Metasurface-based spin-selective optical cavity

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Wednesday, August 17, 2016

Outline

- Motivation
- Introduction to the problem
- Proposed cavity design
- Metasurface optics
- ► Conclusion

Goals and motivation

We seek a cavity which **differentiates between left- and right-handed light** within the cavity volume.

$$\begin{array}{c} \left| \mathsf{H} \right\rangle \\ \left| \mathsf{V} \right\rangle \\ \left| \mathsf{L} \right\rangle \\ \left| \mathsf{R} \right\rangle \end{array} \left(\begin{array}{c} - & \\ - & \\ | \mathsf{L} \right\rangle \\ - & \\ | \mathsf{L} \right\rangle \end{array} \right)$$

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- quantum information processing

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Related problem: spin-preserving mirror For incident light normal to a good conductor, we have

$$\left(\frac{E_{0R}}{E_{0I}}\right)_{N} = \frac{Z_{2}\cos\theta_{I} - Z_{1}\cos\theta_{T}}{Z_{2}\cos\theta_{I} + Z_{1}\cos\theta_{T}} \approx -1$$

where Z_1, Z_2 are the impedences of air and the conductor respectively, and $Z_1 \gg |Z_2|$.

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Hence, E_R gains a uniform π phase shift and is "reflected" with no preferred transverse axis.

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It is useful to preserve one handedness in our cavity: hence, we may use a quarter wave plate preceding the mirror to "preserve" spin after reflection.

We use birefringent materials to impose polarization-dependent path lengths.





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For an orthonormal polarization basis $\hat{\imath}$, $\hat{\jmath}$, denote

$$\ket{ \mathsf{u}(z)} = egin{pmatrix} u_1(z) \ u_2(z) \end{pmatrix} := \ket{ u_1(z) \, \hat{\imath}} + \ket{ u_2(z) \, \hat{\jmath}}.$$

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Then we may define the expected *local* rotation operator,

$$R(heta) := egin{pmatrix} \cos heta & -\sin heta \ \sin heta & \cos heta \end{pmatrix}.$$

Hence we define propagation in some birefringent region aligned with our polarization basis

$$\hat{Q}(z_i,z_j) = egin{pmatrix} \hat{U}(z_i) & 0 \ 0 & \hat{U}(z_j) \end{pmatrix}$$

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and the cavity roundtrip operator follows:

$$egin{aligned} \hat{T} &= \hat{Q}(lpha+\delta,lpha) R\left(rac{\pi}{4}
ight) \hat{Q}(2eta,0) R^{\dagger}\left(rac{\pi}{4}
ight) \ldots \ \hat{Q}(lpha+\delta,lpha+\delta) R\left(rac{\pi}{4}
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$$\hat{\mathcal{T}} = rac{1}{2} \left(\hat{\mathcal{U}}(4lpha+2\delta) + \hat{\mathcal{U}}(4lpha+4eta+2\delta)
ight) I_2 + \ rac{1}{2} \hat{\mathcal{U}}(4lpha) \left(\hat{\mathcal{U}}(4eta) - 1
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Hence, we find (normalized) eigenvectors of

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angle$$

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that is,

$$\hat{U}(4\alpha + 4\beta + 2\delta), \quad \hat{U}(4\alpha + 2\delta)$$

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A phase picture of optical elements Phase profile for a thin lens with focal length f:

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If we allow birefringence: Half wave plate:

$$\phi_x = \pi; \ \phi_y = 0$$

Quarter wave plate:

$$\phi_x = \frac{\pi}{2}; \ \phi_y = 0$$

Arbabi *et al.* implement arrays of **elliptical**, **subwavelength high-contrast posts to exhibit birefringence**.

²Arbabi, Horie, Bagheri, & Faraon. Dieletric metasurfaces. *Nature Nano.* **10**, 937-944 (2015).

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- lattice constant
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RCWA is used to determine phase and amplitude for a given parameter set.

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Silicon nitride-based metasurfaces



Figure: low-contrast metasurface optics (SEM).³ (a) lens, (b) vortex beam generator.

³Zhan *et al.*. Low-contrast dielectric metasurface optics. *ACS Photonics.* (2015).

Further work

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 - Transverse modes (cavity as system of coupled harmonic oscillators)
 - Explicit definition of propagation operator and mode functions

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 - Transverse modes (cavity as system of coupled harmonic oscillators)
 - Explicit definition of propagation operator and mode functions
- ► Simulate elements, cavity with FDTD

Acknowledgements

I would like to thank Dr. Majumdar, Alan, and the NOISE Lab group for their guidance and support.

Thank you to the INT REU directors and administrators for all of their time and attention in support of this summer.