

New Developments in Nuclear Supersymmetry



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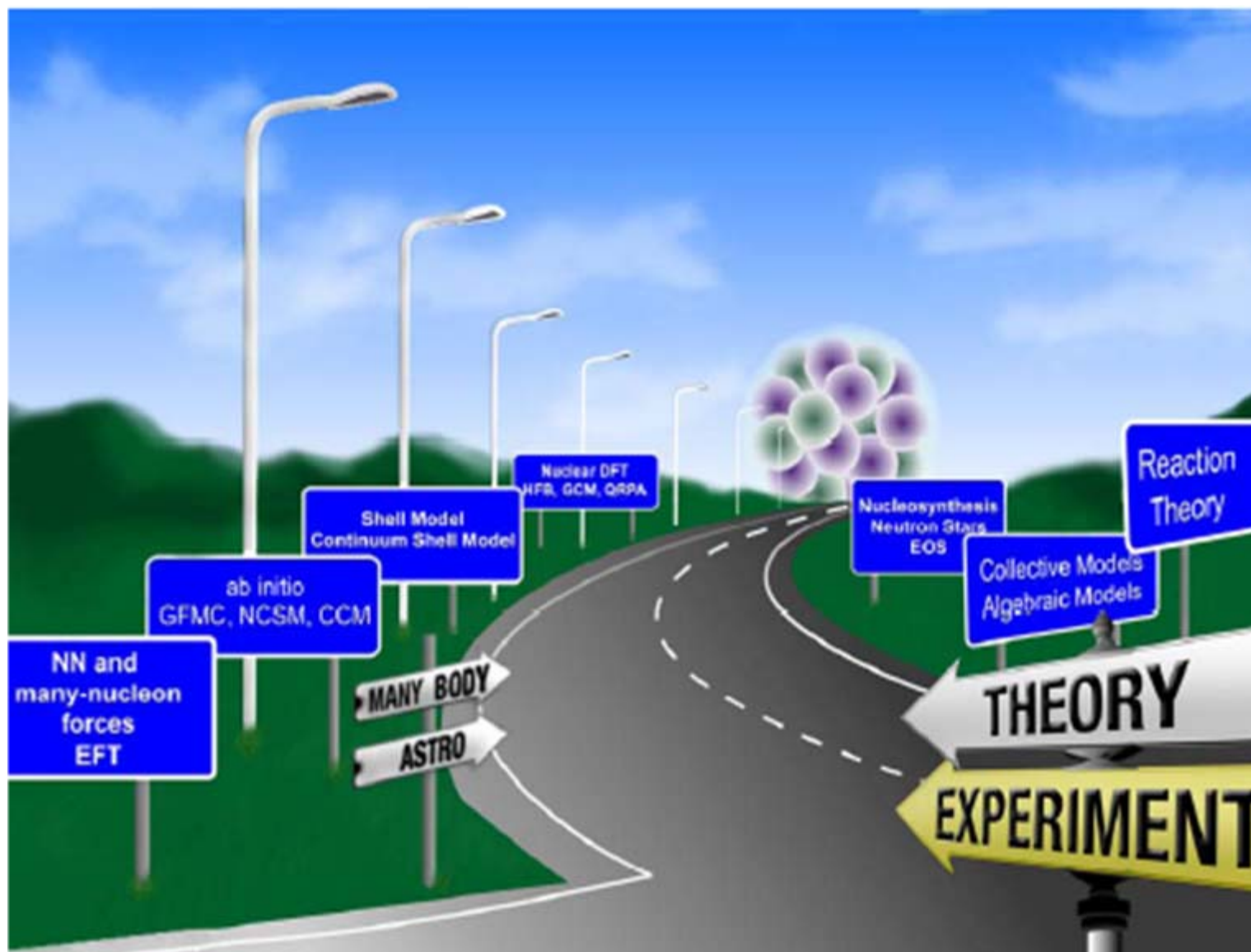
Alejandro Frank

Jan Jolie (Köln)

Gerhard Graw (München)

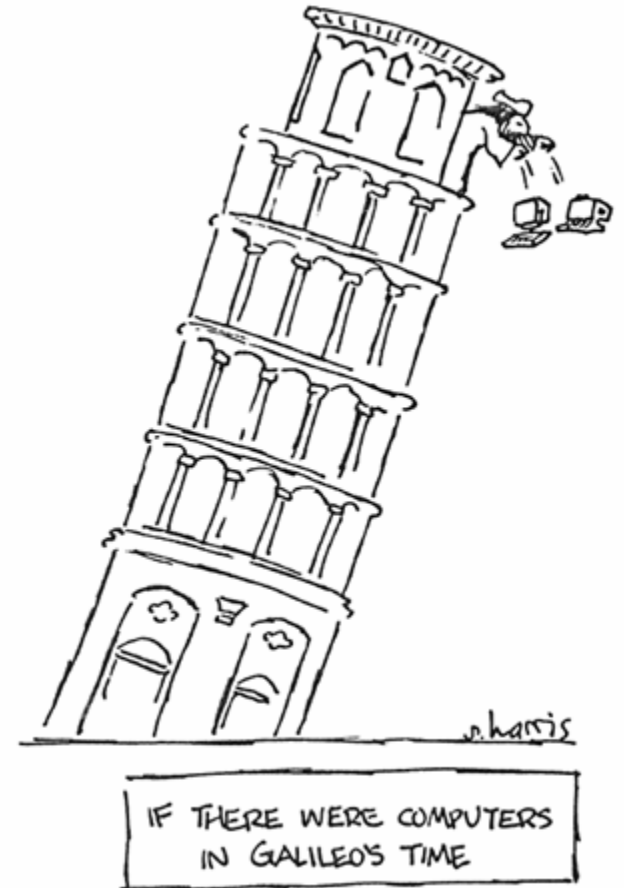
The Nuclear Many-Body Problem

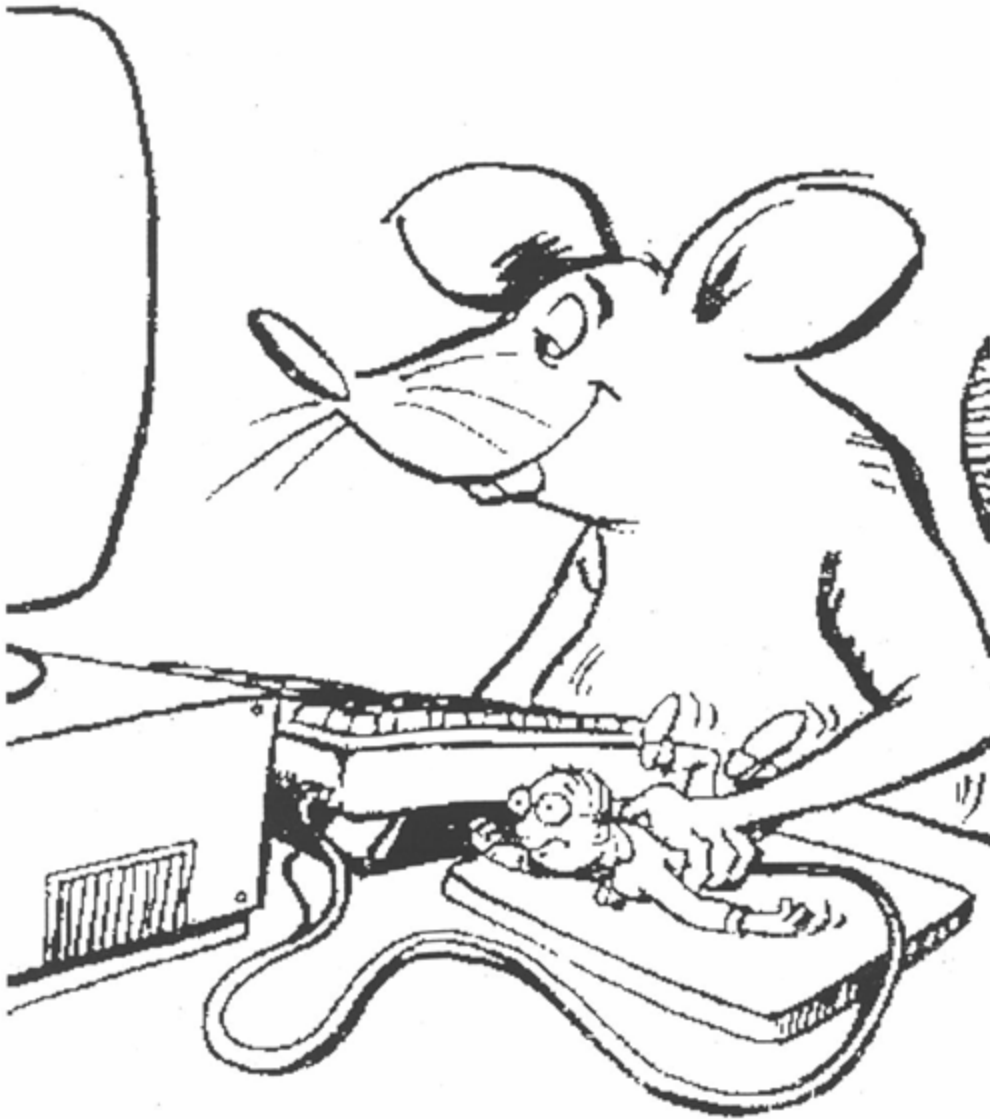
- Ab initio methods: GFMC, NCSM, CCM, ...
- Effective field theory
- Shell model: Monte Carlo, continuum SM, ...
- Mean-field methods: DFT, QRPA, HFB, GCM, ...
- Phenomenological models of collective motion: IBM and its extensions, ...
- Dynamical (super)symmetries
- ...



Motivation

- Large scale calculations
 - Ab initio
 - Shell model
 - Mean field
- Symmetry methods
 - IBM and IBFM with isospin
 - Nuclear supersymmetry





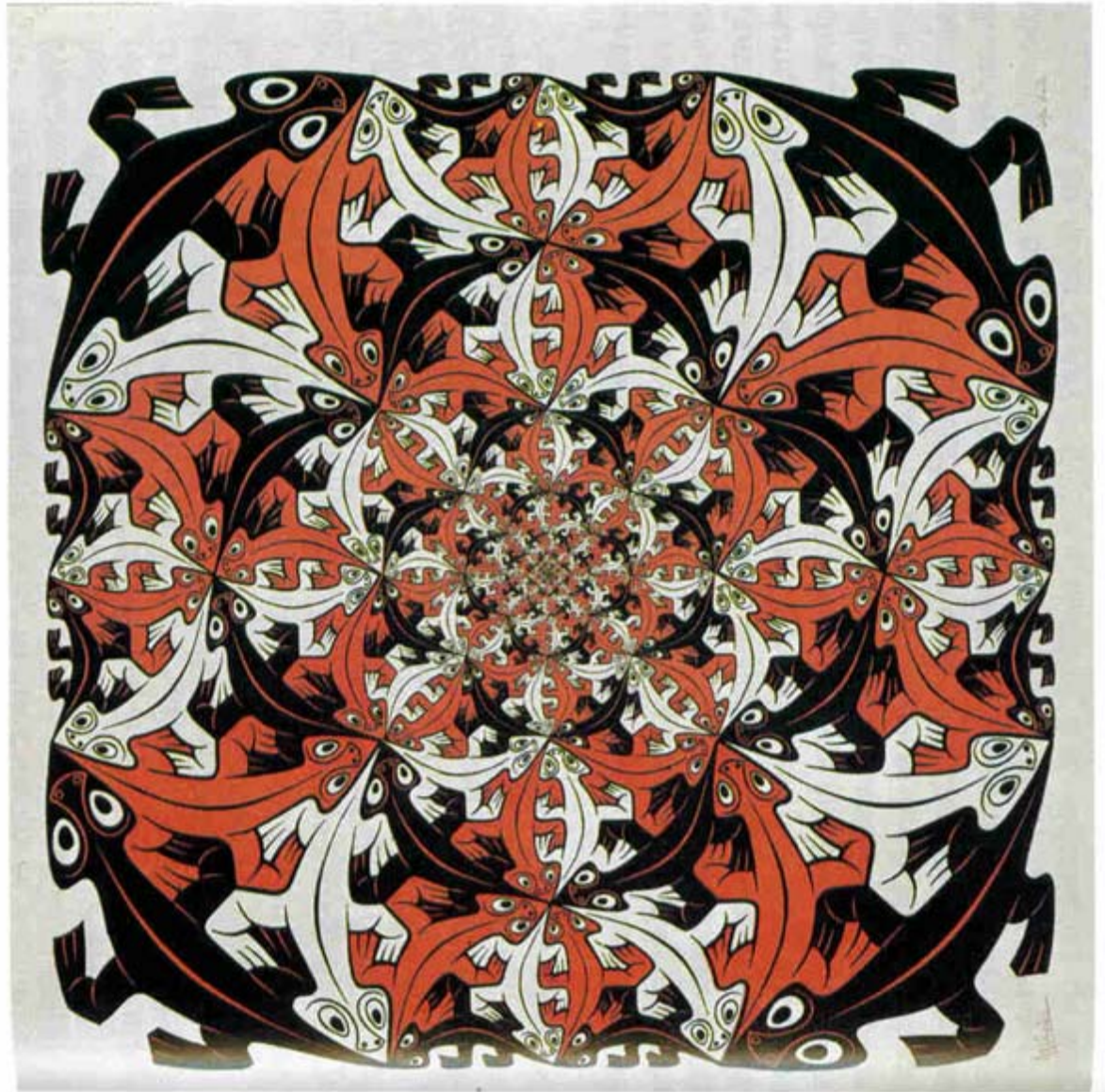
*I am very happy to learn
that the computer
understands the problem,
but I would like to
understand it too*

Eugene Wigner

Motivation

- What are the “effective” degrees of freedom?
- Are there “effective” symmetries?
- Symmetries provide benchmarks
- Examples:
 - special solutions to the Bohr Hamiltonian,
 - dynamical symmetries of the IBM,
 - pseudo-spin symmetries

Smaller and smaller
M.C. Escher



Outline

- Introduction
- Interacting boson models
- Dynamical supersymmetries
- Heavy nuclei: the $A \sim 190$ mass region
- Light nuclei: sd- and pf-shell
- Summary and conclusions

Symmetries

- *Geometric symmetries*

Buckyball with icosahedral symmetry

- *Permutation symmetries*

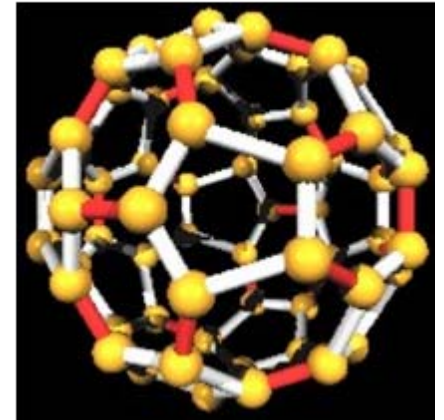
Fermi-Dirac and Bose-Einstein statistics

- *Space-time symmetries*

Rotational invariance in nonrelativistic QM,
Lorentz invariance in relativistic QM

- *Gauge symmetries*

Dirac equation with external electromagnetic field



Dynamical Symmetries

- Hydrogen atom (Pauli, 1926)
- Isospin symmetry (Heisenberg, 1932)
- Spin-isospin symmetry (Wigner, 1937)
- Pairing, seniority (Racah, 1943)
- Elliott model (Elliott, 1958)
- Flavor symmetry (Gell-Mann, Ne'eman, 1962)
- Interacting boson model (Arima, Iachello, 1974)
- Nuclear supersymmetry (Iachello, 1980)

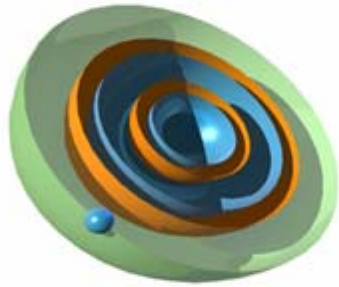
Interacting Boson Model

- The IBM describes collective excitations in even-even nuclei in terms of a system of correlated pairs of nucleons with angular momentum $L=0$ and $L=2$ which are treated as bosons (s and d bosons) (Arima and Iachello, 1974)
- The number of bosons N is half the number of valence nucleons
- Introduce boson creation and annihilation operators

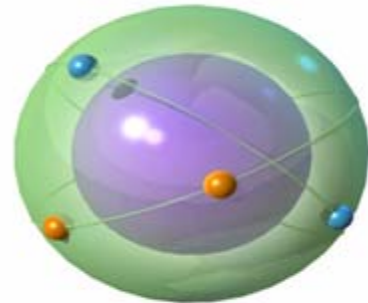
$$b_i^\dagger, b_i, \quad i = l, m \quad (l = 0, 2 \quad -l \leq m \leq l)$$

which satisfy the commutation relations

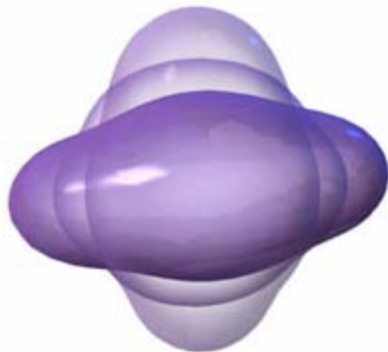
$$[b_i, b_j^\dagger] = \delta_{ij}, \quad [b_i^\dagger, b_j^\dagger] = [b_i, b_j] = 0$$



Shell structure:
valence nucleons



Cooper pairing:
s, d boson system



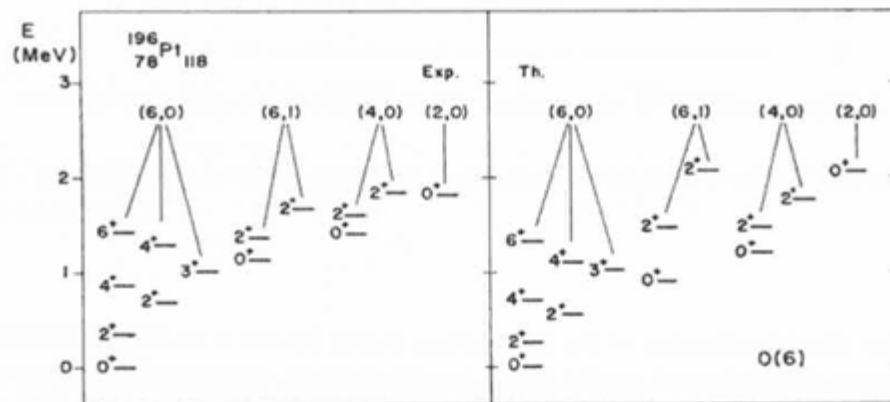
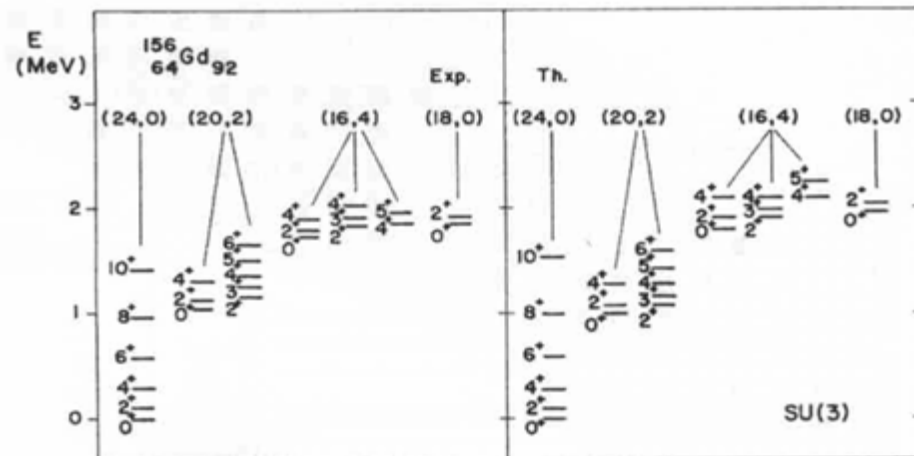
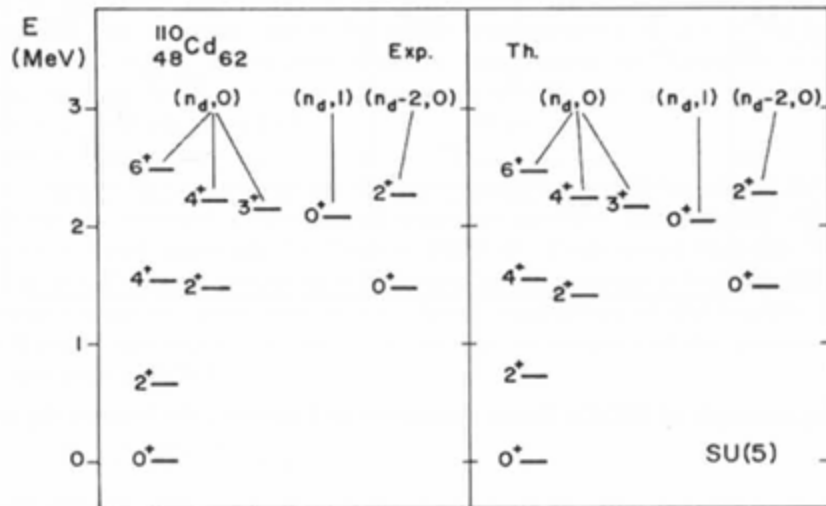
Collective motion:
nuclear shapes

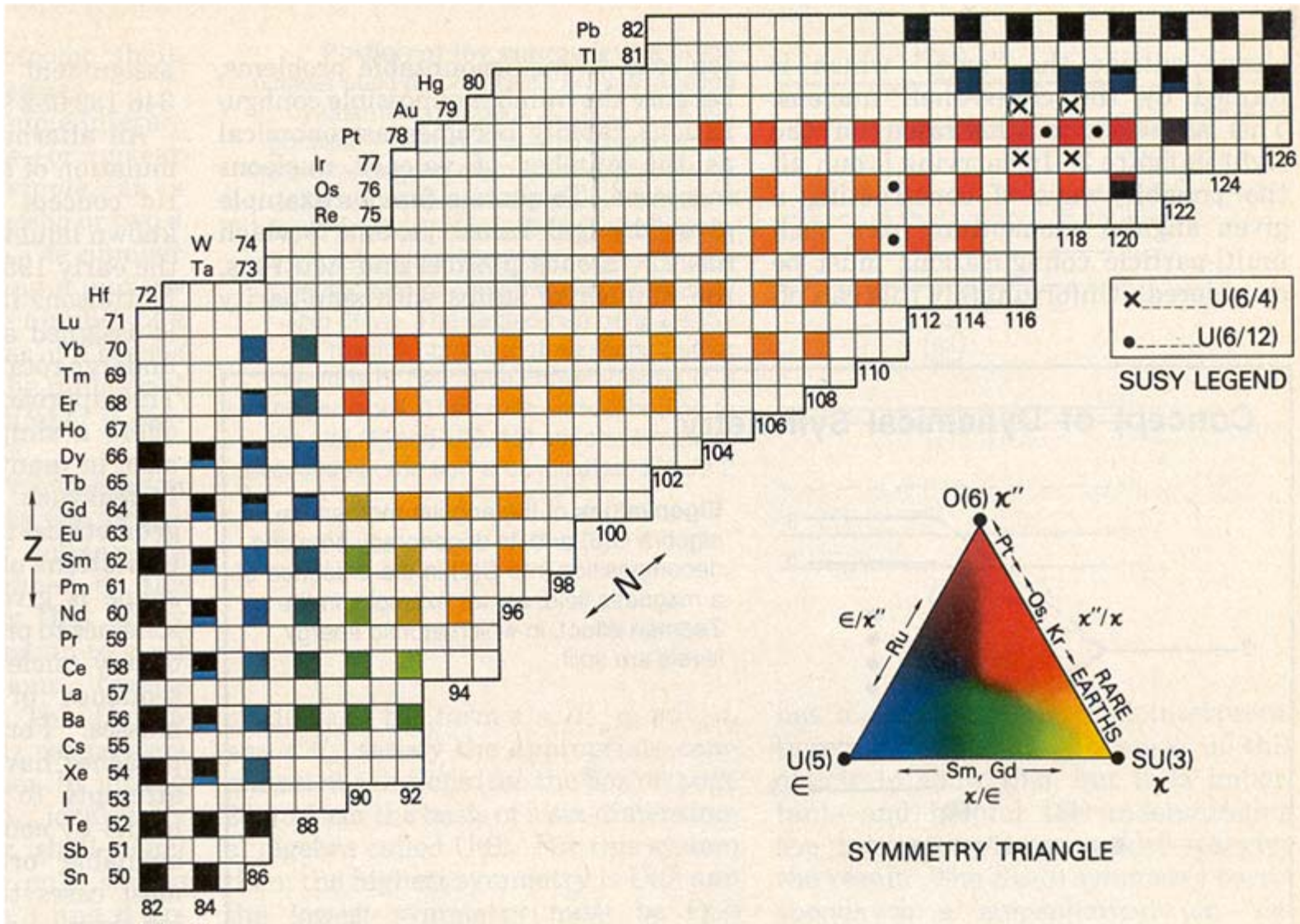
Dynamical Symmetries

$$U(6) \supset \left\{ \begin{array}{lll} U(5) \supset SO(5) \supset SO(3) & \text{vibrational} \\ n_d & \tau & L \\ \\ SU(3) \supset SO(3) & \text{rotational} \\ (\lambda, \mu) & L & \\ \\ SO(6) \supset SO(5) \supset SO(3) & \gamma - \text{unstable} \\ \sigma & \tau & L \end{array} \right.$$

Schematic
Hamiltonian :

$$\begin{aligned}
 H &= \epsilon \hat{n}_d - \kappa \hat{Q}(\chi) \cdot \hat{Q}(\chi) \\
 \hat{n}_d &= \sum_m d_m^\dagger d_m \\
 \hat{Q}(\chi) &= (s^\dagger \tilde{d} + d^\dagger \tilde{s} + \chi d^\dagger \tilde{d})^{(2)}
 \end{aligned}$$





Nuclear Supersymmetry

- Consider an extension of the IBM which includes, in addition to the collective degrees of freedom (bosons), single-particle degrees of freedom of an extra unpaired proton or neutron (fermion with angular momentum $j=j_1, j_2, \dots$)
- For the extra nucleon, introduce fermion creation and annihilation operators satisfy anticommutation relations

$$\{a_\mu, a_\nu^\dagger\} = \delta_{\mu\nu}, \quad \{a_\mu^\dagger, a_\nu^\dagger\} = \{a_\mu, a_\nu\} = 0$$

Building Blocks

$$\begin{array}{ll}
 \text{bosons} & l = 0, 2 \quad \sum_l (2l + 1) = 6 \\
 \text{fermions} & j = j_1, j_2, \dots \quad \sum_j (2j + 1) = \Omega
 \end{array}$$

Model	Generators	Invariant	Algebra
IBM	$b_i^\dagger b_j$	N	$U(6)$
IBFM	$b_i^\dagger b_j, a_\mu^\dagger a_\nu$	N, M	$U(6) \otimes U(\Omega)$
SUSY	$b_i^\dagger b_j, a_\mu^\dagger a_\nu, b_i^\dagger a_\mu, a_\mu^\dagger b_i$	\mathcal{N}	$U(6/\Omega)$

$$\begin{array}{ll}
 N = \sum_i b_i^\dagger b_i & \text{total number of bosons} \\
 M = \sum_\mu a_\mu^\dagger a_\mu & \text{total number of fermions} \\
 \mathcal{N} = N + M & \text{total number of bosons plus fermions}
 \end{array}$$

Algebraic Structure

B_{ij}	$=$	$b_i^\dagger b_j$	boson	\rightarrow	boson
$A_{\mu\nu}$	$=$	$a_\mu^\dagger a_\nu$	fermion	\rightarrow	fermion
$F_{i\mu}$	$=$	$b_i^\dagger a_\mu$	fermion	\rightarrow	boson
$G_{\mu i}$	$=$	$a_\mu^\dagger b_i$	boson	\rightarrow	fermion

Supersymmetry: the total number
of bosons AND fermions is conserved

$$\mathcal{N} = N + M$$

Hamiltonian

$$H = H_B + H_F + V_{BF}$$

$$H_B = \sum_{ij} \epsilon_{ij} B_{ij} + \sum_{ijkl} u_{ijkl} B_{ij} B_{kl}$$

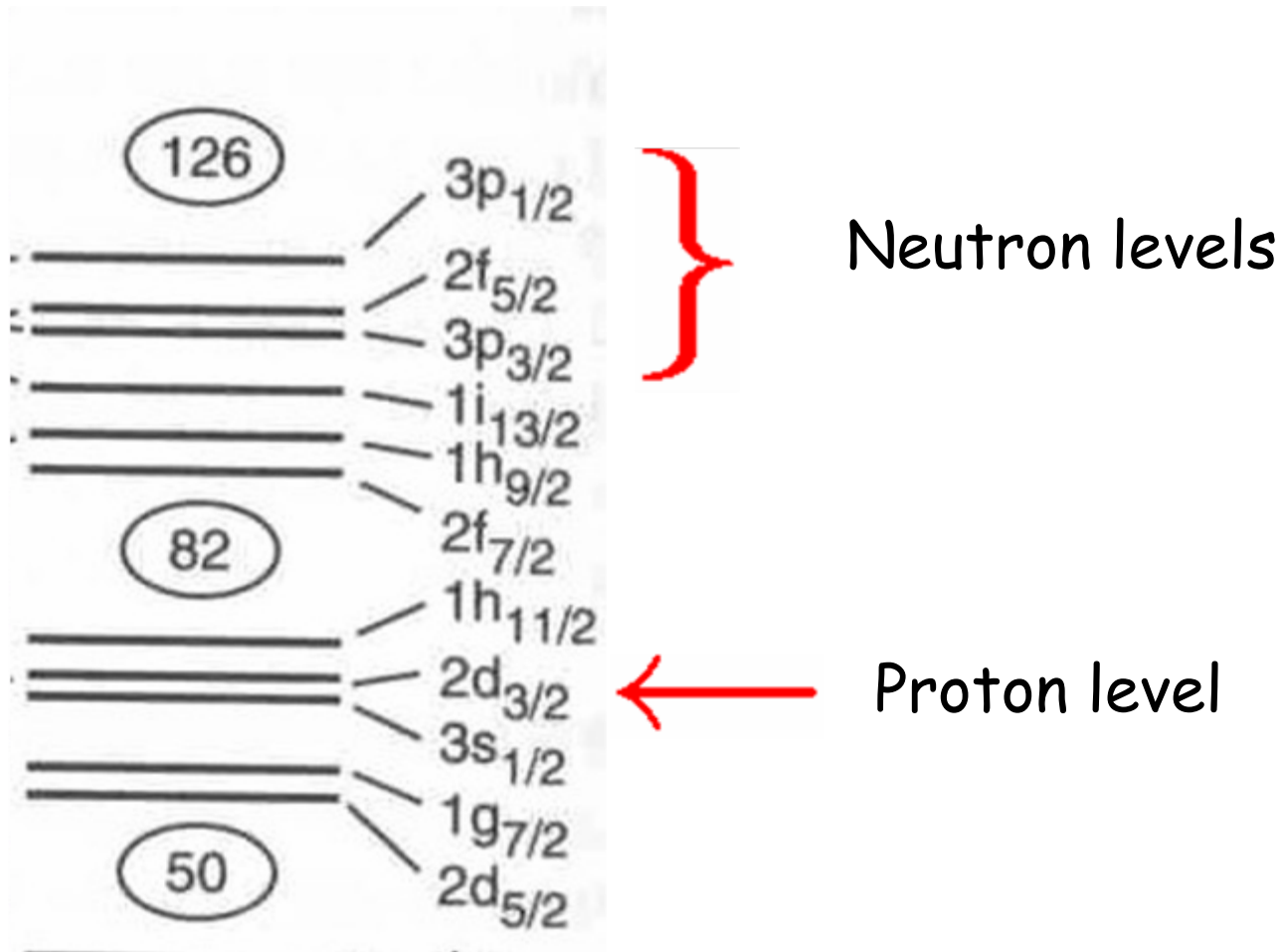
$$H_F = \sum_{\mu\nu} \epsilon_{\mu\nu} A_{\mu\nu} + \sum_{\mu\nu\rho\sigma} v_{\mu\nu\rho\sigma} A_{\mu\nu} A_{\rho\sigma}$$

$$V_{BF} = \sum_{ij\mu\nu} w_{ij\mu\nu} B_{ij} A_{\mu\nu}$$

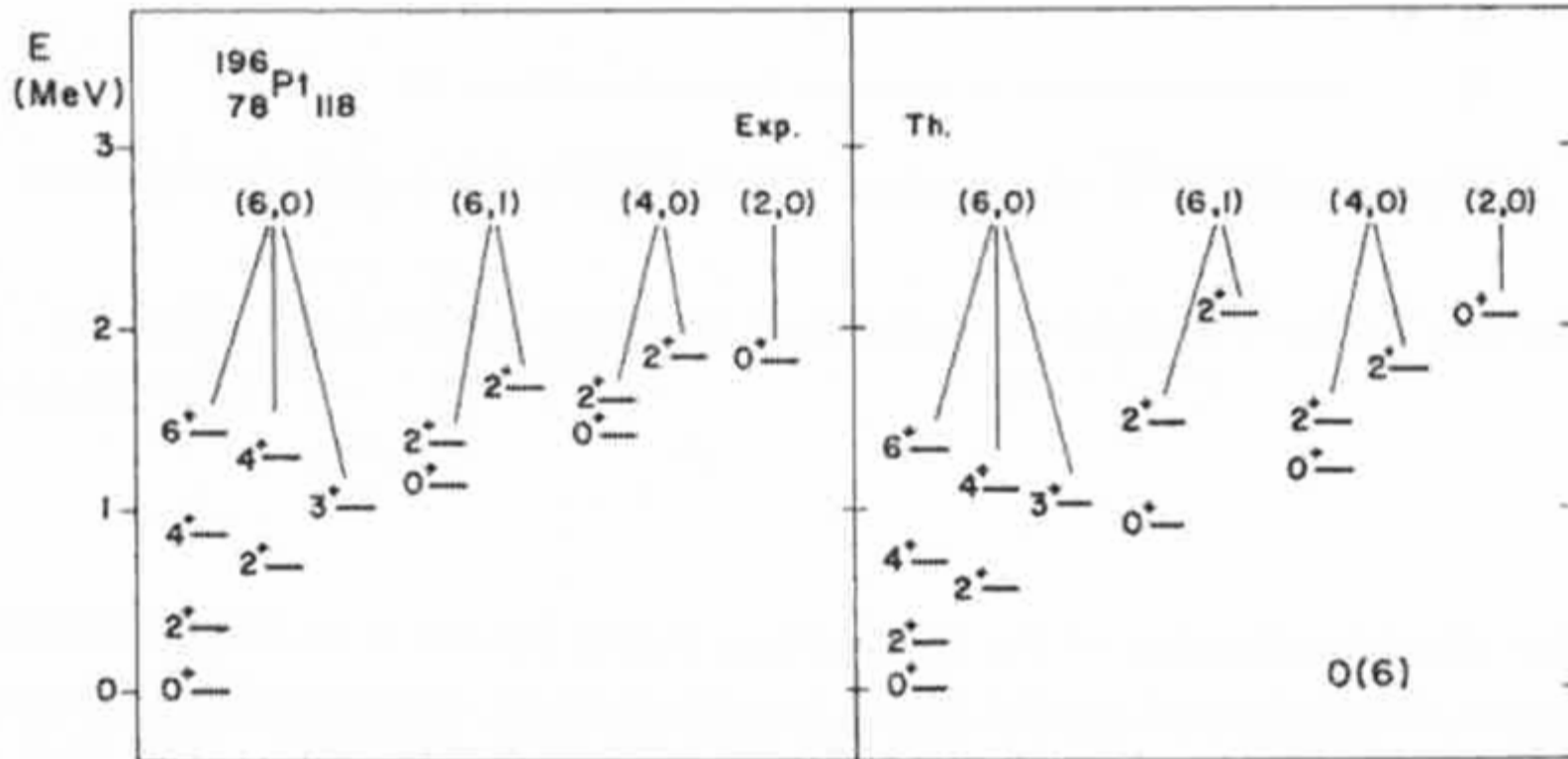
Examples

even-even	s, d	$U_B(6) \supset SO_B(6)$
odd-proton	$j = 2d_{3/2}$	$SU_F(4) \sim SO_B(6)$ spinor
odd-neutron	$j = 3p_{1/2}, 3p_{3/2}, 2f_{5/2}$ $= (\tilde{l} = 0, 2) \otimes (\tilde{s} = \frac{1}{2})$	$U_F(12) \supset U_F(6) \otimes U_F(2)$ $\supset SO_F(6) \otimes U_F(2)$ pseudo-spin

Supersymmetry in Heavy Nuclei

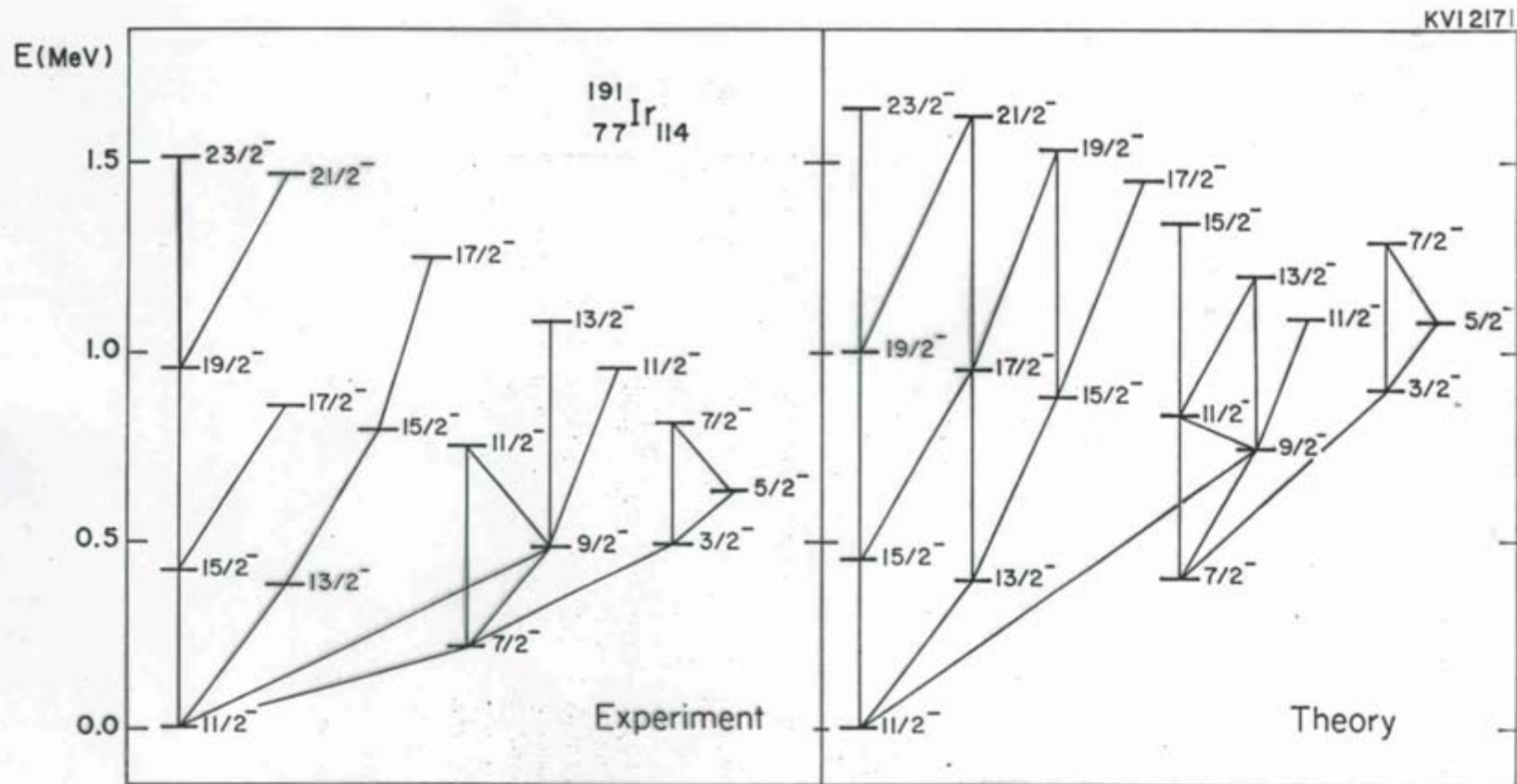


Even-even nucleus



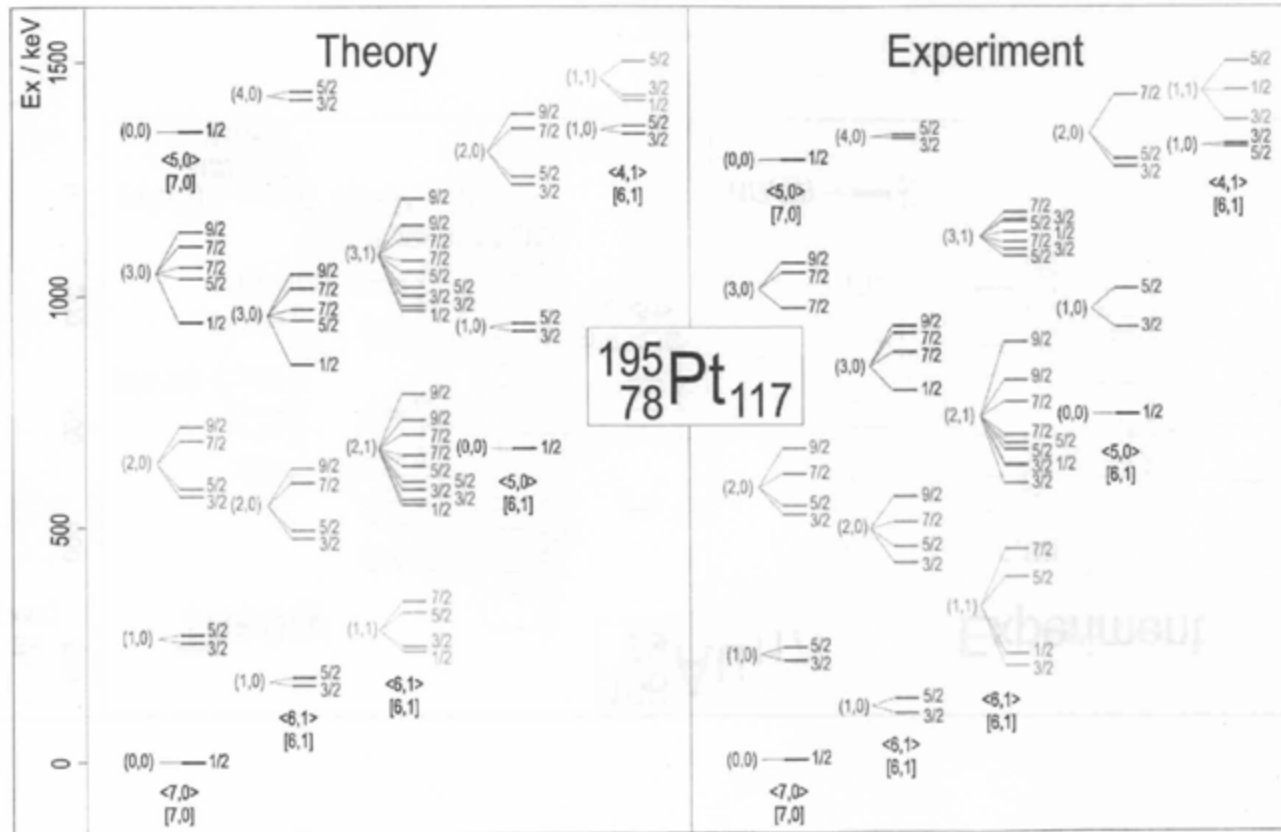
Cizewski et al, PRL 40, 167 (1978)
 Arima, Iachello, PRL 40, 385 (1978)

Odd-proton nucleus



Iachello, PRL 44, 772 (1980)

Odd-neutron nucleus

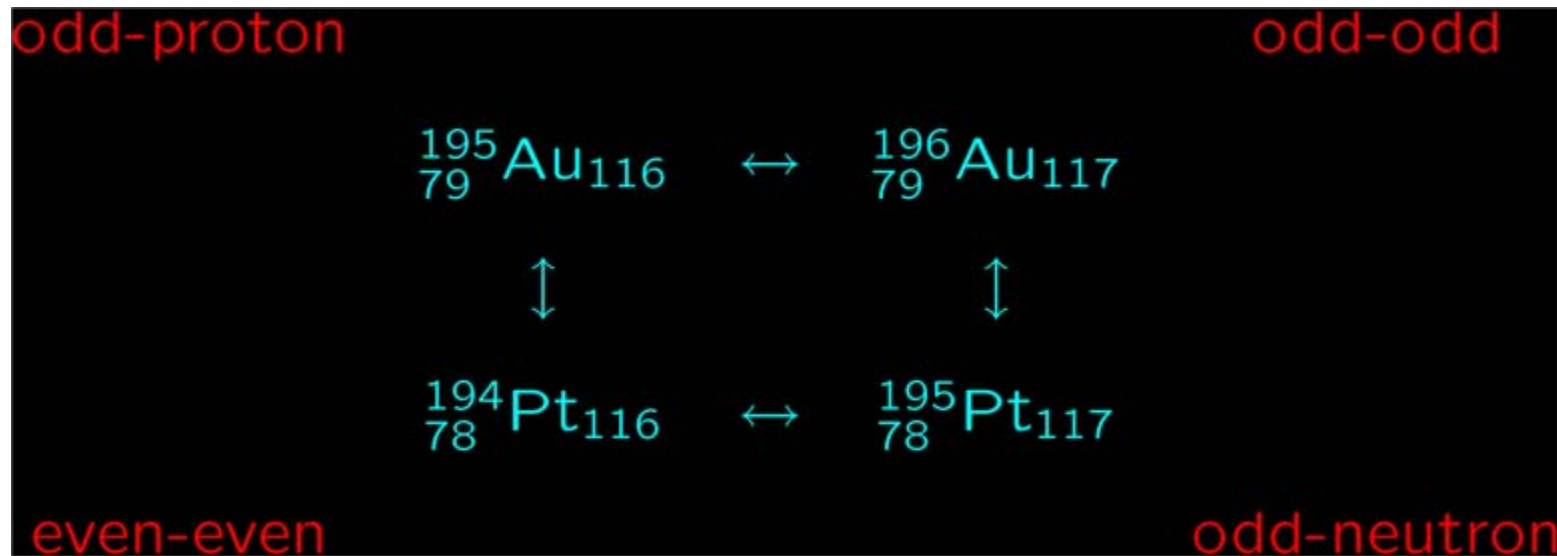


Source: *Journal of Nuclear Energy, Part C*, July 1985

Balantekin, Bars, Bijker, Iachello, PRC 27, 1761 (1983)

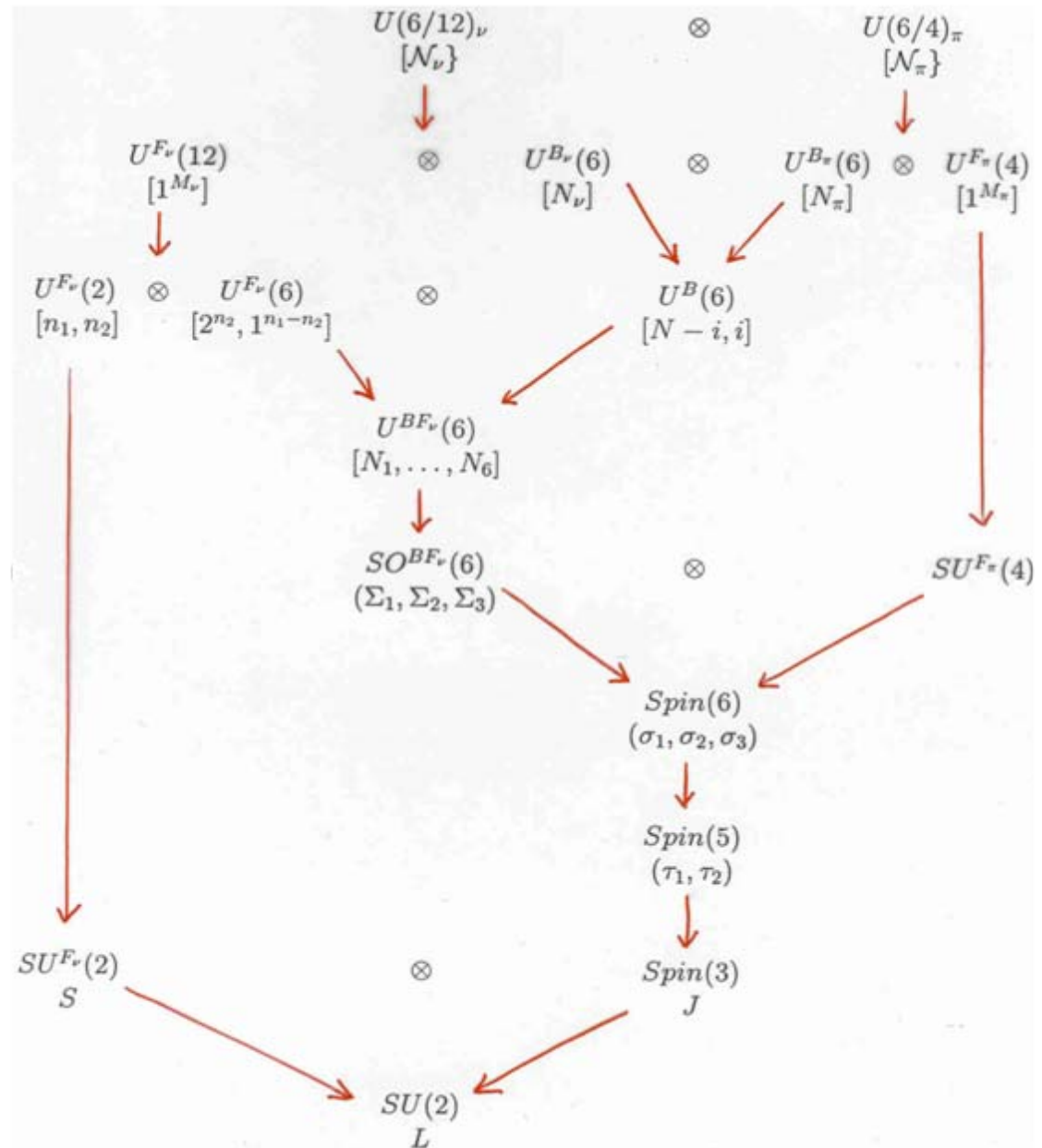
Supersymmetric Quartet of Nuclei

Neutron-proton SUSY : $U(6/12)_\nu \otimes U(6/4)_\pi$



Van Isacker, Jolie, Heyde, Frank, PRL 54, 653 (1985)

Group chain

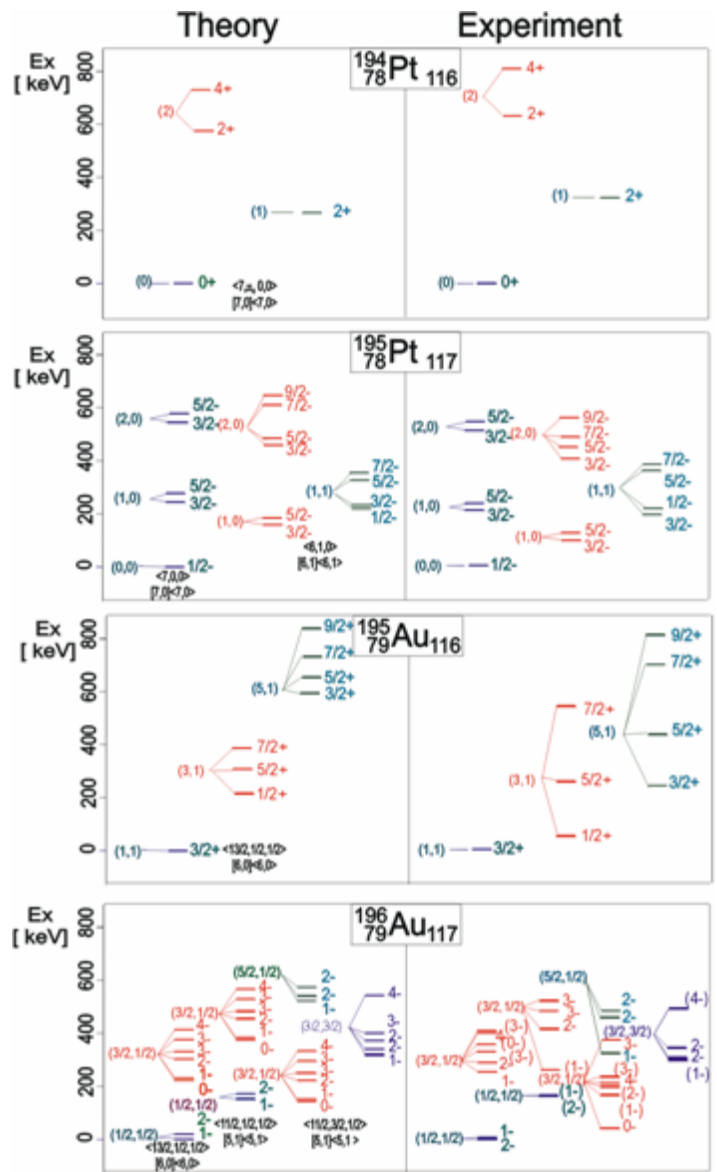


Hamiltonian

$$H = a C_{2UBF_v}(6) + b C_{2SOBF_v}(6) + c C_{2Spin}(6) \\ + d C_{2Spin}(5) + e C_{2Spin}(3) + f C_{2SU}(2)$$

Energies

$$E = a [N_1(N_1 + 5) + N_2(N_2 + 3) + N_3(N_3 + 1)] \\ + b [\Sigma_1(\Sigma_1 + 4) + \Sigma_2(\Sigma_2 + 2) + \Sigma_3^2] \\ + c [\sigma_1(\sigma_1 + 4) + \sigma_2(\sigma_2 + 2) + \sigma_3^2] \\ + d [\tau_1(\tau_1 + 3) + \tau_2(\tau_2 + 1)] \\ + e J(J + 1) + f L(L + 1)$$



One-proton transfer

Test of the fermionic generators of the superalgebra !

$$P_{\pi,1}^{(\frac{3}{2})\dagger} = -\sqrt{\frac{1}{6}} \left(\tilde{s}_{\pi} \times a_{\pi,\frac{3}{2}}^{\dagger} \right)^{(\frac{3}{2})} + \sqrt{\frac{5}{6}} \left(\tilde{d}_{\pi} \times a_{\pi,\frac{3}{2}}^{\dagger} \right)^{(\frac{3}{2})}$$
$$P_{\pi,2}^{(\frac{3}{2})\dagger} = +\sqrt{\frac{5}{6}} \left(\tilde{s}_{\pi} \times a_{\pi,\frac{3}{2}}^{\dagger} \right)^{(\frac{3}{2})} + \sqrt{\frac{1}{6}} \left(\tilde{d}_{\pi} \times a_{\pi,\frac{3}{2}}^{\dagger} \right)^{(\frac{3}{2})}$$

Barea, Bijker, Frank, JPA 37, 10251 (2004)

$^{194}\text{Pt} \rightarrow ^{195}\text{Au}$

$$|\langle f || P_{\pi,1}^{(\frac{3}{2})\dagger} || i \rangle|^2$$

$$|\langle f || P_{\pi,2}^{(\frac{3}{2})\dagger} || i \rangle|^2$$

$$\langle (N + \frac{3}{2}, \frac{1}{2}, \frac{1}{2}), (\frac{1}{2}, \frac{1}{2}), \frac{3}{2} |$$

$$\frac{2(N_{\pi}+1)}{3}$$

$$\frac{8(N+6)^2(N_{\pi}+1)}{15(N+3)^2}$$

$$\langle (N + \frac{1}{2}, \frac{1}{2}, -\frac{1}{2}), (\frac{1}{2}, \frac{1}{2}), \frac{3}{2} |$$

$$0$$

$$\frac{6(N+1)(N+5)(N_{\pi}+1)}{5(N+3)^2}$$

$^{195}\text{Pt} \rightarrow ^{196}\text{Au}$

$$|\langle f || P_{\pi,1}^{(\frac{3}{2})\dagger} || i \rangle|^2$$

$$|\langle f || P_{\pi,2}^{(\frac{3}{2})\dagger} || i \rangle|^2$$

$$\langle (N + \frac{3}{2}, \frac{1}{2}, \frac{1}{2}), (\frac{1}{2}, \frac{1}{2}), \frac{3}{2}, L |$$

$$\frac{2(N_{\pi}+1)}{3} \frac{2L+1}{4}$$

$$\frac{8(N+6)^2(N_{\pi}+1)}{15(N+3)^2} \frac{2L+1}{4}$$

$$\langle (N + \frac{1}{2}, \frac{1}{2}, -\frac{1}{2}), (\frac{1}{2}, \frac{1}{2}), \frac{3}{2}, L |$$

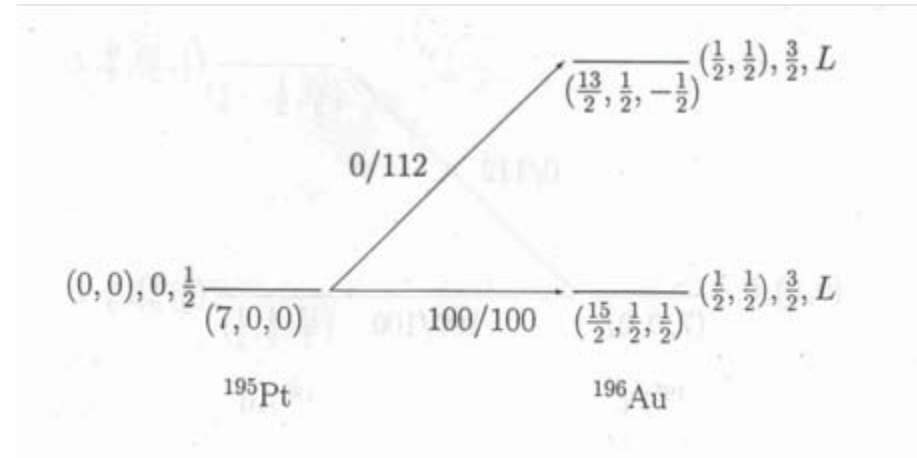
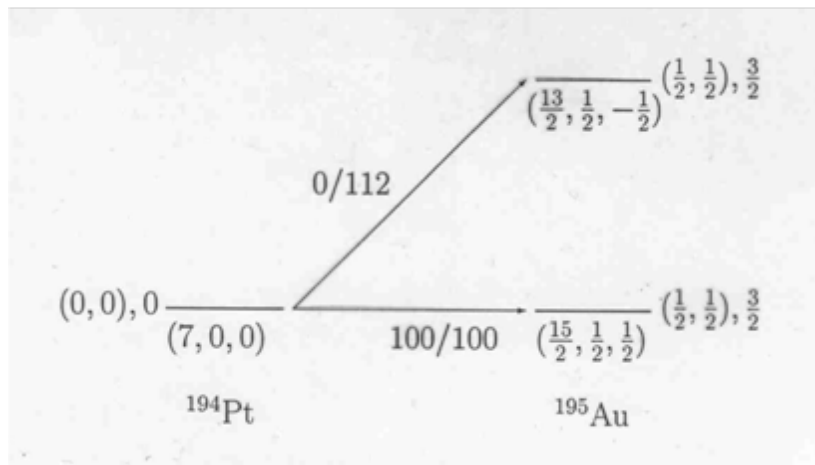
$$0$$

$$\frac{6(N+1)(N+5)(N_{\pi}+1)}{5(N+3)^2} \frac{2L+1}{4}$$

Correlations

$$R_1 = \frac{I_{gs \rightarrow exc}}{I_{gs \rightarrow gs}} = 0$$

$$R_2 = \frac{I_{gs \rightarrow exc}}{I_{gs \rightarrow gs}} = \frac{9(N+1)(N+5)}{4(N+6)^2}$$



Correlations

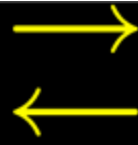
- One-proton transfer reactions

$$\frac{S_i(^{195}\text{Pt} \rightarrow ^{196}\text{Au})}{S_i(^{194}\text{Pt} \rightarrow ^{195}\text{Au})} = \frac{2L + 1}{8}$$

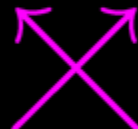
- One-neutron transfer reactions

$$\frac{S_i(^{195}\text{Pt} \rightarrow ^{194}\text{Pt})}{S_i(^{194}\text{Pt} \rightarrow ^{195}\text{Pt})} = \begin{cases} \frac{1}{2} \\ \frac{N_\pi + 1}{2(N + 1)(N_\nu + 1)} \end{cases}$$

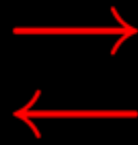
${}_{79}^{195}\text{Au}_{116}$



${}_{79}^{196}\text{Au}_{117}$



${}_{78}^{194}\text{Pt}_{116}$



${}_{78}^{195}\text{Pt}_{117}$

Two-nucleon transfer

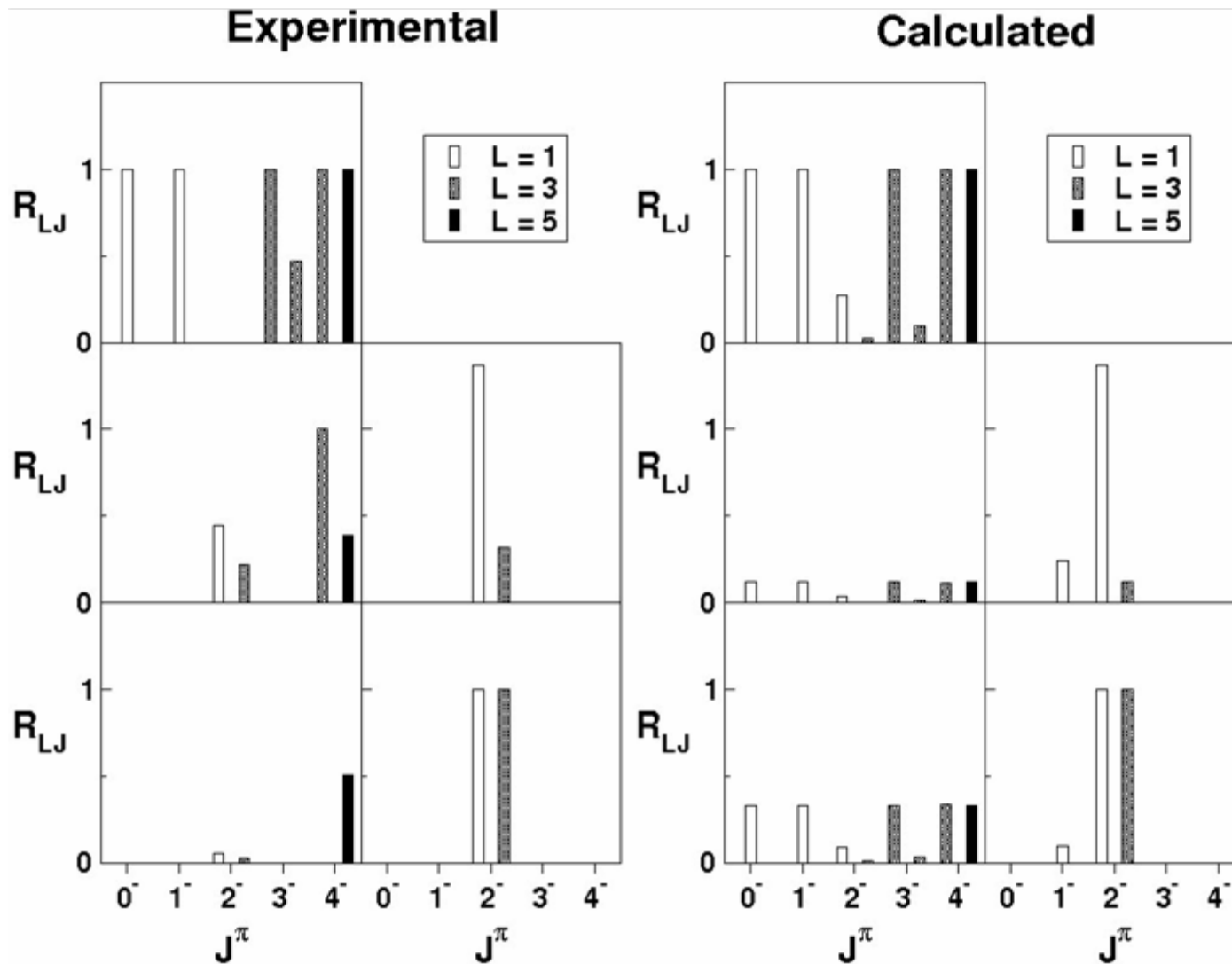
Reaction $^{198}\text{Hg} (d, \alpha) ^{196}\text{Au}$

Spectroscopic factors

$$G_{LJ} = \left| \sum_{j\nu j\pi} g_{j\nu j\pi}^{LJ} \langle ^{196}\text{Au} \| (a_{j\nu}^\dagger a_{j\pi}^\dagger)^{(\lambda)} \| ^{198}\text{Hg} \rangle \right|^2$$

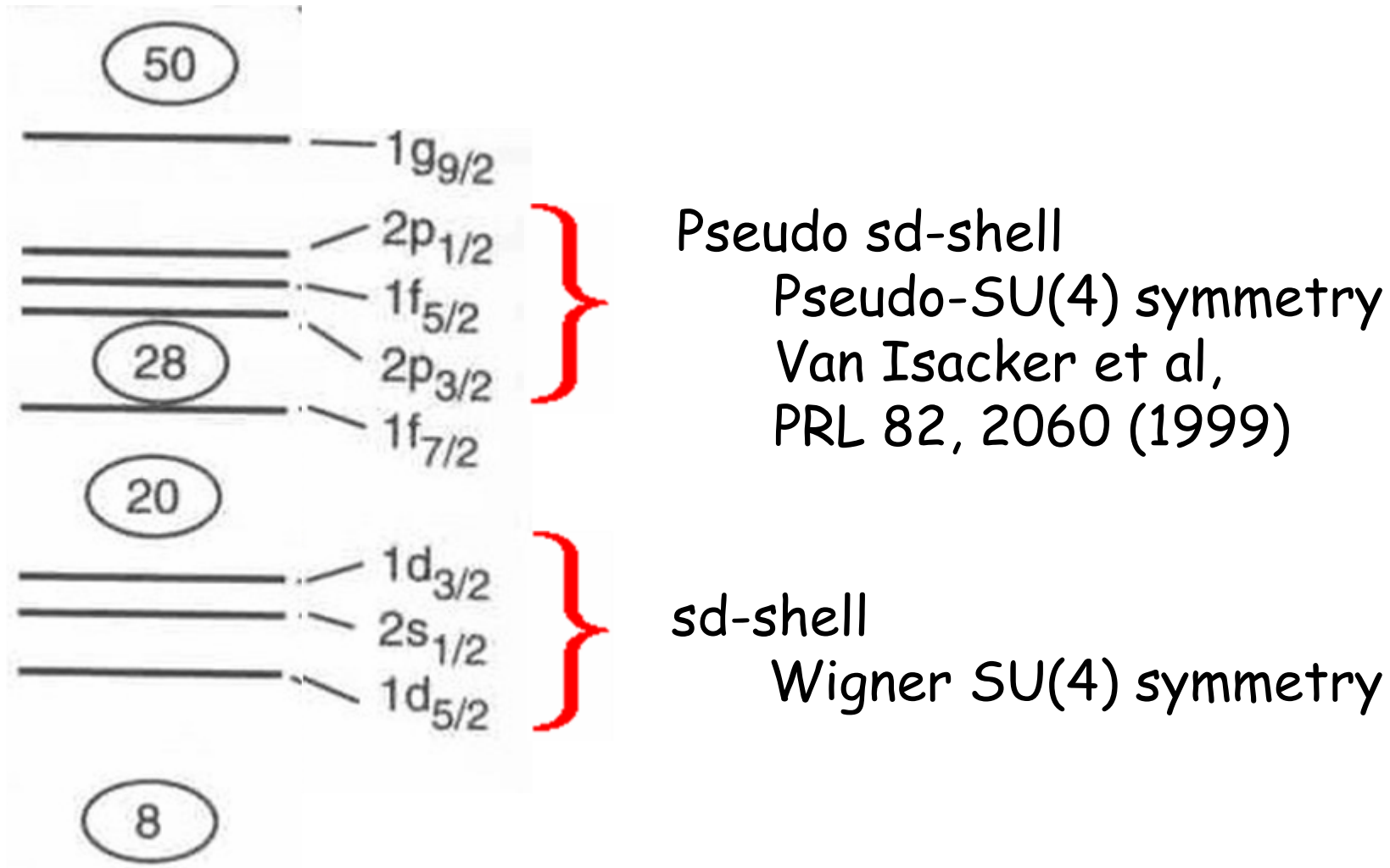
Relative strength

$$R_{LJ} = \frac{G_{LJ}}{G_{LJ}(\text{ref})} = \begin{cases} \frac{N+4}{15N} & = 0.12 \\ \frac{2(N+4)(N+6)}{15N(N+3)} & = 0.33 \end{cases}$$



Barea, Bijker, Frank, PRL 94, 152501 (2005)

Supersymmetry in Light Nuclei



Interacting Boson Models

- Heavy nuclei: protons and neutrons in different major shells

IBM-1	s^\dagger, d^\dagger	$U(6)$
IBM-2	$(s_\nu^\dagger, d_\nu^\dagger) \otimes (s_\pi^\dagger, d_\pi^\dagger)$ $T = 1, M_T = \pm 1$	$U_\nu(6) \otimes U_\pi(6)$

- Light nuclei: protons and neutrons occupy same major shells \Rightarrow isospin invariant IBM

IBM-3	s^\dagger, d^\dagger $(S, T) = (0, 1)$	$U(18)$
IBM-4	s^\dagger, d^\dagger $(S, T) = (0, 1), (1, 0)$	$U(36)$

Elliott, White, PLB 97, 169 (1980)

Elliott, Evans, PLB 101, 216 (1981)

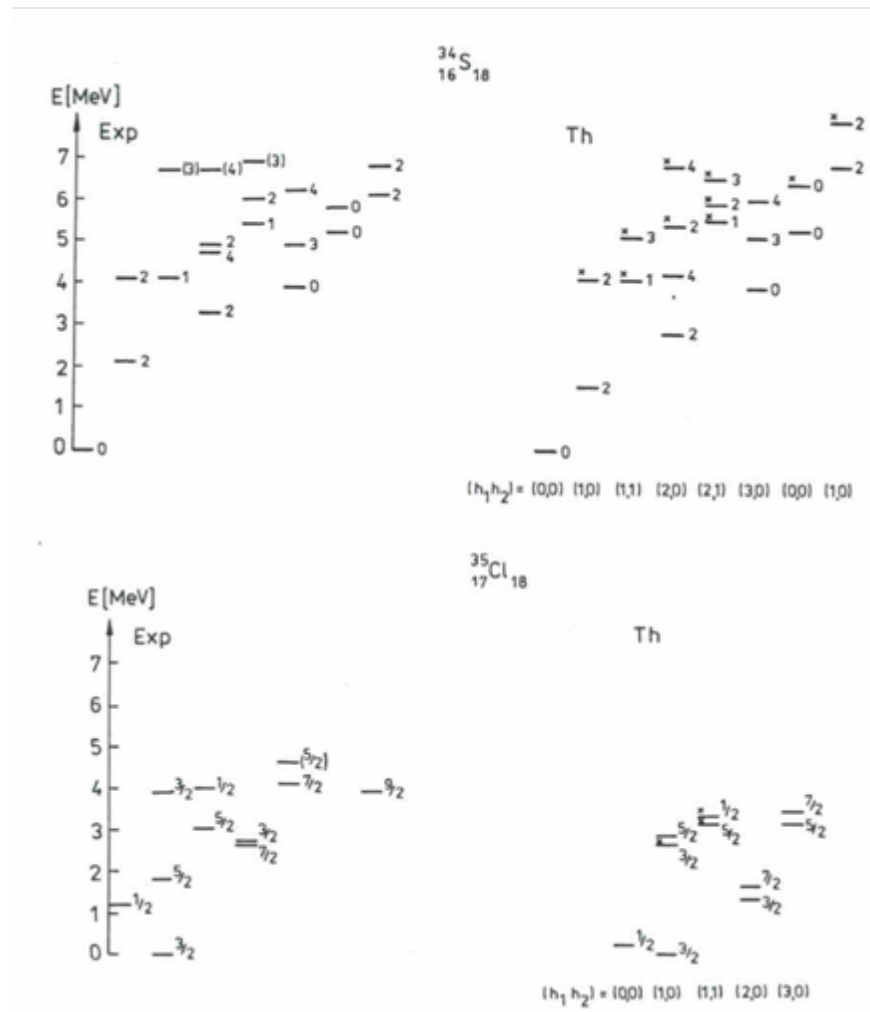
Dynamical Supersymmetry

$$\begin{aligned} \text{bosons : } U^B(36) &\supset U_L^B(6) \otimes SU_{ST}^B(6) \\ &\supset U_L^B(6) \otimes SU_{ST}^B(4) \end{aligned}$$

$$\text{fermions : } U^F(24) \supset U_L^F(6) \otimes SU_{ST}^F(4)$$

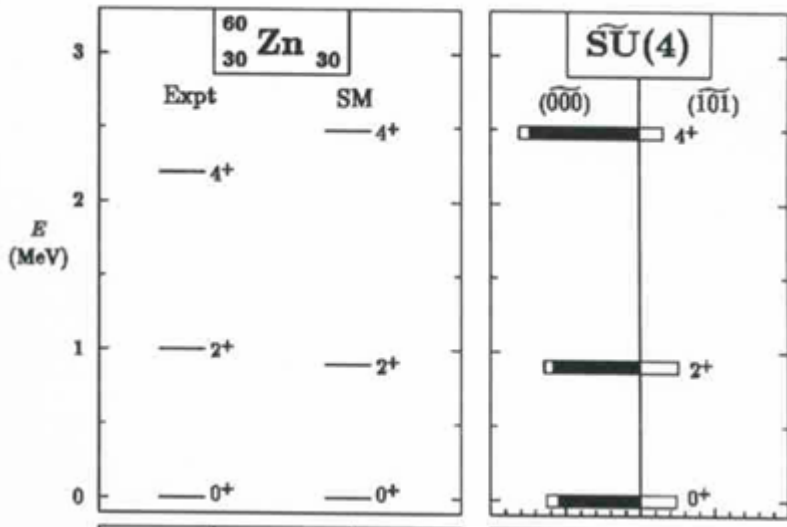
$$\begin{aligned} U(36/24) &\supset U^B(36) \otimes U^F(24) \\ &\supset U_L^B(6) \otimes SU_{ST}^B(4) \otimes U_L^F(6) \otimes SU_{ST}^F(4) \\ &\supset U_L^{BF}(6) \otimes SU_{ST}^{BF}(4) \\ &\supset \dots \end{aligned}$$

Example in the sd-shell



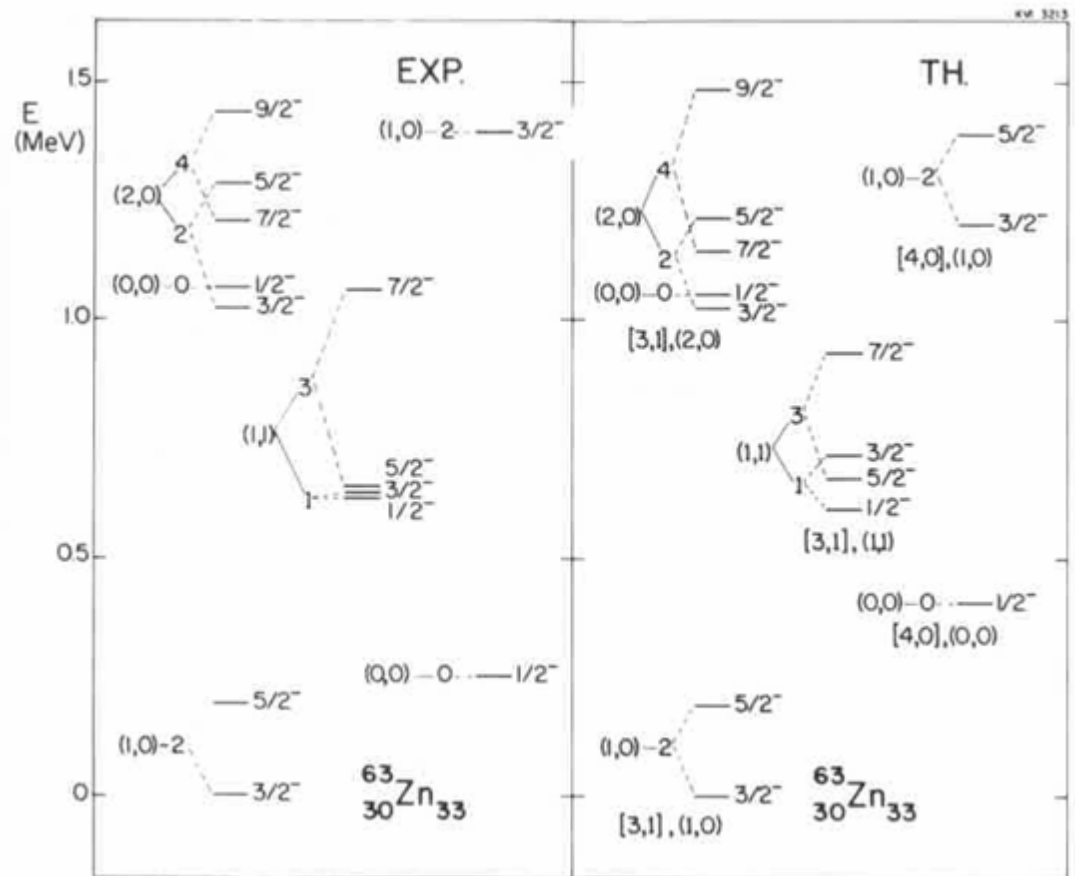
Szpikowski et al,
NPA 487, 301 (1988)

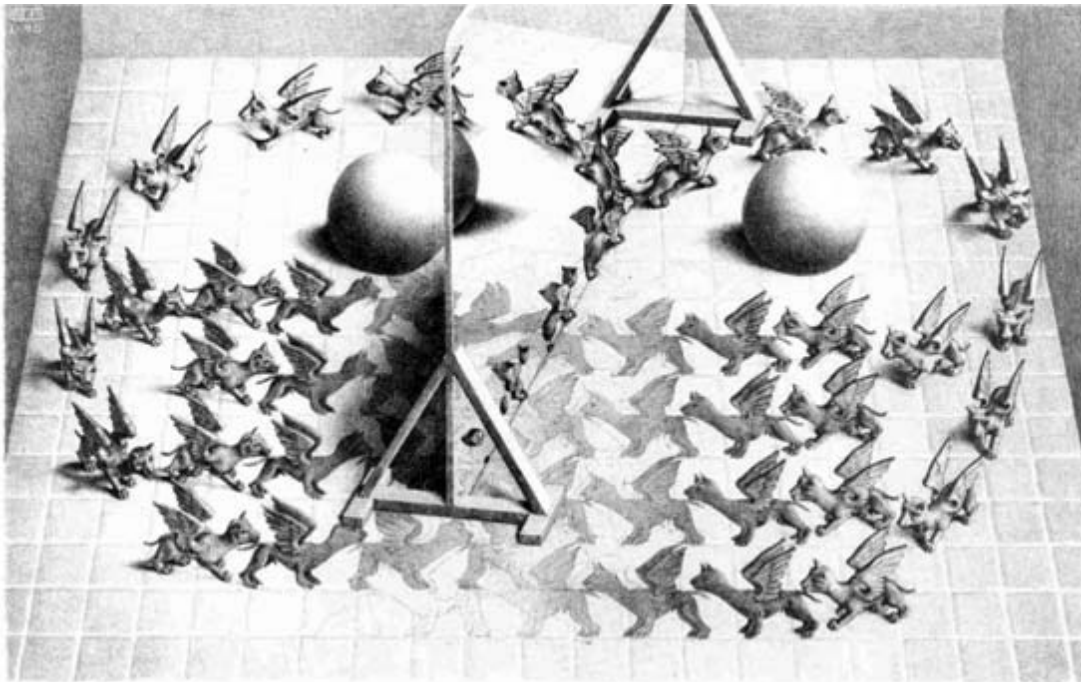
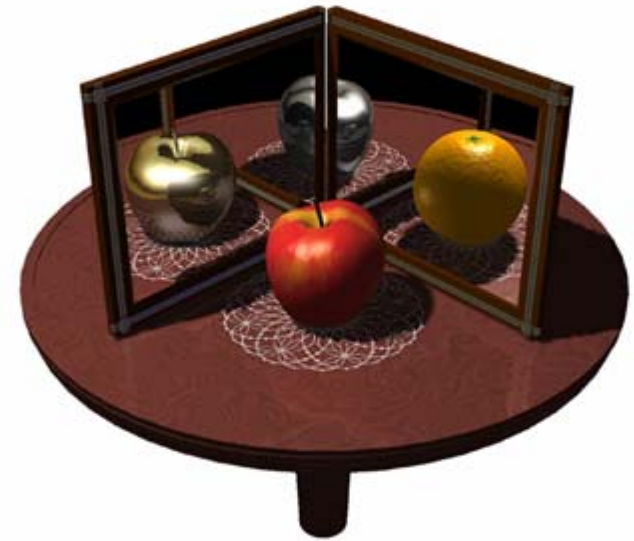
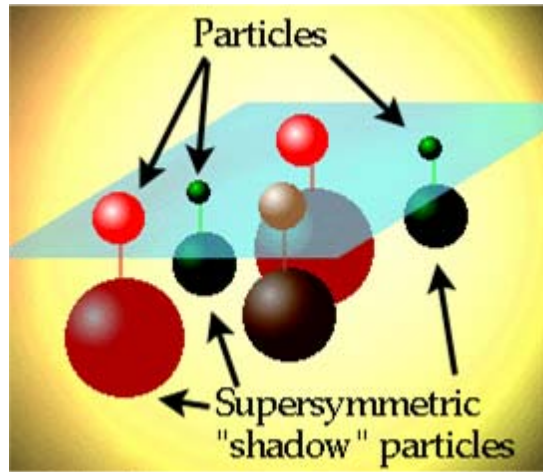
Example in the pf-shell



Van Isacker et al,
PRL 82, 2060 (1999)

R. Bijker,
Ph.D. Thesis, 1984





Magic mirror
M.C. Escher

Summary and Conclusions

- **Nuclear supersymmetry**: energy formulas, selection rules, transition rates, etc.
- **Supersymmetry** leads to correlations between different transfer reactions
- Applications in both heavy and light nuclei
- **Proton-rich nuclei**: dynamical (super)symmetries of isospin invariant IBM and IBFM?
- **Neutron-rich nuclei**: are there additional degrees of freedom (valence protons, valence neutrons, skins), what are the corresponding symmetries?
- SUSY without dynamical symmetry
- Predictability