

Proton-neutron asymmetry in exotic nuclei

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Collective properties of exotic nuclei

Extensive new set of nuclei

- **Proton**-neutron imbalance
- Changes in shell structure?

Theoretical effort: Anticipate new collective phenomena

- Signatures by which phenomena can be recognized
- Estimates of where phenomena may occur

Most low-energy collective phenomena essentially *isoscalar*

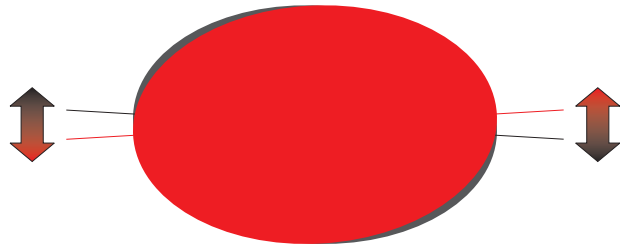
Similar **proton** and neutron distributions in the ground state

Deformation arises from strong **proton**-neutron quadrupole interaction, which couples **proton** and neutron deformations

Proton-neutron asymmetry in collective excitations

Scissors mode

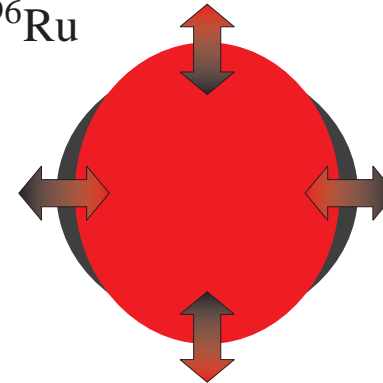
e.g., ^{156}Gd



N. Lo Iudice and F. Palumbo, *Phys. Rev. Lett.* **41**, 1532 (1978).
F. Iachello, *Nucl. Phys. A* **358**, 89c (1981).
D. Bohle *et al.*, *Phys. Lett. B* **137**, 27 (1984).

Mixed symmetry states

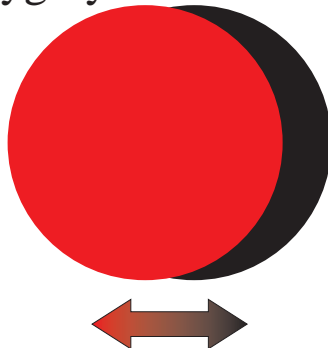
e.g., ^{94}Mo , ^{96}Ru



F. Iachello, *Phys. Rev. Lett.* **53**, 1427 (1984).
N. Pietralla *et al.*, *Phys. Rev. Lett.* **84**, 3775 (2000).

Dipole resonances

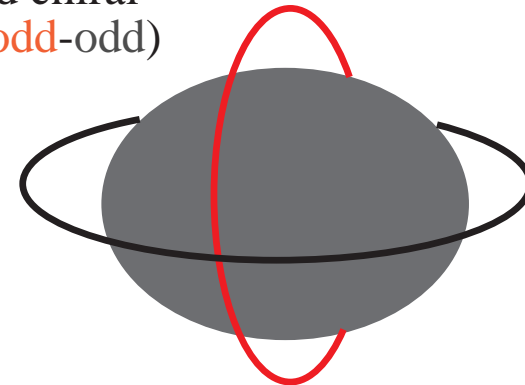
Giant and pygmy resonances



M. Goldhaber and E. Teller, *Phys. Rev.* **74**, 1046 (1948).
A. Zilges *et al.*, *Phys. Lett. B* **542**, 43 (2002).

Asymmetry in coupling

Shears and chiral rotation (odd-odd)



S. Frauendorf, *Rev. Mod. Phys.* **73**, 463 (2001).

Proton-neutron asymmetry in the ground state?

Very neutron-rich nuclei

Well-separated **proton** and neutron valence spaces

⇒ Reduced **proton**-neutron coupling strengths?

⇒ Larger role for **proton**-neutron asymmetry in ground state?

Nuclear structure

Ground state properties, excitation modes, transition radiations (*M1*)

Mechanisms for triaxiality

- Higher-order interactions in one-fluid Hamiltonian

$$(d^\dagger d^\dagger d^\dagger \tilde{d} \tilde{d} \tilde{d} \text{ or } \cos^2 3\gamma)$$

P. Van Isacker and J. Chen, Phys. Rev. C **24**, 684 (1981).

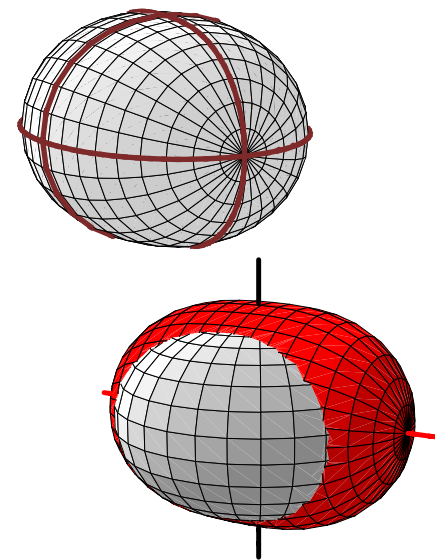
- Higher-multipolarity pairs (hexadecapole)

K. Heyde *et al.*, Nucl. Phys. A **398**, 235 (1983).

- Unaligned **proton** and neutron symmetry axes

A. E. L. Dieperink and R. Bijker, Phys. Lett. B **116**, 77 (1982).

J. N. Ginocchio and A. Leviatan, Ann. Phys. (N.Y.) **216**, 152 (1992).



The interacting boson model (IBM-1)

Truncation to s -wave ($J = 0$) and d -wave ($J = 2$) nucleon pairs

$$s_0 \quad d_{+2} \quad d_{+1} \quad d_0 \quad d_{-1} \quad d_{-2}$$

States: Linear combinations of $(s_0^\dagger)^n (d_{+2}^\dagger)^{n'} \dots |0\rangle$

Operators ($H, \hat{L}, \hat{T}, \dots$): Polynomials in $b^\dagger b$ e.g, $\hat{L} = \sqrt{10}[d^\dagger \times \tilde{d}]^{(1)}$

Algebraic model: Constructed from elements of Lie algebra

$$U(6) : s_0^\dagger s_0 \quad s_0^\dagger d_{+2} \quad \dots \quad d_{-2}^\dagger d_{-2}$$

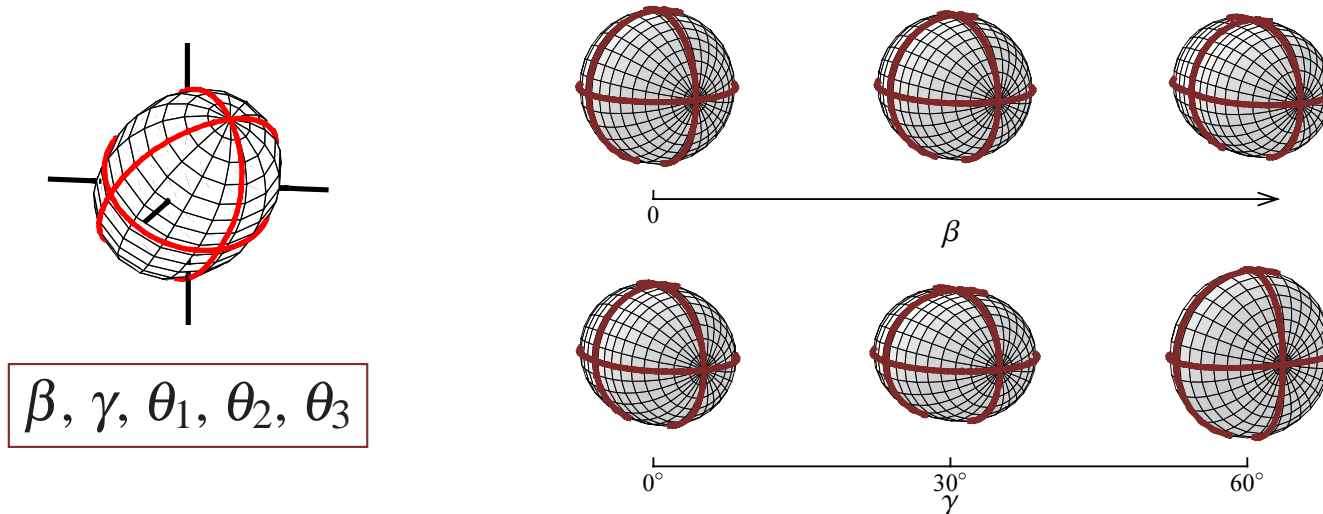
Dynamical symmetry

$$U(6) \supset \left(\begin{array}{c} U(5) \\ SO(6) \\ \underline{SU(3)} \\ \underline{SU(3)} \end{array} \right) \supset SO(5) \quad \supset \underbrace{SO(3) \supset SO(2)}_{\text{Angular momentum}}$$

- H constructed from Casimir (invariant) operators of subalgebra chain
- Eigenstates have good quantum numbers
- Problem exactly soluble (energies, eigenstates, transition MEs)
- Defines distinct form of ground state configuration (“phase”)

Classical limit of the IBM-1

Quadrupole-deformed liquid drop



$$\beta, \gamma, \theta_1, \theta_2, \theta_3$$

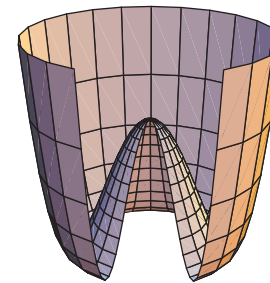
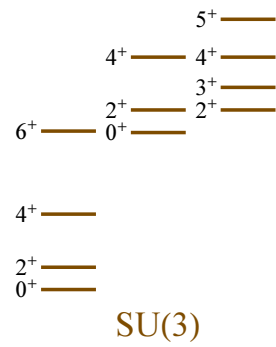
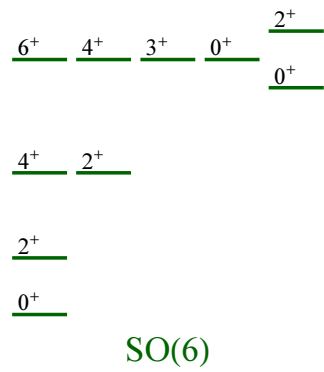
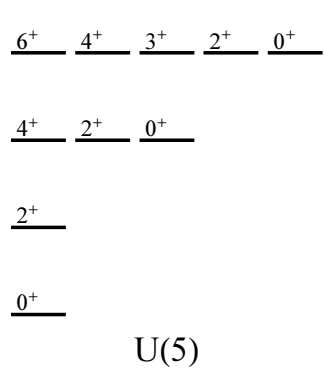
Coherent states $|\beta, \gamma\rangle$

$$|\beta, \gamma\rangle = \left[s_0^\dagger + \beta \cos \gamma d_0^\dagger + \frac{1}{\sqrt{2}} \beta \sin \gamma (d_{+2}^\dagger + d_{-2}^\dagger) \right]^N |0\rangle$$

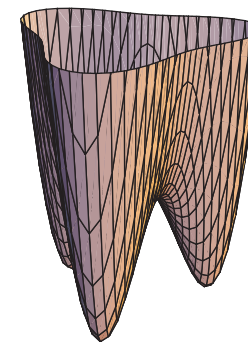
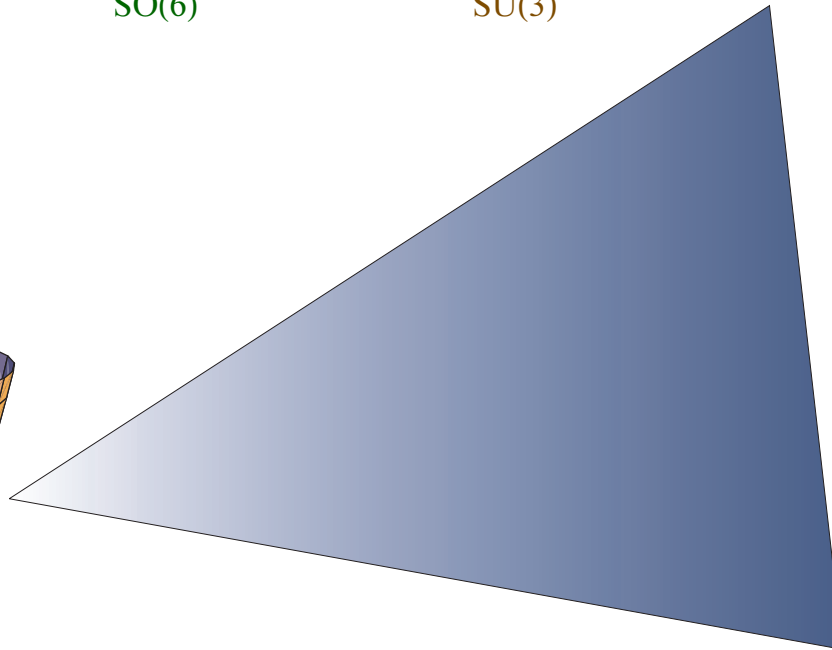
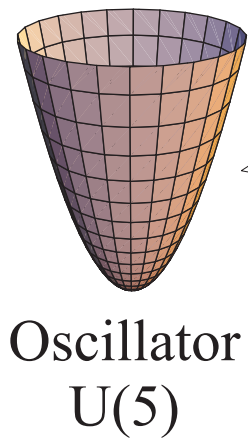
Classical energy surface

$$\mathcal{E}(\beta, \gamma) = \langle \beta, \gamma | H | \beta, \gamma \rangle$$

Minimization of $\mathcal{E} \Rightarrow$ ground state energy, equilibrium coordinate values



Deformed
 γ -soft
 SO(6)



Phase diagram of the IBM-1

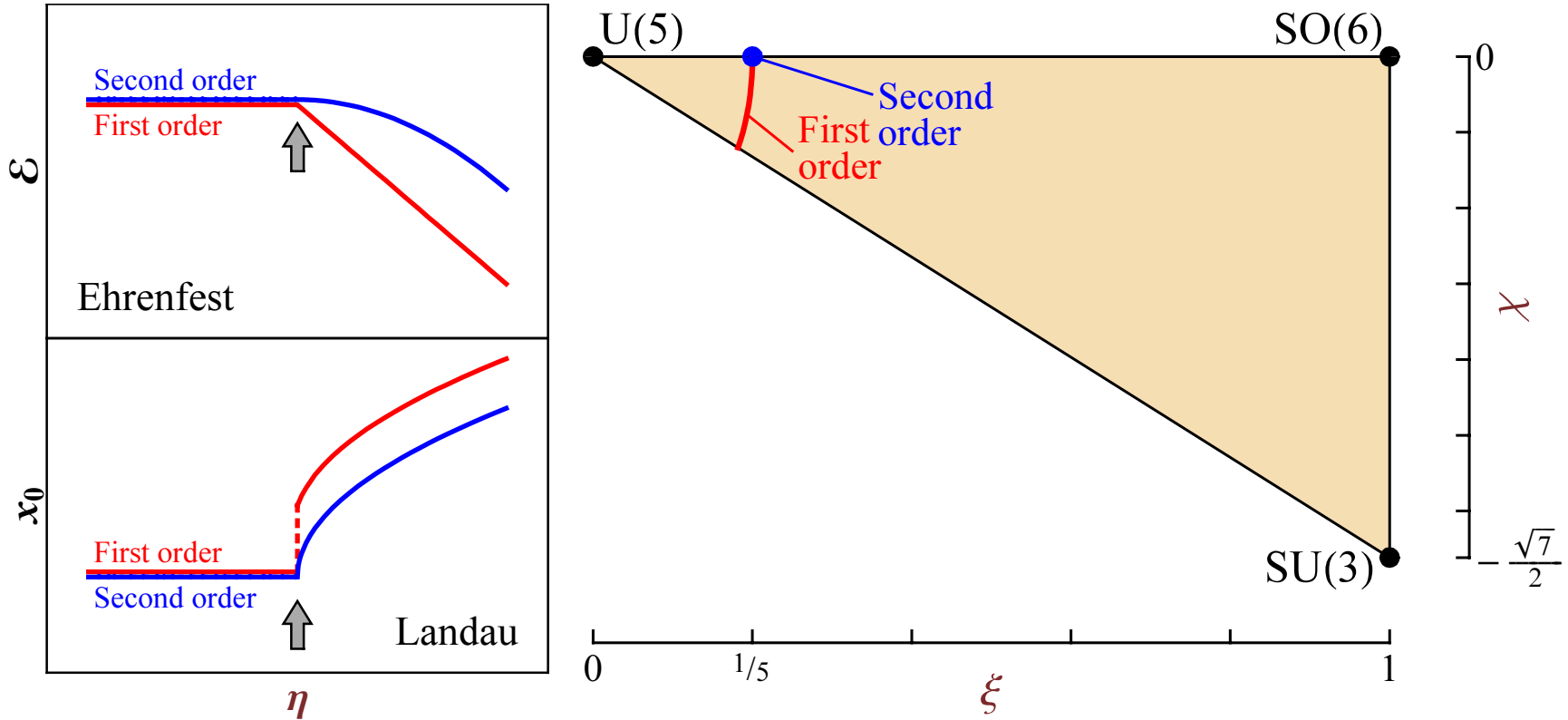
$$H = (1 - \xi) \frac{1}{N} \hat{n}_d - \xi \frac{1}{N^2} \hat{Q}^\chi \cdot \hat{Q}^\chi$$

ξ : (spherical) \leftrightarrow (deformed)

$$\hat{n}_d = d^\dagger \cdot \tilde{d} \quad \hat{Q}^\chi = (s^\dagger \times \tilde{d} + d^\dagger \times \tilde{s})^{(2)} + \chi (d^\dagger \times \tilde{d})^{(2)}$$

χ : (prolate) \leftrightarrow (γ -soft) \leftrightarrow (oblate)

A. E. L. Dieperink, O. Scholten, and F. Iachello, Phys. Rev. Lett. **44**, 1747 (1980).
 D. H. Feng, R. Gilmore, and S. R. Deans, Phys. Rev. C **23**, 1254 (1981).



The proton-neutron interacting boson model (IBM-2)

Proton and neutron pairs as separate boson species

$$\underbrace{s_{\pi,0}^\dagger \ d_{\pi,-2}^\dagger \ d_{\pi,-1}^\dagger \ d_{\pi,0}^\dagger \ d_{\pi,+1}^\dagger \ d_{\pi,+2}^\dagger}_{\text{Proton}} \quad \underbrace{s_{\nu,0}^\dagger \ d_{\nu,-2}^\dagger \ d_{\nu,-1}^\dagger \ d_{\nu,0}^\dagger \ d_{\nu,+1}^\dagger \ d_{\nu,+2}^\dagger}_{\text{Neutron}}$$

$$H = \varepsilon_\pi \hat{n}_{d\pi} + \varepsilon_\nu \hat{n}_{d\nu} + \kappa_{\pi\pi} Q_\pi \cdot Q_\pi + \kappa_{\pi\nu} Q_\pi \cdot Q_\nu + \kappa_{\nu\nu} Q_\nu \cdot Q_\nu + \dots$$

Symmetric dynamical symmetries (isoscalar)

$U_{\pi\nu}(5)$	$SO_{\pi\nu}(6)$	$SU_{\pi\nu}(3)$	$\overline{SU_{\pi\nu}(3)}$
Spherical	γ -soft	Prolate	Oblate

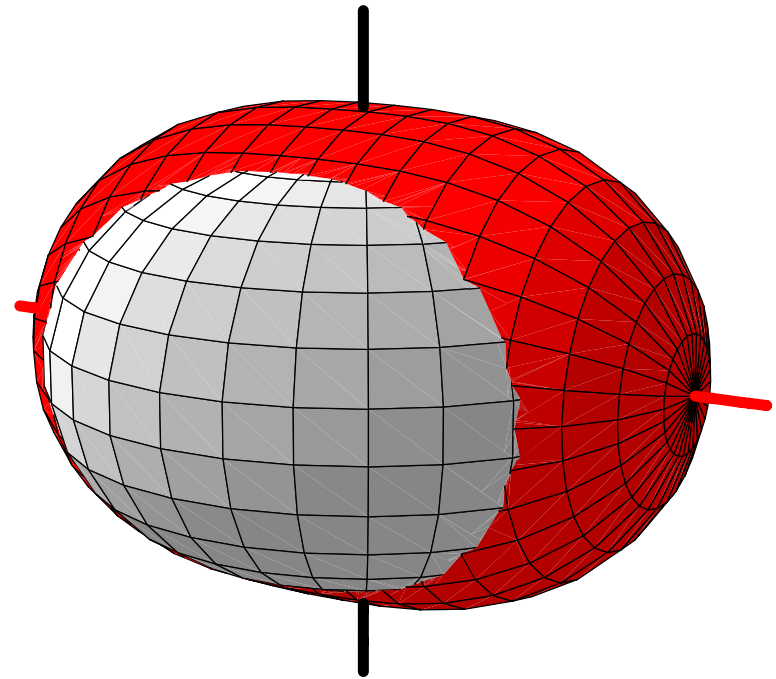
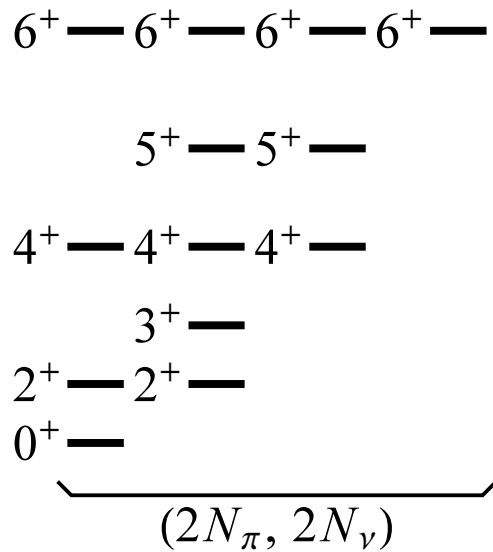
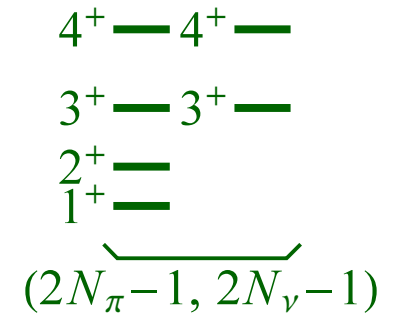
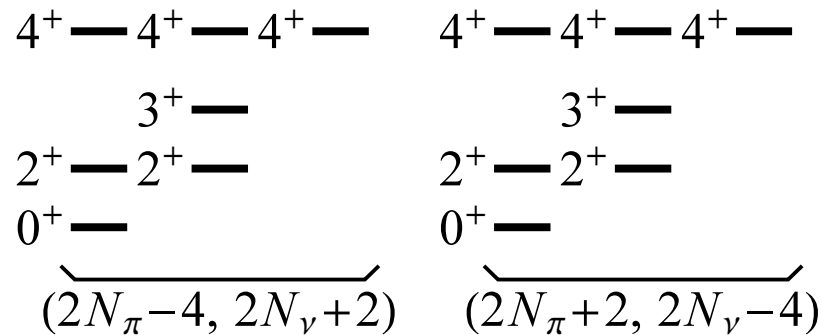
P. Van Isacker, K. Heyde, J. Jolie, and A. Sevrin, Ann. Phys. (NY) **171**, 253 (1986).

Asymmetric dynamical symmetries (isovector)

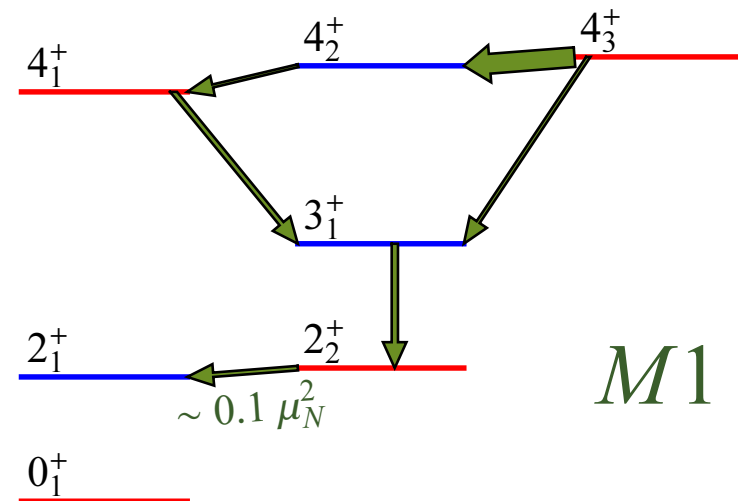
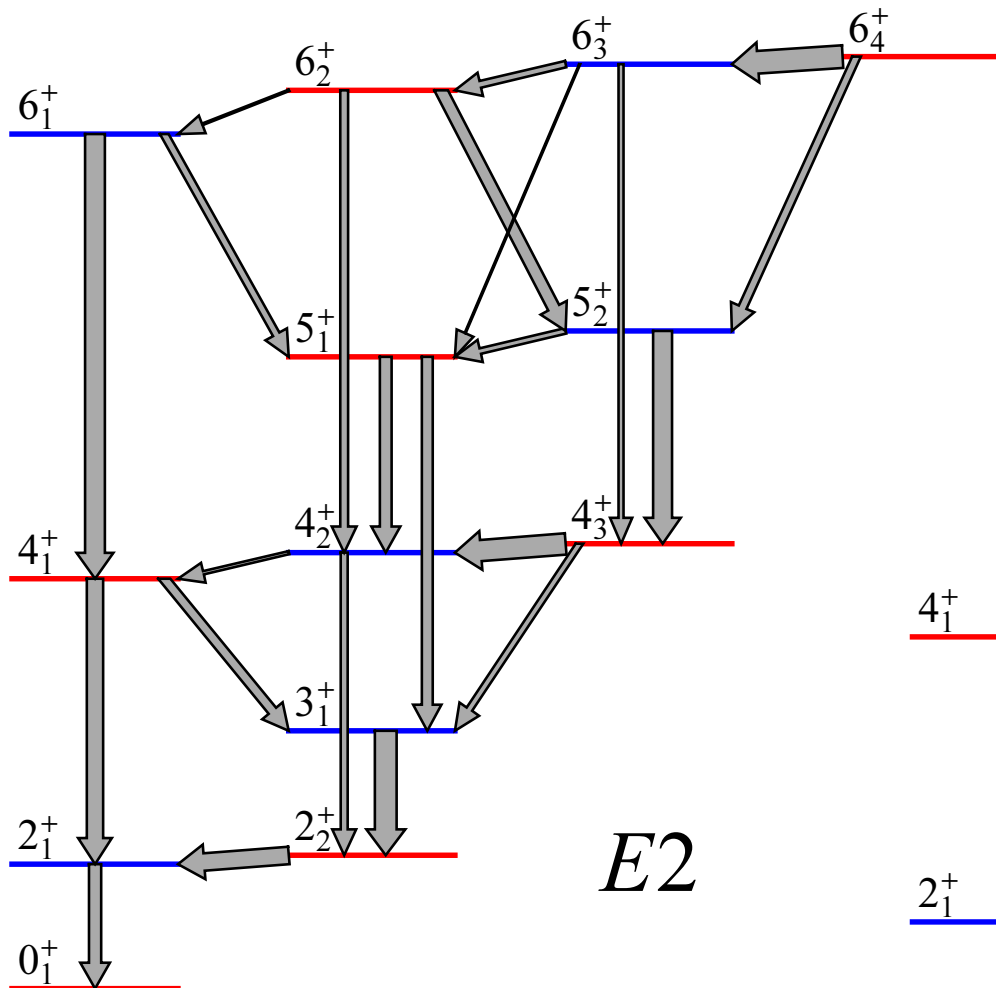
$$U_\pi(6) \otimes U_\nu(6) \supset \left\{ \begin{array}{l} SU_\pi(3) \otimes \overline{SU_\nu(3)} \supset \boxed{SU_{\pi\nu}^*(3)} \\ \overline{SU_\pi(3)} \otimes SU_\nu(3) \supset \boxed{SU_{\pi\nu}^*(3)} \end{array} \right\} \supset SO_{\pi\nu}(3) \supset SO_{\pi\nu}(2)$$

A. E. L. Dieperink and R. Bijker, Phys. Lett. B **116**, 77 (1982).
 A. Sevrin, K. Heyde, and J. Jolie, Phys. Rev. C **36**, 2621 (1987).
 N. R. Walet and P. J. Brussaard, Nucl. Phys. A **474**, 61 (1987).

The $SU_{\pi\nu}^*(3)$ dynamical symmetry: Proton-neutron triaxial structure



The $SU_{\pi\nu}^*(3)$ dynamical symmetry: Electromagnetic transitions



Essential parameters and coordinates for the IBM-2

Collective coordinates (“order parameters”)

$$\left\{ \begin{array}{c} \beta_\pi \ \gamma_\pi \ \theta_{1\pi} \ \theta_{2\pi} \ \theta_{3\pi} \\ \beta_\nu \ \gamma_\nu \ \theta_{1\nu} \ \theta_{2\nu} \ \theta_{3\nu} \end{array} \right\} \Rightarrow \left\{ \begin{array}{c} \beta_\pi \ \gamma_\pi \ \vartheta_1 \ \vartheta_2 \ \vartheta_3 \\ \beta_\nu \ \gamma_\nu \end{array} \right\} \Rightarrow \left\{ \begin{array}{c} \beta_\pi \ \gamma_\pi \\ \beta_\nu \ \gamma_\nu \end{array} \right\}$$

Coherent state energy surface $\mathcal{E}(\beta_\pi, \gamma_\pi, \beta_\nu, \gamma_\nu, \vartheta_1, \vartheta_2, \vartheta_3)$

Four order parameters: $\beta_\pi, \gamma_\pi, \beta_\nu,$ and γ_ν

Hamiltonian parameters (“control parameters”)

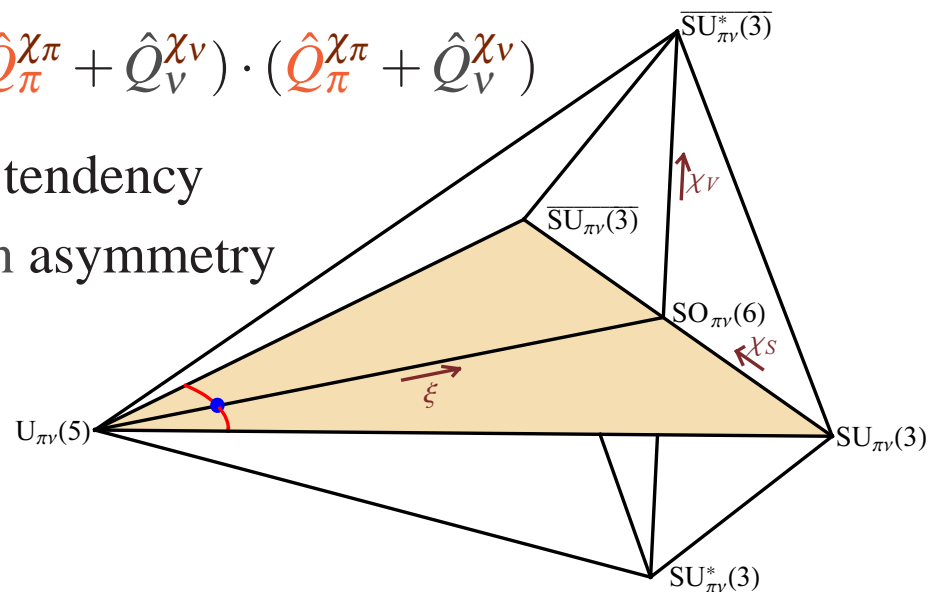
$$H = (1 - \xi) \frac{1}{N} (\hat{n}_{d\pi} + \hat{n}_{d\nu}) - \xi \frac{1}{N^2} (\hat{Q}_\pi^{\chi_\pi} + \hat{Q}_\nu^{\chi_\nu}) \cdot (\hat{Q}_\pi^{\chi_\pi} + \hat{Q}_\nu^{\chi_\nu})$$

$$\chi_S = \frac{1}{2} (\chi_\pi + \chi_\nu) \text{ prolate-oblate tendency}$$

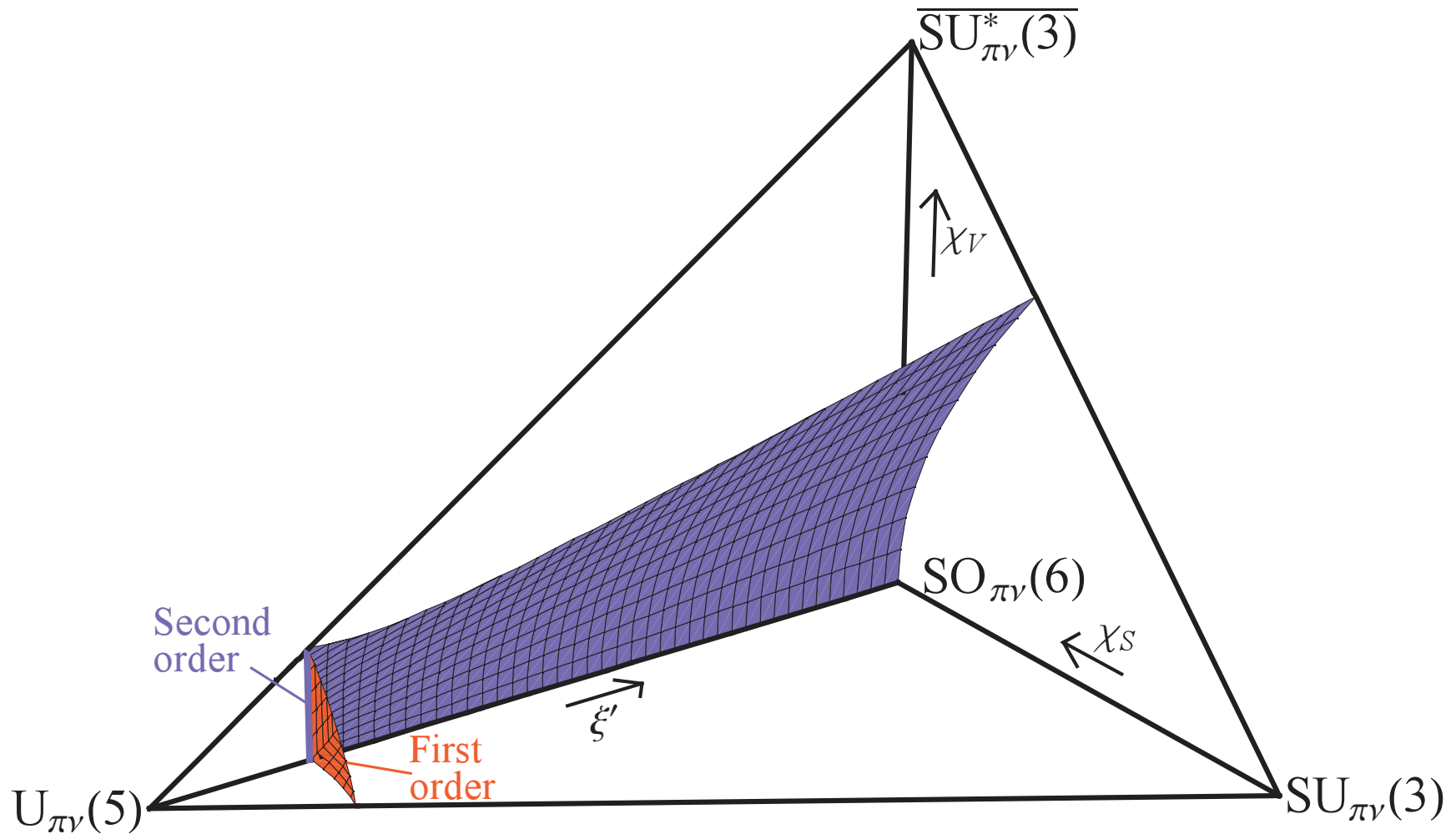
$$\chi_V = \frac{1}{2} (\chi_\pi - \chi_\nu) \text{ proton-neutron asymmetry}$$

Three control parameters:

$$\xi, \chi_S, \text{ and } \chi_V$$



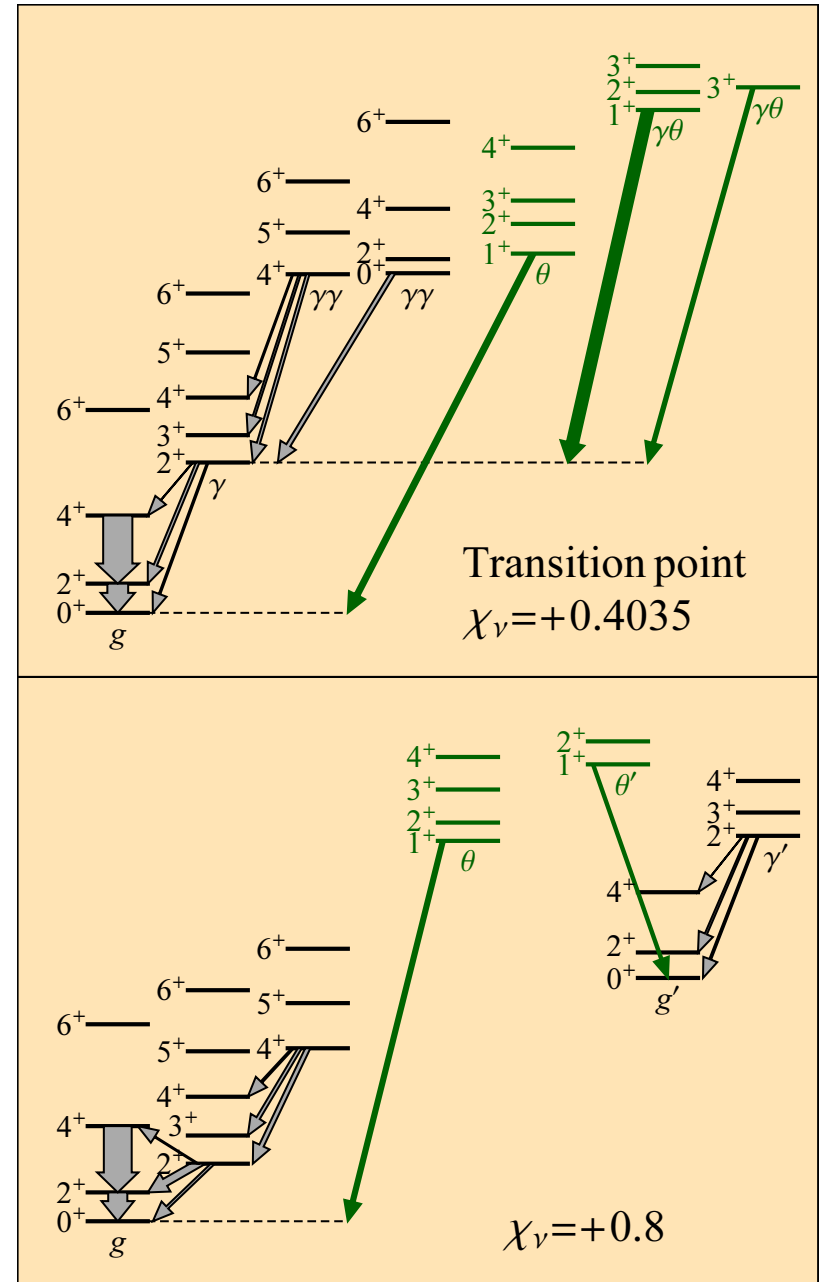
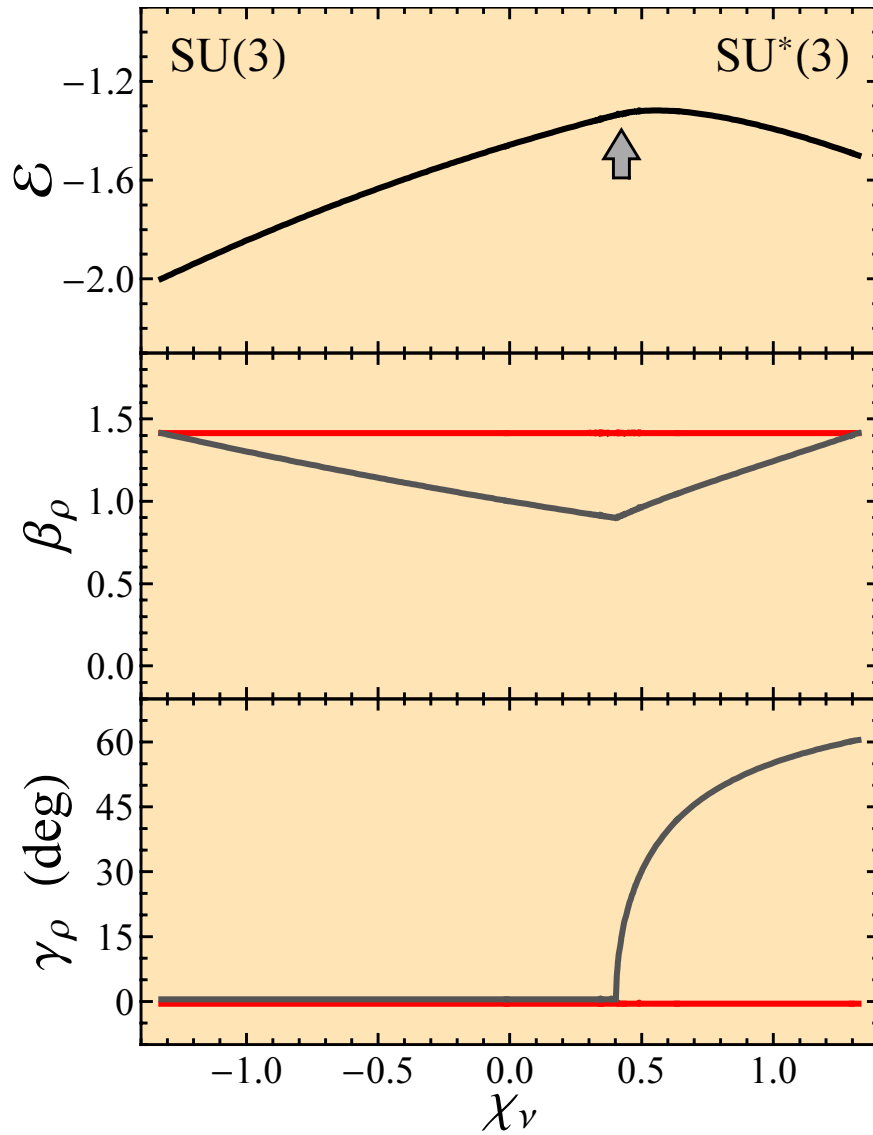
Phase diagram of the IBM-2



M. A. Caprio and F. Iachello, Phys. Rev. Lett. **93**, 242502 (2004).
M. A. Caprio and F. Iachello, Ann. Phys. (N.Y.) **318**, 454 (2005).

$$N_{\pi}/N_{\nu} = 1$$

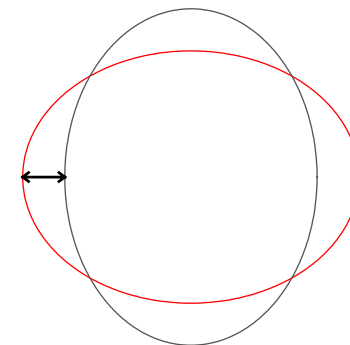
SU_{πν}(3)-SU*_{πν}(3) transition



Proton-neutron symmetry energy (Majorana operator)

Difference between **proton** and neutron deformation tensors

$$\begin{aligned}\hat{M} &\equiv -2 \sum_{k=1,3} (d_{\pi}^{\dagger} \times d_{\nu}^{\dagger})^{(k)} \cdot (\tilde{d}_{\pi} \times \tilde{d}_{\nu})^{(k)} \\ &\quad + (s_{\pi}^{\dagger} \times d_{\nu}^{\dagger} - s_{\nu}^{\dagger} \times d_{\pi}^{\dagger})^{(2)} \cdot (\tilde{s}_{\pi} \times \tilde{d}_{\nu} - \tilde{s}_{\nu} \times \tilde{d}_{\pi})^{(2)} \\ &\approx |\alpha_{\pi} - \alpha_{\nu}|^2\end{aligned}$$



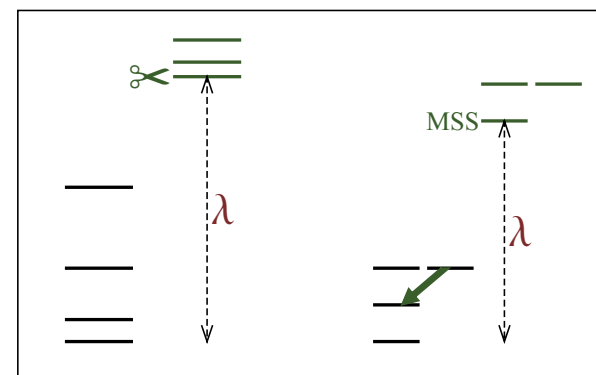
Major ingredient in realistic Hamiltonian

$$H = \underbrace{\varepsilon_{\pi} \hat{n}_{d\pi} + \varepsilon_{\nu} \hat{n}_{d\nu}}_{\text{Pair energy}} + \underbrace{\kappa_{\pi\pi} Q_{\pi} \cdot Q_{\pi} + \kappa_{\pi\nu} Q_{\pi} \cdot Q_{\nu} + \kappa_{\nu\nu} Q_{\nu} \cdot Q_{\nu}}_{\text{Quadrupole}} + \underbrace{\lambda \hat{M}}_{\text{Symmetry}}$$

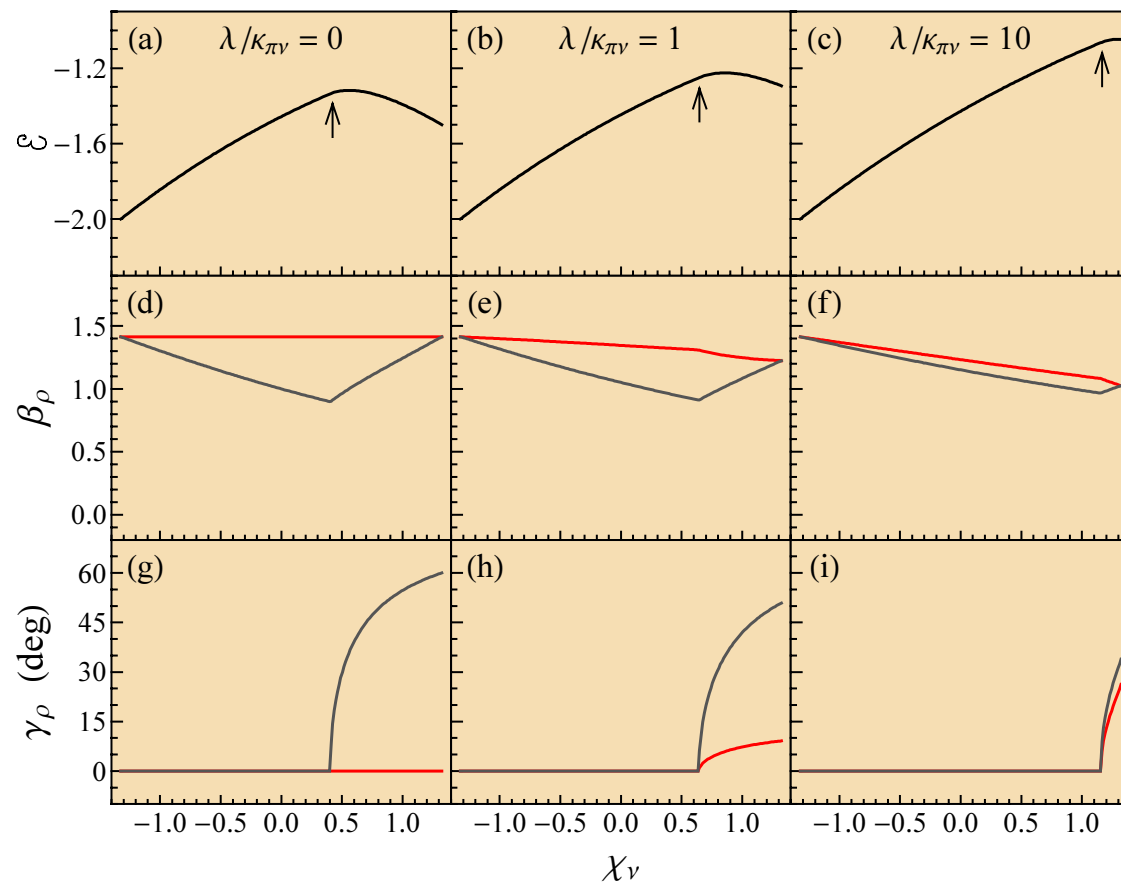
Strength λ approximately known

- From scissors and mixed-symmetry energies
- From $M1$ mixing ratios

$$\frac{\lambda}{\kappa_{\pi\nu}} \approx 5$$



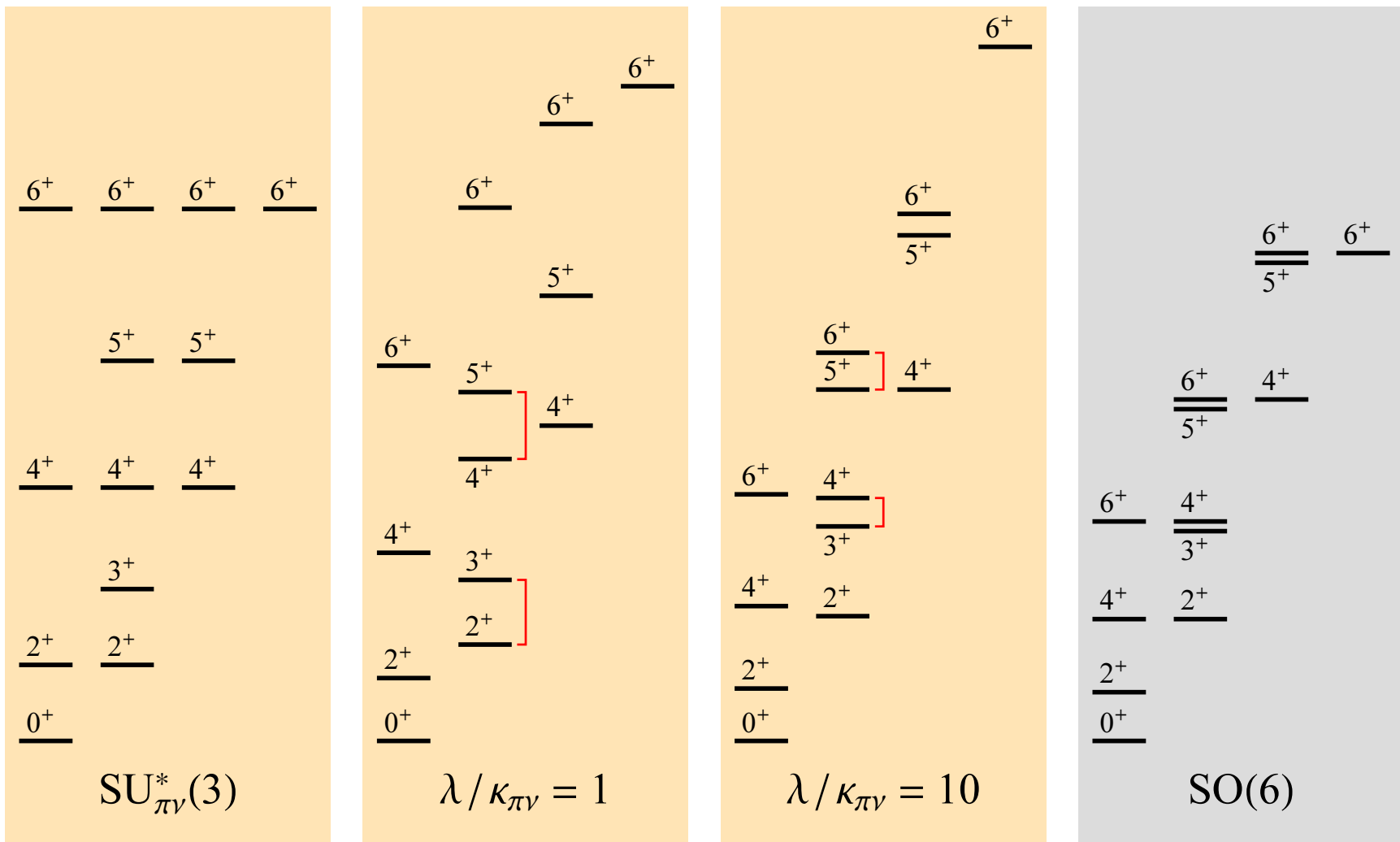
Effect of Majorana operator on phase transition



- Phase transition to triaxiality delayed
- **Proton** and neutron equilibrium coordinates values brought together
- *But also* energy minimum at triaxial deformation shallower

$SU_{\pi\nu}^*(3)$ triaxial \Rightarrow one-fluid triaxial \Rightarrow one-fluid γ -soft



Effect of Majorana operator on $SU_{\pi\nu}^*(3)$ structure



Majorana strength \rightarrow

Proton-neutron triaxiality

Main signatures

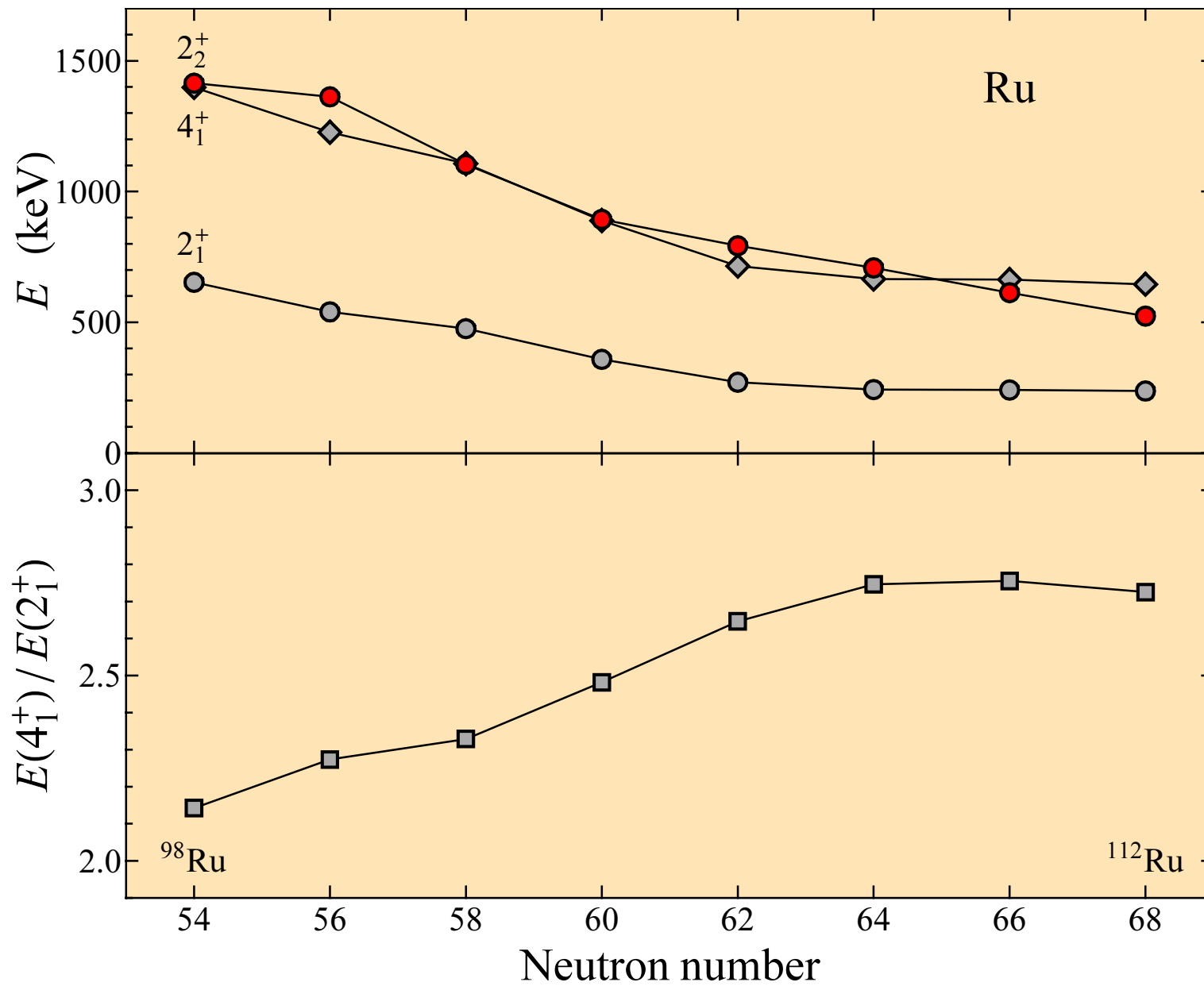
- Low-lying $K = 2$ band
but rotational $L(L + 1)$ energy sequence
- Unusual $B(E2)$ strength pattern
similar to classic rigid triaxial rotor (Davydov)
- Anharmonically low $K = 4$ band
- Strong $M1$ admixtures 
- Orthogonal scissors mode 

But attenuated by Majorana operator

$SU_{\pi\nu}^*(3)$ triaxial

\Rightarrow one-fluid triaxial

\Rightarrow one-fluid γ -soft



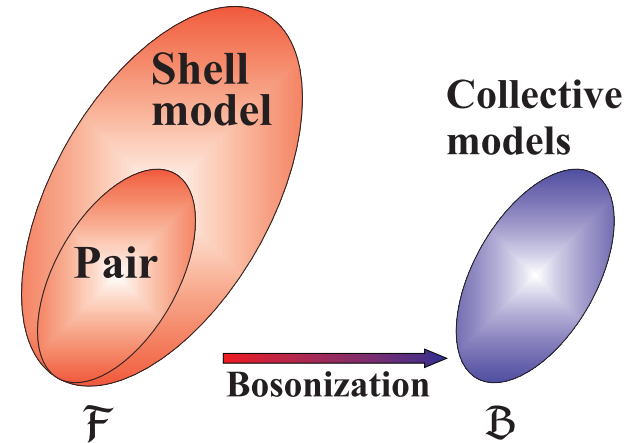
Where might asymmetric structure be expected?

Collective structure depends upon underlying single-particle structure

- Energy spacing (subshell gaps?)
- Ordering of orbitals (low j ? high j ?)
- Radial wave functions (compact? diffuse?)

Manifested in effective interactions

- Pairing interaction (s -wave, d -wave, ...)
- Multipole interaction (quadrupole, ...)
- Symmetry energy (Majorana)



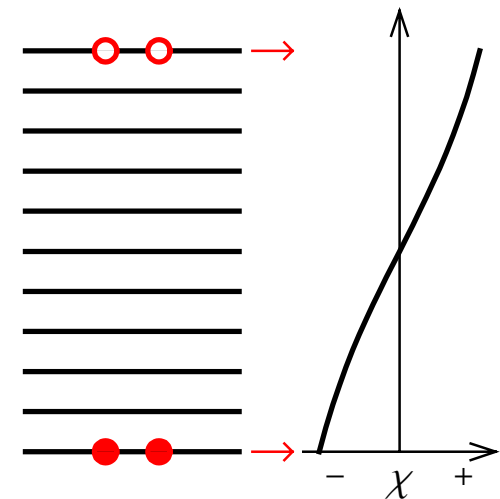
Qualitative estimate

Particle-like bosons \Rightarrow Prolate tendency

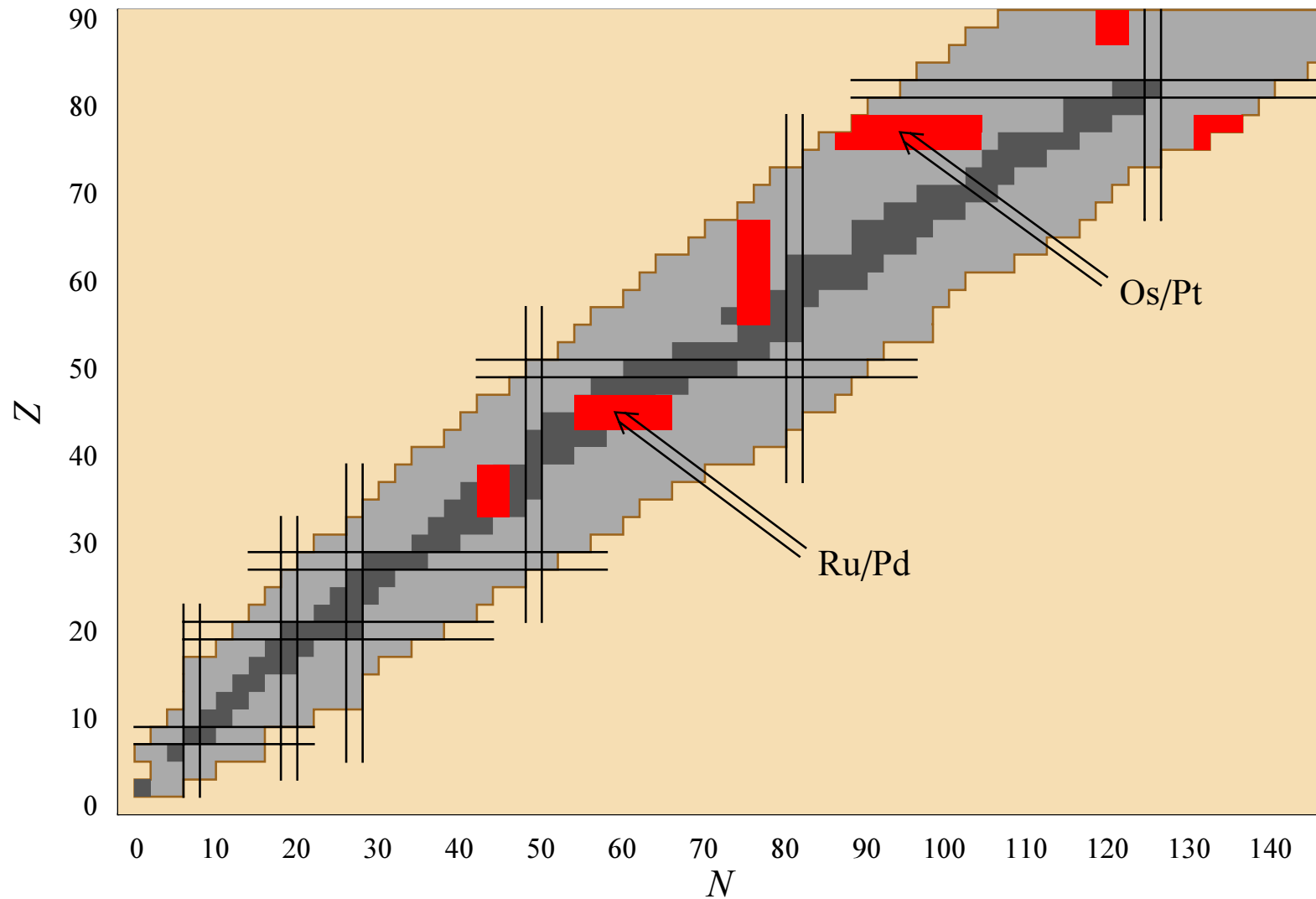
Hole-like bosons \Rightarrow Oblate tendency

But very sensitive to underlying shell structure

A. van Egmond and K. Allaart, Nucl. Phys. A **425**, 275 (1984).
T. Otsuka, Nucl. Phys. A **557**, 531c (1993).



Prospective regions for $SU_{\pi V}^*(3)$ triaxial structure



Conclusions

In preparation for exotic beam facility...

Have investigated **proton**-neutron asymmetric collective structure, within framework of IBM-2

Proton-neutron asymmetry

- Suppressed by Majorana interaction
- But could play role for nuclei far from stability

$SU_{\pi\nu}^*(3)$ *dynamical symmetry*

- Ideal limit, not likely to be reached
- Illustrates basic characteristics of **proton**-neutron triaxiality

Full collective analysis of two-fluid system

- **Phase diagram**
- **Nature of phase transitions**
- **Signatures of asymmetric structure**

Bose-Fermi system

- Odd mass or odd-odd nuclei
 - bosonic core + unpaired nucleons
- Odd nuclei will play major role in shell structure studies
- Coupling to unpaired nucleon significantly influences collective structure of even-even core (core polarization)
- Interacting boson fermion model (IBFM)