

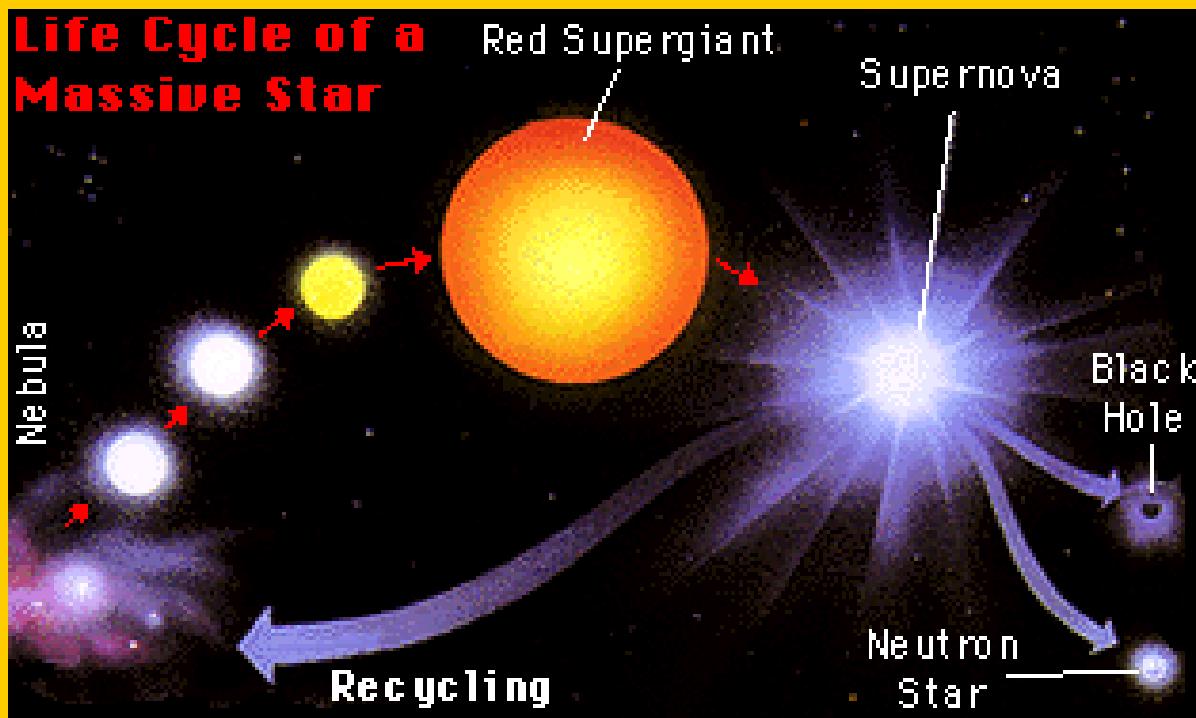


Hans A. Bethe (1906–2005)

Theoretical Research on the
Nuclear Astrophysics of Compact Stars
at the **University of Milano**

Glitches and vortex pinning in superfluid pulsars

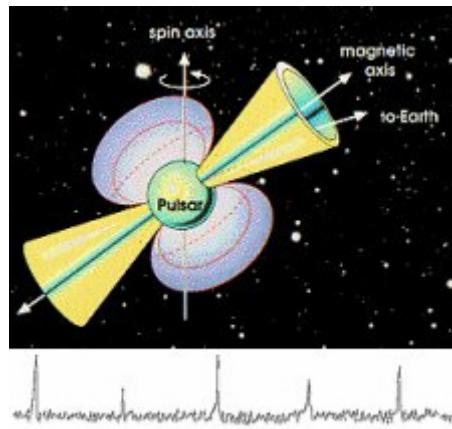
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INT Seattle, 14-17 November 2005

Pulsar glitches

► Pulsar spin-down



emission of e.m. and gravitational waves $\rightarrow \dot{E}_{\text{em}} > 0$

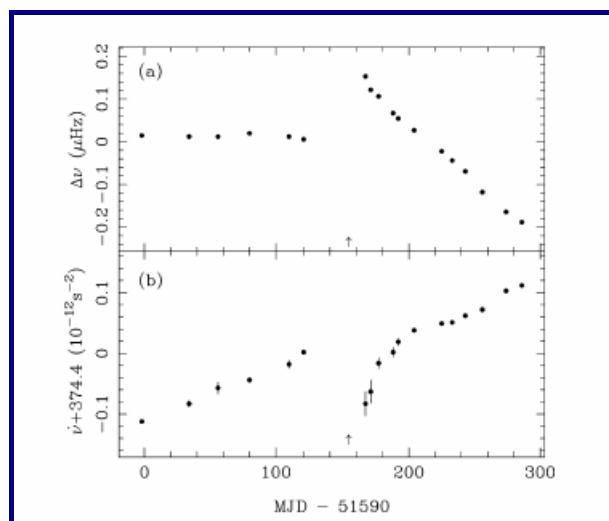
$$\text{energy conservation} \rightarrow \dot{E}_{\text{em}} + \frac{d}{dt} \left(\frac{1}{2} I \Omega^2 \right) = 0$$

\Downarrow

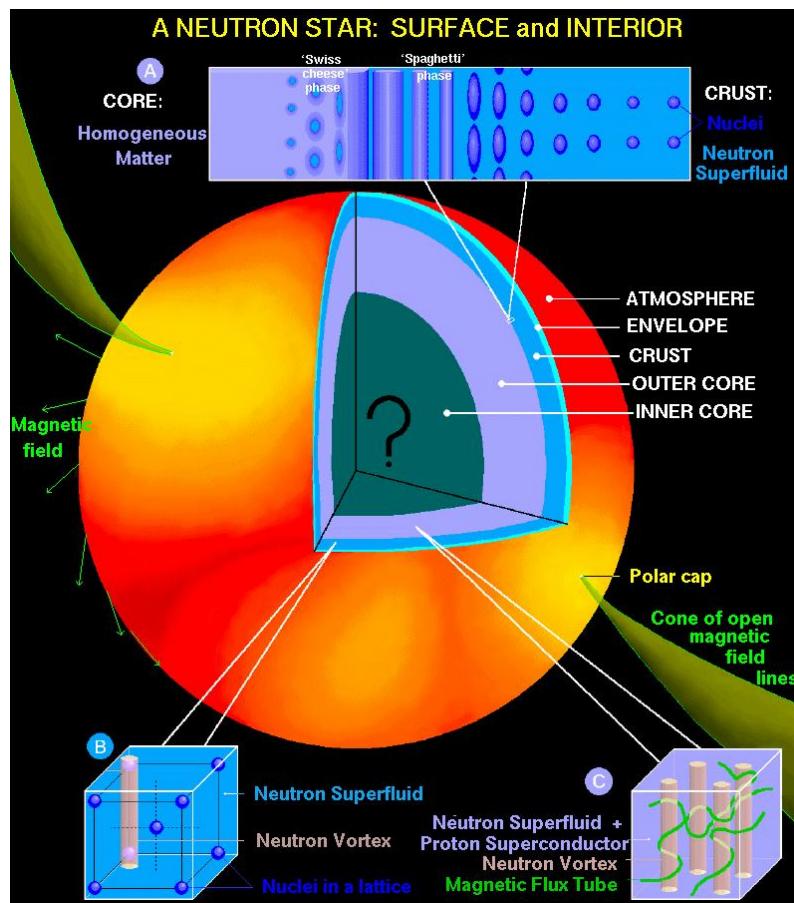
decrease of rotational frequency $\rightarrow \dot{\Omega} < 0$

► Rotational glitches

spin-ups $\rightarrow \frac{\Delta\Omega}{\Omega} \approx 10^{-6}$ (Vela) $\frac{\Delta\Omega}{\Omega} \approx 10^{-8}$ (Crab)



Superfluidity in Neutron Stars



★ **Inner crust** ($\rho_{\text{drip}} < \rho < 0.6 \rho_{\text{nuc}}$)

gas of unbound superfluid neutrons (1S_0 pairing) $\rightarrow n_G$

lattice of neutron-rich nuclei (Wigner-Seitz cells) $\rightarrow R_N, R_{WS}$

ρ_B	1.5×10^{12}	9.6×10^{12}	3.4×10^{13}	7.8×10^{13}	1.3×10^{14}
n_G	4.8×10^{-4}	4.7×10^{-3}	1.8×10^{-2}	4.4×10^{-2}	7.4×10^{-2}
R_{WS}	44.0	35.5	27.0	19.4	13.77
R_N	6.0	6.7	7.3	6.7	5.2
a	0.77	0.83	0.94	1.12	1.25
N	280	1050	1750	1460	950
Z	40	50	50	40	32
N_{bound}	110	110	110	70	40

Negele & Vautherin (1973)

Quantized vortex lines

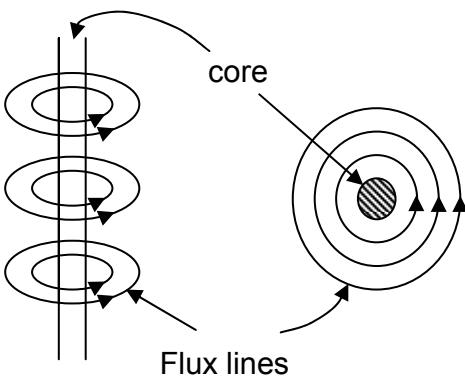
→ Quantized vortex lines

superfluidity → irrotational flow → $\nabla \times \vec{v}_s = 0$

classical vortex → $\vec{v}_s = \frac{C}{r} \hat{e}_\theta \rightarrow \nabla \times \vec{v}_s = 2\pi C \delta^{(2)}(\vec{r})$

\vec{v}_s minimizes $E_\Omega = E_{\text{lab}} - \vec{L}_{\text{lab}} \cdot \vec{\Omega}$ (if $\Omega > \Omega_{\text{cr}}$)

\vec{v}_s singular at $r=0 \rightarrow$ **vortex core** (empty or normal matter)



quantized vortex lines → $C = k \frac{\hbar}{2m_N} \quad (k=1,2,\dots)$

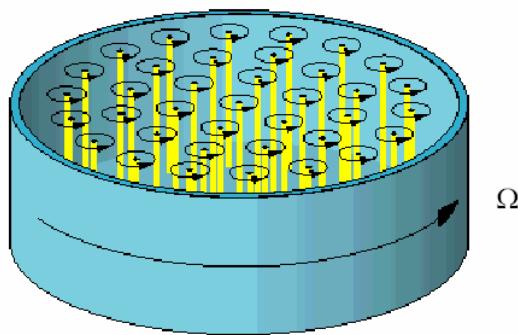
quantized vorticity → $\oint \vec{v}_s \cdot d\vec{l} = k \frac{\hbar}{2m_N} \quad (k=1,2,\dots)$

Rotations of a superfluid

Rotating vessel: Feynman-Onsager formula

$$\nabla \times (\vec{\Omega} \times \vec{r}) = 2\vec{\Omega} \neq 0 \rightarrow \text{no rigid rotations } \vec{\Omega}$$

array of parallel vortices (if $\Omega \gg \Omega_{\text{cr}}$) $\rightarrow \langle \vec{v}_s \rangle = \sum_i \vec{v}_{si}$



$$E_\Omega \text{ minimized by } \langle \vec{v}_s \rangle = \vec{\Omega} \times \vec{r}$$

\Downarrow

$$\text{uniform density of vortices} \rightarrow n_v = \frac{N_v}{\pi R^2} = \frac{4m_N}{h} \Omega$$

$\Downarrow \Downarrow$

superfluid angular momentum quantized in vortices

Vortex theory for glitches

► Vortex pinning and Magnus force

if superfluid vortices are pinned $\rightarrow \dot{n}_v = 0 \rightarrow \dot{\Omega}_s = 0$

but slow-down of normal component $\rightarrow \dot{\Omega}_n < 0$



rotational lag of components $\rightarrow \Delta\Omega = \Omega_s - \Omega_n > 0$

outward drag force on vortex $\rightarrow f_{\text{mag}} \propto \Delta\Omega$

► Vortex un-pinning and glitches

since $\Delta\dot{\Omega} > 0 \rightarrow f_{\text{mag}}$ increases with time

pinning energy \rightarrow maximum pinning force f_{pin}

when $f_{\text{mag}} \geq f_{\text{pin}} \rightarrow$ unpinning of many vortices



transfer of angular momentum to the star surface



normal component spin-up \Rightarrow pulsar glitch

Self-repeating mechanism !

Glitches as evidence of macroscopic superfluidity !

Vortex-nucleus interaction

Why is there an interaction?

energy density is density-dependent $\rightarrow \varepsilon = \varepsilon(n)$

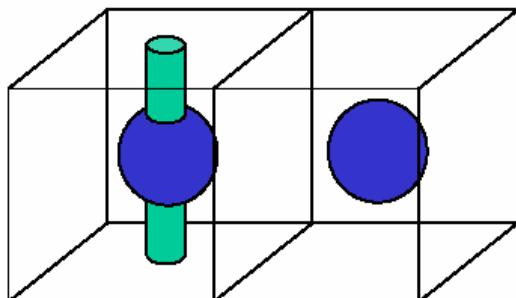
inner crust (nuclei + neutron gas) $\rightarrow n_N \gg n_G$



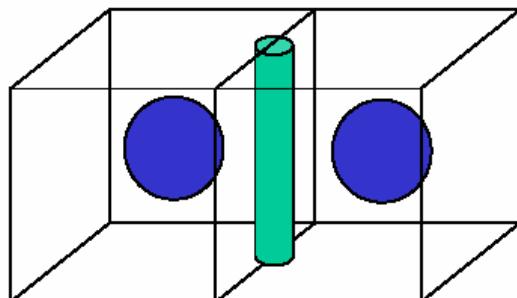
presence and position of nuclei modify $\varepsilon(n)$

Configurations and pinning energy

vortex-nucleus configurations in nuclear lattice



Nuclear pinning



Interstitial pinning



pinning energy $\rightarrow \Delta E_{\text{pin}} = E_{\text{NP}} - E_{\text{IP}}$



$\Delta E_{\text{pin}} > 0 \Rightarrow$ *interstitial* pinning \Rightarrow negligible

$\Delta E_{\text{pin}} < 0 \Rightarrow$ *nuclear* pinning \Rightarrow real pinning

Pinning energy: early models

Simple-model assumptions

uniform densities and pairing gap $\rightarrow n_N, n_G, \Delta = \Delta(n)$

superfluid-normal energy difference $\rightarrow \varepsilon_s - \varepsilon_n = \varepsilon_{\text{kin}} + \varepsilon_{\text{cond}}$

$$\text{kinetic term} \rightarrow \varepsilon_{\text{kin}}(n) = \frac{B}{r^2} n$$

$$\text{condensation term} \rightarrow \varepsilon_{\text{cond}}(n) = -\frac{3}{8} \frac{\Delta^2}{A} n^{\frac{1}{3}}$$

$$(A = 200 \text{ MeV fm}^2 \quad B = 5.17 \text{ MeV fm}^2)$$

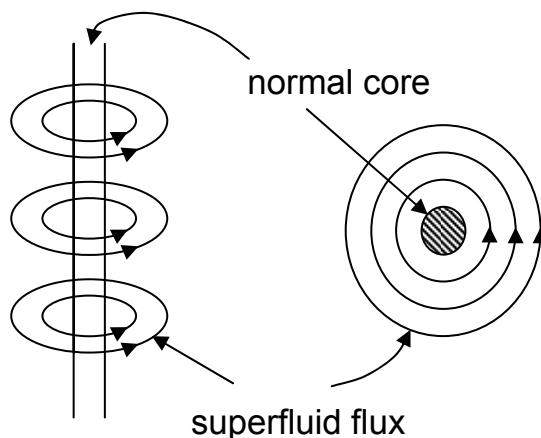
Vortex structure

$$\varepsilon_{\text{kin}}(n) + \varepsilon_{\text{cond}}(n) = 0 \rightarrow \text{core radius} \rightarrow r_c = r_c(n) \quad (= \xi_{\text{BCS}})$$



$$r > r_c \rightarrow \varepsilon_s - \varepsilon_n < 0 \rightarrow \text{superfluid vortex}$$

$$r < r_c \rightarrow \varepsilon_s - \varepsilon_n > 0 \rightarrow \text{normal core at rest}$$



► Naïve pinning energy (Alpar, 1977)

simple assumption → ΔE_{pin} related *only* to ΔE_{cond}

ΔE_{cond} due to $n_G \Rightarrow n_N$ over the nucleus volume V_N



$$\Delta E_{\text{pin}} \equiv [\epsilon_{\text{cond}}(n_N) - \epsilon_{\text{cond}}(n_G)] \cdot V_N$$

► Detailed model (Epstein & Baym, 1988)

realistic density profile for nuclei → $n_N(R)$

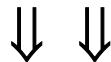
Ginzburg-Landau for neutron pairing (*but* $T \ll T_c$!)

formula for ΔE_{kin} due to $n_G \Rightarrow n_N$



Zone →	1	2	3	4	5
ρ	1.5×10^{12}	9.6×10^{12}	3.4×10^{13}	7.8×10^{13}	1.3×10^{14}
ΔE_{pin}	-	0.4	15	9	5.4
F_{pin}	-	0.11	3.6	1.9	1.7

ρ in g/cm^3 , ΔE_{pin} in MeV, F_{pin} in MeV/fm



nuclear pinning for $\rho \geq 10^{13} \text{ g/cm}^3$

large pinning energies and pinning forces

(*but* large rescaling factors are needed in G-L !)

Semi-classical consistent model

► Model assumptions (Donati & Pizzochero, 2004)

Local Density Approximation ($T = 0$) $\rightarrow p_F(\vec{x}) = \hbar(3\pi^2 n(\vec{x}))^{1/3}$

realistic density profile for nuclei $n_N(R) \rightarrow U_N(R)$

all contributions to energy $\rightarrow \varepsilon_{\text{cond}}(n), \varepsilon_{\text{kin}}(n), \varepsilon_{\text{int}}(n)$

BCS neutron pairing gap (Argonne) $\rightarrow \Delta = \Delta(n)$

require thermodynamical and mechanical equilibrium

► Density profiles determination

find chemical potential $\rightarrow \mu(\vec{x}) = \frac{\delta \varepsilon_{\text{tot}}(n(\vec{x}))}{\delta n(\vec{x})}$

Thomas-Fermi ansatz $\rightarrow \mu(\vec{x})$ constant

at given \vec{x} invert equation $\frac{\delta \varepsilon_{\text{tot}}(n(\vec{x}))}{\delta n(\vec{x})} = \mu \Rightarrow n(\vec{x})$

total neutron number $\rightarrow N = \int_V n(\vec{x}) d^3\vec{x} \Rightarrow \mu = \mu(\rho)$

$\Downarrow \Downarrow$

Thomas-Fermi \Rightarrow mechanical equilibrium

► Vortex core structure

two consistent vortex scenarios → surface Σ_{core}



empty core

$$n_s(\Sigma_{\text{core}}) = 0$$

pure phase

normal matter core

$$P_s(\Sigma_{\text{core}}) = P_n(\Sigma_{\text{core}})$$

mixed phase



Σ_{core} axially symmetric but *not* cylindrical in NP

Model-consistent and unambiguous core surface !

► Equilibrium phase determination

for any given WS-cell and configuration (NP or IP)



calculate for both phases → E, N

at $T = 0$ and fixed μ → potential $\Omega = E - \mu N$



for equilibrium state → minimize $\Omega = E - \mu N$



$\Omega_{\min} \Rightarrow$ thermodynamical equilibrium

Beyond the consistent model

❖ Realistic pinning (Donati & Pizzochero, 2005)

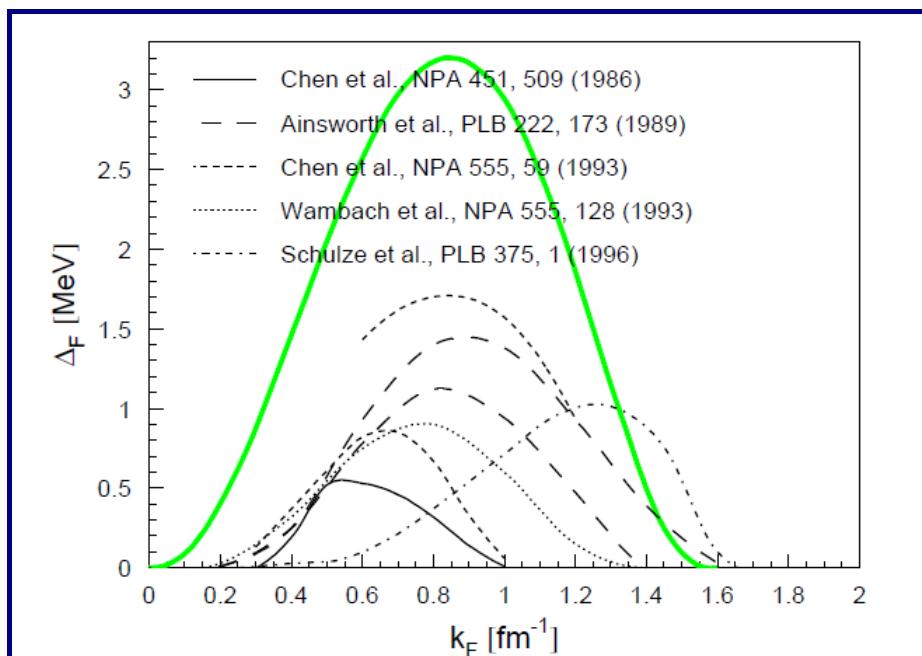
few unphysical approximations in fully-consistent model
supplement model with general properties



more realistic pinning parameters

❖ More realistic in-medium pairing

medium polarization → self-energy & vertex corrections



Lombardo & Schulze (2001)

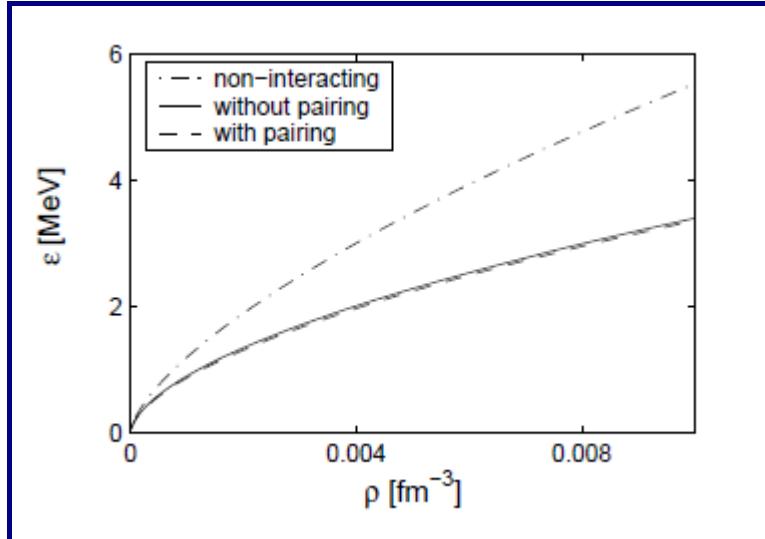
reduction of BCS gap by factor between 2 and 3



rescale Argonne gap → $\Delta_\beta = \frac{1}{\beta} \Delta_{\text{Argonne}}$ with $2 \leq \beta \leq 3$

More realistic scale of energies

LDA → non-interacting neutrons (Fermi model)



Bulgac & Yu (2003)

Fermi systems with $a \rightarrow -\infty$ and $r_0 \ll |a|$

intermediate-density regime → $r_0 < k_F^{-1} < |a|$



general property → $\frac{E}{N} = \alpha \frac{3}{5} E_F$ with $\alpha \cong 0.5$

Zone	1	2	3	4	5
ρ_B	1.5×10^{12}	9.6×10^{12}	3.4×10^{13}	7.8×10^{13}	1.3×10^{14}
n_G	4.8×10^{-4}	4.7×10^{-3}	1.8×10^{-2}	4.4×10^{-2}	7.4×10^{-2}
$k_{F,G}^{-1}$	4.2	1.9	1.2	0.9	0.8

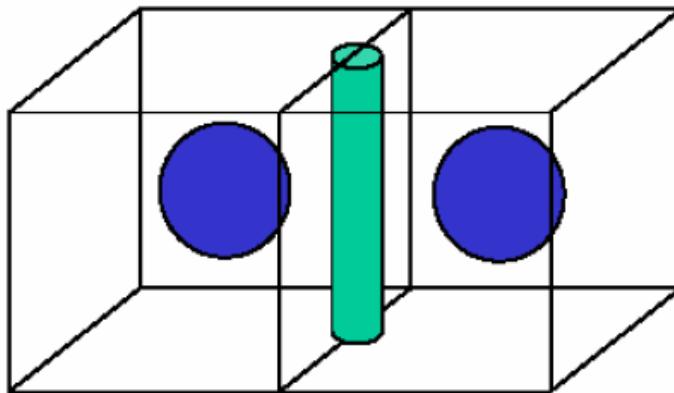
ρ_B in g/cm³, n_G in fm⁻³, $k_{F,G}^{-1}$ in fm



rescale Fermi term → $\varepsilon_\alpha(\vec{x}) = \alpha \varepsilon_{\text{fer}}(\vec{x})$ with $\alpha = 0.5$

◆ More realistic vortex-nucleus kinetics

continuity equation at nucleus → cut-off on vortex flow



approximate analytical formula for nucleus in vortex flow



$$\Delta K_{nucleus} \propto \frac{1}{D^2}$$

◆ Caution: **collective pinning!**

LDA → vortex core *always* normal → $r_{\text{core}} \approx \xi_{\text{BCS}}$

Zone	1	2	3	4	5
$\beta = 2$	13.4	8.7	10.3	22.3	77.6
$\beta = 3$	20.0	13.0	15.4	33.5	116.4
R_{WS}	44.0	35.5	27.0	19.4	13.77

$\xi_{\text{BCS}}(n_G)$ and R_{RW} in fm

at large densities (zones 4 & 5) → $r_{\text{core}} > R_{\text{WS}}$



vortex core contains *several* nuclei

Results of the realistic model

◆ Pinning energies and forces^(*)

only NP significative $\rightarrow \Delta E_{\text{pin}} = E_{\text{NP}} - E_{\text{IP}} < 0$

$$\text{pinning energy} \rightarrow E_p = |\Delta E_{\text{pin}}|$$

$$\text{average pinning force} \rightarrow F_p = \frac{E_p}{R_{\text{ws}}}$$

Zone	E_p				F_p			
	$\beta = 2$	$\beta = 3$	Paper I	Ref. [5]	$\beta = 2$	$\beta = 3$	Paper I	Ref. [5]
1	-	-	-	-	-	-	-	-
2	-	-	-	0.4	-	-	-	0.11
3	3.1	2.7	5.2	15	0.11	0.10	0.19	3.6
4	1.8	0.7	5.1	9	0.09	0.04	0.26	1.9
5	-	-	0.4	5.4	-	-	0.03	1.7

E_p in MeV, F_p in MeV/fm, Paper I $\rightarrow \alpha=\beta=1$, [5] \rightarrow Epstein & Baym, 1988



nuclear pinning only for $10^{13} < \rho < 10^{14} \text{ g/cm}^3$

small pinning energies and pinning forces



limited region of weak pinning

(*) P. Donati & P. Pizzochero, Phys.Rev.Lett. **90**(2003)211101;
Nucl.Phys. **A742**(2004)363;
Phys.Lett. **B**(2005)submitted.