

Pairing properties of nucleonic matter with dressed nucleons

nucl-th/0508035, PRC in press

Herbert Müther
Universität Tübingen

&

Willem Dickhoff

Washington University in St. Louis

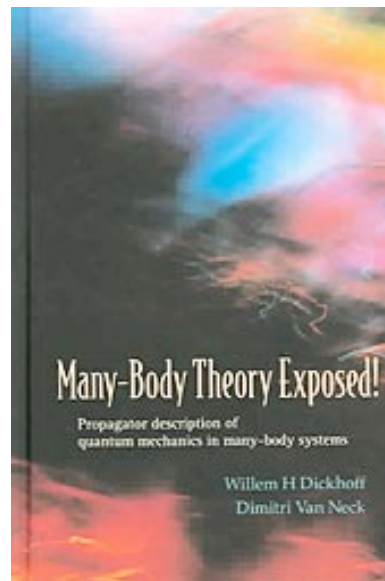
- Pairing issues in matter
- Short-range correlations and data!
- Consequences of SRC
- Self-consistent Green's functions,
dressed nucleons due to SRC
- BCS: a reminder of standard *pairing*
- SCGF plus BCS
- Results and Outlook

Description of the nuclear many-body problem

Ingredients: Nucleons interacting by "realistic interactions"
Nonrelativistic many-body problem

Method: Green's functions (Propagators)
⇒ amplitudes instead of wave functions
keep track of all nucleons, including the high-momentum ones

Book:



Review: Prog. Part. Nucl. Phys. **52**, 377 (2004)

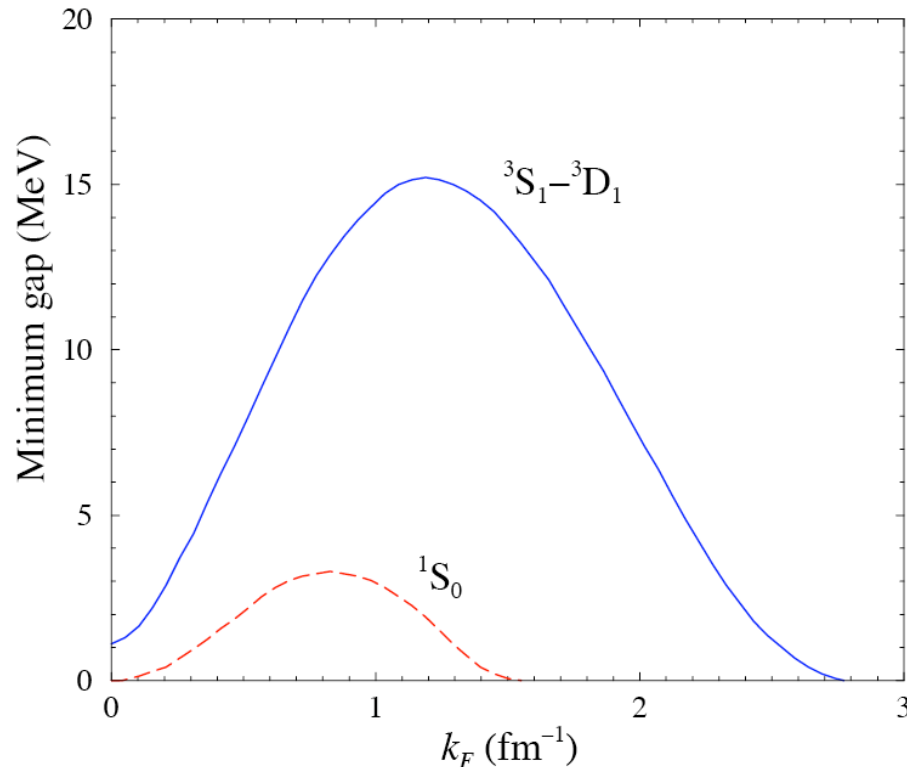
Lecture notes: <http://www.nsl.msui.edu/~brown/theory-group/lecture-notes.html>

Some pairing issues in infinite matter

- Gap size in nuclear matter & neutron matter
- Density & temperature range of superfluidity
- Resolution of 3S_1 - 3D_1 puzzle
- Influence of short-range correlations (SRC)
- Influence of polarization contributions
- Relation of infinite matter results & finite nuclei

Review: e.g. Dean & Hjorth-Jensen, RMP75, 607 (2003)

Puzzle related to gap size in 3S_1 - 3D_1 channel



Mean-field particles

Early nineties: BCS gaps ~ 10 MeV

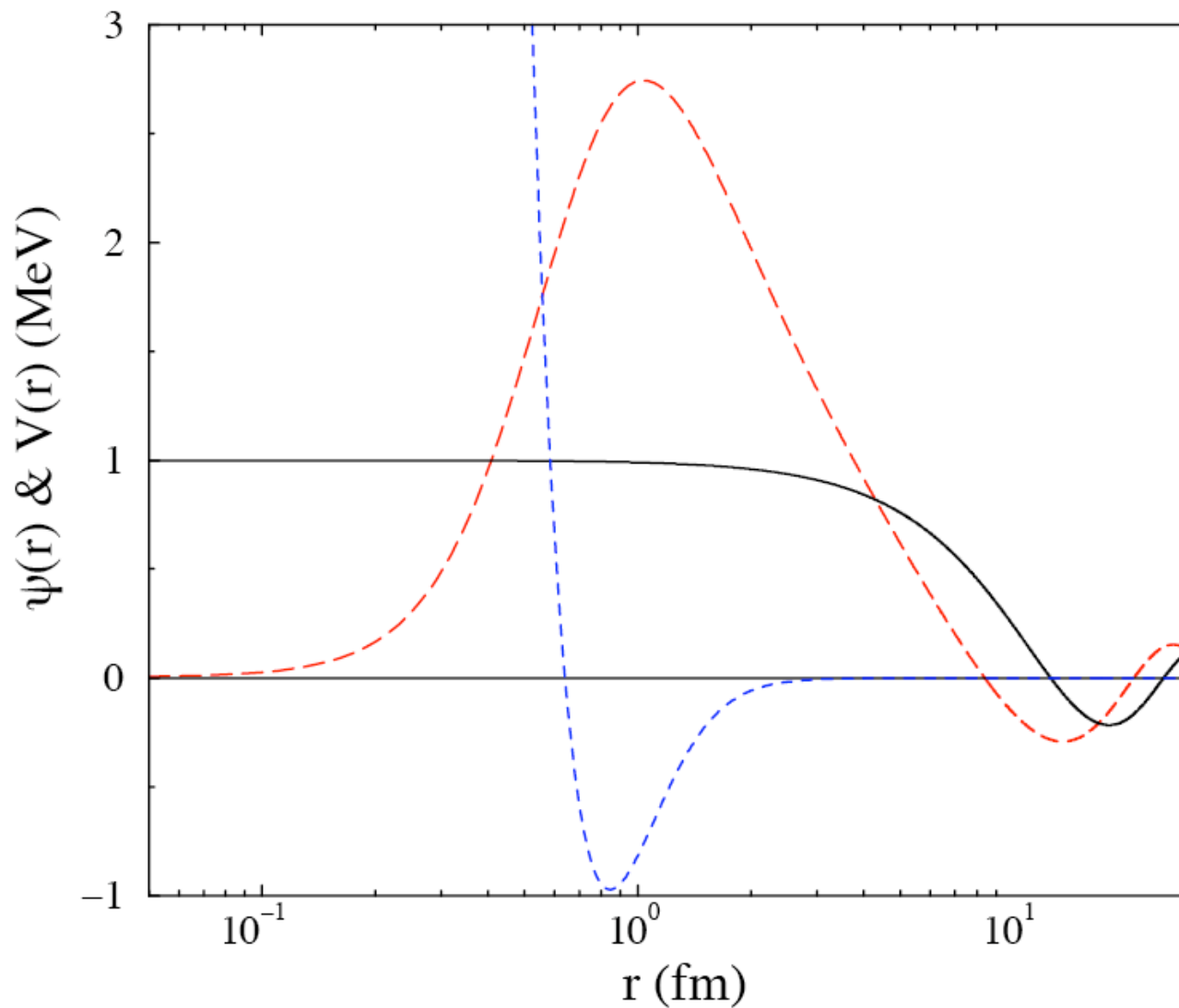
Alm et al. Z.Phys.A337,355 (1990)

Vonderfecht et al. PLB253,1 (1991)

Baldo et al. PLB283, 8 (1992)

Dressing nucleons is expected to reduce pairing strength as suggested by in-medium scattering

Relative wave function and potential



Reid 1S_0
 $k_0 = 0.25 \text{ fm}^{-1}$

$V(r)/100$

$\sin(k_0 r)/k_0 r$

correlated wf

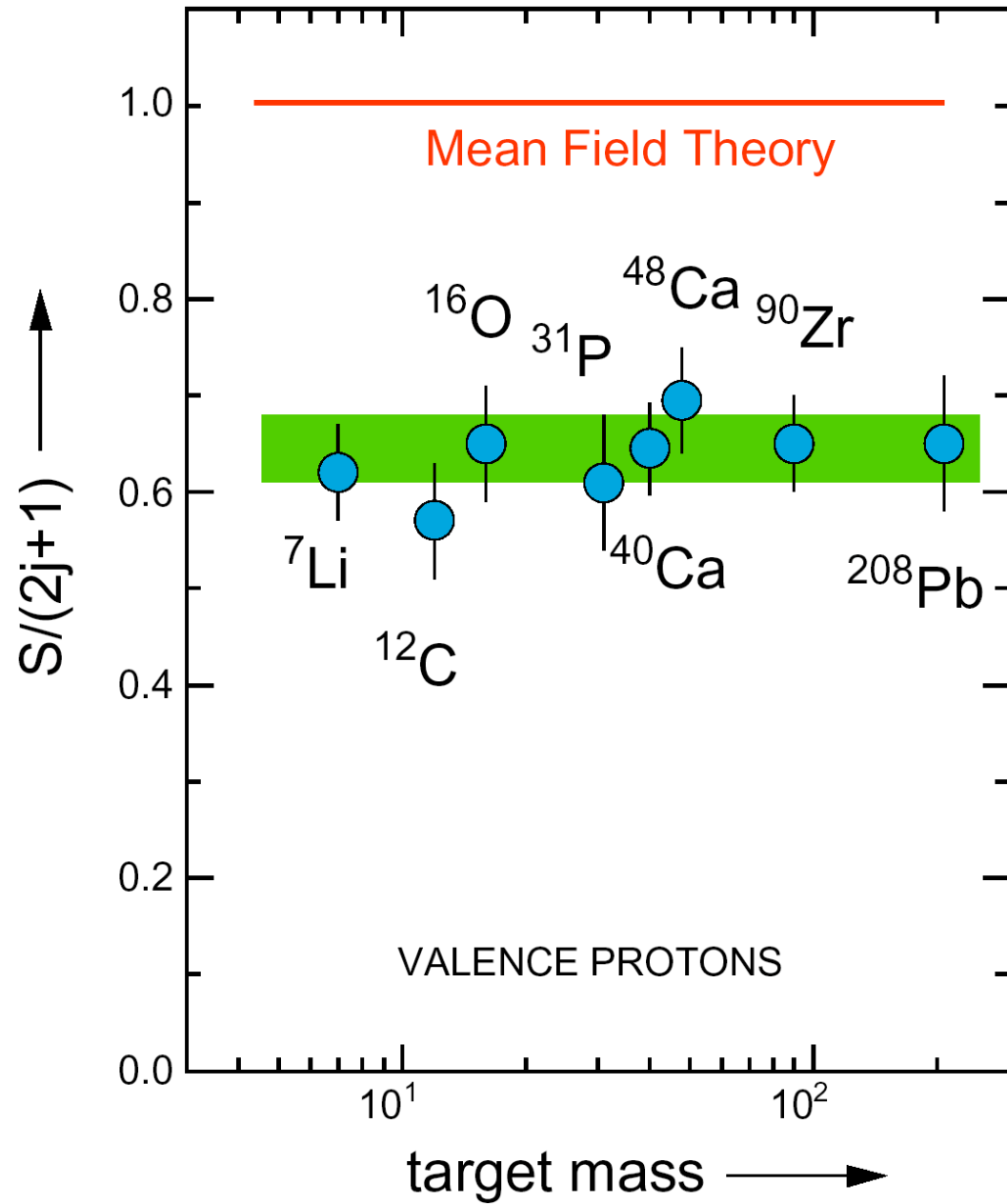
Two effects associated with short-range correlations

- Depletion of the Fermi sea
- Admixture of high-momentum components

Recent data confirm both aspects (predicted by nuclear matter results)

Removal probability for
valence protons
from
NIKHEF data
Lapikás,
NPA553,297c(1993)

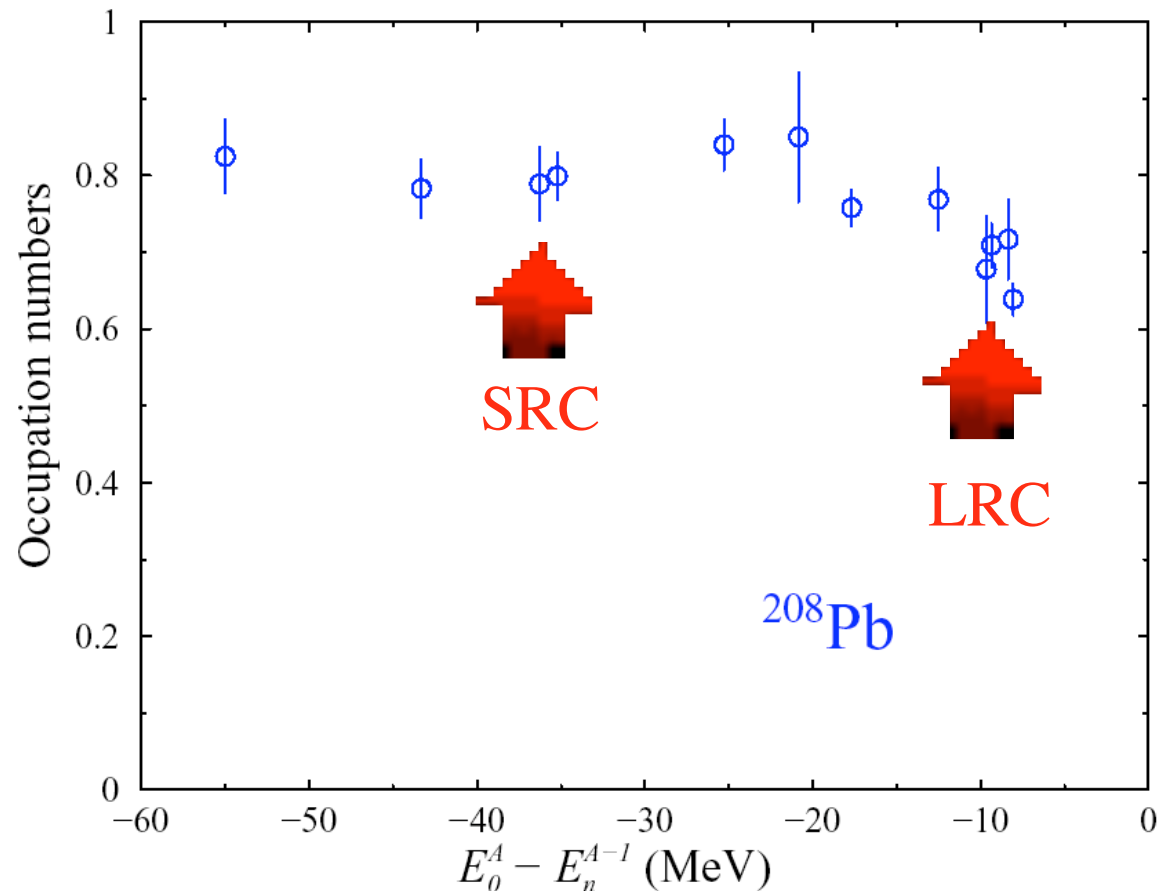
Note:
We have seen mostly
data for removal of
valence protons



M. van Batenburg & L. Lapikás from $^{208}\text{Pb} (e, e' p) ^{207}\text{Tl}$

NIKHEF in preparation

Occupation of deeply-bound proton levels from EXPERIMENT



Up to 100 MeV
missing energy
and
270 MeV/c
missing momentum

Covers the whole
mean-field domain
for the FIRST time!!

Confirms predictions
for depletion

What are the last protons doing? Answer is coming!

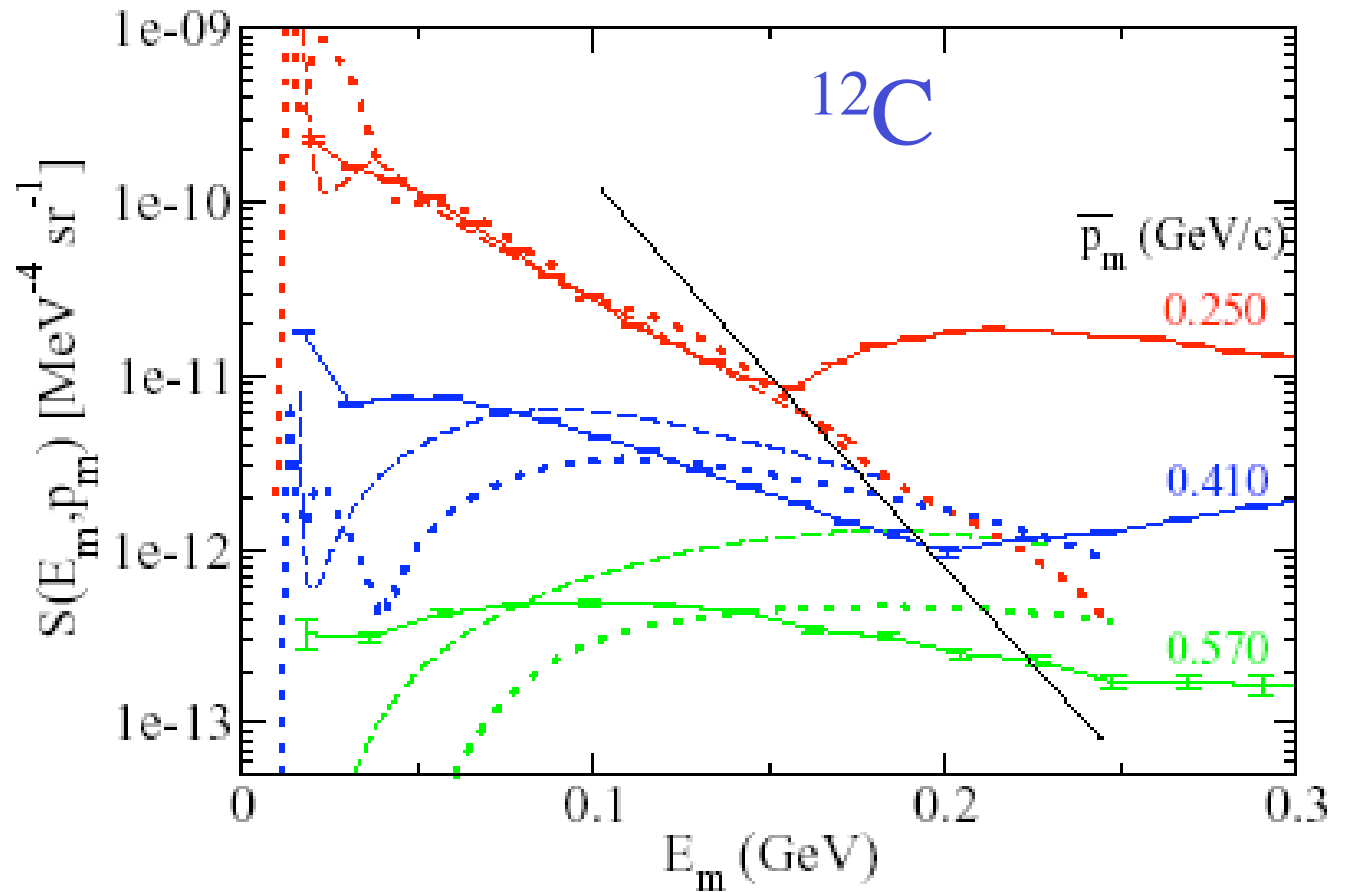
• Jlab E97-006

PRL 93,
182501(2004)

Rohe et al.

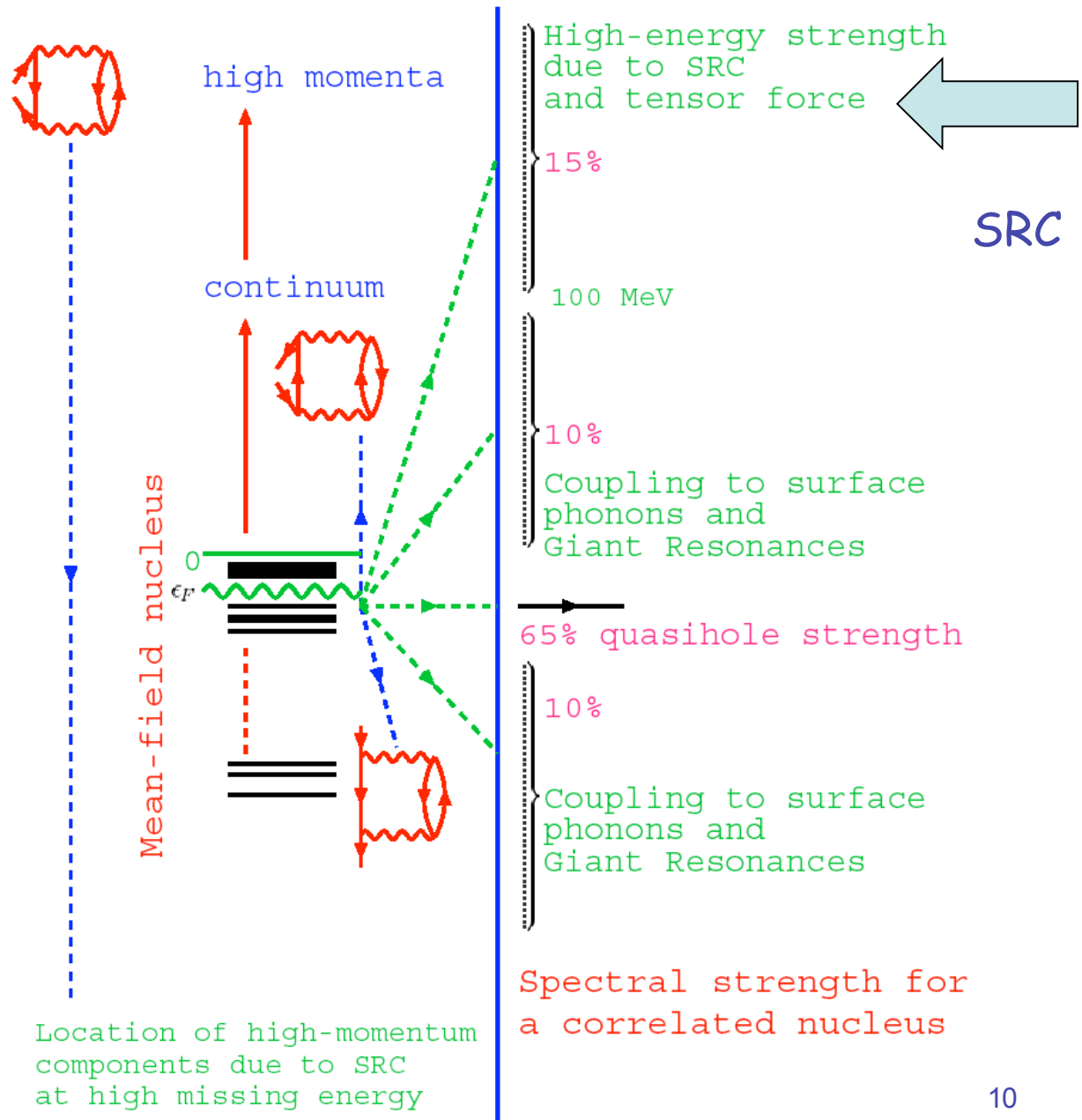
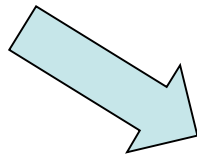
• Location of
high-momentum
components

• Integrated
strength in
agreement with
theoretical
predictions!



Location of
single-particle
strength in nuclei

SRC



We now essentially know what all the protons are doing in the ground state of a “closed-shell” nucleus !!!

- Unique for a **correlated** many-body system
- Information available for electrons in atoms (Hartree-Fock)
- **Not** for electrons in solids
- **Not** for atoms in quantum liquids
- **Not** for quarks in nucleons

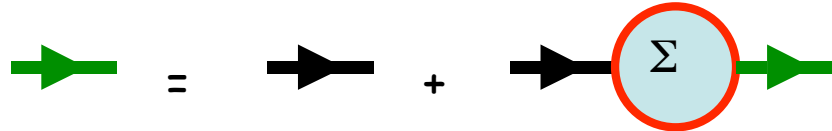
⇒ **Study nucleus
for its intrinsic interest
as a quantum many-body problem!**

Green's function and Γ -matrix approach (ladders)

Single-particle Green's function

$$G(k, t_1, t_2) = -i \langle T c_k(t_1) c_k^\dagger(t_2) \rangle$$

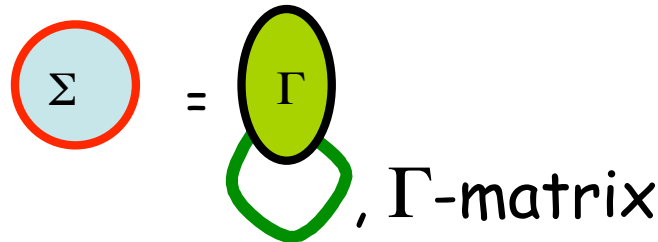
Dyson equation:



$$G(k, \omega) = G^{(0)}(k, \omega) + G^{(0)}(k, \omega) \Sigma(k, \omega) G(k, \omega)$$

$$G(k, \omega) = \frac{1}{\omega - k^2 / 2m - \Sigma(k, \omega)} \Rightarrow S(k, \omega) = -2 \text{Im} G(k, \omega)$$

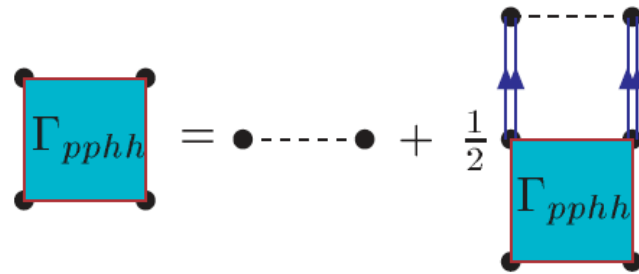
Self-energy



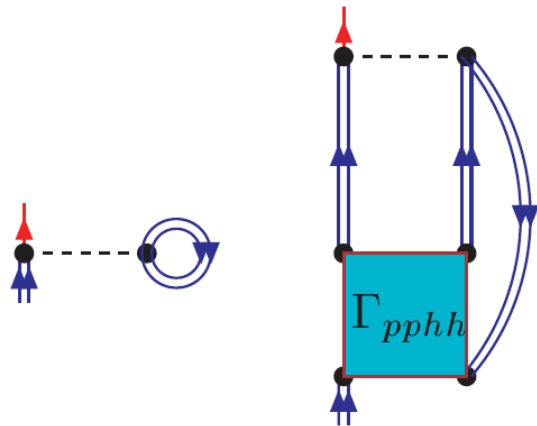
Γ -matrix

- Pairing instability possible
- Finite temperature calculation can avoid this

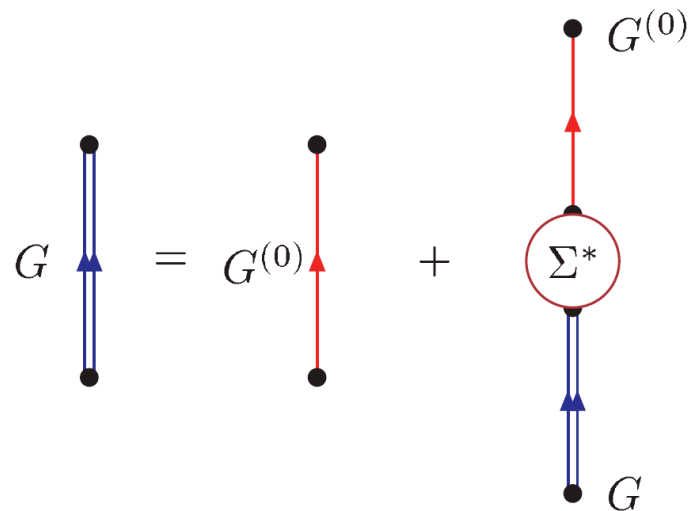
Self-consistency
& SRC



Interaction



Self-energy

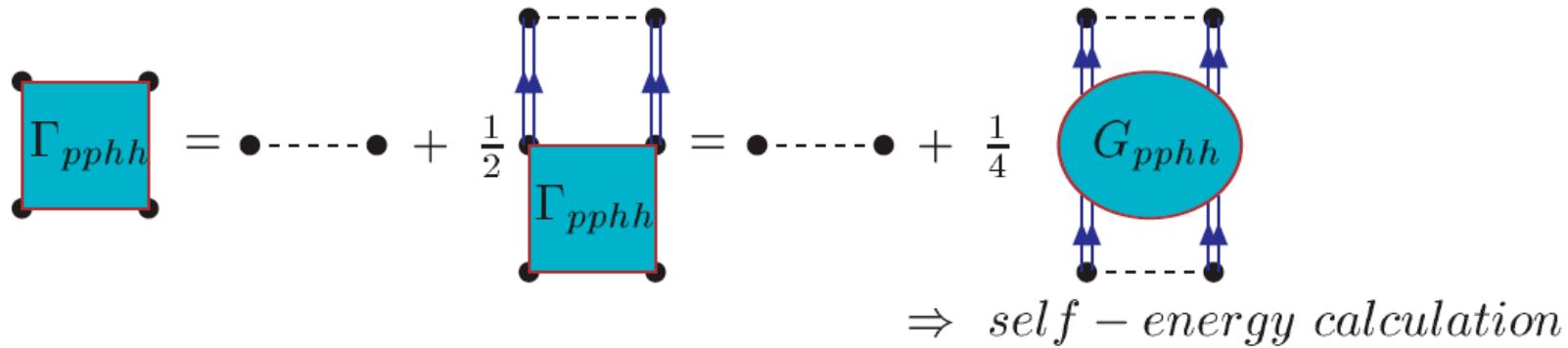


Dyson equation

Pairing of dressed nucleons 13

Ladder diagrams in the medium (options)

Ladders in the medium



$$\langle k\ell | \Gamma_{pphh}^{JST}(K, E) | k' \ell' \rangle = \langle k\ell | V^{JST} | k' \ell' \rangle + \frac{1}{2} \sum_{\ell''} \int_0^{\infty} \frac{dq}{(2\pi)^3} q^2 \langle k\ell | V^{JST} | q\ell'' \rangle G_{pphh}^f(q; K, E) \langle q\ell'' | \Gamma_{pphh}^{JST}(K, E) | k' \ell' \rangle$$

G_{pphh}^f has different form depending on the level of sophistication

Nuclear matter:

$$G_{BG}^f(k_1, k_2; E) = \frac{\theta(k_1 - k_F) \theta(k_2 - k_F)}{E - \varepsilon(k_1) - \varepsilon(k_2) + i\eta}$$

Bethe-Goldstone

$$G_{GF}^f(k_1, k_2; E) = \frac{\theta(k_1 - k_F) \theta(k_2 - k_F)}{E - \varepsilon(k_1) - \varepsilon(k_2) + i\eta} - \frac{\theta(k_F - k_1) \theta(k_F - k_2)}{E - \varepsilon(k_1) - \varepsilon(k_2) - i\eta}$$

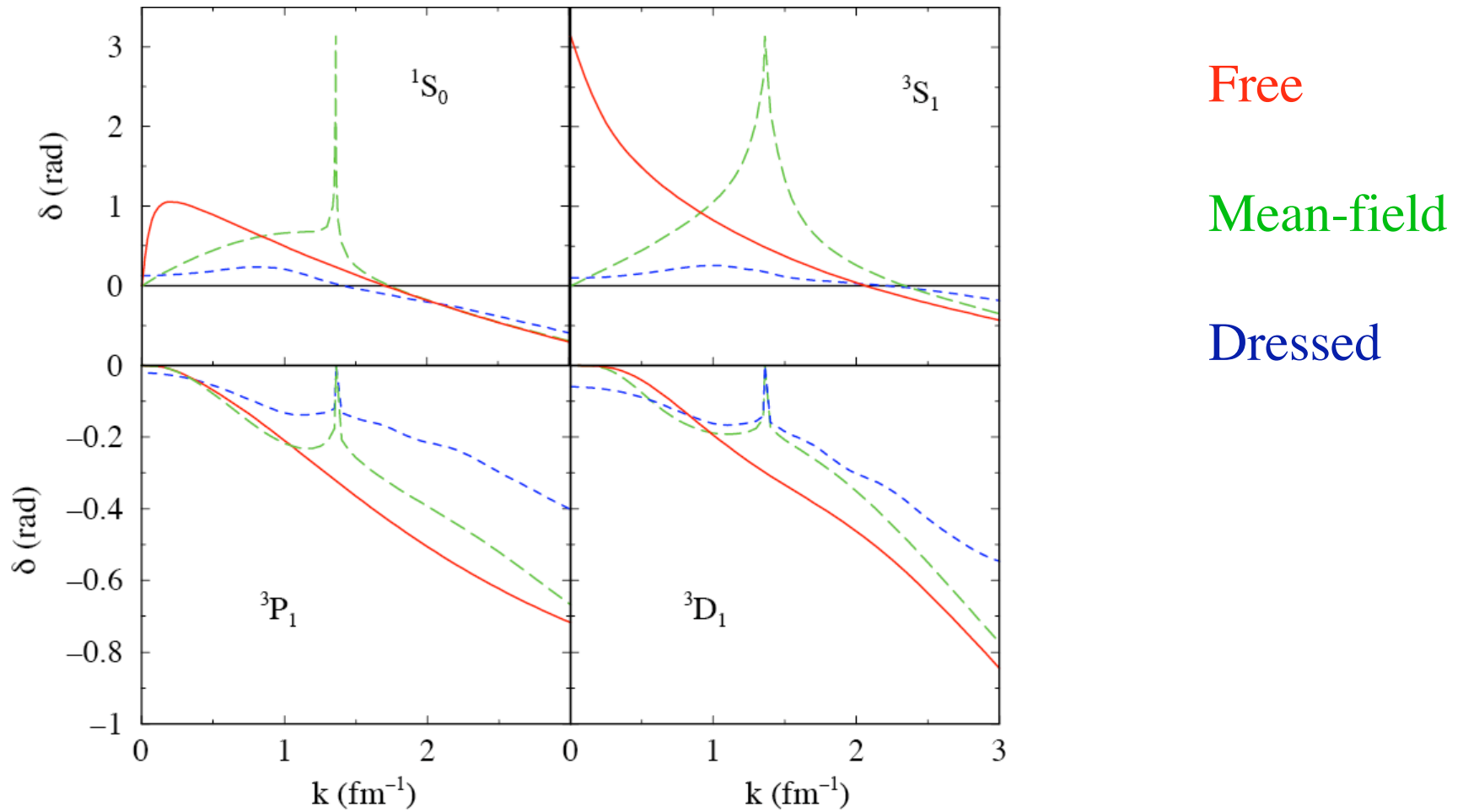
Galitskii-Feynman

$$G_{pphh}^f(k_1, k_2; E) = \int_{\varepsilon_F}^{\infty} dE_1 \int_{\varepsilon_F}^{\infty} dE_2 \frac{S_p(k_1; E_1) S_p(k_2; E_2)}{E - E_1 - E_2 + i\eta} - \int_{-\infty}^{\varepsilon_F} dE_1 \int_{-\infty}^{\varepsilon_F} dE_2 \frac{S_h(k_1; E_1) S_h(k_2; E_2)}{E - E_1 - E_2 - i\eta}$$

SC

Pairing of dressed nucleons 14

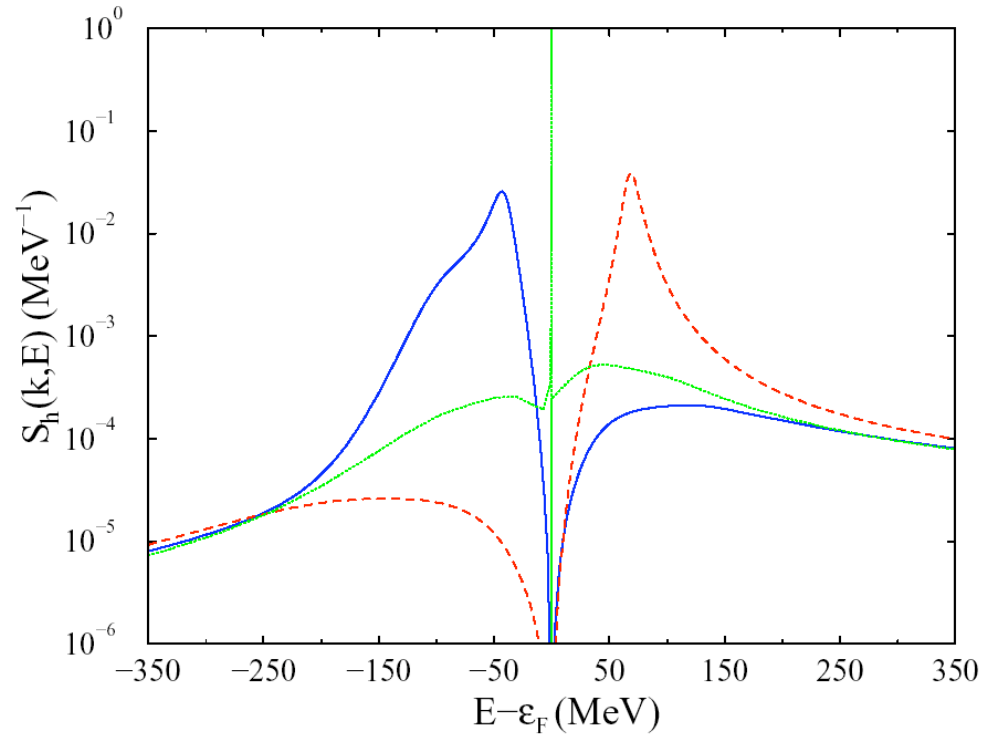
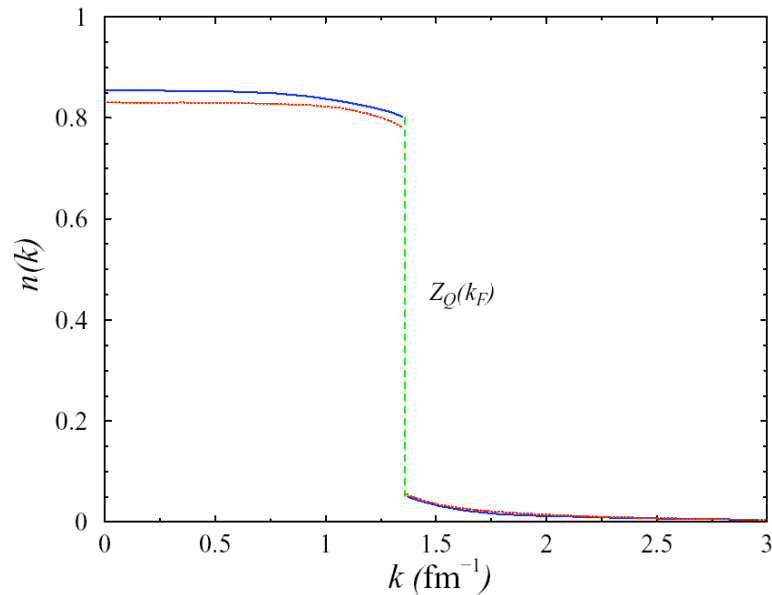
Phase shifts for dressed nucleons



PRC60, 064319 (1999) also PRC58, 2807 (1998)

Results from Nuclear Matter 2nd generation (2000)

- Spectral functions for $k = 0, 1.36, \text{ \& } 2.1 \text{ fm}^{-1}$
- Common tails on both sides of ε_F



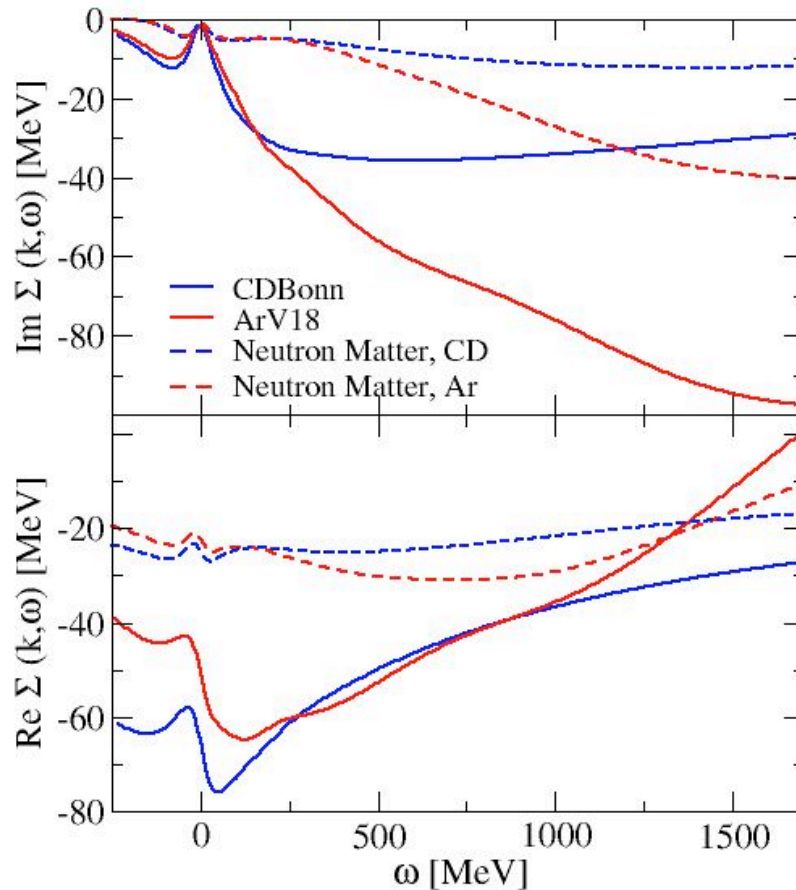
Momentum distribution: only minor changes

Occupation in nuclei: **Depleted similarly!**

Thesis Libby Roth Stoddard (2000)

Self-energy

$$G(k, \omega) = \frac{1}{\omega - k^2/2m - \Sigma(k, \omega)} \Rightarrow S(k, \omega) = -2\text{Im}G(k, \omega)$$



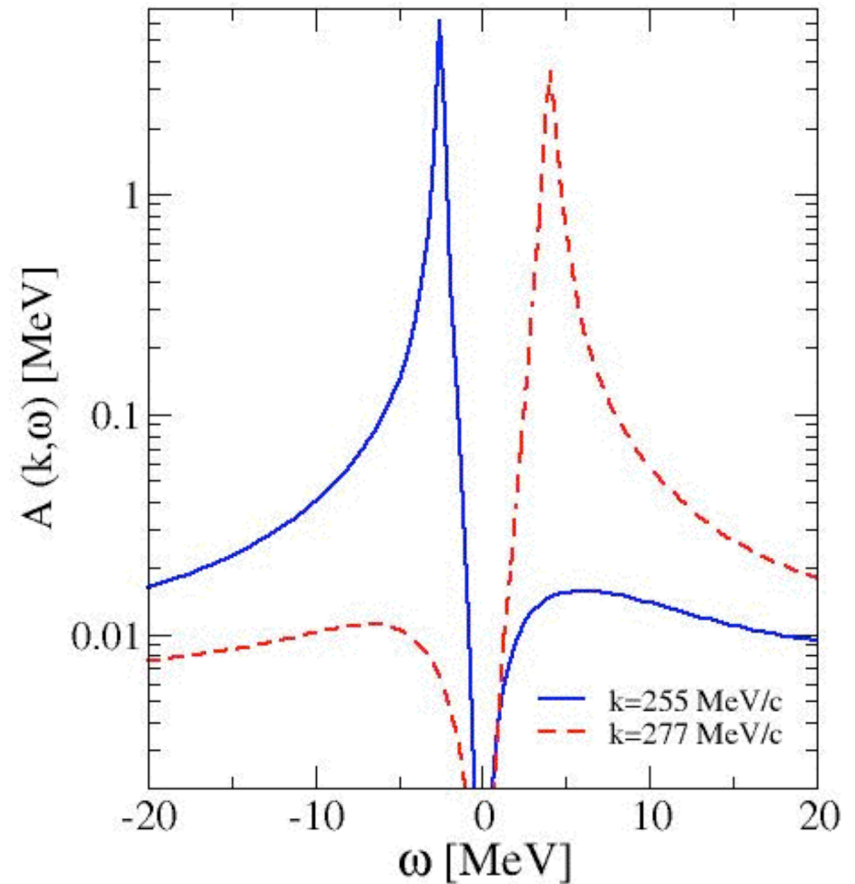
Real and imaginary part of the retarded self-energy

- $k_F = 1.35 \text{ fm}^{-1}$
- $T = 5 \text{ MeV}$
- $k = 1.14 \text{ fm}^{-1}$

Note differences due to NN interaction

Spectral functions

- Strength above and below the Fermi energy as in BCS
- but broad distribution in energy
- BCS not just a cartoon of SCGF but both features must be considered in a consistent way
- CDBonn interaction at $T=0$



BCS: a reminder

NN correlations on top of Hartree-Fock: ε_k, c_k^+

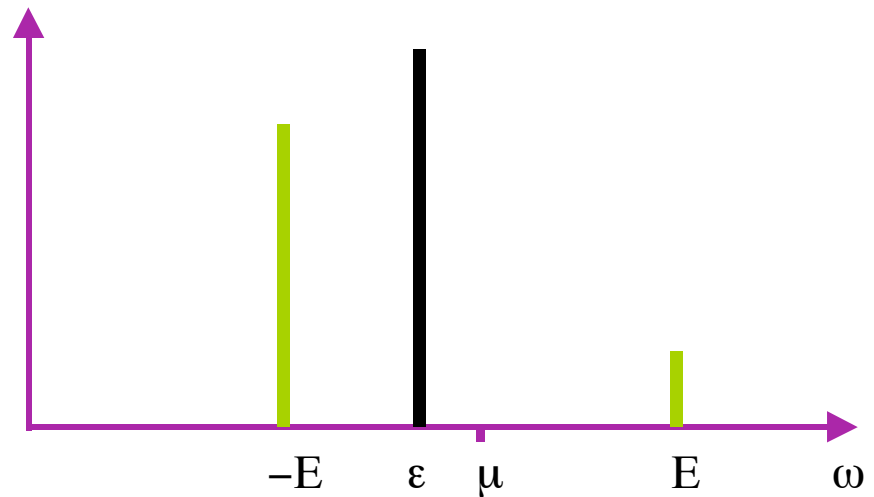
Bogoliubov transformation $a_k^+ = u_k c_k^+ + v_k c_{\bar{k}}$

with
$$\begin{matrix} u_k^2 \\ v_k^2 \end{matrix} = \frac{1}{2} \left[1 \pm \frac{\varepsilon_k - \mu}{\sqrt{(\varepsilon_k - \mu)^2 + \Delta(k)^2}} \right], \quad E(k) = \sqrt{(\varepsilon_k - \mu)^2 + \Delta(k)^2}$$

Gap equation

$$\Delta(k) = \int k'^2 dk' \langle k, \bar{k} | V | k', \bar{k}' \rangle \frac{\Delta(k')}{-2E(k)}$$

Spectral function $S(k, \omega)$



Solution of the gap equation

$$\Delta(k) = \sum_{k'} \langle k, \bar{k} | V | k', \bar{k}' \rangle \frac{\Delta(k')}{\omega - 2E(k)} \quad \text{with} \quad E(k) = \sqrt{(\varepsilon_k - \mu)^2 + \Delta(k)^2} \quad \text{and} \quad \omega=0$$

Define:
$$\delta(k) = \frac{\Delta(k)}{\omega - 2E(k)}$$

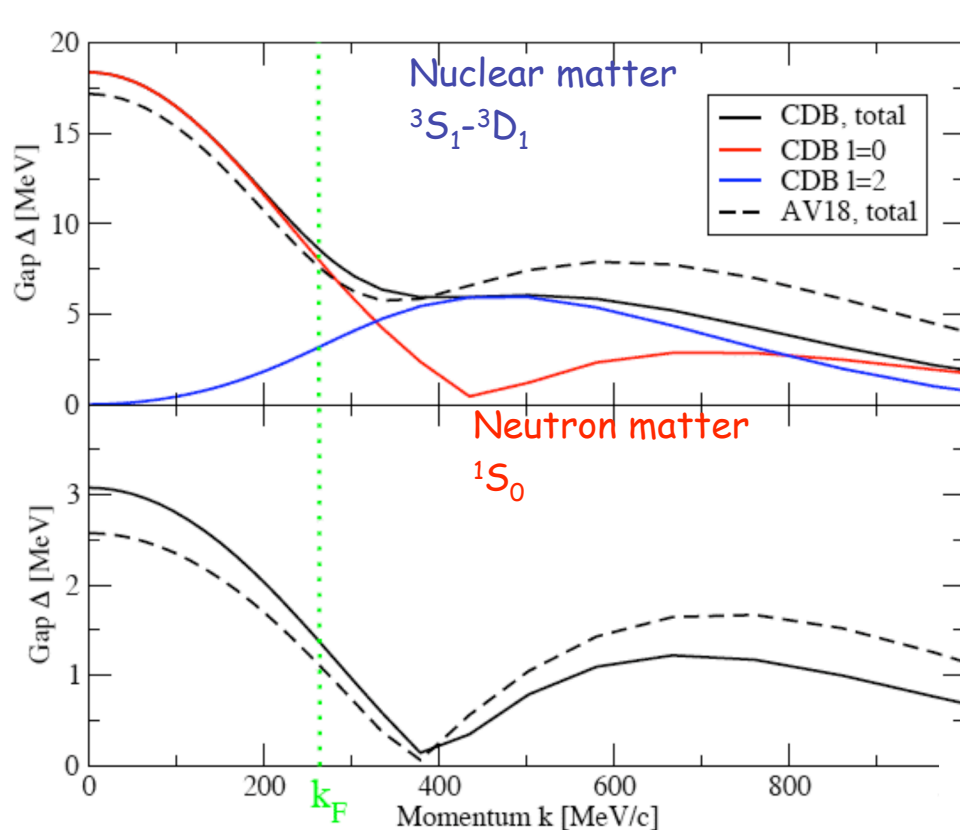
$$\begin{pmatrix} 2E(k) + \langle k | V | k \rangle, & \dots, & \langle k | V | k' \rangle \\ \vdots & \ddots & \vdots \\ \langle k' | V | k \rangle, & \dots, & 2E(k') + \langle k' | V | k' \rangle \end{pmatrix} \begin{pmatrix} \delta(k) \\ \vdots \\ \delta(k') \end{pmatrix} = \omega \begin{pmatrix} \delta(k) \\ \vdots \\ \delta(k') \end{pmatrix}$$

Eigenvalue problem for a pair of nucleons at $\omega=0$

Steps of the calculation:

- Assume $\Delta(k)$ and determine $E(k)$
- Solve eigenvalue equation and evaluate new $\Delta(k)$
 - If lowest eigenvalue $\omega < 0$ enhance $\Delta(k)$ (resp. $\delta(k)$)
 - If lowest eigenvalue $\omega > 0$ reduce $\Delta(k)$
- Repeat until convergence

Gaps from BCS for realistic interactions



- momentum dependence $\Delta(k)$
- different NN interactions
- very similar to pairing gaps in finite nuclei for like particles...!?
- for np pairing no strong empirical evidence...?!?
- consequences for neutron stars:
 - Neutrino propagation
 - Glitches
 - Cooling

$T = 0$
Mean-field particles

Beyond BCS in the framework of SCGF

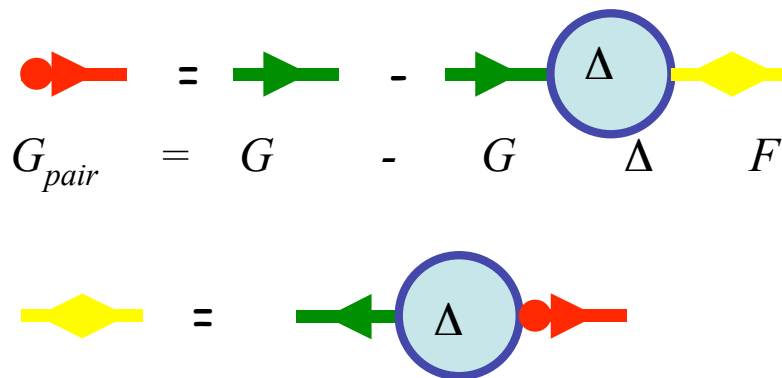
Generalized Green's functions: Extend $G(k, t_1, t_2) = -i \langle T c_k(t_1) c_k^+(t_2) \rangle$

Anomalous propagators $G(k, t_1, t_2) = \begin{pmatrix} -i \langle T c c^+ \rangle & -i \langle T c c \rangle \\ i \langle T c^+ c^+ \rangle & i \langle T c^+ c \rangle \end{pmatrix} = \begin{pmatrix} G & F \\ F^+ & \bar{G} \end{pmatrix}$

Generalized Dyson equation: Gorkov equations

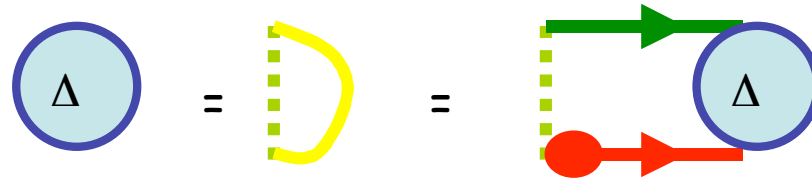
$$\begin{pmatrix} \omega - t_k - \Sigma(k, \omega) & -\Delta(k, \omega) \\ -\Delta^+(k, \omega) & \omega + t_k + \Sigma(k, \omega) \end{pmatrix} \begin{pmatrix} G_{pair}(k, \omega) & F(k, \omega) \\ F^+(k, \omega) & \bar{G}_{pair}(k, \omega) \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Leads to e.g.



G includes all normal self-energy terms

Anomalous self-energy: Δ & generalized Gap equation



$$\Delta(k) = \int k'^2 dk' \langle k|V|k' \rangle \int d\omega \int d\omega' \frac{1 - f(\omega) - f(\omega')}{-\omega - \omega'} S(k', \omega) S_{pair}(k', \omega') \Delta(k')$$

Fermi function $f(\omega) = \frac{1}{e^{\beta\omega} + 1}$

If we replace $S(k, \omega)$ by "HF" approx. and $S_{pair}(k, \omega)$ by BCS:
 \Rightarrow Usual Gap equation

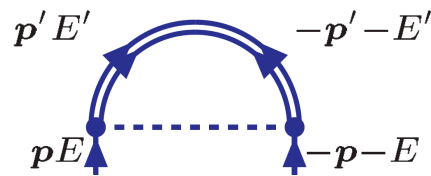
If we take $S_{pair}(k, \omega) = S(k, \omega)$:

\Rightarrow Corresponds to the homogeneous solution of Γ -matrix eq.

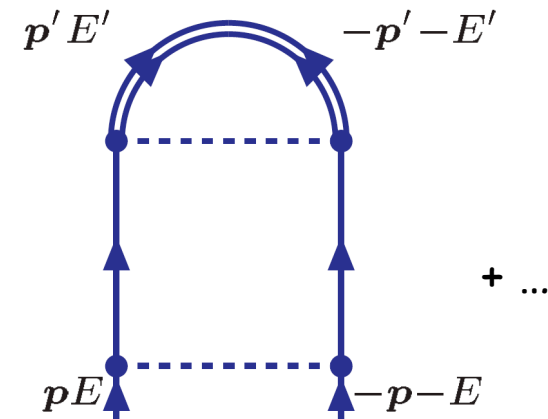
With $S_{pair}(k, \omega)$:

\Rightarrow The above and self-consistency

Consistency of Gap equation (anomalous self-energy) and Ladder diagrams



Iteration of Gorkov equations for anomalous propagator generates



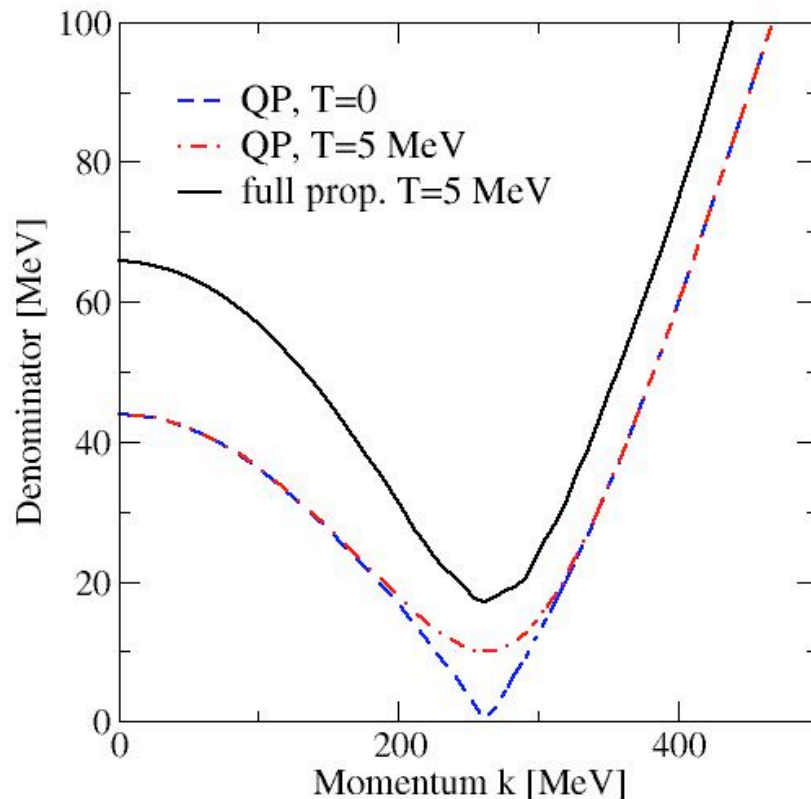
... and all other ladder diagrams at total momentum and energy zero (w.r.t. 2μ) plus anomalous self-energy terms in normal part of propagator

So truly consistent with inclusion of ladder diagrams at other total momenta and energies

Features of generalized gap equation

$$\Delta(k) = \int k'^2 dk' \langle k|V|k' \rangle \int d\omega \int d\omega' \frac{1 - f(\omega) - f(\omega')}{-\omega - \omega'} S(k', \omega) S_{pair}(k', \omega') \Delta(k')$$

$$\frac{1}{-2|\varepsilon_{k'} - \mu|}$$



Dashed:

Spectral strength only at 1 energy

Dashed-dot:

Effect of temperature (5 MeV)

Solid:

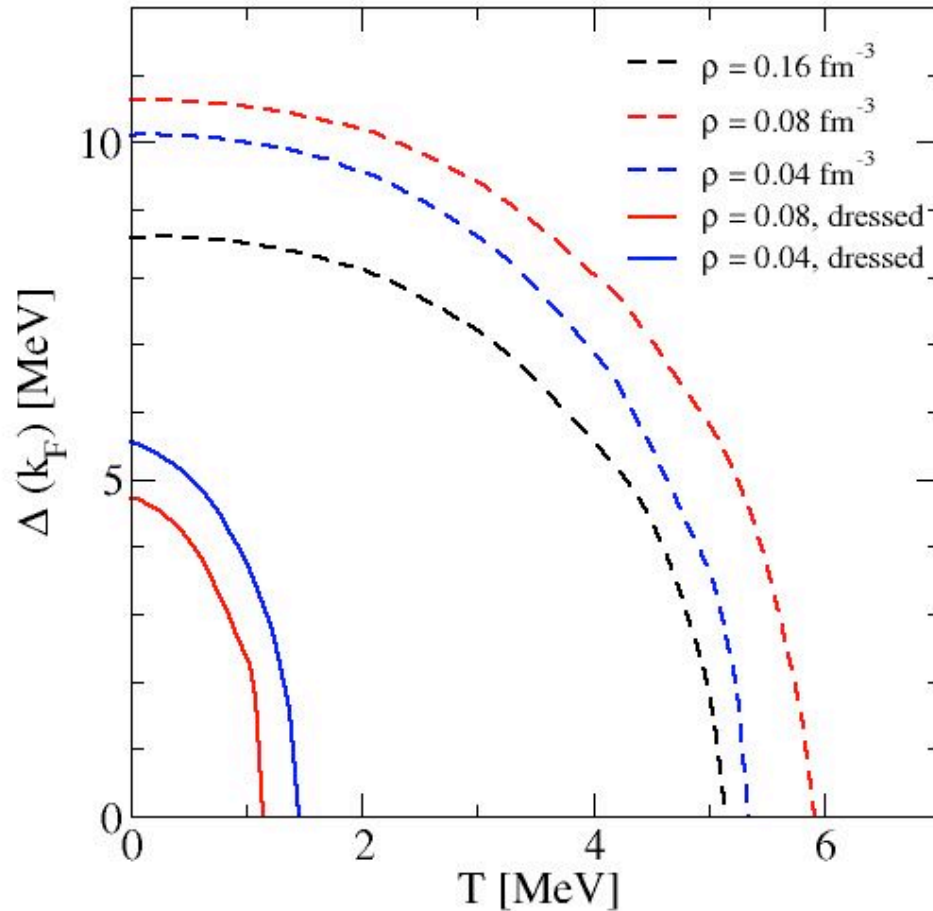
Includes complete strength distribution due to SRC

Related studies by

Baldo, Lombardo, Schuck et al.

use BHF self-energy

Proton-neutron pairing in symmetric nuclear matter

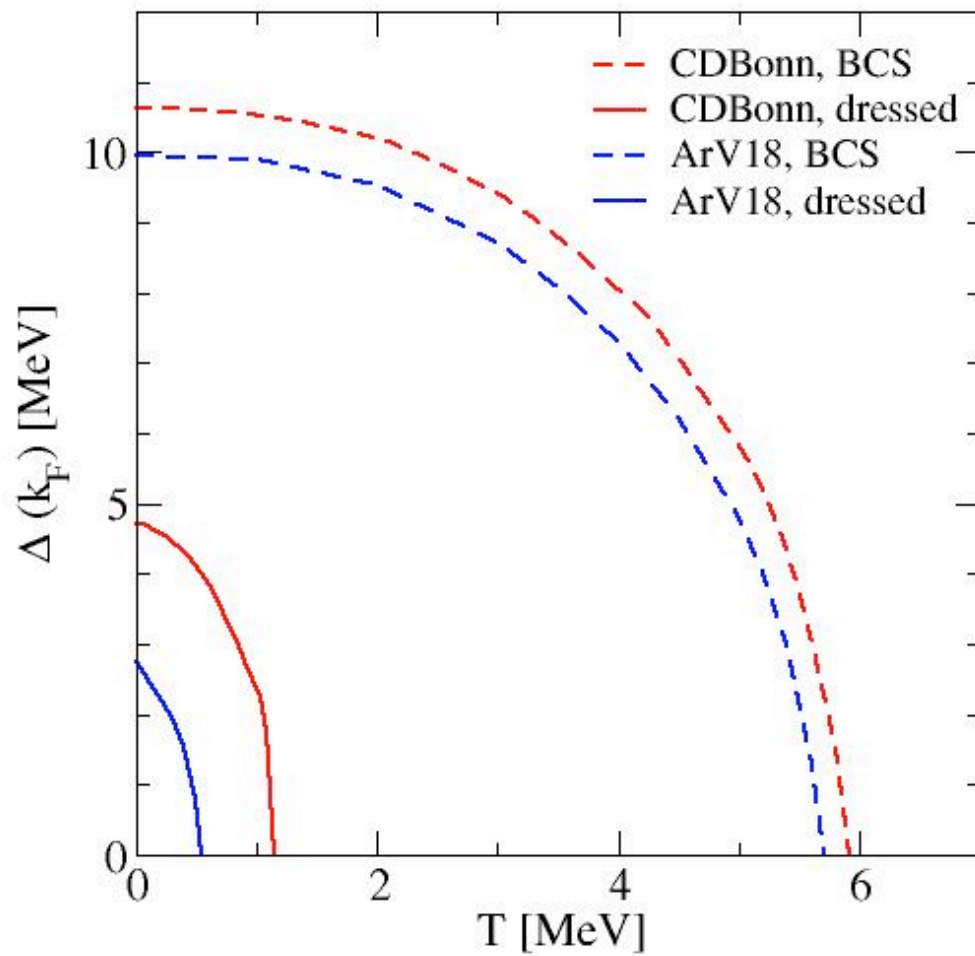


Using CDBonn

Dashed lines:
quasiparticle poles

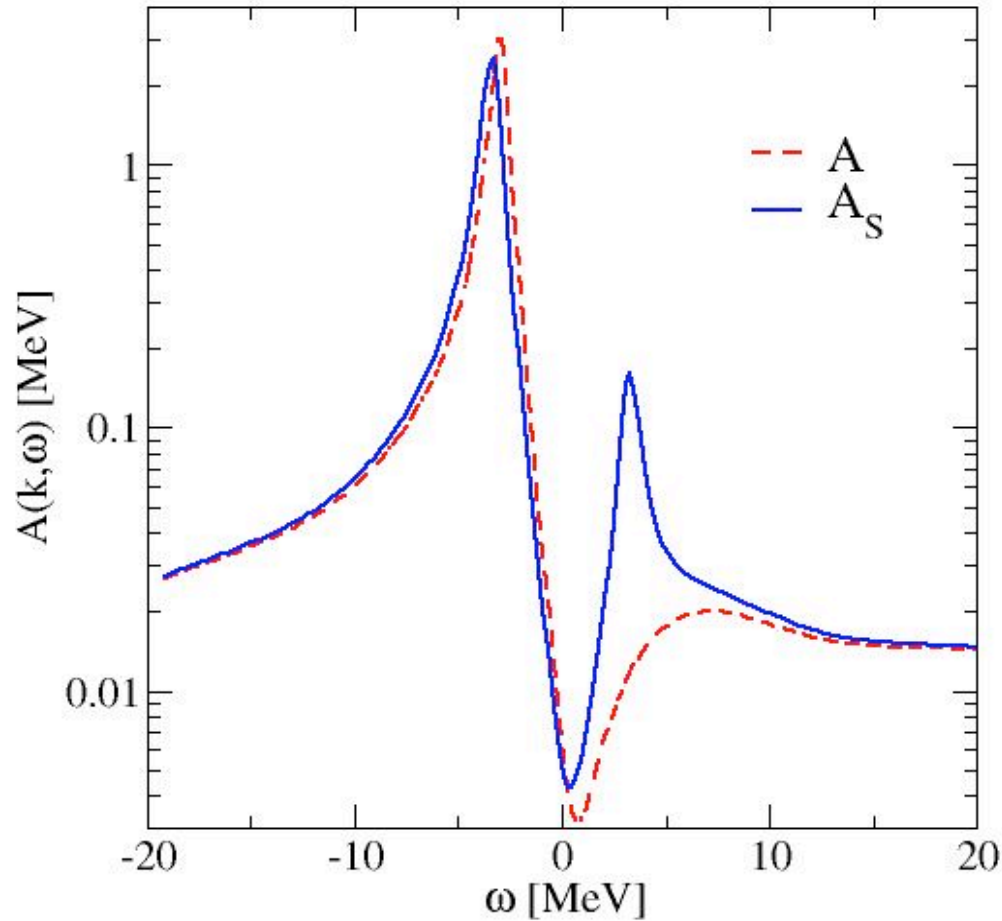
Solid lines:
dressed nucleons

No pairing at saturation
density!



CDBonn yields stronger pairing than ArV18

Pairing and spectral functions



Spectral functions

$S(k, \omega)$ dashed = $A(k, \omega)$

$S_{\text{pair}}(k, \omega)$ solid = $A_S(k, \omega)$

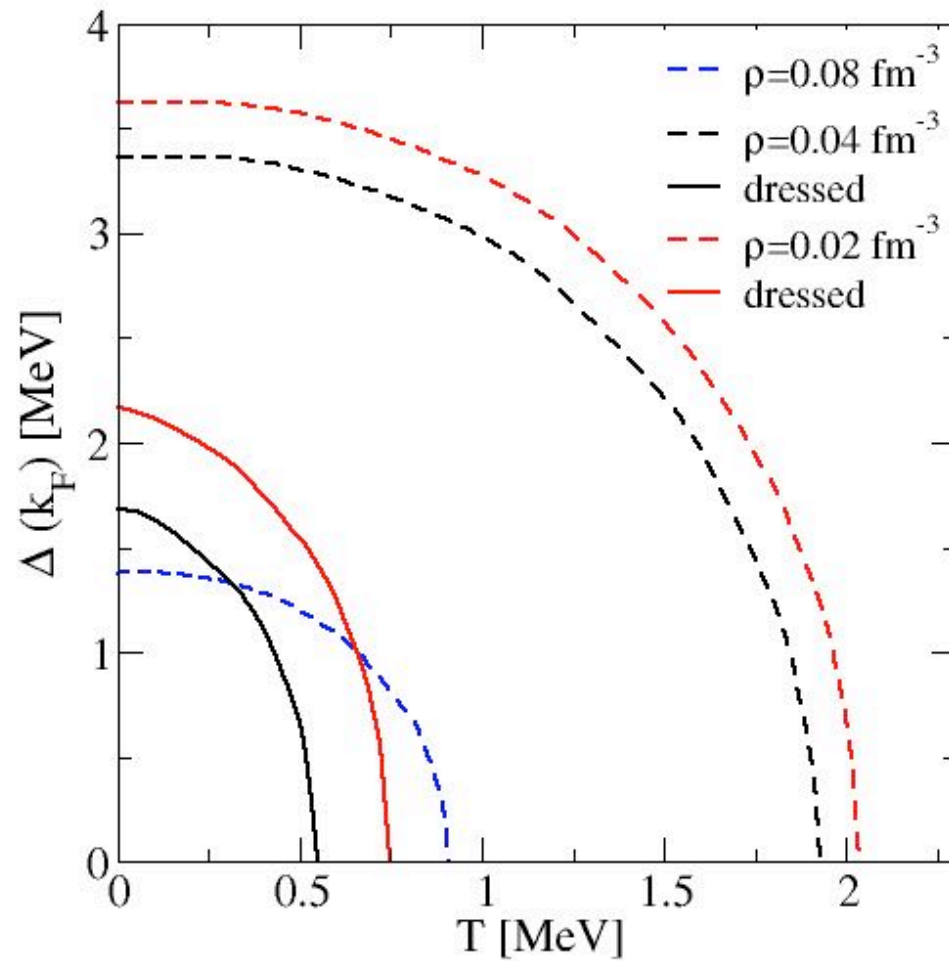
$\rho = 0.08 \text{ fm}^{-3}$

$T = 0.5 \text{ MeV}$

$k = 193 \text{ MeV}/c$ $0.9 k_F$

Expected effect

Pairing in neutron matter



*Dressing effects weaker,
but non-negligible
CDBonn*

Conclusions

- **Consistent** treatment of short-range and pairing correlations
- Short-range correlations reduce pairing **significantly**
- **No** proton-neutron pairing at normal density!
- Correlation effects are **weaker** in neutron matter but nonnegligible

Outlook

- Pairing and screening of interaction (polarization effects)...
- Complex and energy-dependent Δ
- Effects of pairing plus correlations on physics of neutrino propagation ...