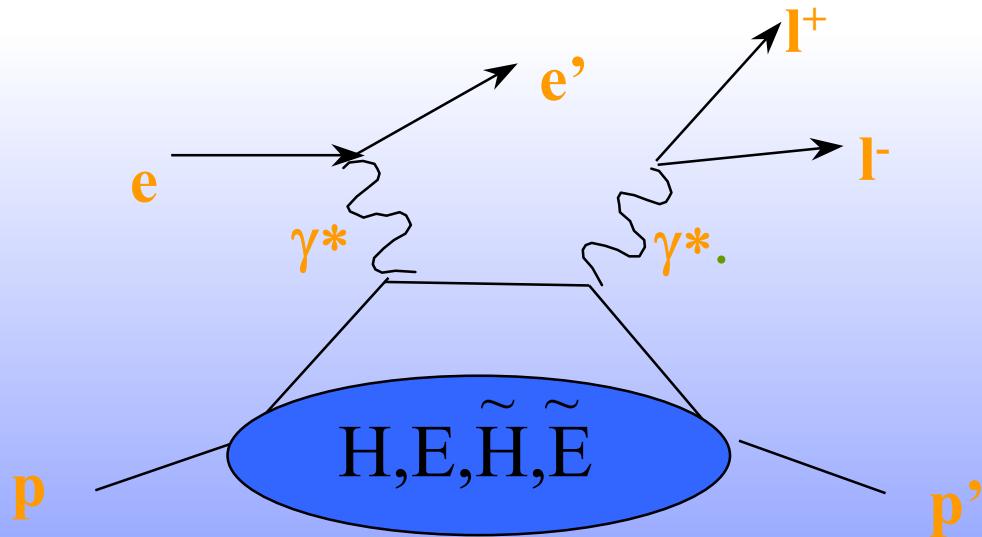


1/ DDVCS

(+ 2/ first study on the (x,t) correlations of the GPDs)



Seattle, 10/07/03
M. Guidal, IPN Orsay

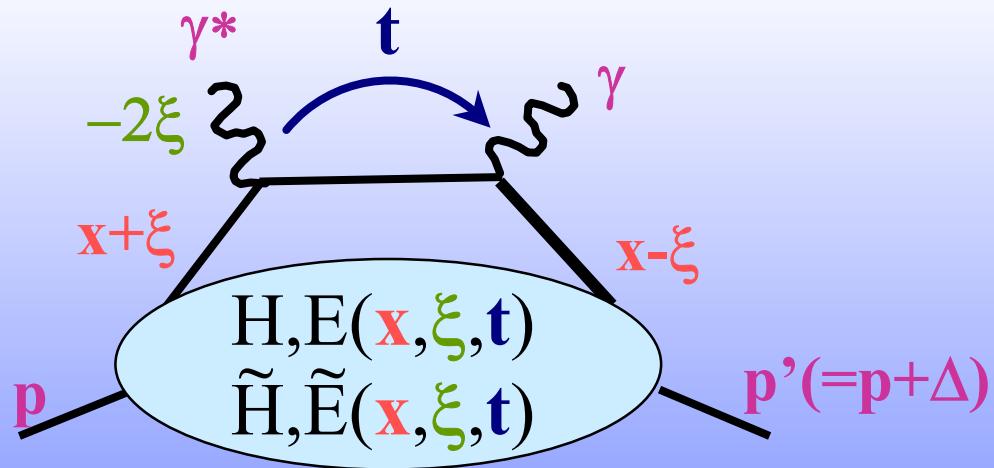
1/ DDVCS

(Double Deeply Virtual Compton Scattering)

**Direct access to GPDS
(no deconvolution)**

The handbag diagram

(Ji, Radyushkin, Mueller,)



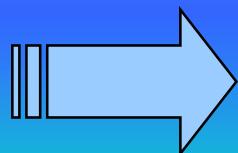
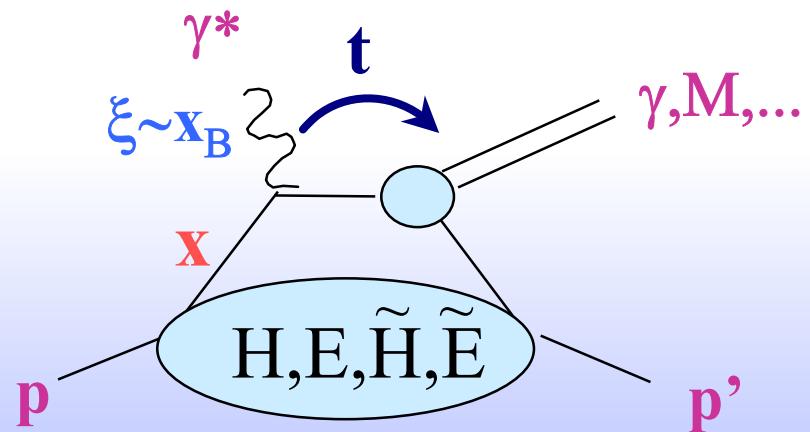
$$\Rightarrow \left\{ \gamma^- [H^q(x, \xi, t) \bar{N}(p') \gamma^+ N(p) + E^q(x, \xi, t) \bar{N}(p') i \sigma^{+\kappa} \underline{\Delta}_{\kappa} N(p)] \right.$$
$$\left. + \gamma^5 \gamma^- [\tilde{H}^q(x, \xi, t) \bar{N}(p') \gamma^+ \gamma^5 N(p) + \tilde{E}^q(x, \xi, t) \bar{N}(p') \gamma^5 \underline{\Delta}^+ N(p)] \right\}$$

$H^q(x, \xi, t, Q^2)$ but only ξ , t and Q^2 accessible experimentally

$$\xi = \frac{x_B/2}{1-x_B/2} \quad t = (p-p')^2$$

$$Q^2 = -(e-e')^2$$

$x \neq x_B !$



x : mute variable

Deconvolution needed!

$$\frac{d\sigma}{dQ^2 dx_B dt} \sim \left| A \int_{-1}^1 \frac{H^q(x, \xi, t, Q^2)}{x - \xi + i\epsilon} dx + B \int_{-1}^1 \frac{E^q(x, \xi, t, Q^2)}{x - \xi + i\epsilon} dx + \dots \right|^2$$

$H^q(x, \xi, t)$ is *real* but amplitude is *complex*

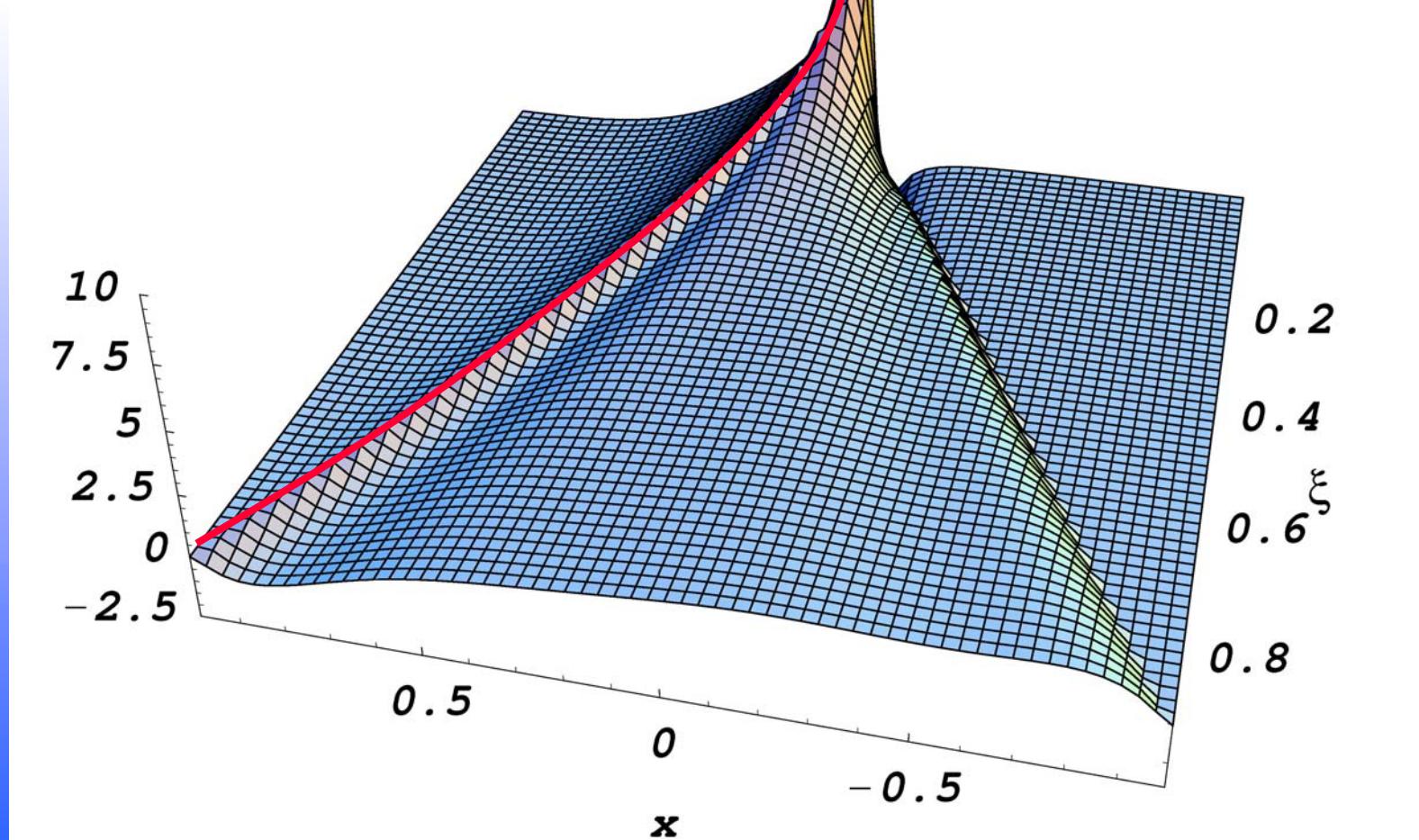
$$\int_{-1}^1 \frac{H^q(x, \xi, t)}{x - \xi + i\epsilon} dx \longrightarrow \text{Differential cross section}$$

$$= \int_{-1}^1 \frac{H^q(x, \xi, t)}{x - \xi} dx \longrightarrow \text{Real part} \\ (\text{charge asymmetry})$$

$$- i \pi \delta(x - \xi) H^q(x, \xi, t) dx \longrightarrow \text{Imaginary part} \\ (\text{spin asymmetries})$$

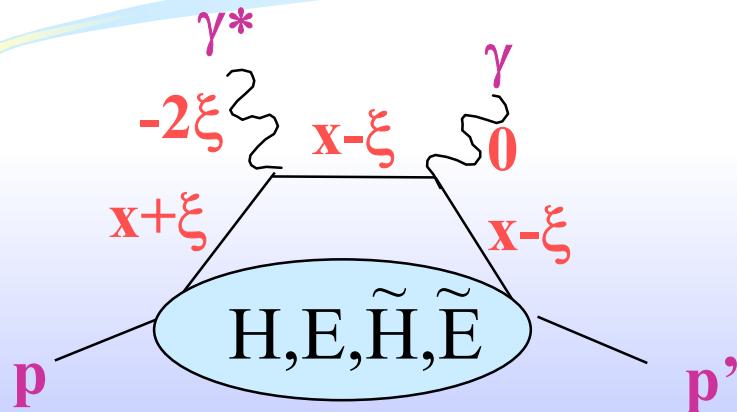
$\underbrace{}$
 $H^q(\xi, \xi, t)$

$H^u(x, \xi, t=0)$



(Goeke, Polyakov, Vanderhaeghen)

DVCS



$$\sigma \sim \left| \int_{-1}^1 \frac{H(x, \xi, t)}{x - \xi + i\epsilon} dx \right|^2$$

Im $\sim H(\xi, \xi, t)$

$$Re \sim \left| \int_{-1}^1 \frac{H(x, \xi, t)}{x - \xi} dx \right|^2$$

DDVCS



$$\sigma \sim \left| \int_{-1}^1 \frac{H(x, \xi, t)}{x - (2\xi' - \xi) + i\epsilon} dx \right|^2$$

Im $\sim H(2\xi' - \xi, \xi, t)$

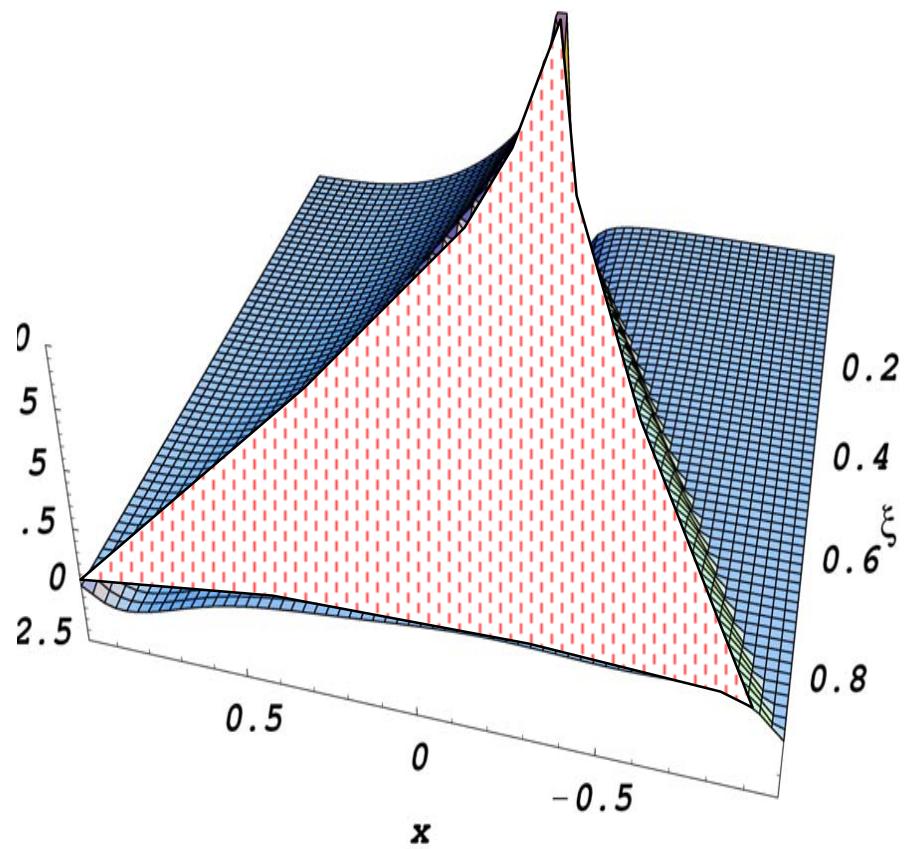
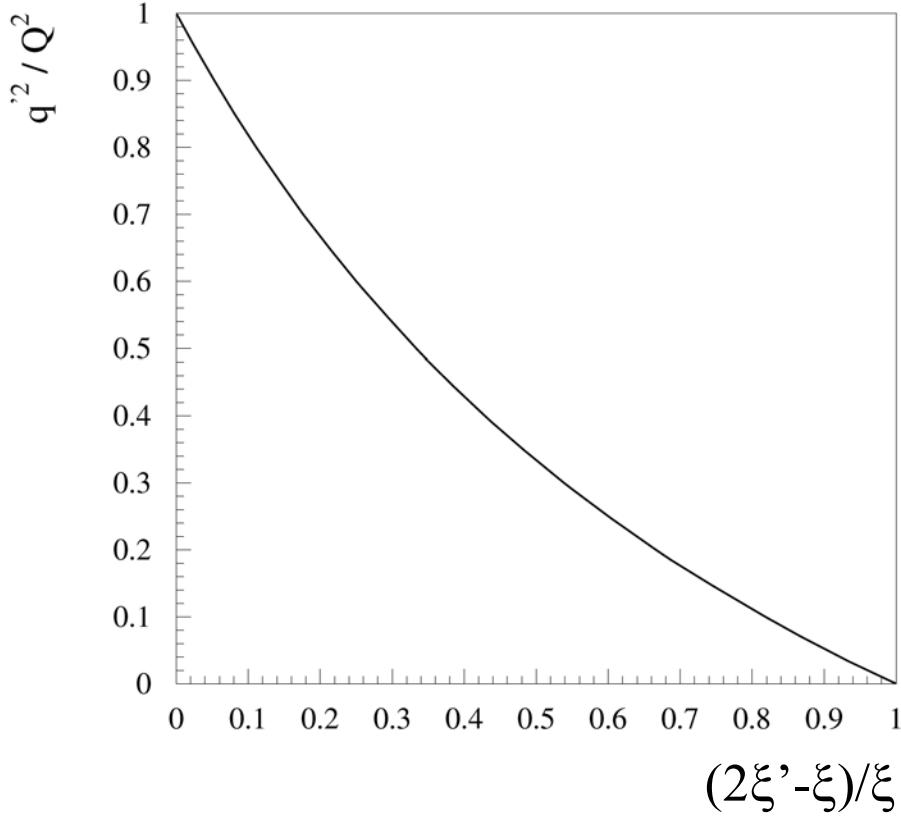
$$Re \sim \left| \int_{-1}^1 \frac{H(x, \xi, t)}{x - (2\xi' - \xi)} dx \right|^2$$

$$q'^2=0 \text{ GeV}^2 \quad q'^2=Q^2$$

$0 < 2\xi' - \xi < \xi$

Numerical example :
 $Q^2=5 \text{ GeV}^2, q'^2=2 \text{ GeV}^2$
 $(2\xi' - \xi) = .43\xi$

SSA in $\gamma^* + p \rightarrow l^+l^- + p$

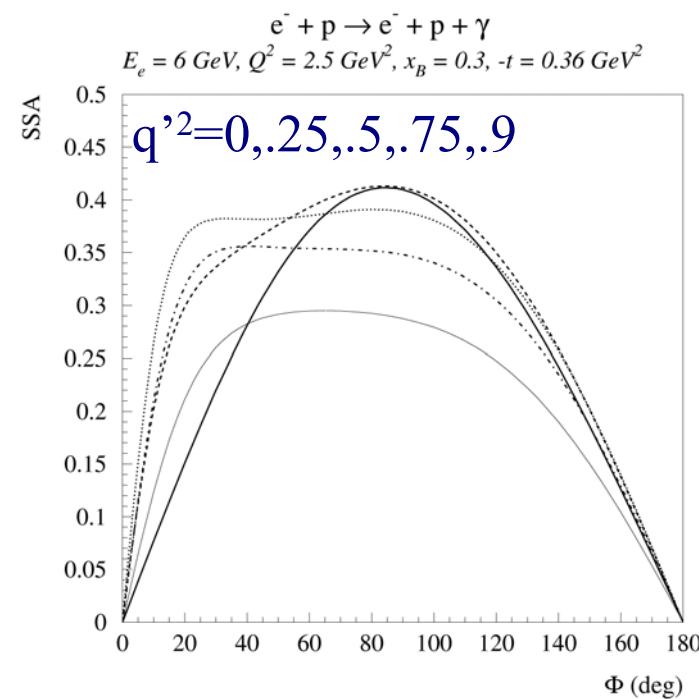
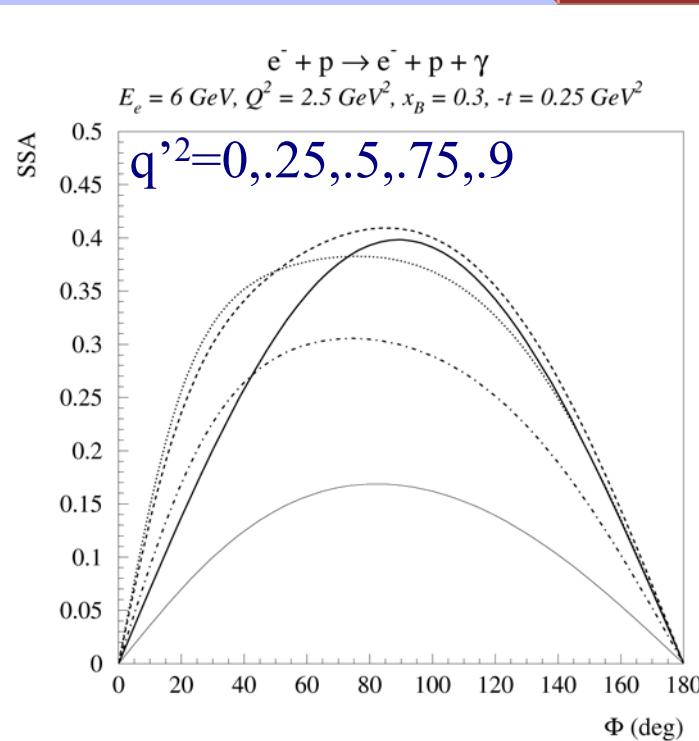
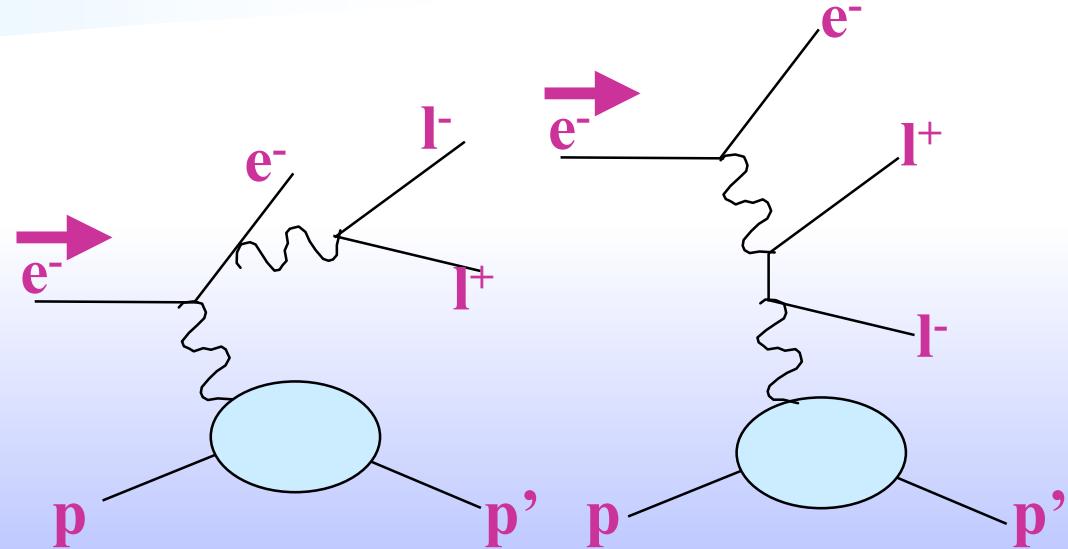
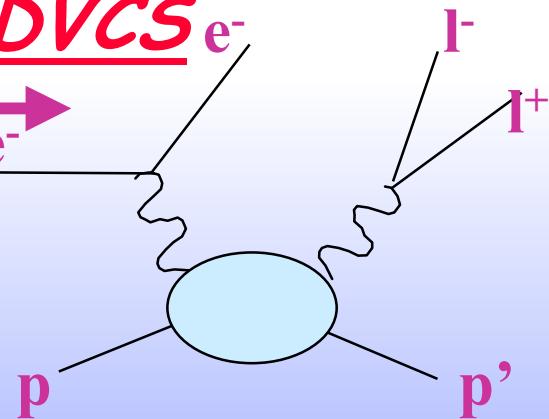


Beam Spin asymmetry for

ep ep γ^* l⁺l⁻

Bethe-Heitler

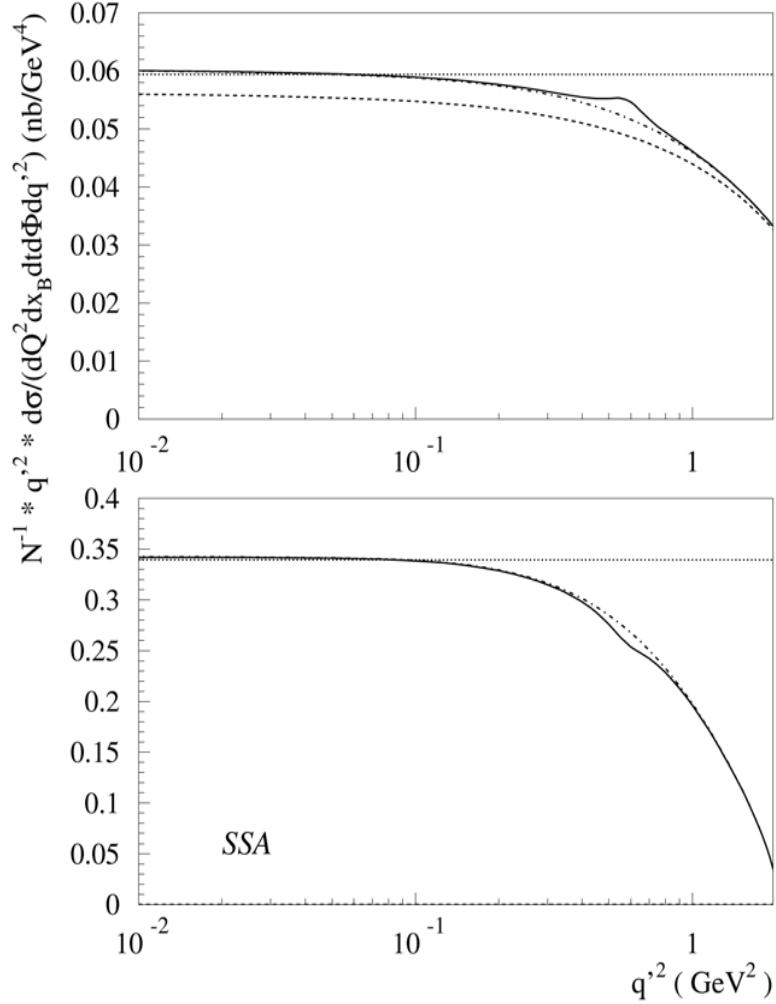
DDVCS e⁻



ρ influence

ep ep ρ $l^+ l^-$

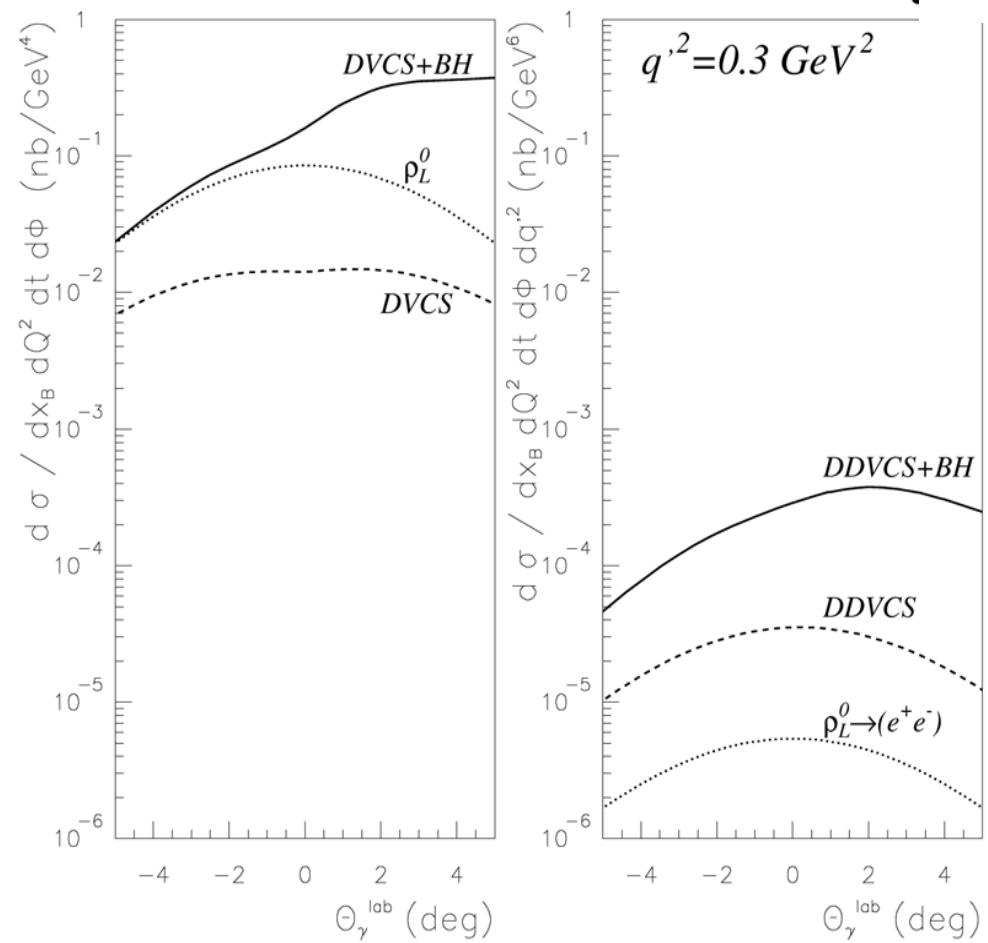
$$E_e = 11 \text{ GeV}, Q^2 = 4 \text{ GeV}^2, x_B = 0.25, t = -0.2 \text{ GeV}^2, \Phi = 90^\circ$$



$$E_e = 6 \text{ GeV}, Q^2 = 2.5 \text{ GeV}^2, x_B = 0.3, \Phi = 0 \text{ deg.}$$

$$e^- + p \rightarrow e^- + p + \gamma, \rho_L^0$$

$$e^- + p \rightarrow e^- + p + (\gamma, \rho_L^0)$$

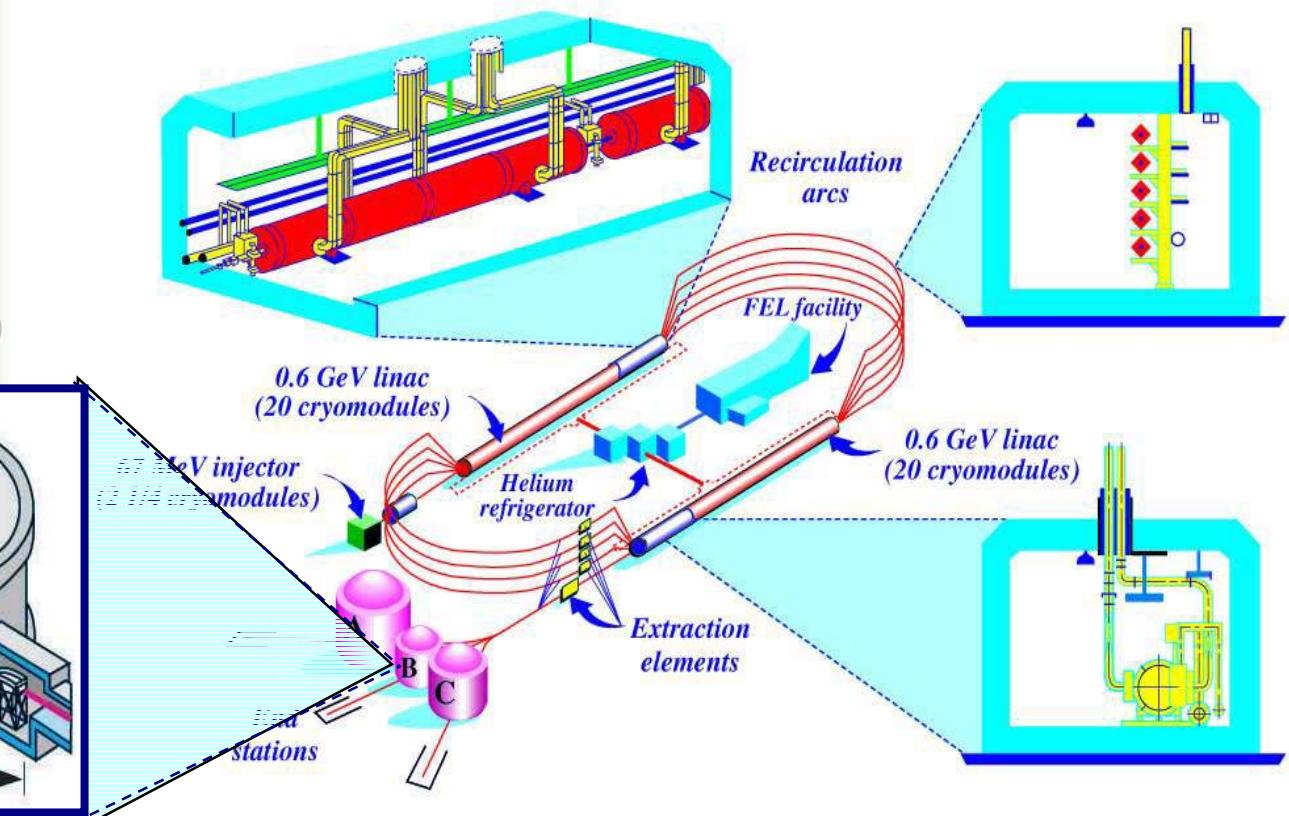
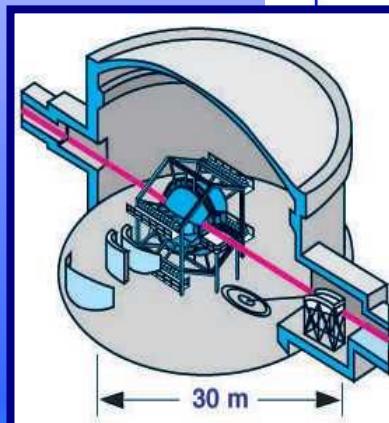


First experimental look

MACHINE CONFIGURATION

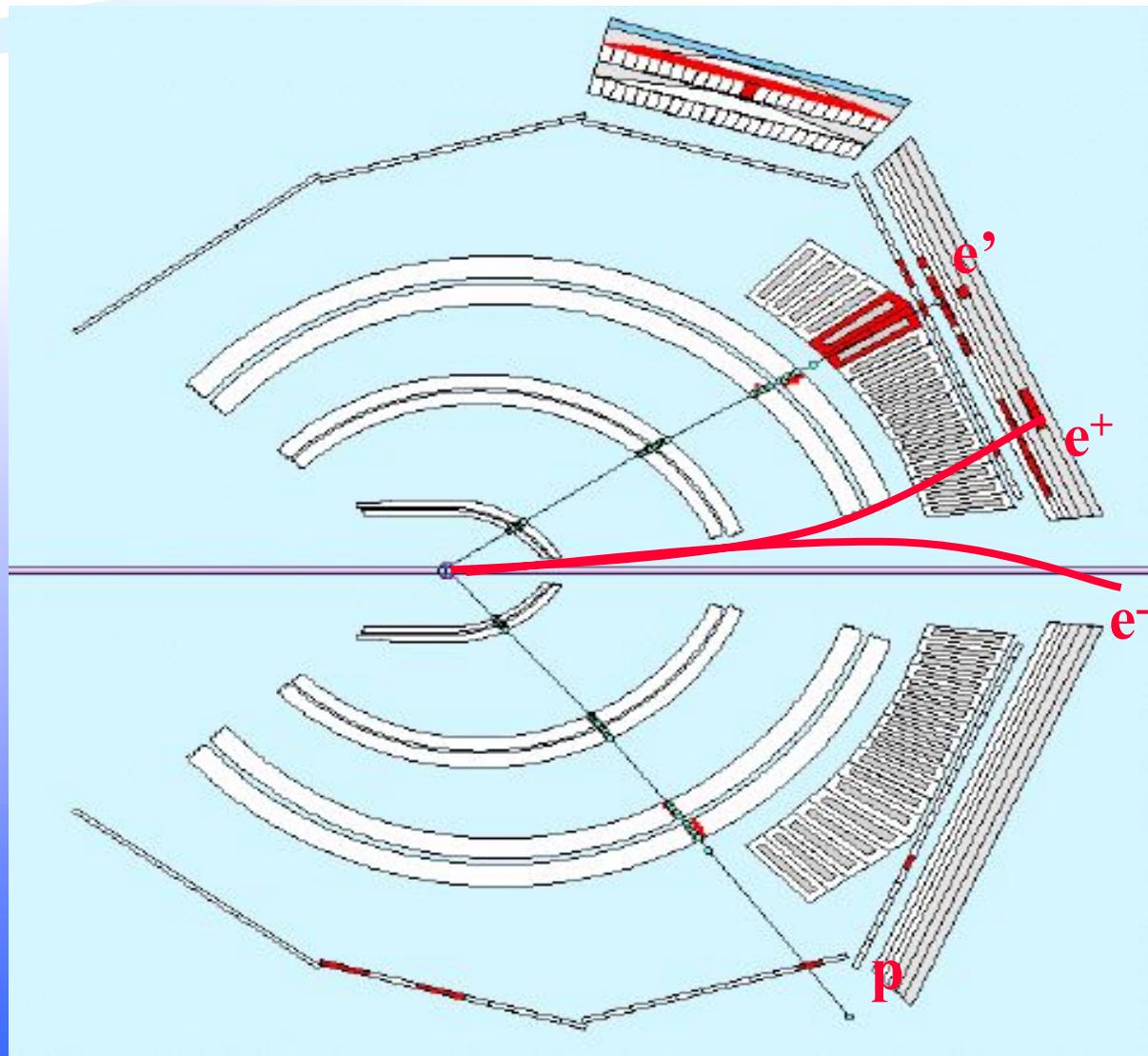


CLAS, JLab



Drawing: Machine Configuration/JLab

**Any
DDVCS
event in
CLAS ?**

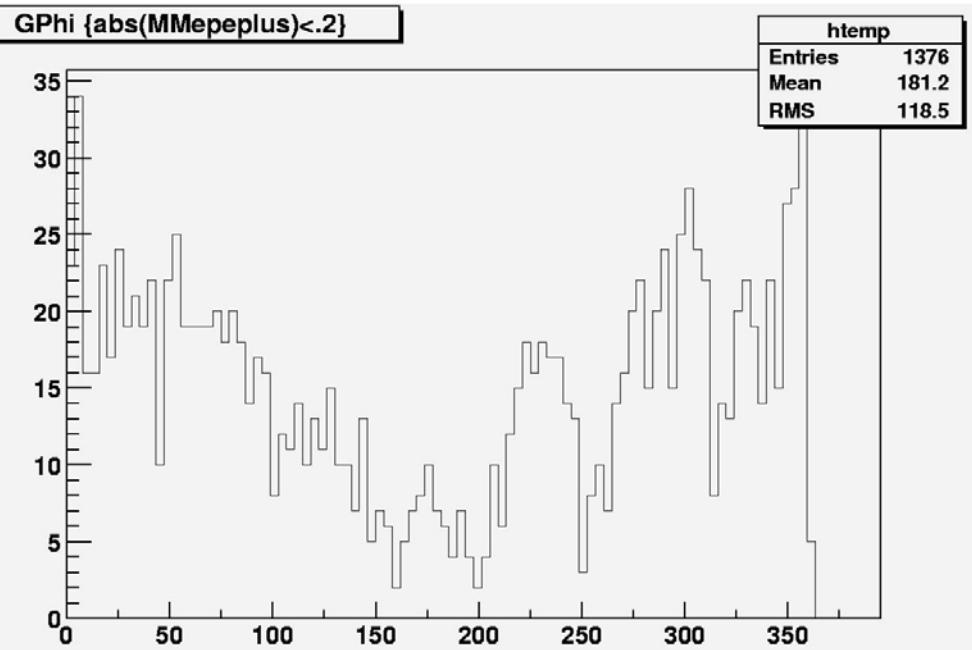
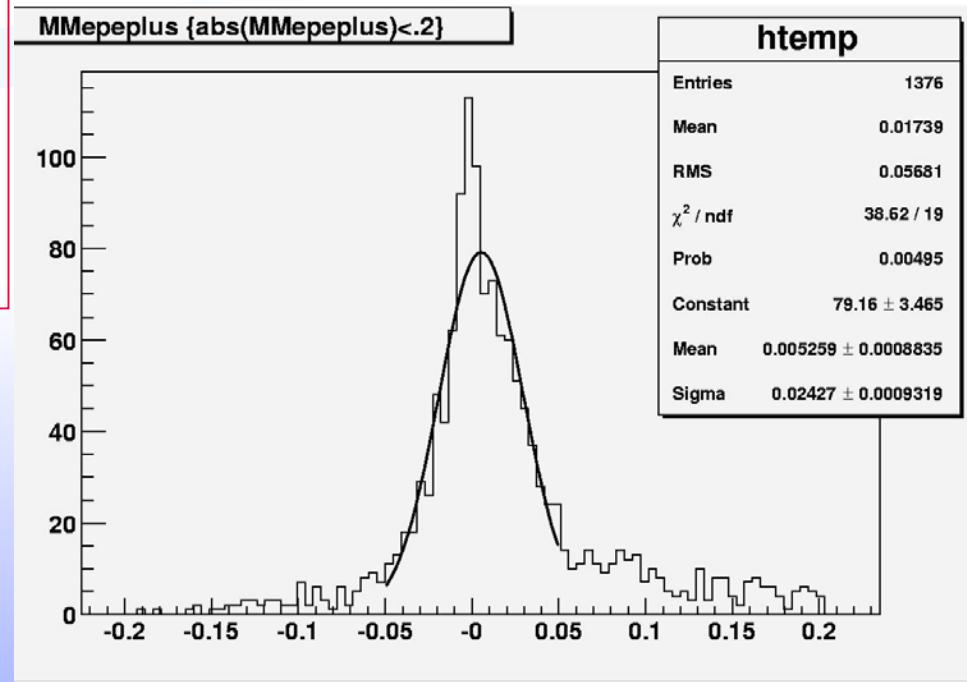


Detect e' , e^+ and p \rightarrow missing e^-

First experimental investigations

CLAS, 30 days beam time,
1/3 statistics, Ee=5.75 GeV

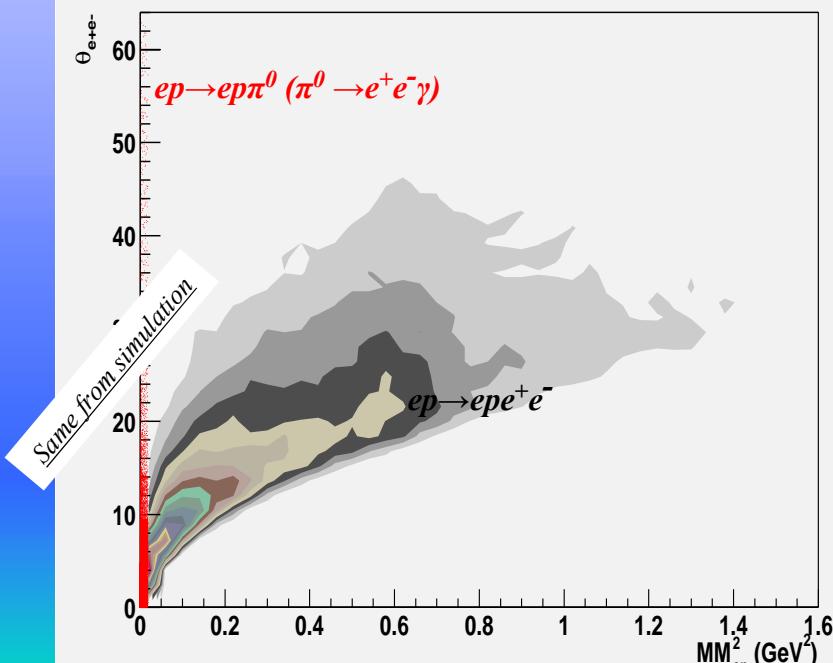
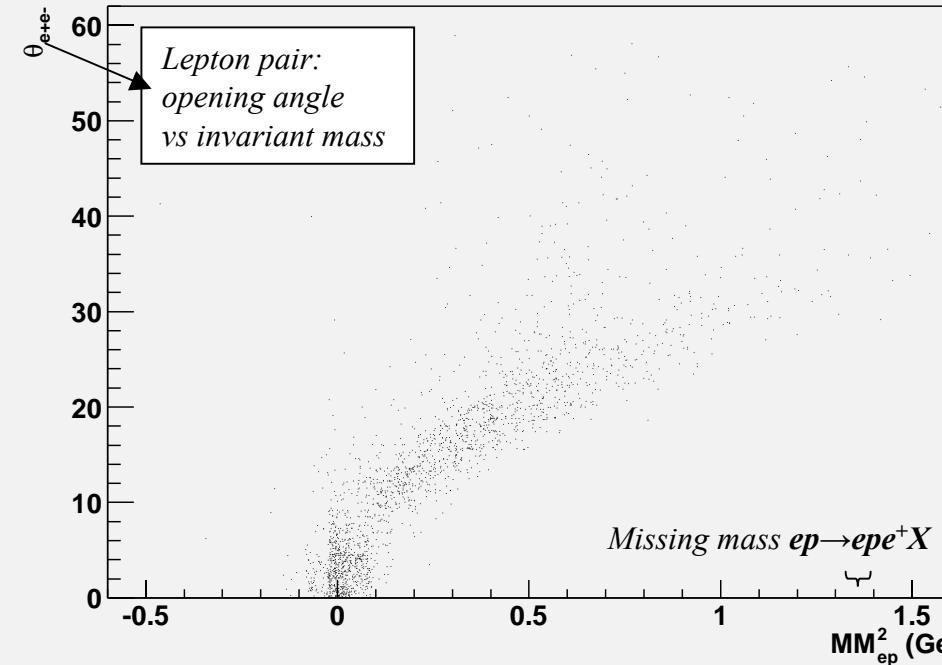
Detect $e^+pe^- \rightarrow$ missing e^-



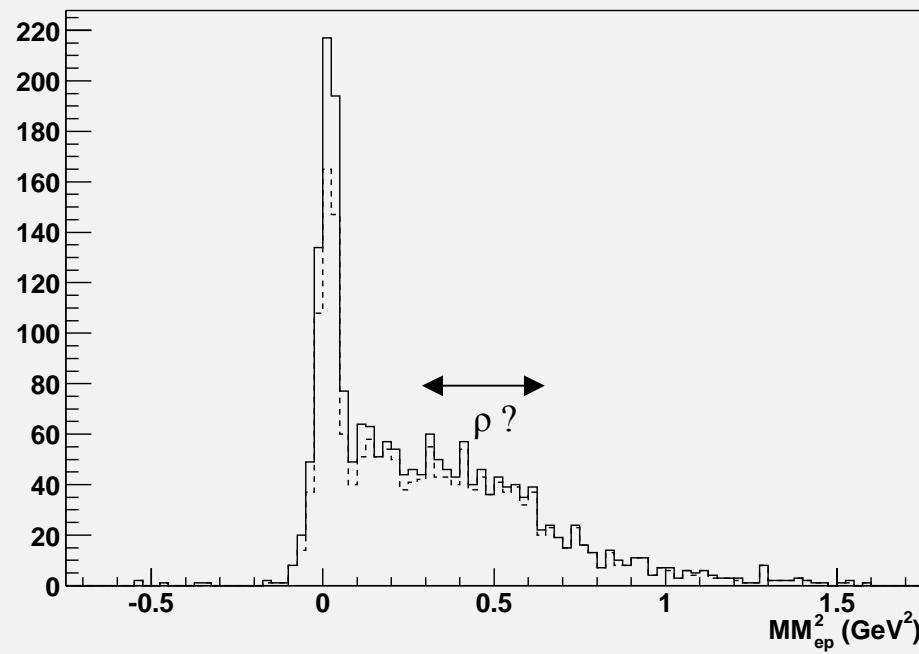
Missing Mass 2 ($e^+pe^- X$) (GeV^2)

Analysis S. Morrow (Saclay)

Φ distribution



Distributions of counts
(not corrected from acceptance)



(q^2')

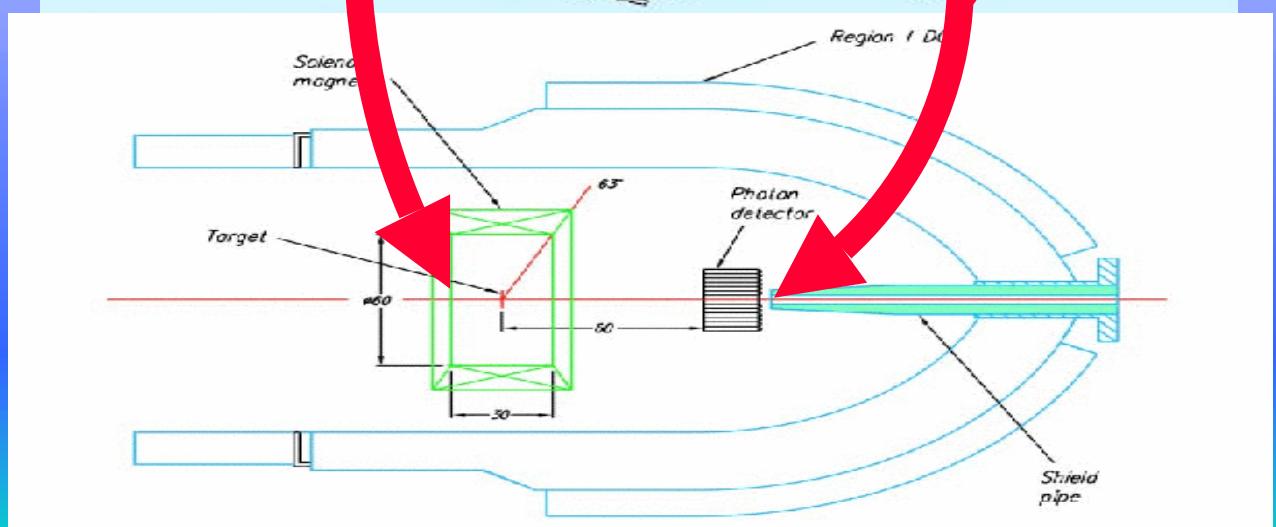
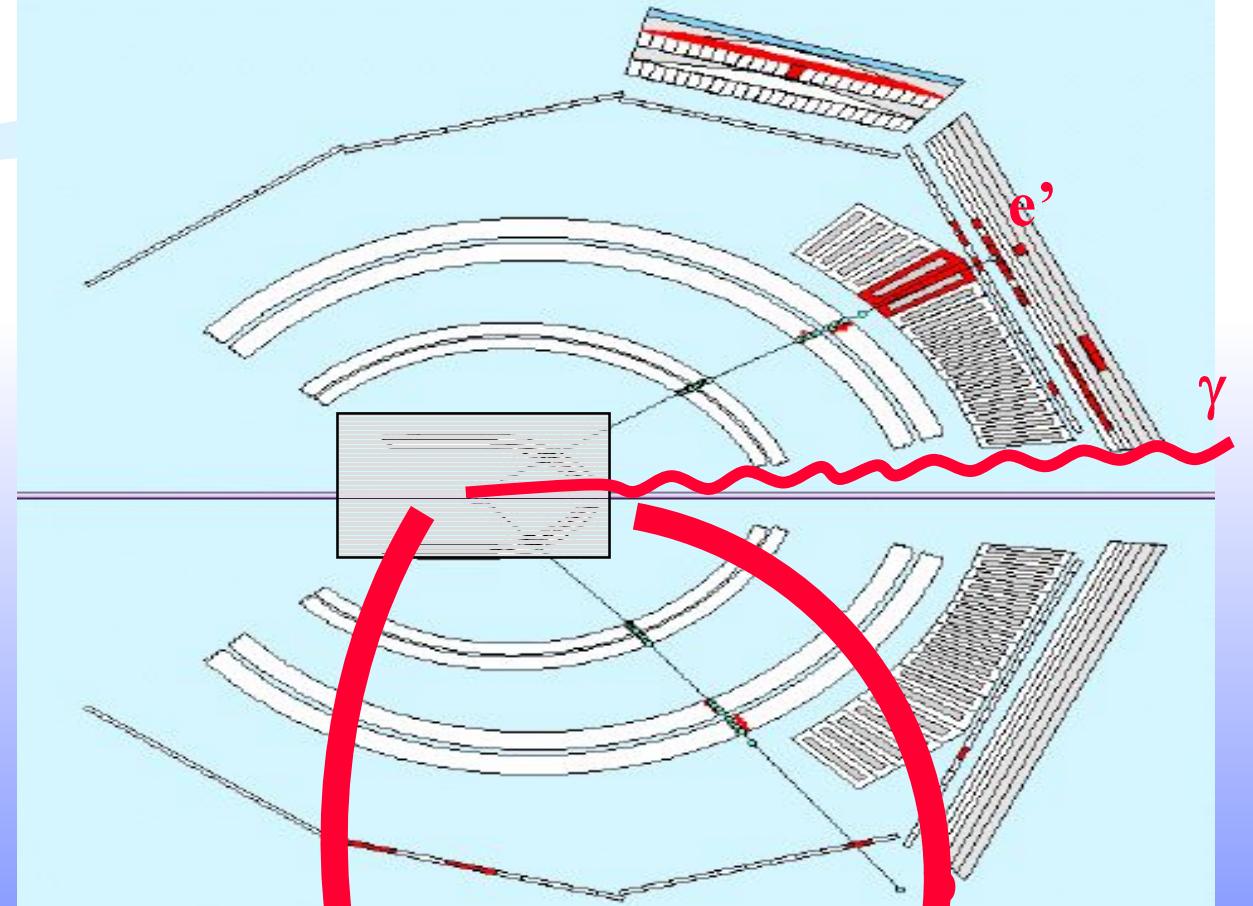
Next step: extract
beam polarization asymmetry

Figures, courtesy of M. Garcon
& S. Morrow (Saclay)

A typical
 $e p \Rightarrow e p \gamma$
event in
CLAS

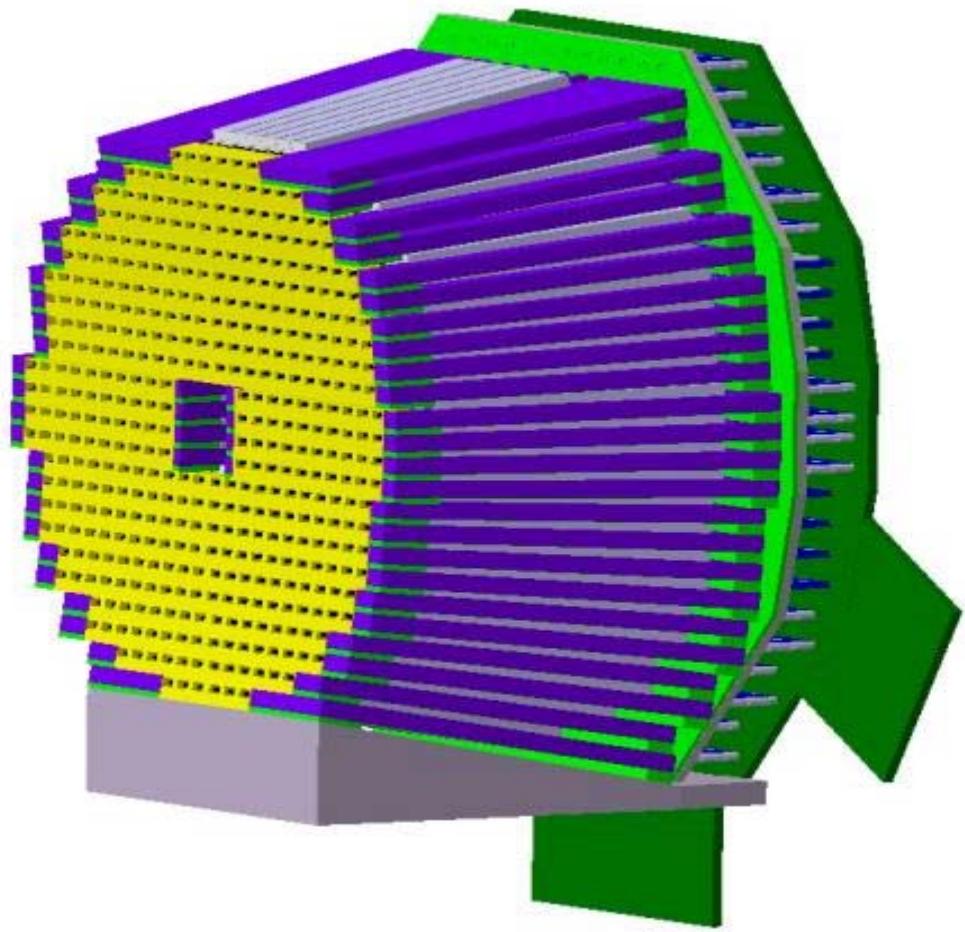
Add EM
calorimeter
at forward
angles

Add solenoid
Moller shield
around target





JLab/ITEP/ Orsay/Saclay/ UVA collaboration

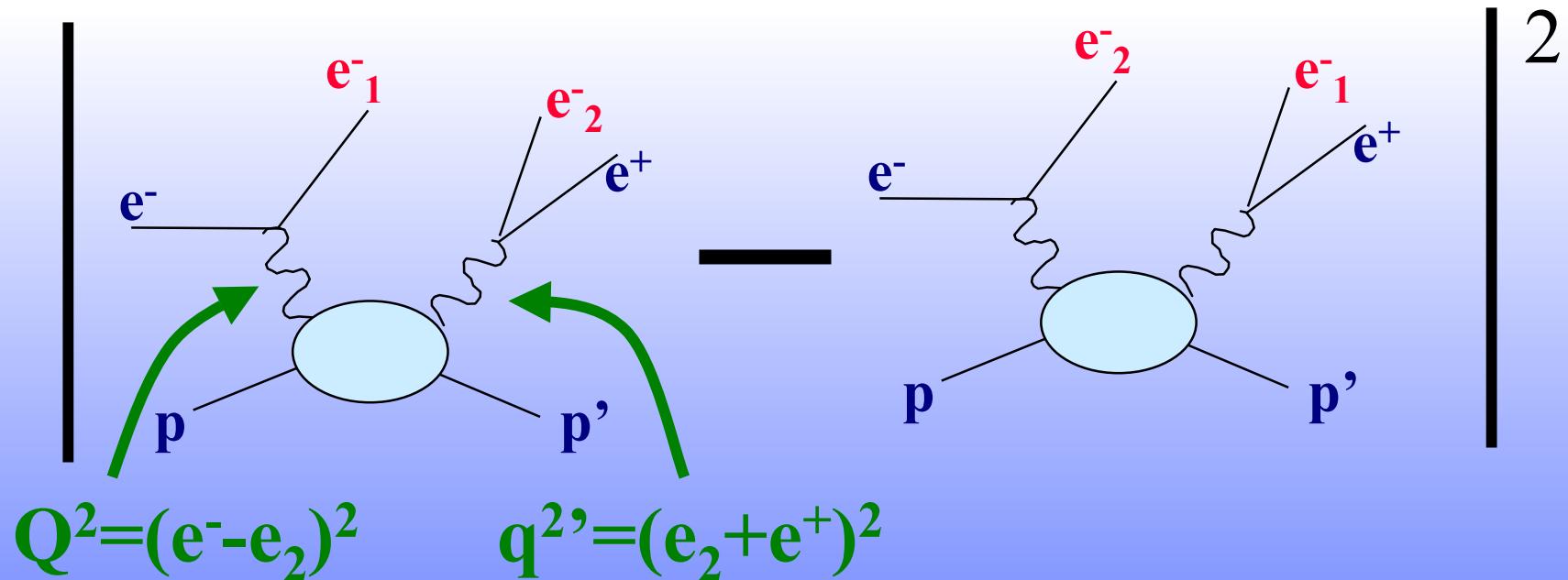


ROSIE

- 😊 Dynamical range : $50 \text{ MeV} < E_\gamma < 5 \text{ GeV}$ ($\sigma \sim 5\%/\sqrt{E_\gamma}$)
- Counting rates $\sim 1 \text{ MHz}$
- Magnetic field environment : $B \sim 1 \text{ T}$

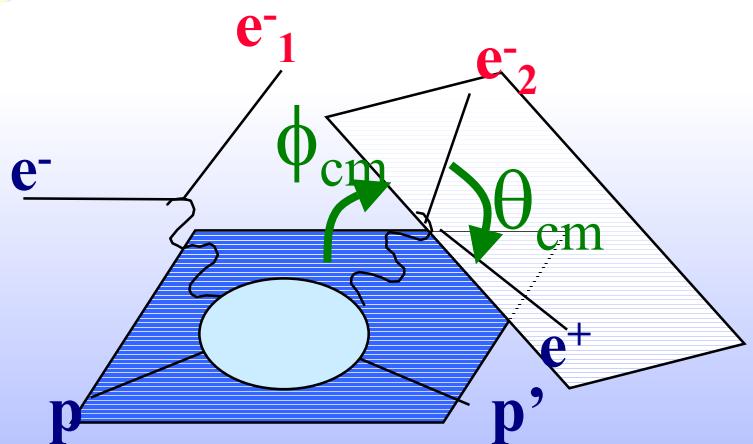
- ➡ ~400 PbWO₄ crystals : $\sim 10 \times 10 \text{ mm}^2$, $l = 160 \text{ mm}$ (18 λ 's)
Read-out : APDs + preamps

The particular case of e^-/e^+ pairs



- 1/ Theoretically, one has to anti-symmetrize
- 2/ Experimentally, ambiguity on defining $Q^2, q^{2'}, \dots$

A bit of kinematics



2 possibilities :

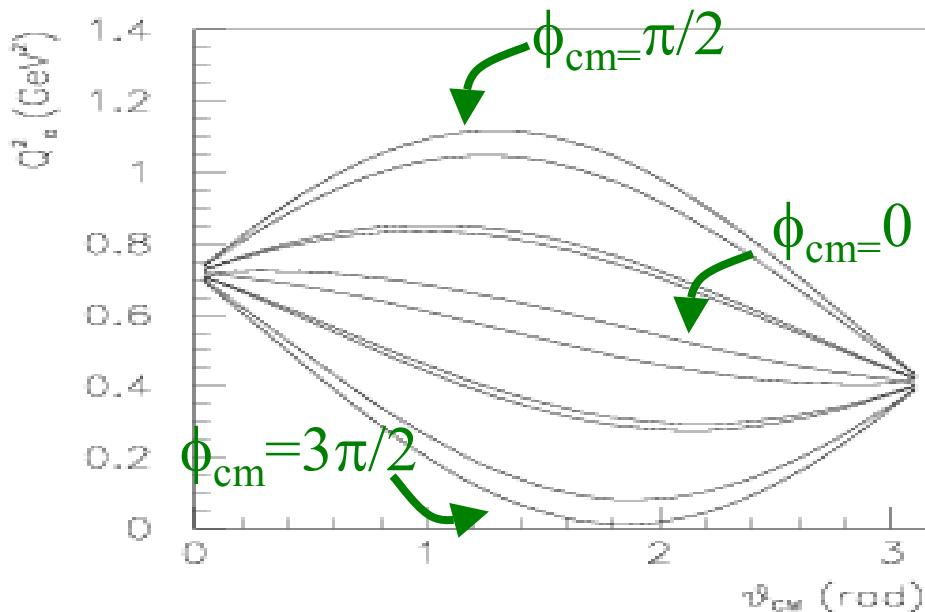
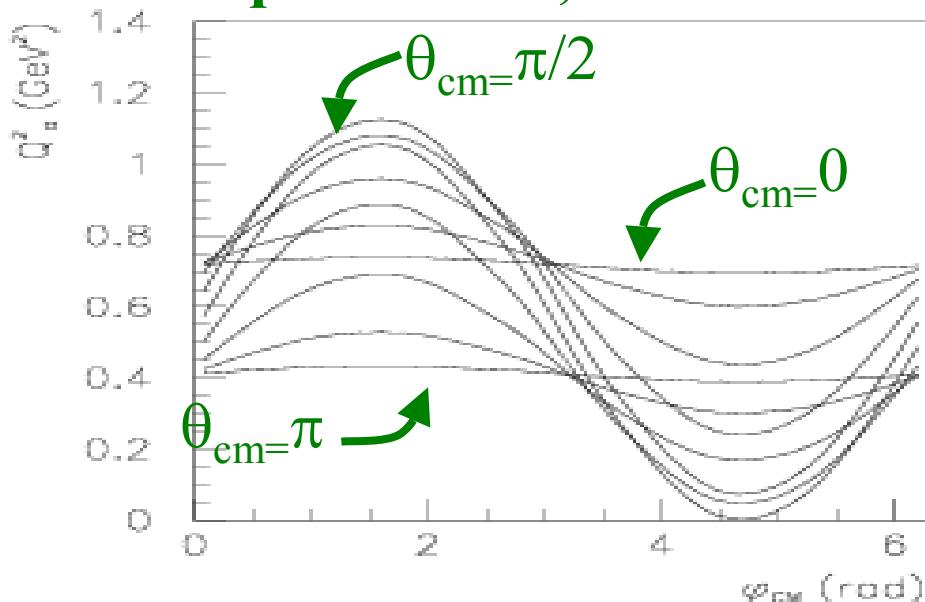
$$Q^2 = (e - e_1)^2$$

$$q'^2 = (e_2 + e^+)^2$$

$$Q^2_a = (e^- - e_1)^2$$

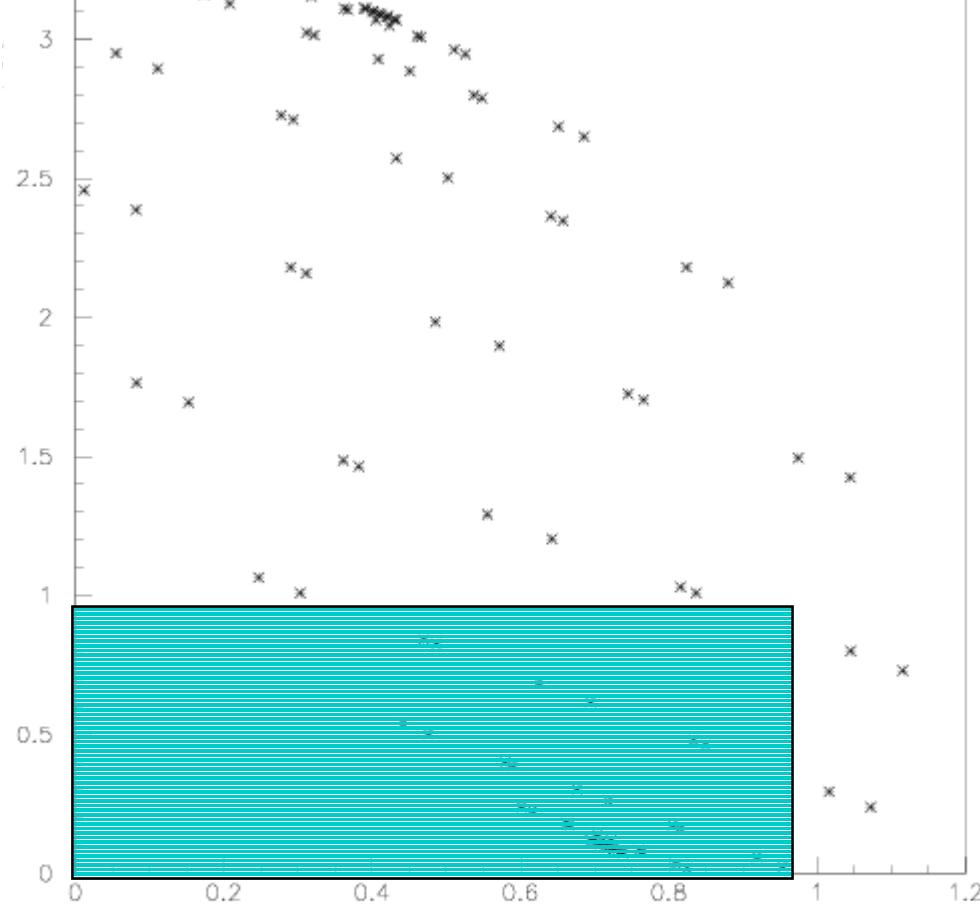
$$q'^2_a = (e_1 + e^+)^2$$

$E_e = 6 \text{ GeV}, Q^2 = 2.5 \text{ GeV}^2,$
 $q'^2 = .3 \text{ GeV}^2, t = -.14 \text{ GeV}^2$



$$q'^2_a = (e_1 + e^+)^2$$

$E_e = 6 \text{ GeV}, Q^2 = 2.5 \text{ GeV}^2, x_B = .3, Q_p = .3 \text{ GeV}^2, t = -.14 \text{ GeV}^2$



$$Q^2_a = (e^- - e_1)^2$$

At least, one virtuality $> 1 \text{ GeV}^2$



2/ Study on the (x,t) correlation of the GPDs

(in coll. with M. VdH & M. Polyakov)

Link between GPDs & FFs

$$F_1^q = \int_{-1}^1 H^q(x, \xi, t) dx \quad F_2^q = \int_{-1}^1 E^q(x, \xi, t) dx$$

Low t ($-t < 1$ GeV 2) : **Regge ansatz** (Goeke, Poliakov, VdH)

$$F_1^u(t) \int_0^1 u_v(x) 1/(x^{\alpha'_1 t}) dx \quad F_1^d(t) \int_0^1 d_v(x) 1/(x^{\alpha'_2 t}) dx$$

$$\mathbf{F_1^p} = e_u F_1^u(t) + e_d F_1^d(t) \quad \mathbf{F_1^n} = e_u F_1^d(t) + e_d F_1^u(t)$$

Similarly $F_2^q(t) = \int_0^1 \kappa_v q_v(x) 1/(x^{\alpha'_2 t}) dx$

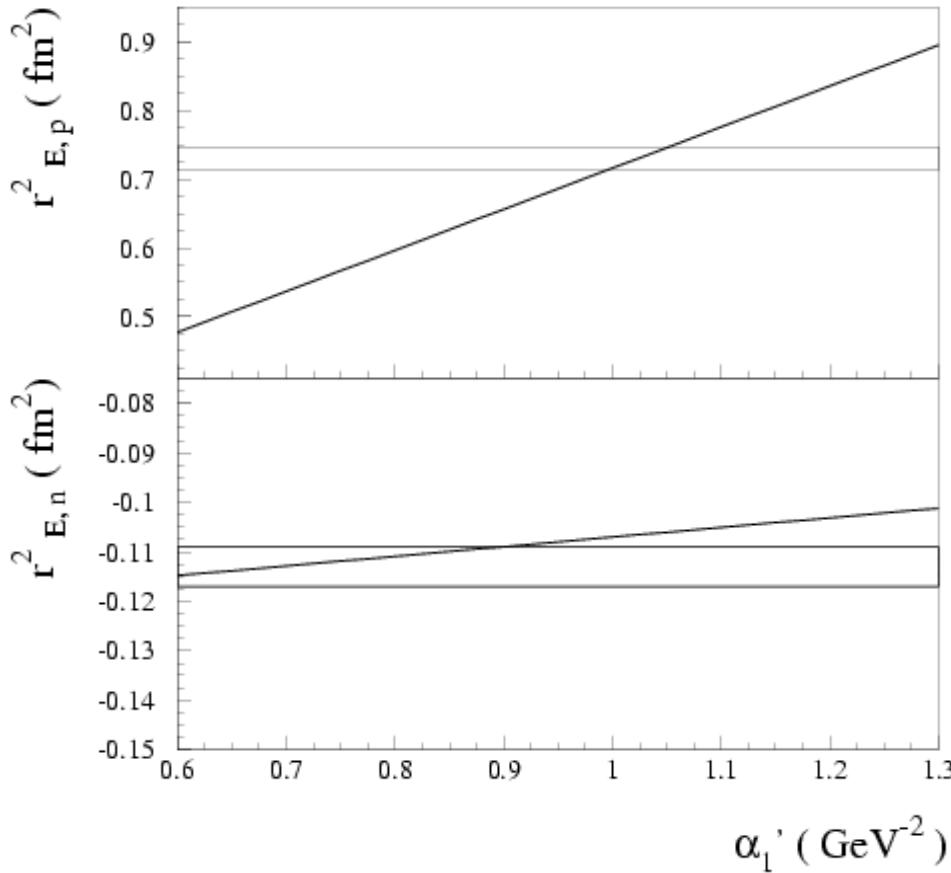
→ 2 free parameters : α'_1, α'_2 (constrained by Regge)

→ Fit 4 form factors : $G_{E,M}^{p,n}$

Proton & neutron electric charge radii

$$F_1^u = \int_0^1 u_v(x) 1/(x^{\alpha'_1} t) dx \quad \Rightarrow \quad r^2_{1,p} = -6\alpha'_1 \ln x (e_u u_v + e_d d_v) dx$$

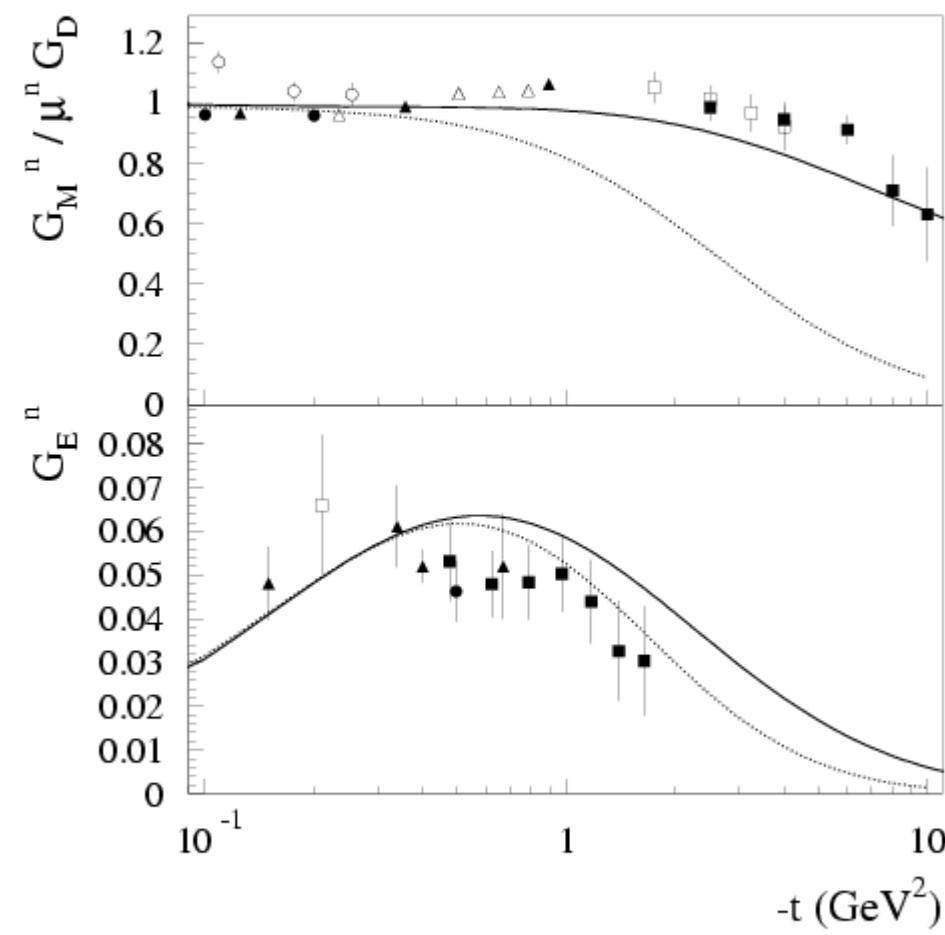
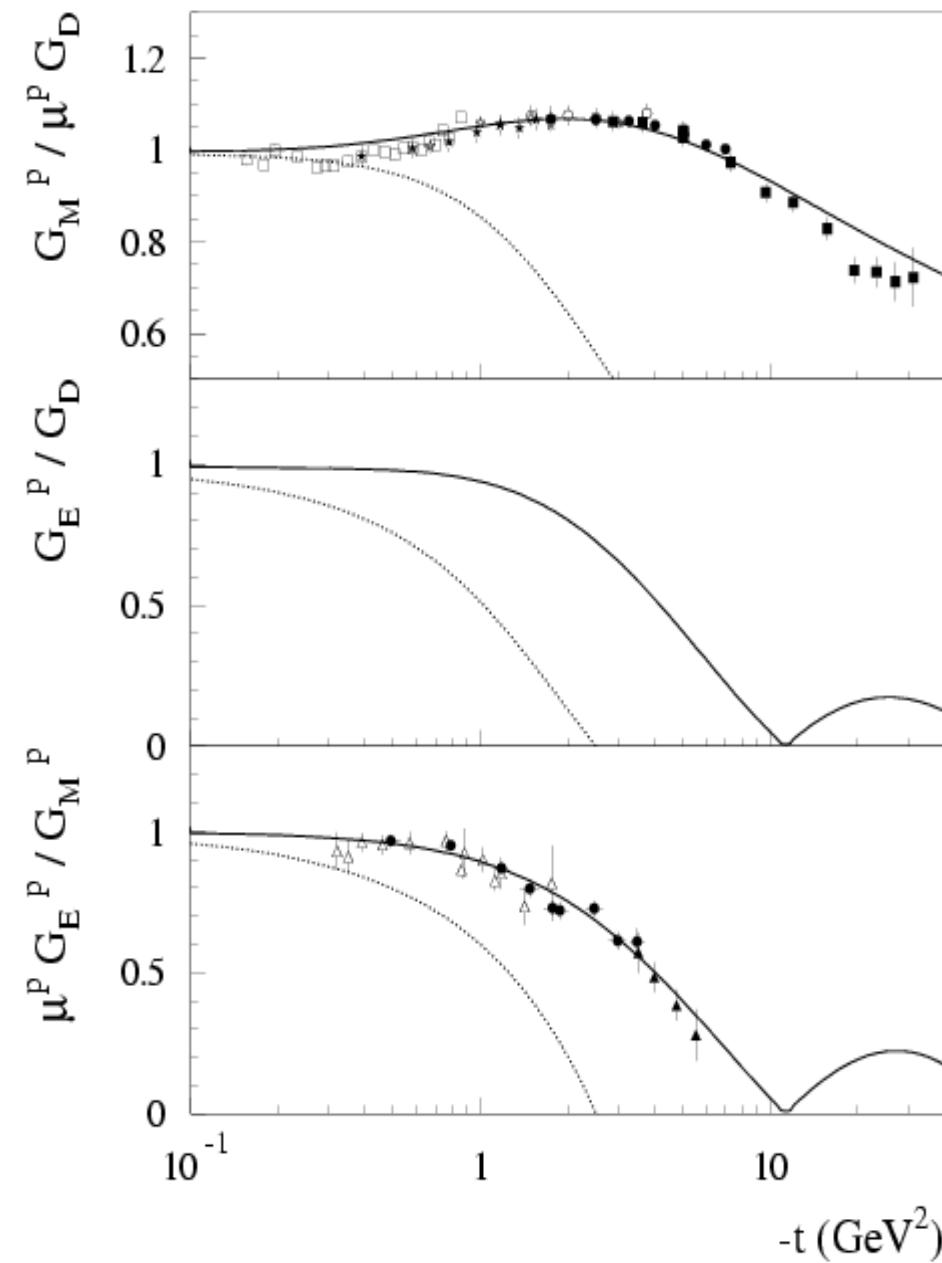
$$r^2_{E,N} = r^2_{1,N} + 3/2 \kappa_N / m_N^2$$



$\Rightarrow \alpha'_1 \sim 1 \text{ GeV}^{-2}$

PROTON electric and magnetic form factors

NEUTRON electric and magnetic form factors



Extending the t domain

pQCD : $F_1(t) \sim 1/t^2$ at large t ($F_2 \sim 1/t^3$)

1/ Introduce non-linear Regge trajectories :

$$\alpha(t) = \alpha(0) + \alpha' 2T [1 - \sqrt{1 - t/T}]$$

(Brisudova, Burakovsky, Goldman)

T : (free) non-linearity parameter



2/ a/ x-dependence of Eq is not constrained

b/ Large t power behavior is fixed by large x behavior

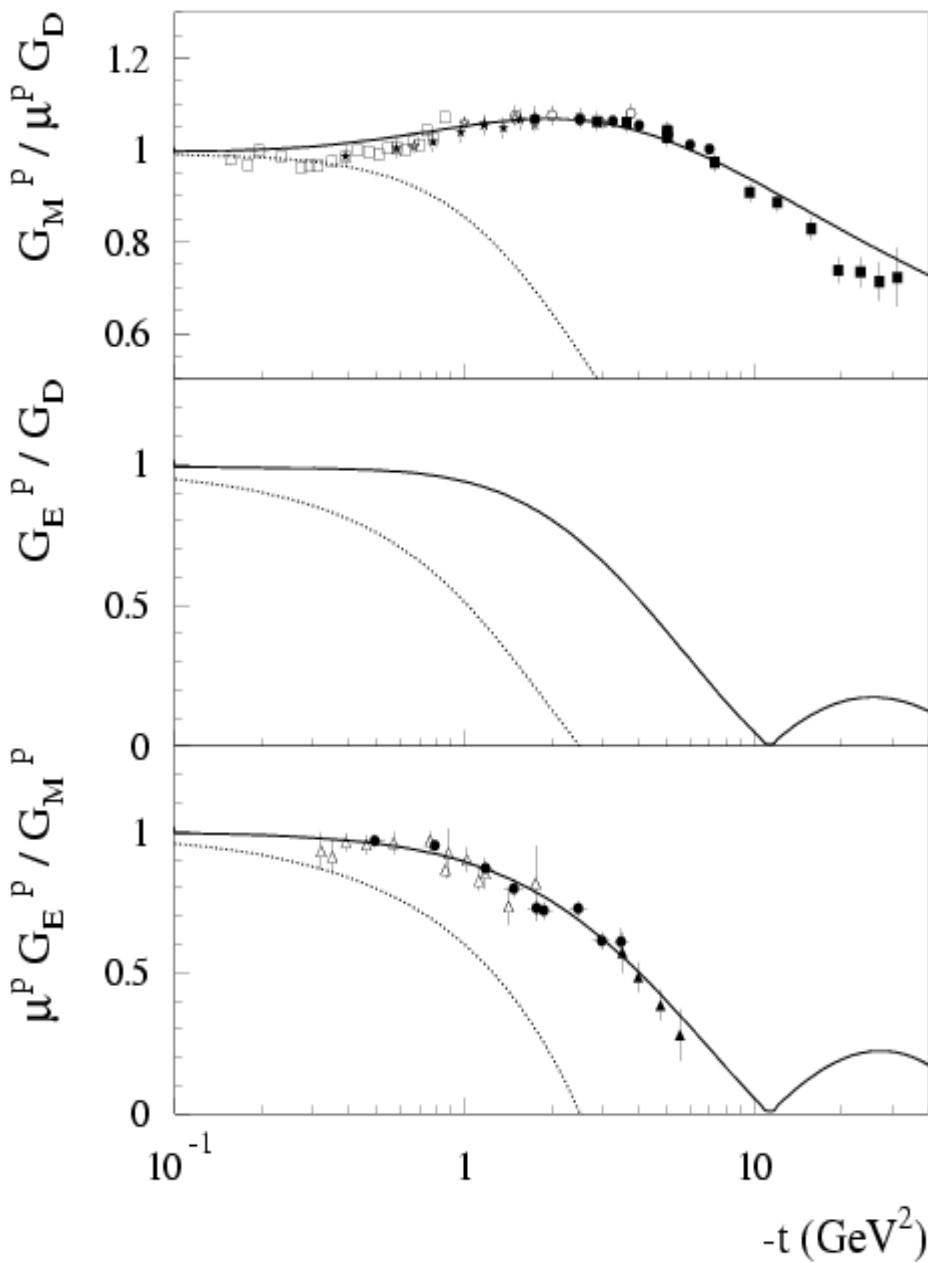
→ Large x behavior of E should be different from H

$$F_2^q(t) = \int_0^1 \kappa_v (1-x)^{\eta q} q_v(x) / (x^{\alpha' 2^t}) dx$$

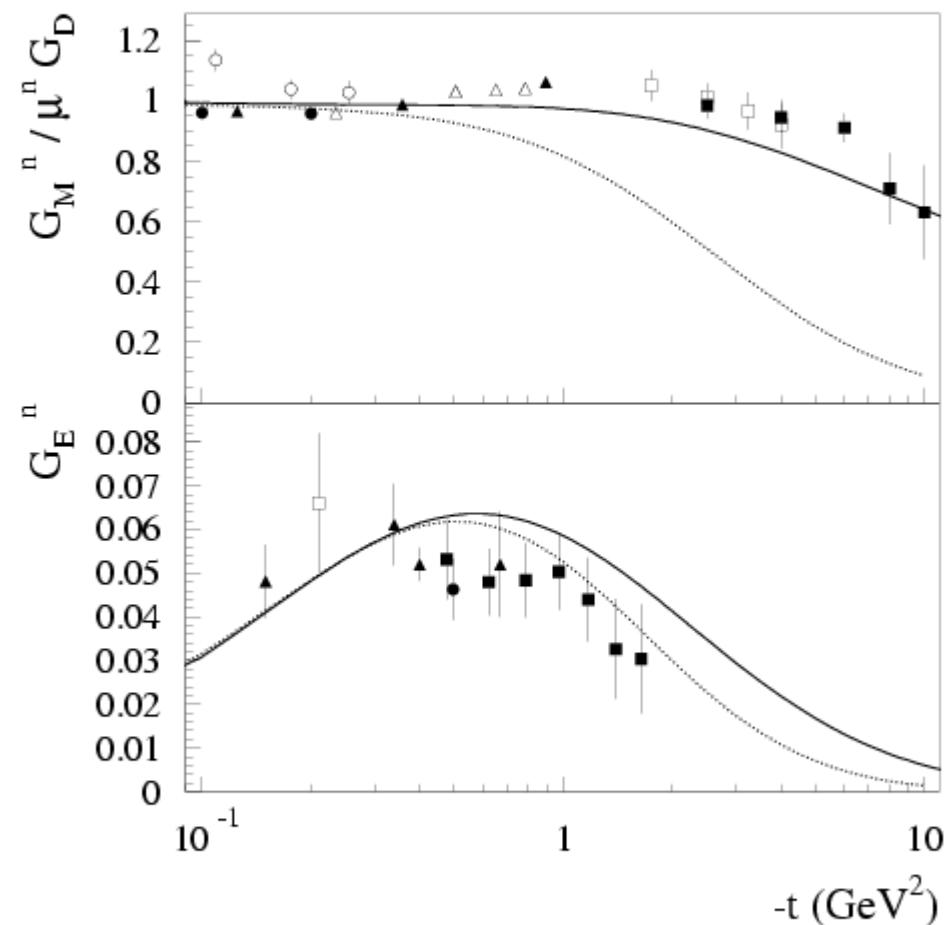
→ 6 parameters : $\alpha'_1, \alpha'_2, \eta^u, \eta^d, T_1, T_2,$

→ Fit 4 form factors : $G_{E,M}^{p,n}$ on whole t domain !

PROTON electric and magnetic form factors



NEUTRON electric and magnetic form factors

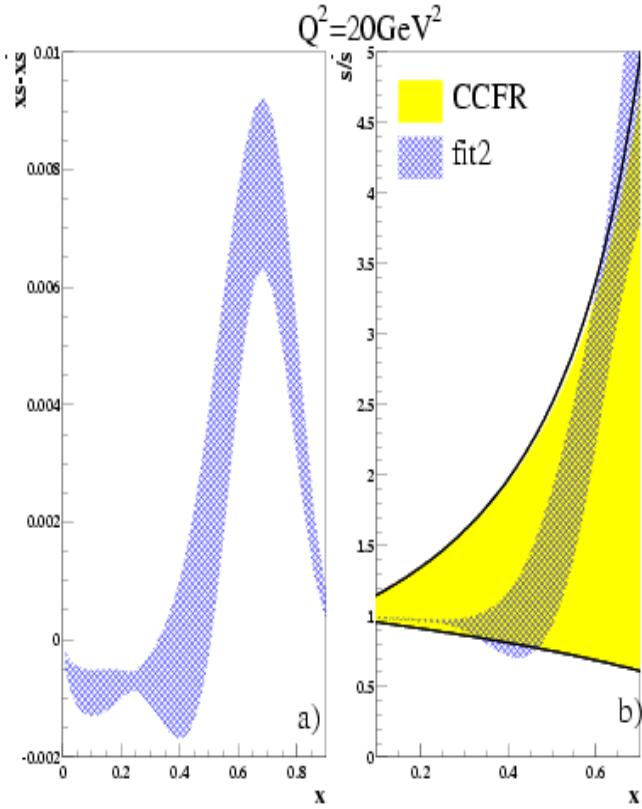


Nucleon strangeness : F_1^s

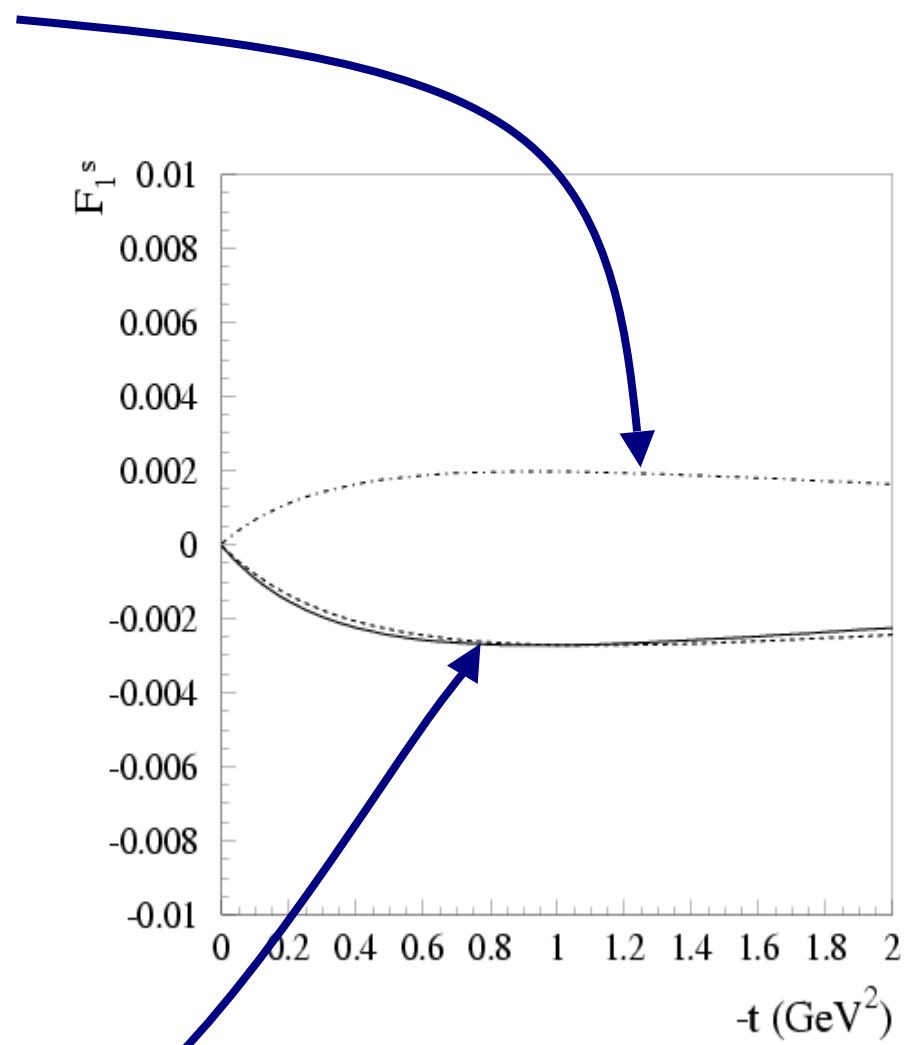
$$\int_{-1}^1 [s(x) - \bar{s}(x)] dx = 0$$

But : $F_1^s(t) = \int_{-1}^1 [s(x) - \bar{s}(x)] / (x^{\alpha \cdot t}) dx \neq 0$

- extracted from $\nu, \bar{\nu}$ data
- DIS fit performed by Barone, Pascaud, Zomer (2000)
fit $s \neq \bar{s}$ (CDHS data) : $\langle xs \rangle - \langle x\bar{s} \rangle \simeq +2 \cdot 10^{-3}$



- fit based on $\nu, \bar{\nu}$ CCFR + NuTeV (2001) data :
 $\langle xs \rangle - \langle x\bar{s} \rangle \simeq -(2.7 \pm 1.3) \cdot 10^{-3}$



SUMMARY

1/DDVCS

- ➡ The (continuously) varying virtuality of the outgoing photon allows to “tune” the kinematical point (x, ξ, t) at which the GPDs are sampled
- ➡ Currently working out the anti-symmetrisation issues
- ➡ Experimental feasibility : maybe not so far away....

2/ (x,t) dep. of GPDs

- ➡ α' should depend on Q^2 to match evolution of PDF/GPDs