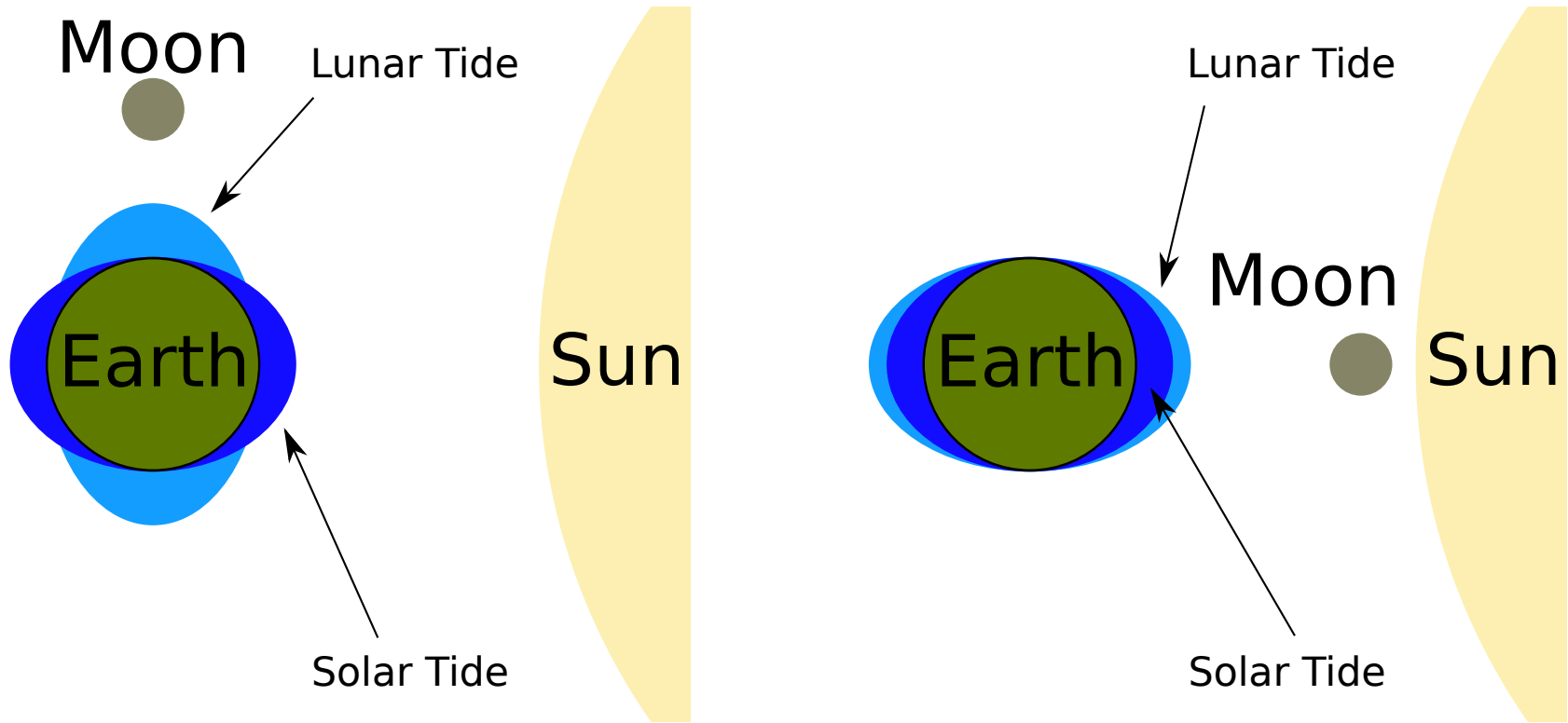


# Why are there three-body forces?

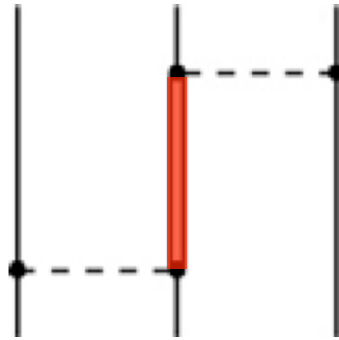
tidal effects lead to 3-body forces in earth-sun-moon system



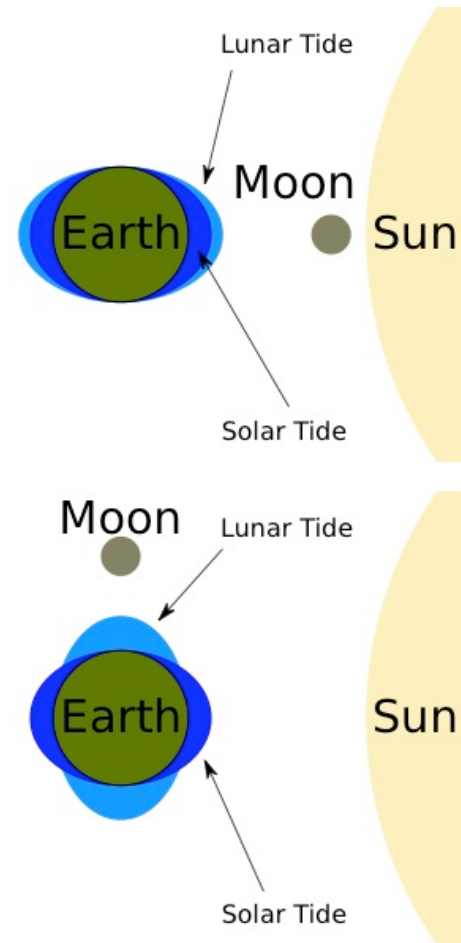
# Why are there 3N forces?

Nucleons are finite-mass composite particles,  
can be excited to resonances

dominant contribution from  $\Delta(1232 \text{ MeV})$

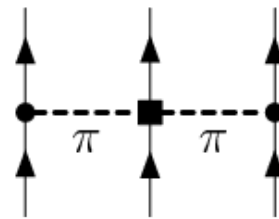


+ many shorter-range parts



chiral effective field theory (EFT)

Delta-less ( $\Delta$  is treated as heavy):

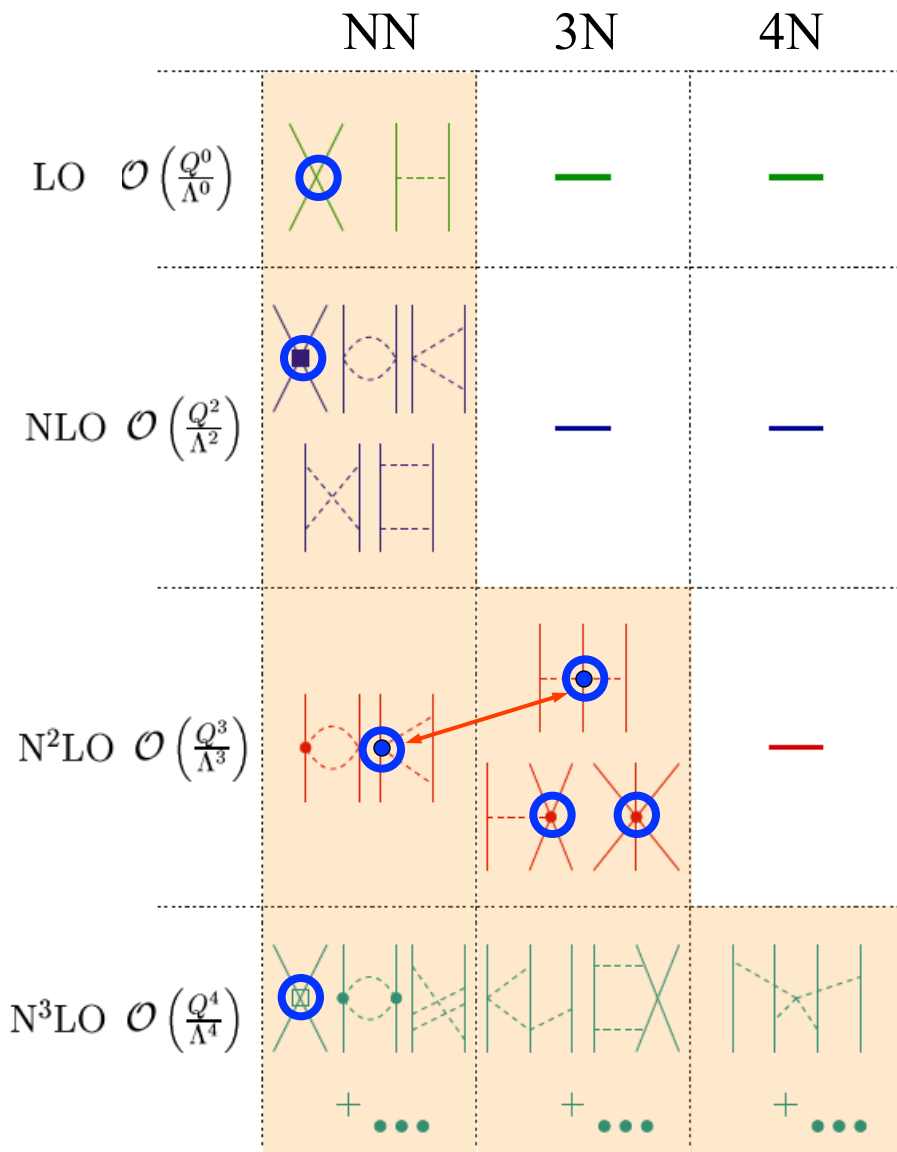


+ shorter-range parts

**EFT provides a systematic and powerful approach for 3N forces**

# Chiral effective field theory for nuclear forces

Separation of scales: low momenta  $\frac{1}{\lambda} = Q \ll \Lambda_b$  breakdown scale  $\sim 500$  MeV



include long-range pion physics

few short-range couplings,  
fit to experiment once

systematic: can work to desired  
accuracy and obtain **error estimates**

consistent **electroweak interactions**  
and **matching to lattice QCD**

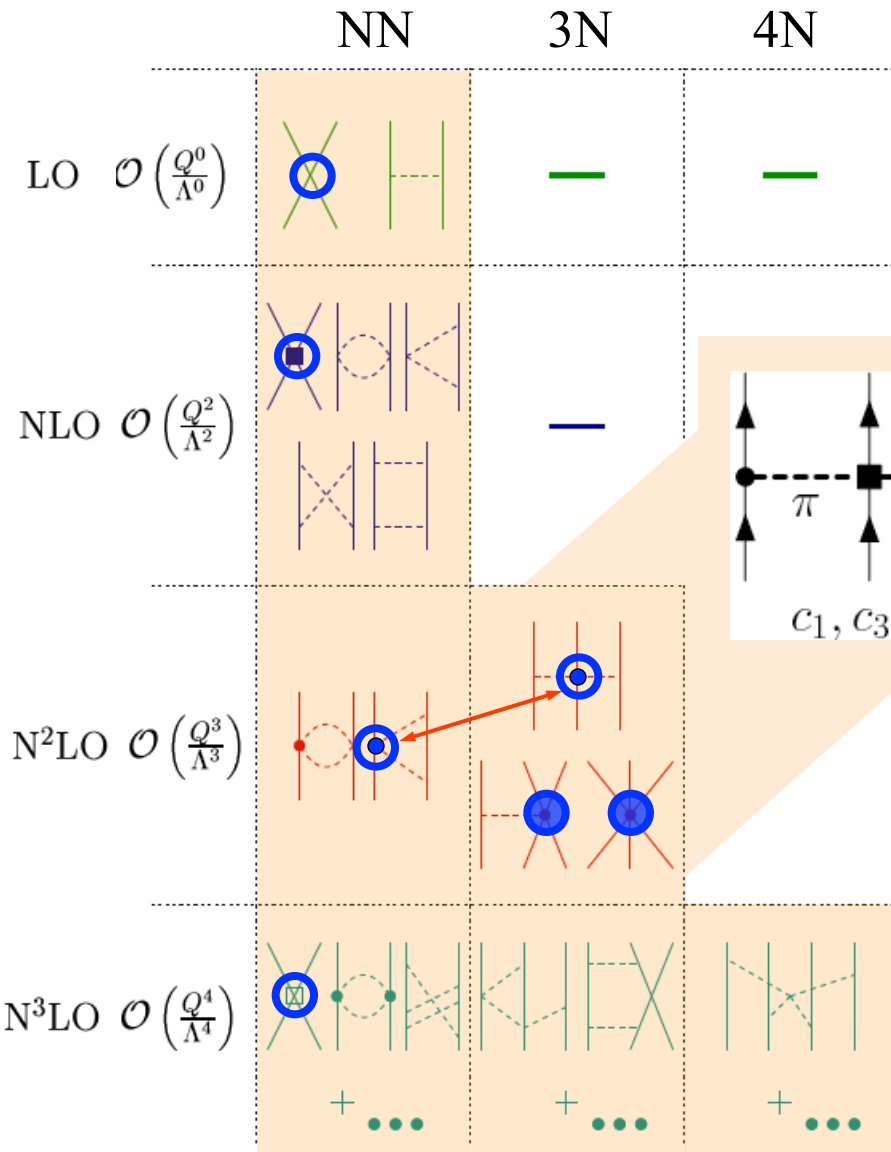
# 3N forces in different EFTs

	pionless	chiral	chiral+ $\Delta$
LO			
NLO			
N <sup>2</sup> LO			

FIG. 23 Order of 3NF contributions in pionless and chiral EFT and in EFT with explicit  $\Delta$  degrees of freedom (chiral+ $\Delta$ ). Open vertices in the last column indicate the differences of the low-energy constants in chiral and chiral+ $\Delta$  EFT.

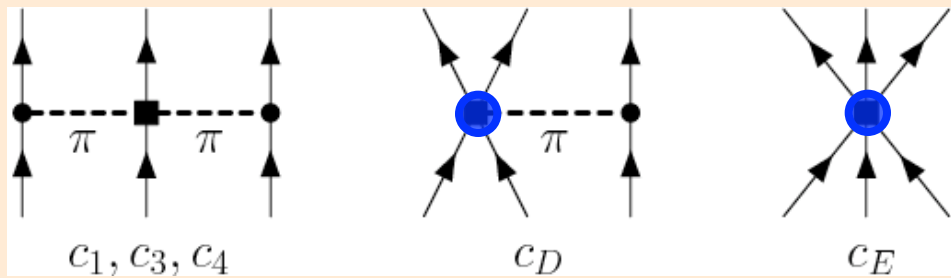
# Chiral effective field theory and many-body forces

Separation of scales: low momenta  $\frac{1}{\lambda} = Q \ll \Lambda_b$  breakdown scale  $\sim 500$  MeV



consistent NN-3N-4N interactions

3N,4N: 2 new couplings to N<sup>3</sup>LO



$c_i$  from  $\pi$ N and NN [Meissner et al. \(2007\)](#)

$$c_1 = -0.9_{-0.5}^{+0.2}, \quad c_3 = -4.7_{-1.0}^{+1.2}, \quad c_4 = 3.5_{-0.2}^{+0.5}$$

single- $\Delta$ :  $c_1=0, c_3=-c_4/2=-3 \text{ GeV}^{-1}$

$c_D, c_E$  fit to  ${}^3\text{H}, {}^4\text{He}$  properties only

## Range of $c_i$ couplings

	$c_1$	$c_3$	$c_4$	
Fettes <i>et al.</i> (1998) (Fit 1)	-1.2	-5.9	3.5	$\pi$ N
Büttiker and Meißner (2000)	-0.8	-4.7	3.4	$\pi$ N
Meißner (2007)	-0.9	-4.7	3.5	$\pi$ N
Rentmeester <i>et al.</i> (2003)	-0.8	-4.8	4.0	NN
Entem and Machleidt (2002)	-0.8	-3.4	3.4	NN
Entem and Machleidt (2003)	-0.8	-3.2	5.4	NN
Epelbaum <i>et al.</i> (2005)	-0.8	-3.4	3.4	NN
Bernard <i>et al.</i> (1997)	-0.9	-5.3	3.7	res

## High-order analysis Krebs *et al.* (KGE) (2012)

	$c_1$ [GeV <sup>-1</sup> ]	$c_3$ [GeV <sup>-1</sup> ]
N <sup>2</sup> LO/N <sup>3</sup> LO EGM NN [31, 32]	-0.81	-3.40
N <sup>3</sup> LO EM NN [33, 34]	-0.81	-3.20
N <sup>2</sup> LO KGE [39]	-(0.26 - 0.58)	-(2.80 - 3.14)
'N <sup>2</sup> LO' KGE (recom.) [39]	-(0.37 - 0.73)	-(2.71 - 3.38)
N <sup>3</sup> LO KGE [39]	-(0.75 - 1.13)	-(4.77 - 5.51)

# Subleading chiral 3N forces

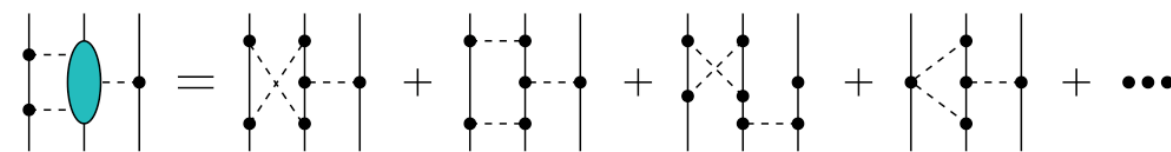
parameter-free  $N^3LO$  Bernard et al. (2007,2011), Ishikawa, Robilotta (2007)

one-loop contributions:

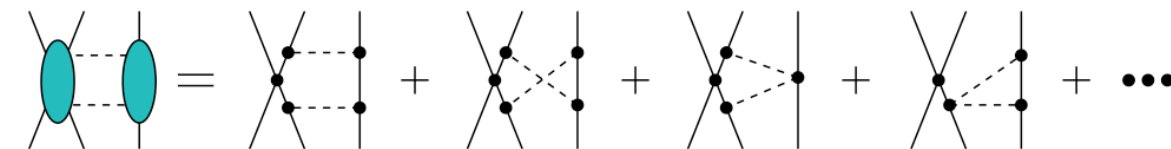
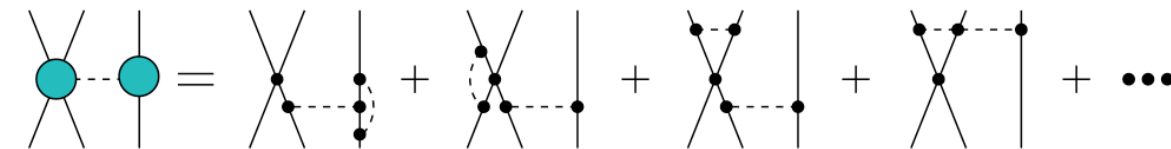
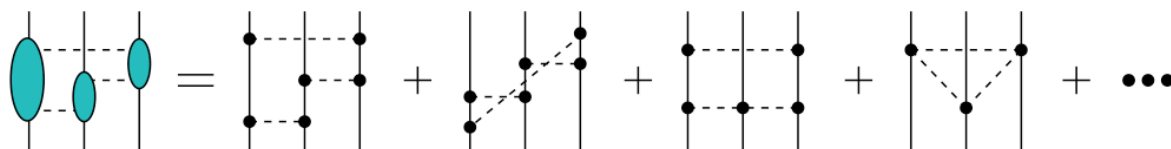
$2\pi$ -exchange,  $2\pi$ - $1\pi$ -exchange, rings, contact- $1\pi$ -, contact- $2\pi$ -exchange



decrease  $c_i$  strengths  
 $\delta c_3 = -\delta c_4 = 1 \text{ GeV}^{-1}$

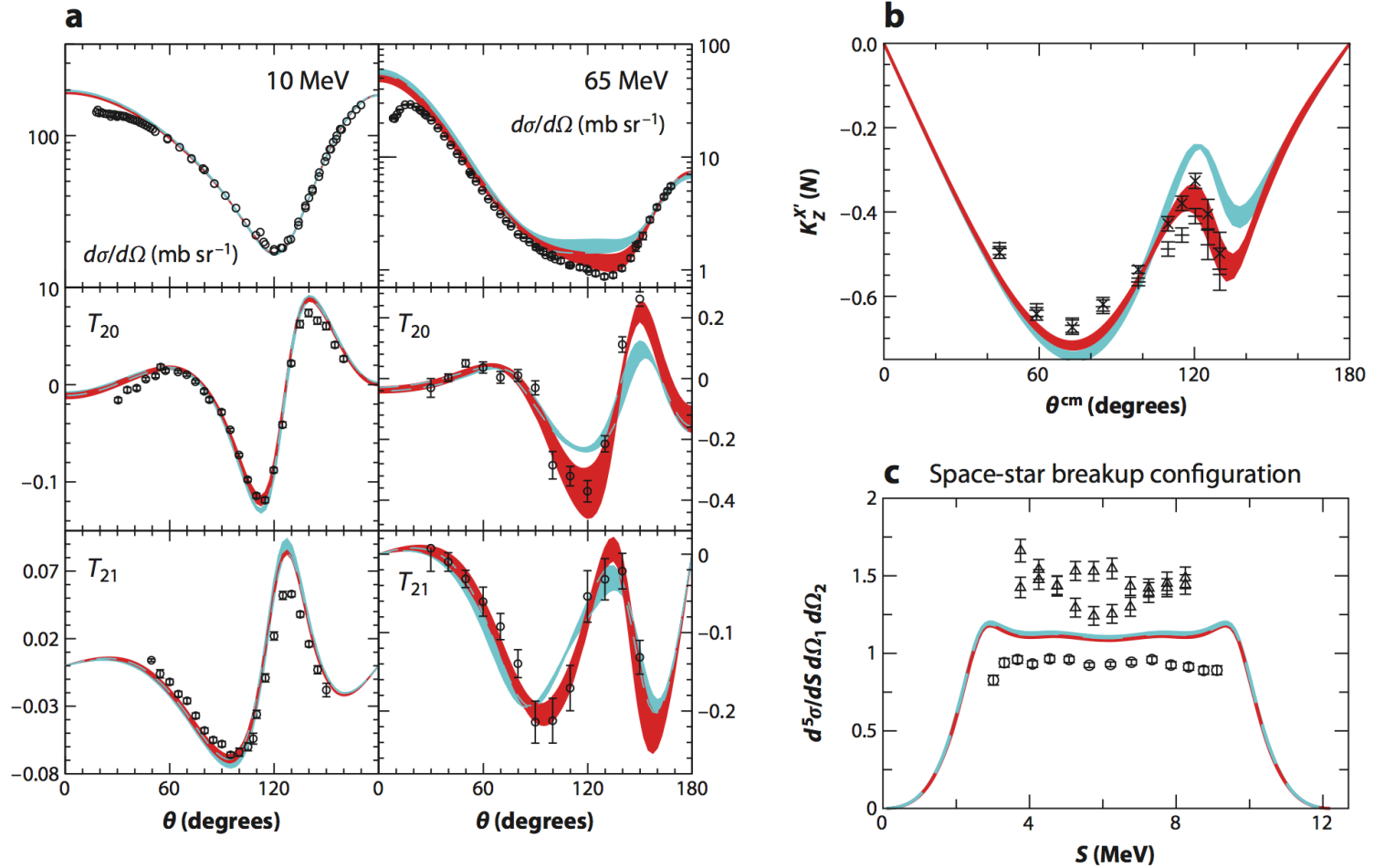


comparable to  
 $N^2LO$  uncertainty



$1/m$  corrections: spin-orbit parts, interesting for  $A_y$  puzzle

# neutron-deuteron scattering at NLO and N<sup>2</sup>LO



**Figure 6**

(a) Differential cross section and tensor analyzing powers  $T_{20}$  and  $T_{21}$  for elastic nucleon-deuteron ( $Nd$ ) scattering at  $E_{\text{lab}}^N = 10$  and  $65$  MeV. (b) The nucleon-to-nucleon polarization transfer coefficient in elastic  $Nd$  scattering at  $E_{\text{lab}}^N = 22.7$  MeV [the proton-deuteron ( $pd$ ) data are from Reference 72]. (c)  $Nd$  breakup cross section in the space-star configuration (upper sets of data,  $nd$ ; lower sets of data,  $pd$ ). The blue and red shaded bands show the results from the chiral effective field theory at next-to-leading order and next-to-next-to-leading order, in order. The precise kinematical description and references to data can be found in Reference 70.

figure from E. Epelbaum and U.-G. Meißner, experiment in a and b is Coulomb corrected p-d



# Importance of 3N forces for light nuclei

Quantum Monte-Carlo calculations [Pieper et al. \(2010\)](#).

based on phenomenological potentials: NN: Argonne  $v_{18}$  + 3N: Illinois-7

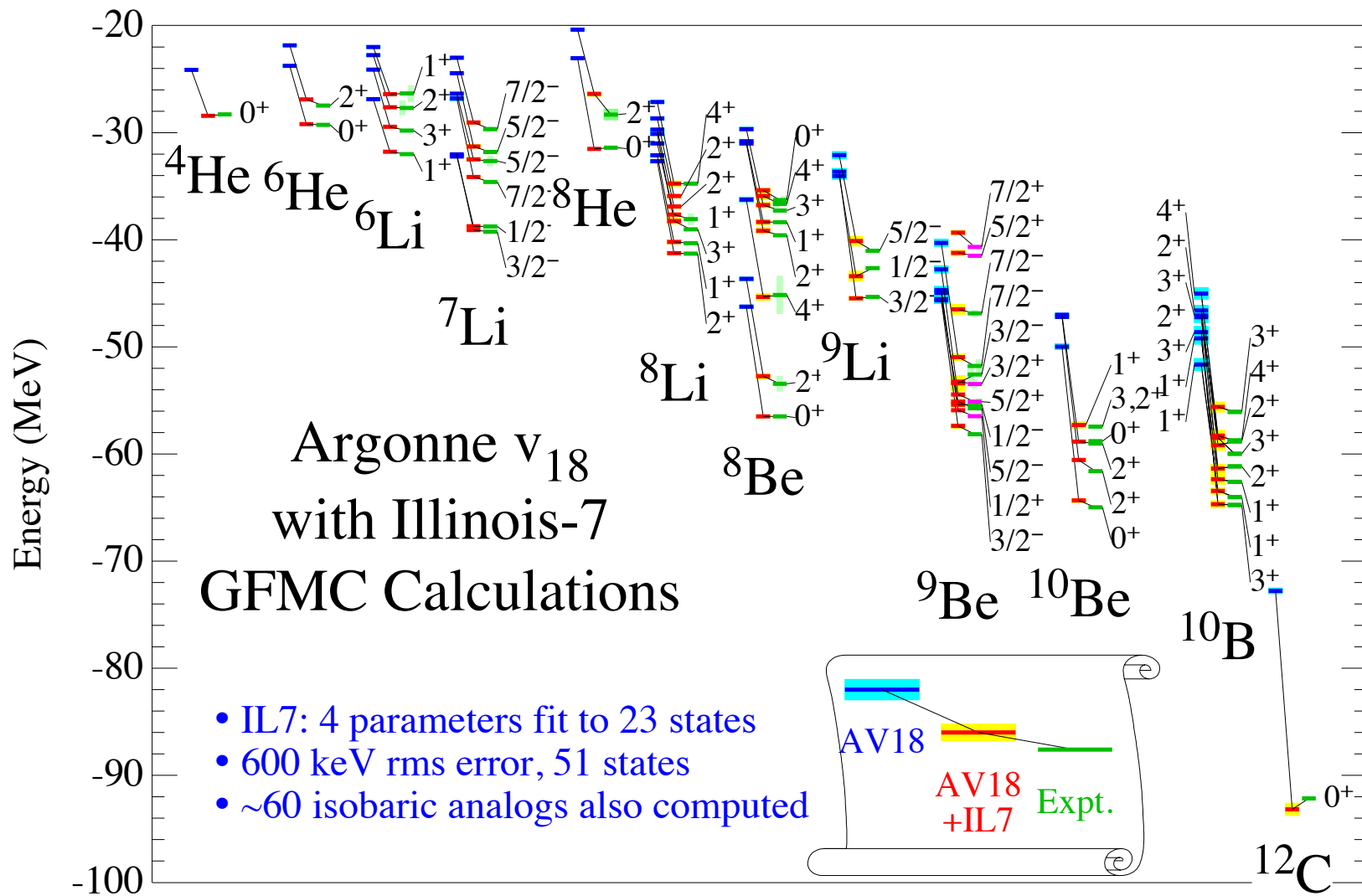
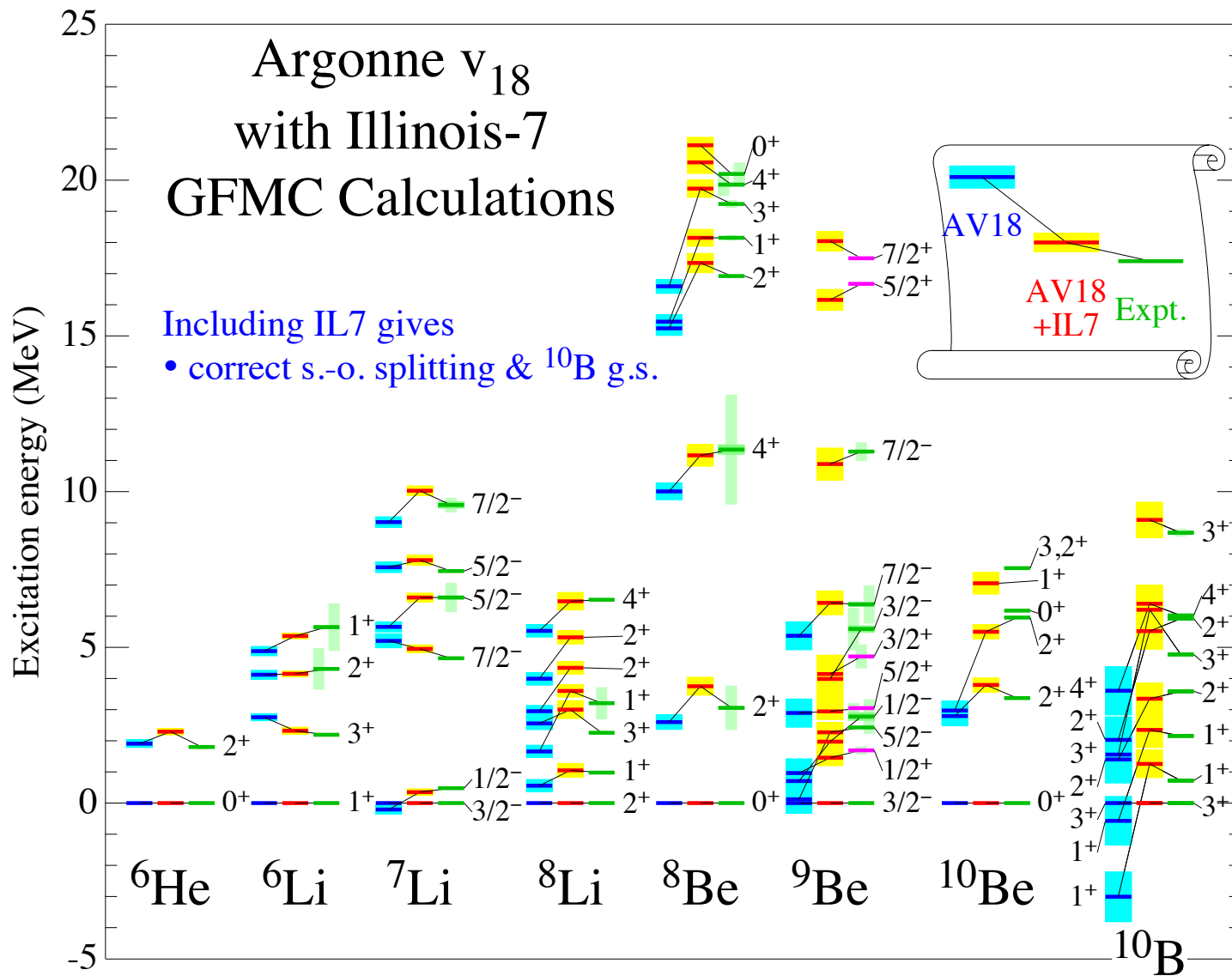


figure from R. Wiringa

# Importance of 3N forces for spectra

spectra too compressed without 3N forces,  $^{10}\text{B}$   $1^+$  vs.  $3^+$

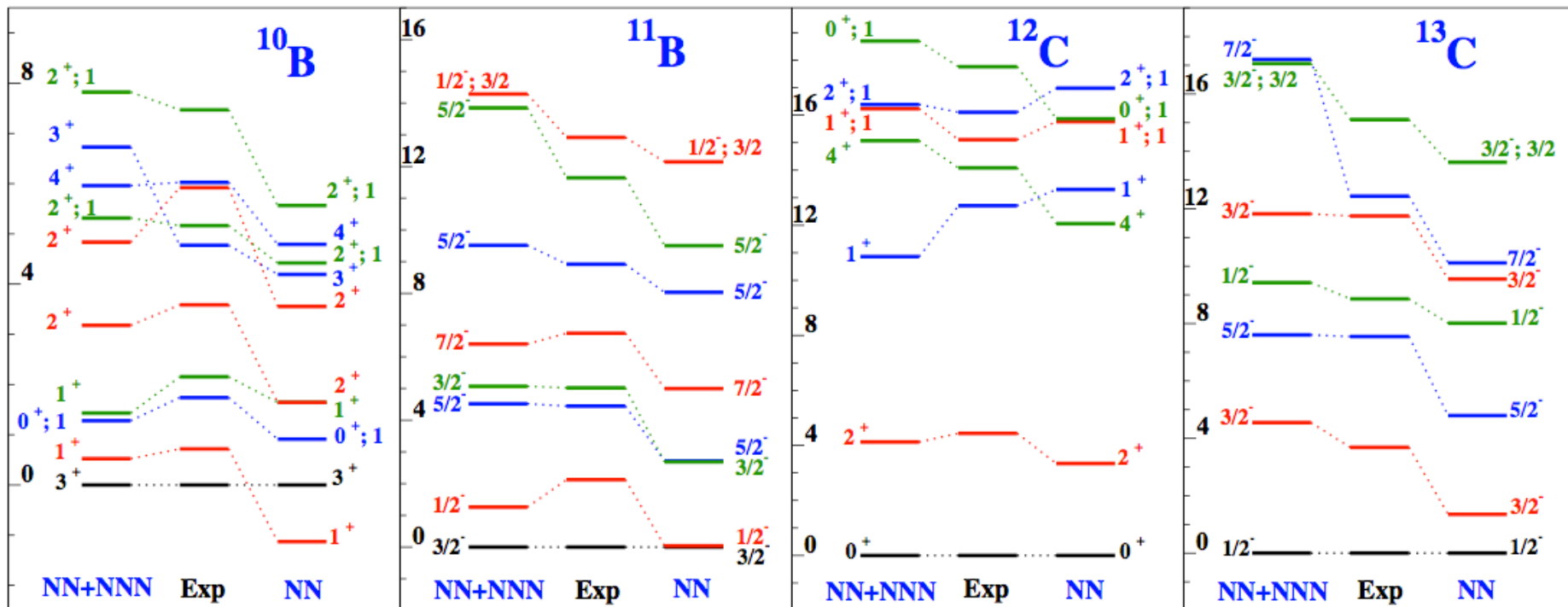


# Importance of 3N forces for spectra of p-shell nuclei

large-basis Hamiltonian diagonalization

using “No-Core Shell Model” Navratil et al., Phys. Rev. Lett. **99**, 042501 (2007).

NN interactions at N<sup>3</sup>LO and 3N interactions at N<sup>2</sup>LO



agreement supports chiral EFT interactions

3N forces: <sup>10</sup>B 1<sup>+</sup> vs. 3<sup>+</sup>, spin-orbit splitting p<sub>3/2</sub>-p<sub>1/2</sub> in <sup>13</sup>C

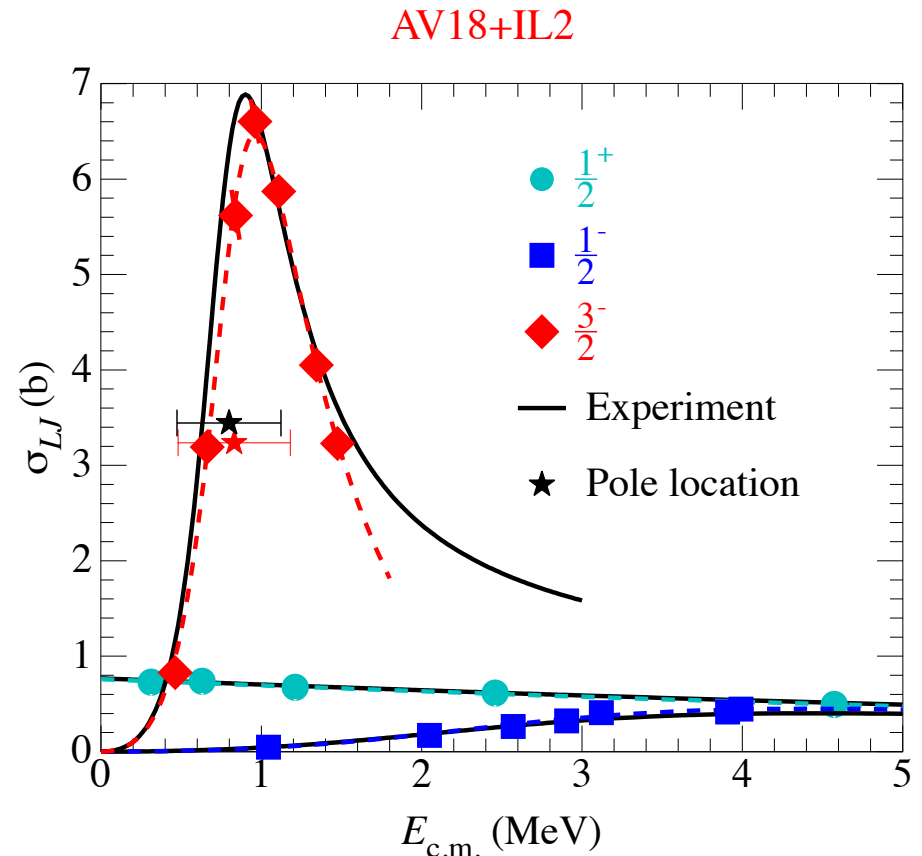
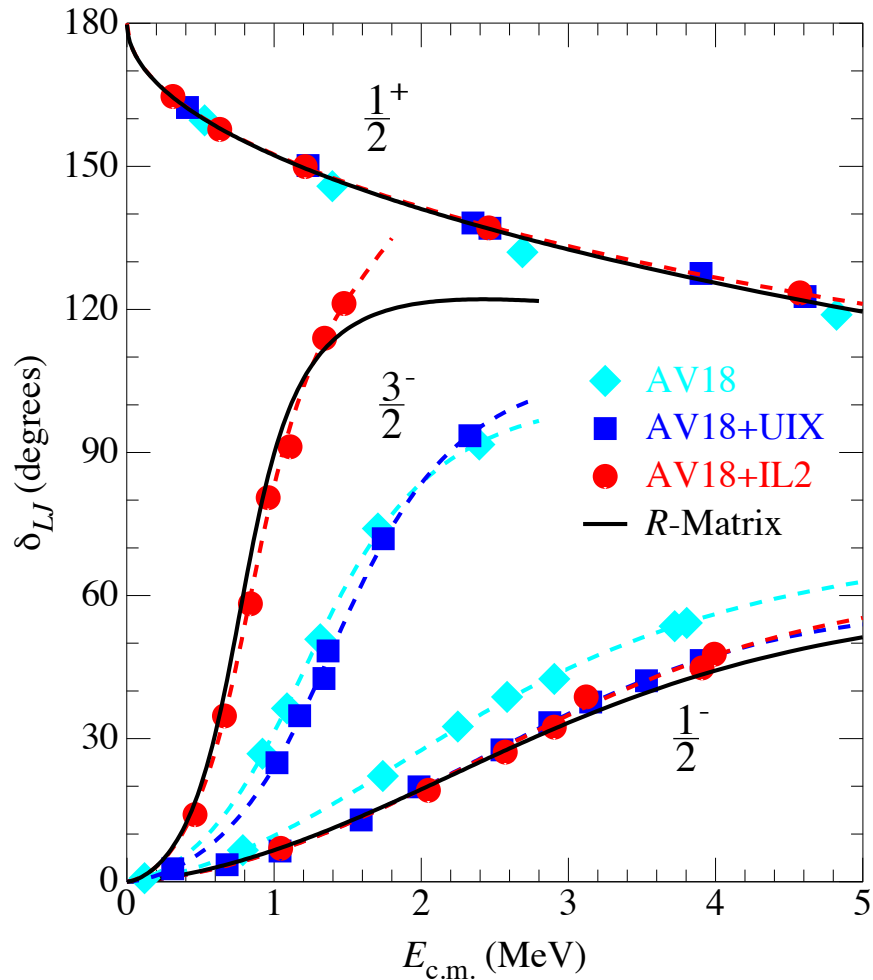
# ${}^5\text{He}$ AS $n+{}^4\text{He}$ SCATTERING

Black curves: Hale phase shifts from  $R$ -matrix analysis up to  $J = \frac{9}{2}$  of data

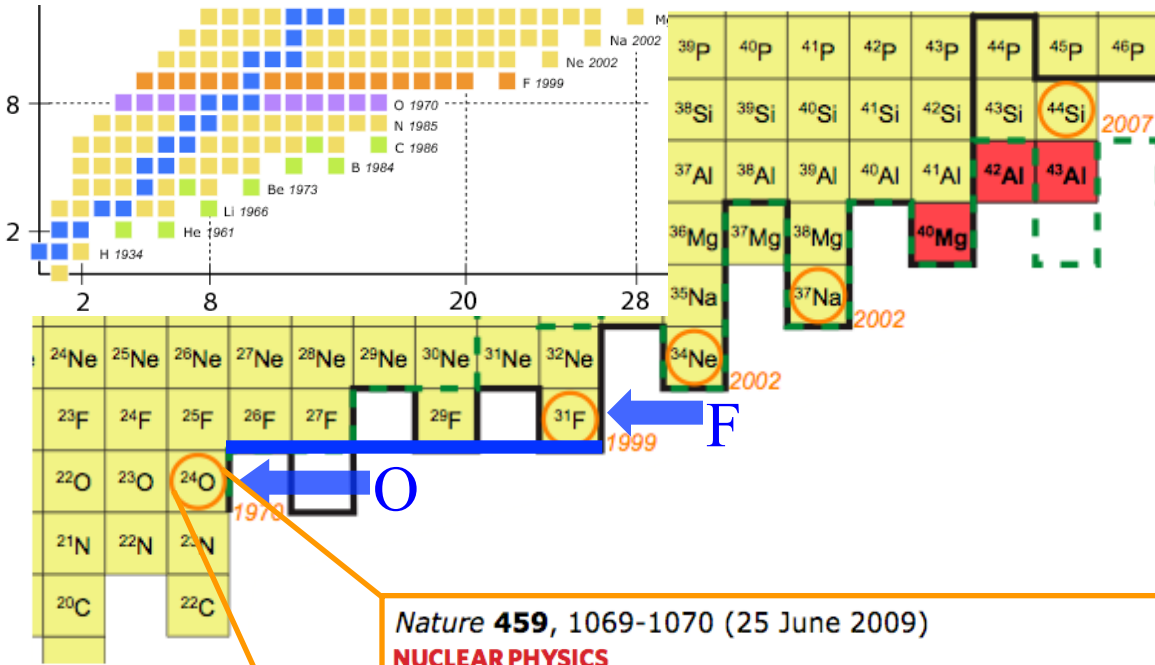
AV18 with no  $V_{ijk}$  underbinds  ${}^5\text{He}(3/2^-)$  & overbinds  ${}^5\text{He}(1/2^-)$

AV18+UIX improves  ${}^5\text{He}(1/2^-)$  but still too small spin-orbit splitting

AV18+IL2 reproduces locations and widths of both  $P$ -wave resonances



# The oxygen anomaly



*Nature* **459**, 1069-1070 (25 June 2009)

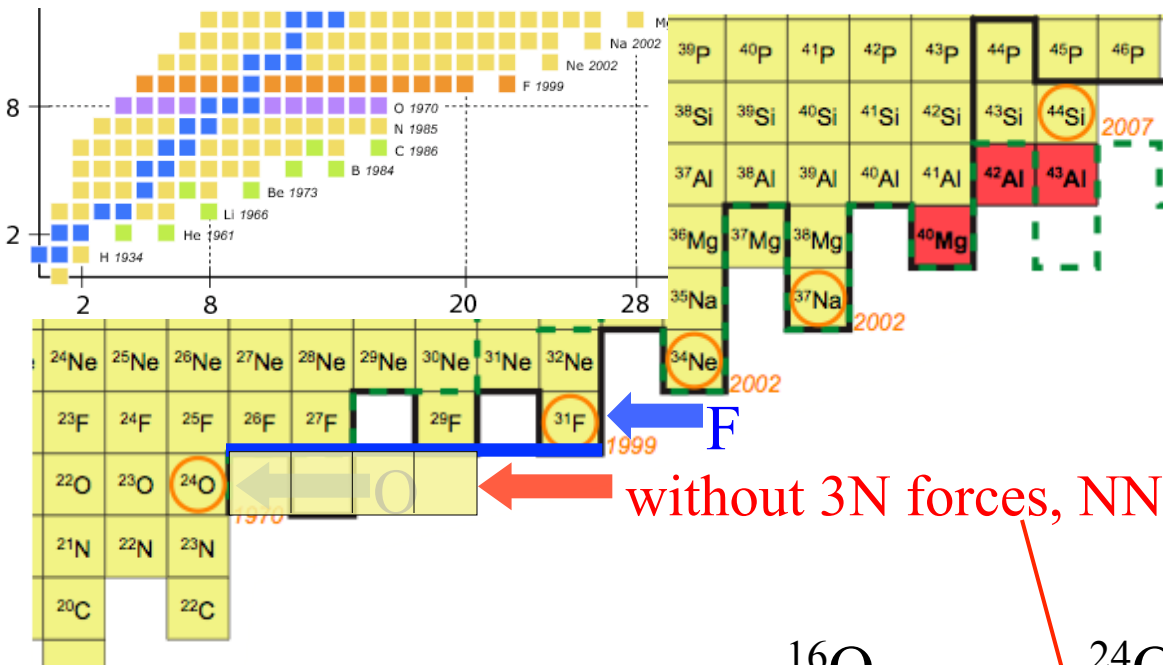
**NUCLEAR PHYSICS**

## Unexpected doubly magic nucleus

Robert V. F. Janssens

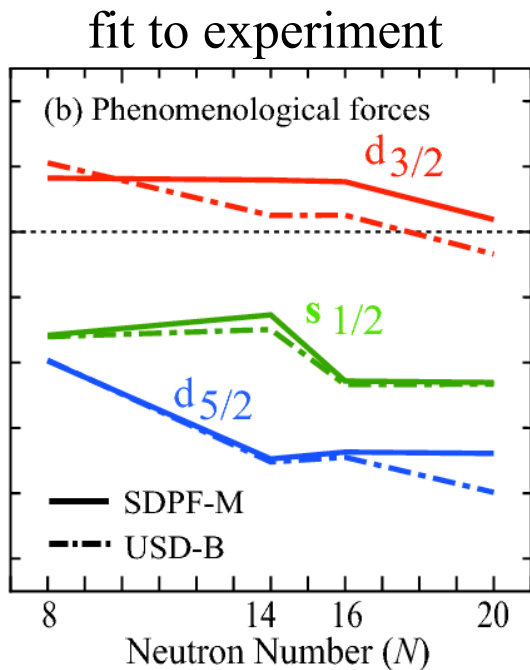
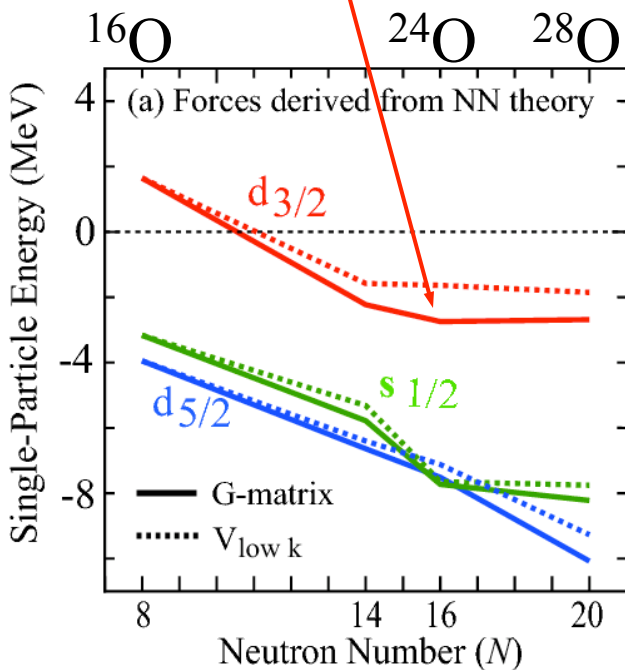
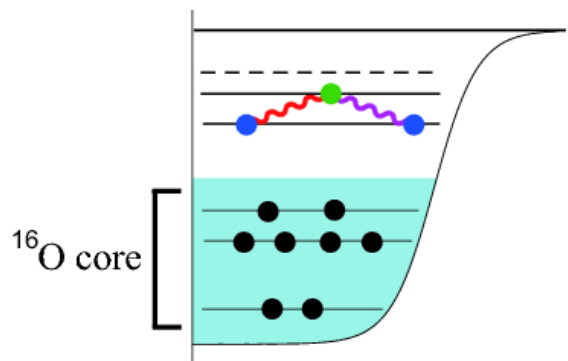
Nuclei with a 'magic' number of both protons and neutrons, dubbed doubly magic, are particularly stable. The oxygen isotope  $^{24}\text{O}$  has been found to be one such nucleus — yet it lies just at the limit of stability.

# The oxygen anomaly - not reproduced without 3N forces



without 3N forces, NN interactions too attractive

many-body theory based on two-nucleon forces:  
drip-line incorrect at  $^{28}\text{O}$



# The shell model - impact of 3N forces

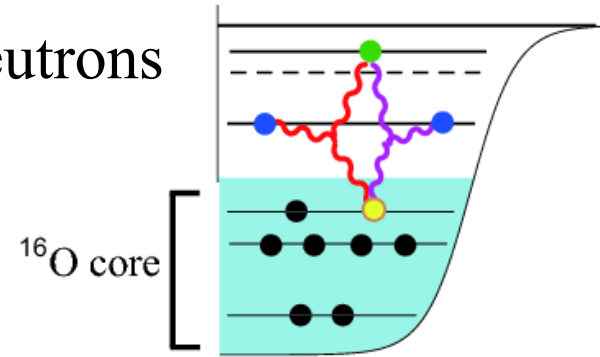
include 'normal-ordered' 2-body part of 3N forces (enhanced by core A)

leads to repulsive interactions between valence neutrons

contributions from residual three valence-nucleon interactions suppressed by  $E_{\text{ex}}/E_{\text{F}} \sim N_{\text{valence}}/N_{\text{core}}$

Friman, AS (2011)

residual 3N amplified in most neutron-rich nuclei



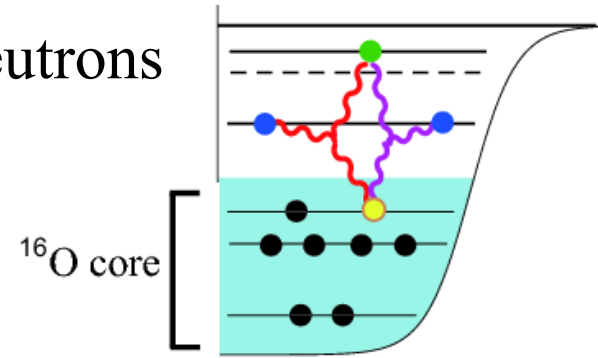
# Oxygen isotopes - impact of 3N forces

include ‘normal-ordered’ 2-body part of 3N forces (enhanced by core A)

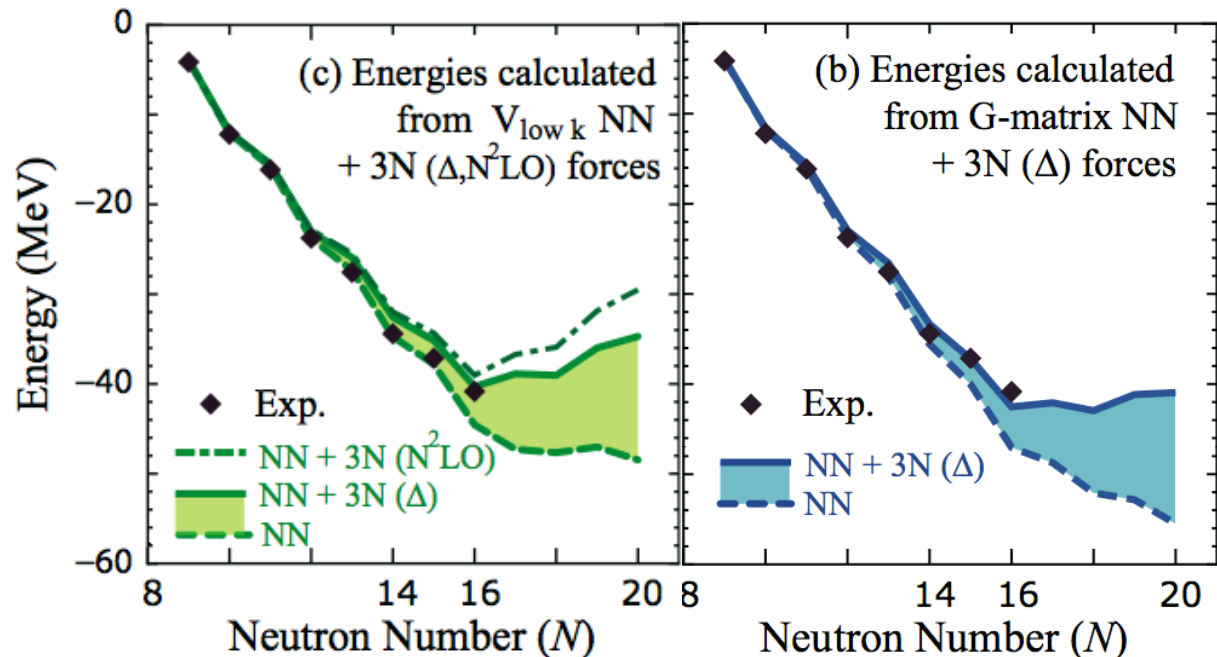
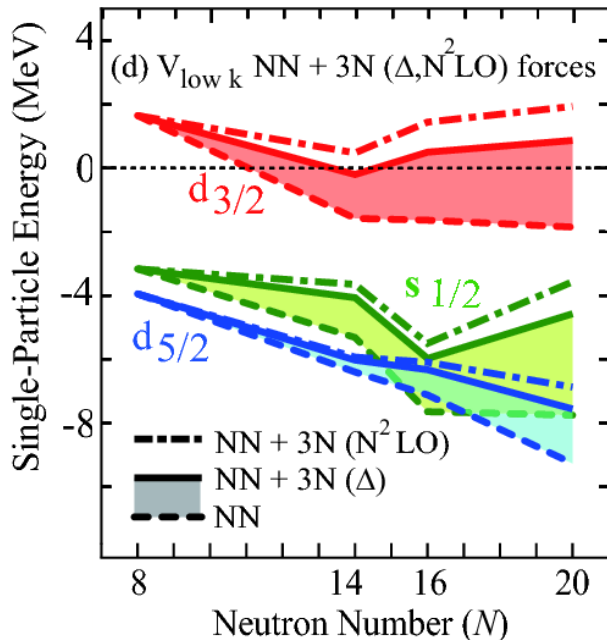
leads to repulsive interactions between valence neutrons

contributions from residual three valence-nucleon interactions suppressed by  $E_{\text{ex}}/E_{\text{F}} \sim N_{\text{valence}}/N_{\text{core}}$

Friman, AS (2011)



$d_{3/2}$  orbital remains unbound from  $^{16}\text{O}$  to  $^{28}\text{O}$



microscopic explanation of the oxygen anomaly Otsuka et al. (2010)



# New ab-initio methods extend reach

impact of 3N forces confirmed in large-space calculations:

Coupled Cluster theory with phenomenological 3N forces [Hagen et al. \(2012\)](#)

In-Medium Similarity RG based on chiral NN+3N [Hergert et al. \(2013\)](#)

Green's function methods based on chiral NN+3N [Cipollone et al. \(2013\)](#)

