

Unbound exotic nuclei studied by projectile fragmentation

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Seattle 2005

Plan of the Presentation

1 Reactions to study halo nuclei

- Transfer to the continuum
- Fragmentation reaction

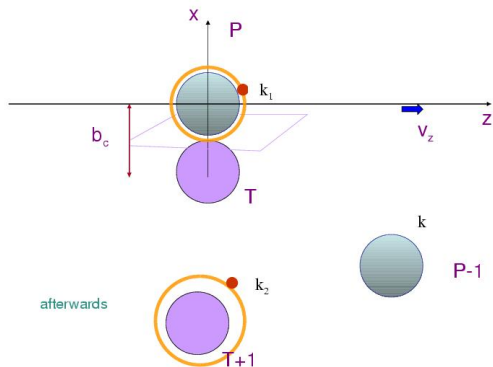
2 Simple time dependent model

- Inelastic excitation
- General wave functions (asymptotic form)
- Simple time dependent model
- Comparison to the transfer to the continuum
- Cross section.
- Determination of the S-matrix.
- Potential correction

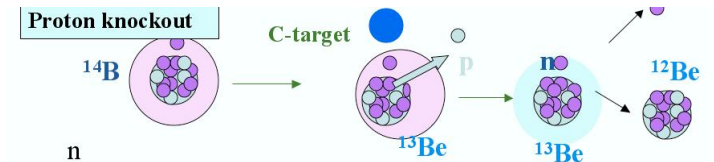
3 Results

- Comparison between time dependent and sudden approximation
- Dependence on the binding energy of the initial state
- Strength of every transition
- Dependence on the scattering length of the final s-state
- Add of a complex part to the potential

Transfer to the continuum reaction



Fragmentation reaction



Transfer to the continuum

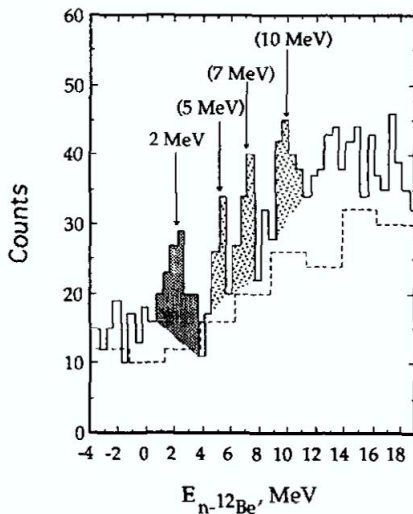


Figure: $^{12}\text{Be}(d,p)$ RIKEN (Korshenninnikov) (1995)

Fragmentation reaction.

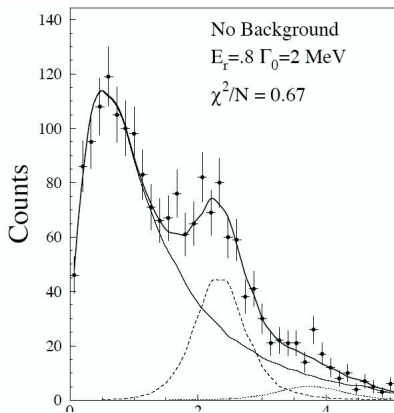


Figure: (a) LPC GANIL (Lecouey, Orr) (2002).

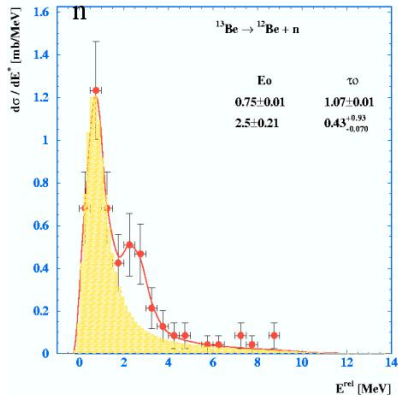


Figure: (b) GSI (U Datta Pramanik) (Surrey conference Jan.2005).



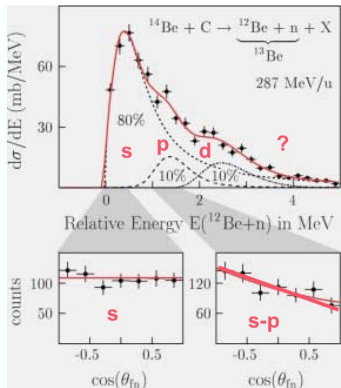


Figure: (a) GSI (Simon Nucl. Phys. A734 (2004) 323-326.

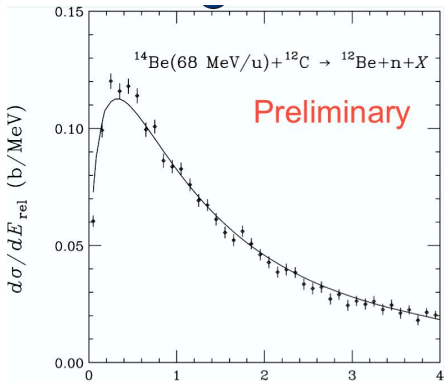


Figure: (b) RIKEN (Nakamura)(ECT* 2004)unpublished.



Inelastic excitation.

Inelastic-like excitations can be described by the first order time dependent perturbation theory amplitude:

$$A_{fi} = \frac{1}{i\hbar} \int_{-\infty}^{\infty} dt \langle \psi_f(t) | V_2(\mathbf{r}) | \psi_i(t) \rangle$$

In order to obtain a simple analytical formula we consider the special case in which $V_2(r) = v_2 \delta(x) \delta(y) \delta(z)$.

$$A_{fi} = \frac{v_2}{i\hbar v} \int_{-\infty}^{\infty} dz \psi_f^*(b_c, 0, z) \psi_i(b_c, 0, z) e^{iqz}$$



General wave functions (asymptotic form).

For the initial state:

$$\psi_i(b_c, 0, z) = -C_i i^l \gamma h_{l_i}^{(1)}(i\gamma r) P_{l_i}(z/r).$$

For the final continuum state:

$$\psi_f(b_c, 0, z) = C_f k \frac{i}{2} (h_{l_f}^{(+)}(kr) - S_{l_f} h_{l_f}^{(-)}(kr)) P_{l_f}(z/r).$$



Simple time dependent model.

$$A_{fi} = \frac{v_2}{i\hbar v} \int_{-\infty}^{\infty} dz \psi_f^*(b_c, 0, z) \psi_i(b_c, 0, z) e^{iqz}$$

$$I(k, q) = I_R + iI_I = |I|e^{i\alpha}$$

$$\bar{S} = S e^{2i\alpha} = e^{2i(\delta+\alpha)}$$

$$|A_{fi}|^2 = C^2 |I|^2 |1 - \bar{S}|^2.$$



Comparison to the transfer to the continuum.

Fragmentation:

$$\frac{dP_{in}}{d\varepsilon_f} = \frac{2}{\pi} \frac{v_2^2}{\hbar^2 v^2} C_i^2 \frac{m}{\hbar^2 k} \sum_{l_f} (2l_f + 1) |1 - \bar{S}_{l_f}|^2 |I_{l_f}|^2.$$

Transfer:

$$\frac{dP_t(b_c)}{d\varepsilon_f} \approx \frac{4\pi}{k^2} \sum_{l_f} (2l_f + 1) |1 - S_{l_f}|^2 B_{l_f, l_i}$$

(G. Blanchon, A. Bonaccorso and N. Vinh Mau, *Nucl. Phys. A739* (2004) 259.) \otimes

Cross section.

$$\frac{d\sigma}{d\varepsilon_f} = C^2 S \int_0^\infty d\mathbf{b}_c \frac{dP_{in}(b_c)}{d\varepsilon_f} P_{ct}(b_c),$$

The core survival probability:

$$P_{ct}(b_c) = |S_{ct}|^2$$



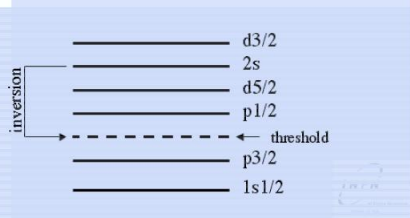
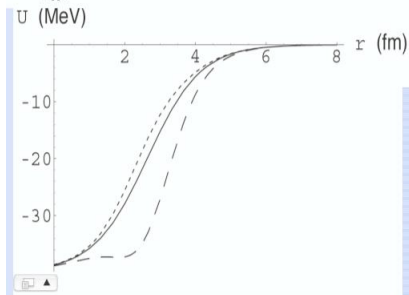
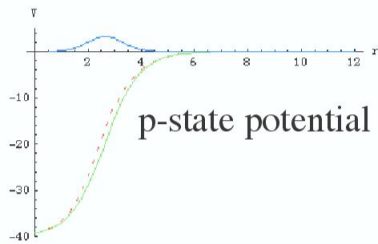
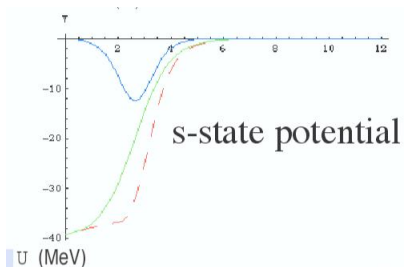
Determination of the S-matrix.

$$\begin{aligned}h &= t + U. \\U(r) &= V_{WS} + \delta V. \\ \delta V(r) &= 16\alpha e^{2(r-R)/a} / (1 + e^{(r-R)/a})^4.\end{aligned}$$

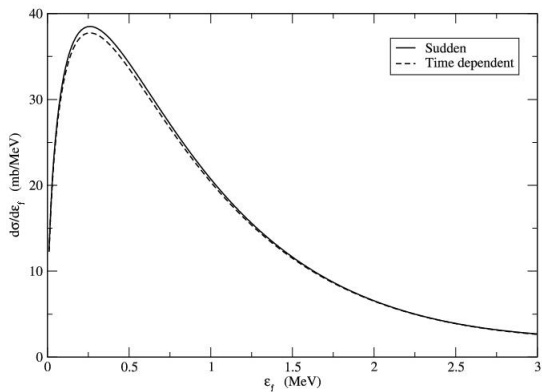
V_{WS} = Woods-Saxon potential plus spin-orbit. δV = correction which originates from particle-vibration couplings.

(*N. Vinh Mau and J. C. Pacheco, Nucl. Phys. A607 (1996) 163.*) \otimes

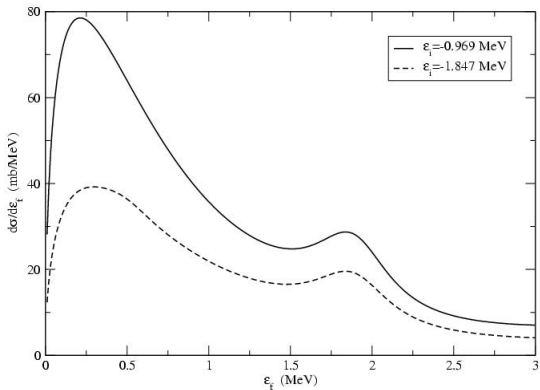
Potential correction.



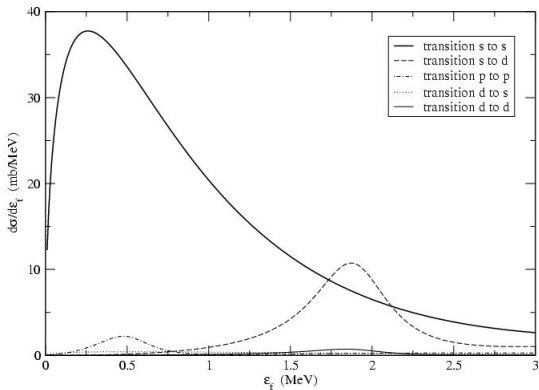
Comparison between time dependent and sudden approximation.



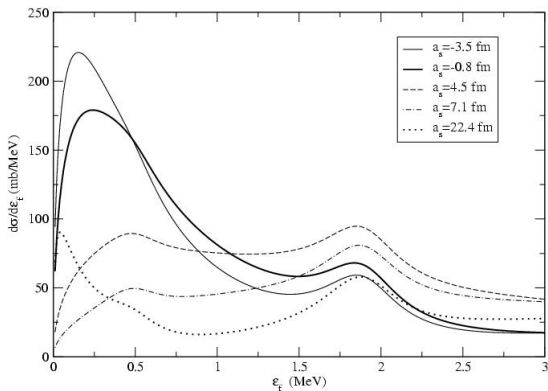
Dependence on the binding energy of the initial state.



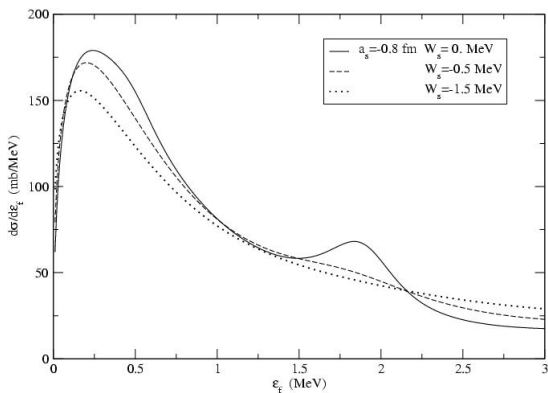
Strength of every transition.



Dependence on the scattering length of the final s-state.



Introduction of an imaginary part to the potential.



Conclusions and outlooks

- ^{13}Be is a signature of the halo state of the neutron.
- Possibility to use s resonances.
- One or two step calculation.

