Nuclei as Laboratories: Nuclear Tests of Fundamental Symmetries



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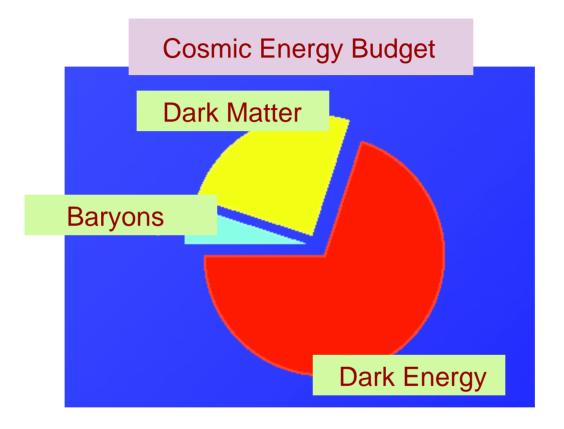
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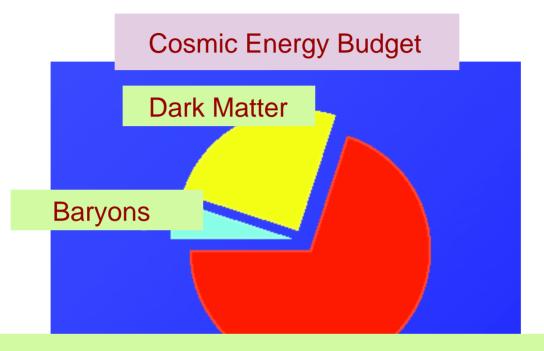
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Nuclear Science



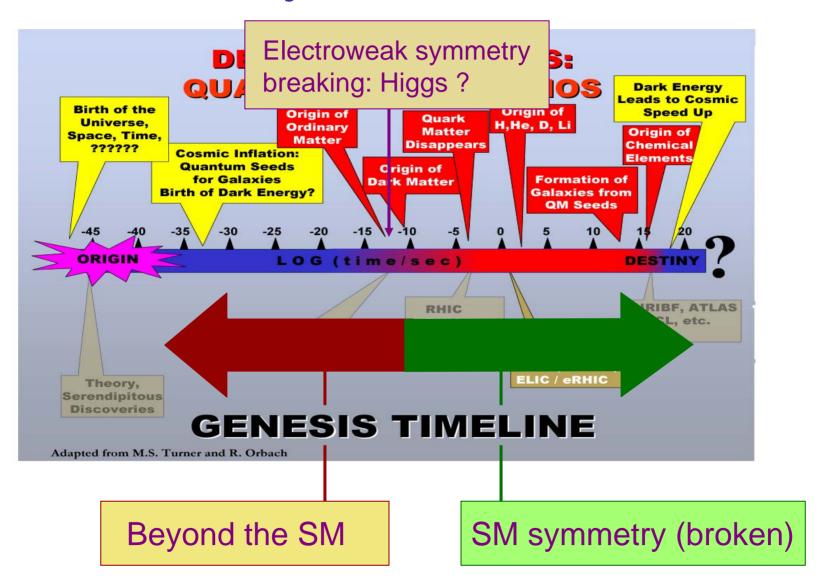
The mission: Explain the origin, evolution, and structure of the baryonic matter of the Universe

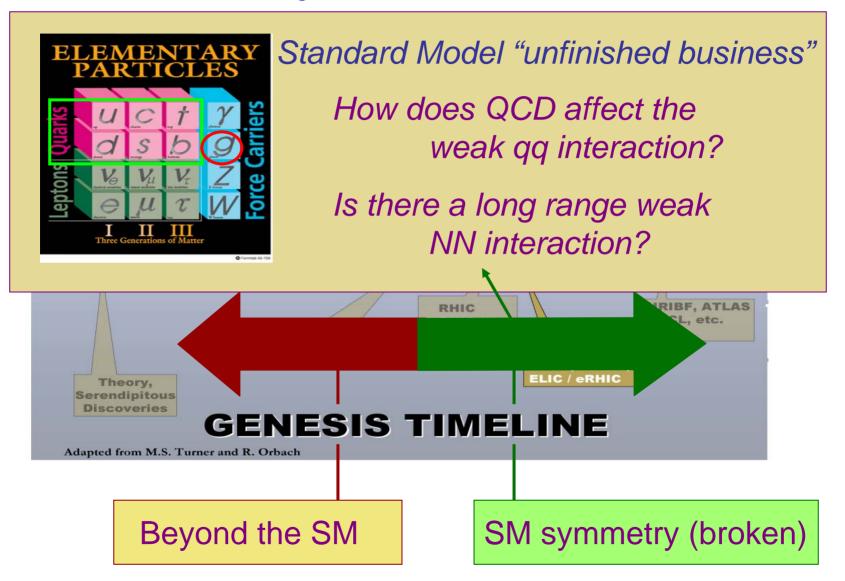
Nuclear Science

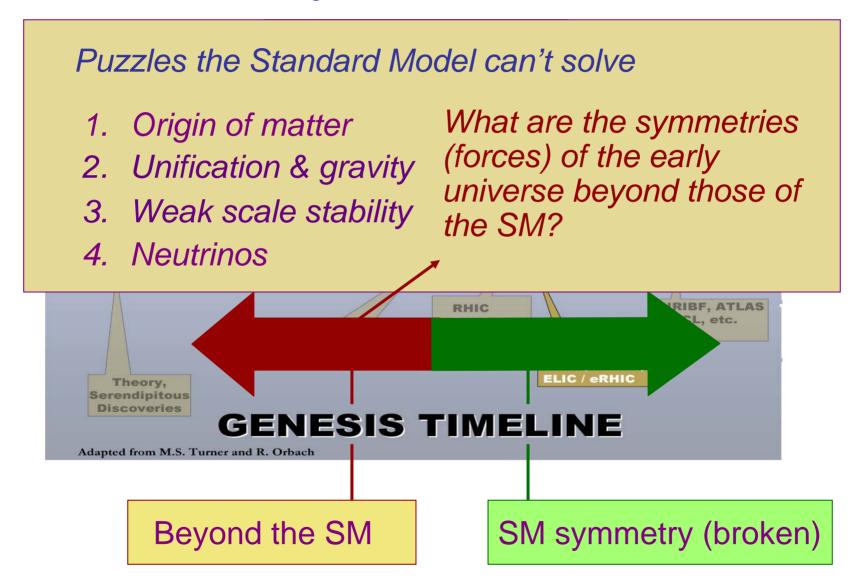


Three frontiers:

- Fundamental symmetries & neutrinos
- Nuclei and nuclear astrophysics
- QCD





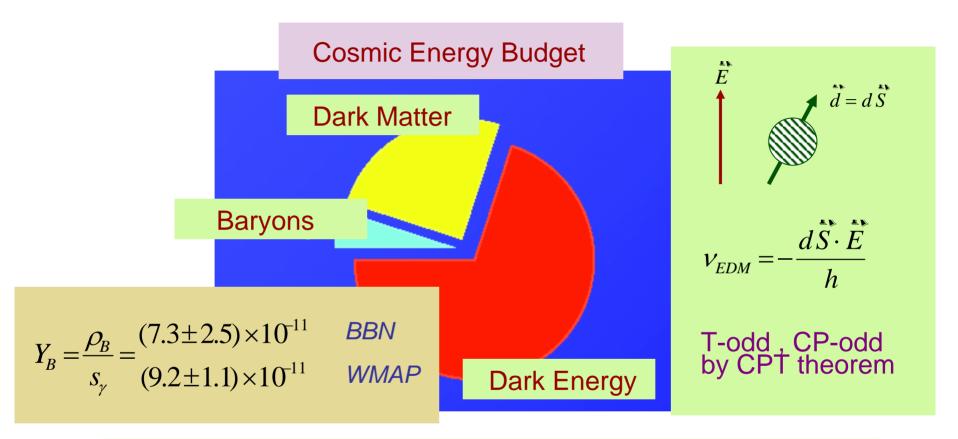


What are the new fundamental symmetries?

- Why is there more matter than antimatter in the present universe?
 - Electric dipole moment searches
- What are the unseen forces that disappeared from view as the universe cooled?
 - Precision electroweak: weak decays, scattering, LFV
- What are the masses of neutrinos and how have they shaped the evolution of the universe?

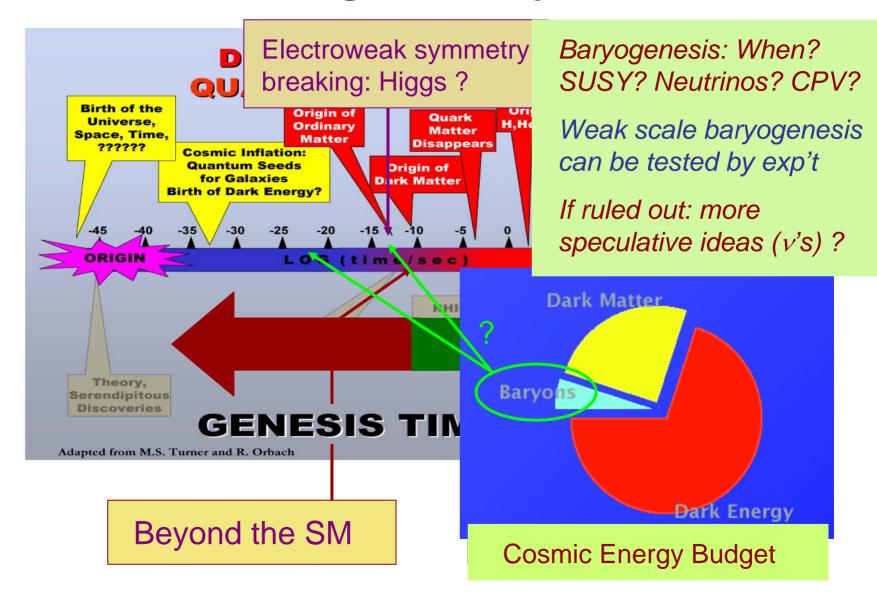
Neutrino oscillations, $0\nu\beta\beta$ -decay, θ_{13} , ...

What is the origin of baryonic matter?



What are the quantitative implications of new EDM experiments for explaining the origin of the baryonic component of the Universe?

What is the origin of baryonic matter?



EW Baryogenesis: Standard Model

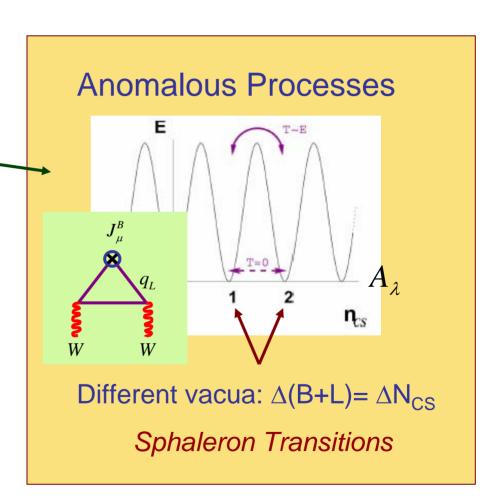
Sakharov:

Weak Scale Baryogenesis

- B violation
- C & CP violation
- Nonequilibrium dynamics

Sakharov, 1967

Kuzmin, Rubakov, Shaposhnikov McLerran,...



EW Baryogenesis: Standard Model

Shaposhnikov

Weak Scale Baryogenesis

- B violation
- C & CP violation
- Nonequilibrium

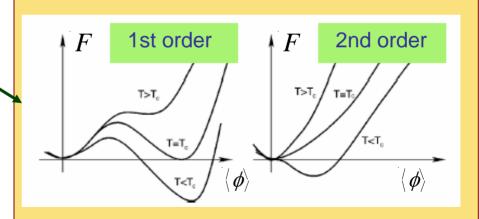
dynamics Sakharov, 1967

- CP-violation too weak
- EW PT too weak

$$J = s_{12} s_{13} s_{23} c_{12} c_{13}^{2} c_{23} \sin \delta_{13}$$

$$= (2.88 \pm 0.33) \times 10^{-5}$$

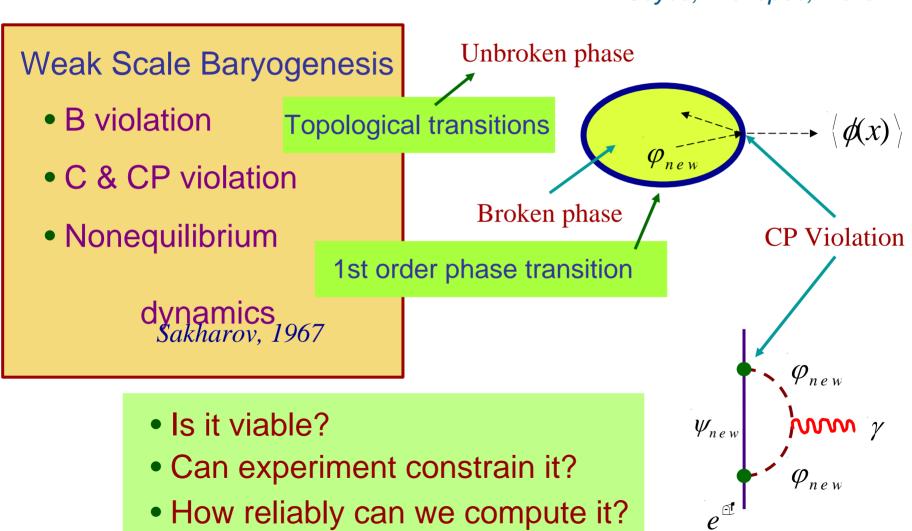
$$\frac{m_{t}^{4}}{M_{W}^{4}} \frac{m_{b}^{4}}{M_{W}^{4}} \frac{m_{c}^{2}}{M_{W}^{2}} \frac{m_{s}^{2}}{M_{W}^{2}} \approx 3 \times 10^{-13}$$



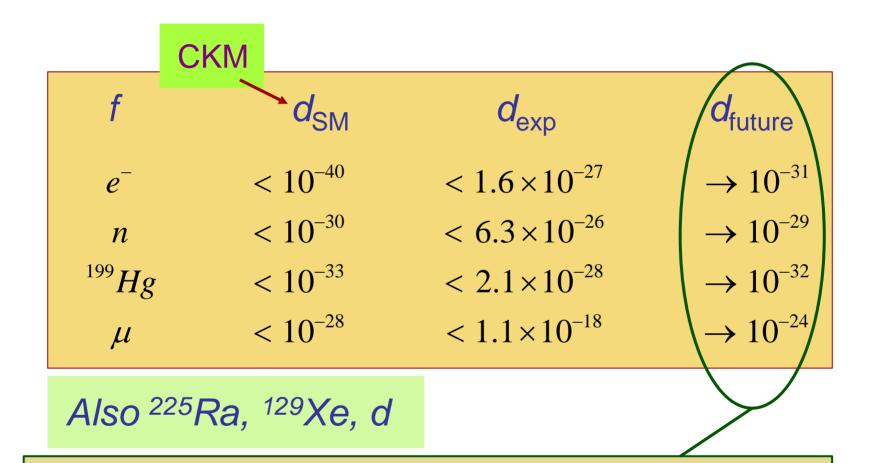
Increasing m_h

Baryogenesis: New Electroweak Physics

90's: Cohen, Kaplan, Nelson Joyce, Prokopec, Turok

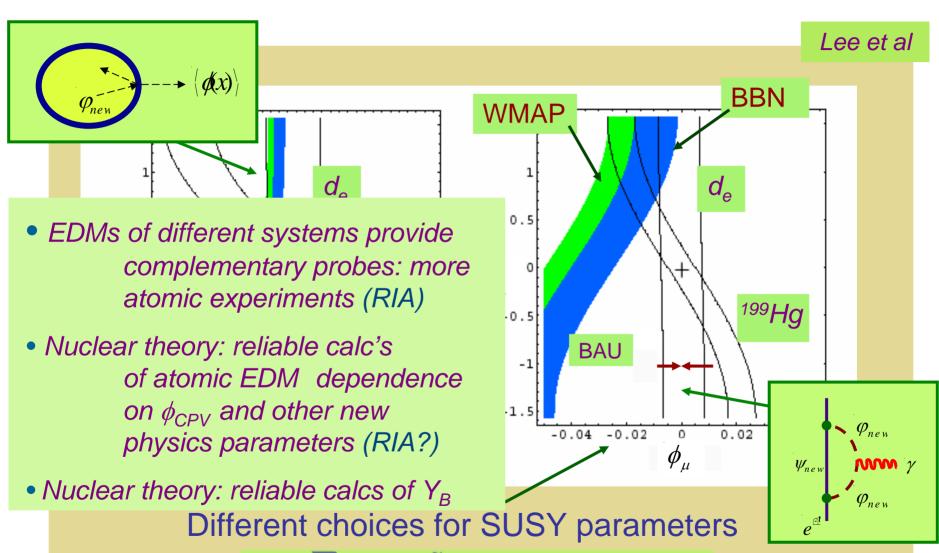


EDM Probes of New CP Violation

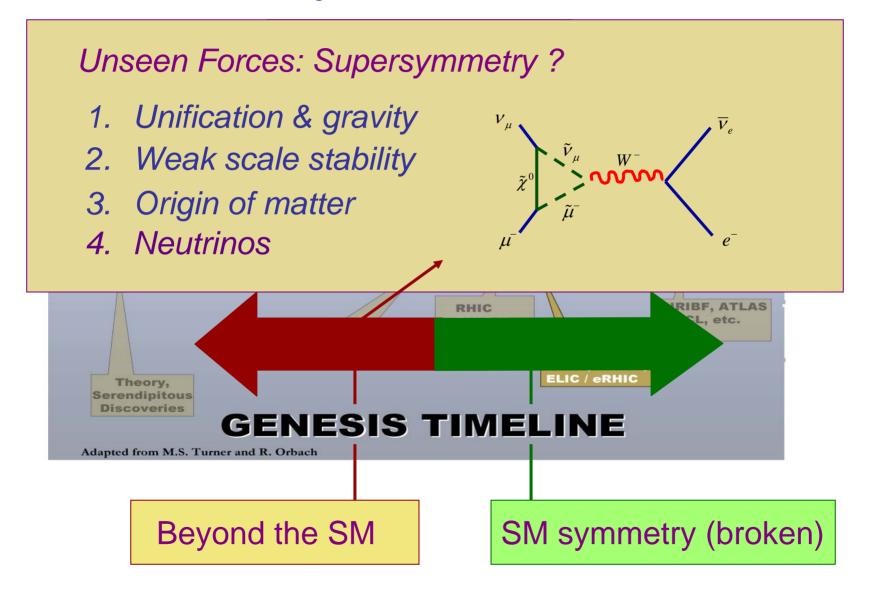


If an EDM is seen, can we identify the new physics?

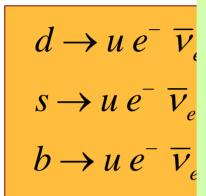
EDM constraints & SUSY CPV

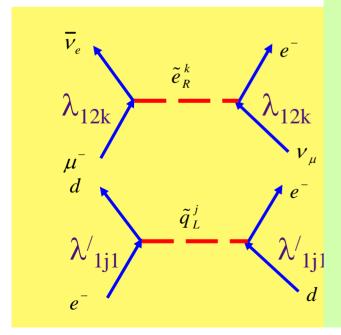


$$Y_B = \sum_k F_k(g_i, \tilde{M}_i; T, v_w, L_w, \dots) \sin \phi_k$$



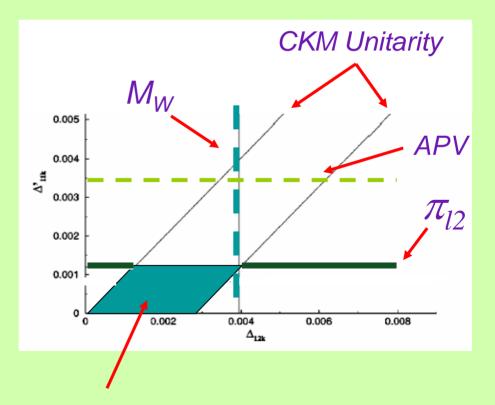
Weak decays & new physics





R Parity Violation Kuryl

Kurylov, R-M, Su



No long-lived LSP or SUSY DM

Weak decays

$$\frac{G_F^{\beta}}{G_F^{\mu}} = |V_{ud}| \left(1 + \Delta r_{\beta} - \Delta r_{\mu}\right)$$

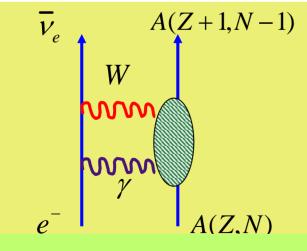
β-decay

$$n \rightarrow p e^{-} \overline{\nu}_{e}$$

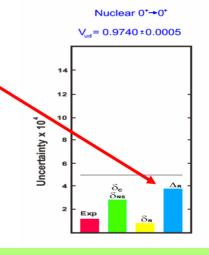
$$A(Z,N) \rightarrow A(Z-1,N+1) e^{+} \nu_{e}$$

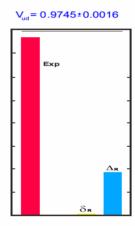
$$\pi^{+} \rightarrow \pi^{0} e^{+} \nu_{e}$$

SM theory input

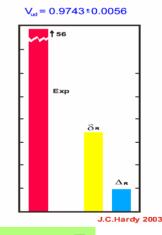


Nuclear structure effects?





Neutron



Pion beta decay

$$M_{W\gamma} = \frac{G_F}{\sqrt{2}} \frac{\hat{\alpha}}{2\pi} \left[\ln \left(\frac{M_Z^2}{\Lambda^2} \right) + C_{\gamma W}(\Lambda) \right]$$

Weak decays & new physics

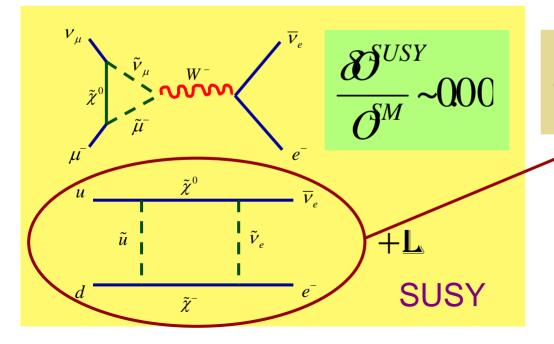
Correlations

$$d \to u e^- \overline{V}_e$$

$$s \to u e^- \overline{V}_e$$

$$b \to u e^- \overline{V}_e$$

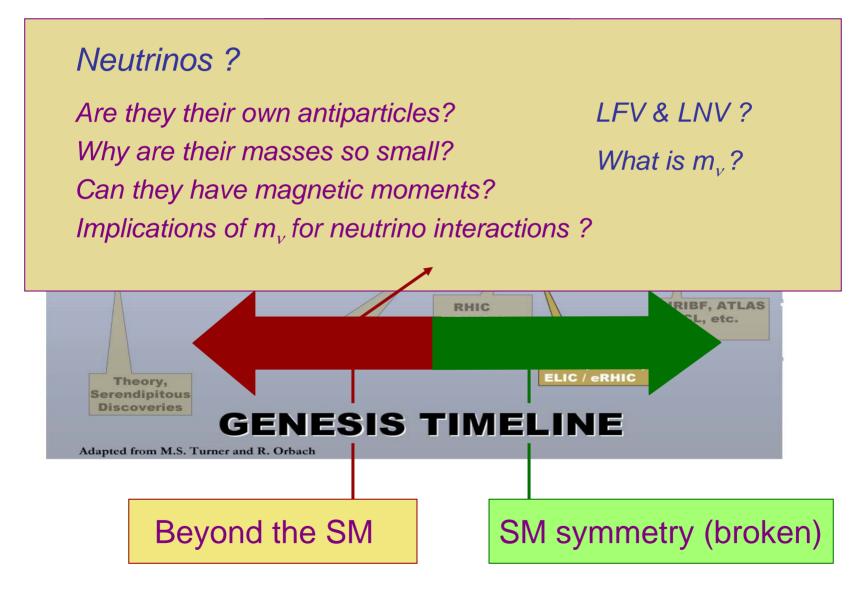
$$\begin{pmatrix} u & c & t \end{pmatrix} \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$



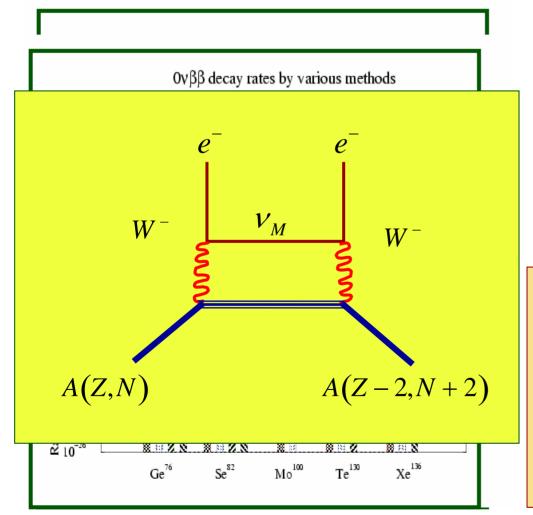
$$dW \propto 1 + a \frac{\overrightarrow{p_e} \cdot \overrightarrow{p_v}}{E_e E_v} + A \overrightarrow{\sigma_n} \cdot \frac{\overrightarrow{p_e}}{E_e} + \mathbf{L}$$

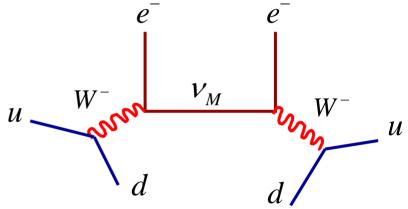
Non (V-A) x (V-A) interactions: m_e/E

 β -decay at RIA?



0ν ββ - decay probes the charge conjugation properties of the neutrino





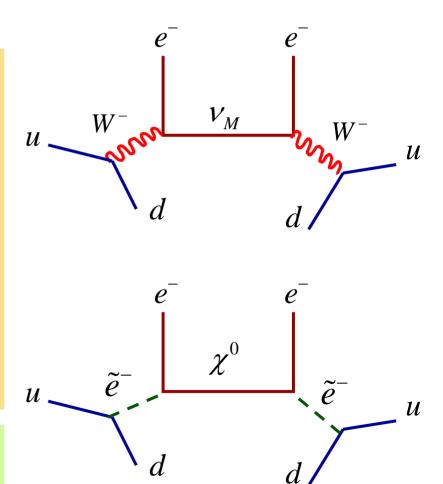
Light $v_{\rm M}$: Nuclear matrix elements difficult to compute

$$\langle m_{\nu} \rangle^{EFF} = \sum_{k} |U_{ek}|^2 m_k e^{2i\delta}$$

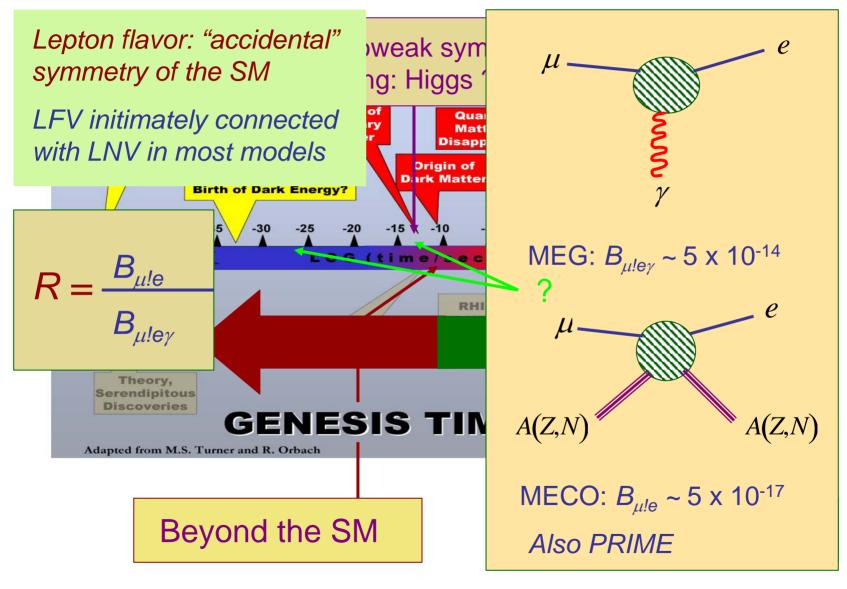
$0v \beta\beta$ - decay: heavy particle exchange

$$egin{aligned} rac{M_H}{M_L} \sim & rac{M_W^4 ar{k}^2}{\Lambda_{eta eta}^5 m_{eta eta}} \sim O(1) \ m_{eta eta} \sim & 0.1 \; eV \ \Lambda_{eta eta} \sim & 1 \; TeV \ ar{k} \sim & 50 \; MeV \end{aligned}$$

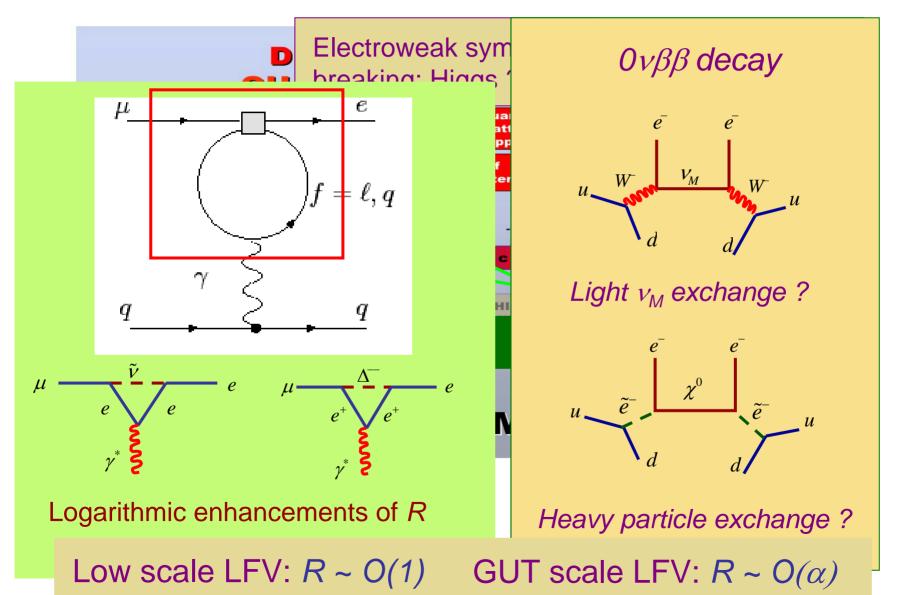
How do we compute & separate heavy particle exchange effects?



LF and LN: symmetries of the early universe?

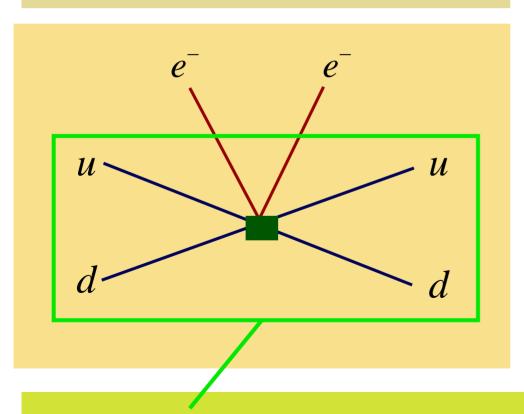


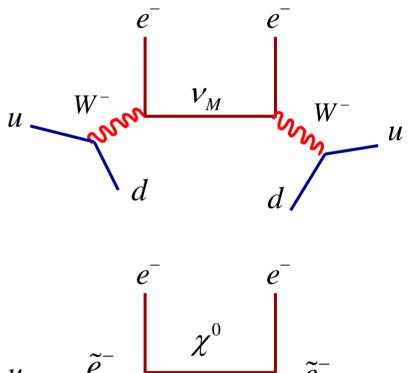
LF and LN: symmetries of the early universe?



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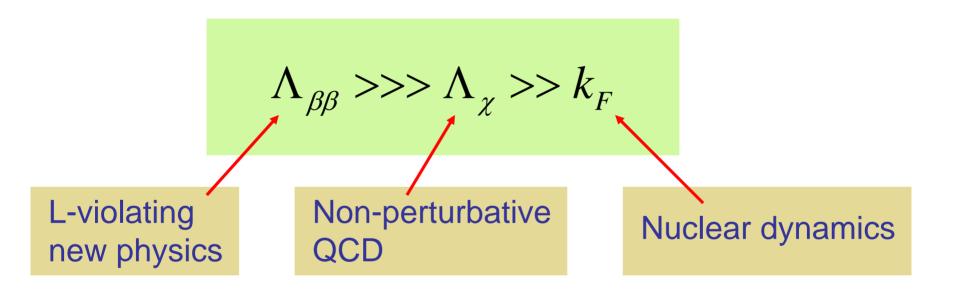




4 quark operator, as in hadronic PV

$0v \beta\beta$ - decay: effective field theory

We have a clear separation of scales



$0v \beta\beta$ - decay in effective field theory

Operator classification

$$\mathcal{L}(q,e)$$
! $\mathcal{L}(\pi,N,e)$

$$\mu = M_{WEAK}$$

 $\mu = M_{HAD}$





Spacetime & chiral transformation properties

$0v \beta\beta$ - decay in effective field theory

Operator classification

$$\mu = M_{WEAK}$$

$$\mathcal{L}(q,e) = \frac{G_F^2}{\Lambda_{\beta\beta}} \sum_{j=1}^{14} C_j(\mu) \, \hat{O}_j^{++} \, \overline{e} \Gamma_j e^c + h.c.$$

$$\hat{O}_{1+}^{ab} = \overline{q}_L \gamma^\mu \tau^a q_L \ \overline{q}_R \gamma_\mu \tau^b q_R$$

$$0$$
ν ββ - decay: $a = b = +$

$\partial v \beta \beta$ - decay in effective field theory

Operator classification

$$\mu = M_{WEAK}$$

$$\hat{O}^{ab}_{1+} = \overline{q}_L \gamma^{\mu} \tau^a q_L \ \overline{q}_R \gamma_{\mu} \tau^b q_R$$

Chiral transformations: $SU(2)_L \times SU(2)_R$

$$\begin{array}{ccc} q_L \to L \, q_L & L \\ q_R \to R \, q_R & R \end{array} = \exp \left(i \, \theta_L \cdot \frac{\tau}{2} \, P_L \right) & \hat{O}_{1+}^{ab} \in (3_L, 3_R) \end{array}$$

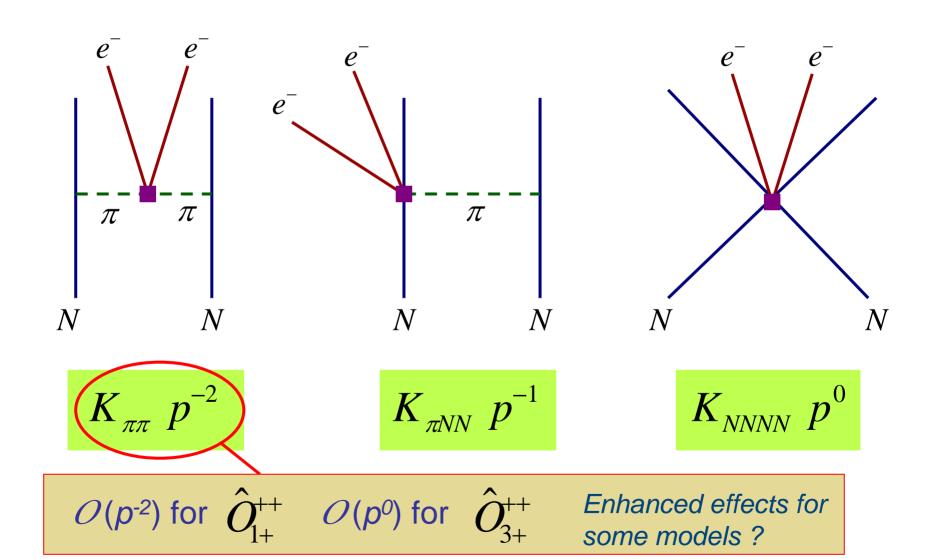
$$\hat{O}^{ab}_{1+} \in (3_L, 3_R)$$

Parity transformations: $q_i \ \ q_R$

$$0ν ββ - decay: $a = b = +$$$

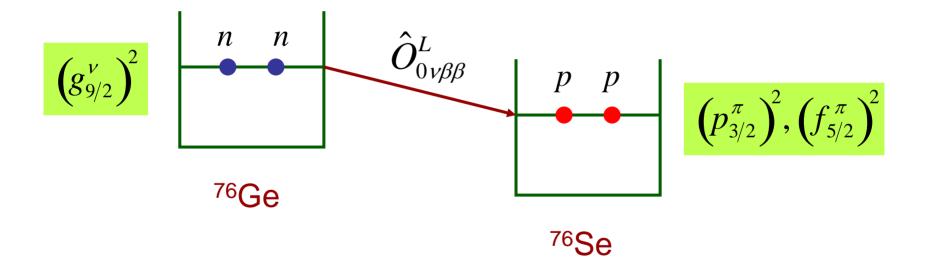
$$\hat{O}_{1+}^{\scriptscriptstyle ++} \leftrightarrow \hat{O}_{1+}^{\scriptscriptstyle ++}$$

$\partial v \beta \beta$ - decay in effective field theory



An open question

Is the power counting of operators sufficient to understand weak matrix elements in nuclei?



$$L = 0, K, 9$$

$$L' = 0, K, 5$$

An open question

Is the power counting of operators sufficient to understand weak matrix elements in nuclei?

$$L = 0, K, 9$$

$$\hat{O}^L_{0
uetaeta}$$

$$L' = 0, K, 5$$

naive

$$M_{fi} \sim p^0$$

$$\mathbf{L} = \mathbf{L'} = \mathbf{0}$$

$$\hat{O}_{0
uetaeta}^{L=0}$$

$$M_{fi} \sim p^2$$

$$L = 2, L' = 0$$

$$\hat{O}_{0
uetaeta}^{L=2}$$

$$M_{fi} \sim p^2$$

$$L = 0, L' = 2$$

$$\hat{O}_{0
uetaeta}^{L=2}$$

$$M_{fi} \sim p^4$$

$$L=4, L'=0$$

$$\hat{O}_{0
uetaeta}^{L=4}$$
 et

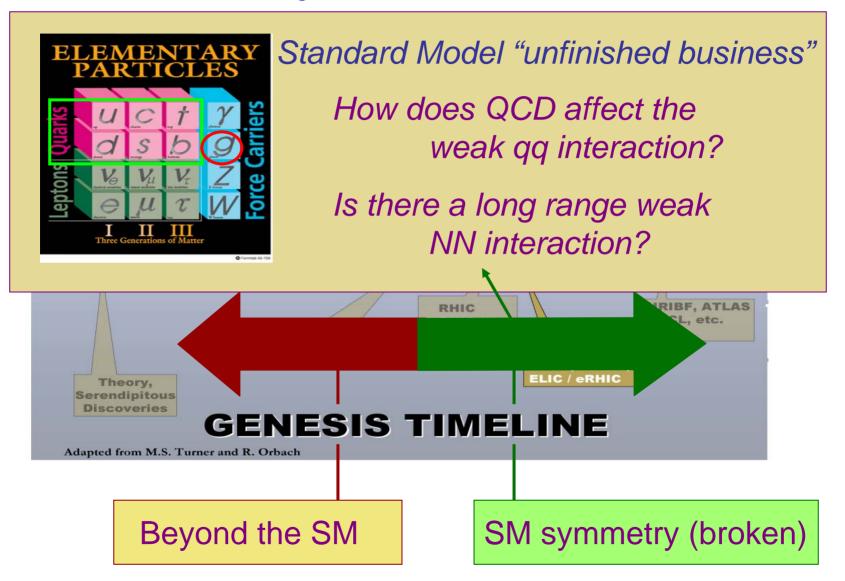
An open question

Complications:

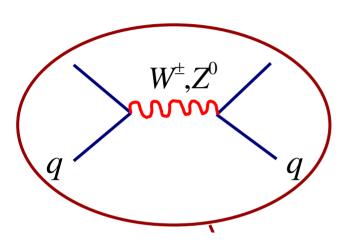
- Bound state wavefunctions (e.g., h.o.) don't obey simple power counting
- Configuration mixing is important in heavy nuclei

Is the power counting of operators sufficient to understand weak matrix elements in nuclei?

- More theoretical study required (RIA)
- Hadronic PV may provide an empirical test

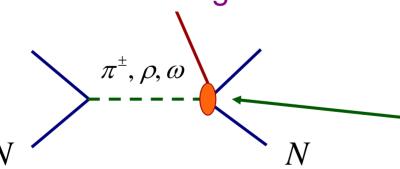


The weak qq force is short range



 $\lambda_{W,Z}$ ~0.00 2m < R_{COR}

Meson-exchange model

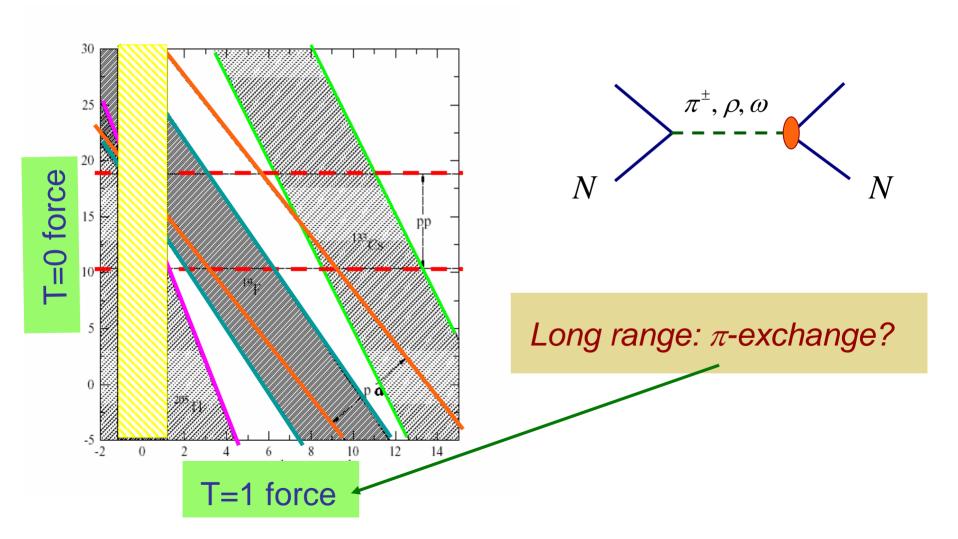


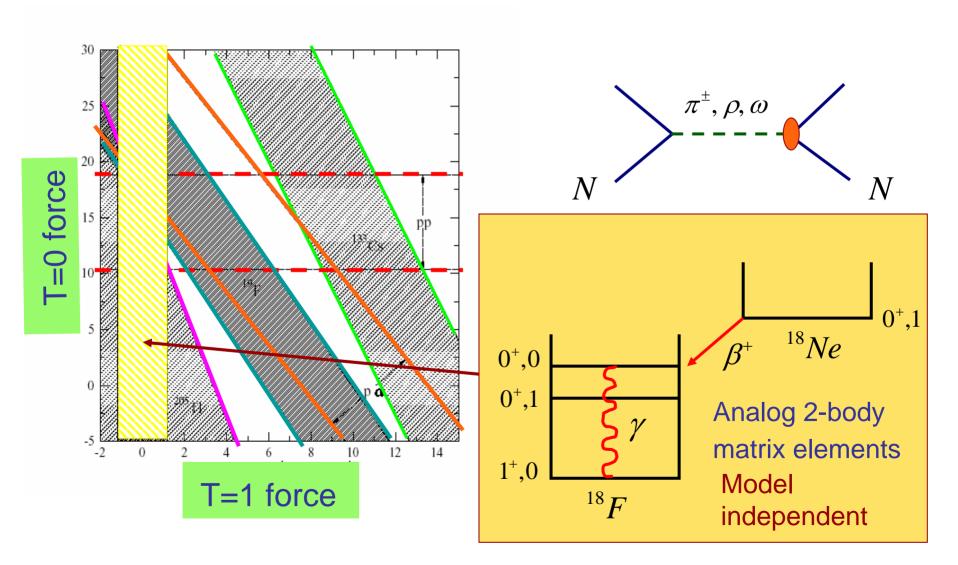
Seven PV mesonnucleon couplings

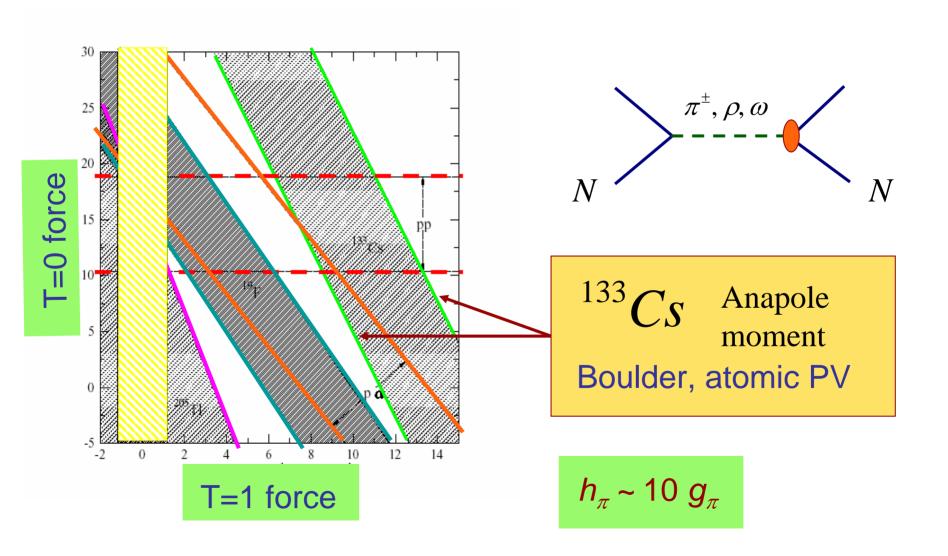
$$h_{\pi}^{1}, h_{\rho}^{0,1,2}, h_{\omega}^{0,1}, h_{\rho}^{1}'$$

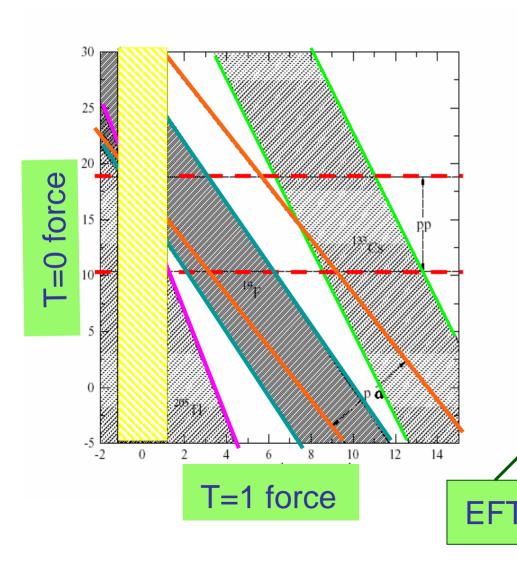
Use parity-violation to filter out EM & strong interactions

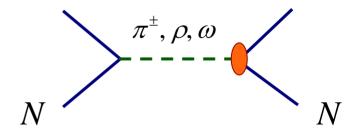
Desplanques, Donoghue, &Holstein (DDH)







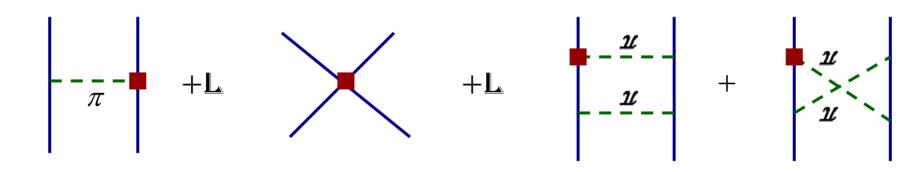




- Problem with expt's
- Problem with nuc th'y
- Problem with model
- No problem (1σ)

Hadronic PV: Effective Field Theory

PV Potential



Long Range

Short Range

Medium Range

$$h^1_{\pi\!N\!N}$$

$$k_{\pi NN}^{1a}$$

$$\lambda_s^{1,2,3}, \lambda_t, \rho_t$$

$$h^1_{\pi\!N\!N}$$

$$O(p^{-1})$$

A program of few-body measurements

Pionless theory

Ab initio few-body calcs

$$m_{N}\lambda_{pp} = -1.22A_{L}(pp) \qquad \text{Done} \qquad \text{LANSCE}$$

$$m_{N}\rho_{t} = -9.35A_{L}(np \rightarrow d\gamma) \qquad \text{HARD*}$$

$$m_{N}\lambda_{pn} = 1.6A_{L}(pp) - 3.7A_{L}(p\alpha) + 37A_{L}(np \rightarrow d\gamma) - 2P_{L}(np \rightarrow d\gamma)$$

$$m_{N}\lambda_{t} = 0.4A_{L}(pp) - 0.7A_{L}(p\alpha) + 7A_{L}(np \rightarrow d\gamma) + P_{L}(np \rightarrow d\gamma)$$

$$m_{N}\lambda_{nn} = 1.6A_{L}(pp) - 0.7A_{L}(p\alpha) + 333A_{L}(np \rightarrow d\gamma) - 1.08P_{L}(np \rightarrow d\gamma) + 0.83\frac{d\phi_{n\alpha}}{dz}$$

$$\lambda_{pp} = \lambda_{s}^{0} + \lambda_{s}^{1} + \lambda_{s}^{2} / \sqrt{6}$$

$$\lambda_{nn} = \lambda_{s}^{0} - \lambda_{s}^{1} + \lambda_{s}^{2} / \sqrt{6}$$

$$\lambda_{pn} = \lambda_{s}^{0} - 2 \lambda_{s}^{2} / \sqrt{6}$$

NIST

*HIGS
$$A_L(\gamma d \rightarrow np)$$

A program of few-body measurements

Complete determination of PV NN & γ NN interactions through $\mathcal{O}(p)$

Attempt to understand the λ^i , h_π etc. from QCD

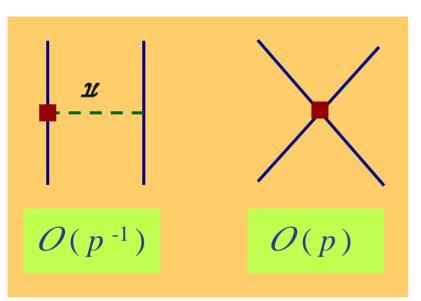
Are the PV LEC's "natural"?

Attempt to understand nuclear PV observables systematically

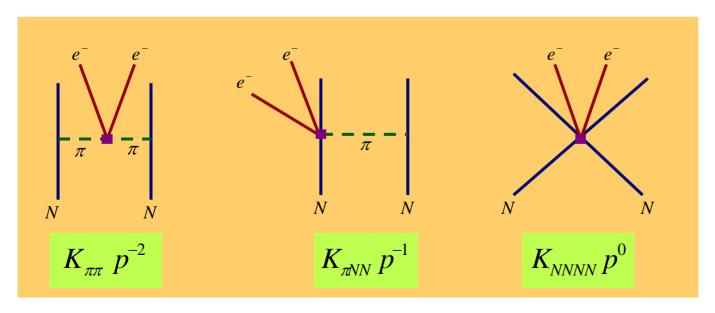
Does EFT power counting work in nuclei?

Hadronic PV in n-rich nuclei?

Hadronic PV as a probe



- Determine V_{PV} through $\mathcal{O}(p)$ from PV low-energy *few-body* studies where power counting works
- Re-analyze *nuclear* PV observables using this V_{PV}



Conclusions

- Nuclei provide unique and powerful laboratories in which to probe the fundamental symmetries of the early universe
- RIA will provide opportunities to carry out new and complementary experiments whose impact can live on well into the LHC era
- A number of theoretical challenges remain to be addressed at the level of field theory, QCD, and nuclear structure
- New experimental and theoretical efforts in nuclear structure physics are a key component of this quest