# Prospects for a Search for Sterile Neutrinos Using the Sudbury Neutrino Observatory

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REU Presentations Presents "Memories of the Summer", appearing on the bottom of this presentation.

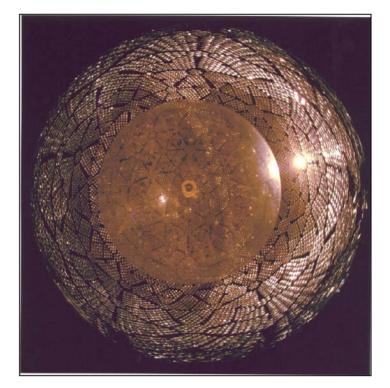
## Outline

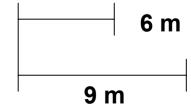
## Goal: To determine whether or not it is possible to use SNO to search for sterile neutrinos

- Sudbury Neutrino Observatory
- Neutrino Oscillations
- Sterile Neutrinos
- Detector Yield Calculations
- Day-Night Asymmetry
- Results
- Conclusions

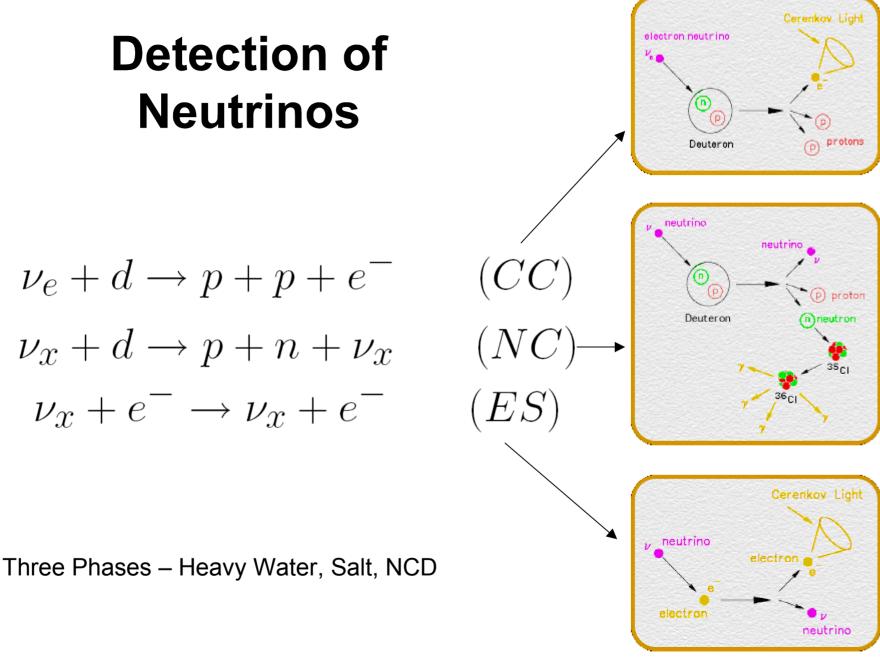


Sudbury Neutrino Observatory





- Heavy Water Čerenkov Detector
- 2 km underground, INCO Ltd. Creighton mine
- 12 m diameter AV, filled with heavy water
- Light water extending from the AV to the edges of the SNO cavity
- •18 m diameter PMT support structure



"CORE DUMPED!"

## **Neutrino Oscillations**

#### Equations? AHHHHH!

- Three flavors of neutrinos, called the "weak states"  $v_e^{},\,v_\mu^{},\,v_\tau^{}$
- Massless in the Standard Model
- "Mass states" v<sub>1</sub>, v<sub>2</sub>, v<sub>3</sub>. Weak states are a linear combination of Mass States:

$$|\nu_x\rangle = \sum_m U_{xm} |\nu_m\rangle$$

- U is the neutrino mixing matrix, similar to the CKM matrix for quark mixing
- Assume only two neutrino generations:

$$|\nu_e\rangle = \cos\theta |\nu_1\rangle + \sin\theta |\nu_2\rangle$$
$$|\nu_\mu\rangle = -\sin\theta |\nu_1\rangle + \cos\theta |\nu_2\rangle$$

## Neutrino Oscillations, Cont.

Otherwise known as "Oh god, not more equations!"

• Time evolution of a state is determined by its energy:

 $|\nu_e(t)\rangle = \cos\theta \exp[-i(E_1/\hbar)t]|\nu_1\rangle + \sin\theta \exp[-i(E_2/\hbar)t]|\nu_2\rangle$ 

- -Energy of a state is related to its mass :  $E_k = \sqrt{p^2 c^2 + m_k^2 c^4}$
- Starting from the fact that the probability an electron neutrino will oscillate into a muon neutrino is given by the absolute square of the amplitude of  $\langle \nu_{\mu} | \nu_{e}(t) \rangle$ :

$$P(\nu_e \to \nu_\mu) = 1 - \sin^2 2\theta \sin^2(\frac{1.27\Delta m^2 L}{E_\nu})$$

$$P(\nu_e \to \nu_e) = 1 - P(\nu_e \to \nu_\mu)$$

Neutrino oscillation can be understood in a very simple way. What happens is that the lighter mass states in the original  $v_x$  travel faster than the heavier ones, and get ahead of the latter. Thus, the various  $v_m$  components of the beam get out of phase with one another, and do not add up to a  $v_x$  anymore. Thus, as it travels, the beam picks up components corresponding other flavors.

- B. Kayser, "The Physics of Massive Neutrinos"

"Accessing Random Insult Generator... 'You suck at life.' "

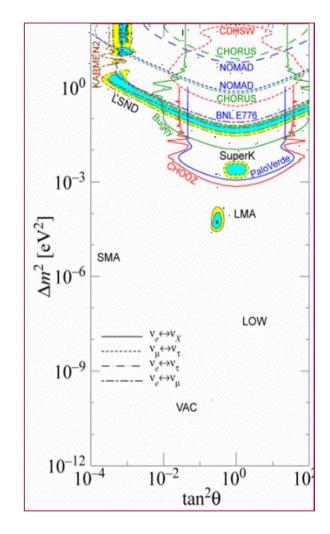
## **Sterile Neutrinos**

#### Why am I wasting your time talking about them?

- Three neutrinos -> Three Mass Scales
- Three different types of experiments measure the three different mass scales
- Experimentally established limits:

 $\Delta m_{12}^2 < 10^{-5}$  $\Delta m_{23}^2 < 10^{-3}$  $\Delta m_{13}^2 > 10^{-1}$ 

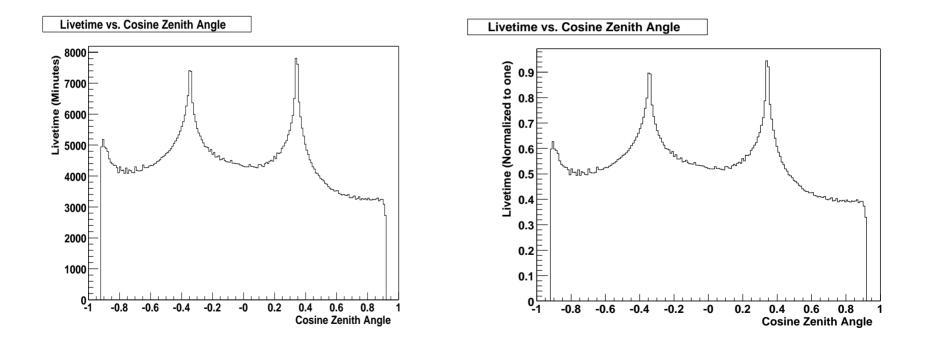
- $\Delta m_{12}^2$  +  $\Delta m_{23}^2$  =  $\Delta m_{13}^2$ . Doesn't add up!
- Need at least one more neutrino
- LEP says there are only three active neutrinos
- Any additional neutrinos must not interact with the Z, must be sterile



### **Detector Yield**

In other words, just how many of these neutrinos will we see?

$$Y = U\epsilon_n \int dt \int_0^V dV \int_{T_{th}}^\infty R_{NC}(T) dT \int_0^\infty \int_{-1}^1 W(\cos\theta_z) \Phi_{SSM}(E_\nu) \sigma_{\nu d}^{NC}(E_\nu) P(\nu_a \to \nu_s) d\cos\theta_z dE_\nu$$



### **Detector Yield**

#### **Still Explaining the Big Equation**

$$Y = U\epsilon_n \int dt \int_0^V dV \int_{T_{th}}^\infty R_{NC}(T) dT \int_0^\infty \int_{-1}^1 W(\cos \theta_z) \Phi_{SSM}(E_\nu) \sigma_{\nu d}^{NC}(E_\nu) P(\nu_a \to \nu_s) d\cos \theta_z dE_\nu$$

"Who is your Daddy, and what does he do?"

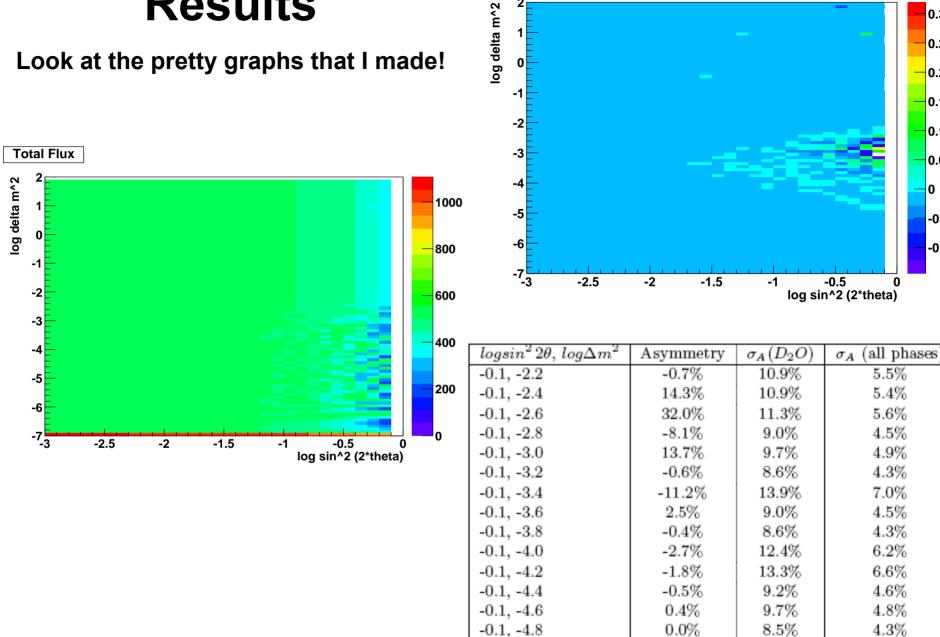
## Day – Night Asymmetry

How we actually "see" the sterile neutrinos in SNO

$$N = U\epsilon_n \int dt \int_0^V dV \int_{T_{th}}^\infty R_{NC}(T) dT \int_0^\infty \int_{-1}^0 W(\cos\theta_z) \Phi_{SSM}(E_\nu) \sigma_{\nu d}^{NC}(E_\nu) P(\nu_a \to \nu_s) d\cos\theta_z dE_\nu$$
$$D = U\epsilon_n \int dt \int_0^V dV \int_{T_{th}}^\infty R_{NC}(T) dT \int_0^\infty \int_0^1 W(\cos\theta_z) \Phi_{SSM}(E_\nu) \sigma_{\nu d}^{NC}(E_\nu) P(\nu_a \to \nu_s) d\cos\theta_z dE_\nu$$
$$A = 2 \frac{N - D}{N + D}$$

- Different values of the length of travel of the neutrino during night and day change the probability of oscillating into a sterile state
- Non-zero asymmetries should be observed if sterile neutrinos exist
- Can SNO see them?

### **Results**



-0.1, -5.0

0.1%

9.0%

4.5%

Asymmetry

0.3

0.25

0.2

0.15

0.1

0.05

-0.05

-0.1

0

2

"I'm a cop, you idiot!"

## Conclusions

#### Last Slide. Hooray!

- Limits on the parameters for detecting steriles in SNO.
- Currently, these limits would be:

 $10^{-2} < \sin^2 2\theta$ 

 $10^{-5} < \Delta m^2 < 10^{-2}$ 

• Can sterile neutrinos be detected at all through a day – night asymmetry?

Yes, but whether or not a non-zero asymmetry can be measured to  $3\sigma$  is not yet known

• Much more work is required

Non-constant Earth-Sun Distance

Matter effects

Extension of  $\sigma_A$  to the entire data set is only an approximation. N and D should be higher in second and third phases

Generalize model to include mixing between sterile and all three active neutrinos

Additional sterile neutrinos?

# Acknowledgements

I lied, but this really is the last slide

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