

Baryogenesis and EDMs: Constraints on CP Violation Beyond the Standard Model

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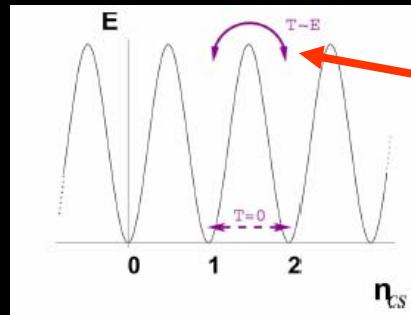
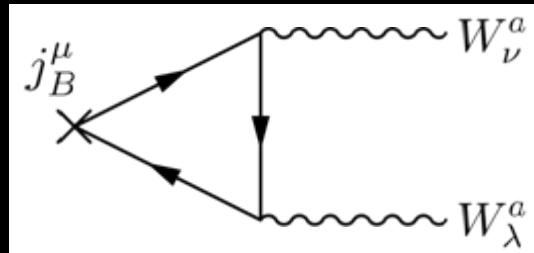
Outline

- Electroweak Baryogenesis
 - *Model scenario: MSSM Baryogenesis*
 - *Transport Equations from QFT*
 - *Resonantly-Enhanced Sources*
 - *Higgs-to-baryon conversion*
- EDM Constraints
 - *MSSM Loops*
 - *Constraints on EWB*

Baryogenesis during Bubble Nucleation

Kuzmin, Rubakov, Shaposhnikov;
Cohen, Kaplan, Nelson;
Joyce, Prokopec, Turok

- “Bubble” nucleation of regions of broken EW phase
- Baryon number violating sphalerons active outside bubbles

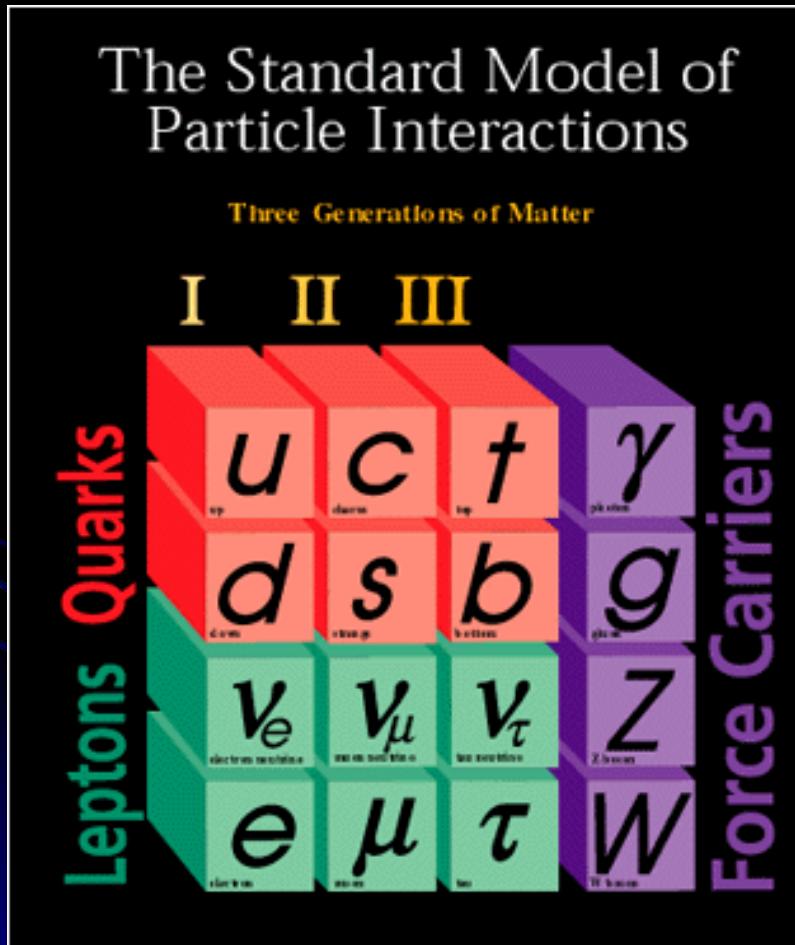


$$\begin{aligned}\Delta B &= \Delta L \\ &= N_F \Delta n_{CS}\end{aligned}$$

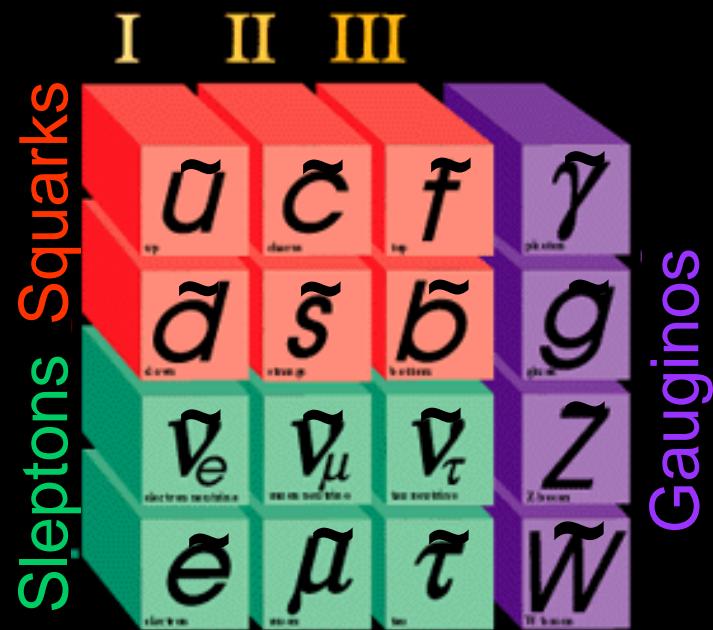
- CP-violating interactions active inside bubble walls
- n_L created in wall, diffuses outside, induces baryon number violation

$$\partial_t \rho_B(x) - D_q \nabla^2 \rho_B(x) \propto -n_F \Gamma_{\text{sph}} [n_L(x) + R \rho_B(x)]$$

Supersymmetry



The Minimal Supersymmetric Extension of the Standard Model (MSSM)



Viability of MSSM Baryogenesis

- First-order phase transition for $m_H \lesssim 120$ GeV
 - Hints of higher limit?
- LEP limit $m_H \gtrsim 114$ GeV
- First-order PT also requires one light top squark
 - Precision electroweak requires heavy \tilde{t}_L
 - So \tilde{t}_R must be light
 - Difficult to generate large CPV squark source with disparate masses
 - Rely on Higgsino source instead...
- Extensions of MSSM with extra light scalar (e.g. NMSSM) can strengthen phase transition and allow degenerate squarks
 - enhancing CP-violating squark source

Densities in MSSM Baryogenesis

- Densities for supermultiplets:

$$Q = n_{t_L} + n_{\tilde{t}_L} + n_{b_L} + n_{\tilde{b}_L}$$

$$T = n_{t_R} + n_{\tilde{t}_R}$$

$$B = n_{b_R} + n_{\tilde{b}_R}$$

$$H = n_{H_u} + n_{\tilde{H}_u} - n_{H_d} - n_{\tilde{H}_d}$$

$$h = n_{H_u} + n_{\tilde{H}_u} + n_{H_d} + n_{\tilde{H}_d}$$

Assume supergauge interactions keep superpartners in chemical equilibrium

- Relate particle densities to chemical potentials:

$$\begin{aligned} n_i &= g_i \int \frac{d^3k}{(2\pi)^3} [n_{B,F}(\omega_{\mathbf{k}}, \mu_i) - n_{B,F}(\omega_{\mathbf{k}}, -\mu_i)] \\ &\Rightarrow n_i = \frac{k_i T^2}{6} \mu_i \end{aligned}$$

Transport Equations

- Coupled transport equations for particle densities:

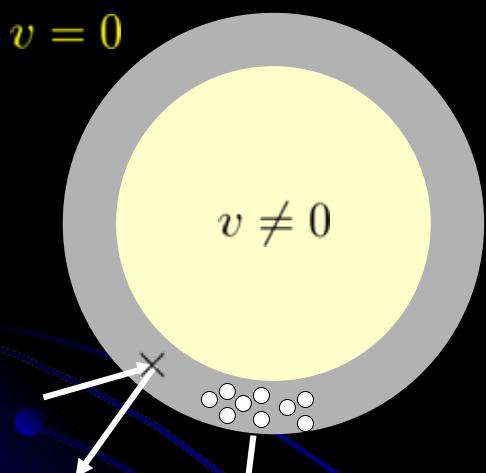
Huet,
Nelson
(1995):

$$\partial_\mu J_i^\mu = S_i^{CP}(\mu_i) + S_i^{C\bar{P}} + S_i^{\text{s.s.}}$$

↓

$v = 0$

$v \neq 0$



$$v_w T' - D_q T'' = S_{\tilde{t}_R}^{CP} - S^{\text{s.s.}} + \Gamma_M^-(\mu_Q - \mu_T) + \Gamma_Y(\mu_Q - \mu_T + \mu_H)$$

$$v_w Q' - D_q Q'' = -S_{\tilde{t}_R}^{C\bar{P}} + 2S^{\text{s.s.}} - \Gamma_M^-(\mu_Q - \mu_T) - \Gamma_Y(\mu_Q - \mu_T + \mu_H)$$

$$v_w H' - D_h H'' = S_{\tilde{H}}^{CP} - (\Gamma_H)\mu_H - (\Gamma_Y)(\mu_Q - \mu_T + \mu_H)$$

Higgs-(s)quark
flavor transfer

$$S_{\tilde{H}}^{CP} \gg S_{\tilde{t}_R}^{CP}$$

$$D_h \gg D_q$$

CP violating sources

Relaxation rates

Higgsino source dominates in MSSM, Γ_Y transfers it to (s)quark sector

Derive from Field Theory

- Derive transport equations for particle densities:

Thermal expectation values of currents:

$$\langle \partial_\mu j^\mu(x) \rangle \quad \text{where} \quad \langle \mathcal{O}(x) \rangle = \frac{1}{Z} \text{Tr} \left[e^{-\beta \mathcal{F}} \mathcal{O}(x) \right]$$

Bosons $j_\phi^\mu = i(\phi^* \partial^\mu \phi - \phi \partial^\mu \phi^*)$

Fermions $j_\psi^\mu = i\bar{\psi} \gamma^\mu \psi$

Closed Time Path QFT

- Expectation value in given "in" state:

$$\langle n | S_{\text{int}}^\dagger T \{ \mathcal{O}(x) S_{\text{int}} \} | n \rangle$$

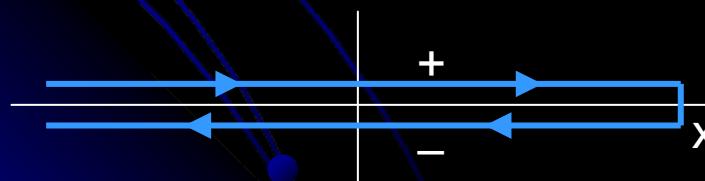
$$\Downarrow \quad S_{\text{int}} = T \exp \left(i \int d^4 x \mathcal{L}_{\text{int}} \right)$$

$$\langle n | \mathcal{P} \left\{ \mathcal{O}_+(x) \exp \left(i \int d^4 x \mathcal{L}_+ - i \int d^4 x \mathcal{L}_- \right) \right\} | n \rangle$$

"path-ordering"

$$\Downarrow \quad \langle n | \mathcal{P} \left\{ \mathcal{O}(x) \exp \left(i \int_C d^4 x \mathcal{L} \right) \right\} | n \rangle$$

Four Green's functions:

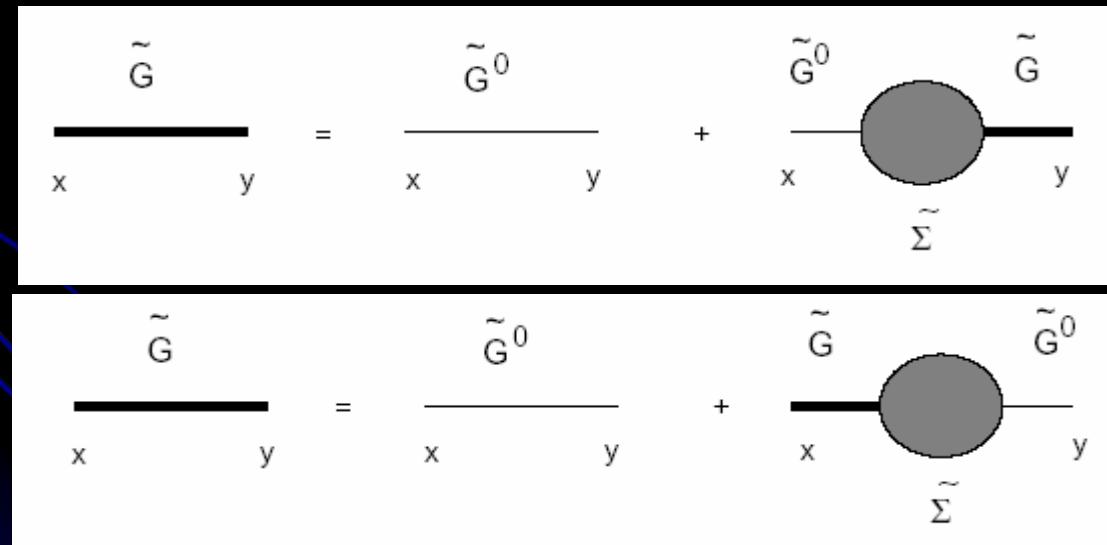


$$\Rightarrow \tilde{G}(x, y) = \begin{pmatrix} G^{++} & -G^{+-} \\ G^{-+} & -G^{--} \end{pmatrix}$$

Schwinger-Dyson Equations

$$\tilde{G}(x, y) = \tilde{G}^0(x, y) + \int d^4z \int d^4w \tilde{G}^0(x, z) \tilde{\Sigma}(z, w) \tilde{G}(w, y)$$

$$\tilde{G}(x, y) = \tilde{G}^0(x, y) + \int d^4z \int d^4w \tilde{G}(x, z) \tilde{\Sigma}(z, w) \tilde{G}^0(w, y)$$



Quantum Transport Equations

- Schwinger-Dyson equations imply transport equations:

$$\frac{\partial n}{\partial X^0} + \nabla \cdot \mathbf{j}(X)$$

$$= \int d^3z \int_{-\infty}^{X^0} dz^0 [\Sigma^>(X, z) G^<(z, X) - G^>(X, z) \Sigma^<(z, X)]$$

integral over entire history of system

CP-violating and conserving interactions enter self-energies

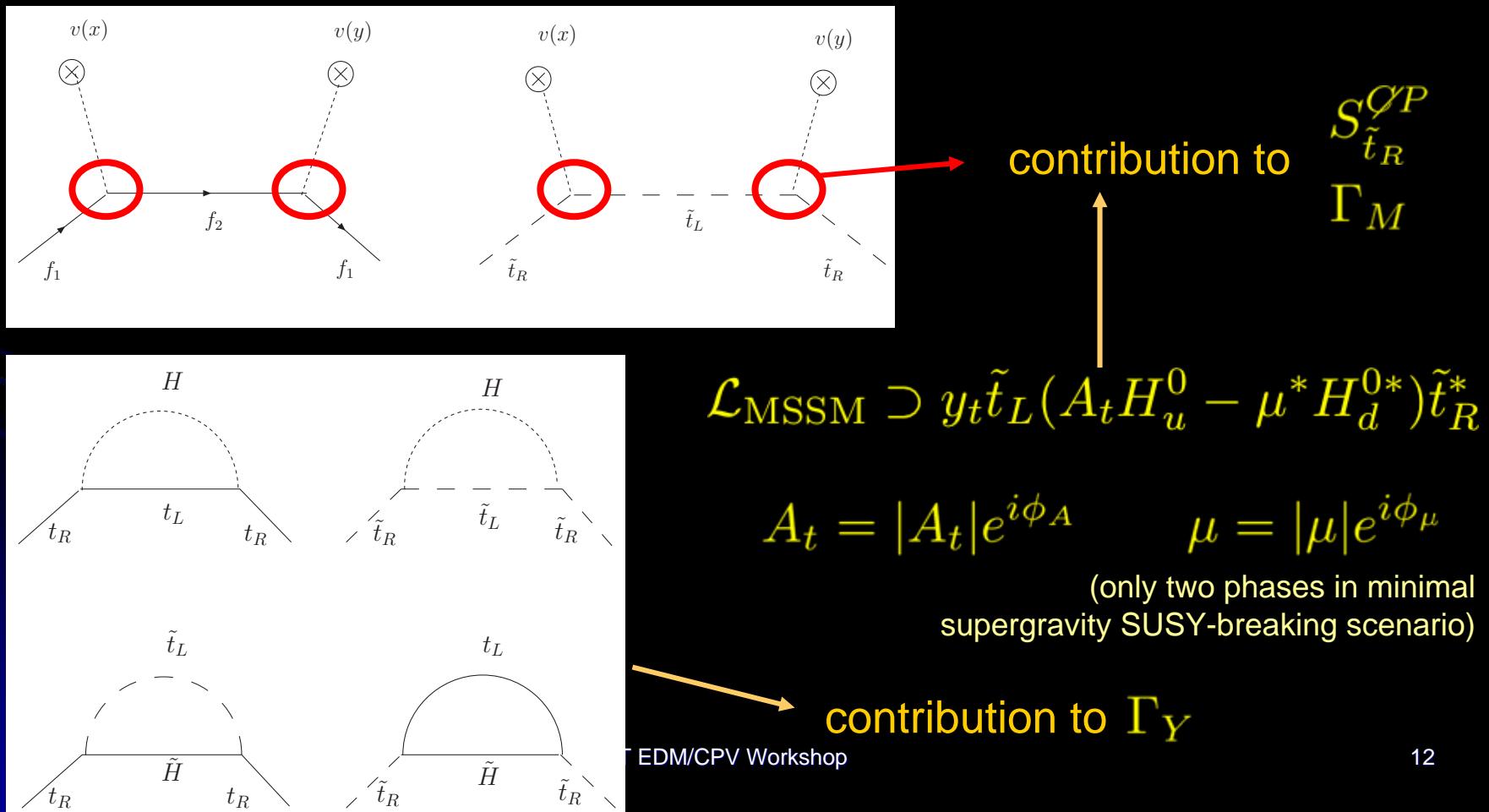
$$- \Sigma^<(X, z) G^>(z, X) + G^<(X, z) \Sigma^>(z, X)]$$

slowly varying in X : $v(X), \mu_i(X)$

$$G_{\tilde{t}_R}^>(x, y) = \int \frac{d^4k}{(2\pi)^4} e^{-ik \cdot (x-y)} [1 + n_B(k^0 - \mu_{\tilde{t}_R})] \rho_{\tilde{t}_R}(k)$$

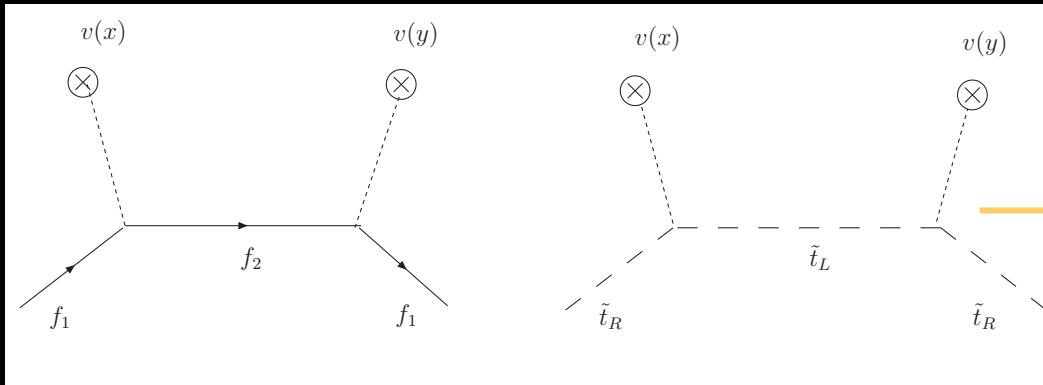
Examples of Self-Energies

- Fermion and scalar interactions with Higgs vevs, Higgs particles, and Higgsinos



Resonances and Enhancements

- Quantum Effects give resonances in sources and relaxation rates, enhanced Yukawa transfer rate

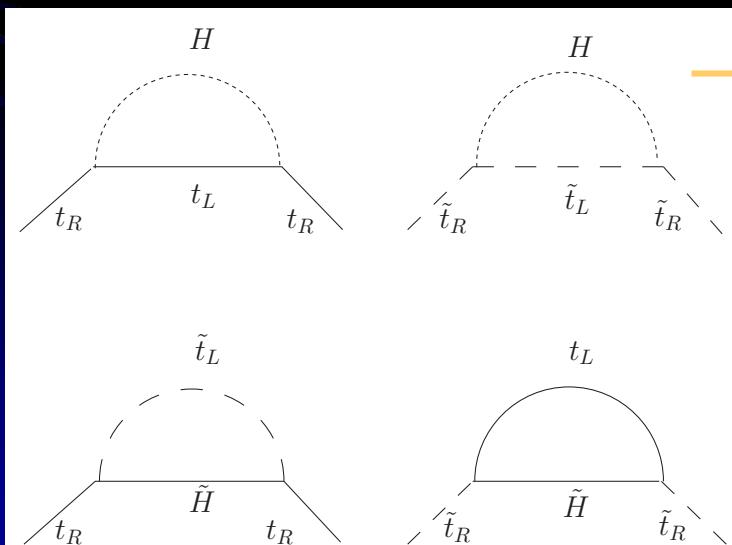


$$S_{\tilde{t}_R}^{QGP} \sim \frac{1}{(m_{\tilde{t}_L} - m_{\tilde{t}_R})^2}$$

Riotto (1998)

$$\Gamma_M^- \sim \frac{1}{m_{\tilde{t}_L} - m_{\tilde{t}_R}}$$

Cirigliano, CL, Ramsey-Musolf (2005)



First calculation of Γ_Y at $\mathcal{O}(\alpha_s^0)$

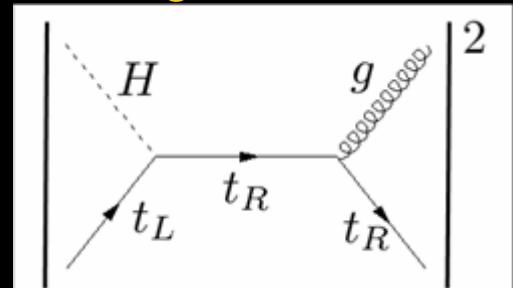
Cirigliano, CL, Ramsey-Musolf, Tulin (2006)

cf. gluon scattering

Joyce, Prokopec, Turok (1994)

$$\Gamma_Y \propto$$

|



|

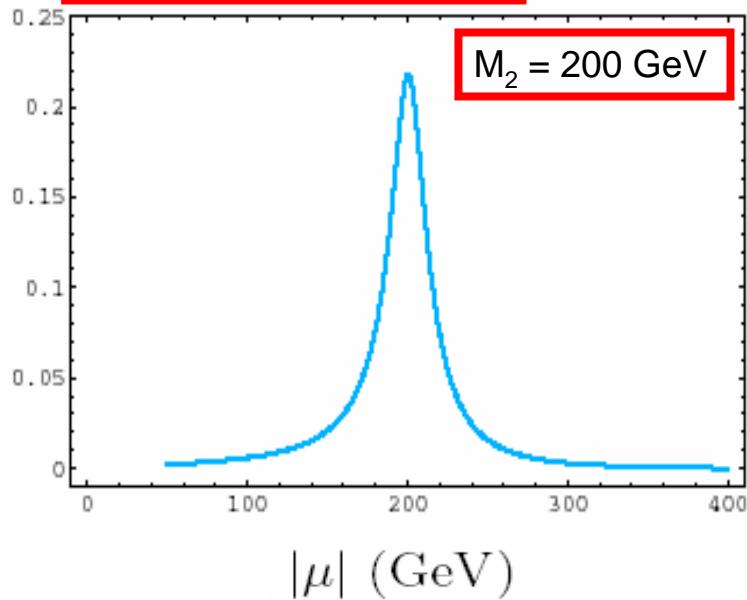
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EDM/CPV Workshop

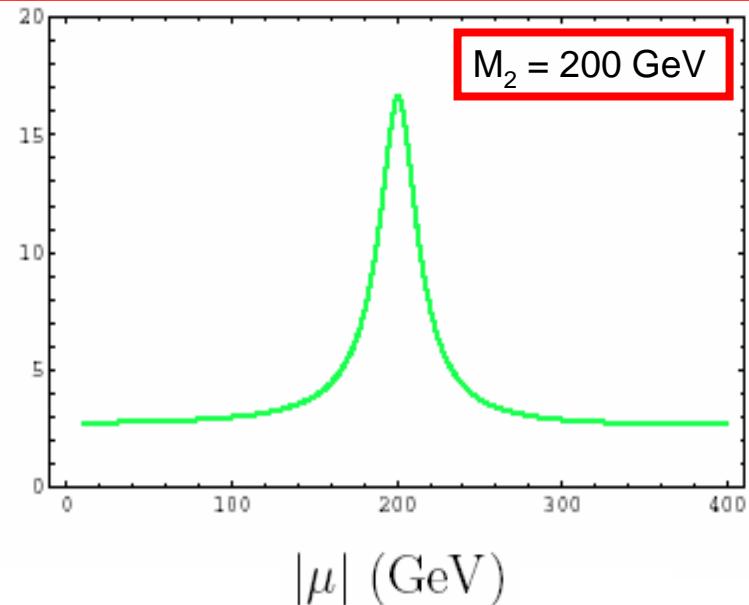
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Higgsino Source and Relaxation

$$-S_{\tilde{H}}^{CP}/(v^2 \dot{\beta} \sin \phi_\mu)$$



$$(\Gamma_h + \Gamma_M^-)/(\Gamma_h + \Gamma_m)_{\text{Huet-Nelson}}$$



$$\rho_B(\bar{z} > 0) \propto \frac{S_{\tilde{H}}^{CP} + S_{\tilde{t}_R}^{CP}}{\sqrt{\Gamma_h + \Gamma_M^-}}$$

21 Mar 2007

INT EDM/CPV Workshop

Resonant enhancements of relaxation terms mitigate but do not cancel out those of CP-violating sources

Transfer of Higgs Source to Baryons

- CP-violating Higgsino source transferred to baryons via Yukawa or tri-scalar interactions

$$v_w H' - D_h H'' = S_{\tilde{H}}^{CP} - \Gamma_H \mu_H - \Gamma_Y (\mu_Q - \mu_T + \mu_H)$$

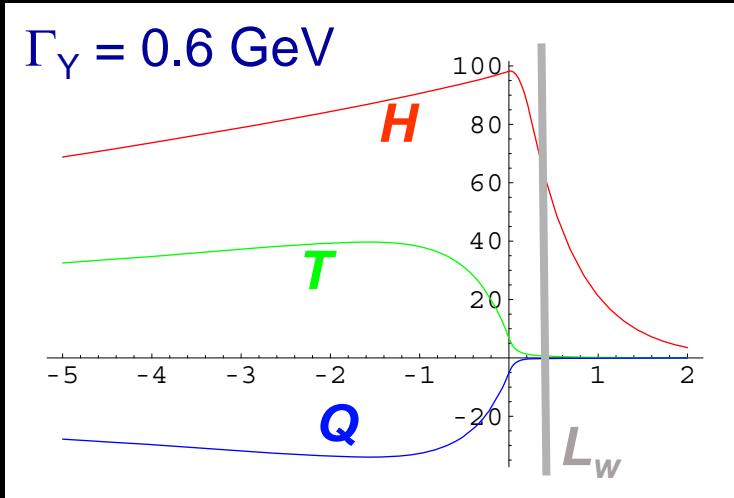
- Efficient diffusion of Higgs number into sphaleron-active region

$$D_h \sim 100/T \gg D_q \sim 6/T$$

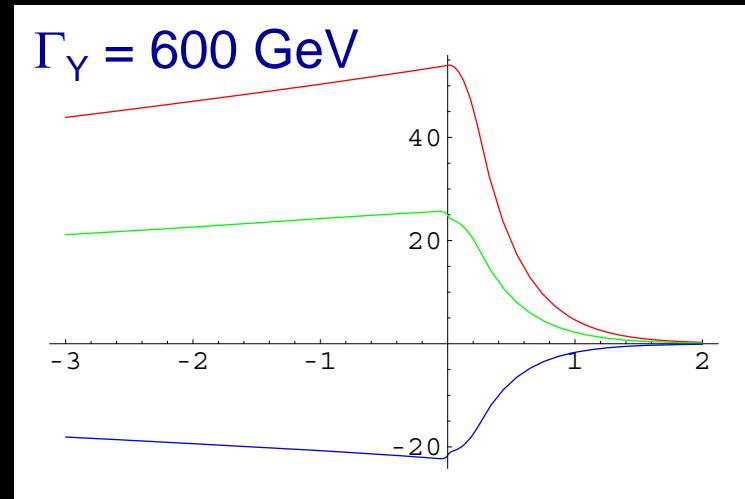
- Conversion to baryons unsuppressed outside bubble wall

$$v_w Q' - D_q Q'' \approx -\Gamma_M^{\cancel{CP}} (\mu_Q - \mu_T) - \Gamma_Y (\mu_Q - \mu_T + \mu_H)$$

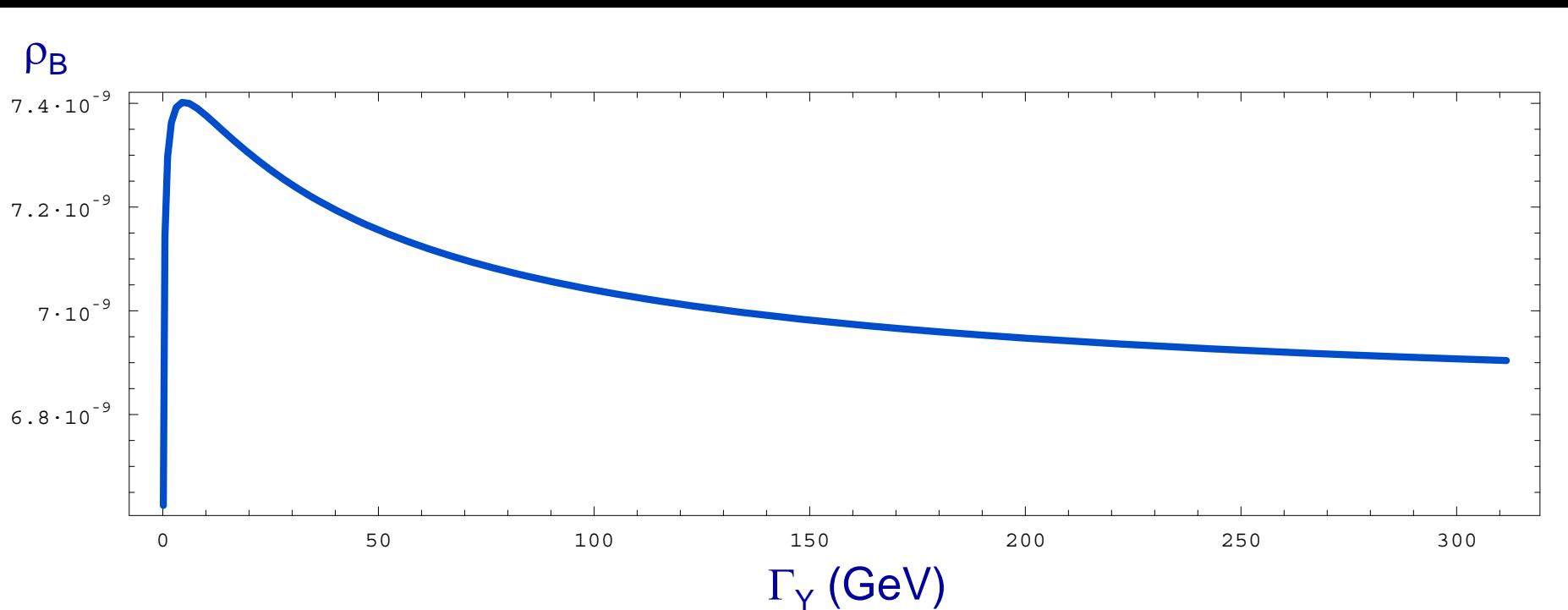
Solution for Densities



$$\begin{aligned}\Gamma_M &= 46 \text{ GeV} \\ \Gamma_H &= 25 \text{ GeV}\end{aligned}$$

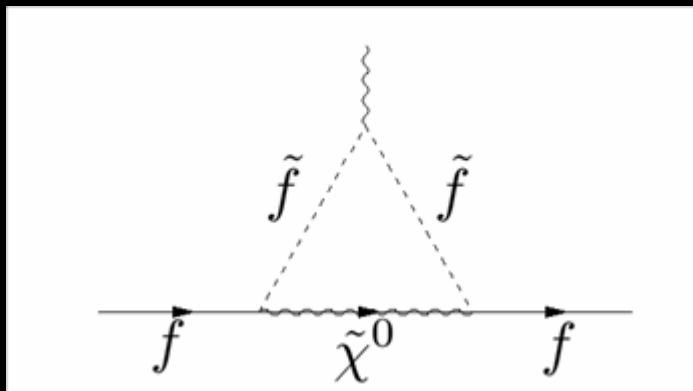


Solution for Baryon Density



Electric Dipole Moments

- One-loop MSSM graphs generating EDMs:

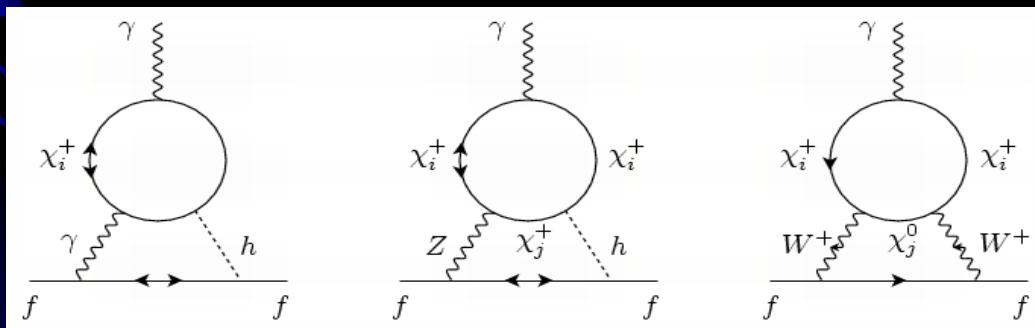


Ibrahim, Nath

Suppressed for large
values of sfermion masses $m_{\tilde{f}}$

- Two-loop graphs independent of sfermion masses

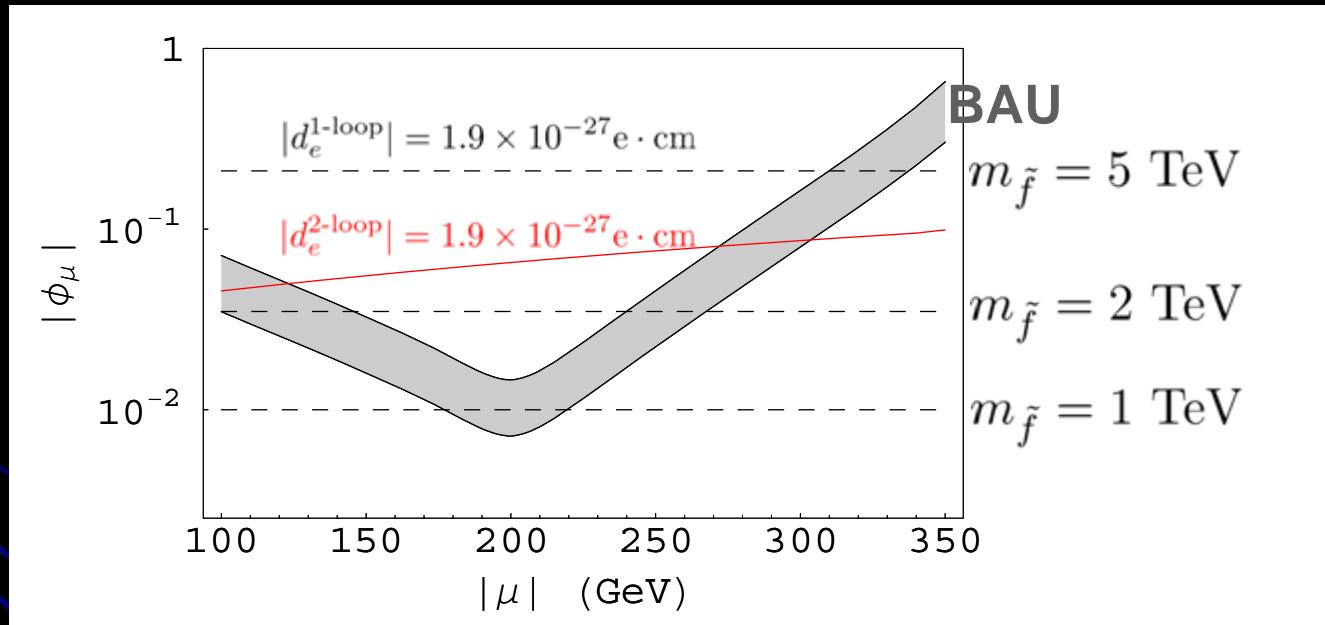
Chang,
Chang,
Keung;
Pilaftsis;
Giudice,
Romanino



Independent of $m_{\tilde{f}}$
Dominant for
 $m_{\tilde{f}} \gtrsim 2 - 3$ TeV

Combined BAU & EDM constraints

$$\mathcal{L}_{\text{SUSY}} \supset |\mu| e^{i\phi_\mu} (\tilde{H}_d^0 \tilde{H}_u^0 - \tilde{H}_d^- \tilde{H}_u^+) - \frac{1}{2} M_2 \bar{\psi}_W \psi_W$$



$M_2 = 200 \text{ GeV}$

Combined BAU & EDM constraints

(1-loop)

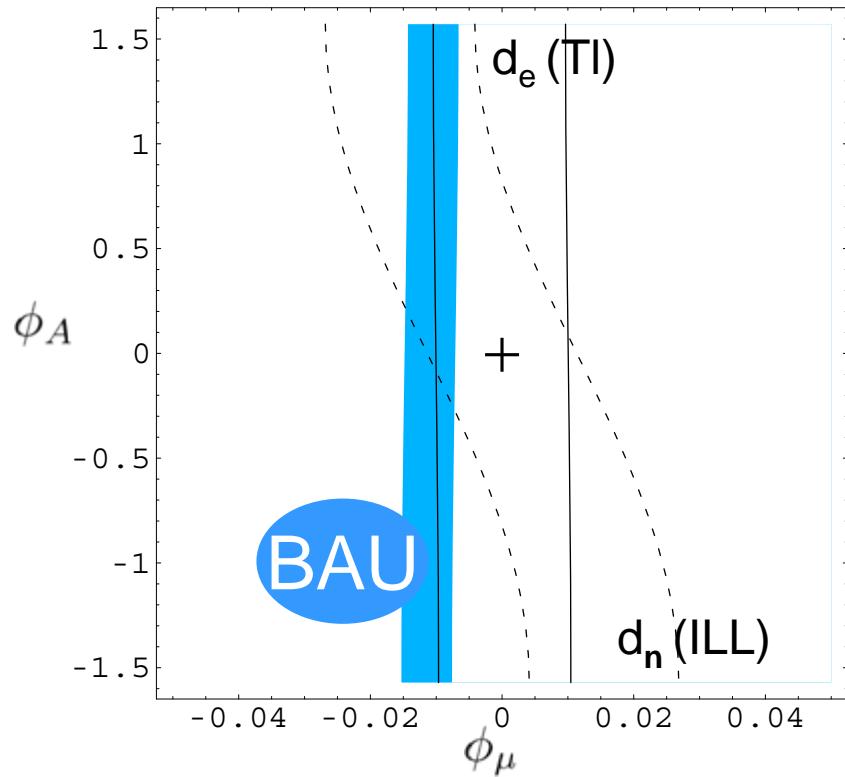
$$\mathcal{L}_{\text{SUSY}} \supset y_t \tilde{t}_L (A_t H_u^0 - \mu^* H_d^{0*}) \tilde{t}_R^* + \mu (\tilde{H}_d^0 \tilde{H}_u^0 - \tilde{H}_d^- \tilde{H}_u^+)$$

$$A_t = |A_t| e^{i\phi_A}$$

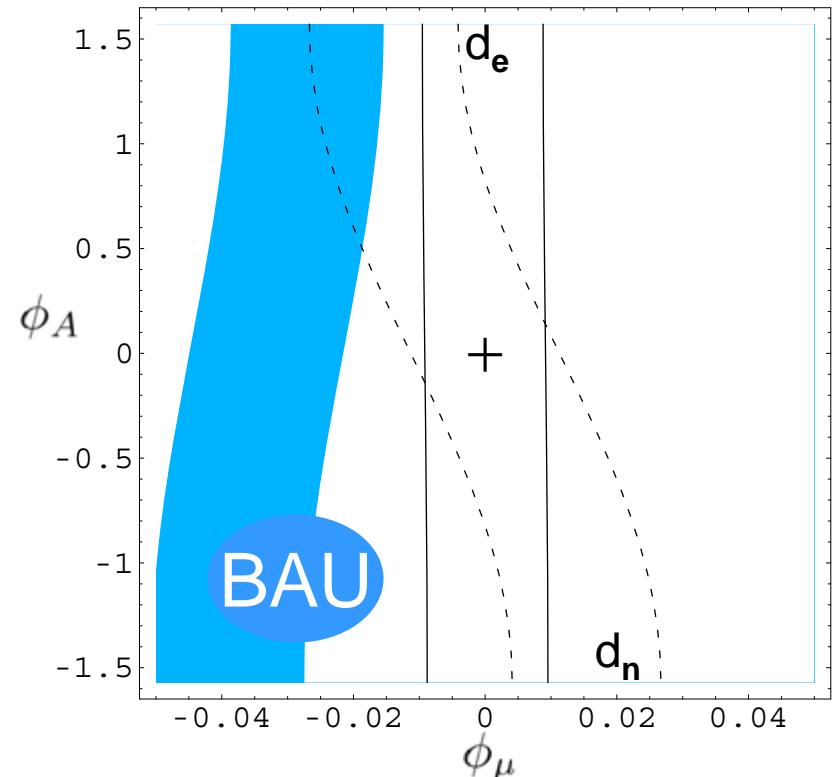
$$\mu = |\mu| e^{i\phi_\mu}$$

neutron edm calc.:
Pospelov, Ritz

$$\mu = M_2 = 200 \text{ GeV}$$



$$M_2 = 200 \text{ GeV}, \mu = 250 \text{ GeV}$$



The Future of Baryogenesis

- More complete, consistent calculation of BAU with non-equilibrium QFT. Still in progress...
 - Resummation of Higgs vev insertions, mixing of mass/flavor eigenstates (cf. Carena et al., Konstandin et al.)
 - Relax approximation of gauge or superpartner equilibrium (keep track of gauge, gaugino densities)
- Non-minimal supersymmetric extensions
 - More CP violating parameters
 - Extra scalar field, stronger phase transition
 - Enhanced squark source possible
- Collider and EDM searches may point to the correct scenario for baryogenesis