

# **Quantum calculation of vortices in the inner crust of neutron stars**

P. Avogadro, R.A. Broglia, E. Vigezzi

**Milano University and INFN**

F. Barranco

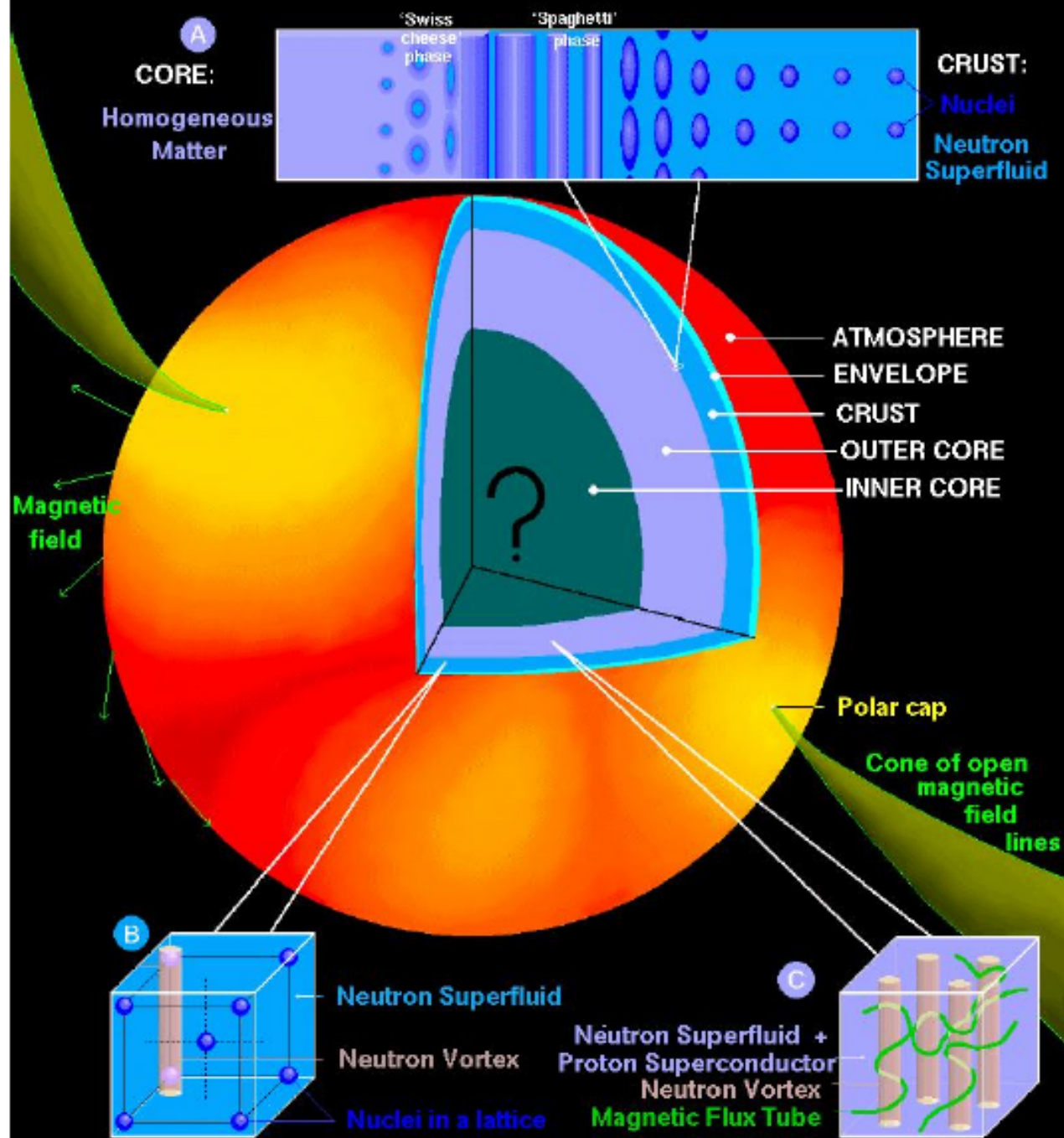
**Sevilla University**

Subtitle: HFB calculation at positive Fermi energies and  
vortex constrain on the pairing field

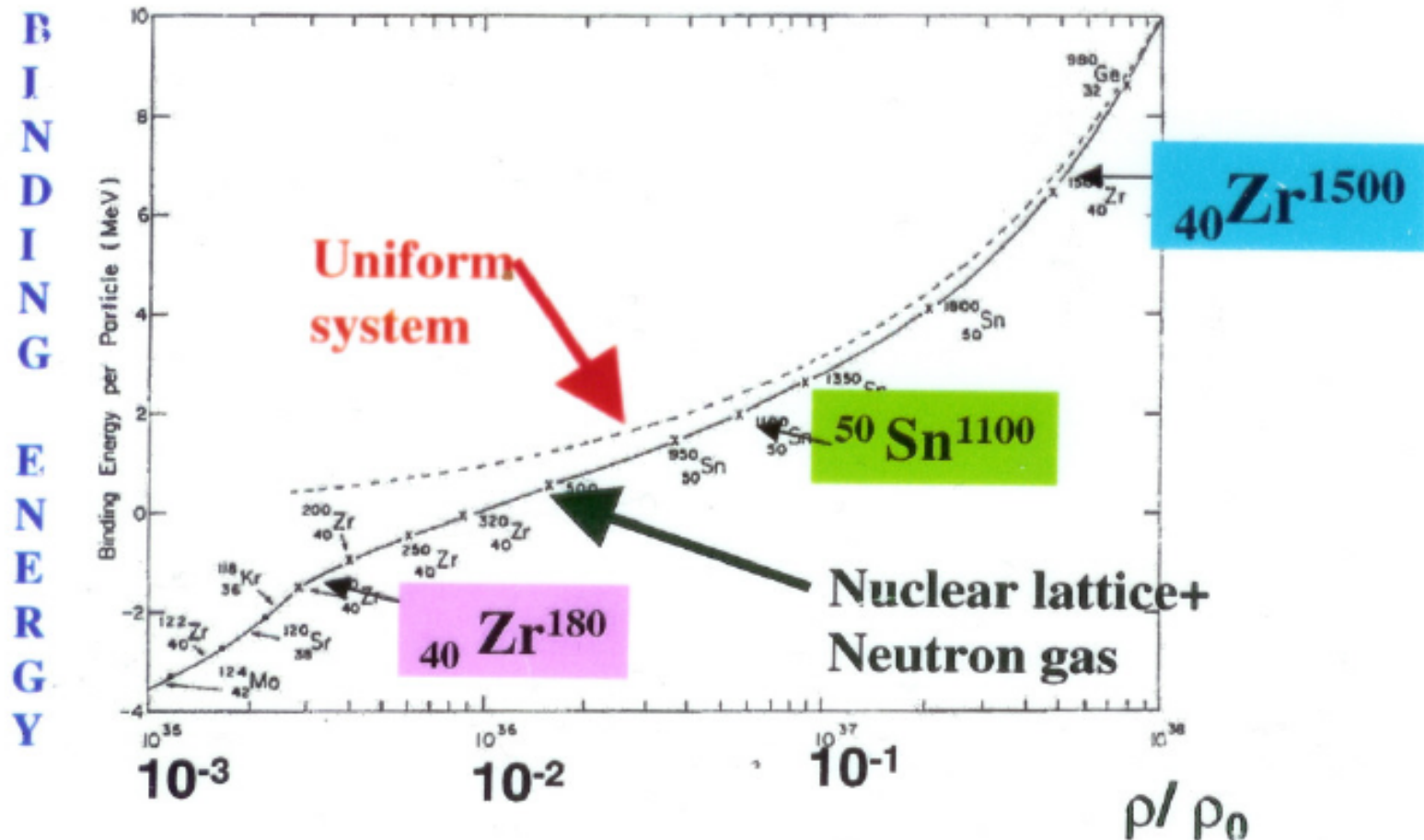
### Outline of the talk

- The scenario: the inner crust of a neutron star
- The model
- Results
- Comparison with other approaches
- Conclusions and perspectives

# A NEUTRON STAR: SURFACE and INTERIOR



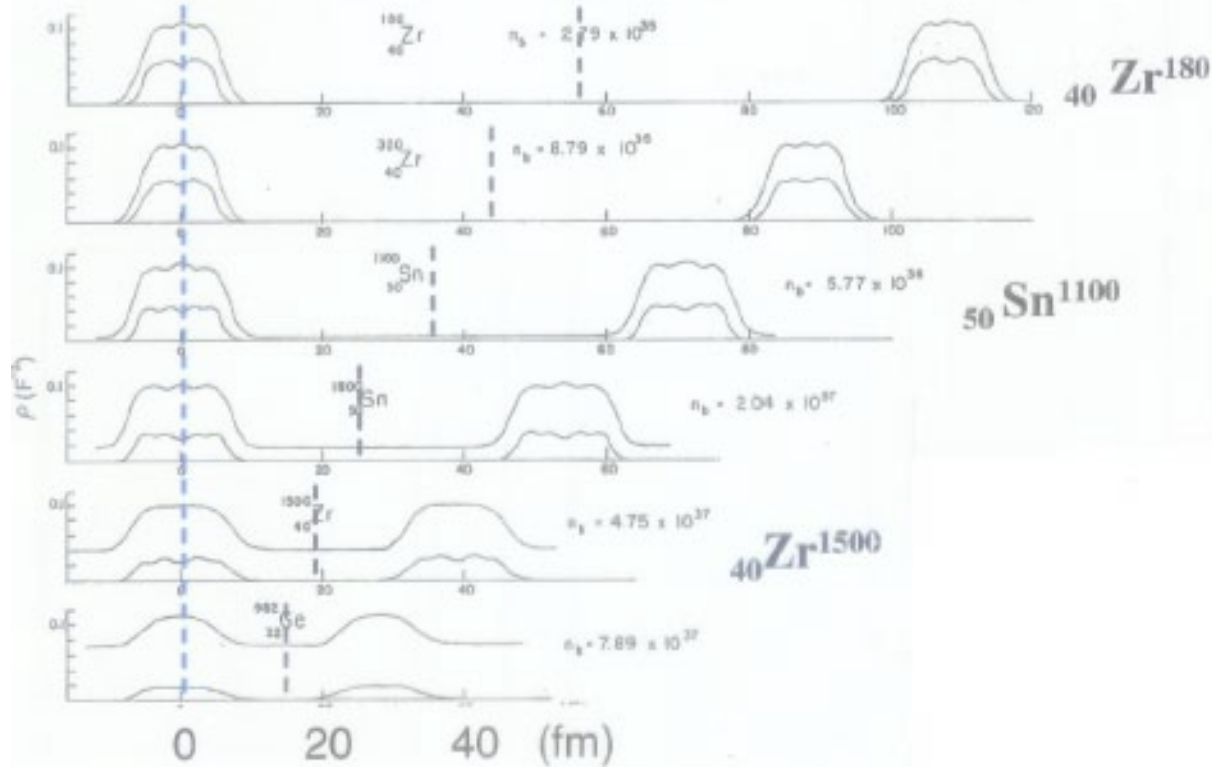
The inner crust: coexistence of a Coulomb lattice of finite nuclei with a sea of free neutrons

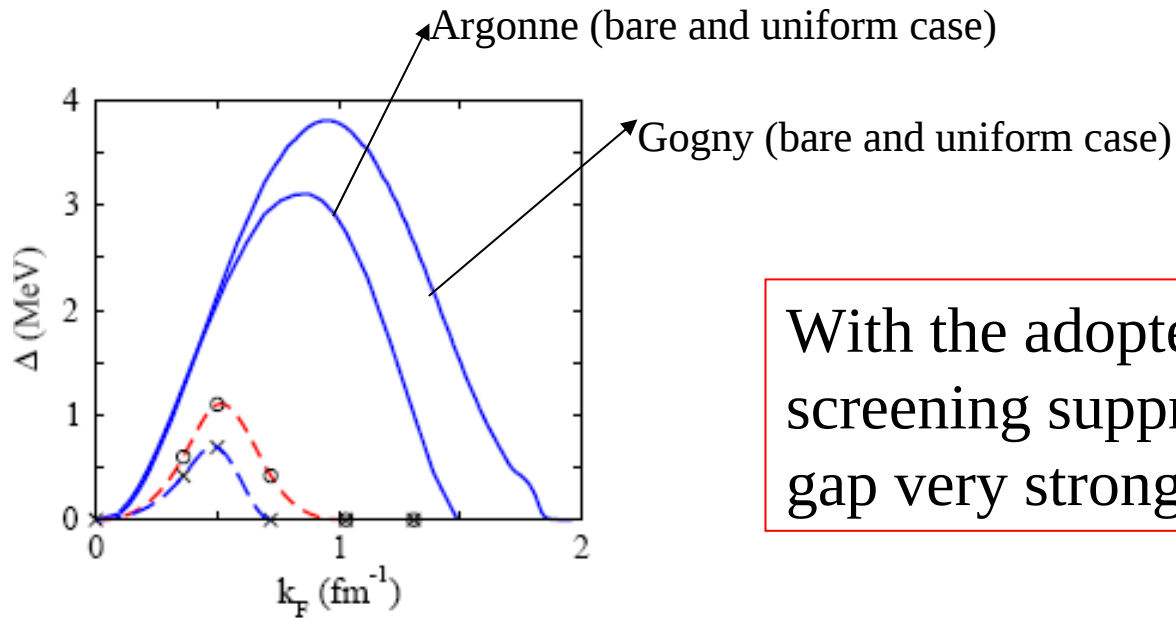


J. Negele, D. Vautherin  
Nucl. Phys. A207 (1974) 298

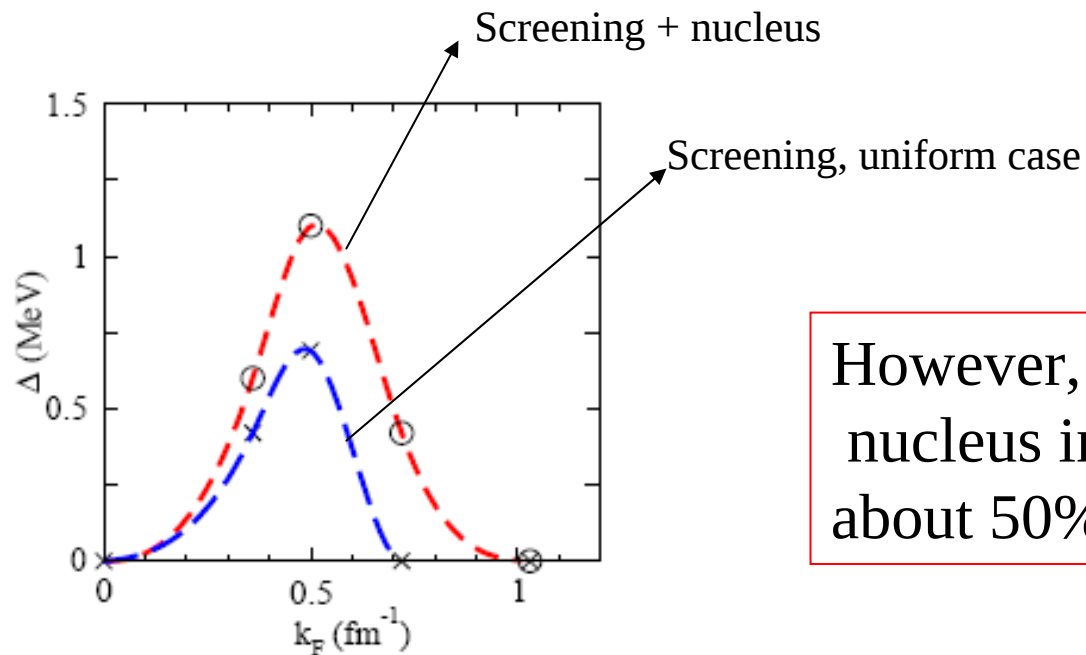
M. Baldo et al  
Nucl. Phys. A750 (2005) 409

Wigner-Seitz cell radius



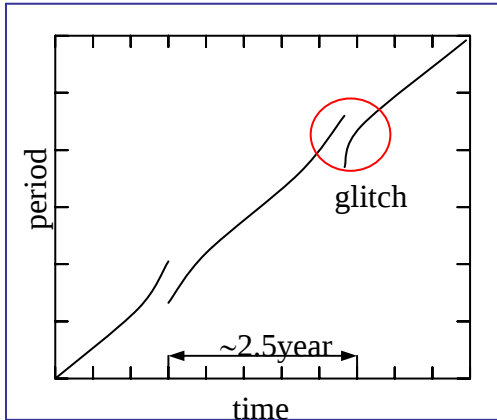


With the adopted interaction, screening suppresses the pairing gap very strongly for  $k_F > 0.7$  fm<sup>-1</sup>



However, the presence of the nucleus increases the gap by about 50%

# Glitches



As a rule, rotational period of a neutron star slowly increases because the system loses energy emitting electromagnetic radiation.

Sudden spin ups are measured, at regular intervals

One of the accredited explanations



Superfluid nature of nucleons in the inner crust

P. W. Anderson and N. Itoh, Nature 256(1975)25

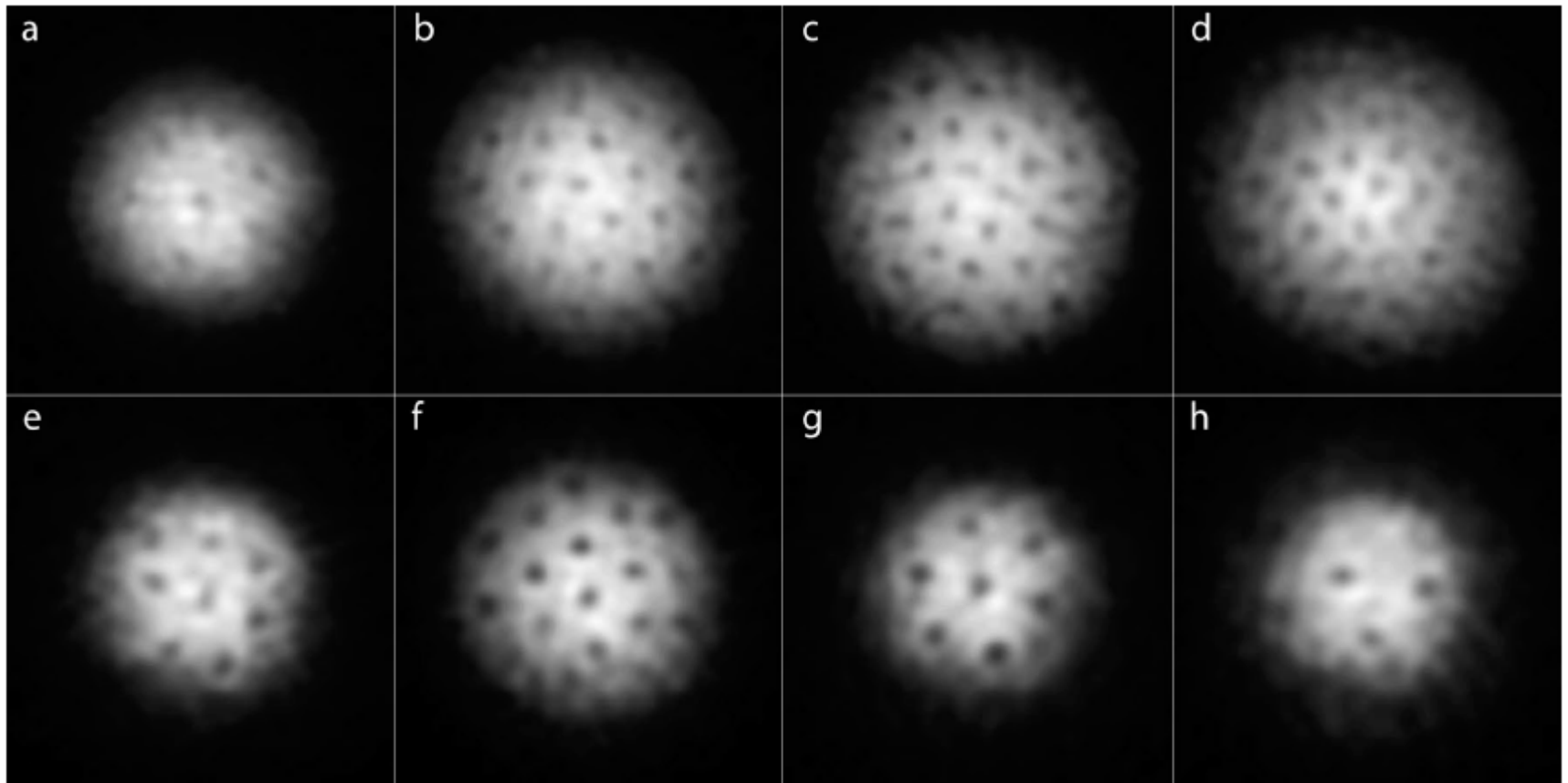
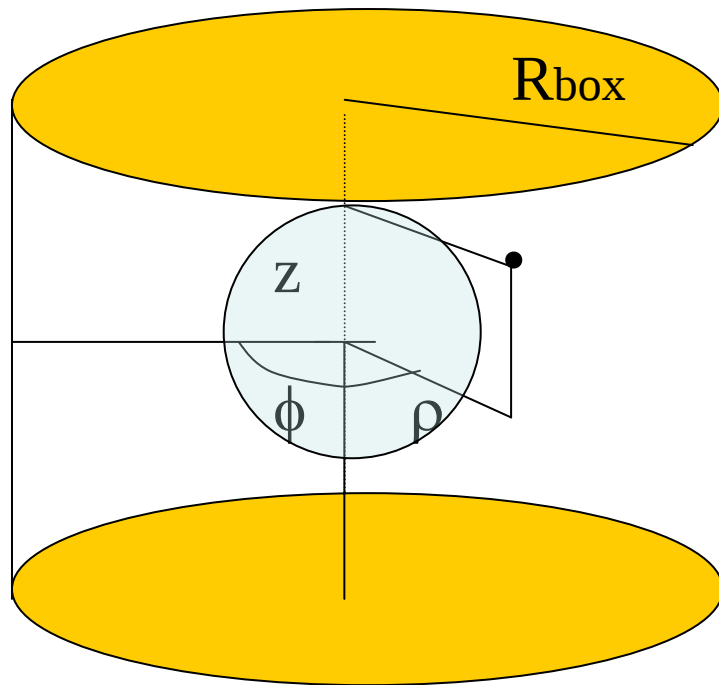


Fig. 2: Vortices in a strongly interacting gas of fermionic atoms on the BEC- and the BCS-side of the Feshbach resonance. At the given field, the cloud of lithium atoms was stirred for 300 ms (a) to 500 ms (b-h) followed by an equilibration time of 500 ms. After 2 ms of ballistic expansion, the magnetic field was ramped to 735 G for imaging (see text for details). The magnetic fields were (a) 740 G, (b) 766 G, (c) 792 G, (d) 812 G, (e) 833 G, (f) 843 G, (g) 853 G and (h) 863 G. The field of view of each image is  $880\text{ }\mu\text{m} \times 880\text{ }\mu\text{m}$ .

Corrente – fase – delta

# Microscopic quantum calculation of the vortex-nucleus system



We solve the  
HFB (De Gennes)  
equations expanding on  
a single-particle basis in  
cylindrical coordinates

$$u_{\alpha}(\rho, \phi, z) \sim \sum_{nk} J_{nm}(\rho) \sin(kz) e^{im\phi} u_{nk;m;\alpha}$$
$$v_{\alpha}(\rho, \phi, z) \sim \sum_{nk} J_{nm-\nu}(\rho) \sin(kz) e^{i(m-\nu)\phi} v_{nk;m;\alpha}$$

Using a zero-range pairing interaction,

only local quantities  
are needed



$$\begin{aligned}\eta(\rho, z) &= \sum_{\alpha} v_{\alpha}(\rho, \phi, z) v_{\alpha}^{*}(\rho, \phi, z) \\ V(\rho, z) &= \text{Skyrme Density Functional} \\ \kappa(\rho, \phi, z) &= \sum_{\alpha} u_{\alpha}(\rho, \phi, z) v_{\alpha}^{*}(\rho, \phi, z) \\ \Delta(\rho, \phi, z) &= \Delta(\rho, z) e^{i\nu\phi} = \\ &\quad \frac{g}{2} \sum_{\alpha} u_{\alpha}(\rho, \phi, z) v_{\alpha}^{*}(\rho, \phi, z)\end{aligned}$$

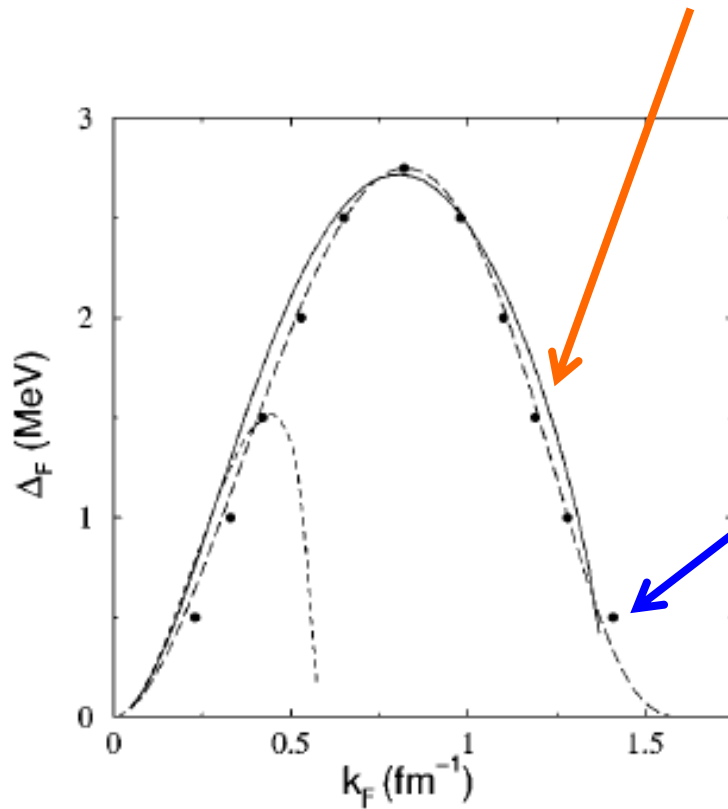
- The equations are solved self-consistently
- SI Skyrme interaction (Brink-Vautherin)
- Protons are constrained to have a spherical geometry
- No spin-orbit interaction

$$\begin{bmatrix} \hat{K} + V(\rho, z) - \lambda & \Delta(\rho, \phi, z) \\ \Delta^{*}(\rho, \phi, z) & -(\hat{K} + V(\rho, z) - \lambda) \end{bmatrix} \begin{bmatrix} u_{\alpha}(\rho, \phi, z) \\ v_{\alpha}(\rho, \phi, z) \end{bmatrix} = E_{\alpha} \begin{bmatrix} u_{\alpha}(\rho, \phi, z) \\ v_{\alpha}(\rho, \phi, z) \end{bmatrix}$$

$$V = -480(1 - 0.7(\rho/\rho_0)^{0.45})\delta(r_1 - r_2)\text{MeVfm}^3$$

**Ecut = 60 MeV**

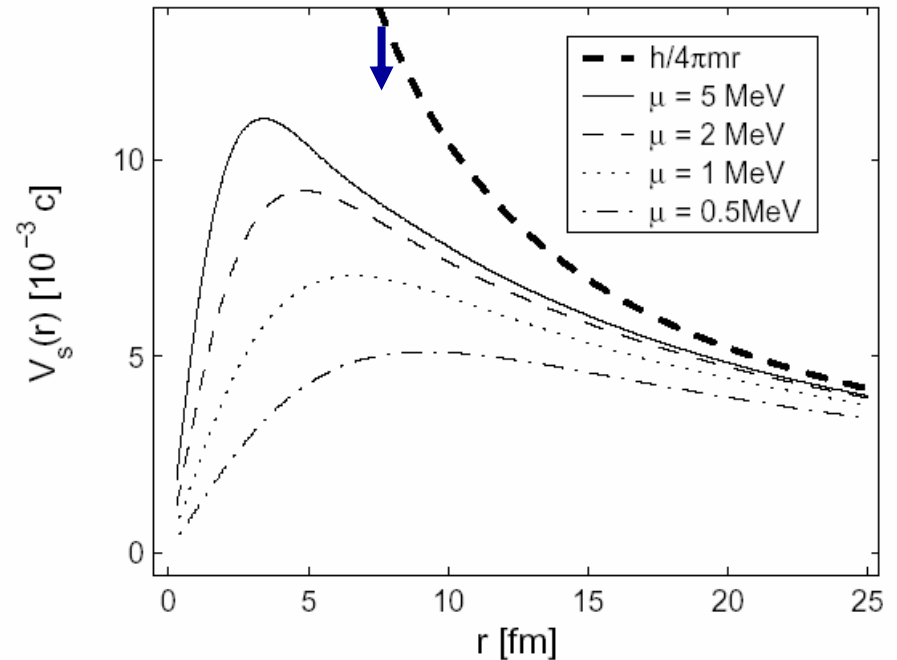
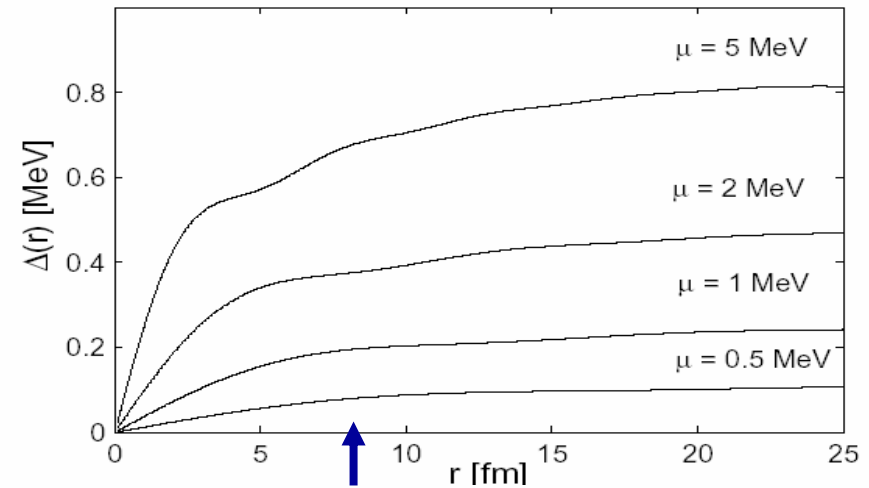
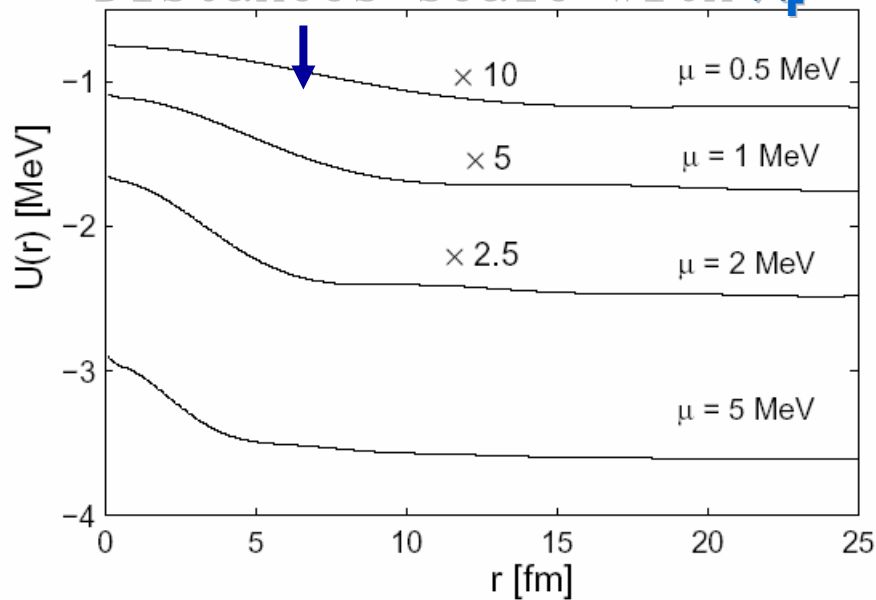
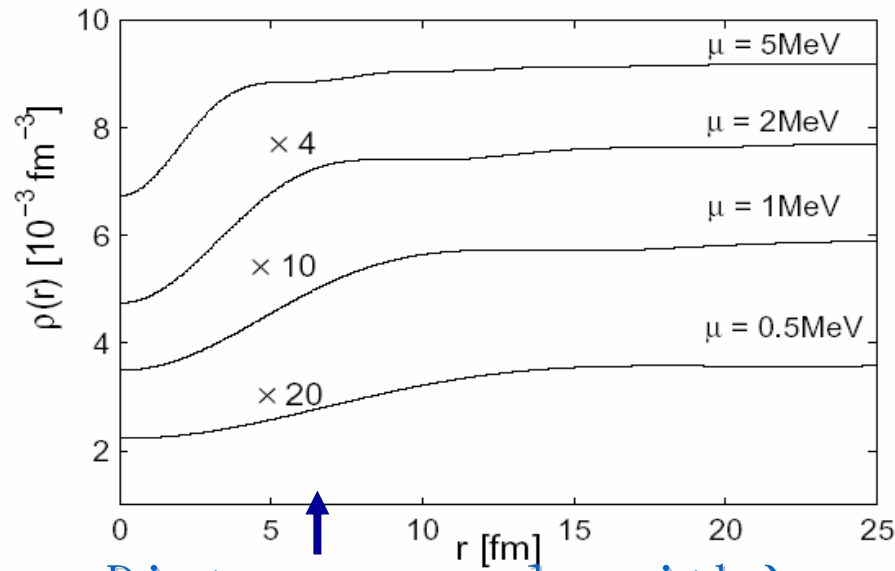
Ecut = 50 MeV in our calculations



**Gogny force  
( gap essentially  
identical to the  
bare force at  
 $k_F = 0.7 \text{ fm}^{-1}$  )**

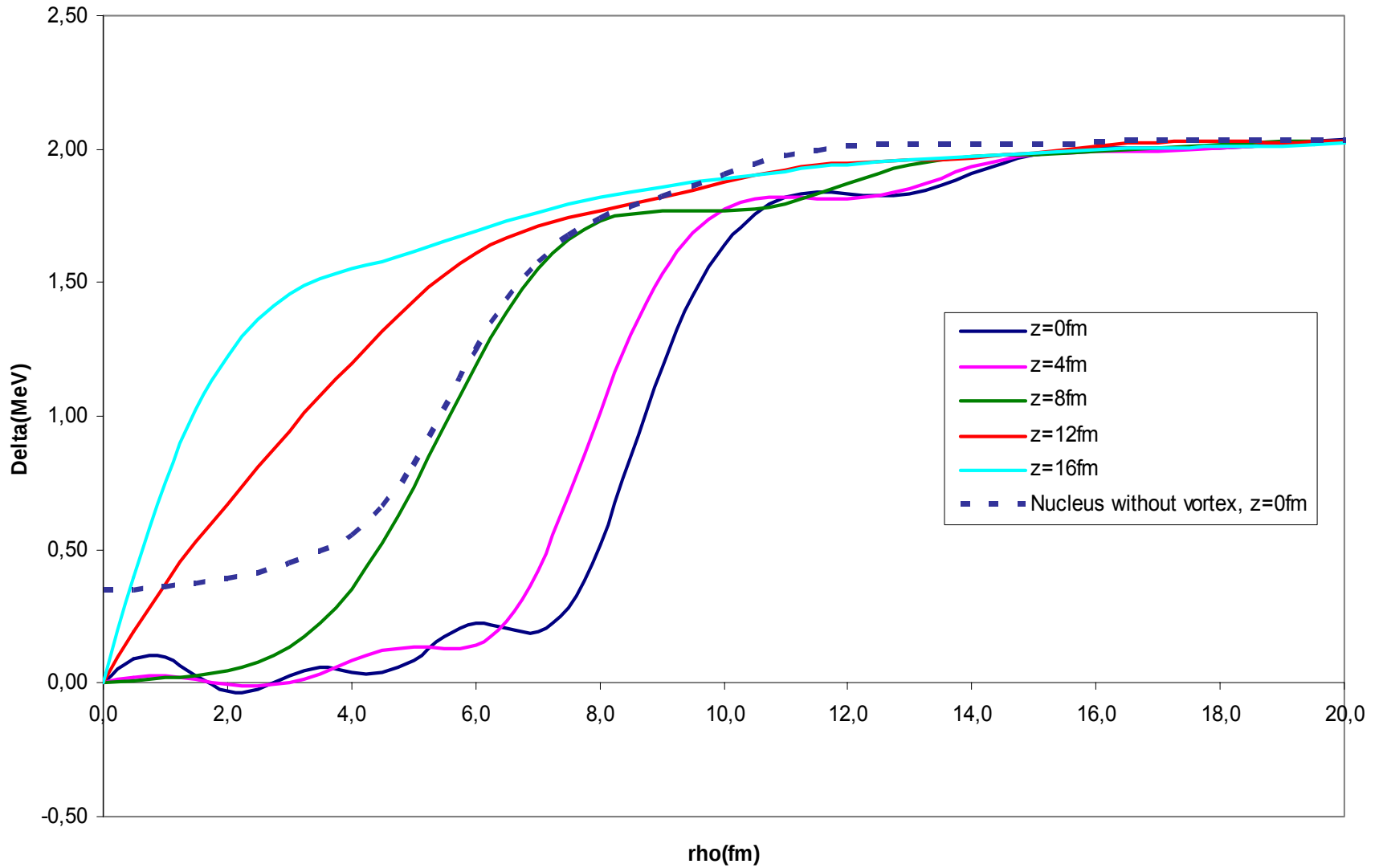
**E. Garrido et al. Phys. Rev. C60(1999)64312**

# Vortex in uniform matter: Y. Yu and A. Bulgac, PRL 90, 161101 (2003)

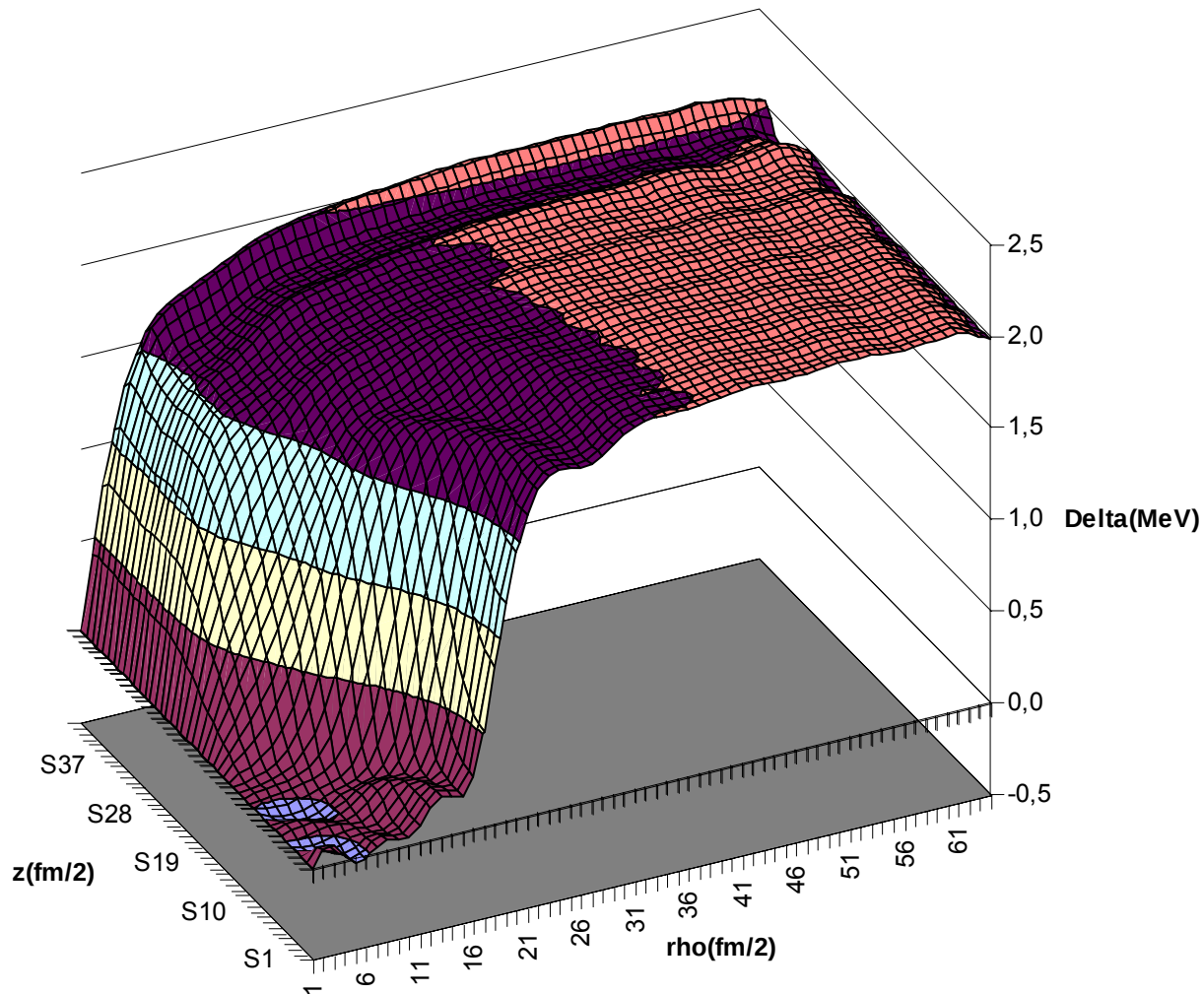


# Pairing gap of pinned vortex

Pairing of Pinned Vortex

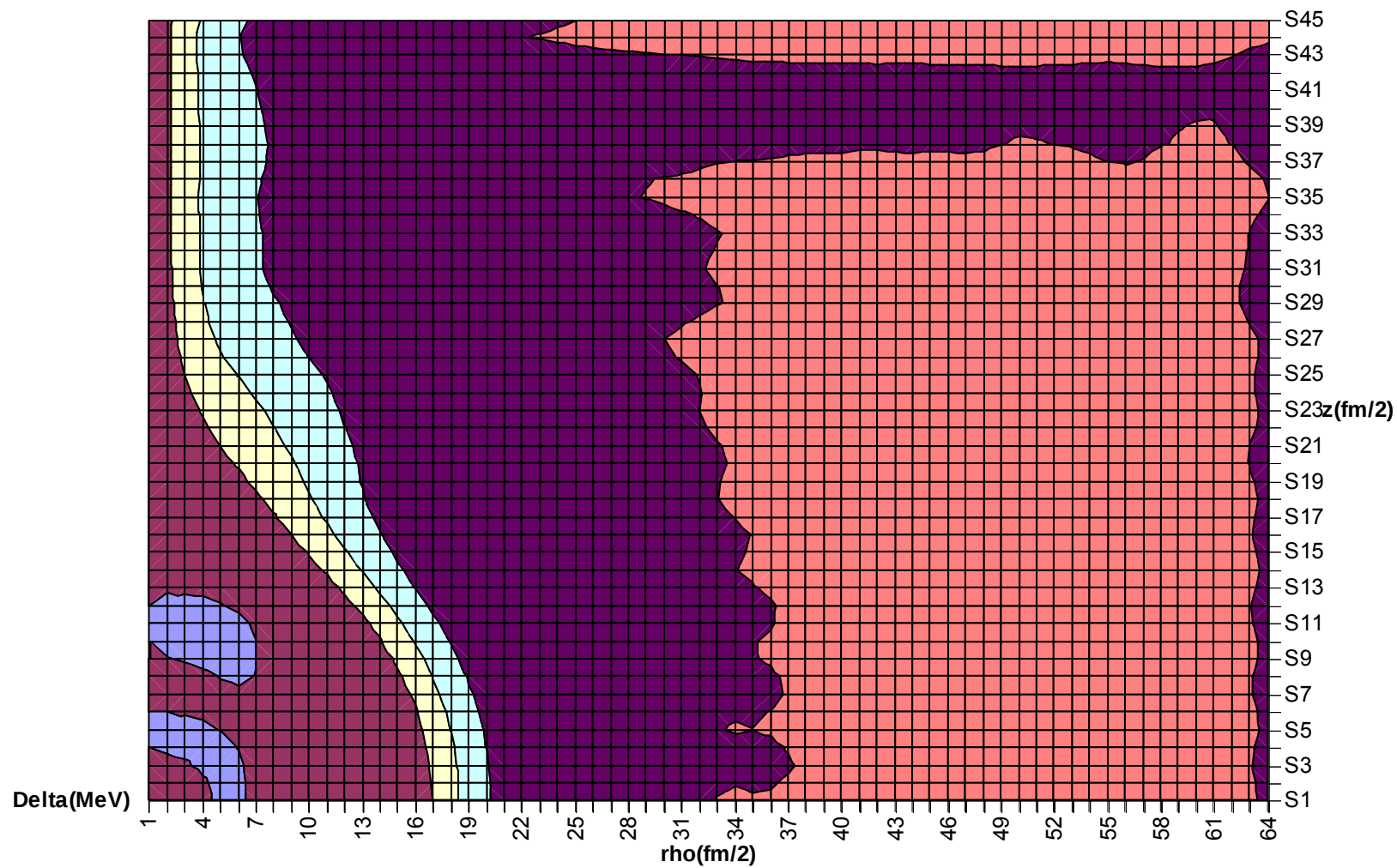


### Pairing of Pinned Vortex



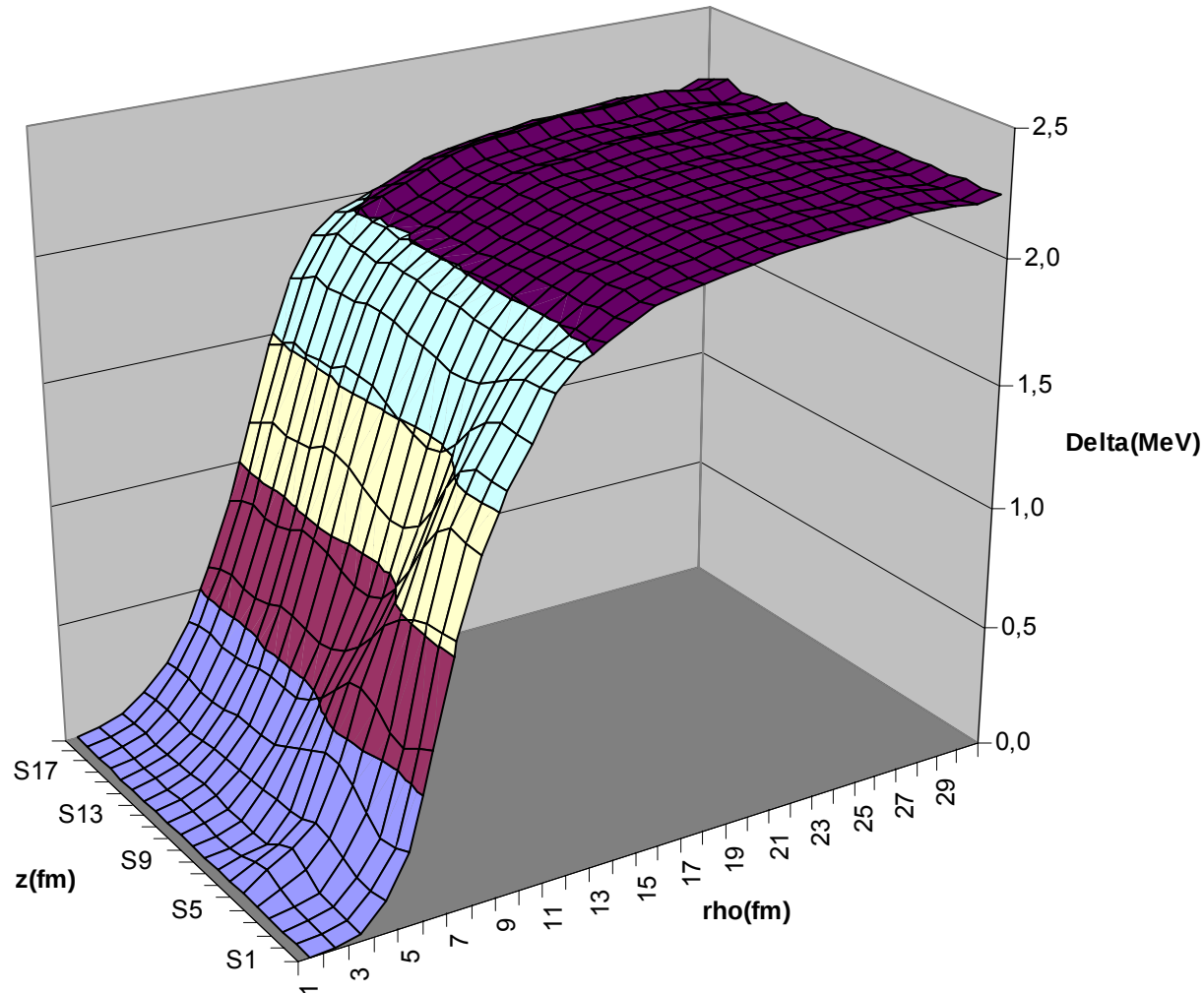
The vortex is expelled from the nuclear volume, because there are no levels available which can satisfy the parity and angular momentum conditions.

# Pairing of Pinned Vortex



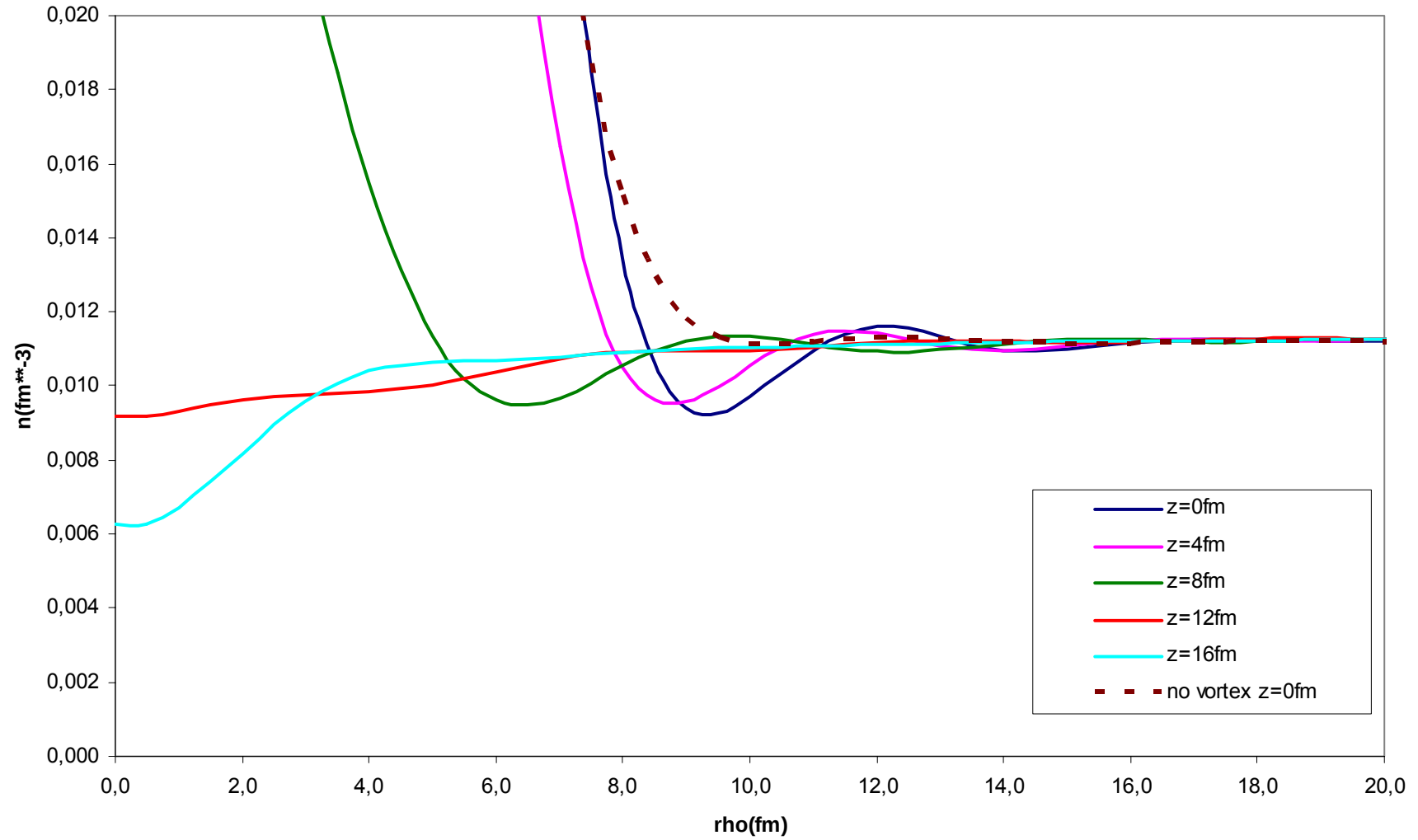
## Pairing gap of pinned vortex, $\nu = 2$

Pairing field of pinned vortex,  $\nu=2$

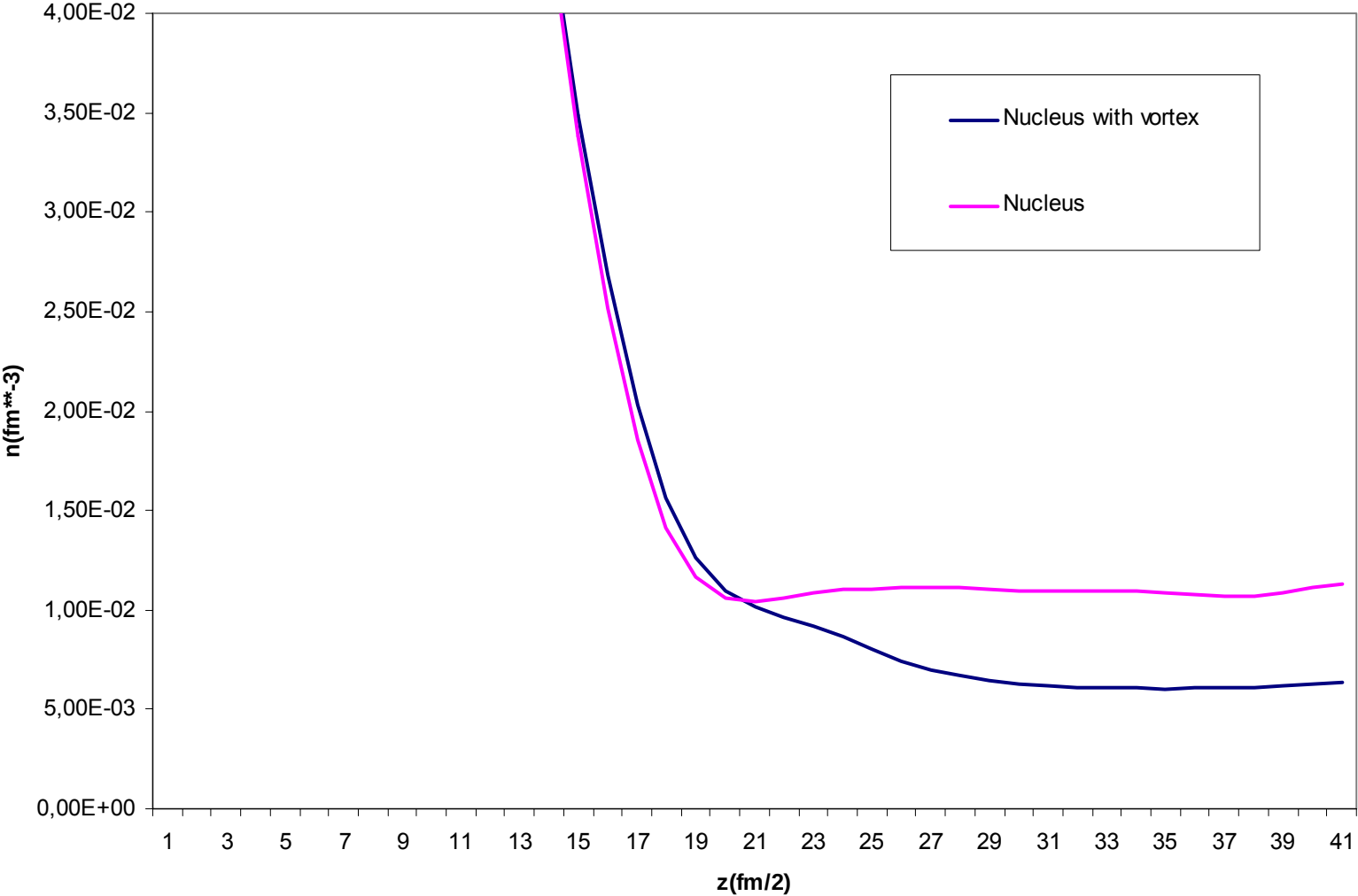


In this case the vortex goes through the nuclear volume, essentially undisturbed, because the levels can satisfy the parity and angular momentum conditions.

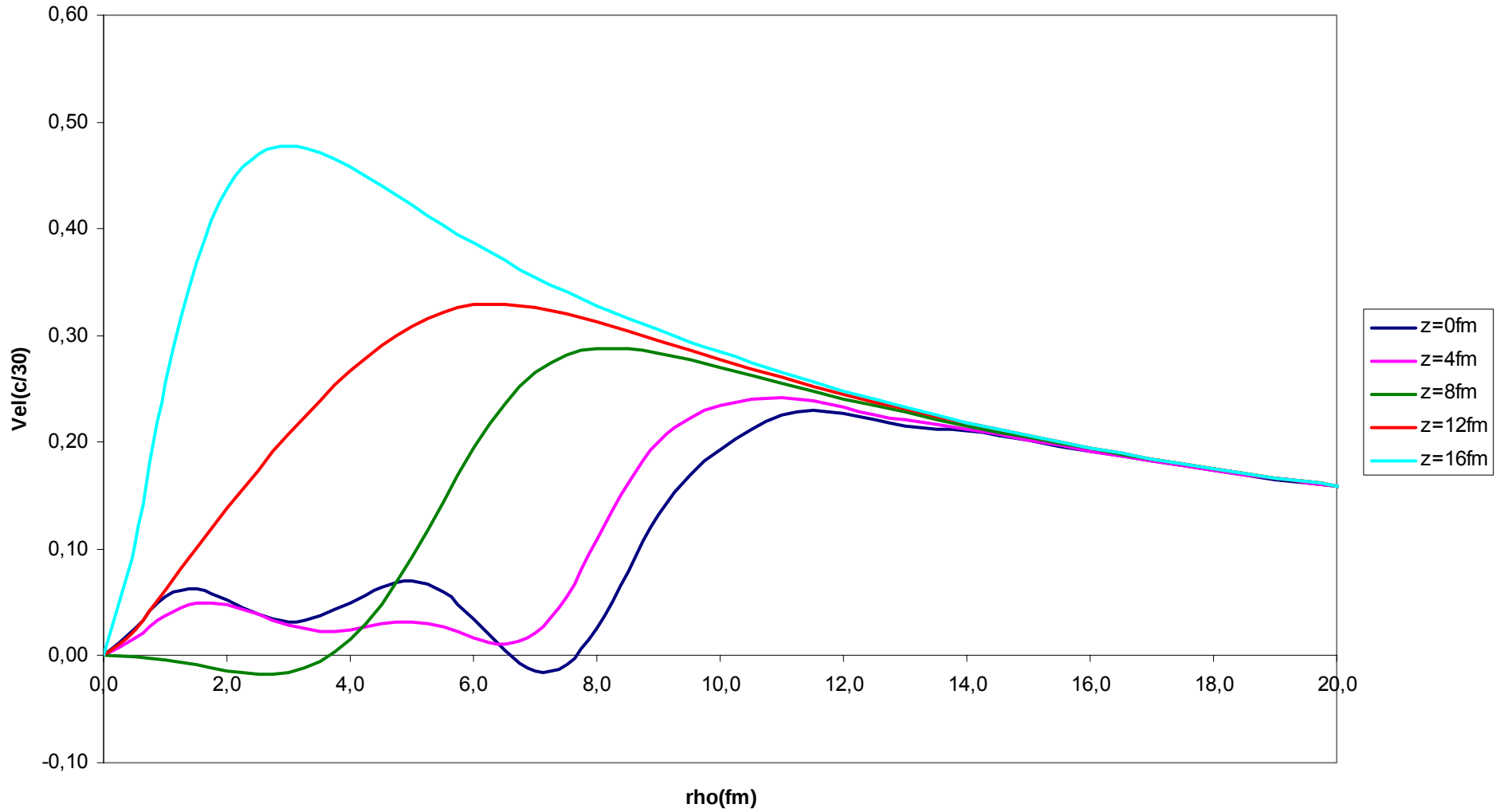
# Density of Pinned Vortex



Density along the z-axis



## Velocity of Pinned Vortex



**The energy to create a vortex** is obtained taking the difference between two calculations in the cylindrical box, one with and the other without the vortex, each with the same number of particles.

In each calculation, the energy has three different contributions:

- Kinetic energy
- Mean field potential energy
- Pairing field potential energy

|       | Uniform | Vortex  | Diff.           |                                                |
|-------|---------|---------|-----------------|------------------------------------------------|
| Ekin  | 13871.1 | 13815.4 | - 55.7          |                                                |
| Epot  | -3660.2 | -3700.1 | - 39.9          |                                                |
| Epair | -2289.6 | -2125.0 | 164.6           |                                                |
|       |         |         | <u>69.0</u> MeV | Energy cost for a vortex in uniform matter     |
|       | Nucleus | Vortex  | Diff.           |                                                |
| Ekin  | 16279.5 | 16185.8 | - 93.7          |                                                |
| Epot  | -7077.9 | -7126.7 | - 48.8          |                                                |
| Epair | -2400.6 | -2193.9 | 206.7           |                                                |
|       |         |         | <u>64.2</u> MeV | Energy cost for a vortex pinned to the nucleus |

Pinning energy:  $64.2 - 69.0 = - 4.8$  MeV

Preliminary result: accuracy is being checked

### Previous calculations of pinned vortices:

-R. Epstein and G. Baym, *Astrophys. J.* 328(1988)680

Analytic treatment based on the Ginzburg-Landau equation

-F. De Blasio and O. Elgaroy, *Astr. Astroph.* 370,939(2001)

Numerical solution of De Gennes equations with a fixed nuclear mean field and imposing cylindrical symmetry (spaghetti phase)

-P.M. Pizzochero and P. Donati, *Nucl. Phys. A*742,363(2004)

Semiclassical model.

## Conclusions and perspectives

-We have solved the HFB equations for a single vortex in the crust of neutron stars, considering explicitly the presence of the nucleus, generalizing previous studies in uniform matter.

-We have found that ( $\nu=1$ ) vortex stay outside of the nuclear volume, where the pairing goes to zero.

-The nuclear surface acquires a slight prolate deformation

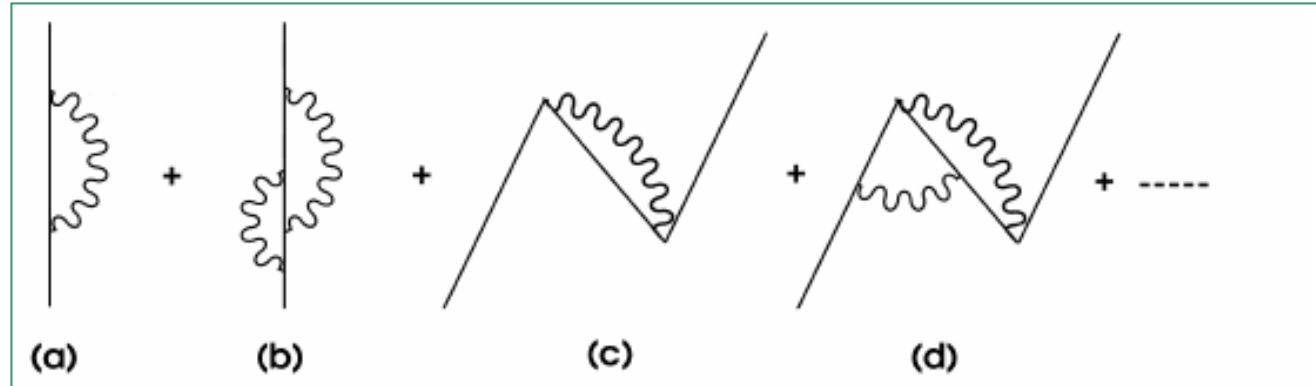
-Preliminary numerical results at  $k_F = 0.7\text{fm}^{-1}$  indicate that the pinning energy is of the order of a few MeV.

Many open questions. Among them:

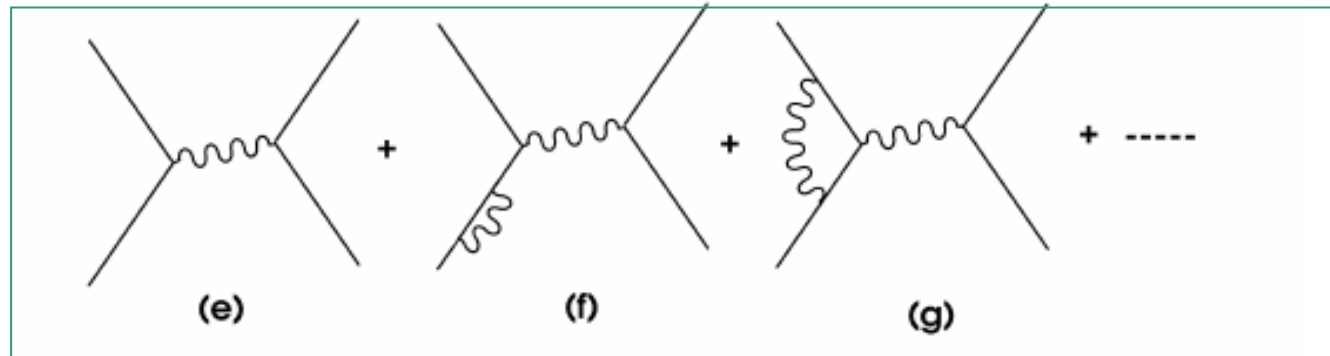
- Include medium polarization effects
- 3D calculations
- Vortex dynamics

## Going beyond mean field: medium polarization effects

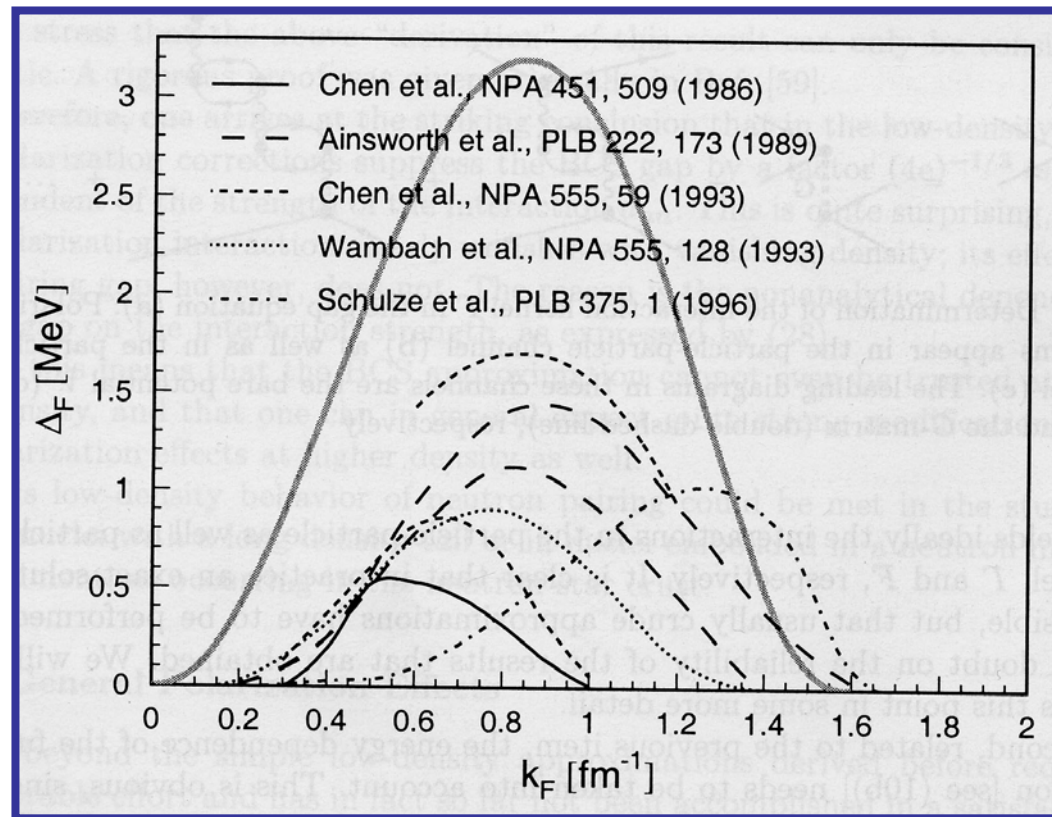
Self-energy



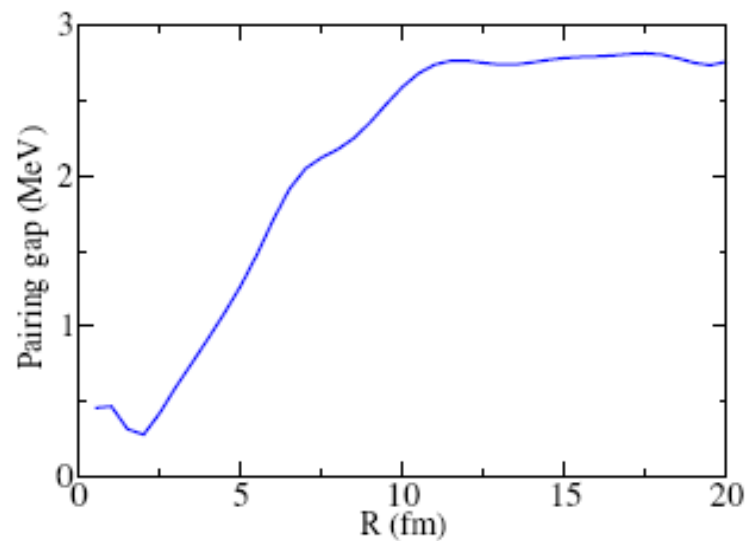
Induced interaction  
(screening)



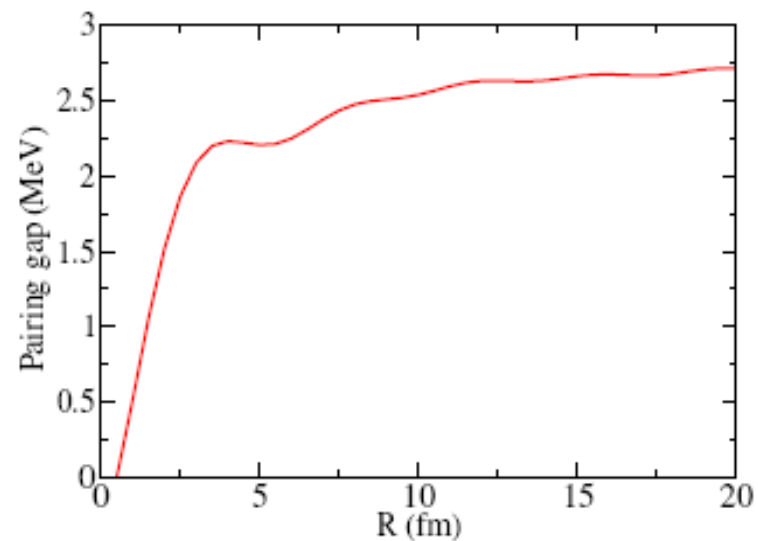
The results of various calculations in neutron matter can be understood in terms of Landau parameters



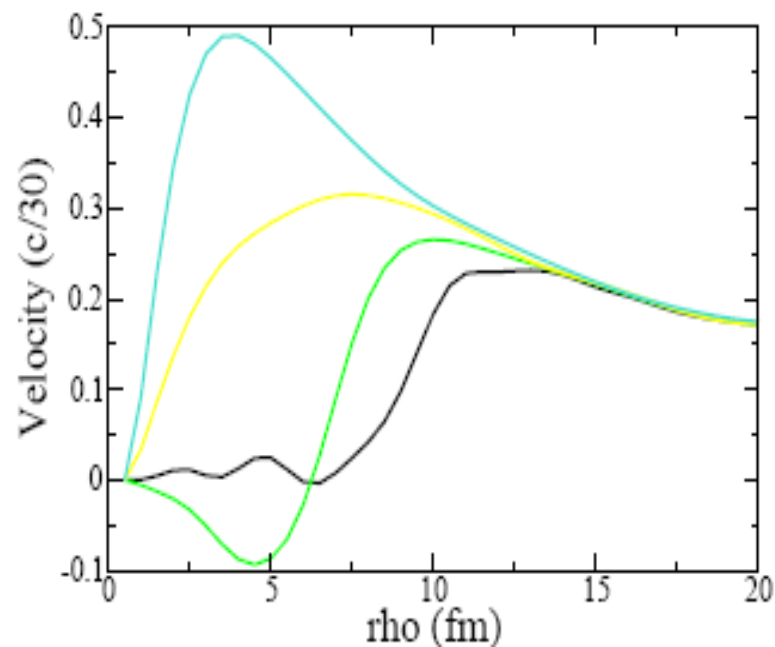
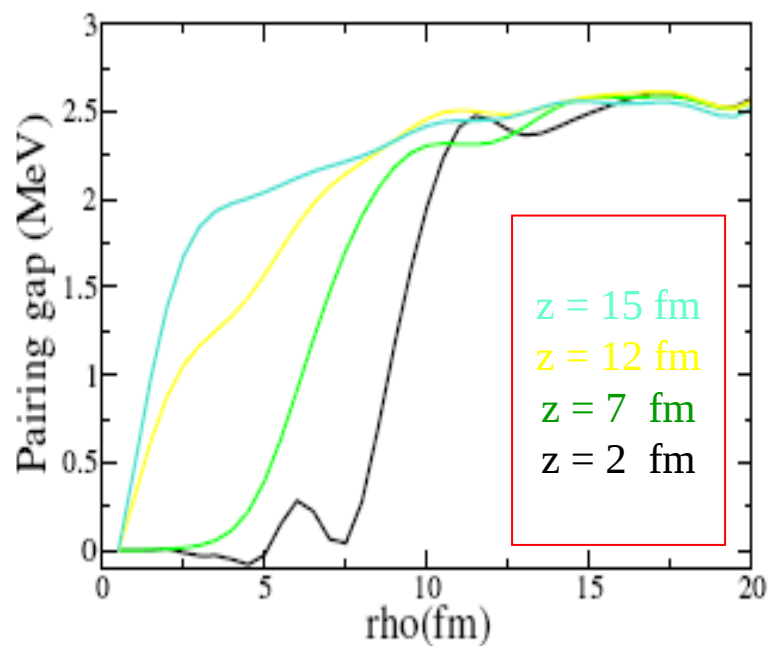
Pairing gap in the Wigner cell (no vortex)



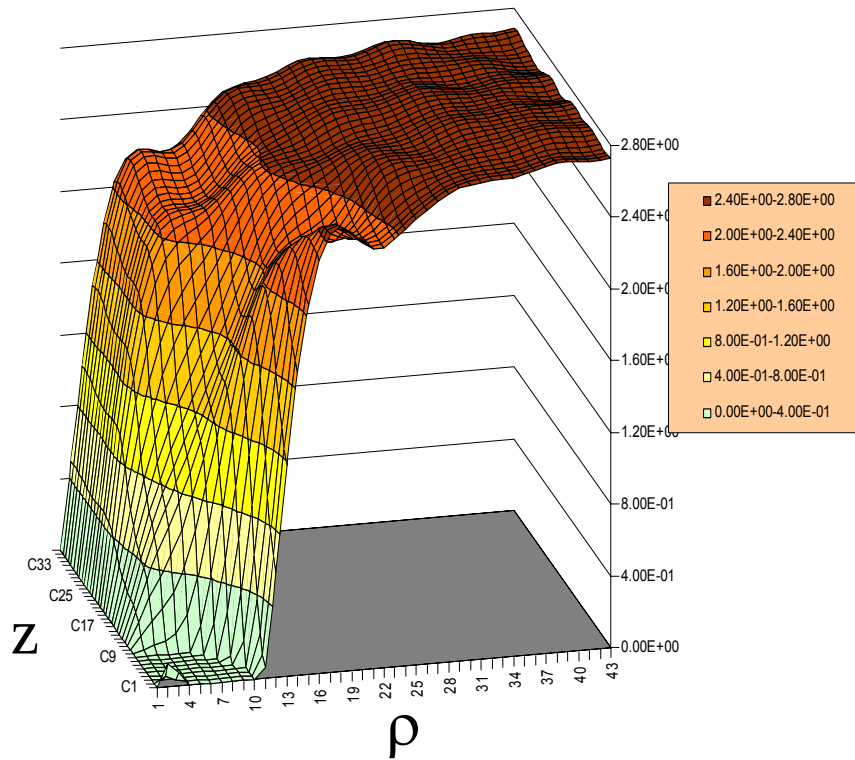
Pairing gap for a vortex in uniform matter



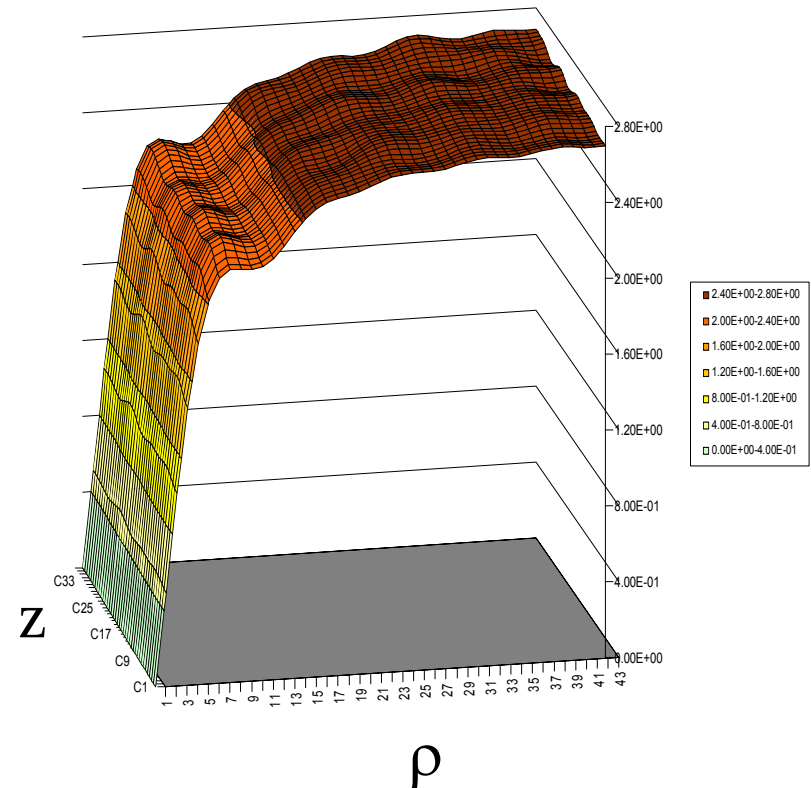
Pairing gap and velocity field for vortex in the Wigner cell



## Vortex pairing field



Proximity effects in the presence of the nucleus



Vortex in uniform neutron matter

neutron matter of the  
inner crust



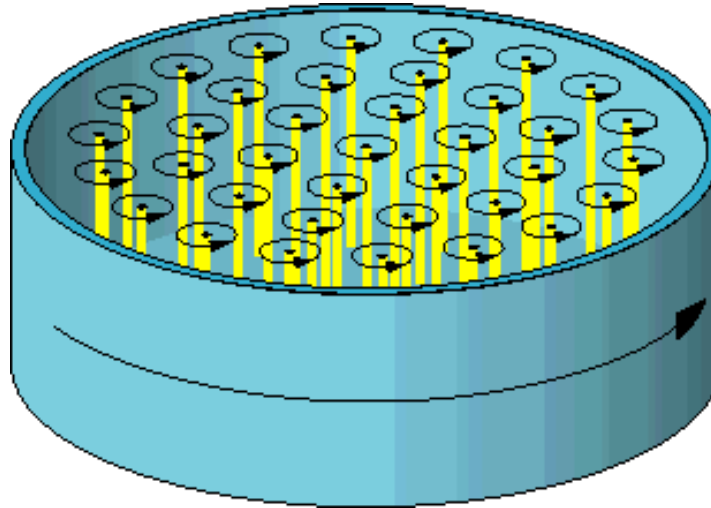
superfluid in a rotating  
container

in order to minimize  
the energy of the system



it develops an array of  
microscopic linear vortices of  
superfluid matter

Modificare!!



Each vortex line carries one unit of angular momentum