

Summary lecture 1

- Examined toy QM model with SUSY.
- Showed how naive discretization fails – lattice theory picks up new U.V contributions - fine tuning problem.
- Avoid by correcting action with counterterm computed at 1-loop.
- Generic way to deal with $D < 4$ theories (and maybe even $\mathcal{N} = 1$ in 4d with eg DWF)
- Better. Find modified action which is invariant under linear combination of SUSYs.
- New symmetry $Q^2 = 0$ and action Q-exact $S = Q\Lambda$.
Twisted formulation

Two dimensions

- How can we generalize previous construction to field theory ?
- In QM needed 2 SUSYs. Also single scalar field and two real fermions. Try simplest 2D with these features:
 $\mathcal{N} = 2$ Wess-Zumino model ...

$$S_{\text{WZ}} = \int d^2x \partial_\mu \phi \partial_\mu \bar{\phi} + W'(\phi) W'(\bar{\phi}) + \bar{\psi} \gamma_\mu \partial_\mu \psi + \\ + \bar{\psi} \left(\frac{1}{2} (1 + \gamma_5) W''(\phi) + \frac{1}{2} (1 - \gamma_5) W''(\bar{\phi}) \right) \psi$$

Nicolai Map

Fermion operator can be rewritten (chiral basis):

$$M_F = \begin{pmatrix} W''(\phi) & \partial_{\bar{z}} \\ \partial_z & W''(\bar{\phi}) \end{pmatrix}$$

Associated Nicolai Map is

$$\mathcal{N} = \partial_{\bar{z}}\bar{\phi} + W'(\phi)$$

Action $\mathcal{N}\mathcal{A}$ again differs by cross term from continuum (total derivative)

Lattice action

$$S_L = \sum_x \mathcal{N} \bar{\mathcal{N}} + \bar{\omega} (\Delta_z^s \lambda + W_L''(\phi) \omega) + \bar{\lambda} (\Delta_z^s \omega + W_L''(\bar{\phi}) \lambda)$$

where $\psi = \begin{pmatrix} \omega \\ \lambda \end{pmatrix}$, $\bar{\psi} = \begin{pmatrix} \bar{\omega} \\ \bar{\lambda} \end{pmatrix}$.

Symmetric difference $\Delta_z^s = \Delta_1^s + i\Delta_2^s$.

Doublers removed via:

$$W_L'(\phi) = W'(\phi) + \frac{1}{2} D_\mu^+ D_\mu^- \phi$$

Equivalent in 1D

Supersymmetries

The single exact SUSY follows by analogy to QM:

$$Q\phi = \omega$$

$$Q\bar{\phi} = \lambda$$

$$Q\lambda = 0$$

$$Q\omega = 0$$

$$Q\bar{\omega} = \bar{\mathcal{N}}$$

$$Q\bar{\lambda} = \mathcal{N}$$

Note that $Q^2 = 0$ with E.O.M

Homework Problem. Check this invariance.

Additional SUSYs

In continuum this model has 3 additional SUSYs.

Go back to QM:

Continuum bosonic action same if \mathcal{N} replaced by

$$\hat{N} = \Delta^- \phi - P'.$$

Equal to transpose of original fermion op.

Fermion action invariant if interchange $\psi \rightarrow \bar{\psi}$ – Gives 2nd SUSY!

Similar for WZ model: 3 other ways of writing \mathcal{N} – corresponding to $\phi \rightarrow \bar{\phi}$ and flipping sign.

Three new fermion ops. Make fermion act invariant by exchanging $\omega \rightarrow \lambda$ and $\bar{\omega} \rightarrow \bar{\lambda}$ etc.

Twisted/topological form

Not hard to verify that action can be rewritten (see QM) as

$$S_L = Q \sum_x \bar{\omega} \left(\mathcal{N} + \frac{1}{2} B \right) + \bar{\lambda} \left(\overline{\mathcal{N}} + \frac{1}{2} \overline{B} \right)$$

with additional B, \overline{B}

$$Q\bar{\omega} = \overline{B}$$

$$Q\bar{\lambda} = B$$

$$QB = 0$$

$$Q\overline{B} = 0$$

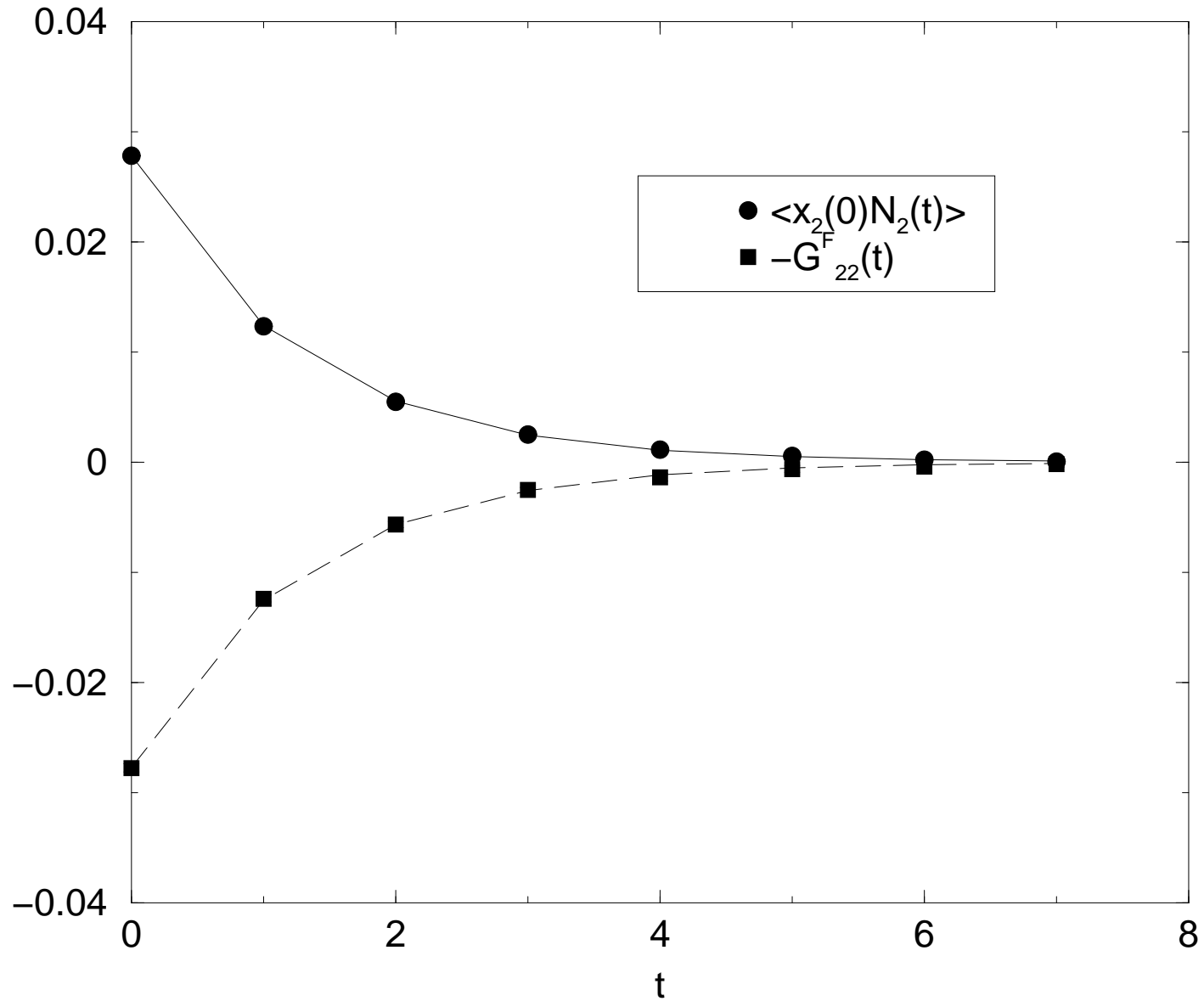
Ward identities

Again, existence of exact SUSY leads to Ward identities:

$$\langle S_B \rangle = \frac{1}{2} N_{\text{d.o.f}}$$

L	$\langle S_B \rangle$	$\frac{1}{2} N_{dof}$
4	31.93(6)	32
8	127.97(7)	128
16	512.0(3)	512
32	2046(3)	2048

WZ Ward identity



Twisting in 2D

Process of exposing this nilpotent supercharge can be made systematic

Fields transform as spinors under $SO(2)_R$ and $SO(2)_{\text{rot}}$

Choose to decompose fields under **twisted** symmetry

$$SO(2)' = \text{Diag} (SO(2)_R \times SO(2)_{\text{rot}})$$

Regard spinors/supercharges as **matrices**

Decompose on products of gamma matrices

$$q = QI + Q_\mu \gamma_\mu + Q_{12} \gamma_1 \gamma_2$$

Scalars, vectors, tensors

Note $\text{Tr}(\Gamma^i \Gamma^j) = 2\delta^{ij}$ where $\Gamma^i = (I, \gamma_\mu, \gamma_5$

Algebra

Original supersymmetry algebra $\{q, q\} = \gamma_\mu p_\mu$

$$\begin{aligned}\{Q, Q\} &= \{Q_{12}, Q_{12}\} = \{Q, Q_{12}\} = \{Q_\mu, Q_\nu\} = 0 \\ \{Q, Q_\mu\} &= p_\mu \quad \{Q_{12}, Q_\mu\} = \epsilon_{\mu\nu} p_\nu\end{aligned}$$

Note that p_μ is Q -variation of something.

Makes it plausible that entire $T_{\mu\nu}$ is Q -exact.

(why?: $p_\mu p_\nu = Q\Lambda_\mu Q\Lambda_\nu = Q(\Lambda_\mu Q_\nu)$)

Hence **twisted theories have Q -exact actions!** (remember

$$T_{\mu\nu} = \frac{\delta S}{\delta g_{\mu\nu}})$$

Kähler-Dirac fermions

Typical Dirac action can be rewritten as:

$$S_F = \text{Tr} [\bar{\Psi} \gamma_\mu \partial_\mu \Psi]$$

where

$$\Psi = \frac{\eta}{2} I + \psi_\mu \gamma_\mu + \chi_{12} \gamma_5$$

Components $(\eta/2, \psi_\mu, \chi_{12})$ constitute a Kähler-Dirac field. In real (why real?) components:

$$S_F = \frac{1}{2} \eta \partial_\mu \psi_\mu + \chi_{12} (\partial_1 \psi_2 - \partial_2 \psi_1)$$

Kähler-Dirac action. Basis indep. Geometrical treatment of fermions.

KD treatment equivalent to twisting

Discretize KD equations

Can discretize such geometrical actions without introducing doubles (Rabin, Joos)

- Replace $\partial_\mu \rightarrow \Delta_\mu^+$ in curl
- Replace $\partial_\mu \rightarrow \Delta_\mu^-$ in div.

Prohibits doubles!

$$S_F = \begin{pmatrix} \eta/2 & \chi_{12} \end{pmatrix} \begin{pmatrix} \Delta_1^- & \Delta_2^- \\ -\Delta_2^+ & \Delta_1^+ \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

Relation to staggered fermions

Natural to place scalars on sites, vectors on links and tensors on diagonal links.

Introduce a lattice with half spacing – all fields now site fields.

Free KD action reduces to staggered action – fermion phases now arise from alternate use of Δ^- , Δ^+ .

Usual U(1) symmetry: rotate $\psi_{1,2} \rightarrow e^\alpha \psi_{1,2}$ and

$\eta, \chi_{12} \rightarrow e^{-\alpha} \eta, \chi_{12}$.

Homework problem 4. Check this

Wess-Zumino using KD fermions

If we choose

$$\begin{aligned}\omega &= \frac{1}{2}\eta + i\chi_{12} \\ \lambda &= \psi_1 + i\psi_2\end{aligned}$$

Find KD action

$$\omega^\dagger D_{\bar{z}}\lambda + \text{h.c.}$$

Technically uses **self-dual** parts of KD field

Duality $f_p \xrightarrow{*} f_{d-p}$ with

$$*f_{\mu_1\dots\mu_{d-p}} = \epsilon_{\mu_1\dots\mu_d} f_{\mu_{d-p+1}\dots\mu_d}$$

$$P^+ = \frac{1}{2}(I + i*)$$