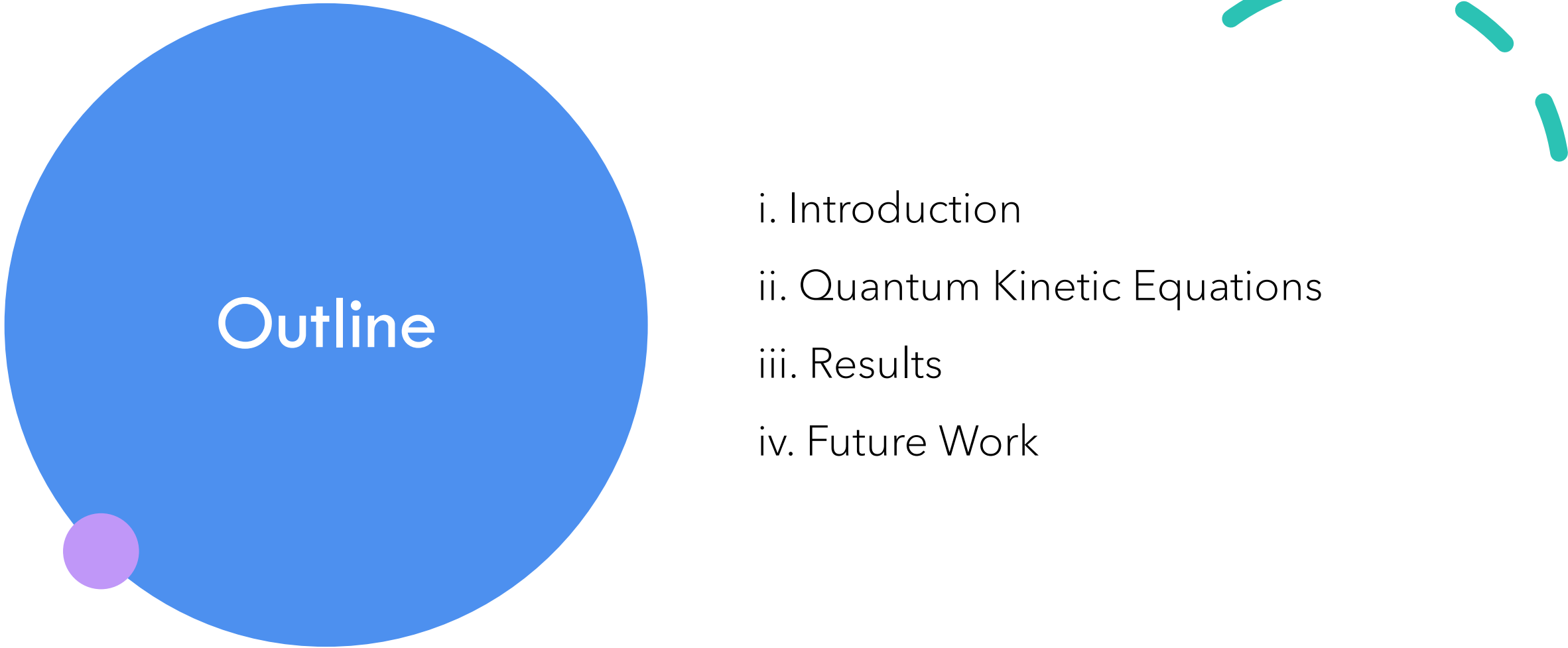




Neutrino Oscillations in Dense Media

Ellen Gates –Department of Physics and Astronomy, University of Missouri
Advisor: Vincenzo Cirigliano –Institute for Nuclear Theory, University of Washington



Outline

- i. Introduction
- ii. Quantum Kinetic Equations
- iii. Results
- iv. Future Work



Outline

i. Introduction

ii. Quantum Kinetic Equations

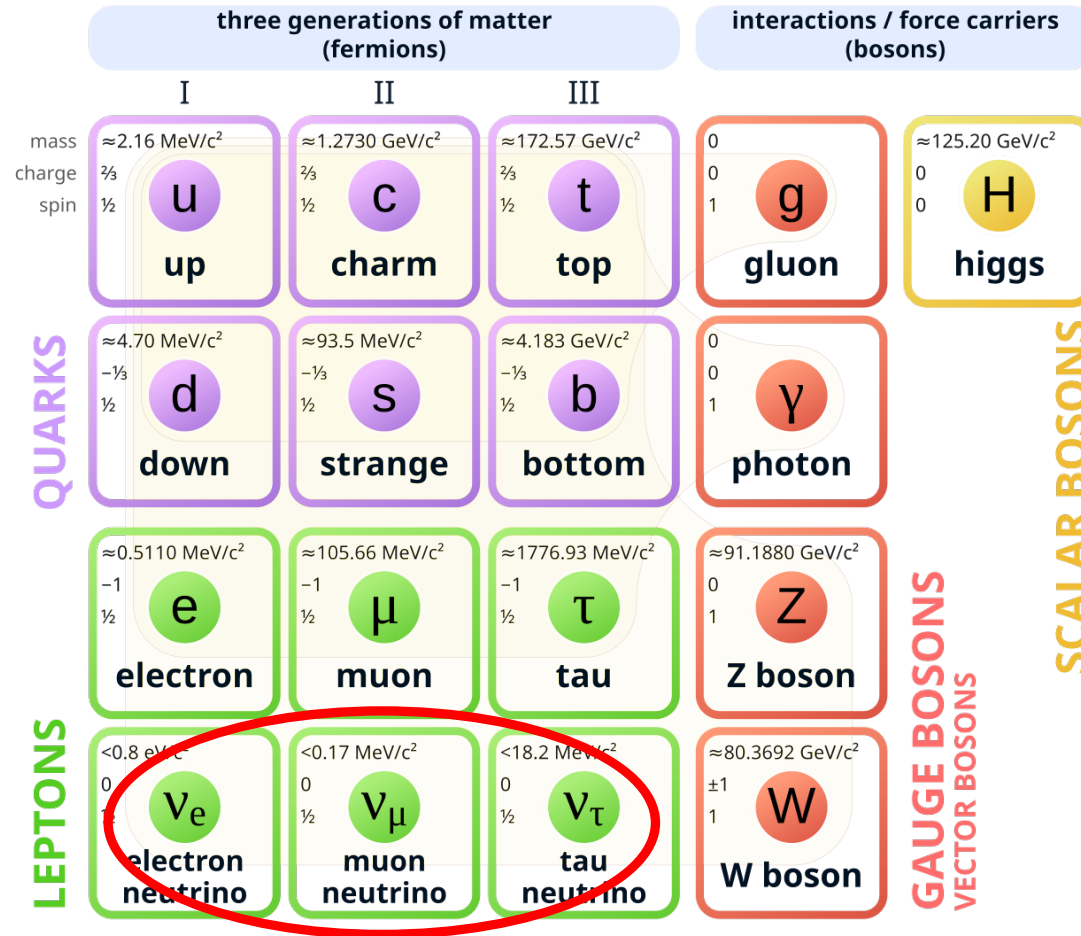
iii. Results

iv. Future Work

Neutrinos

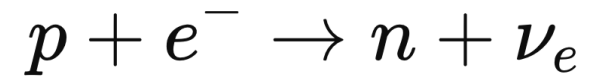
- Neutral leptons
- Weakly interacting
- Three flavors

Standard Model of Elementary Particles



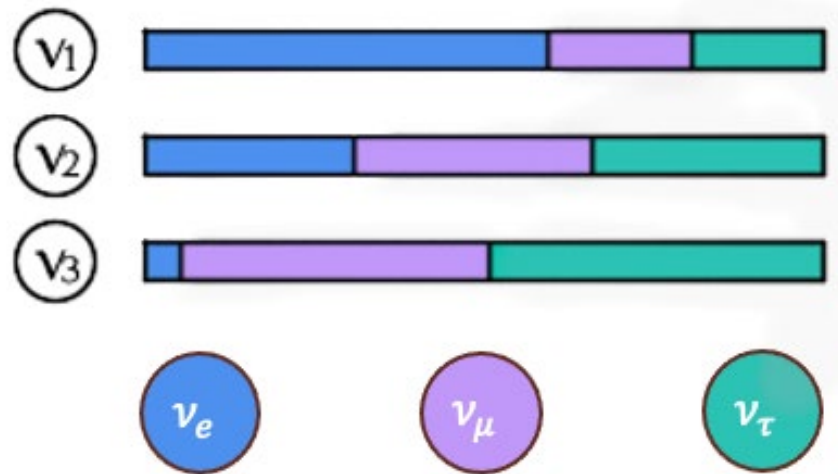
Importance

- Most abundant particle in the universe
- Present in many astronomical settings
 - Early Universe/nucleosynthesis
 - Stars
 - Supernova Explosions
 - Black holes and neutron stars



Neutrino Oscillations

- Massless? no
- Flavors composed of mass eigenstates
- If they have different masses, they can change flavor



$$|\nu_\alpha\rangle = \sum_i U_{\alpha i}^* |\nu_i\rangle$$

$$\alpha = e, \mu, \tau$$

$$i = 1, 2, 3$$

e.g.
$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

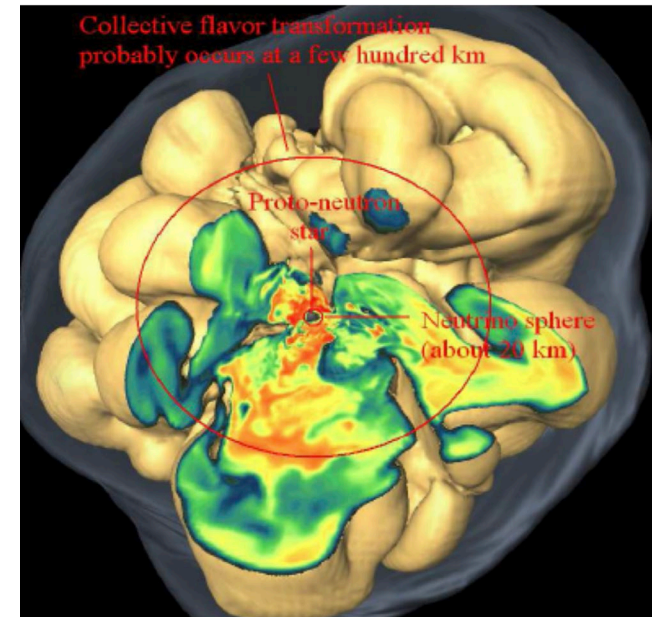
$$|\nu_i(t)\rangle \approx e^{-m_i^2 t / 2E_i} |\nu_i\rangle$$

Neutrinos in Dense Media

- Cannot neglect neutrino interactions in astrophysical setting – too dense
- Want to know the nature of oscillation over time



NASA, ESA, J. Hester and A. Loll (Arizona State University)



V. Cirigliano

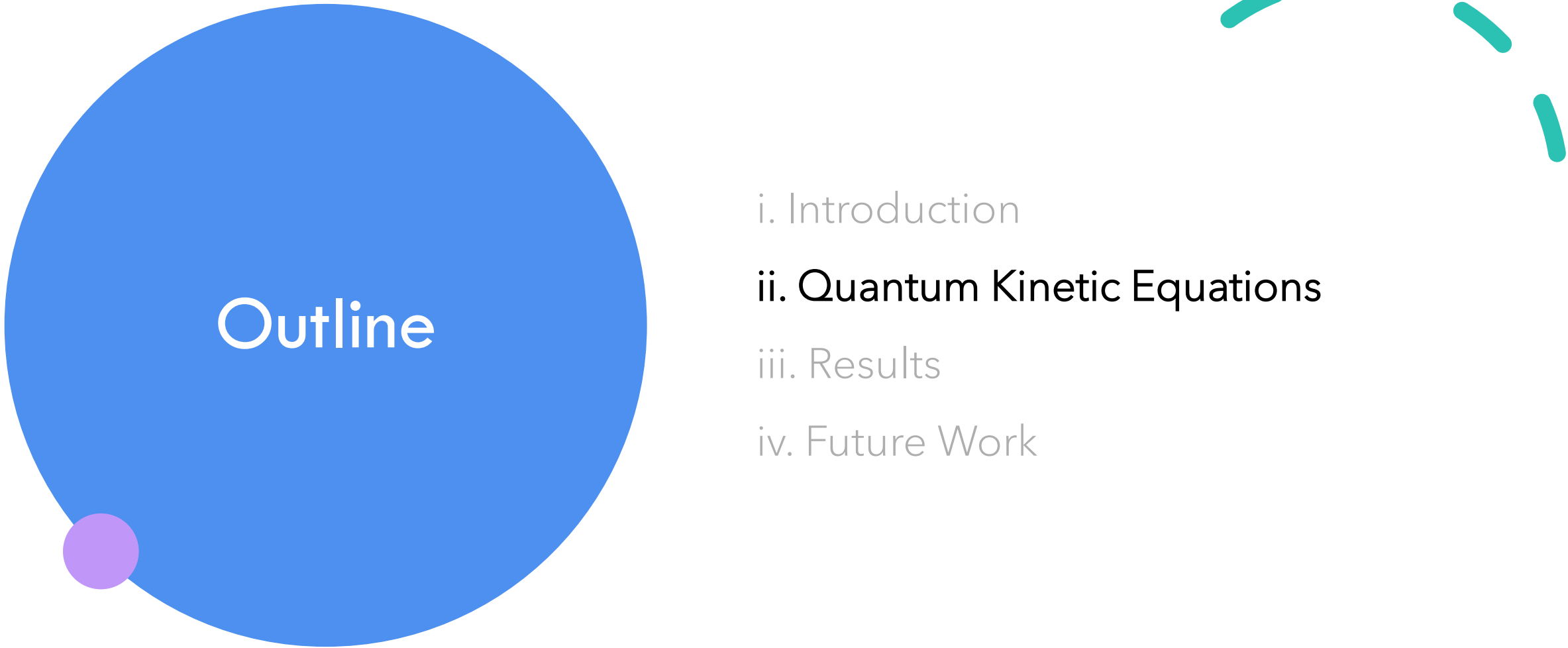
Two Approaches

Quantum Many
Body Approach

vs.

Quantum Kinetic
Equations





Outline

i. Introduction

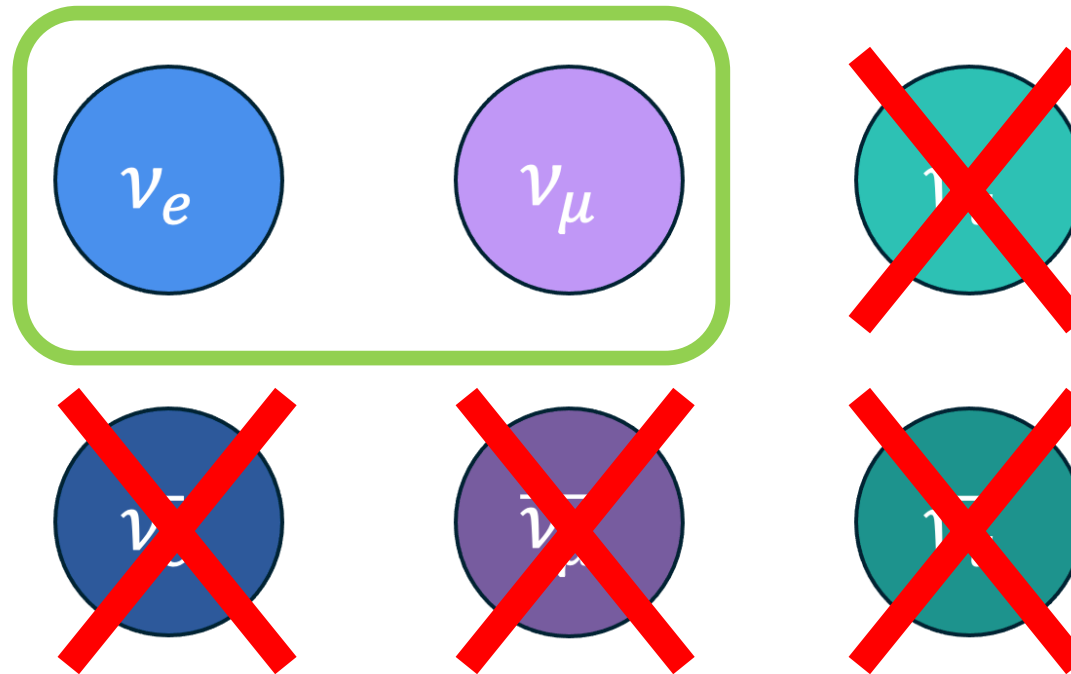
ii. Quantum Kinetic Equations

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iv. Future Work

Simplifications

- Two flavor model
- No antineutrinos



Toy Model

- Supernova Environment
- 35 available momentum bins – discretized
- 10 neutrinos: $6 \nu_e, 4 \nu_\mu$

$$\bar{p} = (p_x, p_y, p_z)$$

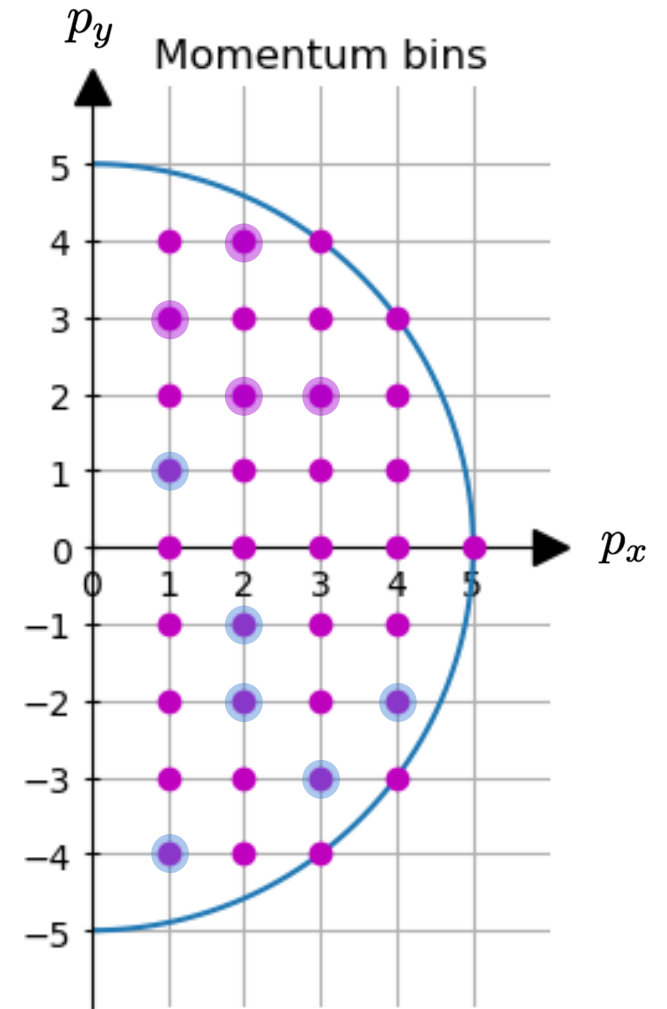
$$p_x = \frac{2\pi}{L} n_x$$

$$p_y = \frac{2\pi}{L} n_y$$

$$p_z = 0$$

$$n_{max} = 5$$

$$x > 0$$



Quantum Kinetic Equations (QKEs)

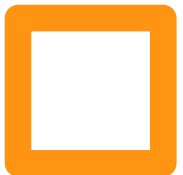
$$F = \begin{pmatrix} f_{LL} & f_{LR} \\ f_{RL} & f_{RR} \end{pmatrix}$$

$$iDF = [H, F] + iC$$

drift and
force term

vacuum
mass and
forward
scattering
terms

inelastic
collision
term



QKEs

$$\dot{f}_{n_p}(t) = -i \left[\frac{\Delta m}{4|\vec{p}_n|} R, f_{n_p}(t) \right] - i \sqrt{2} \frac{G_f}{V} \sum_i \kappa_{ij} [f_{n_q}(t), f_{n_p}(t)] + \text{collision term}$$

time
evolution

vacuum
mass

forward
scattering

inelastic
collisions

$$f(t) = \begin{pmatrix} f_{ee}(t) & f_{e\mu}(t) \\ f_{\mu e}(t) & f_{\mu\mu}(t) \end{pmatrix} = \begin{pmatrix} f_1(t) & f_3(t) + if_4(t) \\ f_3(t) - if_4(t) & f_2(t) \end{pmatrix}$$

Vacuum Mass Term

$$R = \begin{pmatrix} -\cos(2\theta) & \sin(2\theta) \\ \sin(2\theta) & \cos(2\theta) \end{pmatrix}$$

$$\dot{f}_{n_p}(t) = -i \left[\frac{\Delta m}{4|\bar{p}_n|} R, f_{n_p}(t) \right]$$



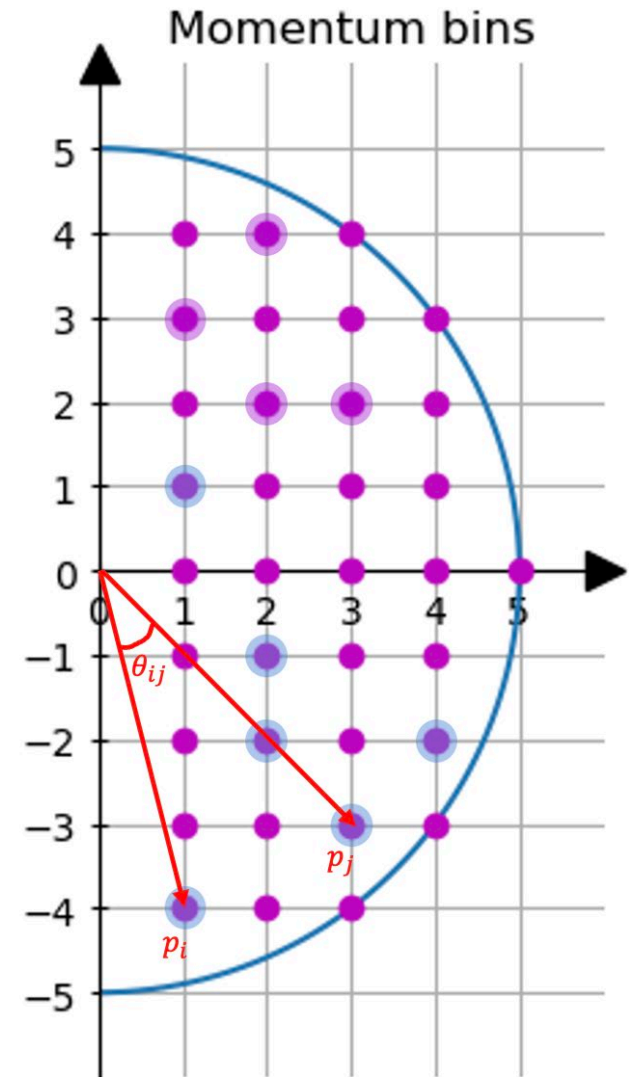
$$f_{n_p}(t) = \begin{pmatrix} 1 - \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2}{4|\bar{p}|} t\right) & -\sin^2(2\theta) \sin^2\left(\frac{\Delta m^2}{4|\bar{p}|} t\right) \\ -\sin^2(2\theta) \sin^2\left(\frac{\Delta m^2}{4|\bar{p}|} t\right) & \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2}{4|\bar{p}|} t\right) \end{pmatrix}$$



Forward Scattering Term

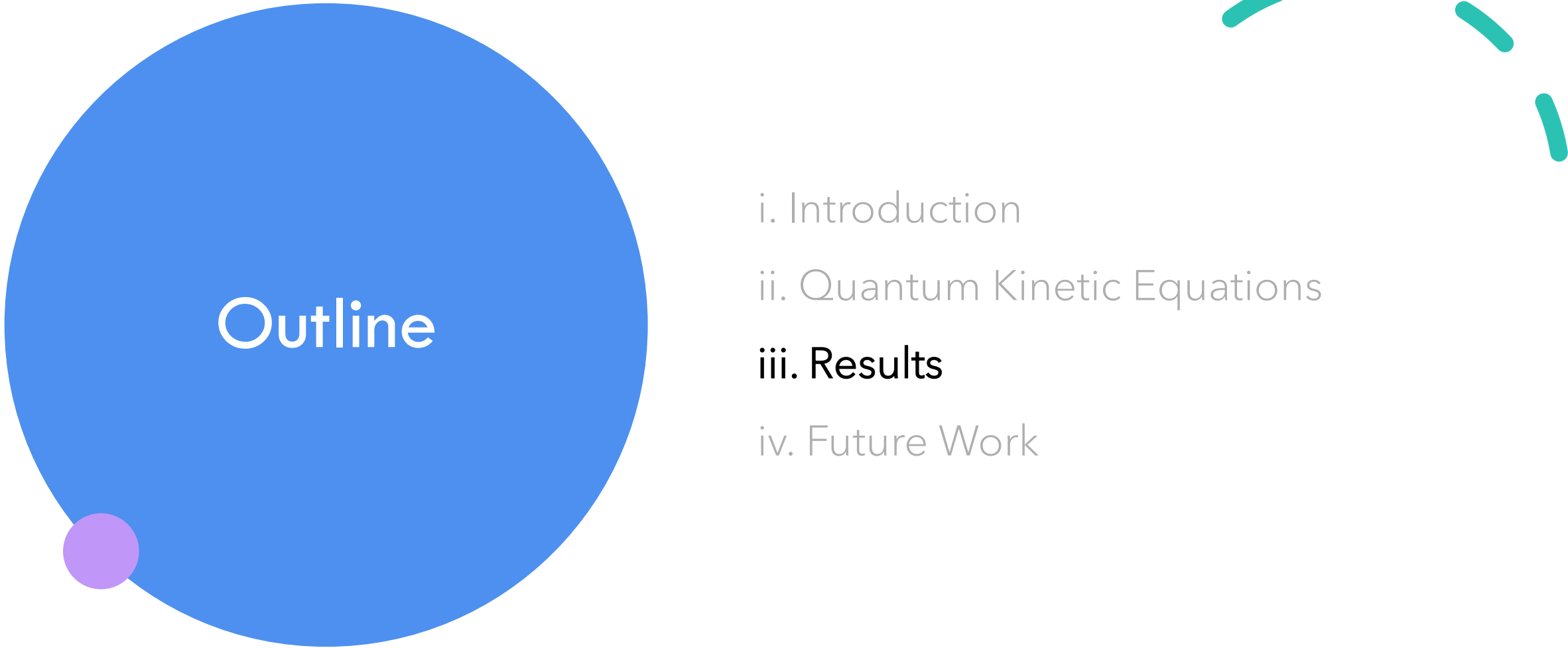
$$\dot{f}_{n_p}(t) = -i\left[\frac{\Delta m}{4|\bar{p}_n|}R, f_{n_p}(t)\right] - i\sqrt{2}\frac{G_f}{V}\sum_i \kappa_{ij}[f_{n_q}(t), f_{n_p}(t)]$$

$$\kappa_{ij} = 1 - \cos \theta_{ij}$$
$$\cos \theta_{ij} = \frac{\bar{p}_i \cdot \bar{p}_j}{|\bar{p}_i||\bar{p}_j|}$$



Collision Term

$$\begin{aligned}
 \mathcal{C}(\vec{n}_p) = & \left(\frac{G_F}{V} \right)^2 V^{1/3} \sum_{\vec{n}_2} \sum_{\vec{n}_3} \sum_{\vec{n}_4} \delta_{\vec{n}_p + \vec{n}_2, \vec{n}_3 + \vec{n}_4} (1 - \hat{n}_p \cdot \hat{n}_2)(1 - \hat{n}_3 \cdot \hat{n}_4) \\
 & \times \frac{|\vec{n}_p| + |\vec{n}_2| - |\vec{n}_3|}{\sqrt{(|\vec{n}_p| + |\vec{n}_2| - |\vec{n}_3|)^2 - n_{4x}^2 - n_{4y}^2}} \theta((|\vec{n}_p| + |\vec{n}_2| - |\vec{n}_3|)^2 - n_{4x}^2 - n_{4y}^2) \\
 & \times \left(\bar{\delta}_{n_{4z}, \sqrt{(|\vec{n}_p| + |\vec{n}_2| - |\vec{n}_3|)^2 - n_{4x}^2 - n_{4y}^2}} - \bar{\delta}_{n_{4z}, -\sqrt{(|\vec{n}_p| + |\vec{n}_2| - |\vec{n}_3|)^2 - n_{4x}^2 - n_{4y}^2}} \right) \\
 & \times \left\{ \left[f_4(I - f_2) + \text{Tr}(f_4(I - f_2)) \right] f_3(I - f_p) + (I - f_p)f_3 \left[(I - f_2)f_4 + \text{Tr}((I - f_2)f_4) \right] \right. \\
 & \left. - \left[(I - f_4)f_2 + \text{Tr}((I - f_4)f_2) \right] (I - f_3)f_p - f_p(I - f_3) \left[f_2(I - f_4) + \text{Tr}(f_2(I - f_4)) \right] \right\}
 \end{aligned}$$



Outline

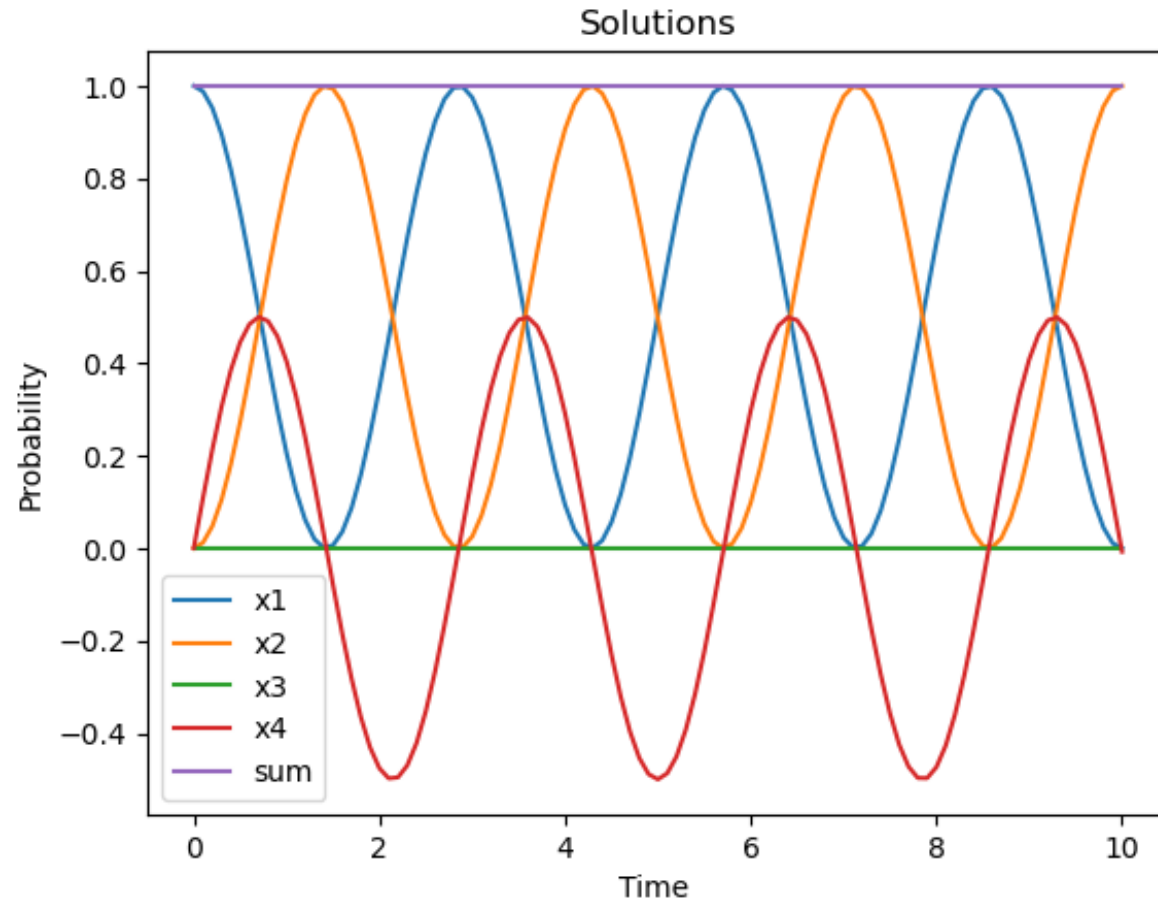
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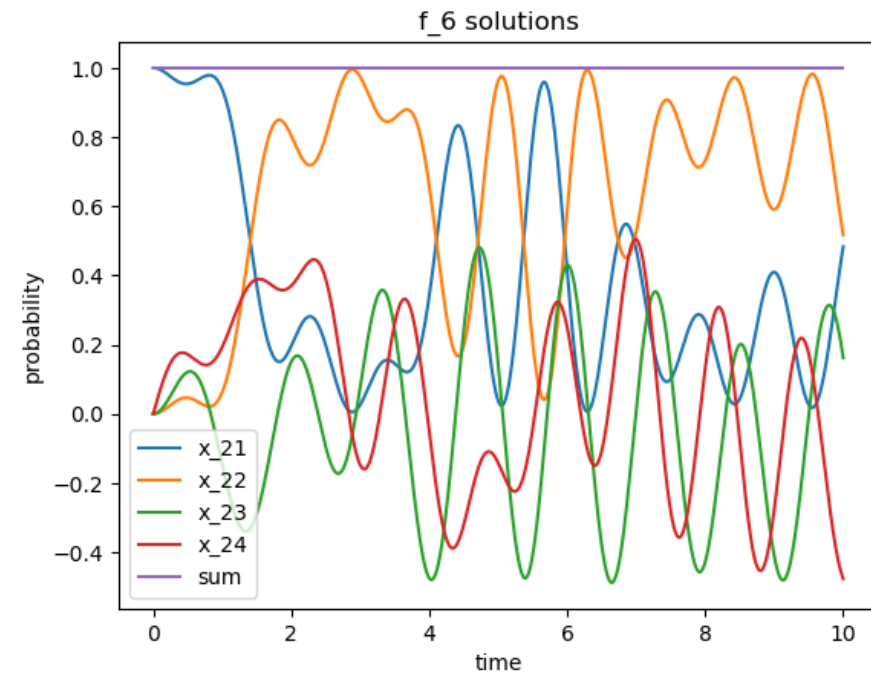
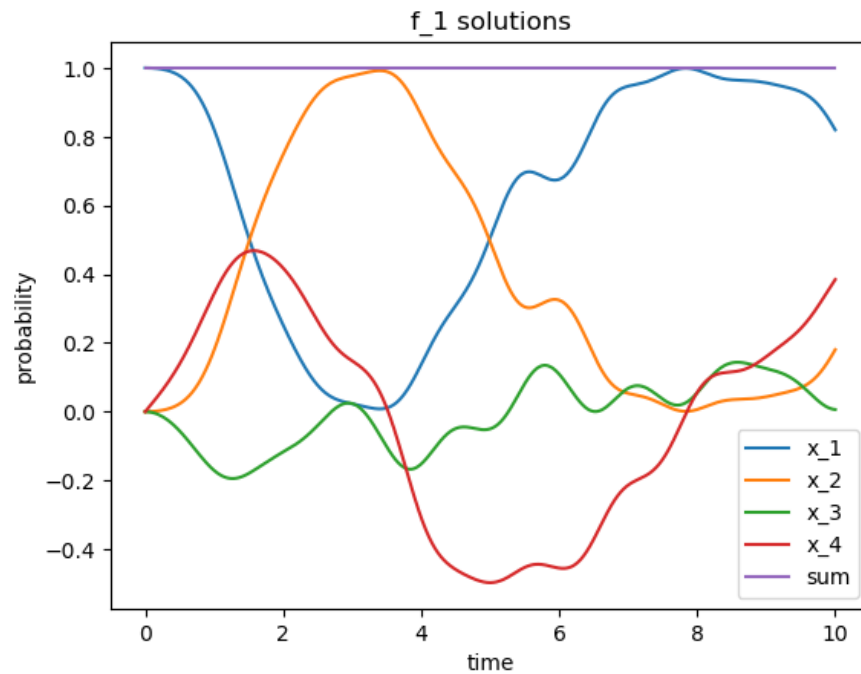
iv. Future Work

Results: No Coherent Evolution

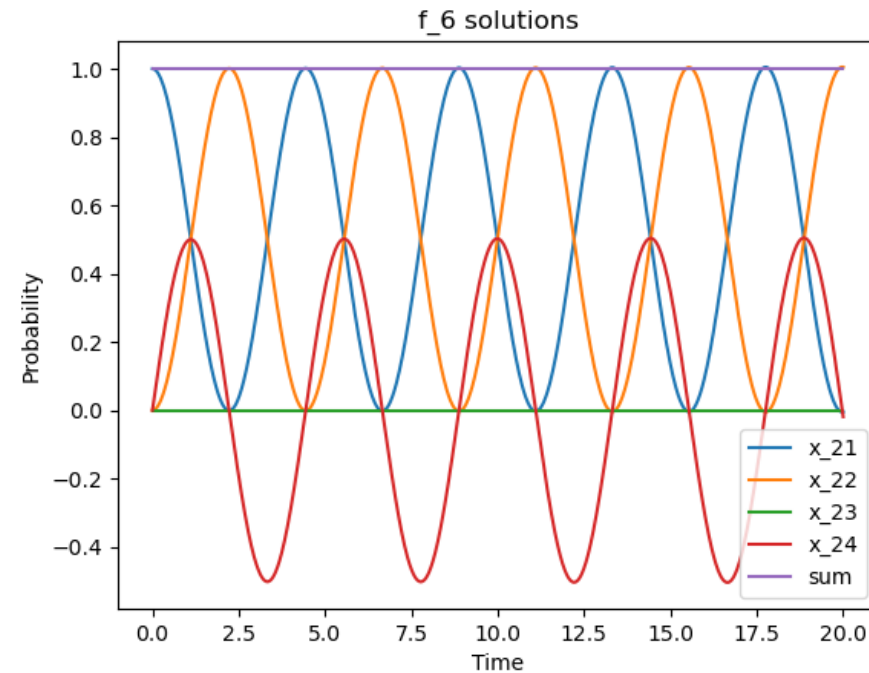
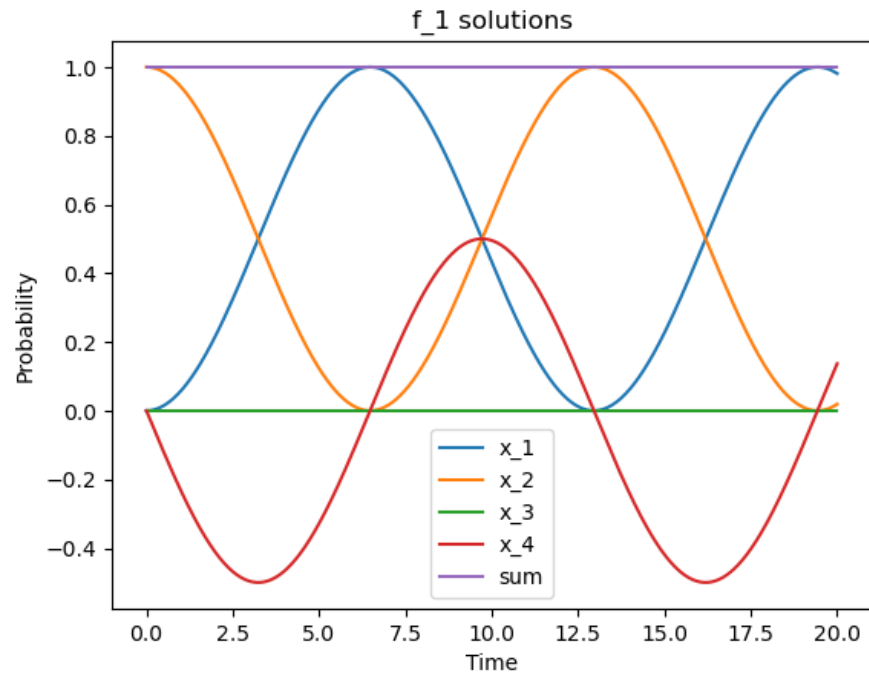


$$f_{n_p}(t) = \begin{pmatrix} 1 - \sin^2(2\theta)\sin^2(\frac{\Delta m^2}{4|\vec{p}|}t) & - \\ - & \sin^2(2\theta)\sin^2(\frac{\Delta m^2}{4|\vec{p}|}t) \end{pmatrix}$$

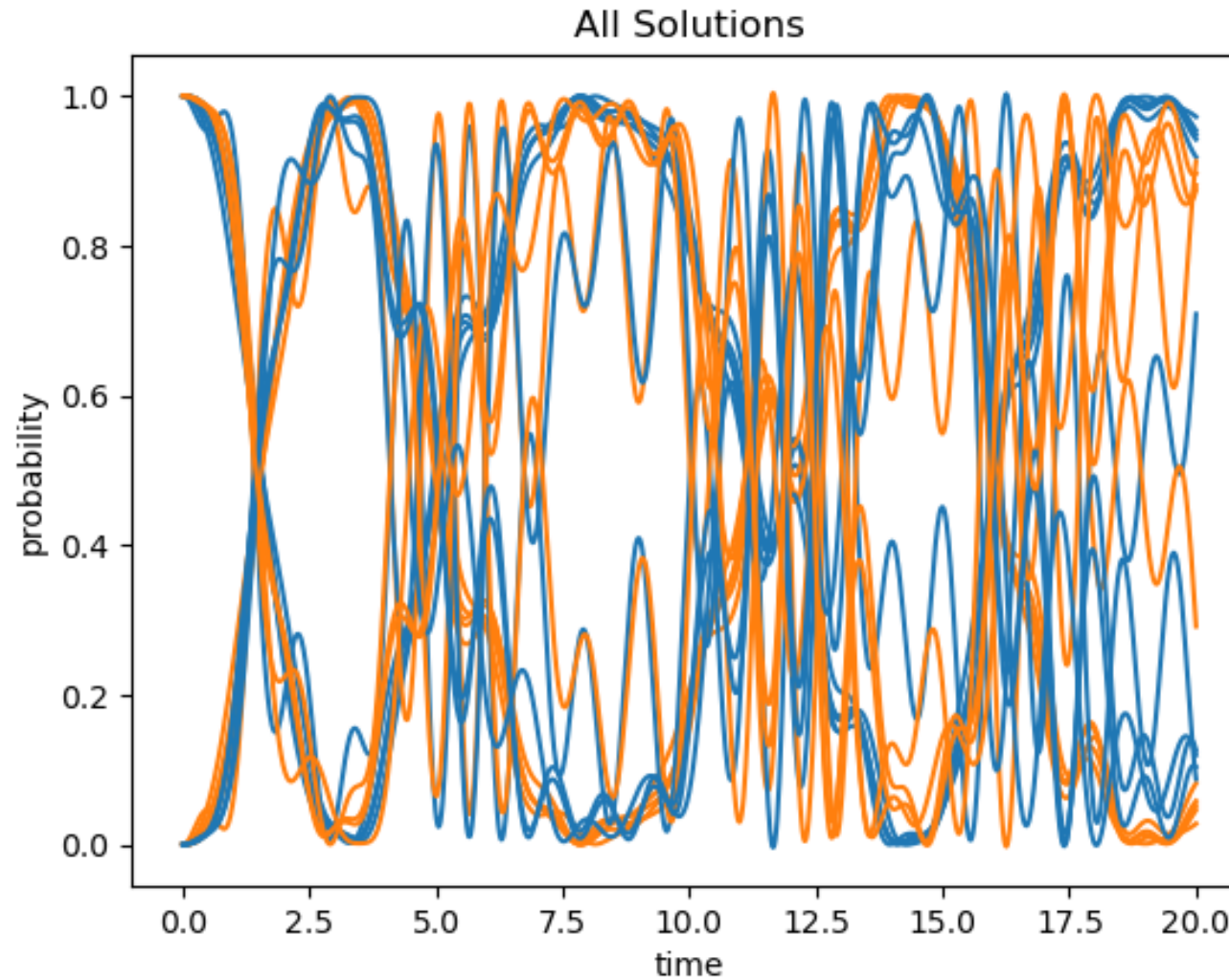
Results: Coherent Evolution



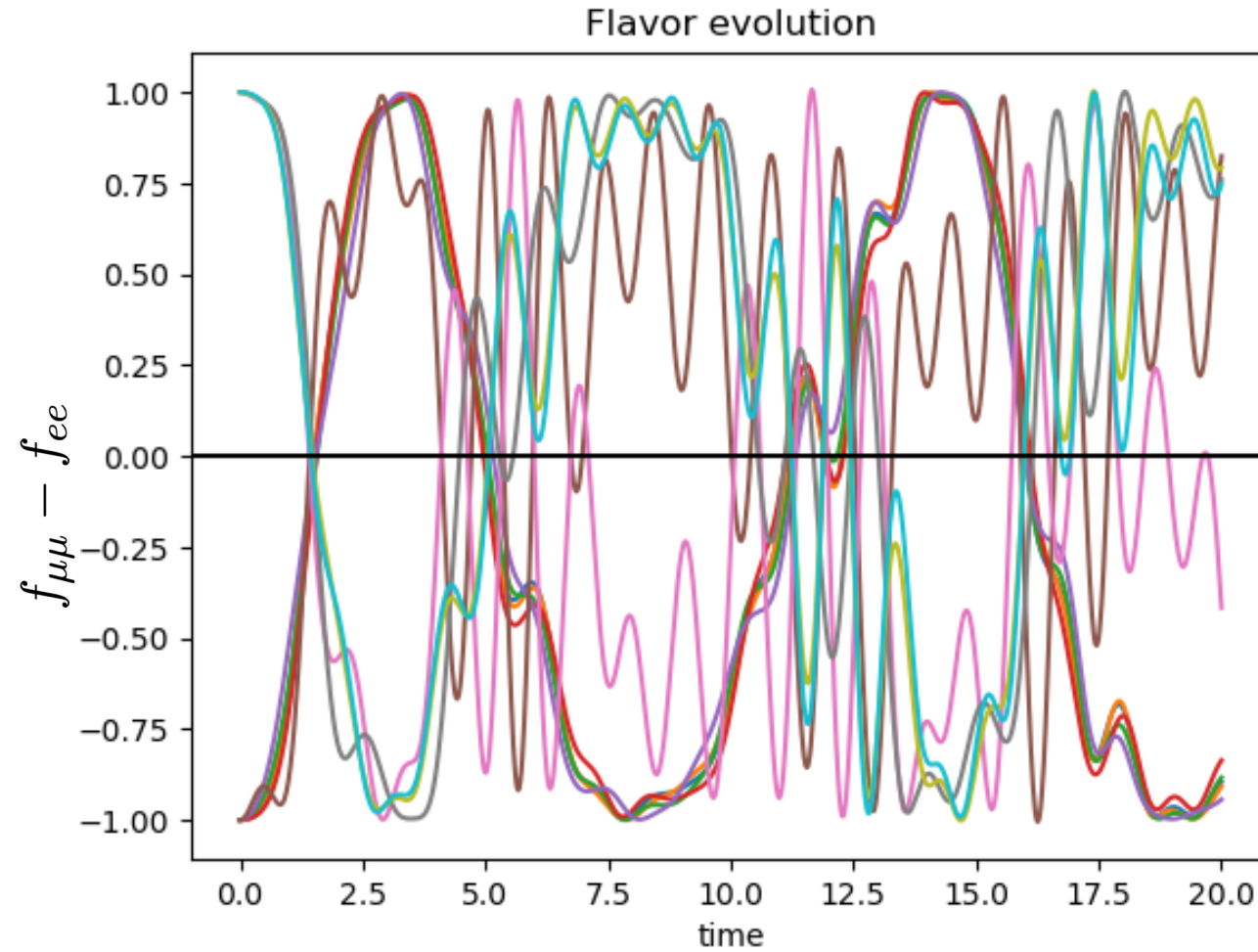
Comparison Between Solutions



Results: Coherent Evolution

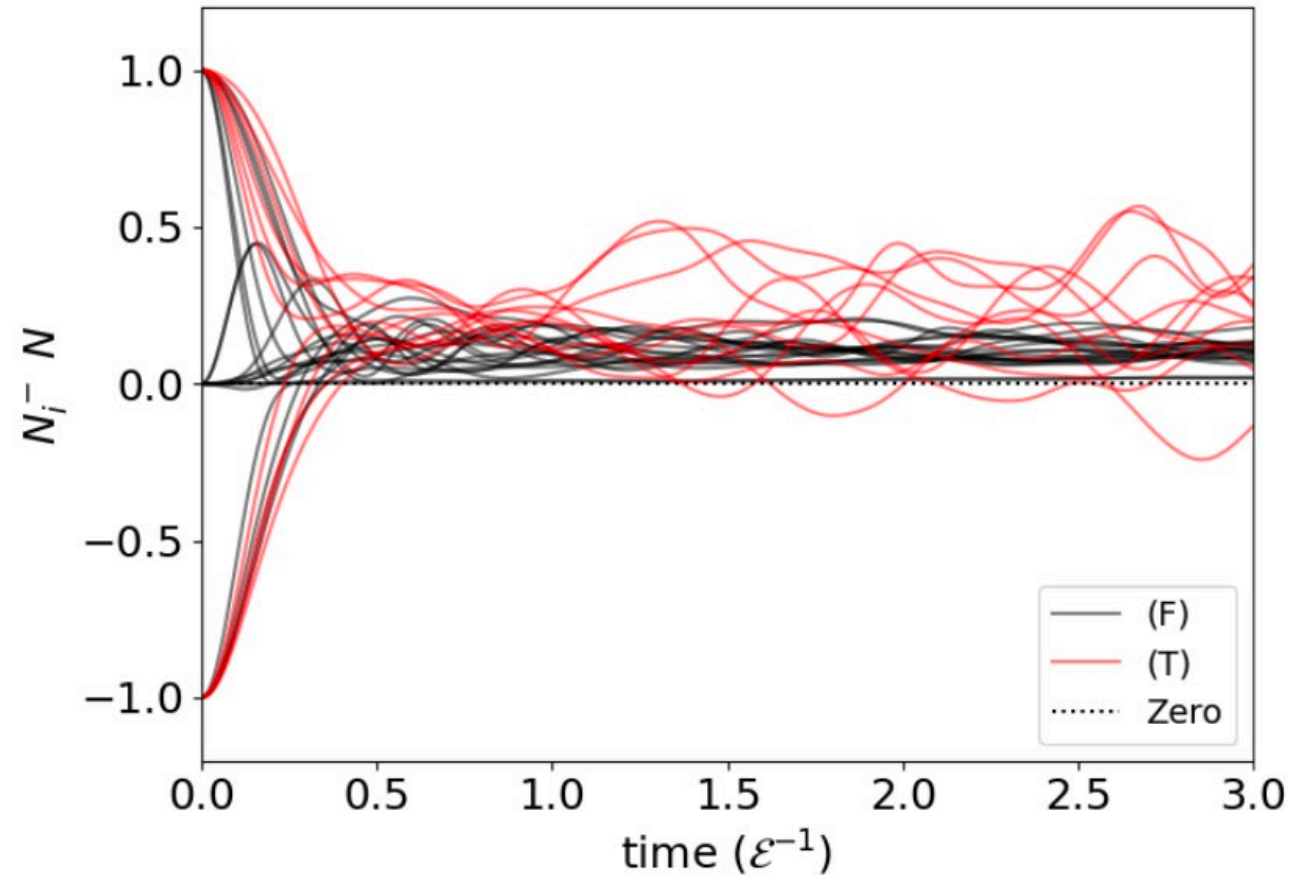


Results: Flavor Evolution



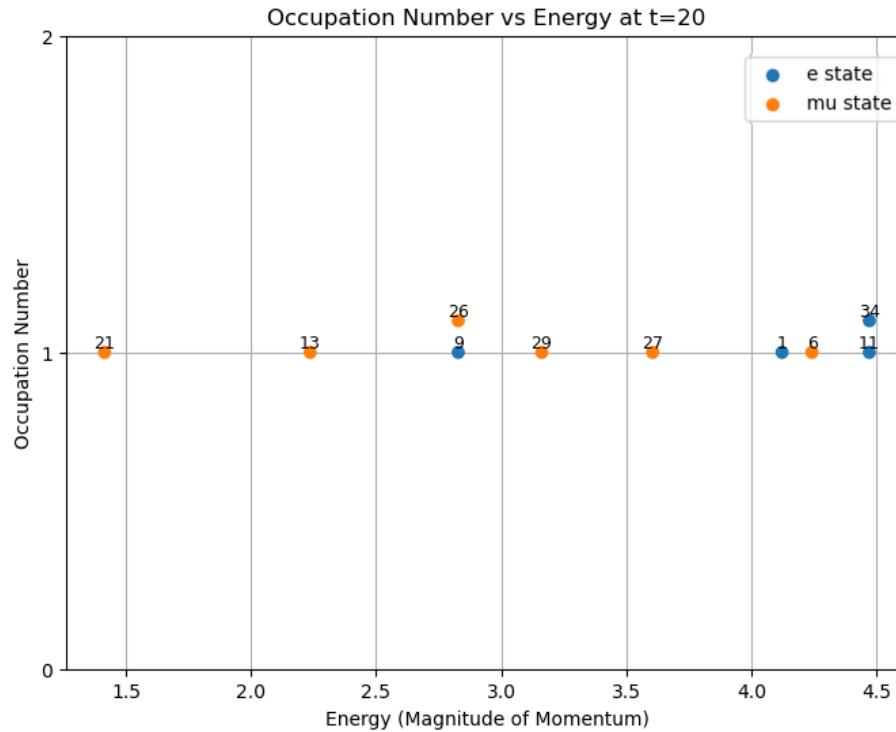
Comparison: Flavor Evolution

- Many body
- With collisions



V. Cirigliano, S. Sen, Y. Yamauchi (2024)

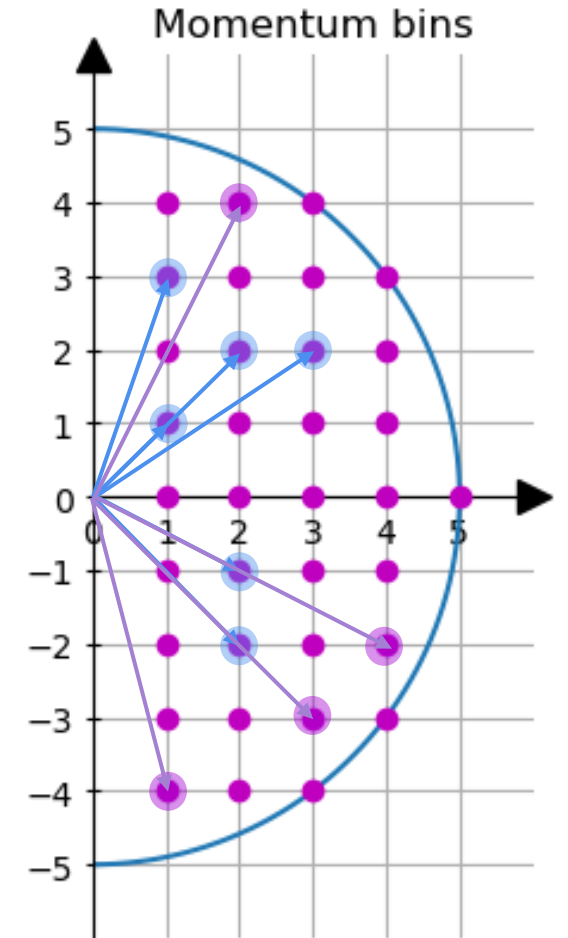
Occupation Number vs. Energy



$$E = |\vec{p}|$$

$$f_{ee} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$f_{\mu\mu} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$





Outline

i. Introduction

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Future Work

- Collision term
- Compare with quantum many body approach

Thank you!



Special thanks to my advisor,
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Arthur Barnard; and my
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