

Hadronic weak interaction

David Bowman

ORNL

Institute for Nuclear Theory

June 5, 2007

Essential references

- [DDH]-Desplanques, Donoguhe, and Holstein, *Ann. Phys.* 124:449(1980)
- Adelberger and Haxton, *Ann. Rev. Nucl. Part. Sci.* 1985. 35:501
- Desplanques, *Physics Reports* 297(1988)1

DDH Theory

- Two-body Meson-exchange potential

$$V = \sum_{k=\pi,\rho,\omega} \sum_{\Delta l} h_{k,\Delta l} Y(m_k r) Q_k(\rho, r, \sigma, \tau)$$

- 6 free parameters $f_\pi, h_{\rho,0}, h_{\rho,1}, h_{\rho,2}, h_{\omega,0},$ and $h_{\omega,1}$
- DDH give reasonable ranges. reduce 6 to 4,

$$f_\pi, h_{\rho,0}, h_{\rho,2}, h_{\omega,0}$$

- Nuclear PV is determined by one-body potentials

$$X_N^{p \text{ or } n} = \pm 5.5 f_\pi - 1.13 h_{\rho,0} - 0.91 h_{\omega,0}$$

- The expressions for observables depend on the N-N PC potential used. (AV18 consistently used here.)

DDH ranges and best values

DDH limits and best

$$\begin{pmatrix} f_{\pi} \\ f_{\rho 0} \\ f_{\rho 1} \\ f_{\rho 2} \\ f_{\omega 0} \\ f_{\omega 1} \end{pmatrix} \begin{pmatrix} 0. & 4.6 & 11.4 \\ -30.8 & -11.4 & 11.4 \\ -0.38 & -0.19 & 0. \\ -11. & -9.5 & 7.6 \\ -10.3 & -1.9 & 5.7 \\ -1.9 & -1.1 & -0.8 \end{pmatrix}$$

10 existing precise experiments and 4 constrained quantities

- p-p scattering at 15, 45 MeV
 - linear combination of $h_{\rho,0} + h_{\rho,2}/\sqrt{6}$ and $h_{\omega,0}$
- p-p scattering at 220 MeV
 - $h_{\rho,0} + h_{\rho,2}/\sqrt{6}$
- ^{18}F
 - f_{π}
- p- α , ^{19}F , ^{41}K , ^{175}Lu , ^{181}Ta asymmetries and ^{133}Cs anapole moment
 - X_N^p

p-p consistency check

- p-p at 15 and 45 MeV measure S-P interference and depend on the same linear combination of couplings
 - $f_{\rho,0} + f_{\rho,1} + f_{\rho,2} / \sqrt{6} + .92(f_{\omega,0} + f_{\omega,1})$
- predicted ratio is .56 and measured ratio is $.59 \pm .27$

One-body PV potential

- If the model-space of a nucleus consists of proton (or neutron) excitations then we are interested in matrix elements between these states, $\langle \psi_2 | V_{PNC} | \psi_1 \rangle$.
Although V_{PNC} can change the state of two nucleons, the amplitudes where only 1 nucleon changes state add coherently and dominate the matrix element for large A .

Discussion of ^{18}F (^{19}F)

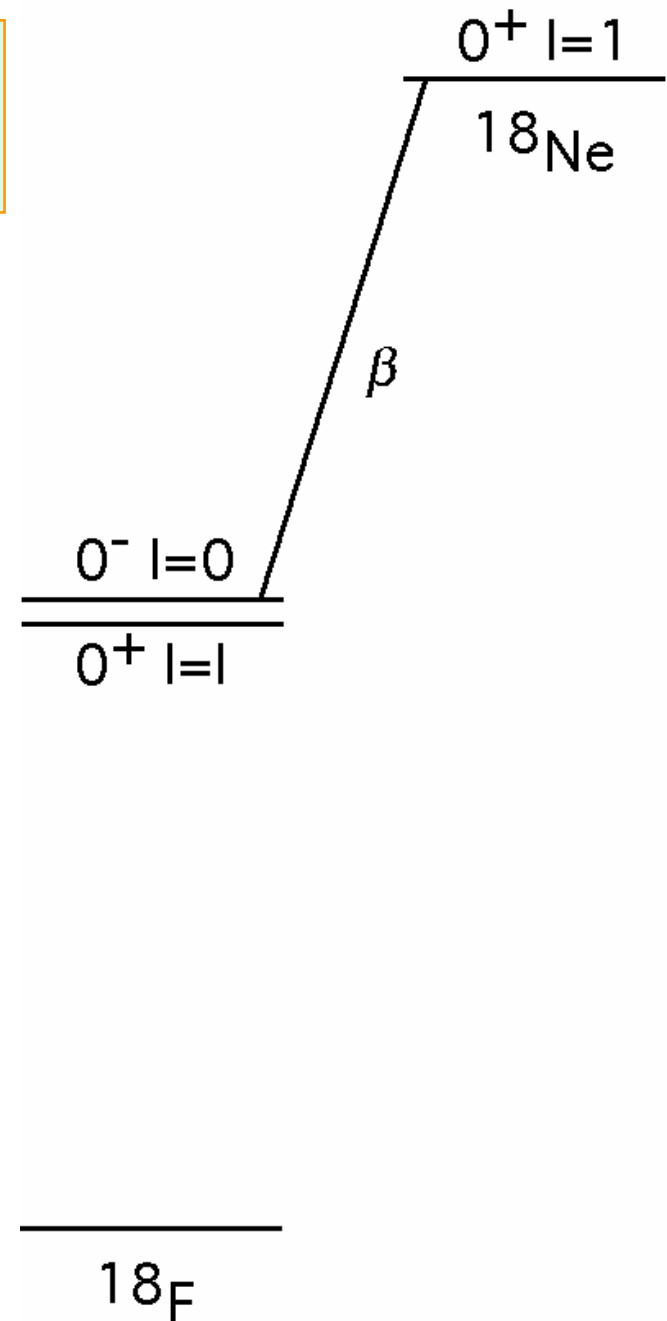
$$P_\gamma = \frac{2}{39 \text{ keV}} \left(\frac{\tau_- E_-^3}{\tau_+ E_+^3} \right)^{1/2} h_\pi \langle + | V_\pi | - \rangle$$

P_γ measured. $\langle + | V_\pi | - \rangle$ is needed.

^{18}Ne G. S. is the IAS of $0^+ I=1$.

β decay amplitude = $aV_\pi(\tau_z \rightarrow \tau_\pm) + b\sigma \cdot p \tau_\pm$

a and b from CVC and PCAC. The ratio of the matrix elements of the 1-body and 2-body matrix elements is independent of the details of the wave function. The measured lifetime determines $\langle + | V_\pi | - \rangle$.



Discussion of ^{21}Ne

- PV circular polarization of the 2789 keV γ in the odd-neutron nucleus ^{21}Ne is consistent with 0.
- One would expect a large asymmetry based on $^{18}\text{F}=0$ and $X_{Np}\neq 0$.
- Both neutron and proton states are active in ^{21}Ne . PV asymmetries involve a theoretically unstable combination of X_{Np} and X_{Nn} . The combination depends on the residual interaction chosen. (A. Brown)
- Although some calculations have $^{21}\text{Ne} \sim X_{Np} + X_{Nn}$, ^{21}Ne can't be used to constrain the HWI.

- The repulsive short-range nucleon-nucleon potential reduces the contributions of the ρ and ω mesons. Deplanques has evaluated the linear combinations of couplings, X 's, that enter in PV in heavy nuclei using nuclear-matter theory for different nucleon-nucleon forces.

X_{Np} (for X_{Nn} change sign of $\Delta l=1$)

| Force | f_{π} | $f_{\rho,0}$ | $f_{\rho,1}$ | $f_{\rho,2}$ | $f_{\omega,0}$ | $f_{\omega,1}$ |
|--------|-----------|--------------|--------------|--------------|----------------|----------------|
| AV18 | 5.5 | -1.13 | -.48 | 0 | -.91 | -.77 |
| RSC | 5.5 | -.89 | -.45 | 0 | -.75 | -.67 |
| T-S | 5.5 | -1.98 | -.81 | 0 | -1.51 | -1.26 |
| Haxton | 5.5 | -1.91 | -.58 | 0 | -1.12 | -.99 |

Consistency of 6 odd-proton nuclei

- p - α , ^{19}F , ^{41}K , ^{175}Lu , ^{181}Ta , and ^{133}Cs are all odd-proton nuclei. All are therefore $\sim X_{N,p}$.

Expt. X_{Np}

$p - \alpha$ 6.1 ± 1.7

^{19}F 7.8 ± 2.0 (coefficient of f_π from experiment)

^{41}K 7.8 ± 1.6

^{133}Cs 13.2 ± 2.3

^{175}Lu $5.9 \pm .5$

^{181}Ta $6.3 \pm .6$

| | Ave. | χ^2/DoF | Pran |
|--------------------|--------------|---------------------|------|
| All 6 | $6.4 \pm .4$ | 11.0/5 | .05 |
| $X^{133}\text{Cs}$ | $6.2 \pm .4$ | 2.1/4 | .73 |

6 parameters and 4 constraints

- We need more independent experiments and/or additional theoretical information
- More theoretical information
 - Take DDH reasonable ranges at face value
 - $f_{\rho,1}$ and $f_{\omega,1}$ enter the expressions for observables with very small coefficients
- Fix $f_{\rho,1}$ and $f_{\omega,1}$ at DDH best values and add the DDH reasonable rang in quadrature to the experimental errors. Now 4 parameters and 4 constraints

Results of 4 parameter fit to 10 measurements

fitted parameters and DDH range

| par | value | error | DDH |
|----------------|-----------|---------|------------------|
| f_{π} | -0.456387 | 0.91383 | 0. 4.6 11.4 |
| $f_{\rho 0}$ | -43.3029 | 8.75909 | -30.8 -11.4 11.4 |
| $f_{\rho 2}$ | 37.0889 | 12.8566 | -11. -9.5 7.6 |
| $f_{\omega 0}$ | 13.698 | 9.38951 | -10.3 -1.9 5.7 |

$\chi^2/\text{DOF} = 7.48286/7$

probability of random occurrence = 0.380391

The AV18 potential was used.

10 Experiments consistently described by three couplings

| Coupling | Value | DDH Range |
|--------------|--------------|----------------------|
| f_π | $-.5 \pm .9$ | $0 \rightarrow 11.4$ |
| $h_{\rho,0}$ | -33 ± 5 | $-31 \rightarrow 11$ |
| $h_{\rho,2}$ | 41 ± 13 | $-11 \rightarrow 8$ |

3 parameter fit, $h_{\omega,0} = 0$, $\chi^2 / \text{DOF} = 9.6 / 8$

Conclusions

- The fit is consistent with the data
- f_{π} is small (we already knew that from ^{18}F)
- $f_{\rho,0}$ is large (at the limit of the DDH range)
- $f_{\omega,0} \neq 0$ is not necessary to describe data
- $f_{\rho,2}$ may be large (2.2 σ outside the DDH range)
- Although $\Delta I=1$ is Cabibbo allowed and $\Delta I=0$ and 2 are Cabibbo suppressed, the fits show the opposite pattern
- It is desirable to determine more linear combinations of couplings

Future work

- Measure anapole moments in closed shell odd proton and odd neutron
 - Check theory of anapole moments - ^{133}Cs is a very complex nucleus
 - Check the one-body approximation. $^{209}\text{Bi} \sim X_{Np}$ and $^{207}\text{Pb} \sim X_{Nn}$
- Measure γ circular polarization in $n+p \rightarrow d + \gamma$
 - Constrains $f_{\rho,2}$
- Measure PV observables in neutron reactions
 - Asymmetry in $n+p \rightarrow d + \gamma$, constrains f_{π}
 - 2-body and few-body asymmetries
 - Use few-body methods to evaluate the asymmetries.

Absolutely necessary to plan and interpret experiments!

Critique of DDH

- The possible spin-isospin structure of the PV interactions is fixed (Herczeg)
- DDH theory is a model. Assumes a particular momentum dependence for the interactions
- No demonstration that the model is complete or that the terms correspond to the physical light mesons
- $2-\pi$ exchange neglected
- Interpretation of many-body systems involves nuclear models (except for ^{18}F and ^{19}F)

EFT

- In principle couplings can be calculated using QCD
- A theory based on systematic expansion in low-energy constants. Early version had 10-12 constants
- C. P. Liu theory has f_π and 5 LEC's corresponding to S-P scattering amplitudes (Danilov parameters).
- Theory applies for energies < 40 MeV (can't use p-p 220 MeV)
- Liu has calculated all two-body PV observables
- Greens' function Monte Carlo method can reliably calculate PV matrix elements for few-body systems. Calculations are essential to plan and interpret experiments

Feasible two-body experiments

- p-p $s_p \bullet k_p$ done
- n+p \rightarrow d+ γ $s_n \bullet k_\gamma$ phase 1 at LANL done
phase 2 proposed at SNS
- n+p \rightarrow d+ γ γ CP FEL or intense n source +
improved CP polarimeter
- n+p s_n rot. next-generation

Feasible few-body experiments

- $p+\alpha$ $s_p \bullet k_p$ done
- $n+\alpha \rightarrow n+\alpha$ s_n rot. preparing NIST
- $n+d \rightarrow t+ \gamma$ $s_n \bullet k_\gamma$ consideration at SNS
- $n+d \rightarrow t+ \gamma$ γ CP see 2-body
- $n+{}^3\text{He} \rightarrow p+t$ $s_n \bullet k_n$ consideration at SNS
- $n+{}^3\text{He} \rightarrow p+t$ $s_3 \bullet k_n$ next-generation
- $n+d$ s_n rot. preliminary expts. done

EFT

- Work out and publish spin-isospin content of EFT
 - Which couplings determine $\Delta I=0, 1, \text{ and } 2$
- Work out one-body approximation for EFT in order to include nuclear PV constraints
 - Expect that $\Delta I=2$ is absent for nuclei
 - Expect that X 's depend on nuclear force
 - Isospin and density dependence?
- Work out few-body observables in order to plan and interpret experiments

- EFT provides a rigorous framework for understanding HWI
 - 6 parameters
 - Valid for E Less than 40 MeV
- Two-body calculations done
- Need few-body calculations to design and interpret experiments
- Need existing and proposed 2 and few-body experiments to constrain f_{π} and 5 LEC's
- couplings reveal short-range structure of q-q correlations

Conclusions

- Experiments require small f_π and large $h_{\rho,0}$ and $h_{\rho,2}$. $\Delta I=1$ is Cabibbo allowed and $\Delta I=0$ and 2 are suppressed!
- The small $\Delta I=1$ is solid. More assumptions are required for $\Delta I=0$ and 2 . $\Delta I=0$ is large and $\Delta I=2$ may be large.
- W and Z exchange are short-range. Above pattern is telling us something about the short-range behavior of q - q correlations in the non-perturbative regime.