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Breakup and the Spectroscopy of Continuum States

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Which Continuum?

- <u>The continuum</u> appears in several ways:
 - Part of expansion of bound states;
 - eg needed in RPA for weakly bound states
 - Dominated by <u>resonances;</u>
 - These `unbound states' identified eg with shell model eigenstates above threshold
 - In <u>non-resonant</u> continuum;
 - eg in breakup reactions, or low-energy capture.
- ALL important parts of nuclear structure!!

'Overlap' Challenges

- Reaction models need few-body degrees of freedom in structure models.
 - Solve a few-body model directly, or
 - Extract few-particle dof from microscopic model
 - Difficult for: GFMC,
 - for HFB, QRPA and RMF structure models
 - Do we transfer quasi-particles, or particles?

Unis What a good theory needs:

- Recoil & Finite Range of projectile vertex.
- Final-state (partial wave) interference
- Nuclear and Coulomb mechanisms
- Core excitation (initial and/or dynamic)
- Final-state interactions:
 - between halo fragments (needed if resonances)
 - between fragments and target (needed if close in)
- <u>Multistep Processes</u> (higher order effects)

CDCC: Coupled Discretised Unis Continuum Channels

Try CDCC:

Coupled Discretised Continuum Channels

- Proposed by Rawitscher, developed by Kamimura group.
- Treat Coulomb and Nuclear mechanisms
 - Need to check convergence of long-range Coulomb process!
- All higher-order effects with a (r,R,L) reaction volume
- Can calculate fragment coincident angular distributions: Predict e.g. $d^3\sigma/dE_1d\Omega_1d\phi_{12}$ and fold with detector apertures & efficiencies

CDCC Hamiltonian & model space UniS

The Hamiltonian for the reaction of a projectile on a target

$$\mathbf{H} = \mathbf{h}_{\text{proj}} + \mathbf{h}_{\text{targ}} + \mathbf{T}_{\alpha} + \mathbf{V}_{\alpha}$$

$$\Rightarrow \mathbf{h}_{\text{proj}} = \mathbf{h}_{\text{core}} + \mathbf{h}_{\text{frag}} + \mathbf{T}_{\text{cf}} + \mathbf{V}_{\text{cf}}$$

$$\Rightarrow \mathbf{V}_{\alpha} = \mathbf{V}_{\text{core-targ}} + \mathbf{V}_{\text{frag-targ}}$$

$$\psi_{\text{JM}}^{\text{CDCC}}(\mathbf{r}, \mathbf{R}) = [\phi_0(\mathbf{r}) \otimes \mathbf{Y}_{\text{L}}(\hat{\mathbf{R}})]_{\text{JM}} \chi_{0,\text{L}}^{\text{J}}(\mathbf{R}) + \sum_{l=0}^{\text{Imax}} \sum_{\mathbf{L}} \sum_{i=1}^{N} [\phi_{i,l}(\mathbf{r}) \otimes \mathbf{Y}_{\text{L}}(\hat{\mathbf{R}})]_{\text{JM}} \chi_{i,l,\text{L}}^{\text{J}}(\mathbf{R})$$

(neglect the internal structure of the target)

CDCC Formalism

The CDCC basis consists of scattering wavefunctions averaged over an energy interval

$$\mathbf{h}_{\text{proj}} \mathbf{\phi}_{\mathbf{k}} = \mathbf{\varepsilon}_{\mathbf{k}} \mathbf{\phi}_{\mathbf{k}}$$
$$\mathbf{\phi}_{\mathbf{i},\mathbf{l}} = \sqrt{\frac{2}{\sqrt{2\pi}}} \int_{\mathbf{k}_{i-1}}^{\mathbf{k}_{i}} \mathbf{w}_{\mathbf{i}}(\mathbf{k}) \mathbf{\phi}_{\text{lm}}(\mathbf{k},\mathbf{r}) d\mathbf{k}$$



$$\mathbf{N}_{i} = \int_{\mathbf{k}_{i-1}}^{\mathbf{k}_{i}} |\mathbf{W}_{i}(\mathbf{k})|^{2} d\mathbf{k}$$
$$\mathbf{N}_{bins} = \frac{\mathbf{k}_{max}}{\Delta \mathbf{k}}$$

Coupling potentials

$$V_{il,i'l'}^{CDCC}(R) = \langle \phi_{il}(r) | V_{\alpha}(r,R) | \phi_{i'l'}(r) \rangle$$

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Uni<mark>S</mark> Testing CDCC Convergence

 Compare, in Adiabatic Few-Body Model, with Bremstrahlung integral



 Compare, in first-order PWBA model, with semiclassical theory



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Adiabatic CDCC: compare with Exact 3-body model



d+²⁰⁸Pb at 50 MeV, nuclear only



Absolute errors in CDCC for d+²⁰⁸Pb at 50 MeV, Nuclear+Coulomb

¹⁵C + 9Be breakup at 54 MeV/u Uni<mark>S</mark>



FIG. 4. Diagrammatic representation of the CDCC model space calculation for ¹⁵C. The left side shows the physical bound states and continuum and the right hand side the included continuum bins (10) in each $n + {}^{14}$ C partial wave. The dashed arrows are representative of the one-way couplings included in the DWBA. The solid arrows show representative couplings for the full CDCC calculations which connect all bins, including diagonal bin couplings, with two-way couplings to all orders. Relative *h* waves were found to make negligible contributions.



FIG. 11. The nucleon-removal parallel momentum distributions $d\sigma/dp_{\parallel}$, for the ¹⁵C+ ⁹Be reaction at 54 MeV/nucleon to the ¹⁴C ground state, shown on a more familiar linear scale. The solid curves assume the stripping contributions have the same form as that calculated using the CDCC. The dashed curves assume the stripping contributions have a parallel momentum distribution at all angles of the residue given by the eikonal calculation shown in Fig. 2.

Tostevin et al, PRC 66, 024607 (2002)

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¹¹Be + ¹²C breakup at 67 MeV/u UniS



Angular distributions of ¹¹Be* <u>left</u>: low-energy continuum <u>right</u>: region of d_{5/2} resonance

CDCC calculations of Howell, Tostevin, Al-Khalili, J. Phys. G **31** (2005) S1881

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Energy excitation spectrum dashed line: multiplied by 0.8



CDCC: sub-Coulomb ⁸B+⁵⁸Ni (26 MeV)

 Multistep Coulomb only Multistep Nuclear Only



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Coulomb+nuclear

Effect of continuumcontinuum couplings



Green lines: no continuum-continuum couplings

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Convergence: max bin E_{rel}



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Elastic Breakup: $\sigma(\theta)$

⁸B breakup on ⁵⁸Ni (E_{beam}=26 MeV)

3-body observables

sensitivity to ⁸B
 structure: overall
 normalisation

 sensitivity to p-target optical potential at larger angles



[Tostevin, Nunes and Thompson, PRC (2001) 024617]

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Breakup reactions CDCC ⁸B + ⁵⁸Ni → ⁷Be+p + ⁵⁸Ni (E_b=26 MeV)



E1 & E2 breakup of ⁸B

- One-proton bound state known:
 - $-^{7}\text{Be}\otimes(0p_{3/2}+0p_{1/2})|_{2+}$ at -0.137 MeV
- Need spectroscopy of non-resonant continuuum!
 - B(E1) & B(E2) for transition p→s,d need to be accurately known
 - E1 and E2 amplitudes interfere in p_{||}(⁷Be) momentum distribution
 - so measure relative E2/E1 amplitudes from asymmetries.

⁸B + ²⁰⁸Pb → ⁷Be parallel momentum distributions



from Mortimer et al., Phys Rev C 65 (2002) 64619

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Extensions started

- Core excitation (static, dynamic)
 - Glauber: Batham et al
 - CDCC bins of particle+core coupled states, Summers & Nunes at MSU
- Three-cluster projectiles

 (e.g. two-neutron halo nuclei):
 - Gaussian expansions: Kamimura et al.
 - Transformed Harmonic Oscillator: Rodriguez-Gallardo et al

Wave functions of ⁶He

- Ground state wave function:
- Solution of coupled equations for E ~ -0.97 MeV.



Nuclei such as ⁶He have highly correlated cluster structures



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1 Neutron stripping from three-body Borromean Nuclei

- Removal of a neutron from ⁶He, ¹¹Li, ¹⁴Be,
 - populates states of ⁵He, ¹⁰Li or ¹³Be.
 - Experiments measure decay spectrum of 5 He = 4 He + n, 13 Be = 12 Be + n, etc
- Can we predict any energy and angular correlations by Glauber model?
- Can we relate these correlations to the structure of the A+1 or the A+2 nucleus?

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1N stripping from ⁶He g.s.

- Calculate overlaps: <⁵He(E_{α-n}) |⁶He(gs)> for a range of ⁵He(E_{α-n}) bin states,
- smooth histogram of Glauber bin cross sections.
- GSI data (H.Simon)

Promising technique!

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Theory: σ_{str} =137 mb, σ_{diff} =38 mb Expt: σ_{str} =127±14 mb, σ_{diff} =30±5 mb from T. Tarutina thesis (Surrey)

s₁₀ component

d_{5/2} component
 -- d_{3/2} component

smoothed bins

--- total bins

E_{12Be−n}, Me∛

Experiment

1N stripping from ¹⁴Be g.s.

200

dσ/dE_{12Bθ-n}, mb/MeV ¹²⁰ ¹²¹ ¹²¹

- Calculate overlaps: <¹³Be(E_{α-n}) |¹⁴Be(gs)>
- Inert-core ^{13,14}Be wfs.
- GSI data (H.Simon)
- See softer data, and <u>not</u> pronounced virtual-s and resonant-d peaks.



0

0

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Elastic Breakup of 2N halo

- Elastic Breakup = <u>Diffraction</u> <u>Dissociation:</u>
 - all nuclear fragments survive along with the target in its ground state,
 - probes continuum excited states of nucleus.
- Need correlations in the three-body continuum of Borromean nuclei.

Continuum Spatial Correlations

from B. Danilin, I. Thompson, PRC 69, 024609 (2004)

 Now average scattering wave functions over angles of k_x and k_y, to see spatial correlations in continuum states in ⁶He:

T-basis 2+ plane wave T-basis 2+ resonance







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`True' 3-body resonances?

 Expect continuum wave functions like:

$$\psi(
ho \ \Omega_5^{
ho}, E \ \Omega_5^{\kappa})$$

 $\propto rac{1}{(\kappa
ho)^{5/2}} \sum_{K,\gamma} C_{K\gamma}(E) \ \psi_{K\gamma}^R(
ho) \ Y_{K\gamma}(\Omega_5^{
ho}) \ Y_{K\gamma}(\Omega_5^{\kappa})$

with

$$|C_{K\gamma}(E)|^2 = \frac{\Gamma_{K\gamma}}{(E - E_0)^2 + \Gamma^2/4}$$

Continuum Energy Correlations

- Now average scattering wave functions over angles of k_x and k_y, for fixed three-body energy E.
- Obtain similar plots for continuum energies.
- (Continuum momentum and angular correlations for later)

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Continuum three-body wave Unis functions

• Three-body scattering at energy E:

hypermomentum $\kappa=\sqrt{k_x^2+k_y^2}=\sqrt{2mE/\hbar^2}$, hyperangle $lpha_\kappa = {
m atan}(k_x/k_y)$

• Plane wave 3-3 scattering states: $(2\pi)^{-3} \exp[i(\mathbf{k}_{\mathbf{x}} \cdot \mathbf{x} + \mathbf{k}_{\mathbf{y}} \cdot \mathbf{y})]$ $= (\kappa\rho)^{-2} \sum i^{K} J_{K+2}(\kappa\rho) \mathcal{Y}_{KLM_{L}}^{l_{x}l_{y}}(\Omega_{5}^{\rho}) \mathcal{Y}_{KLM_{L}}^{l_{x}l_{y}}(\Omega_{5}^{\kappa})^{*}$ $KLM_L l_x l_y$ Dynamical solutions for scattering states:

$$\Psi_{\kappa JM}^{T}(\mathbf{x}, \mathbf{y}, \hat{\mathbf{k}}_{x}, \hat{\mathbf{k}}_{y}, \alpha_{\kappa}) = (\kappa \rho)^{-5/2} \sum_{K\gamma, K'\gamma'} \psi_{K\gamma, K'\gamma'}^{J}(\kappa \rho) \Upsilon_{JM}^{K\gamma}(\Omega_{5}^{\rho})$$

$$\sum_{M'_{L}M'_{S}} \langle L'M'_{L}S'M'_{S} \mid JM \rangle \mathcal{Y}_{K'L'M'_{L}}^{l'_{x}l'_{y}}(\Omega_{5}^{\kappa}) X_{T}$$
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Effect of n-n 'resonance' in E(c-n), E(cn-n) coordinates

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Virtual n-n pole

6.7 0.8 0.9 E(CB-B) [MeV]

0.5 0.6

⁶He excitations & resonances UniS





Pronounced 2+ resonance

No pronounced 1⁻ resonance

Four-body dynamics

- High Energies (first order & all order):
 - T-matrix multiple scattering (Crespo)
 - Eikonal+Adiabatic (Tostevin, Al-Khalili)
 - Eikonal (Exact fragment) (Brooke, Tostevin, Al-Khalili)
 - Adiabatic (Johnson, Christley et al)
- All Energies (all orders), new challenges:
 - 4-body pseudo-state CDCC (Kamimura)
 - 4-body bin-states CDCC
 - <u>"Two-nucleon states in deformed nuclei"</u>

T-matrix expansions for breakup

¹¹Li on protons at 68 MeV/A

- Preliminary Method:
 - First-order expansion on fragment-target T-matrices
 - Pseudo-state continuum, smoothed.
- Strong sensitivity on the structure models for ¹¹Li: S, P0, P2



Crespo et al., PRC66 (2002) 021002 Data: **RIKEN**.

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σ(θ) for ¹¹Li(p,p') at 68 MeV/u



- (a) Comparison of the theoretical calculations with experimental data
- Solid, dashed and dotted lines show the total, monopole and dipole angular distributions, respectively.
- In (b) and (c), solid lines show angular distributions for the monopole and dipole excitations, respectively.
- Dashed and dotted lines are contributions from the halo neutrons and the core nucleons.

Conclusions

- CDCC method good for 2-cluster halo nuclei:
 - Finite-range & recoil included
 - <u>Coulomb and nuclear</u> both approach convergence
 - Large radii and partial-wave limits needed, but feasible now
 - Non-adiabatic treatment of Coulomb breakup
 - <u>Multistep</u> effects manifest from <u>all final-state</u> interactions
- Extensions:
 - Deformed cluster models: Summers & Nunes at MSU
 - Three-cluster projectiles (e.g. two neutron halo nuclei): Kamimura et al.

