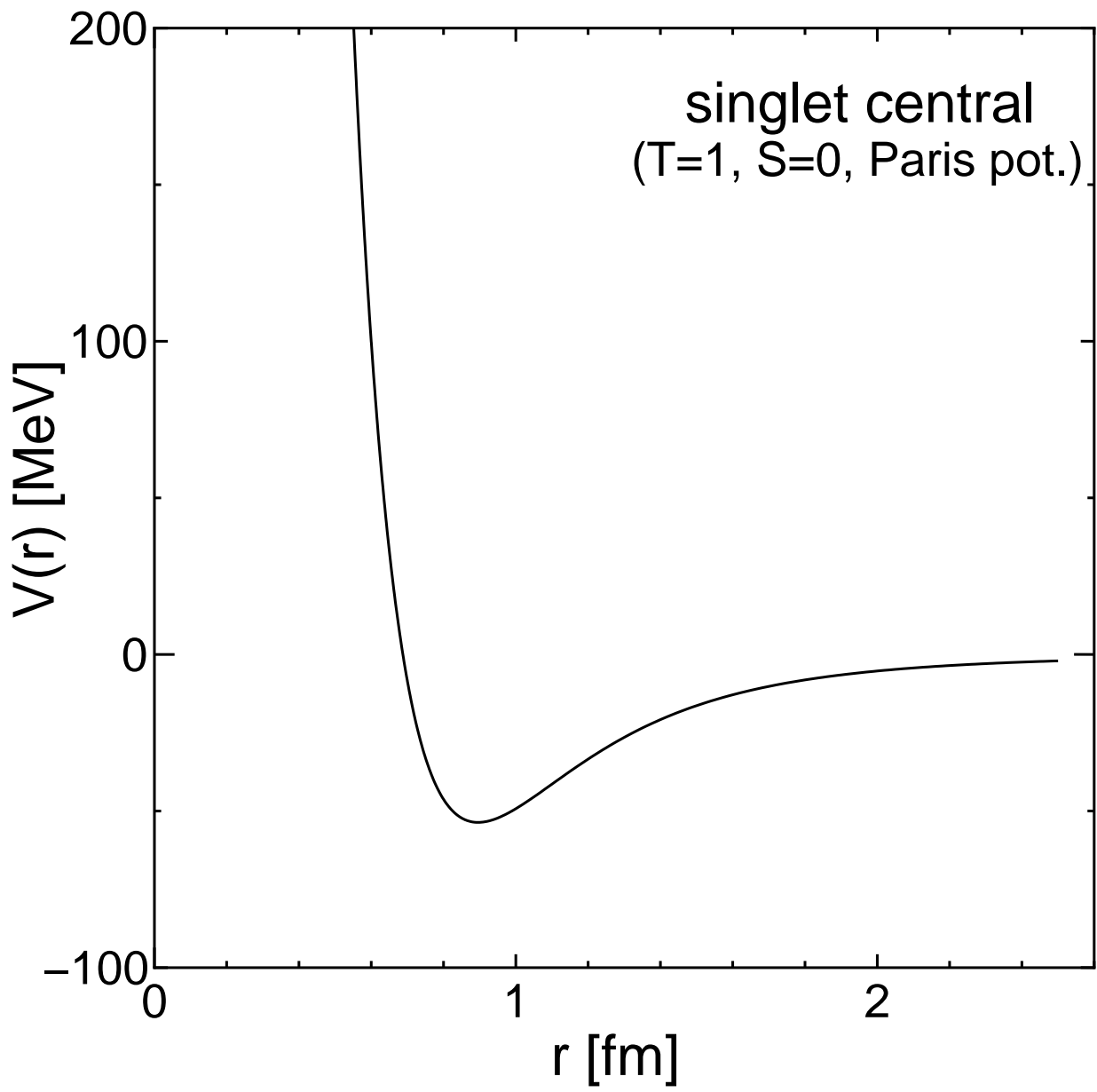


# Microscopic nuclear structure calculations with unitarily transformed effective interactions

Shinichiro Fujii (Kyushu Univ.)

1. Unitary transformation
2. Unitary-model-operator approach (UMOA)
3. “No-core” shell model
4. Results
  - g.s. energies, s.p. levels in  $p$ - and  $sd$ -shell nuclei
  - neutron-rich C isotopes ,  $^{18}\text{O}$
5. Summary



# Derivation of effective interaction (Hamiltonian) by means of unitary transformation

## Hamiltonian

$$H = H_0 + V$$

## Unitary transformation of $H$

$$\tilde{H} = U^{-1}HU$$

$$U = e^S, \quad (S : \text{anti-Hermitian}, S^\dagger = -S)$$

## Decoupling equation

$$Q(e^{-S}He^S)P = 0$$

## Solution

$$S = \text{arctanh}(\omega - \omega^\dagger), \quad \omega = Q\omega P$$

(with the restrictive condition  $PSP = QSQ = 0$ )

K. Suzuki, Prog. Theor. Phys. **68** (1982), 246

## Effective Hamiltonian

$$\underline{H_{\text{eff}} = P\tilde{H}P}$$

## Effective interaction

$$\underline{V_{\text{eff}} = P\tilde{H}P - PH_0P}$$

## Unitary transformation operator $U$ in terms of $\omega$

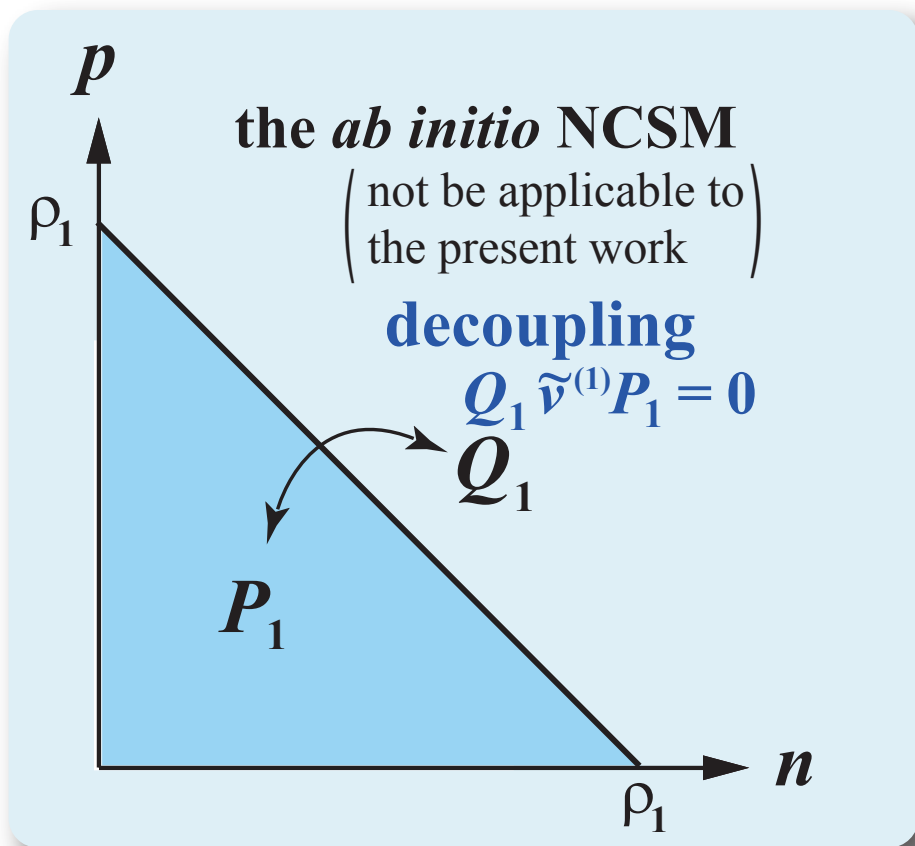
$$U = (1 + \omega - \omega^\dagger)(1 + \omega^\dagger\omega + \omega\omega^\dagger)^{-1/2}$$

$$= \begin{pmatrix} P(1 + \omega^\dagger\omega)^{-1/2}P & -P\omega^\dagger(1 + \omega\omega^\dagger)^{-1/2}Q \\ Q\omega(1 + \omega^\dagger\omega)^{-1/2}P & Q(1 + \omega\omega^\dagger)^{-1/2}Q \end{pmatrix}$$

S. Ōkubo, Prog. Theor. Phys. **12** (1954), 603

# Derivation of effective interaction

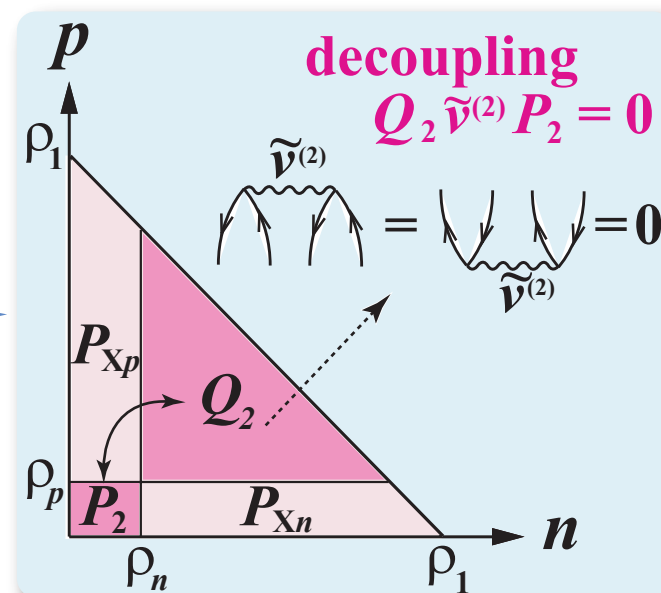
## • Eff. int. in a huge model space



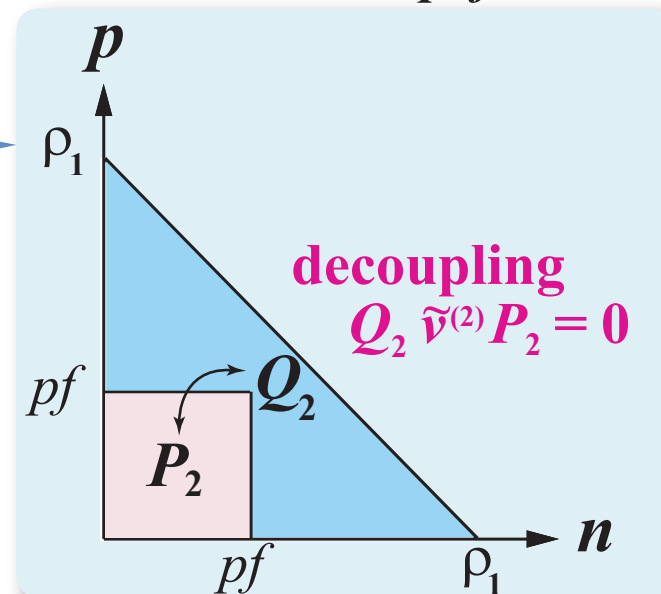
$\rho_1 = 2n_a + l_a + 2n_b + l_b$  ( $\{n_a, l_a\}$  and  $\{n_b, l_b\}$ : sets of h.o. quantum numbers of the two-body states)

- ( For details,
- S. F., T. Mizusaki, T. Otsuka, T. Sebe, and A. Arima, Phys. Lett. **B650**, 9 (2007).
  - S. F., R. Okamoto, and K. Suzuki, Phys. Rev. C **69**, 034328 (2004).

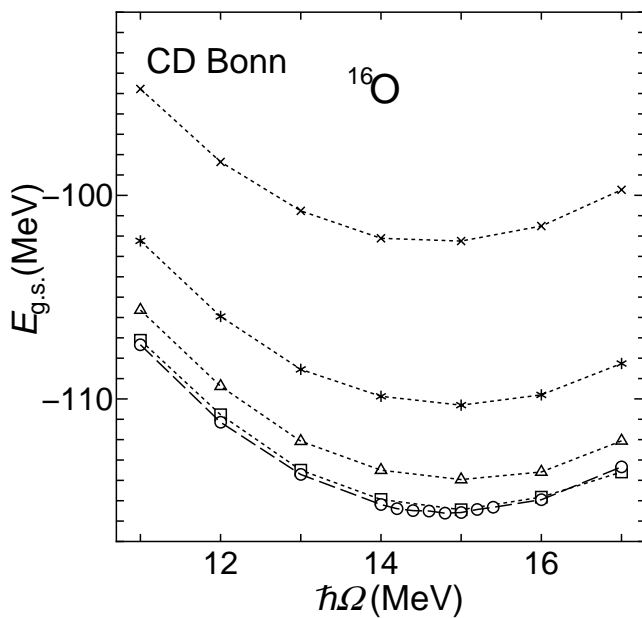
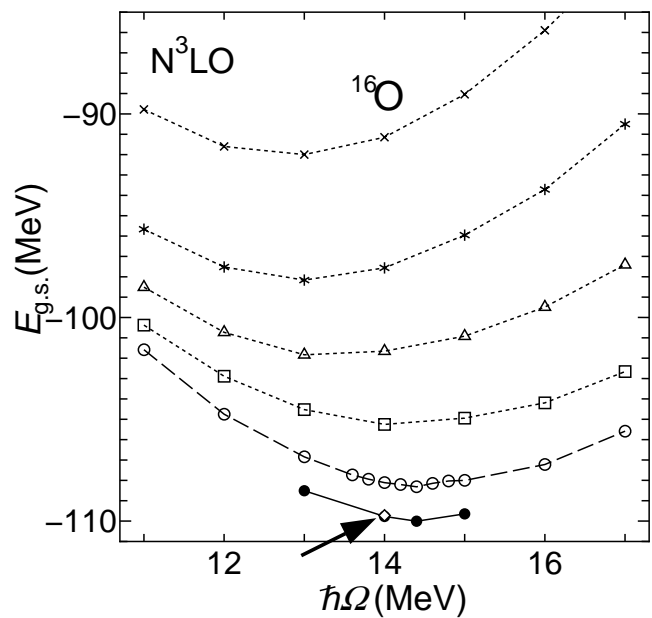
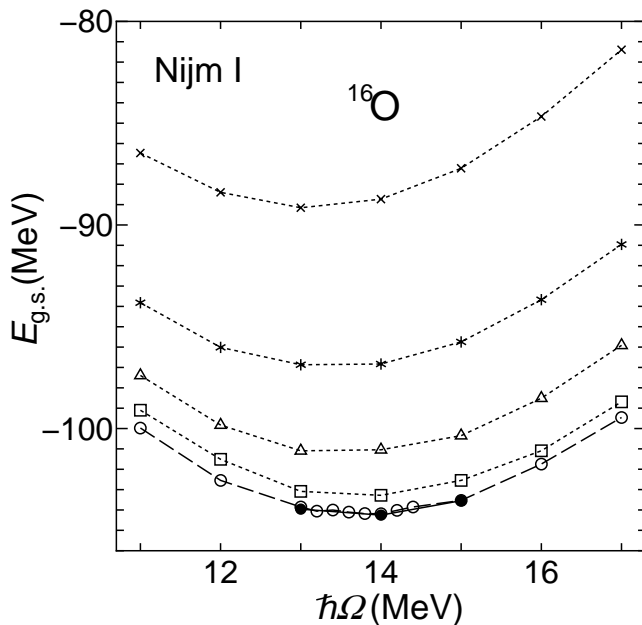
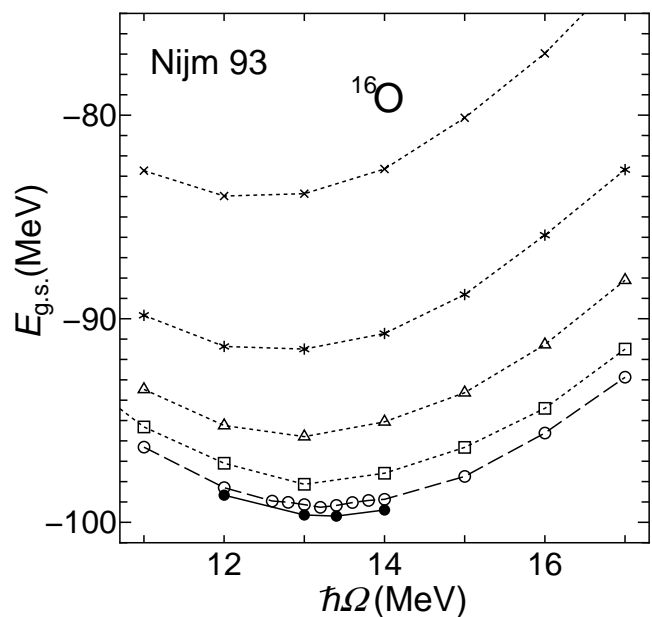
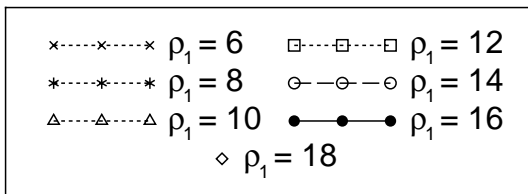
## • Eff. int. for the UMOA



## • Eff. int. for the $0s-1p0f$ shell model

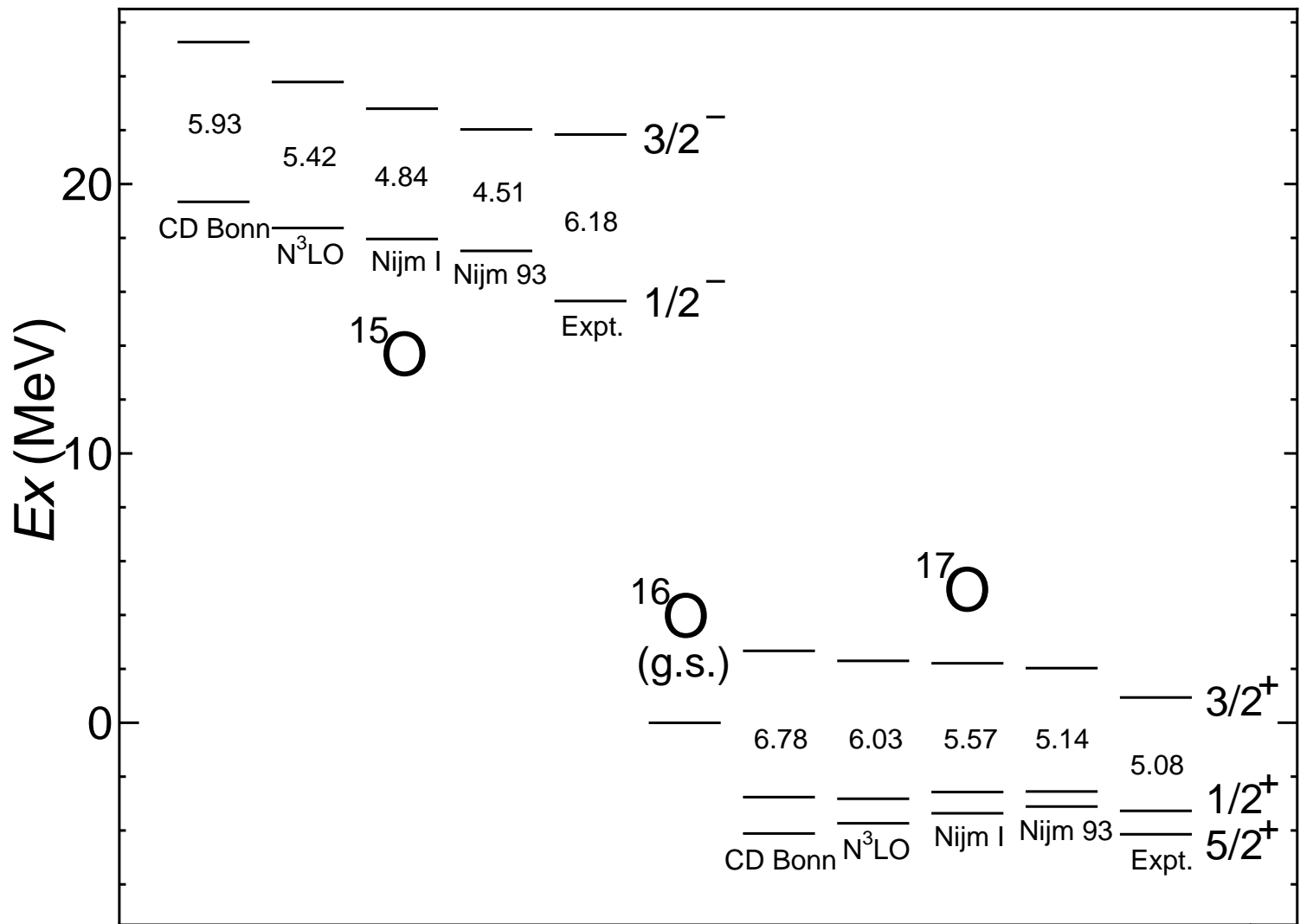


# Ground-state energies of $^{16}\text{O}$

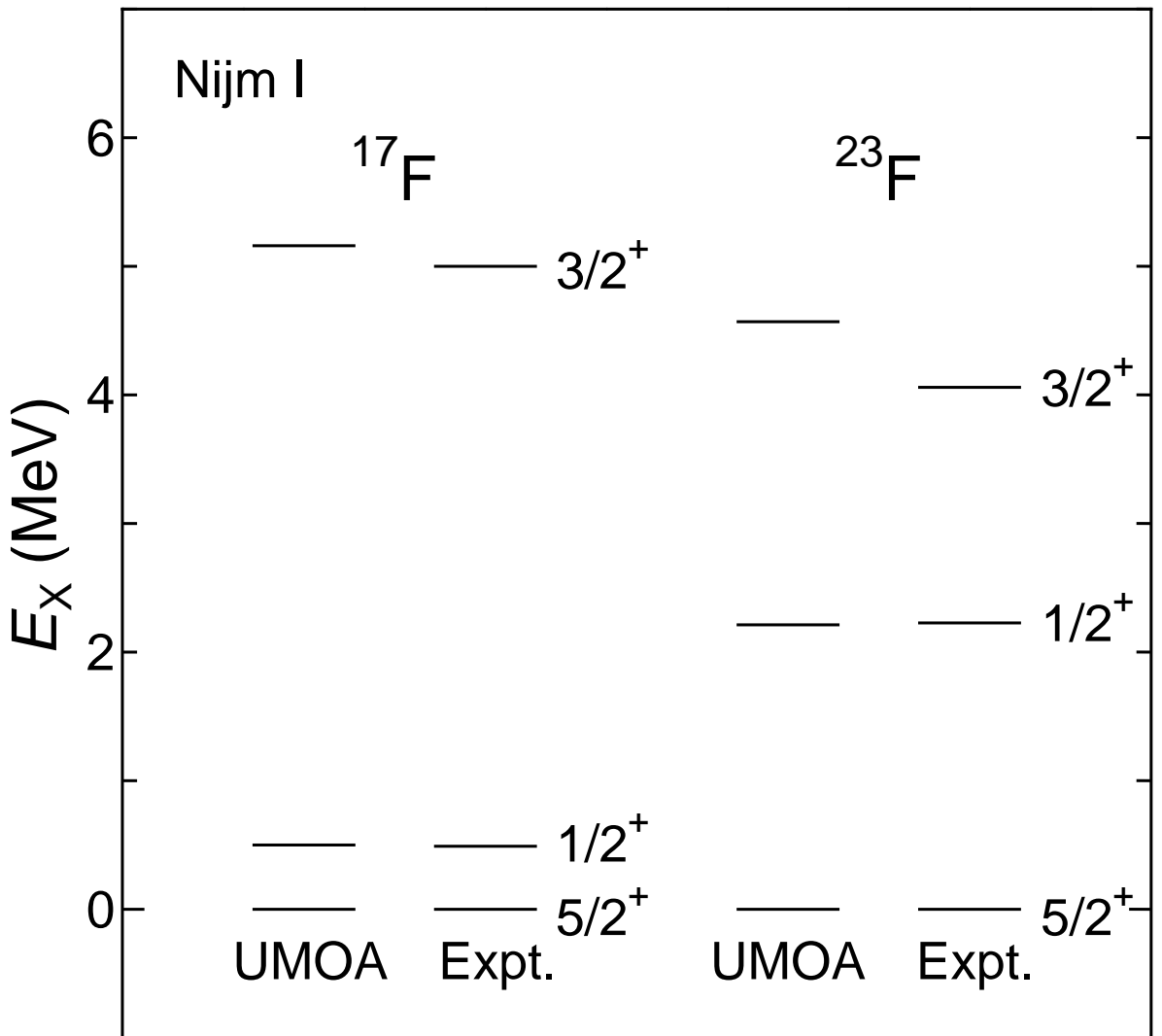


	Nijm 93	Nijm I	$\text{N}^3\text{LO}$	CD Bonn	Expt.	
$E_{g.s.}$	-99.69	-104.25	-110.00	-115.62	-127.62	
BE/A	6.23	6.52	6.88	7.23	7.98	(in MeV)

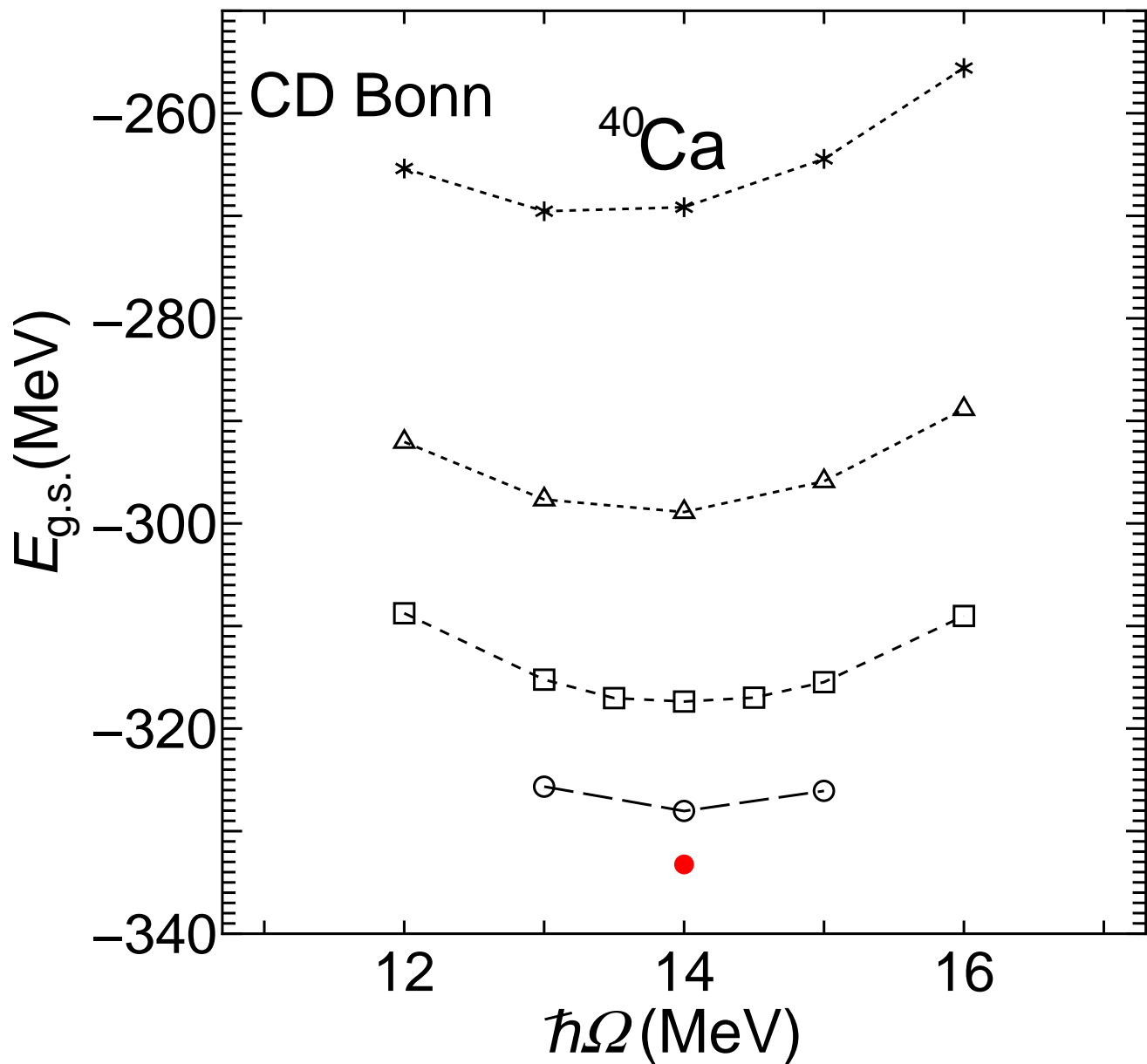
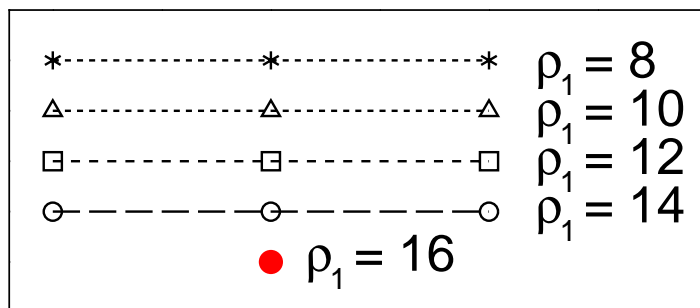
# Comparison of Expt. and UMOA results from modern NN interactions



# Single-particle levels in $^{17}\text{F}$ and $^{23}\text{F}$

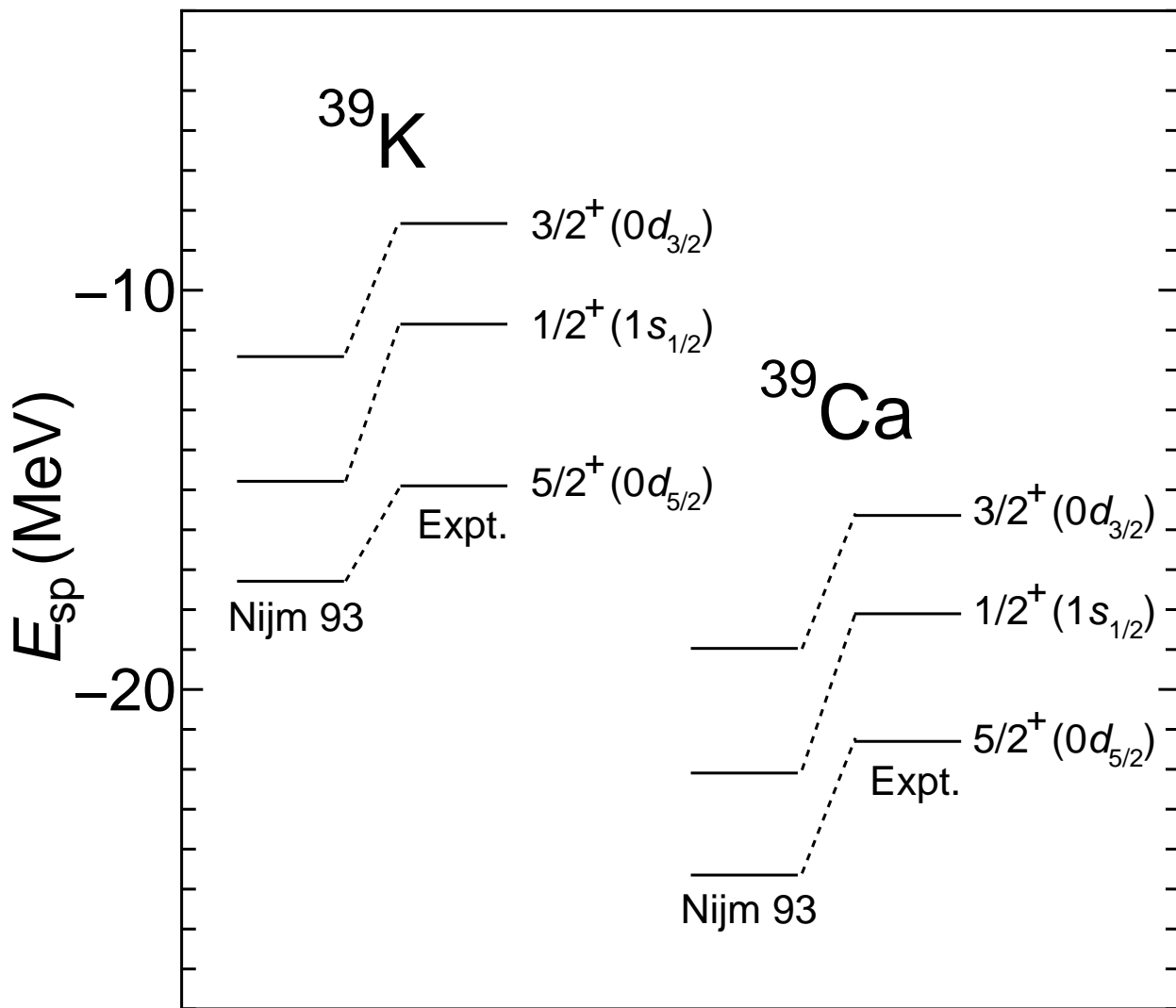


The  $\rho_1$  and  $\hbar\Omega$  dependences of calculated ground-state energies of  $^{40}\text{Ca}$





# Single-particle energies for hole states in $^{40}\text{Ca}$



# New approach to neutron-rich C isotopes

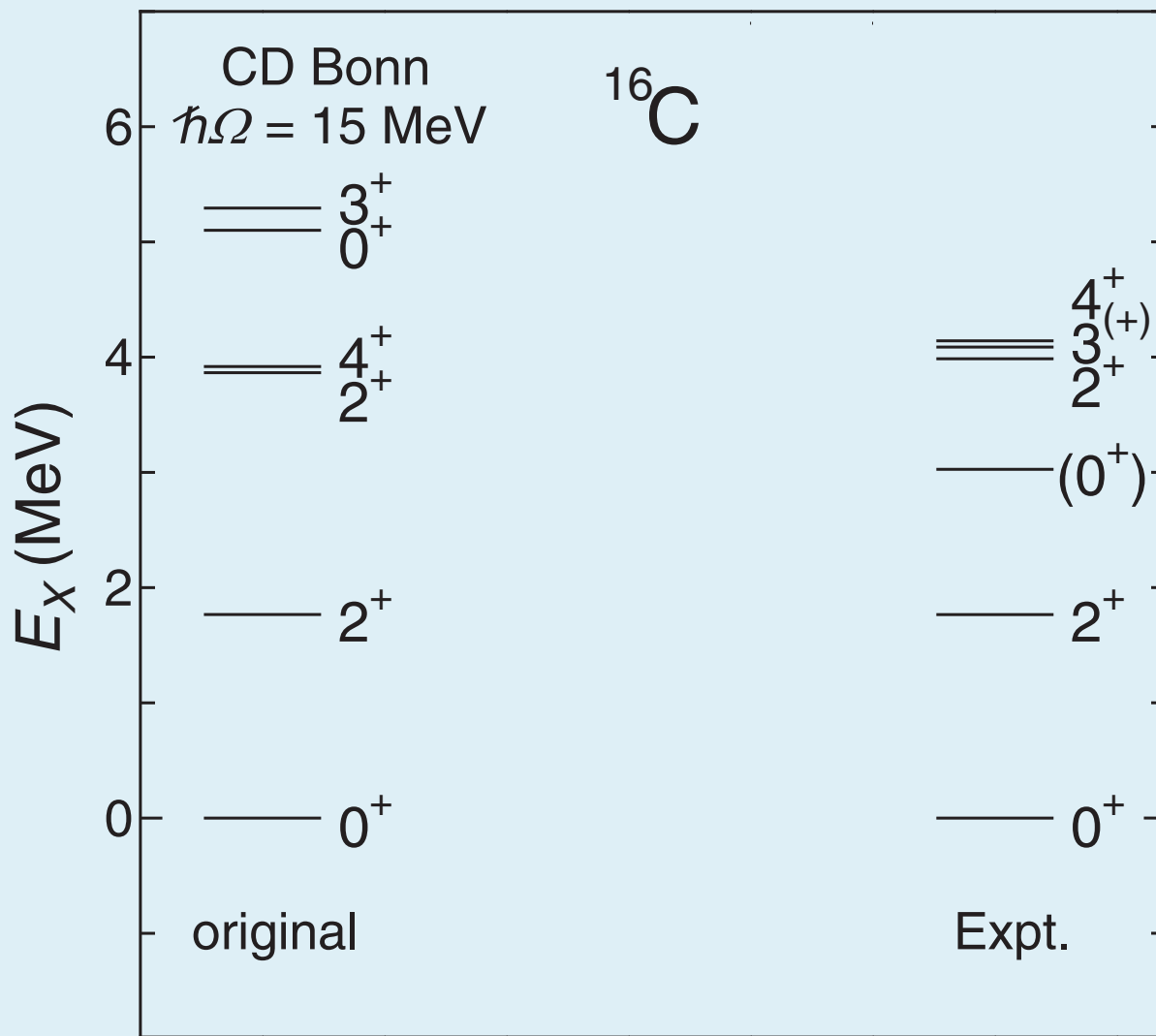
- **Large-scale shell model**

- **Code:** newly developed version of MSHELL
- **Model space:** the  $0s - 1p0f$  shells
- **Nucleon excitation:** up to 2 nucleons from the occupied shells for  $^{14}\text{C}$   
up to 2 nucleons to the  $1p0f$  shells
- **Bare transition operator**

- **Microscopic effective interaction**

Derived from a high-precision NN interaction (CD Bonn, ...) and the Coulomb force in the neutron-proton formalism for the given model space through a unitary-transformation theory

# Low-lying energy levels in $^{16}\text{C}$

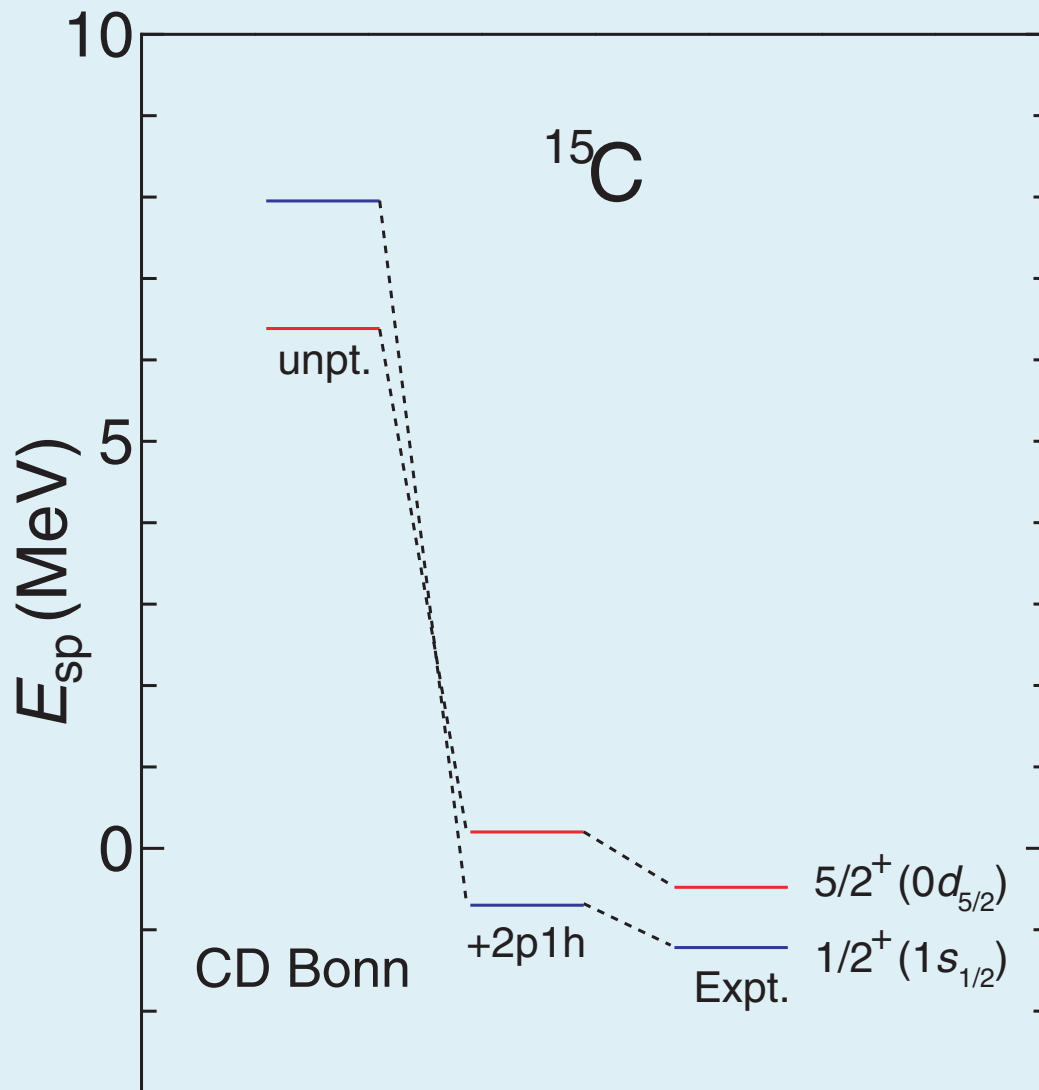


$B(E2; 2_1^+ \rightarrow 0_1^+) \text{ in } e^2\text{fm}^4$

1.30

$0.63_{\pm 0.11(\text{stat})}$   
 $\pm 0.16(\text{syst})$

## Single-particle energies in $^{15}\text{C}$



Calculated results by  
the unitary-model-operator approach  
(UMOA)

In the present shell model without any adjustable parameters

→ wrong ordering for the  $1/2^+$  and  $5/2^+$  states in  $^{15}\text{C}$  due to the *small* model-space size



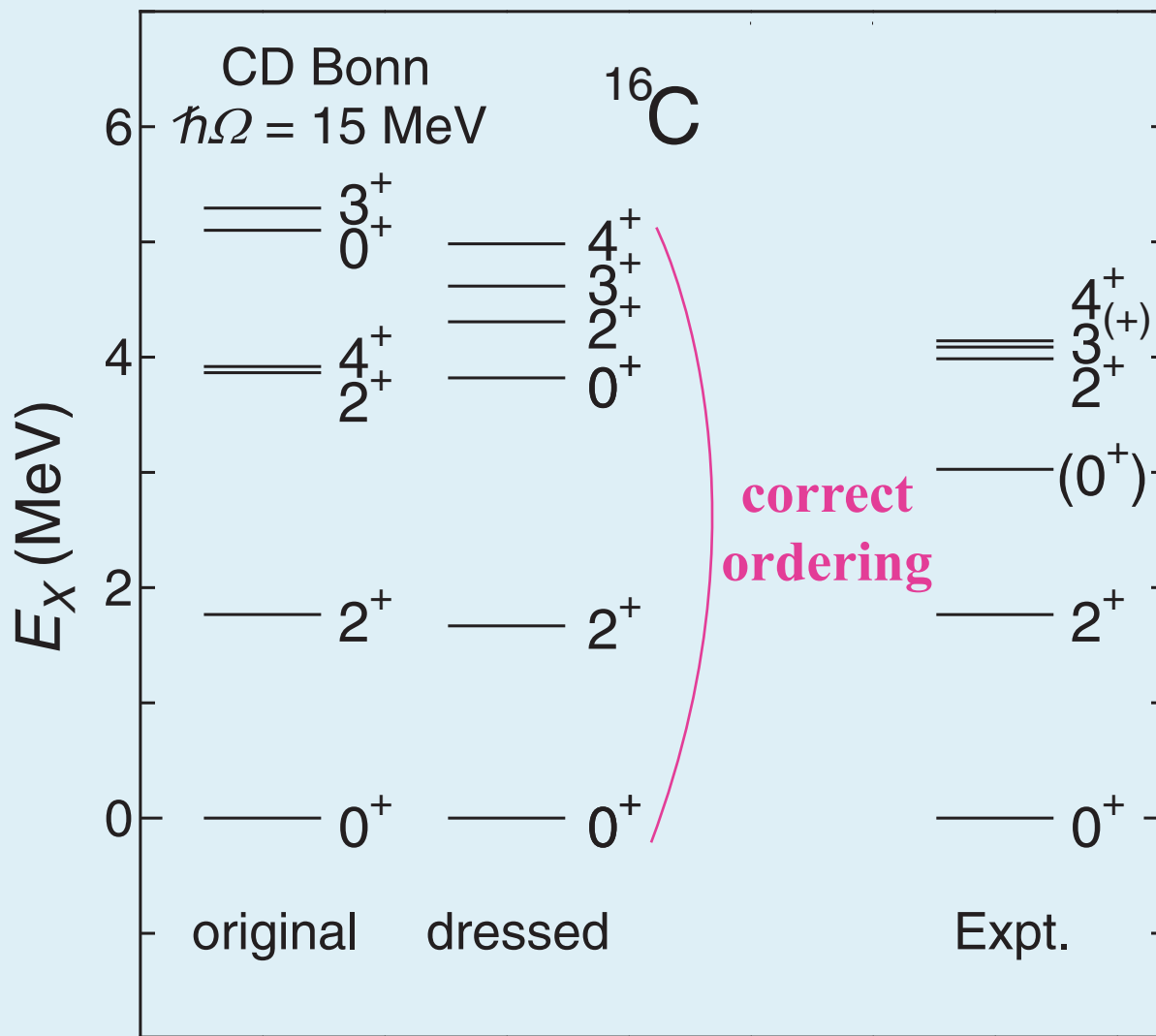
To remedy the wrong ordering and reproduce the binding energies for the  $1/2^+$  and  $5/2^+$  states of the UMOA results

→ introduce a minimal refinement of the one-body energies for the  $0d_{5/2}$  and  $1s_{1/2}$  orbits of the neutron



The calculated results are denoted by  
**"dressed"**

# Low-lying energy levels in $^{16}\text{C}$



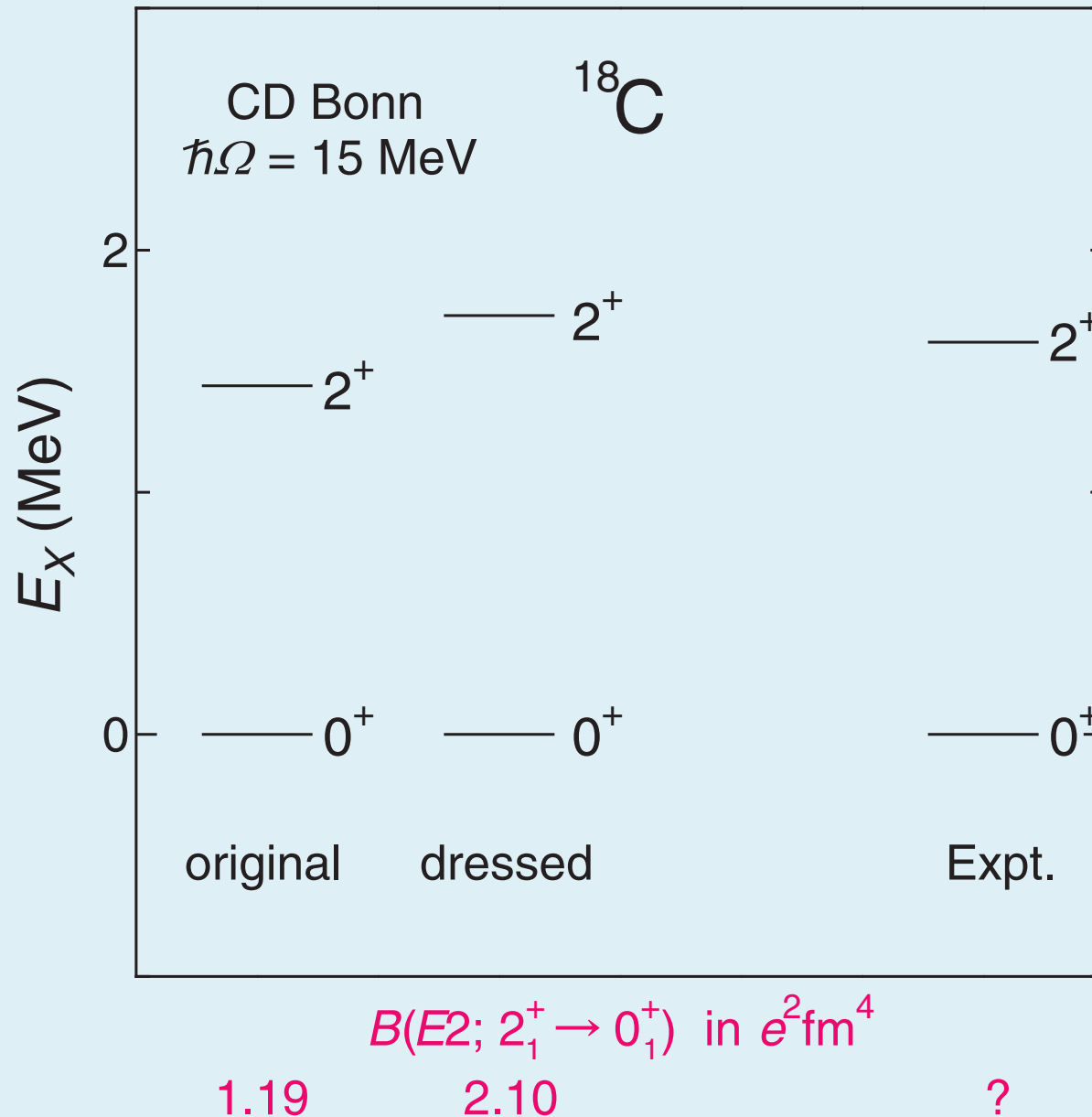
$B(E2; 2_1^+ \rightarrow 0_1^+)$  in  $e^2\text{fm}^4$

1.30

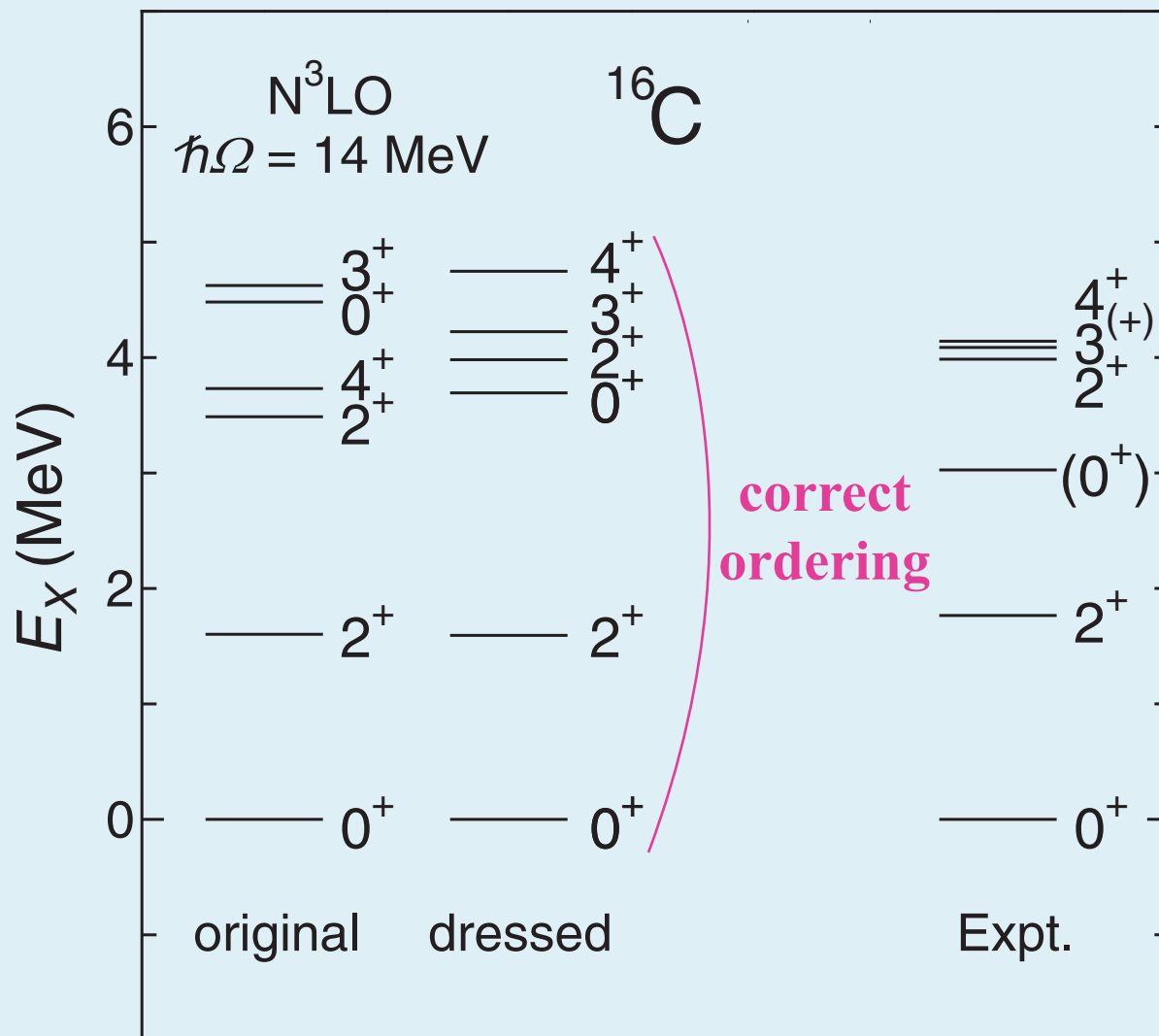
0.84

$0.63^{+0.11(\text{stat})}_{-0.16(\text{syst})}$

# Energy differences between the $2_1^+$ and $0_1^+$ states in $^{18}\text{C}$



# Low-lying energy levels in $^{16}\text{C}$



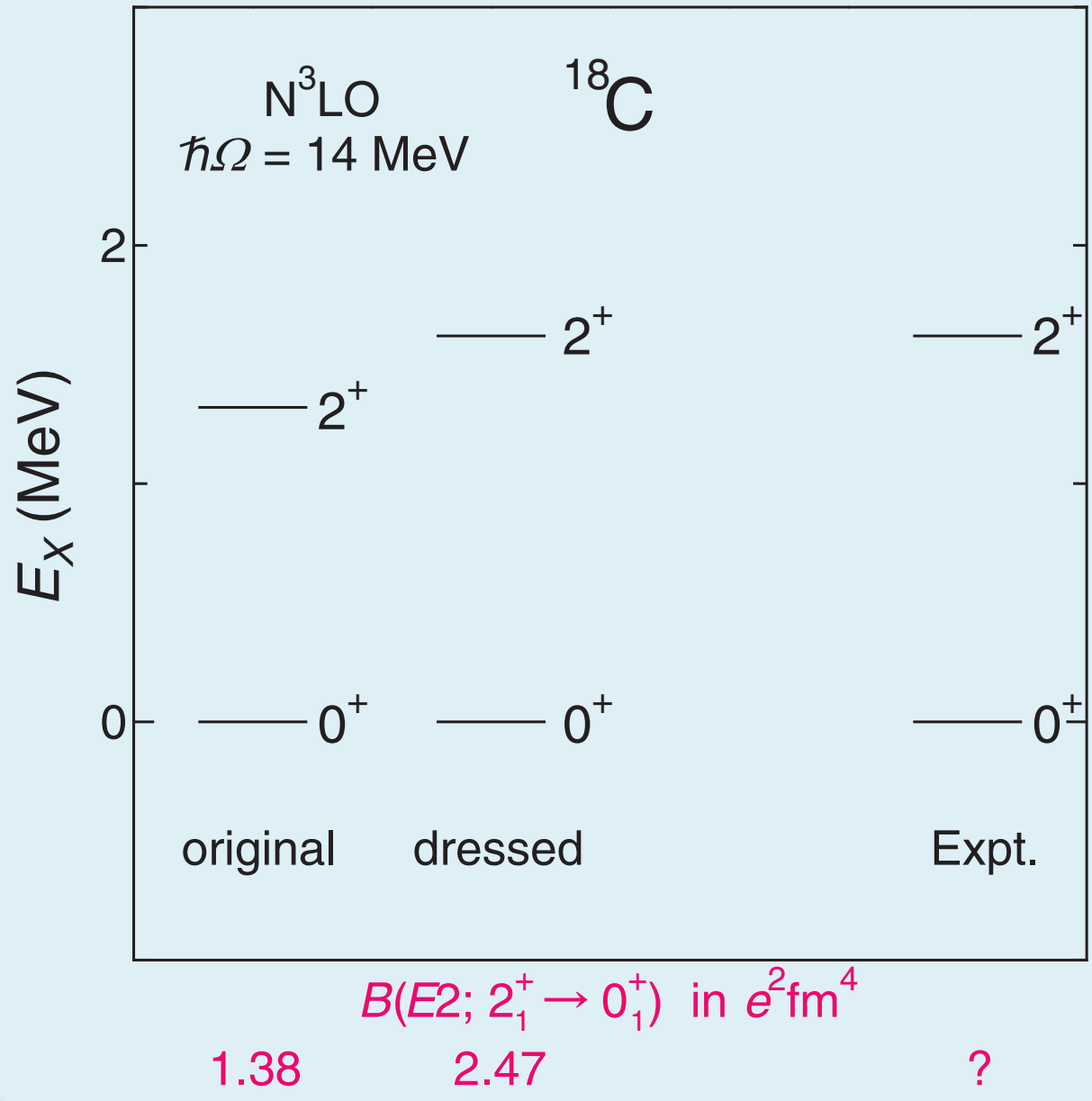
$B(E2; 2_1^+ \rightarrow 0_1^+)$  in  $e^2\text{fm}^4$

1.38

0.91

$0.63^{+0.11(\text{stat})}_{-0.16(\text{syst})}$

# Energy differences between the $2_1^+$ and $0_1^+$ states in $^{18}\text{C}$





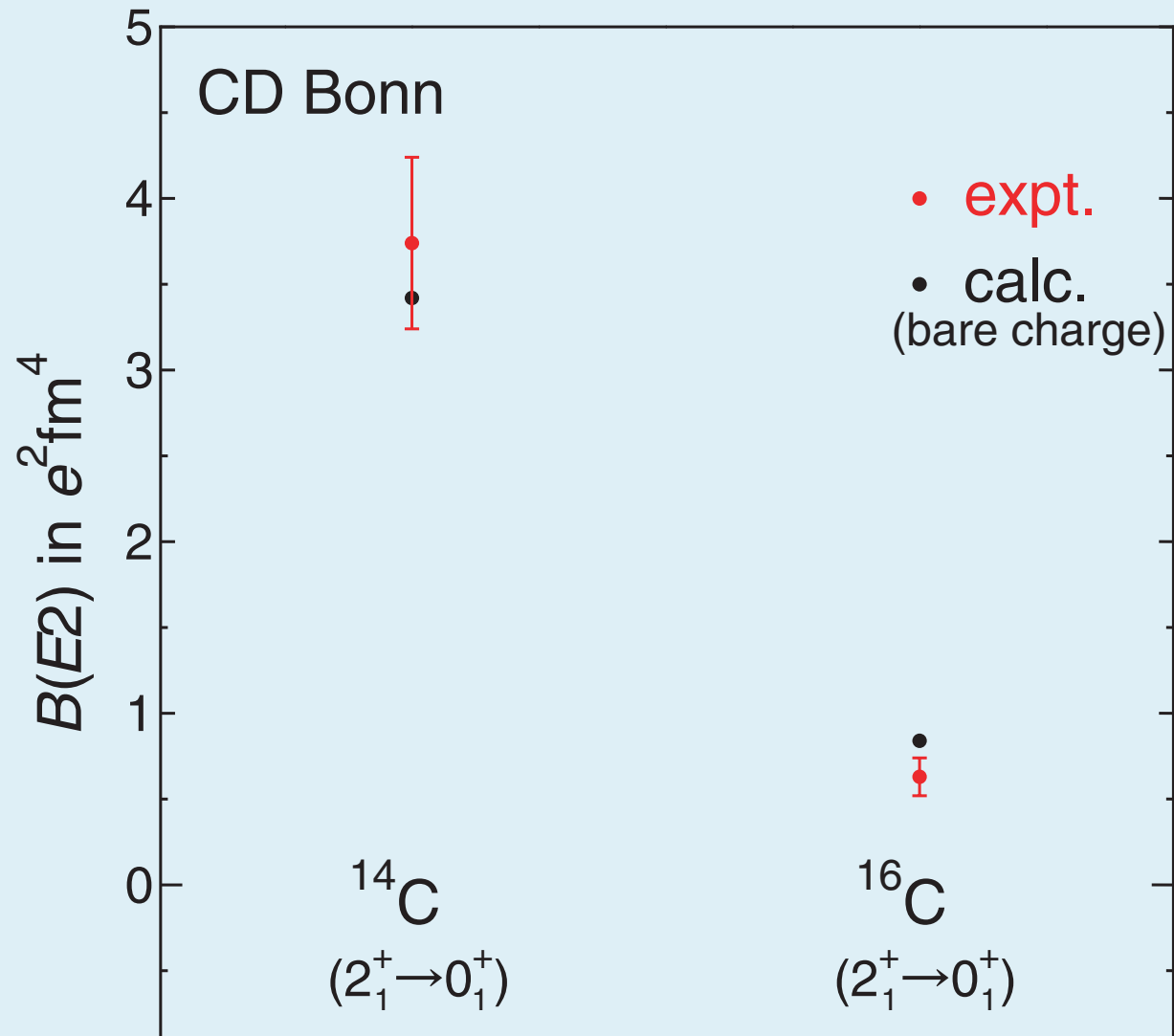
# $B(E2; 2_1^+ \rightarrow 0_1^+) \text{ in } ^{14,16,18}\text{C}$

for "dressed"

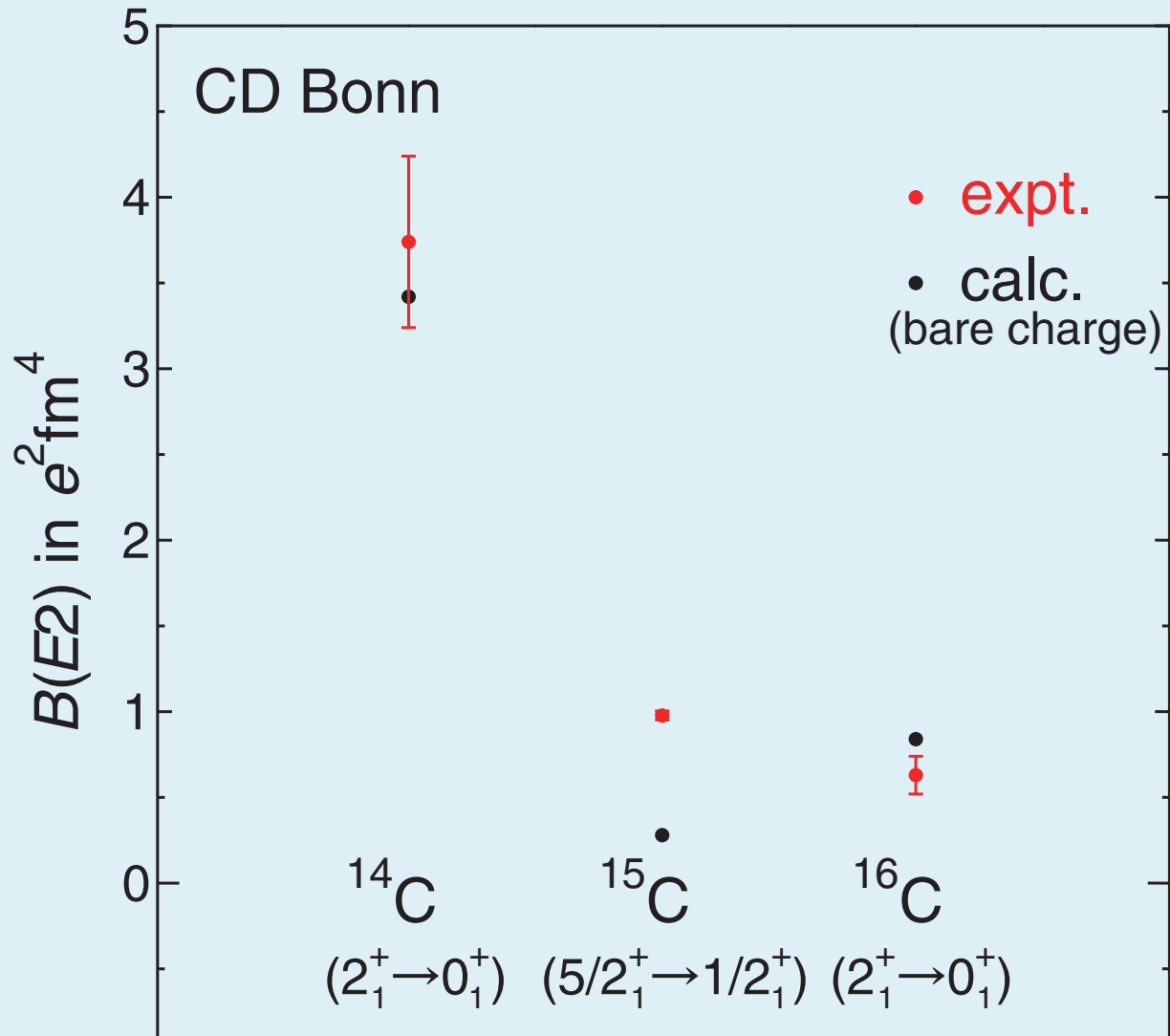
	CD Bonn	N <sup>3</sup> LO	Expt.
<sup>14</sup> C	3.42	4.11	3.74 ± 0.50
<sup>16</sup> C	0.84	0.91	0.63 <sup>± 0.11(stat)</sup> <sub>± 0.16(syst)</sub>
<sup>18</sup> C	2.10	2.47	?

in e<sup>2</sup>fm<sup>4</sup>

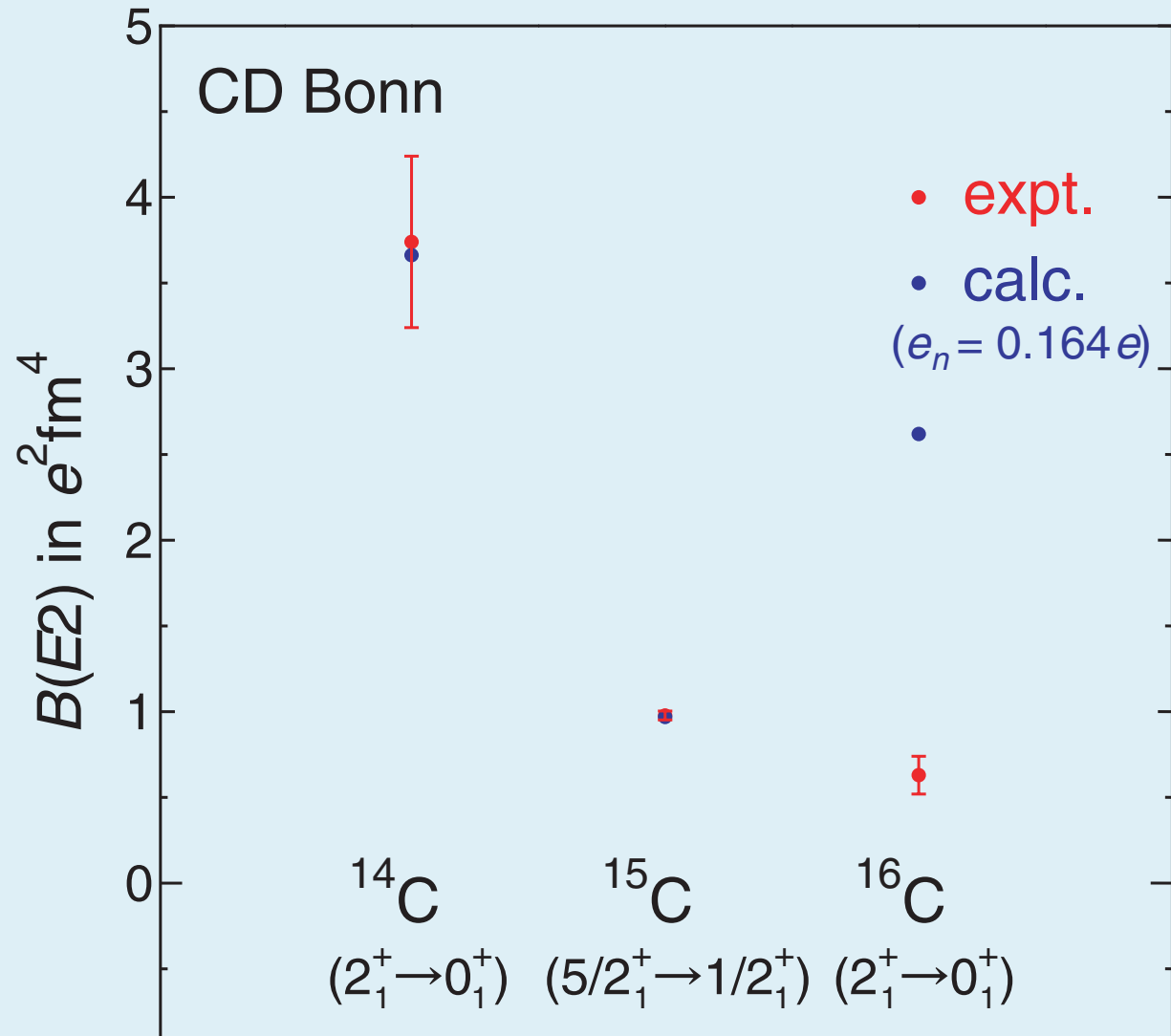
# $B(E2)$ of $^{14-16}\text{C}$



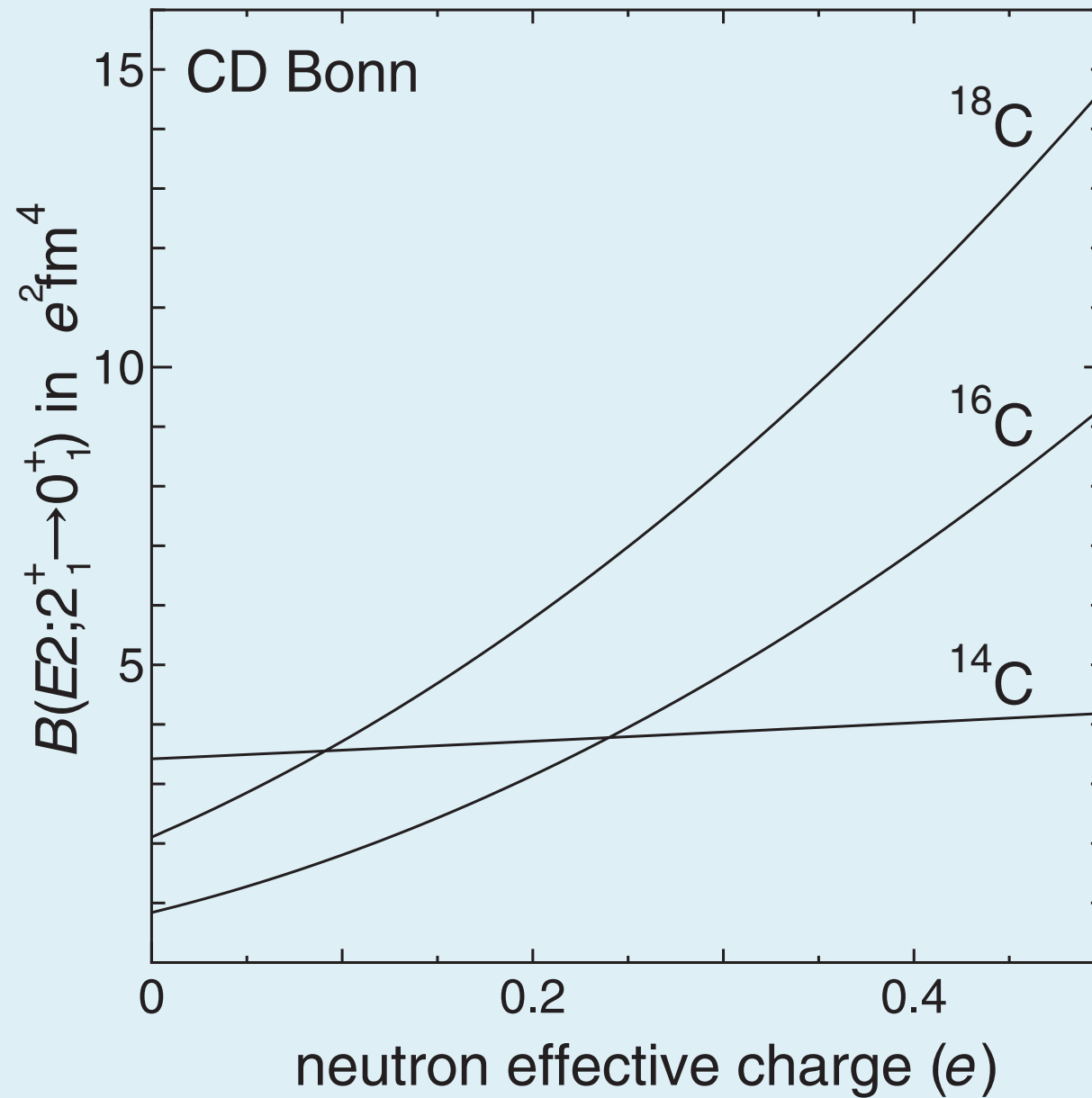
# $B(E2)$ of $^{14-16}\text{C}$



# $B(E2)$ of $^{14-16}\text{C}$



# Dependence of $B(E2)$ on the neutron effective charge



# $B(E2; 2_1^+ \rightarrow 0_1^+) \text{ in } ^{14,16,18}\text{C}$

CD Bonn

for "dressed"

	$e_n = 0$	$e_n = 0.164e$	Expt.
$^{14}\text{C}$	3.42	3.66	$3.74 \pm 0.50$
$^{16}\text{C}$	0.84	2.62	$0.63 \pm 0.11(\text{stat})$ $\pm 0.16(\text{syst})$
$^{18}\text{C}$	2.10	4.98	?

in  $e^2\text{fm}^4$

# $B(E2; 2_1^+ \rightarrow 0_1^+) \text{ in } ^{14,16,18}\text{C}$

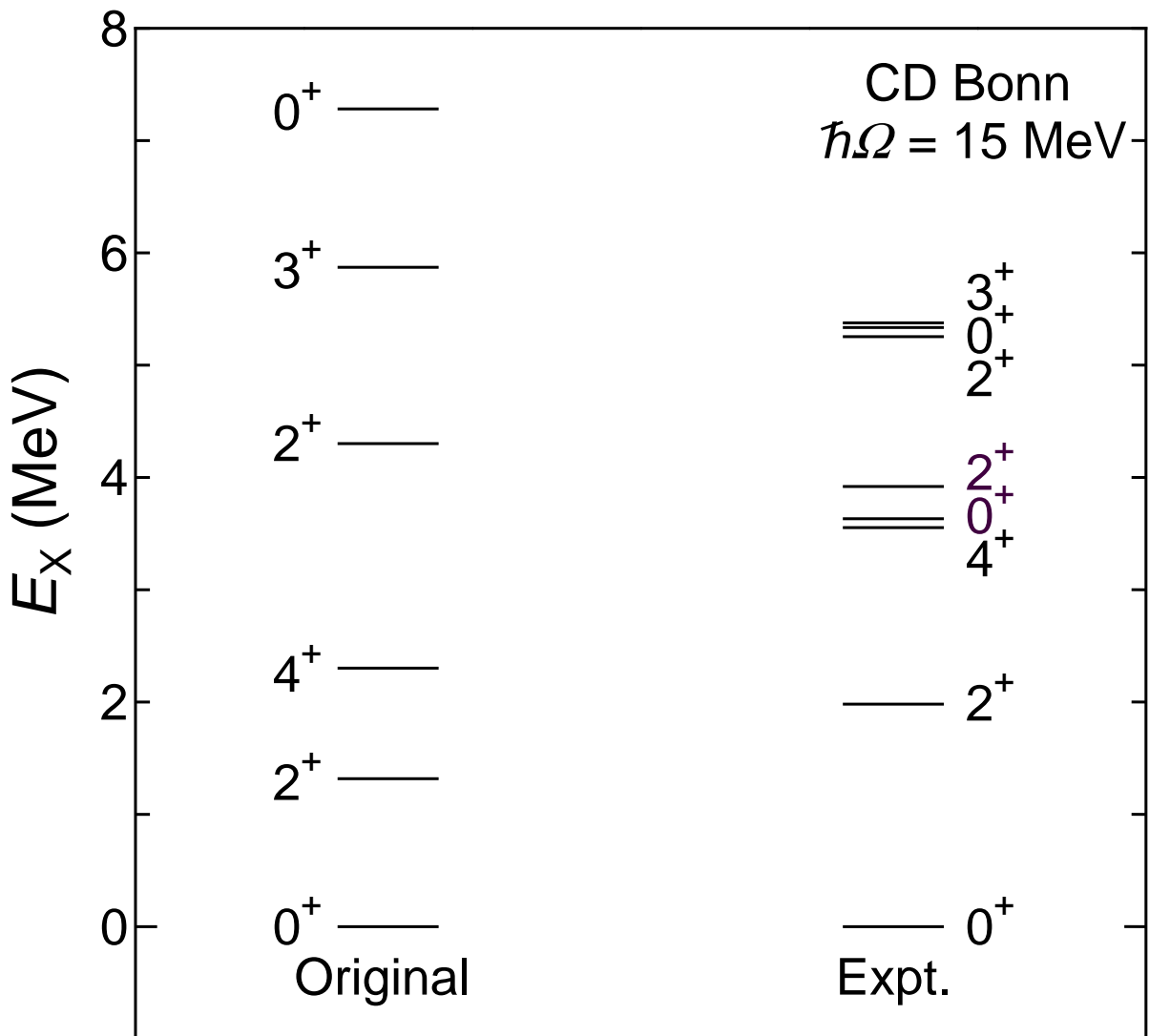
CD Bonn

for "dressed"

	$e_n = 0$	$e_n = 0.164e$	Expt.
$^{14}\text{C}$	3.42	3.66	$3.74 \pm 0.50$
$^{16}\text{C}$	0.84	2.62	$0.63 \pm 0.11(\text{stat}) \pm 0.16(\text{syst})$
$^{18}\text{C}$	2.10	4.98	$4.7^{+0.8}_{-0.6}(\text{stat})^{+1.4}_{-0.9}(\text{syst})$

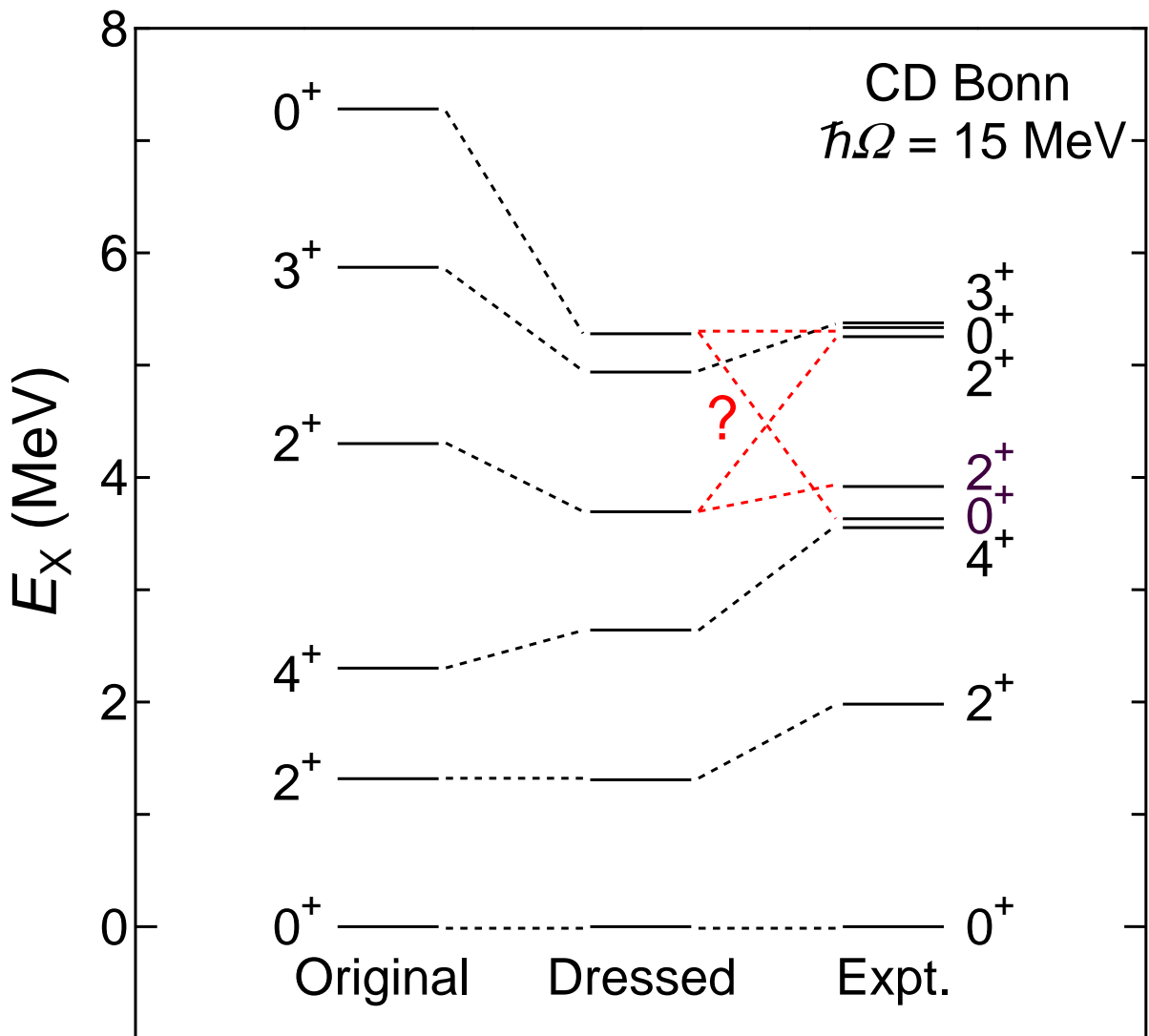
in  $e^2\text{fm}^4$

# Energy levels in $^{18}\text{O}$





# Energy levels in $^{18}\text{O}$



# Summary

- **Derived the effective interaction for a given model space from the high-precision modern NN force through a unitary transformation theory toward the microscopic description of structure of exotic nuclei**
  - **Unitary-model-operator approach (UMOA)**  
g.s. energy for closed-shell nuclei, s.p. energy, ...
  - **Large-scale shell model (with information about s.p. states obtained by the UMOA)**  
complicated structure, transition, ...
- **Including the genuine three-body force and diminishing the approximations in the calculation**

# Collaborators

## UMOA

**Ryoji Okamoto (Kyushu Inst. of Tech.)**

**Kenji Suzuki (Kyushu Inst. of Tech.)**

## **“No-core” shell model**

**Takahiro Mizusaki (Senshu Univ.)**

**Takaharu Otsuka (Univ. of Tokyo)**

**Takashi Sebe (Hosei Univ.)**

**Akito Arima (Japan Science Foundation)**

**Application of the “No-core” shell model to  $^{18}\text{O}$**

**Bruce R. Barrett (Univ. of Arizona)**