# Isospin dependence of the mean field from DBHF

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#### **Outline**



- review DBHF
- nuclear/neutron matter
- effective masses
- optical potentials

## **Saturation of Nuclear Matter**

#### DBHF: realistic NN force, no parameter





Coester line  $\implies$  relativistic!

## Hadronic many-body theory

Relativistic Brueckner: N+OBEP ( $V = \sigma, \omega, \pi, \rho, \eta, \delta$ )  $\implies$  2-N correlations in hole-line expansion  $\implies$  self-consistent sum of ladder diagrams

**Dyson-Equation:**  $G = G_0 + G_0 \Sigma G$ 

Bethe-Salpeter-Equation:  $T = V + i \int VGGQT$ 

Self Energy (Hartree-Fock):  $\Sigma(\rho, k) = \sum_{q \in F} \langle q | T(q, k) | q \rangle = \Sigma_S - \gamma_0 \Sigma_0 + \vec{\gamma} \cdot \vec{k} \Sigma_V$   $\Sigma = \mathbf{T} - \mathbf{T}$ 

## **Determination of** $\Sigma$

• Fit to single-particle potential:



Brockmann/Machleidt 90, Envik et al., Sammarunca et al.,...

$$U(k) = \frac{m^*}{E^*} < k|\Sigma|k\rangle = \frac{m^*}{E^*} \sum_{q \in F} < kq|T(q,k)|kq\rangle$$
$$= \frac{m^*\Sigma_S}{\sqrt{\mathbf{k}^2 + (M + \Sigma_S)^2}} - \Sigma_0$$

• Projection on covariant amplitudes:  $\Sigma_{S,0,V}(k_F,k)$ 

Horowitz/Serot 87, Malfliet et. al., DeJong/Lenske, Tübingen

Include negative energy states

Weigel et al., DeJong/Lenske

# **Isospin dependent self-energy**

$$\Sigma_{n} = -i \int_{F_{n}} (Tr[G_{n}T_{nn}] - G_{n}T_{nn}) - i \int_{F_{p}} (Tr[G_{p}T_{np}] - G_{p}T_{np})$$
  
$$\Sigma_{p} = -i \int_{F_{p}} (Tr[G_{p}T_{pp}] - G_{p}T_{pp}) - i \int_{F_{n}} (Tr[G_{n}T_{np}] - G_{n}T_{np})$$

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- projection method on covariant amplitudes
- averaged mass in n/p channel
- relativistic n/p Pauli operator

van Dalen, C.F., Faessler, NPA 744 (2004) 227.

# **BHF versus DBHF**

#### BHF: 3-body forces necessary (Zuo et al., NPA 706 (2002) 418)



#### All microscopic EOS are soft !

## **Role of correlations**



Lect. Notes Phys. 641 (2004)119 [nucl-th/0309003]

# **Nuclear/neutron matter EOS**



#### neutron matter better controlled.

#### Symmetry energy

$$\begin{split} E(\rho,\beta) &= E(\rho) + E_{\rm sym}(\rho)\beta^2 + \mathcal{O}(\beta^4) + \cdots \\ E_{\rm sym}(\rho) &= \frac{1}{2} \frac{\partial^2 E(\rho,\beta)}{\partial \beta^2} |_{\beta=0} = a_4 + \frac{p_0}{\rho_0^2} (\rho - \rho_0) + \cdots \\ \text{asymmetry parameter}: \ \beta &= Y_n - Y_p \end{split}$$



#### **Isovector potential**

 $U_{\rm iso}(n_B, \mathbf{k}) = (U_n - U_p)/2\beta$ 



BHF: Zuo et al. PRC 72 (2005) 014005 RIA: Chen, Ko, Li, nucl-th/0509009, amplitudes from: McNeil, Shepard, Wallace.

#### **Effective nucleon mass**

Comparison of different approaches  $\implies$  careful ! Many different defnitions of effective masses are used!

Non-relativistic mass:  $m_{NR}^{*} = \left[ M + \frac{1}{k} \frac{d}{dk} U_{s.p.} \right]^{-1}$   $= |\mathbf{k}| [dE/d|\mathbf{k}|]^{-1}$ 

Dirac mass:

$$m_D^* = M + \Sigma_S$$

Relativistic:  $U_{s.p.} \simeq \frac{m_D^*}{E^*} \Sigma_S + \Sigma_0$ 

parameterizes non-locality in space (k-mass) and time (e-mass)

Mahaux et al., Müther, Frick,...

#### **Symmetric nuclear matter**



DBHF/BHF: larger  $m_{NR}^* \implies$  less repulsive

## **Neutron-proton mass splitting**

• BHF:

 $m_{NR,n}^* > m_{NR,p}^*$ 

• RMF:

$$m_{D,n}^* < m_{D,p}^*; \ m_{NR,n}^* < m_{NR,p}^* \quad (\rho + \delta)$$

Baran, Di Toro et al., Phys. Rep. 410 ('05) 335

• DBHF with  $\Sigma$  extracted by fit method:

$$m_{D,n}^* > m_{D,p}^*$$

Alonso & Sammarunca, PRC 67 ('03) 054301

• DBHF with projection method:  $m_{D,r}^*$ 

$$m_{D,n}^* < m_{D,p}^*$$

de Jong & Lenske, PRC 58 ('98) 890, van Dalen, C.F., Faessler, NPA 744 ('04) 227

• non-rel. mass in DBHF:

$$m^*_{NR,n} > m^*_{NR,p}$$

van Dalen, C.F, Faessler, PRL 95 (2005) 022302

## **Dirac mass in asym. matter**



# **Neutron-proton mass splitting**



van Dalen, C.F., Faessler, PRL 95 (2005) 022302

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- Outlook: large scalar/vector fields already generated at tree level => nucl-th/0509049