

Magnetic field characterization and the search for a permanent electric dipole moment in ^{199}Hg

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Outline

Permanent electric dipole moments and symmetry

The experiment at the University of Washington

Characterizing the magnetic field

- Data collection

- Gradients and field reconstruction

- Future work



Purcell and Ramsey (1950)

that this electric dipole would be a polar vector, being the product of the angular momentum (an axial vector) and the magnetic pole strength, which is a pseudoscalar in conformity with the usual convention that electric charge is a simple scalar.

The argument against electric dipoles, in another form, raises directly the question of parity. A nucleon with an electric dipole moment would show an asymmetry between left- and right-handed coordinate systems; in one system the dipole moment would be parallel to the angular momentum and in the other, antiparallel. But there is no compelling reason for excluding this possibility. It would not be the only asymmetry of particles of ordinary experience, which already exhibit conspicuous asymmetry in respect to electric charge. Although magnetic poles were used above as an illustration of a particular mechanism by which a nuclear electric dipole could arise, this is, of course, not the only possibility.

The question of the possible existence of an electric dipole



C, P, and T

- Charge *conjugation*



C, P, and T

- C *Charge conjugation*
- P *Parity*

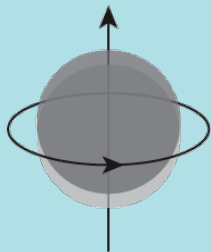


C, P, and T

- Charge conjugation
- Parity
- Time reversal

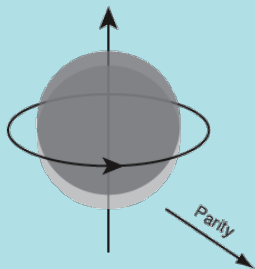


Permanent EDMs and symmetry violation



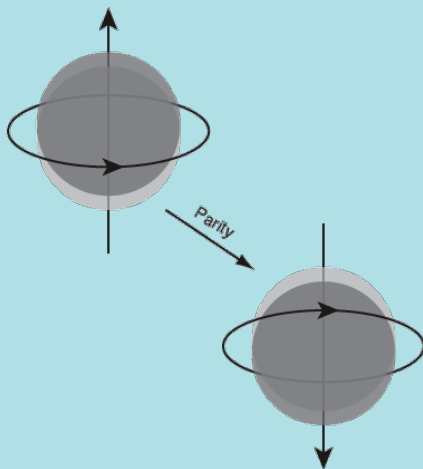


Permanent EDMs and symmetry violation



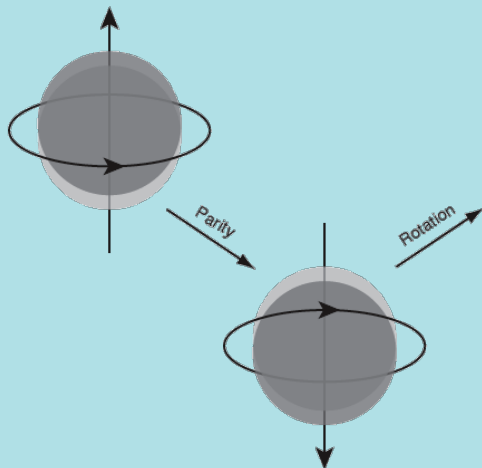


Permanent EDMs and symmetry violation



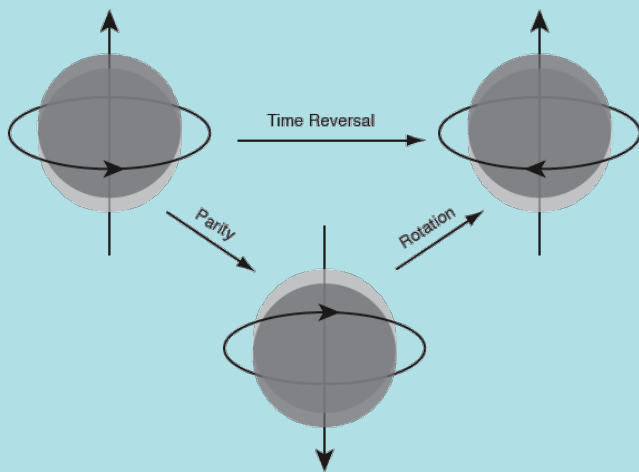


Permanent EDMs and symmetry violation





Permanent EDMs and symmetry violation





Why look for a permanent EDM?

- *CPT* Theorem: $CP \leftrightarrow T$



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- The existence of matter



Why look for a permanent EDM?

- *CPT* Theorem: $CP \leftrightarrow T$
- The existence of matter
- Supersymmetry



Search for a permanent EDM of ^{199}Hg

- ^{199}Hg is diamagnetic (no unpaired electrons) \rightarrow nuclear EDM
- $|d(^{199}\text{Hg})| < 3.1 \times 10^{-29} e \text{ cm}$ (95% C.L.)



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- $|d(^{199}\text{Hg})| < 3.1 \times 10^{-29} e \text{ cm}$ (95% C.L.)
- \uparrow This is exceedingly small



Larmor precession

$$\tau = \mu \times B + d \times E$$



Larmor precession

$$\tau = \mu \times B + d \times E$$

$$\tau_{E \parallel B} - \tau_{E \text{ anti-} \parallel B} = 2dE \sin \theta$$



Larmor precession monitoring

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- Long polarization lifetimes depend on the uniformity of the magnetic field

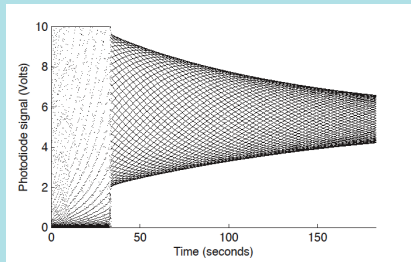


Image: Matthew Swallows



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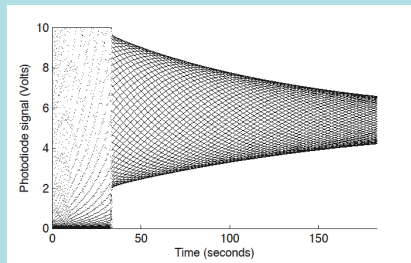
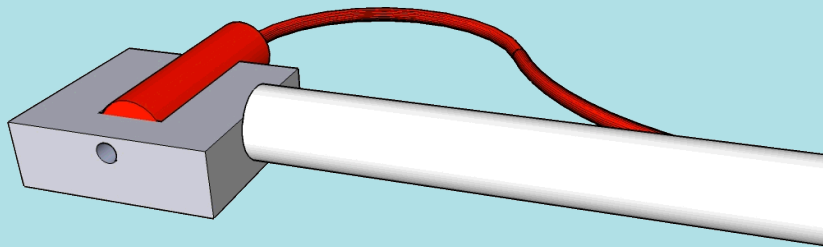


Image: Matthew Swallows

↑ This is where I come in



The flux gate



- .01 milligauss precision
- 10 second measurement intervals, 50 Hz sampling rate



The cosine coil

- Actually three coils: vertical, axial, transverse
- Enclosed in three layers of magnetic shielding

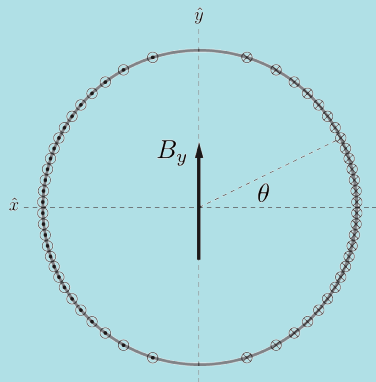


Image: Matthew Swallows



Data collection

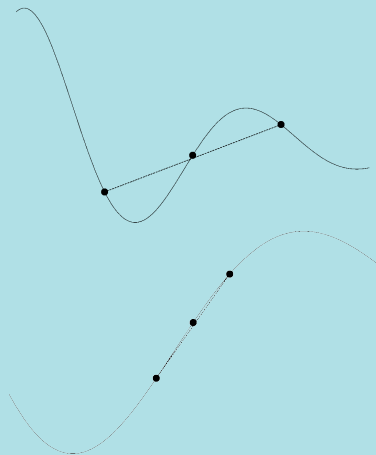
Object: To characterize behavior of field at center of coil

- Cubic lattice of 27 points, 5 mm between adjacent points
- 6 runs: {vertical, axial, transverse} \times 2 orthogonal positions*



Taylor series expansion

- Second order Taylor series expansion about center
- Requires lots of derivatives





Solving for \mathbf{B}

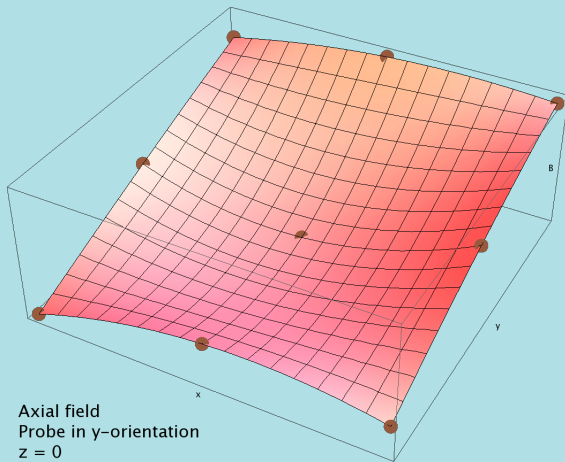
A completely determined* system requires nine linearly independent equations:

$$\nabla \times \mathbf{B} \quad (3)$$

$$\nabla \cdot \mathbf{B} \quad (1)$$

$$\nabla B_a, \quad a = x, y, \text{ or } z \quad (3)$$

$$\nabla B_b, \quad b = x, y, \text{ or } z \text{ and } b \neq a \quad (2)$$





The next step

- Finish data collection on other coil
- Design additional coils to address “worst” gradients



Acknowledgments

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References

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Questions?