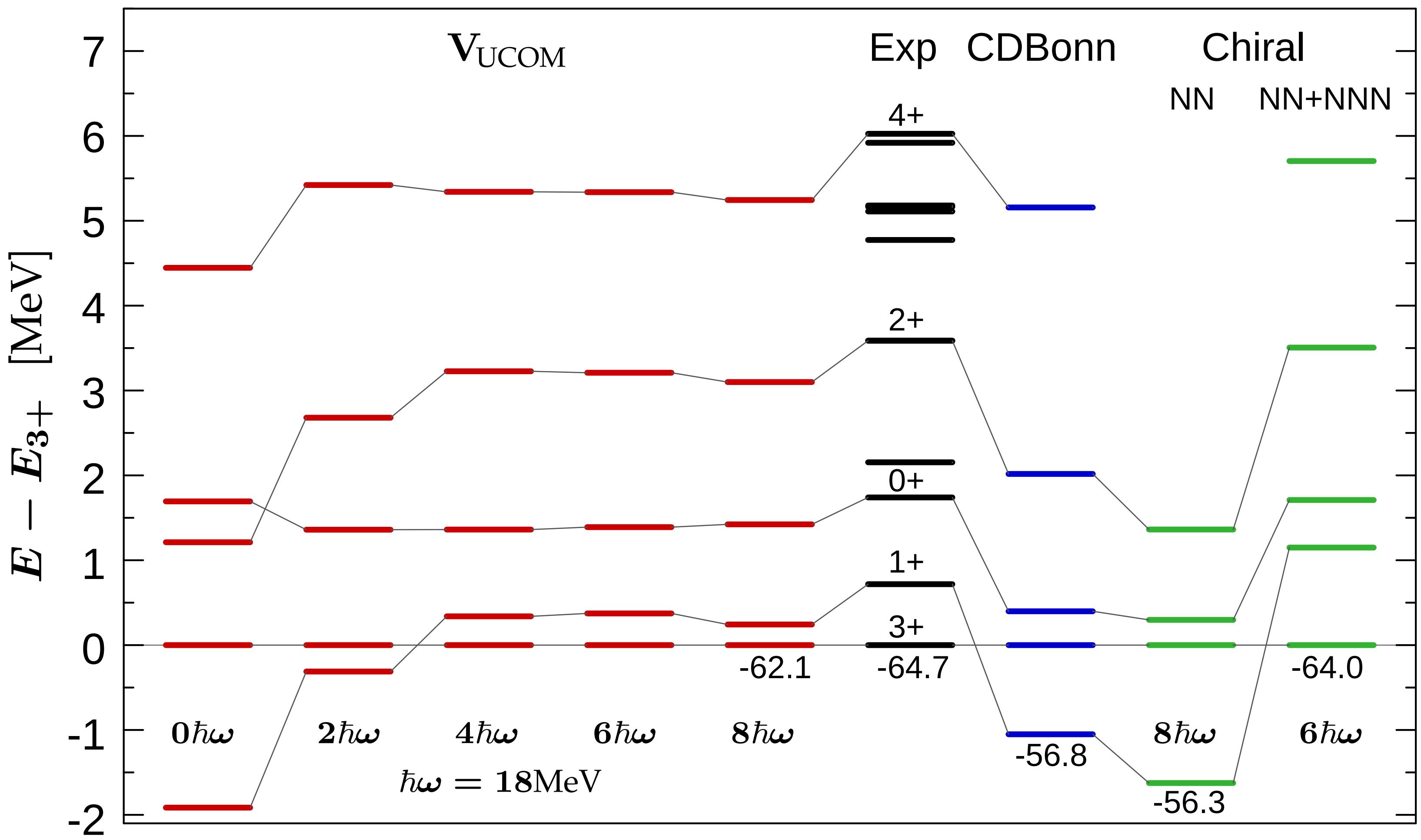
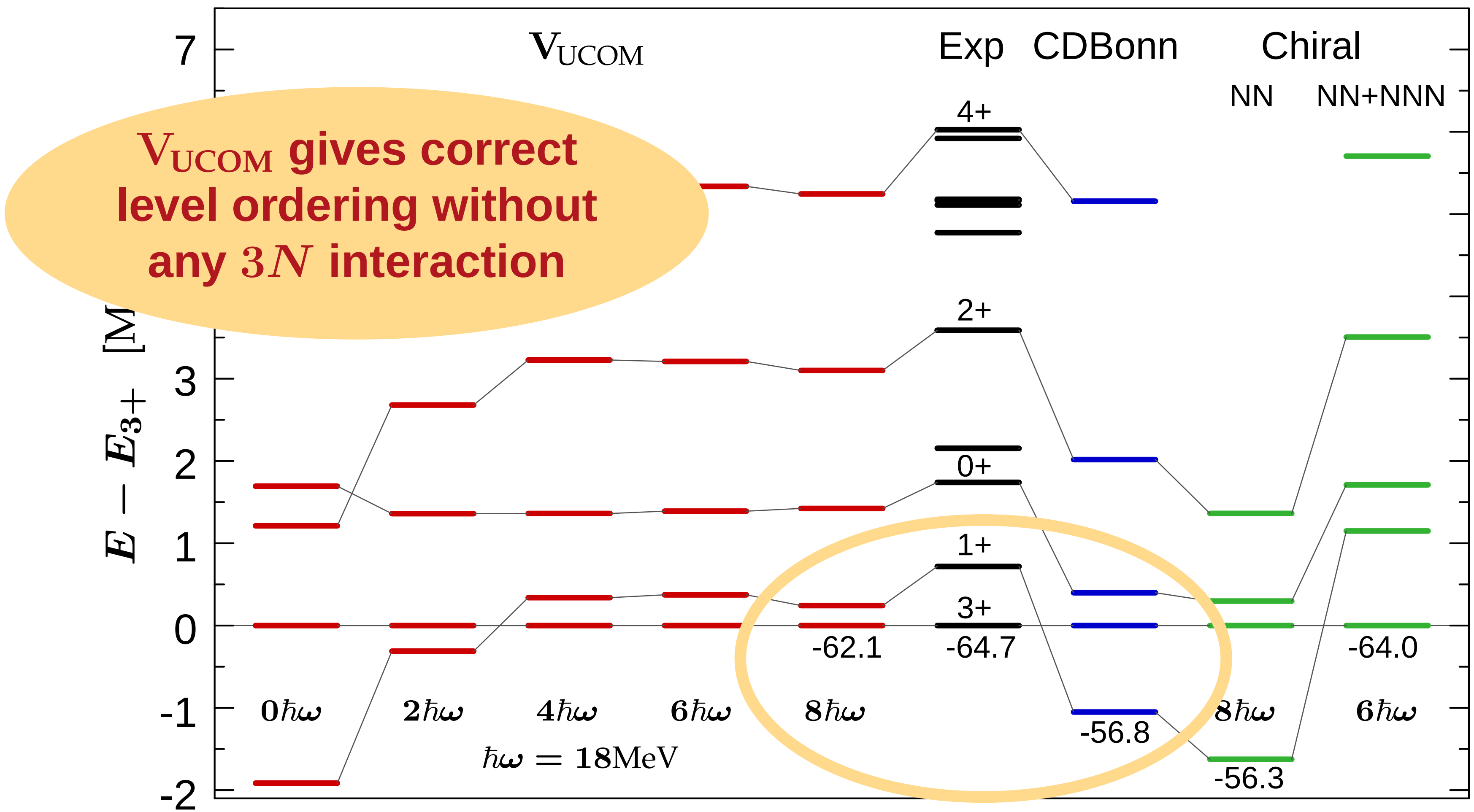


^{10}B: Hallmark of a $3N$ Interaction?

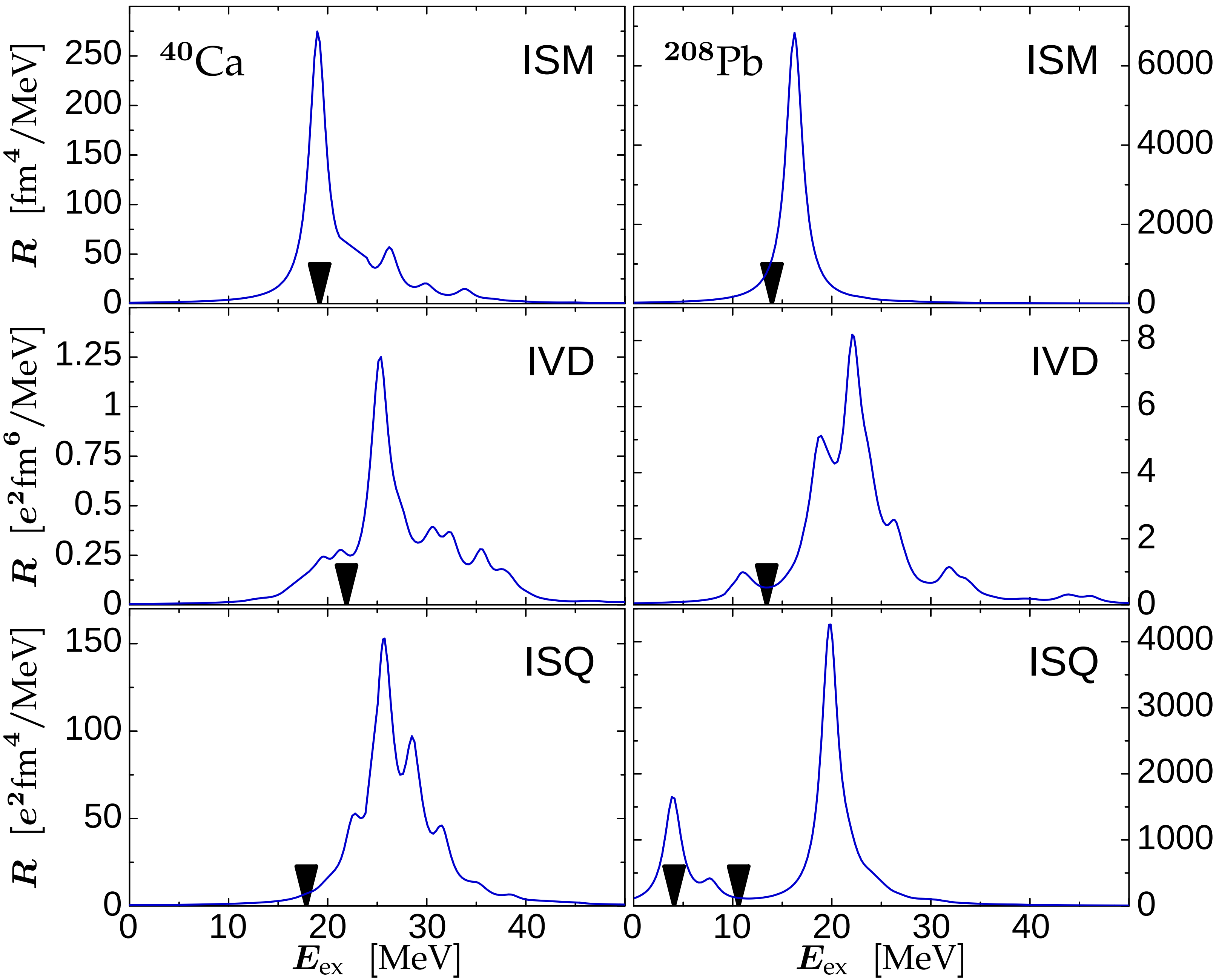


RPA + Realistic Interactions

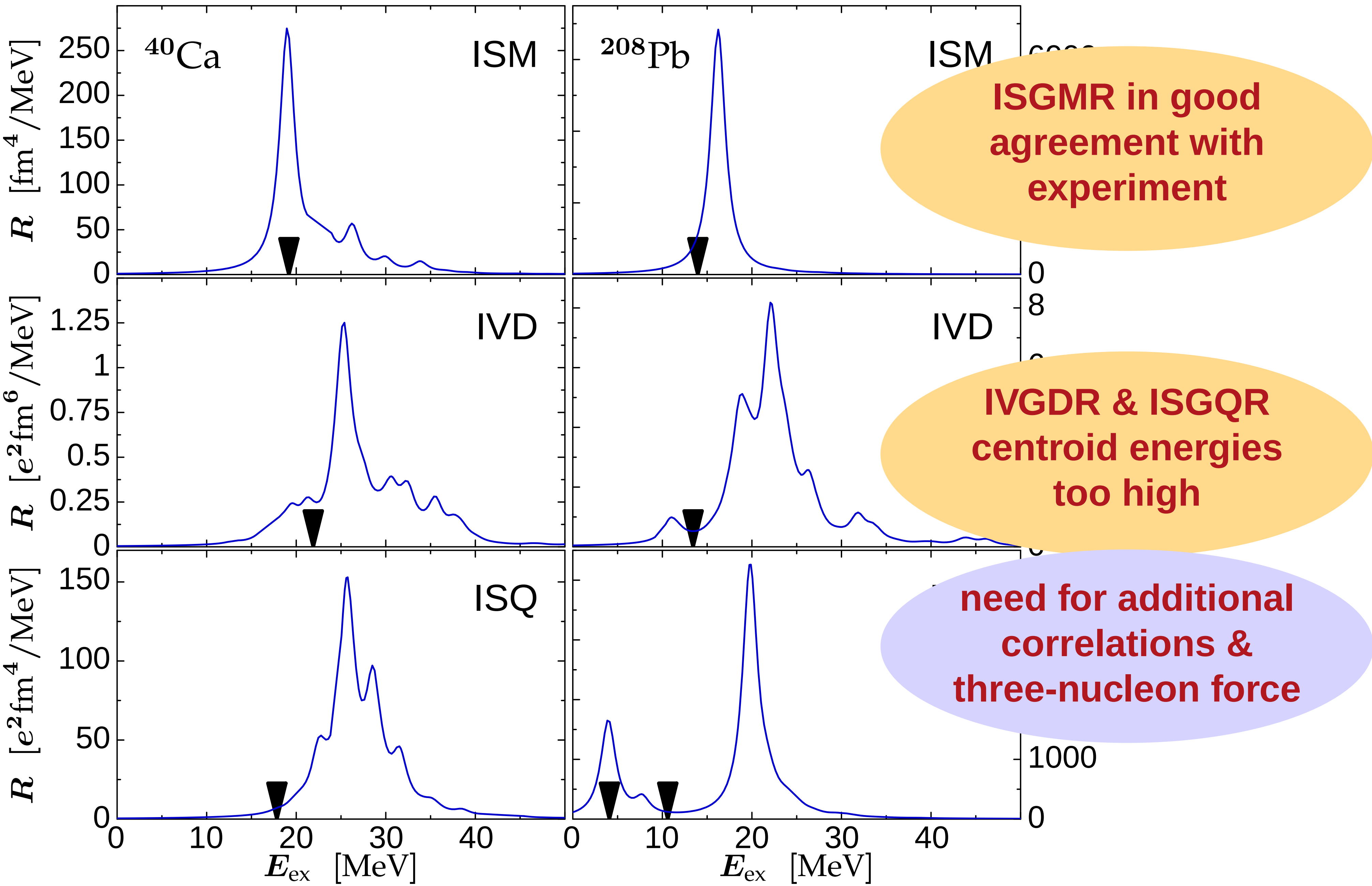
RPA, ERPA & SRPA
+
Matrix Elements of Correlated Realistic Interaction V_{UCOM}

- **fully self-consistent RPA** based on the Hartree-Fock orbits using the same V_{UCOM}
- recovering sum rules with high precision, spurious center-of-mass mode fully decoupled at ~ 10 keV
- **Extended-RPA and Second-RPA** to include effects of ground state correlations and complex configurations

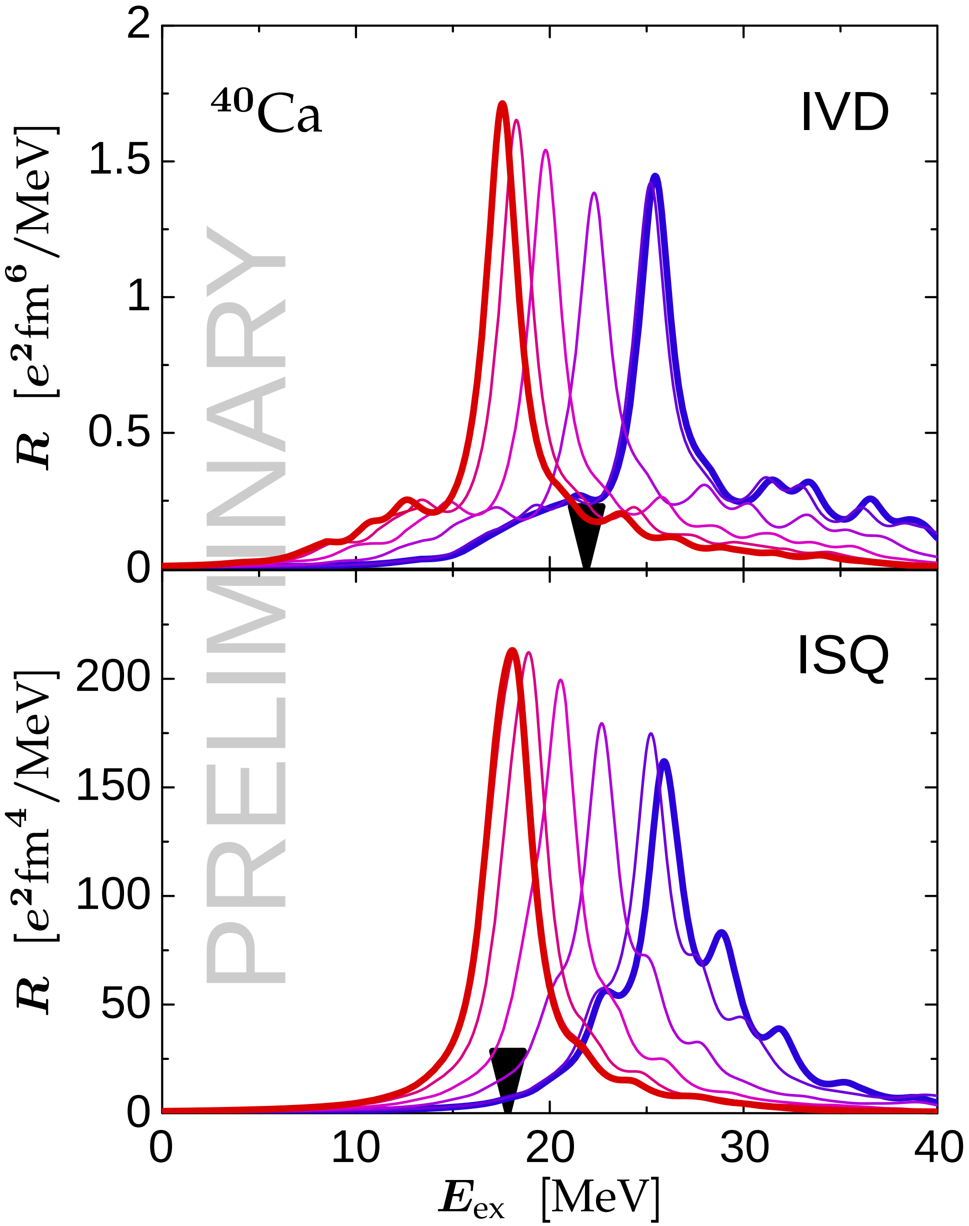
RPA with V_{UCOM}



RPA with V_{UCOM}



SRPA: Complex Configurations



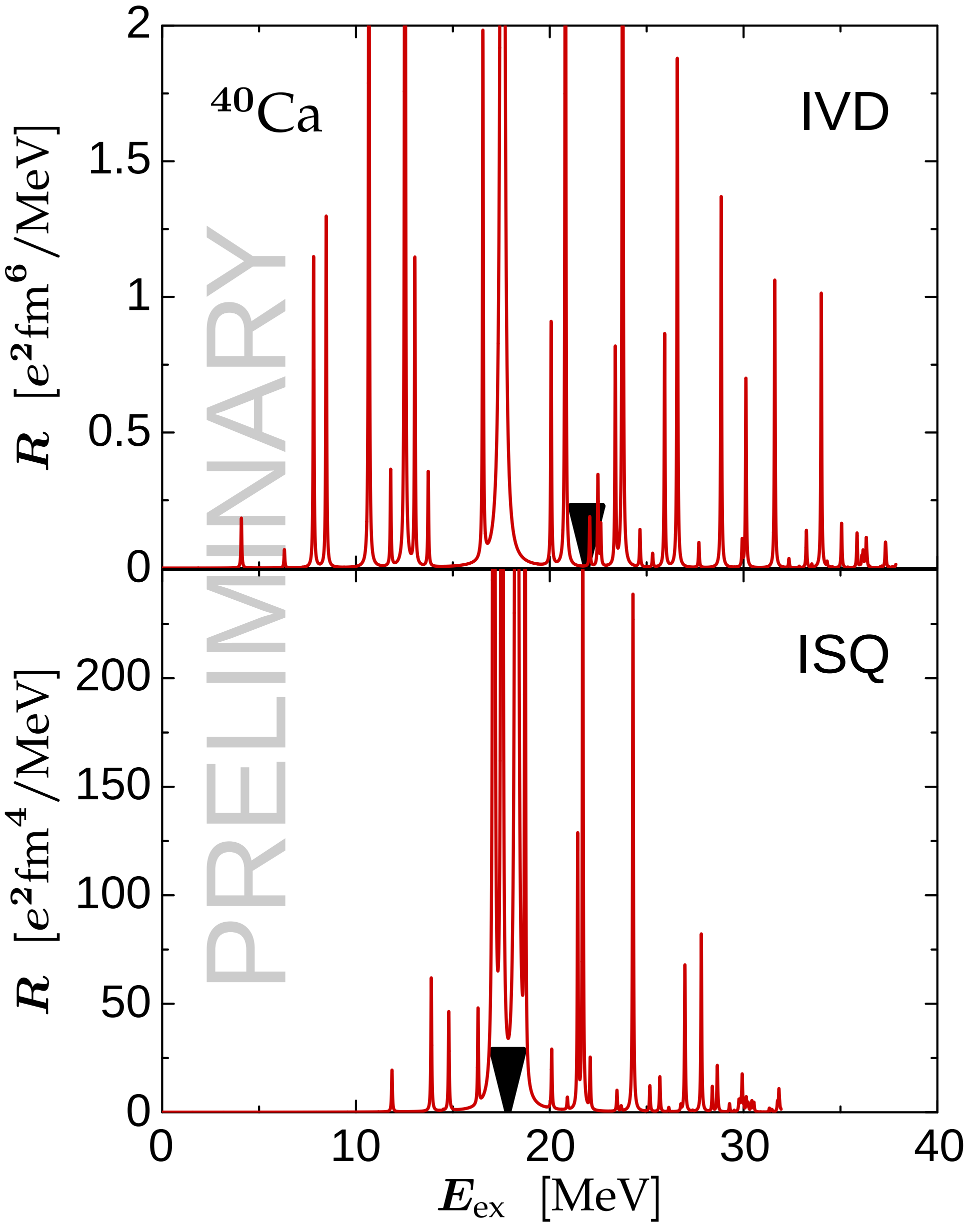
RPA (full 1p1h space)

SRPA (full 1p1h+2p2h space)

(11 major shells, $\Gamma = 2\,\mathrm{MeV}$)

complex configurations have significant impact on response

SRPA: Complex Configurations



RPA (full 1p1h space)

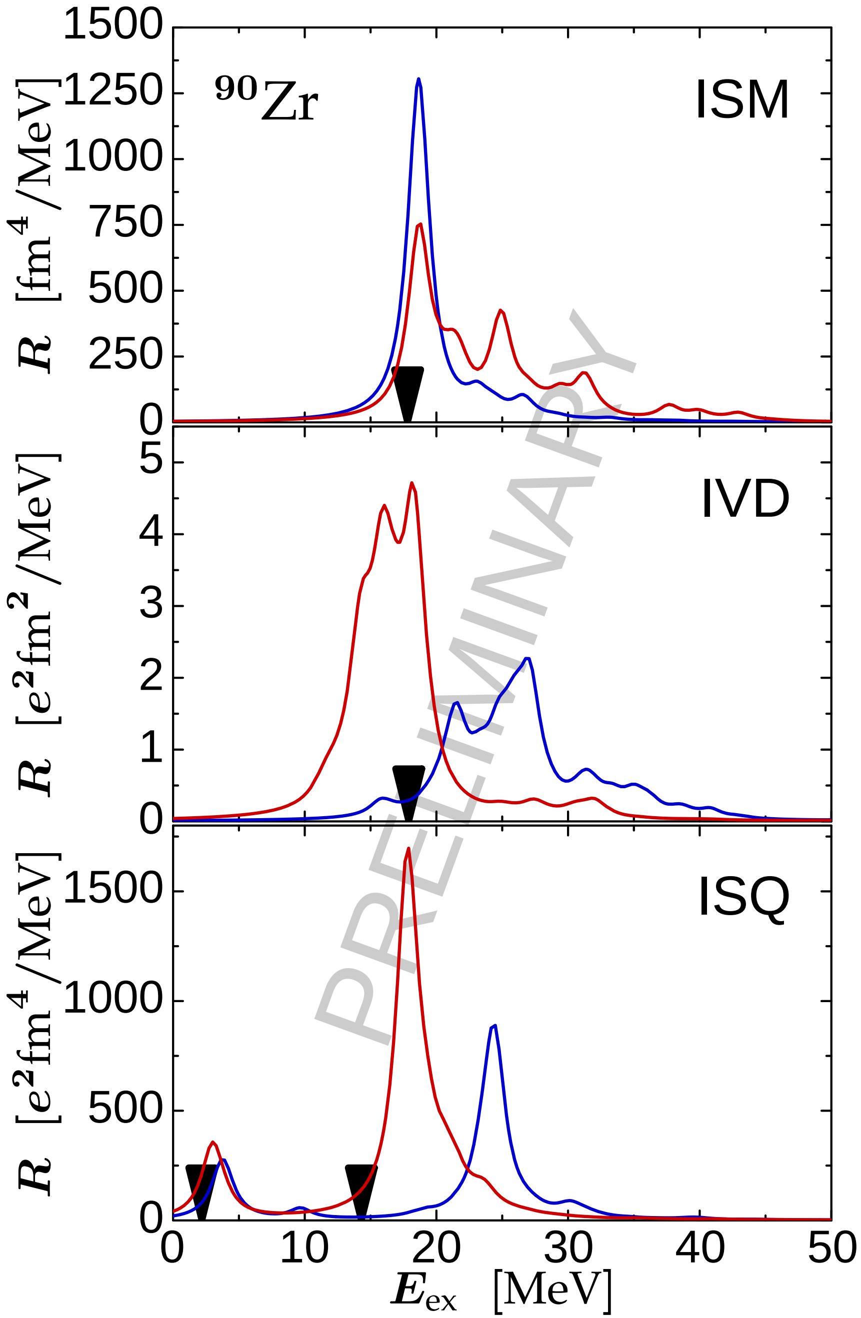
SRPA (full 1p1h+2p2h space)

(11 major shells, $\Gamma = 50\,\text{keV}$)

complex configurations have significant impact on response

possibility to investigate fine structure

Outlook: RPA with Three-Body Forces



- long-range tensor correlator & **repulsive three-body contact interaction**
- systematic improvement of
  - → rms-radii
  - → single-particle spectra
  - → strength distributions

standard V_{UCOM}

V_{UCOM} with long-range tensor & three-body contact interaction