



Stefano Profumo
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Theoretical AstroPhysics Including Relativity

Kellogg Rad Lab

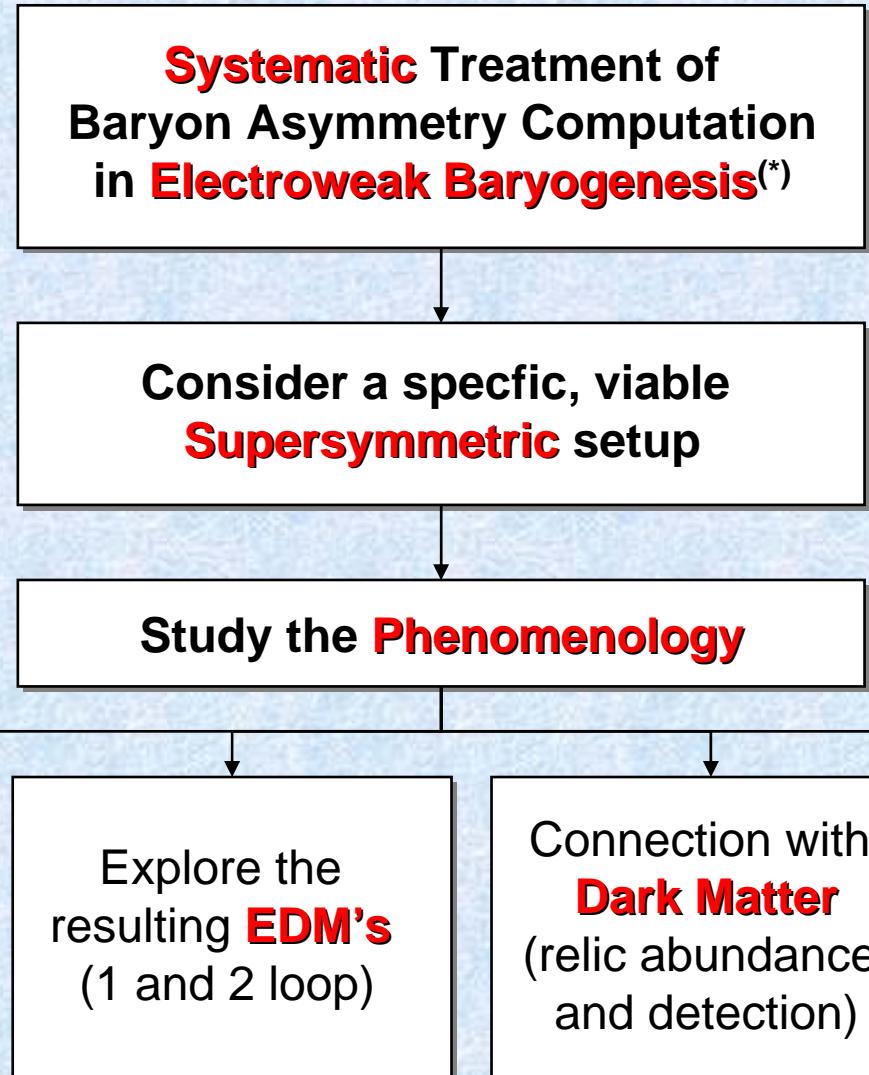
Probing Supersymmetric Baryogenesis: from Electric Dipole Moments to Neutrino Telescopes

Based on: V.Cirigliano, SP and M.Ramsey-Musolf, JHEP07(2006)002

*SP and M.Ramsey-Musolf (Caltech/Madison) [work in progress]
S.Ando (Caltech), V.Barger (Madison), SP, M.Ramsey-Musolf,
G.Shaughnessy (Madison) [work in progress]*

**INT, University of Washington
Seattle, WA, Thursday, March 22, 2007**

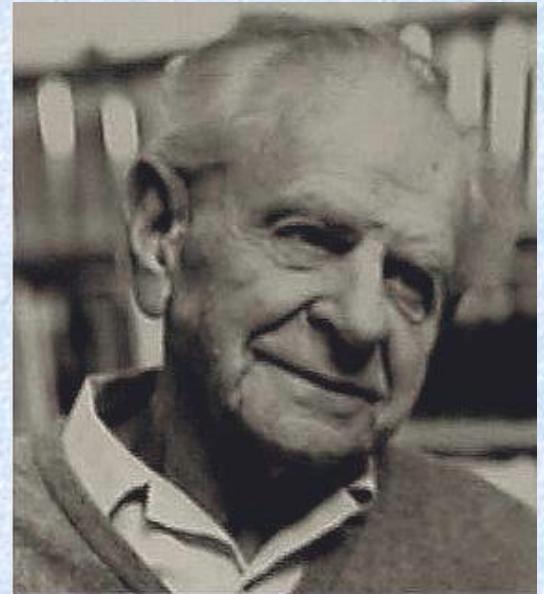
Phenomenology of SUSY EW Baryogenesis



^(*) Cirigliano, Lee, Ramsey-Musolf, Tulin; see C.Lee's and M.J.Ramsey-Musolf's talks

The Nightmare of Popper's Quote

“In so far as a scientific statement
speaks about reality,
it must be falsifiable:
And in so far as it is not falsifiable
it does not speak about reality”



Sir Karl Popper (1902-1994)

Karl Popper, “*The Logic of Scientific Discovery*”

**Supersymmetric Electro-Weak Baryogenesis:
a falsifiable theory**

Supersymmetric EW Baryogenesis

In the **Minimal Supersymmetric Extension of the Standard Model (MSSM)**

- Additional bosonic degrees of freedom couple to the Higgs (e.g. a light scalar top, x6)

*A strongly first order
EW Phase Transition
occurs for larger, LEP-viable
values of m_h*

- The theory features potential additional **CP -violating sources**

*- Gaugino/Higgsino sector
- Scalar quark sector*

BONUS: the MSSM provides ideal candidates for non-baryonic Dark Matter as well !

EWB in the MSSM: Requirements

In the **MSSM**, successful EW baryogenesis requires: ...at **odds** with:

1. Light enough stop

$$m_{\tilde{t}_1} < m_t$$

$$m_{\tilde{t}_1} > 95 \text{ GeV}$$

(LEP-II)

2. Light enough Higgs

$$m_h < 120 \text{ GeV}$$

$$m_h > 114 \text{ GeV}$$

(LEP-II)

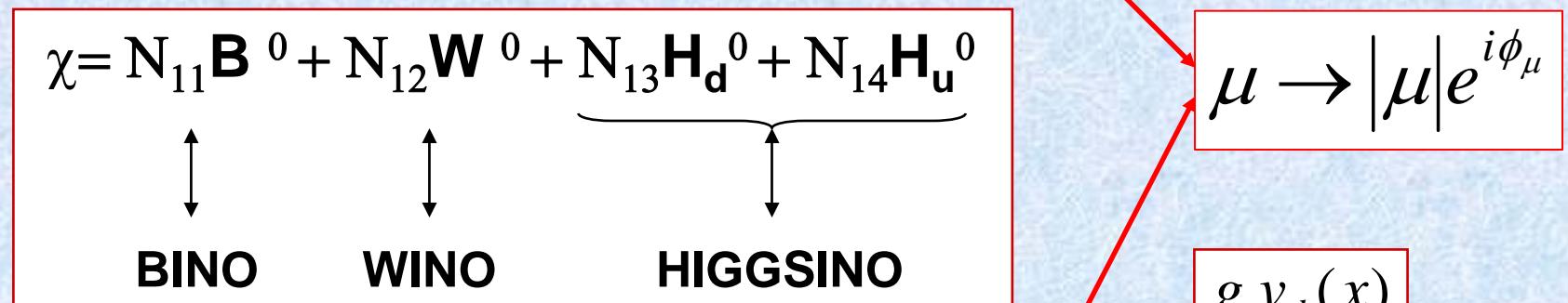
3. Strong enough CP-violating sources
(3rd generation squarks, higgsinos)

e, n and atomic
EDM's

Interludeum: Neutralinos and Charginos

$$M_N = \begin{pmatrix} M_1 & 0 & m_z \sin \beta \sin \theta_W & -m_z \sin \beta \sin \theta_W \\ 0 & M_2 & 0 & -\mu \\ -m_z \cos \beta \sin \theta_W & m_z \cos \beta \cos \theta_W & 0 & 0 \\ m_z \sin \beta \sin \theta_W & -m_z \sin \beta \sin \theta_W & -\mu & 0 \end{pmatrix}$$

Gaugino Mass Relations
($M_1 = \alpha M_2$)



$T \sim T_{EW}$: scattering
of \tilde{H}, \tilde{W} from
background field

$T \ll T_{EW}$: mixing
of \tilde{H}, \tilde{W} to $\tilde{\chi}^+, \tilde{\chi}^0$

$$M_C = \begin{pmatrix} M_2 & m_w \sqrt{2} \cos \beta \\ m_w \sqrt{2} \sin \beta & \mu \end{pmatrix}$$

$\frac{g v_d(x)}{\sqrt{2}}$

$\frac{g v_u(x)}{\sqrt{2}}$

EWB in the MSSM: Setup

- **Free parameters** in the game:

$$M_1, M_2, |\mu|, \phi_\mu$$

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- The second stop must be **heavy**,
and mostly “**left-handed**”

- Increase the **Higgs mass**
- Reduce the SUSY contribution to Δp

$$m_{\tilde{t}_1} \approx m_t, m_{\tilde{t}_2} = 10 \text{ TeV}$$

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CP-violating processes (e.g. 1-loop EDM's)

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- The generated **BAU** also depends on the
heavy MSSM **Higgs** sector through $\Delta \beta$

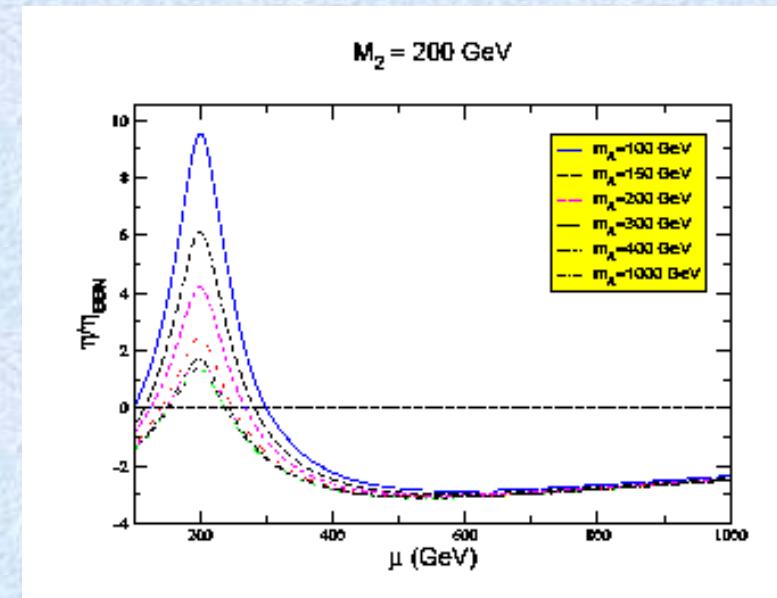
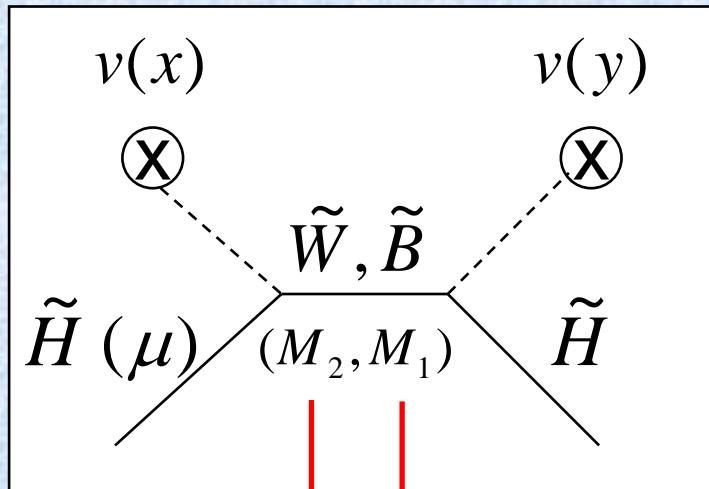
$$m_A = 150 \text{ GeV}, 1000 \text{ GeV}$$

Resonant EW Baryogenesis

...even if all these conditions hold, **CP**-violating sources are large enough if close-to-resonant

conditions are met in the **gaugino-higgsino** sector

$$M_{1,2} \approx \mu$$



Resonant Chargino Source: Well known fact^(*)

Resonant Neutralino Source: **Novelty!** ^(**)

Connection with **Dark Matter** !!

^(*) M.Carena et al., (2003); Lee et al., (2005) ^(**) V.Cirigliano, S.Profumo, M.Ramsey-Musolf (2006)

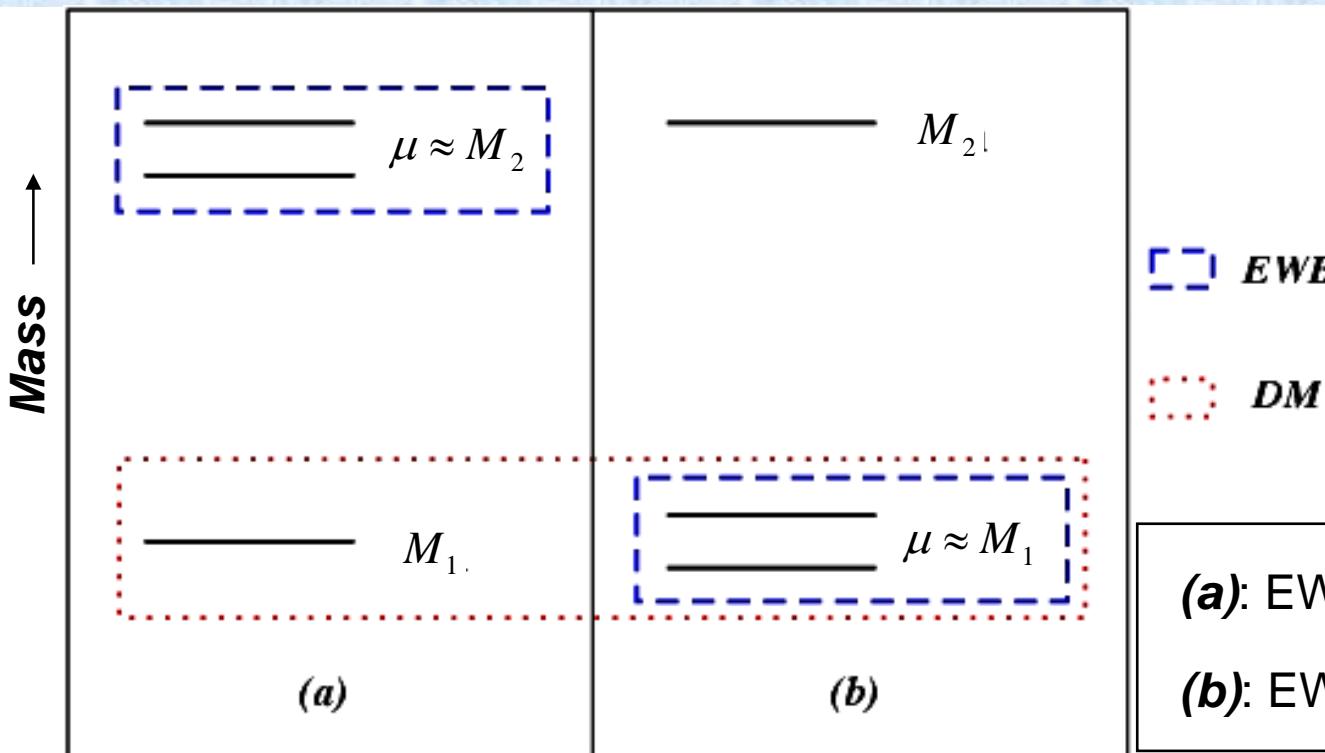
EW Baryogenesis and DM

In the (*R*-parity conserving) MSSM **the LSP is stable**



The LSP must be a phenomenologically viable relic

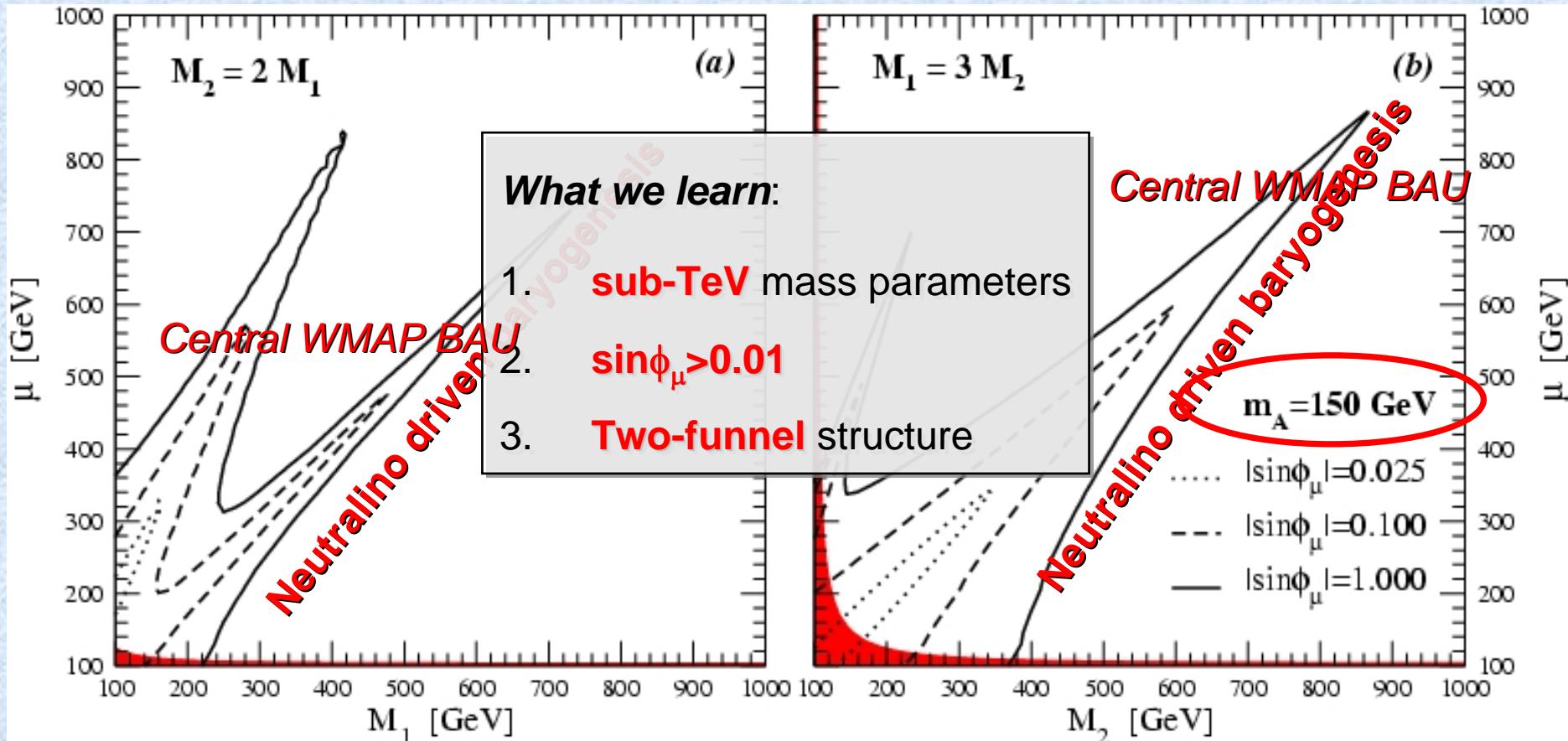
(electrically and color neutral, low enough relic abundance and direct detection rates)



(a): EWB and DM are unrelated

(b): EWB-DM connection

Baryon Asymmetry in the MSSM



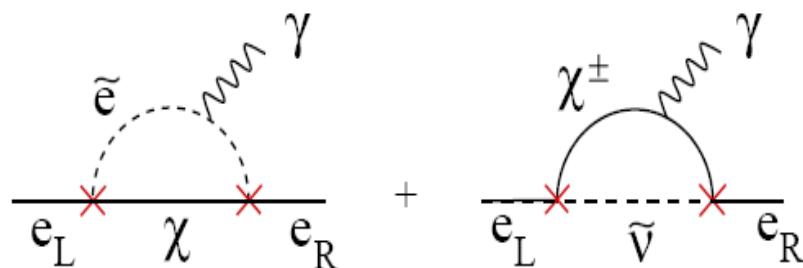
- “**Supergravity**”-like gaugino mass pattern
- “**Anomaly mediation**”-like gaugino mass pattern

$$M_1 = \frac{5}{3} \frac{\sin^2 \vartheta_W}{1 - \sin^2 \vartheta_W} \approx M_2 / 2$$

$$M_1 = \frac{\beta_{g_1}}{\beta_{g_2}} \frac{g_2}{g_1} M_2 \approx 3 M_2$$

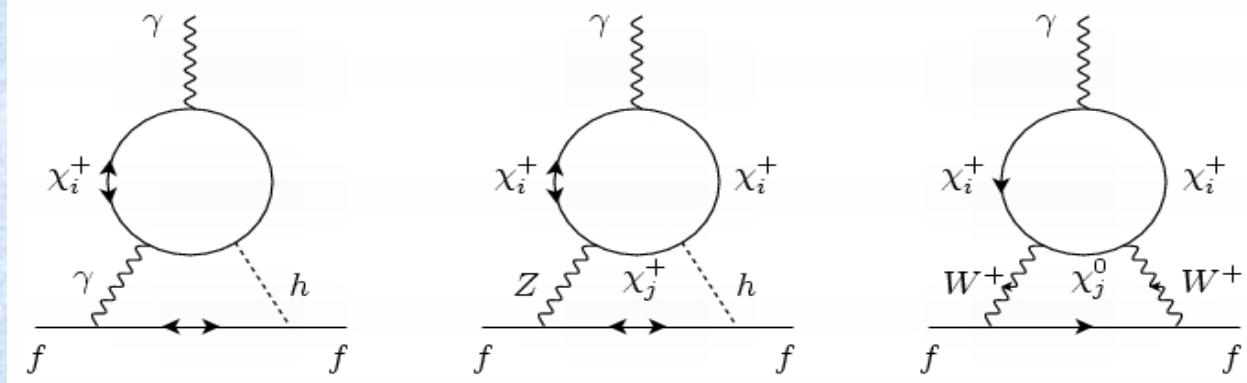
Electric Dipole Moments

CP-violating interactions in the SUSY sector induce EDMs
In the present setup, the best probe is the **electron** EDM



1-loop (electron) EDM

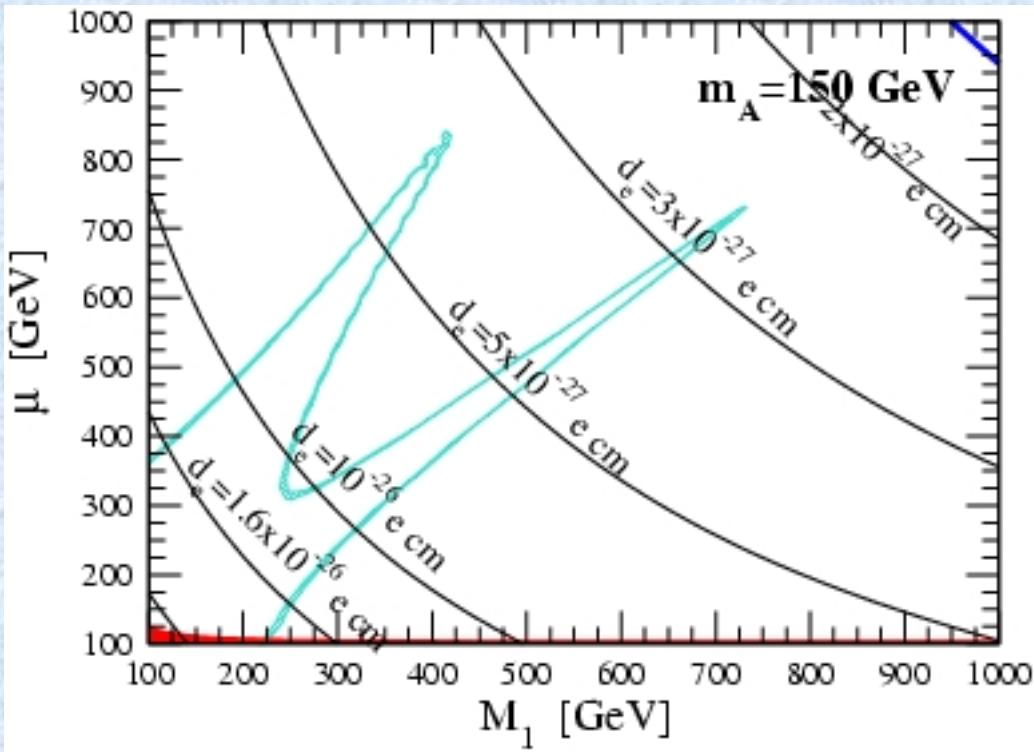
Asymptotically vanish in the limit of large sfermion masses



2-loop EDM

*Only contribution
In, e.g., SplitSUSY*

EDMs and EW Baryogenesis

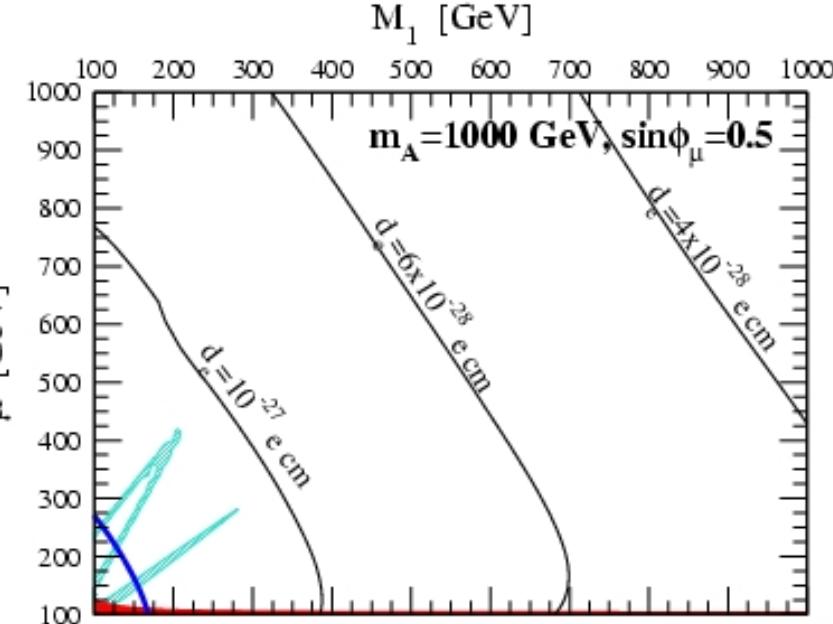


$$\sin \phi_\mu \approx 1$$

- Only **two-loop** EDMs (*heavy sfermions limit*)
- Anomaly mediated case even **worse!**
- **Maximal phases are not compatible with EW Baryogen.**

$$d_e^{\text{exp, cur}} \approx 1.6 \times 10^{-27} \text{ e} \cdot \text{cm}$$

EDMs and EW Baryogenesis



$$\sin \phi_\mu < 1$$

What we learn:

1. EDM and EWB are **compatible**
2. **maximal phases** are **excluded**
by current data
3. there is a **lower bound** on the el. EDM

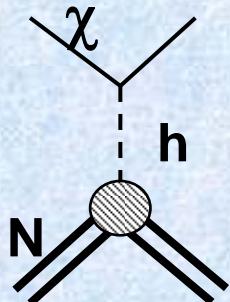
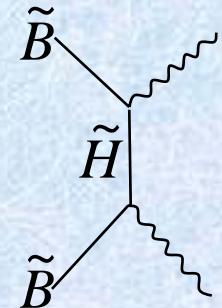
$$d_e \geq 10^{-28} \text{ e} \cdot \text{cm} \gg d_e^{\text{exp, fut}} \approx 10^{-29 \div 30} \text{ e} \cdot \text{cm}$$

**EDM experiments will conclusively
test the EW Baryogenesis scenario!**

EWB and DM: a closer look

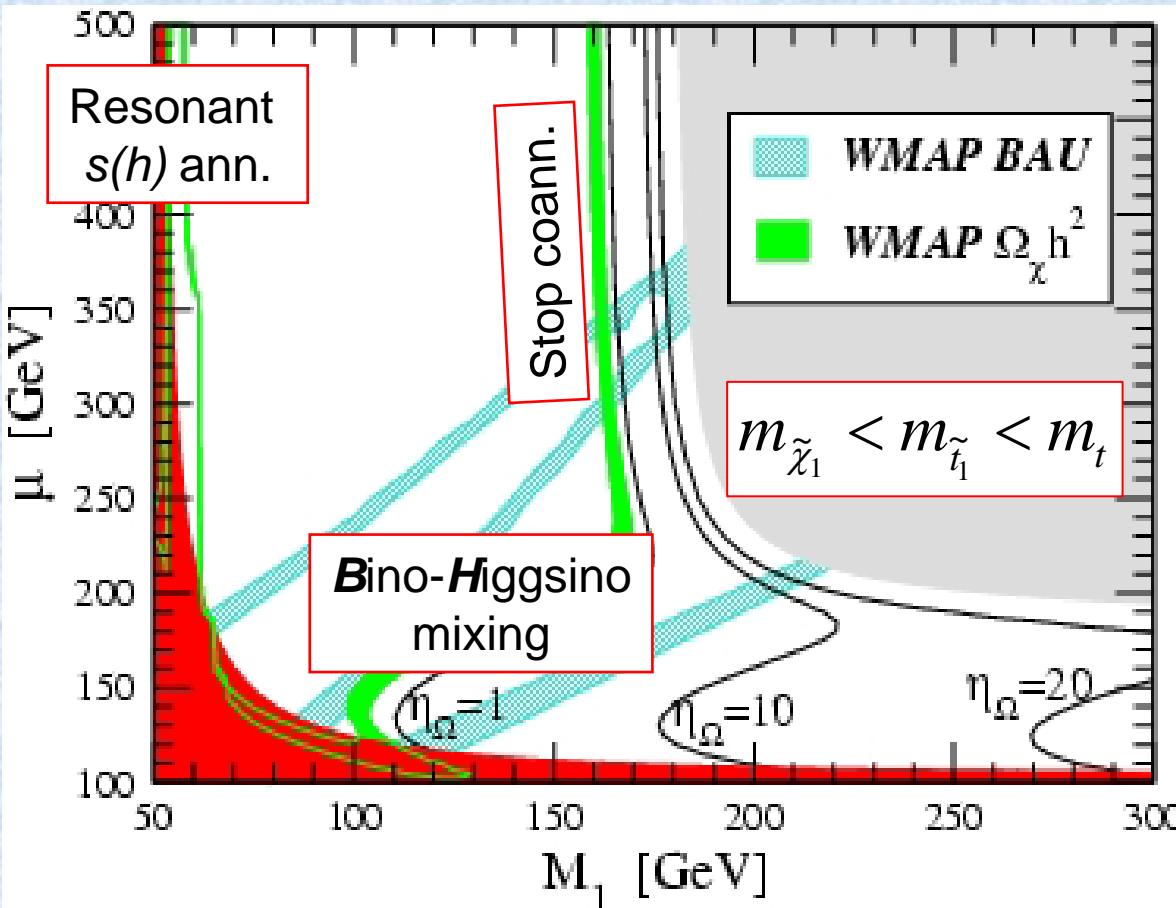
$$M_{1,2} \approx \mu$$

- The **higgsino mixing** is required to have a low enough **relic abundance** in the $M_1 \approx \mu$ case and fulfill EWB + right thermal χ production



- Large higgsino mixing implies large couplings to the Higgses, and hence **large direct detection rates** even with heavy s-quarks
 - Since $m_{\tilde{\chi}_1} < m_{\tilde{t}_1} < m_t$ the DM particle is **light**
- this means that {
- the local DM **number density** is large ($\rho_{\text{DM}} = m_{\text{DM}} \cdot n_{\text{DM}}$)
 - the number of **pair annihilations** is large ($\propto n_{\text{DM}}^2$)

The Neutralino Relic Abundance & EWB



$$m_A = 1 \text{ TeV}, \sin(\phi_\mu) = 0.5$$

- **Excessive** relic abundance regions are **ruled out** (caveat: low- T reheating)
 - **Low** relic abundance regions are viable assuming either **non-thermal** production or cosmological enhancement
- $$\eta_\Omega \equiv (\Omega_{\text{CDM}}^{\text{WMAP}} - \Omega_\chi^{\text{th}})/\Omega_\chi^{\text{th}}$$
- In the anomaly mediated SUSY bkg. case

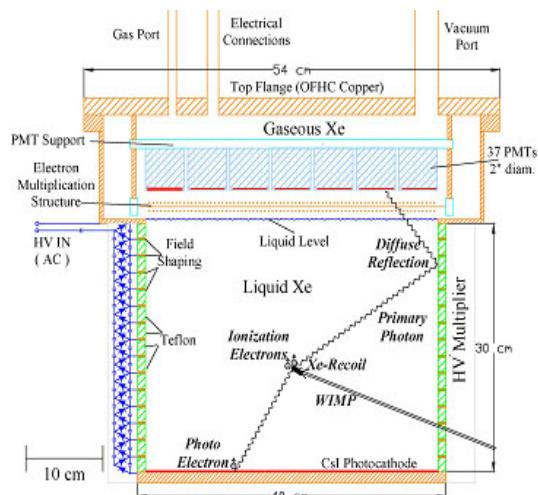
$$30 \leq \eta_\Omega \leq 300$$

Neutralinos can be responsible for both Baryogenesis and Dark Matter

Direct & Indirect Dark Matter Searches

DIRECT DETECTION

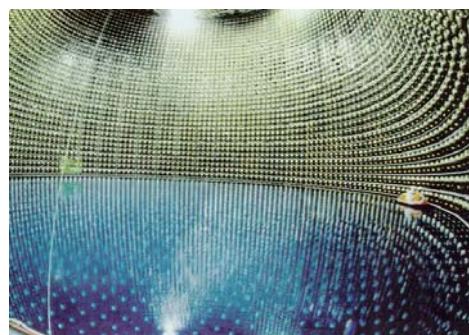
Observation of scattering events of WIMPs off nuclei in low background environments



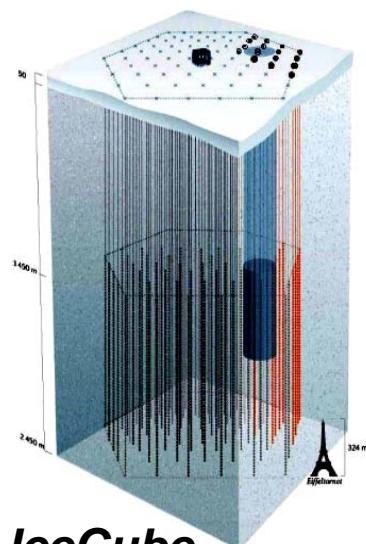
XENON 1-t

HIGH ENERGY NEUTRINOS FROM THE SUN / EARTH

Search for energetic neutrinos produced in $\chi\chi$ pair annihilations in the core of nearby gravitational dips, as the center of the Sun or of the Earth

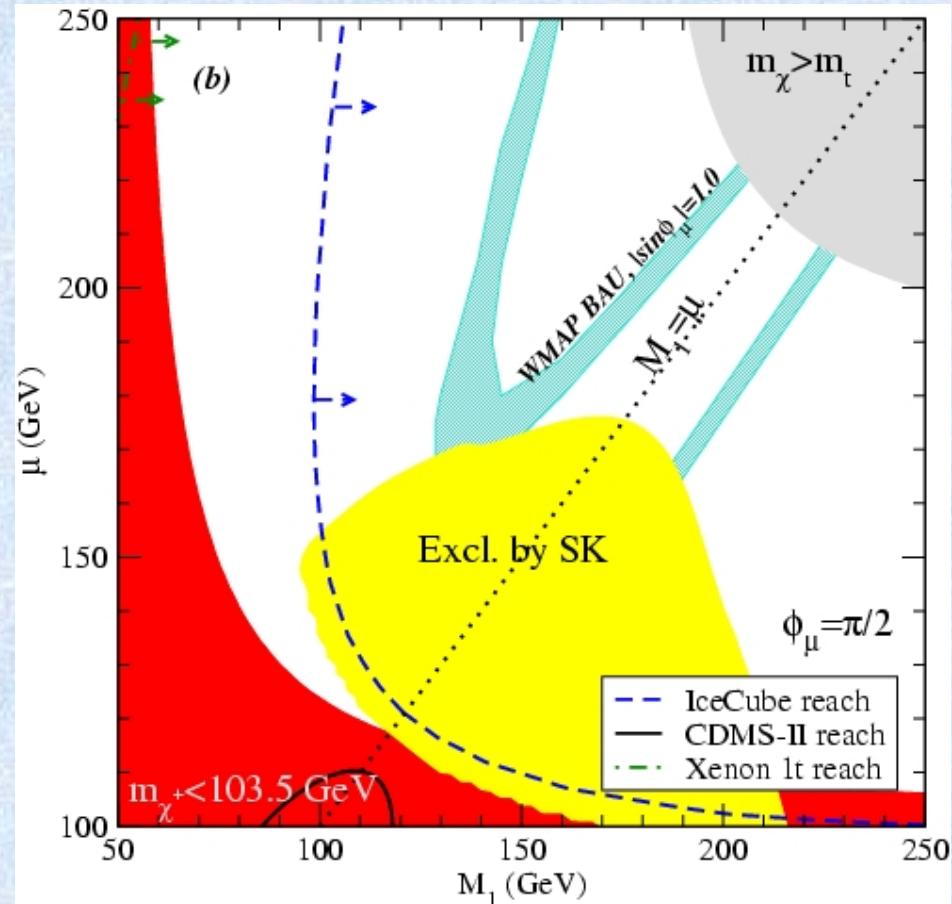
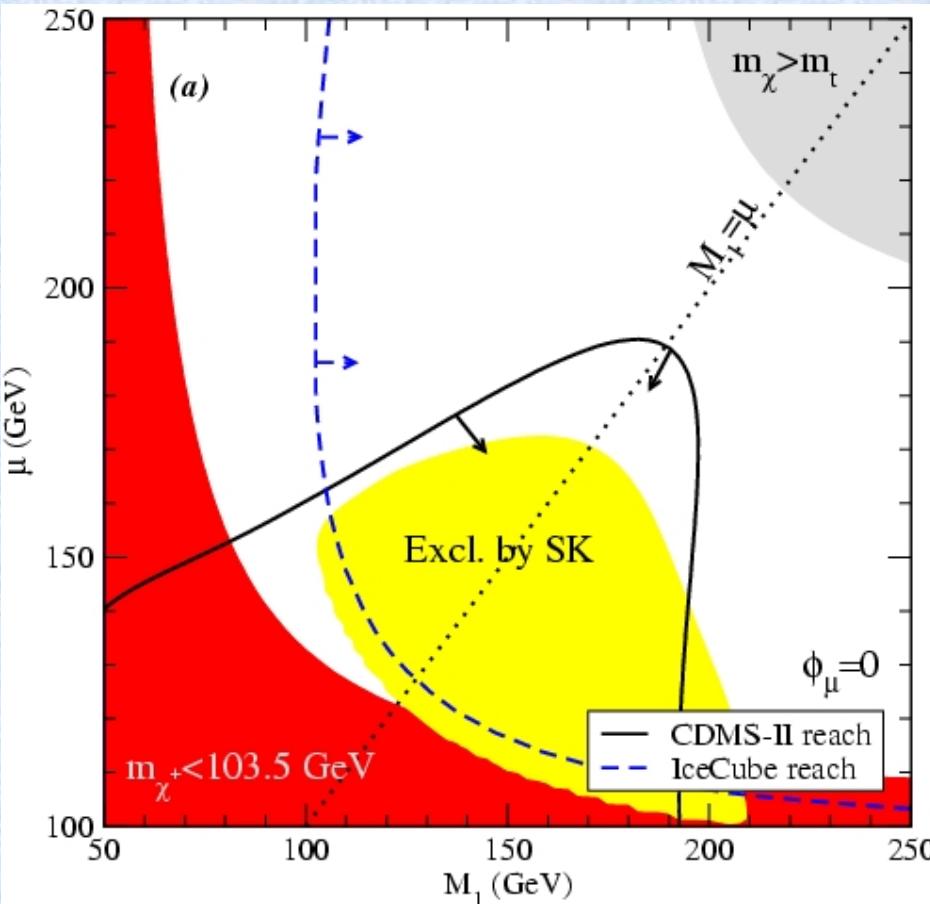


Super-K



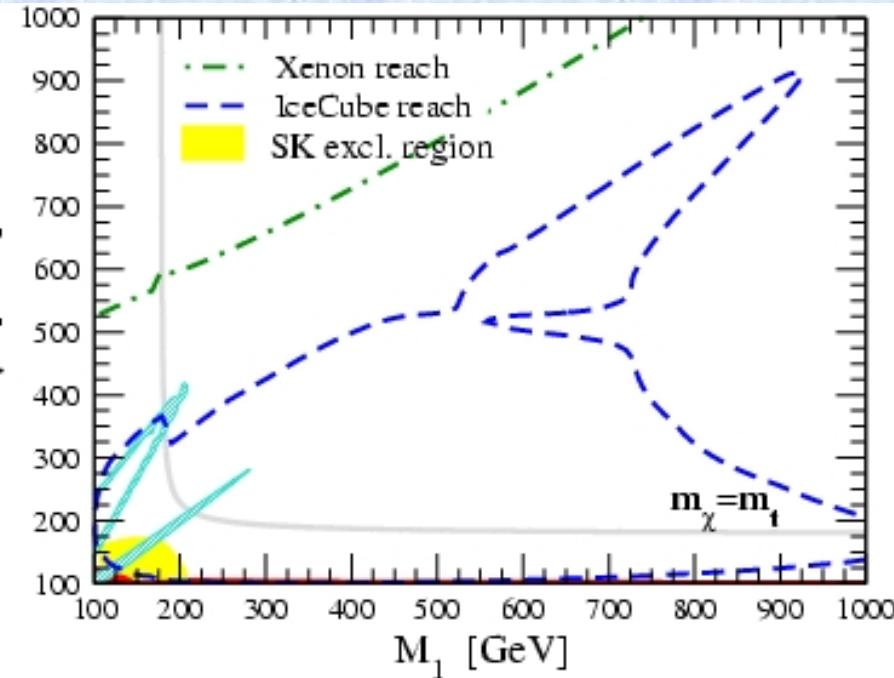
IceCube

Dark Matter searches and EWB



- All **CP conserving case (no EWB)** in reach of neutrino telescopes and/or km⁻ size neutrino telescopes (e.g. IceCube)
- CP violating case (EWB) and supersymmetry
- A sizable portion of the parameter space which we expect to be compatible with EWB is **already ruled out by SuperK** data on the neutrino flux from the Sun
- CP phases **suppress direct detection** and **enhance the neutrino flux f/Sun**

Dark Matter searches and EWB: Zooming Out

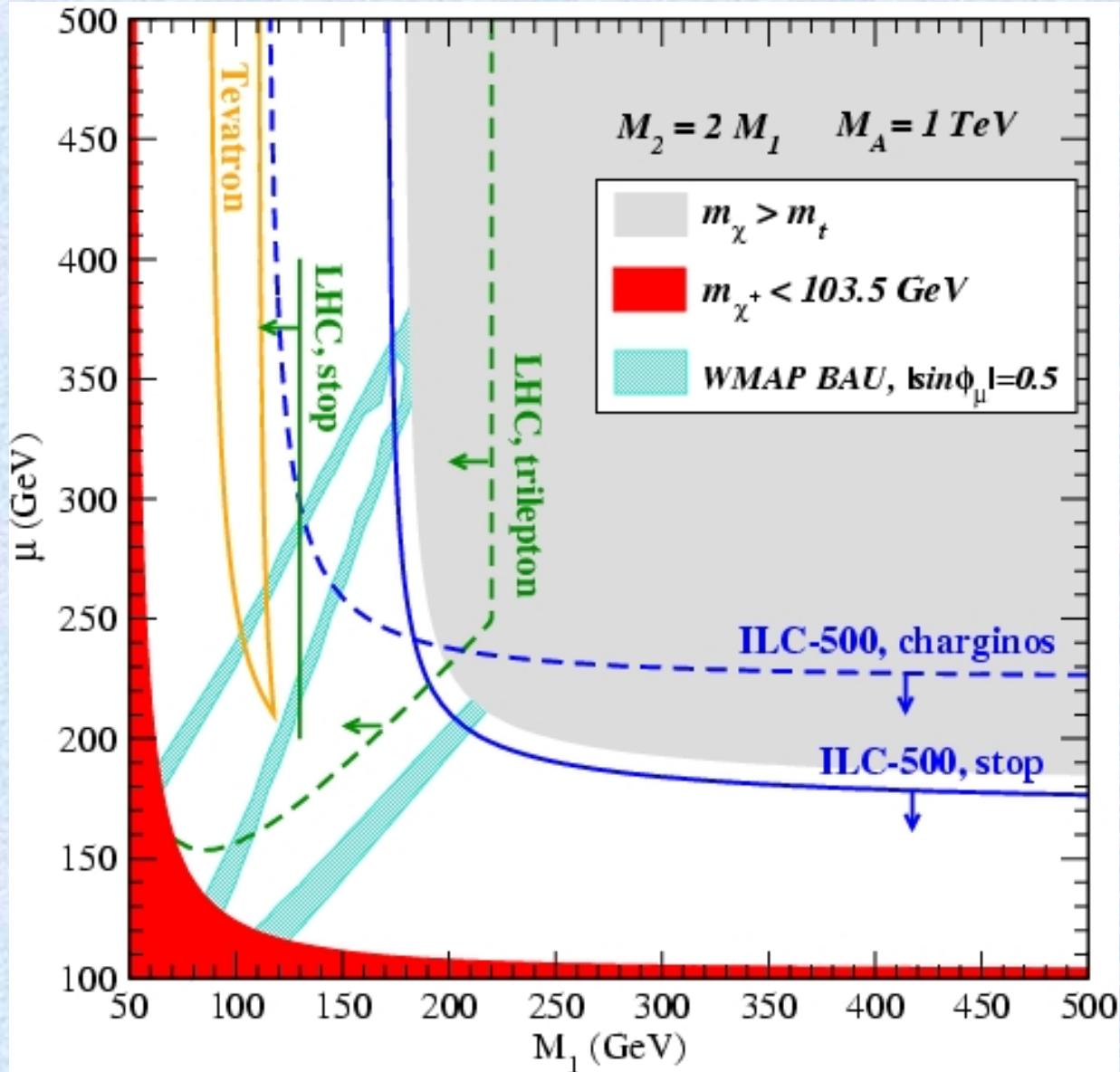


**Supergravity mediated
SUSY breaking**

**Anomaly mediated
SUSY breaking**

**Both ton-sized direct detectors AND neutrino telescopes
will conclusively probe EW Baryogenesis!**

Collider Searches



Tevatron (*)

$$m_{\tilde{t}_1} < m_{\tilde{\chi}_1^+} \quad \tilde{t}_1 \rightarrow c \tilde{\chi}_1^0$$

LHC – same sign top (**)

$$\begin{array}{c} \tilde{g}\tilde{g} \rightarrow \tilde{t}_1^* t \tilde{t}_1^* t \\ \downarrow \qquad \downarrow \\ c\chi \qquad bW \end{array}$$

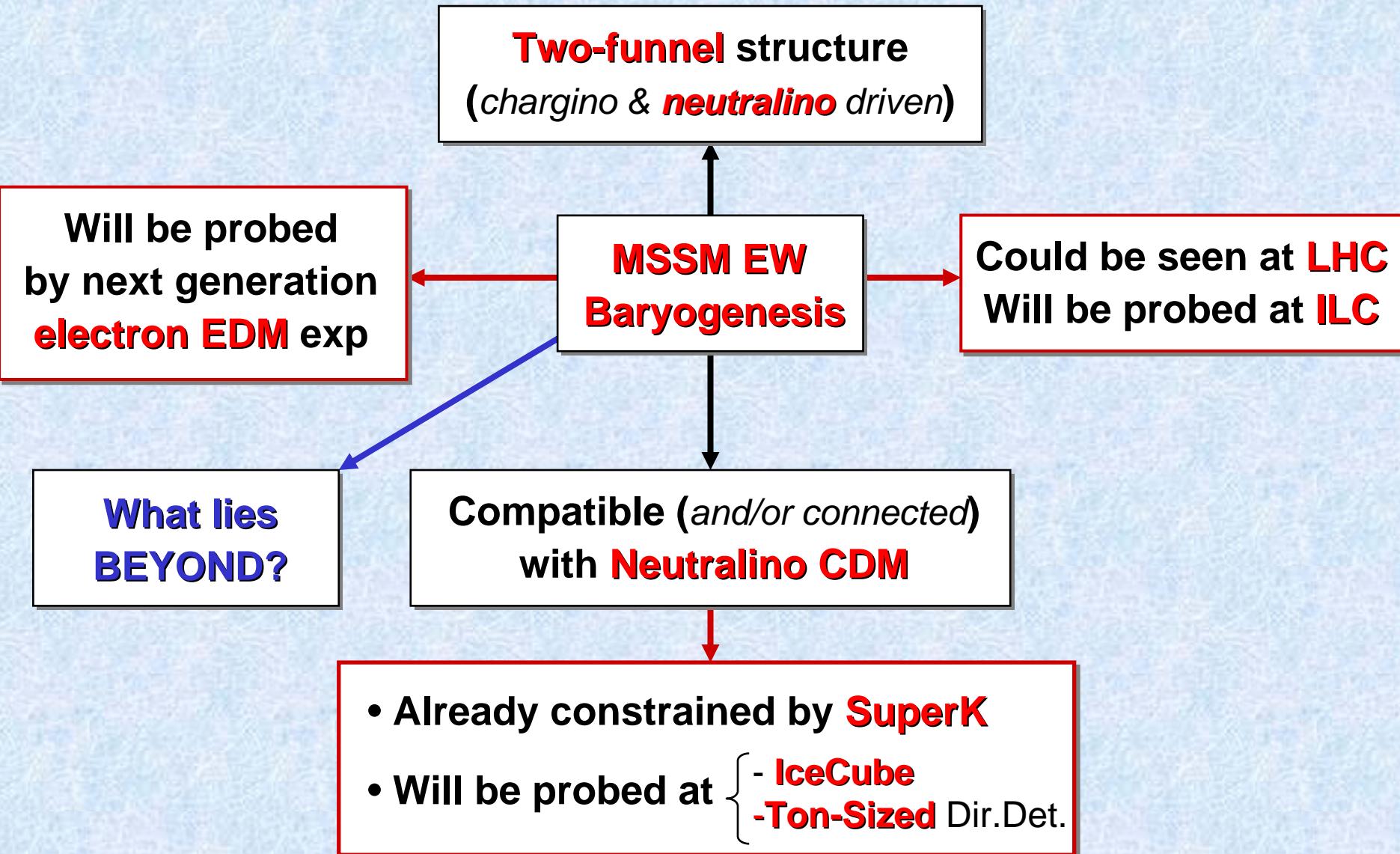
LHC – trilepton (***)

$$\begin{array}{c} \tilde{\chi}_1^\pm \tilde{\chi}_2^0 \\ \tilde{\chi}_1^\pm \rightarrow l\bar{\nu}\chi \text{ and } \tilde{\chi}_2^0 \rightarrow l\bar{l}\chi \end{array}$$

ILC (****)

$$\begin{array}{c} e^+ e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \\ e^+ e^- \rightarrow \tilde{t}_1^* \tilde{t}_1 \end{array}$$

EW Baryogenesis in the MSSM: Summary



Beyond the Minimal SUSY SM

Adding a **Gauge Singlet** Superfield **S** to the Superpotential strongly affects SUSY EWB, adding tree-level cubic terms ^(*)

$$W = \mu H_d H_u + h_s H_d H_u S + \frac{1}{3} k S^3 + \alpha S + (g_u Q u^c H_u + g_d Q d^c H_d g_u + g_e L e^c H_d)$$

(as general as possible, including the **μ -term**, **linear** and **cubic** terms in **S**) ^(**)

The corresponding tree level scalar field potential reads:

$$\begin{aligned} V_F &= |h_s H_d \cdot H_u + \alpha + \kappa S^2|^2 + (|H_d|^2 + |H_u|^2) |\mu^* + h_s S|^2 \\ V_D &= \frac{g_1^2 + g_2^2}{8} (|H_d|^2 - |H_u|^2)^2 + \frac{g_2^2}{2} (H_d^\dagger H_u) (H_d H_u^\dagger) \\ V_{\text{soft}} &= m_d^2 |H_d|^2 + m_u^2 |H_u|^2 + m_s^2 |S|^2 + \\ &\quad \text{---} m_4 (H_d \cdot H_u S + \text{h.c.}) + (b H_d \cdot H_u + \text{h.c.}) + \\ &\quad + m_1^3 (S + \text{h.c.}) + m_2^2 (S^2 + \text{h.c.}) + \text{---} m_3 (S^3 + \text{h.c.}) \end{aligned}$$

^(*) M.Pietroni, Nucl.Phys. **B402** (1993) 27; ^(**) Davies et al., Phys.Lett. **B372** (1996) 88

Beyond the Minimal SUSY SM

1. The **EWPT** is more “naturally” strongly **first order**
(e.g. *if the singlet Higgs is light*)
2. The bound on the **Higgs mass** is alleviated
(*both for EWB and theoretically*)
3. No need for **light stops**
4. Extra possible non-trivial **CP**-structure

Scopes of the projects:

- | | |
|--|---|
| ✓ Study the dynamics of the EWPT
<i>(bubble walls, diffusion, wash-out...)</i> | ✓ Study the CP-structure |
| ✓ Assess the new contribution to
EW precision observables | ✓ Evaluate the new contributions
to Electric Dipole Moments |
| | ✓ DM physics (<i>light singlino...</i>) ^(*) |

^(*) F.Ferrer, L.Krauss and S.Profumo, PRD **74** (2006) 115007

The EW Phase Transition in Singlet Models

A **Toy Model** Warm-up: Minimal **Singlet Extension** of the SM Higgs Sector^(*)

$$\begin{aligned} V(H, S) = & \frac{m^2}{2} H^+ H + \frac{\lambda}{4} (H^+ H)^2 + \\ & a_1 S (H^+ H)/2 + a_2 S^2 (H^+ H)/2 + \\ & b_2 S^2 / 2 + b_3 S^3 / 3 + b_4 S^4 / 4 \end{aligned}$$

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- **Singlet v.e.v.** before the EW Phase Transition

$$R = \frac{b_2 b_4}{b_3^2} (T_c)$$

$$\left\{ \begin{array}{l} \langle \text{Singlet} \rangle_{T > T_c} \neq 0 \Leftrightarrow R < 2/9 \\ \langle \text{Singlet} \rangle_{T > T_c} = 0 \Leftrightarrow R \geq 2/9 \end{array} \right.$$

- Obtain Strongly **First Order** EWPT **without tree-level cubic terms** ($a_1 = b_3 = 0$)
- Connect the “**order parameter**” ϕ_c / T_c to low-energy, collider **observables**

^(*) D O'Connell, M.Ramsey-Musolf, M.Wise, hep-ph/0611014; S.Profumo and M.Ramsey-Musolf

Cosmological probes: Gravitational Waves

- When two or more bubbles **collide**, spherical symmetry is broken; A fraction of their kinetic energy is released in **Gravitational Waves**
- **Turbulent motions** provide another source of GW's
- The **dynamics** of the **EWPT** enter through

$$\alpha = \Delta V / T_*^4$$

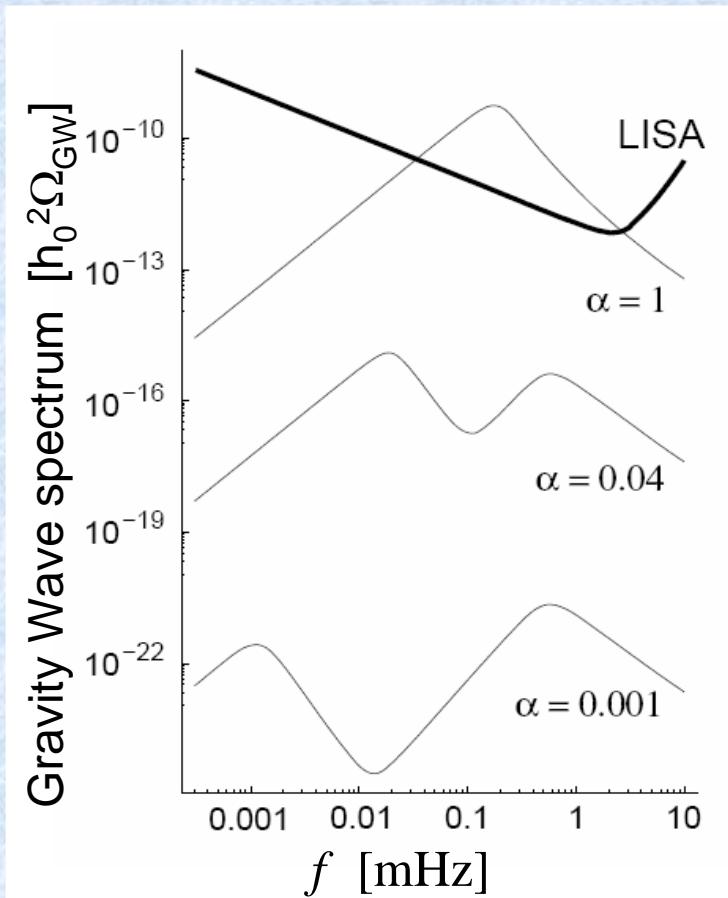
"Latent Heat"
(False vacuum energy)
over Transition T

$$\beta = \dot{\Gamma} / \Gamma \Big|_{T=T_*}$$

Bubble Nucleation
Time Scale
 $\beta \gg H_*$

$$f_{coll} = 5.2 \times 10^{-3} \text{ mHz} \left[\frac{\beta}{H_*} \right] \left[\frac{T_*}{100 \text{ GeV}} \right] \left[\frac{g_*}{100} \right]$$

$$h_0^2 \Omega_{GW}(f_{coll}) \propto \left[\frac{H_*}{\beta} \right]^2 \left[\frac{\alpha}{1+\alpha} \right]^2 \left[\frac{g_*}{100} \right]^{-1/3}$$



^(*) Kamionkowski, Kosowsky, Turner (1994); Nicolis (2003)

Cosmological probes: Heavy Relic Abundances

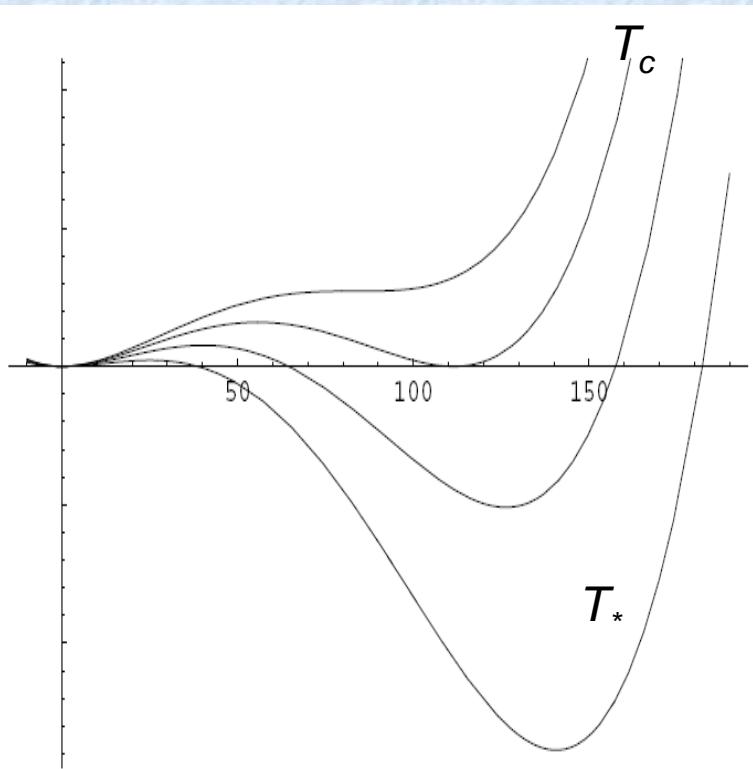
“**Super-Cooling**” dilutes the abundance of **Heavy Relics**

- If $T_{\text{f.o.}} \approx m_\chi / 20 \geq T_c$ the **EWPT** affects the χ **relic density**

Cosmological probes: Heavy Relic Abundances

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- If $T_{\text{f.o.}} \approx m_\chi / 20 \geq T_c$ the **EWPT** affects the χ **relic density**



With a **strongly** first order EWPT, the Universe is trapped in **false vacuum** $\phi_c=0$ until quantum tunneling becomes efficient (T_*)

The **vacuum energy** is then released, **re-heating** the Universe to T_c

$$\frac{s_f}{s_i} \approx \left(\frac{T_c}{T_*} \right)^3$$

$$\frac{(n_\chi)_f}{(n_\chi)_i} \approx \left(\frac{T_*}{T_c} \right)^3$$

Superheavy Relic Density Dilution

