

Lecture Three-body forces and halo nuclei

Considers spinless bosons \rightarrow can easily generalize to nucleons $n1, n2, p1, p2$ and distinguishable particles

$$\mathcal{L} = \Psi^\dagger (i\partial_t + \frac{\vec{\nabla}^2}{2m}) \Psi - C_0 (\Psi^\dagger \Psi)^2 - D_0 (\Psi^\dagger \Psi)^3 + \dots$$

S-wave 2-body scattering amplitude

$$f_2(k) = \frac{1}{-\frac{1}{a} - ik} \left[1 - \frac{r_e k^2}{-\frac{1}{a} - ik} + \dots \right] = X + \alpha + \beta\alpha + \dots$$

large $a > 0$
 bound-state pole
 at $k = i/a$
 $E_2 = -\frac{1}{ma^2}$

$r_e \sim \frac{1}{M_{\text{heavy}}}$
 α : correction $\sim \frac{Q}{M_{\text{heavy}}}$

Naive dimensional analysis: 3-body interaction $\sim Q^3$

breaks down for large a because of Λ -divergence in 3-body system

Simplest 3-body process for 3 bosons: boson-dimer scattering

↓
 dimer
 $\equiv X + \alpha + \beta\alpha + \dots$

Boson-dimer scattering equation = Skorniakov-Ter-Martirosian eqn.

for $l=0$ α : which l ?

$$f_3(k, p; E) = \frac{16}{3a} M(k, p; E) + \frac{4}{\pi} \int_0^\Lambda dq q^2 f_3(k, q; E) \frac{M(q, p; E)}{-\frac{1}{a} + \sqrt{3q^2/4 - mE - i\epsilon}}$$

with inhomogeneous term

$$M(k, p; E) = \frac{1}{2kp} \ln \left(\frac{k^2 + kp + p^2 - mE}{k^2 - kp + p^2 - mE} \right) + \frac{H(\Lambda)}{\Lambda^2}$$

and $E = \frac{3k^2}{4m} - \frac{1}{ma^2} \sim Q^2$



on-shell

$$k \cot \delta = \frac{1}{f_3(k, k, E)} + ik$$

with $D_0(\Lambda) = -4m C_0(\Lambda)^2 \frac{H(\Lambda)}{\Lambda^2}$

For $H=0$ (no 3-body interaction): integral equation has no unique solution for $\Lambda \rightarrow \infty$
 regularized integral equation has unique solution, but 3-body observables show strong dependence on Λ !

\Rightarrow remove cutoff dependence by tuning coupling $D_0(\Lambda) \sim H(\Lambda)$

Bedaque, Hammer, van Kolck (1999) $H(\Lambda)$ has limit cycle running

$$H(\Lambda) = \frac{\cos[s_0 \ln(\Lambda/\Lambda_*) + \arctan s_0]}{\cos[s_0 \ln(\Lambda/\Lambda_*) - \arctan s_0]}$$

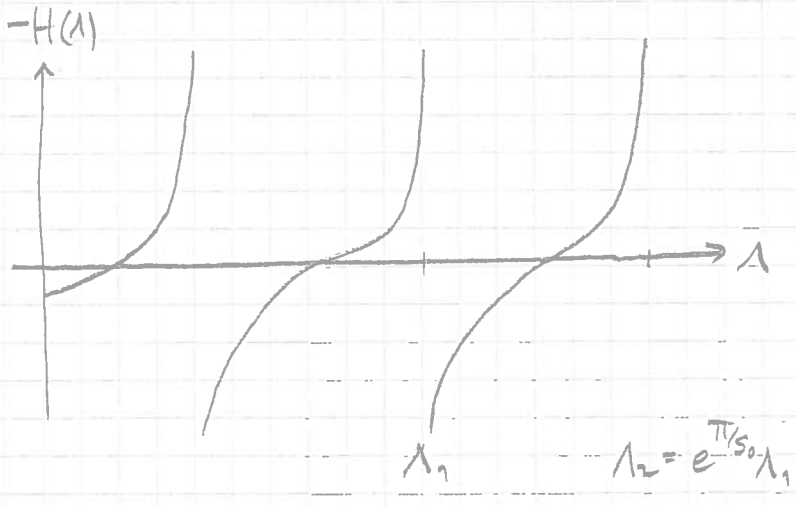
with $s_0 \approx 1.00624$ is a transcendental number

Λ_* dimensional 3-body parameter (\rightarrow dimensional transmutation)
 fix to reproduce our low-energy data

\Rightarrow discrete scale invariance by scaling factor in momentum $e^{\pi/s_0} \approx 22.7$

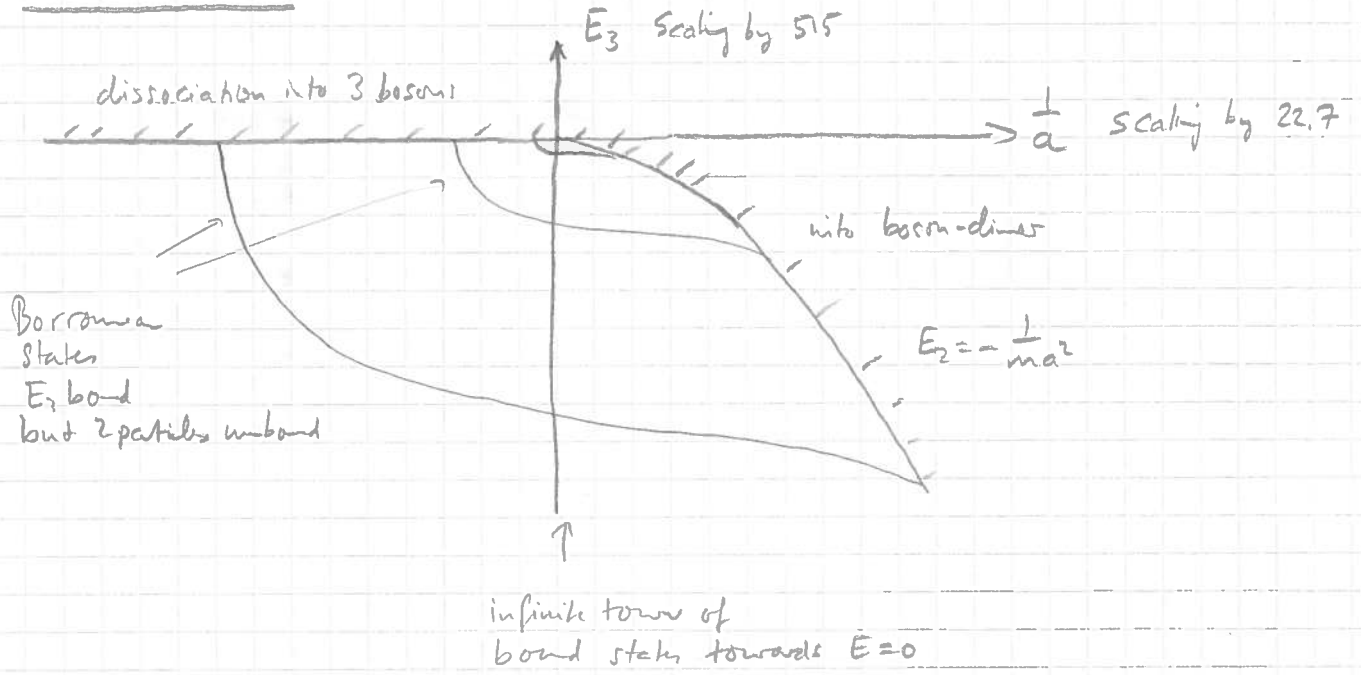
in energy $\sim e^{2\pi/s_0} \approx 515$

Scale invariance in unitary limit is broken by Λ_* leaves 3-body observables unchanged



⇒ discrete scaling symmetry leads to the Efimov effect in 3-body system
 → geometric spectrum of 3-body bound states

Efimov spectrum



fix λ_* to one low-energy E_3