

The G^0 Forward-Angle Measurement of Parity Violating Asymmetries in $\bar{e}p$ Elastic Scattering

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for the G^0 Collaboration

Workshop on Precision ElectroWeak Interactions

College of William and Mary

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Virtually all slides are from Doug Beck's 17-June-2005 data release seminar at JLab.

GO Collaboration

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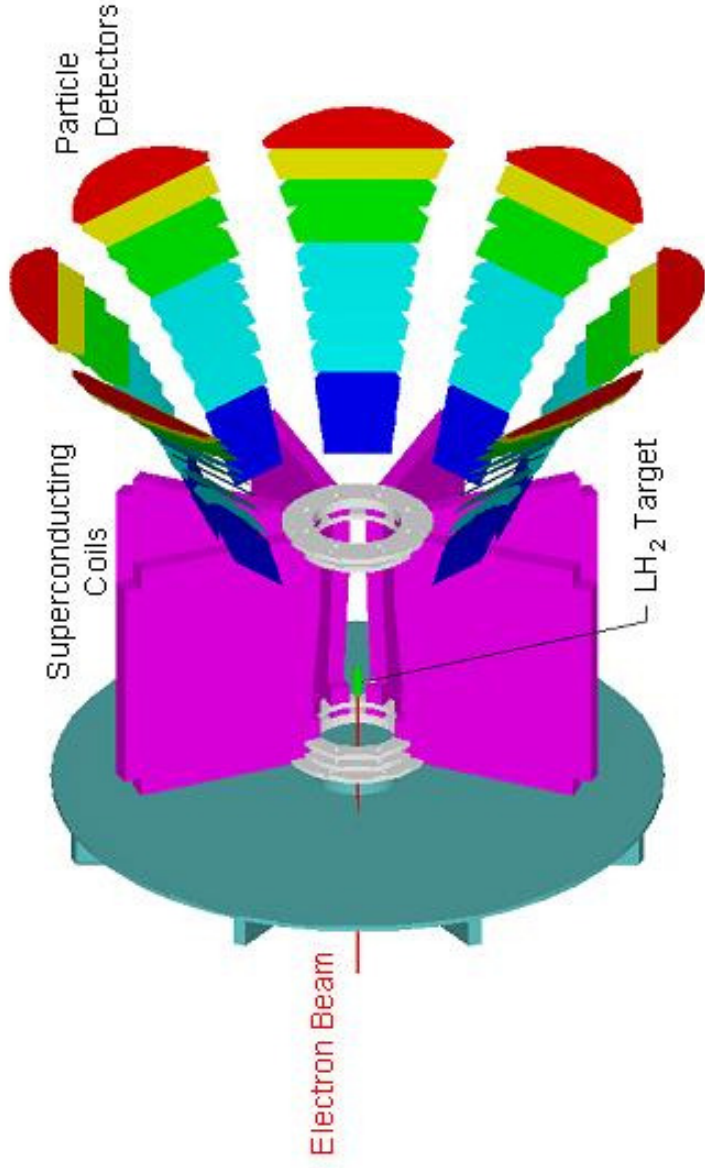
Experimental setup



GO Experiment Overview

- **Measure G_E^Z, G_M^Z**
 - different linear combinations of u, d and s contributions than e.m. form factors
 - strange quark contributions to sea
- **Measure forward and backward asymmetries**
 - recoil protons for forward measurement
 - electrons for backward measurements
 - elastic/inelastic for ^1H , elastic for ^2H
- **Forward measurements complete (101 Coulombs)**

$E_{\text{beam}} = 3.03 \text{ GeV}, 0.33 - 0.93 \text{ GeV}$
$I_{\text{beam}} = 40 \mu\text{A}, 80 \mu\text{A}$
$P_{\text{beam}} = 75\%, 80\%$
$\theta = 52 - 76^0, 104 - 116^0$
$\Delta\Omega = 0.9 \text{ sr}, 0.5 \text{ sr}$
$l_{\text{target}} = 20 \text{ cm}$
$L = 2.1, 4.2 \times 10^{38} \text{ cm}^{-2} \text{ s}^{-1}$
A ~ -1 to -50 ppm, -12 to -70 ppm



GO in Hall C

superconducting magnet (SMS)

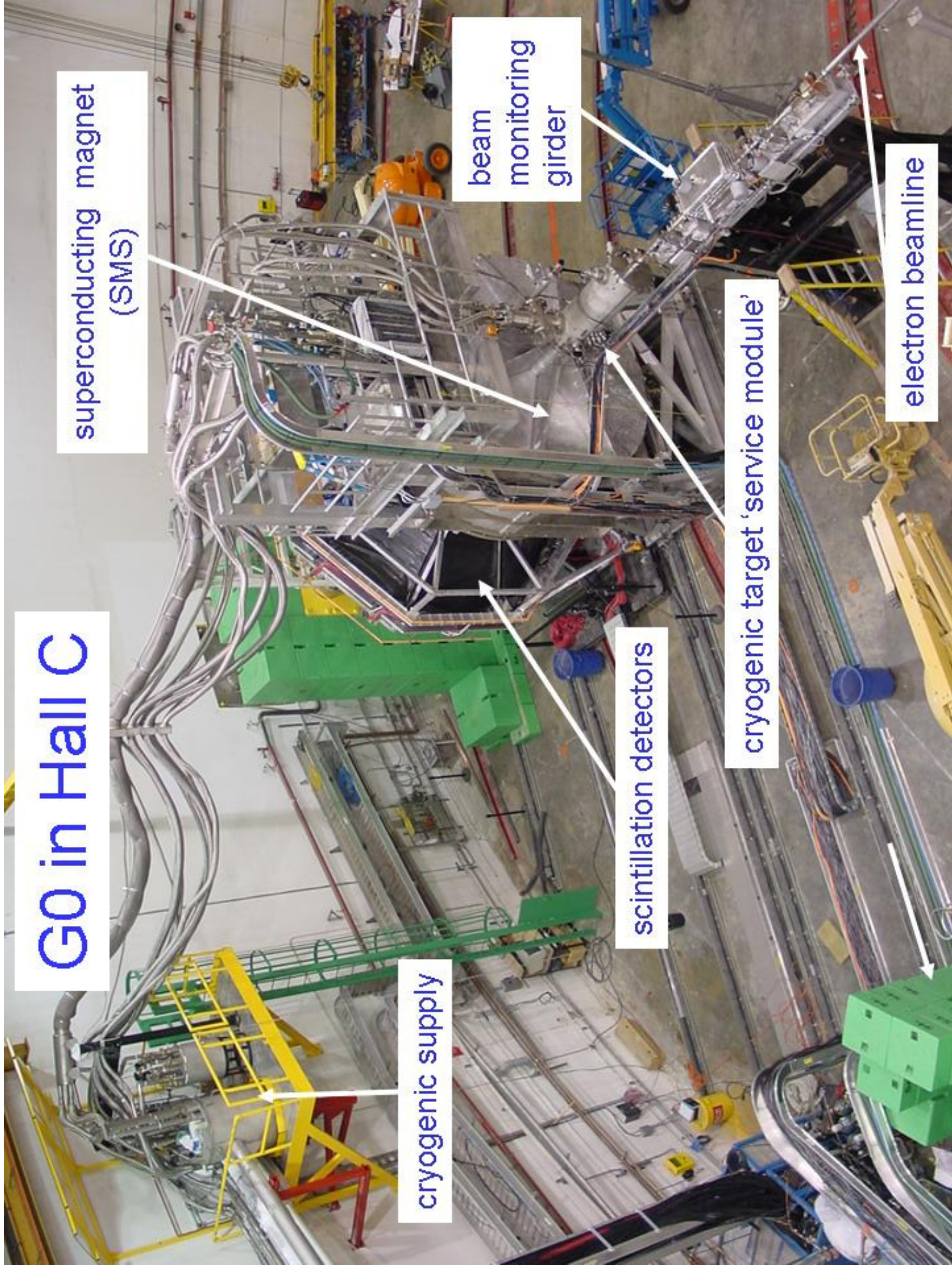
beam monitoring girder

electron beamline

cryogenic target 'service module'

scintillation detectors

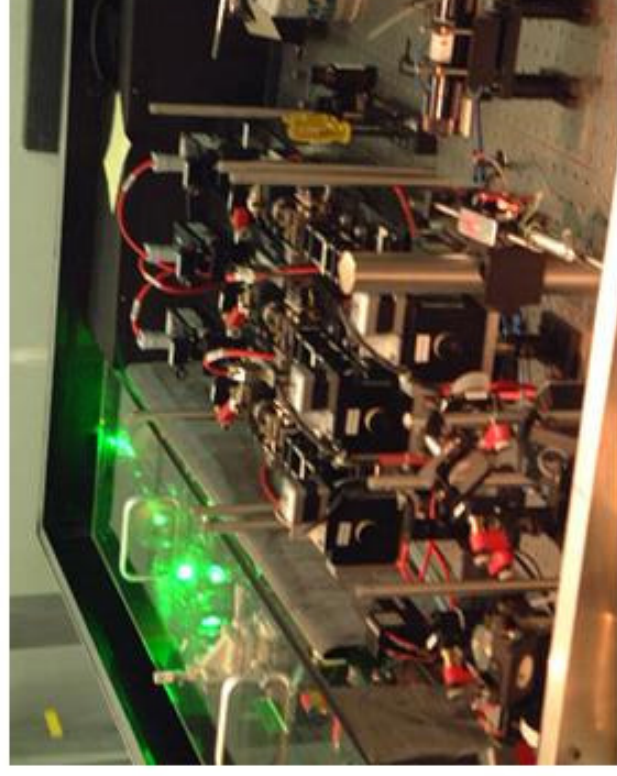
cryogenic supply



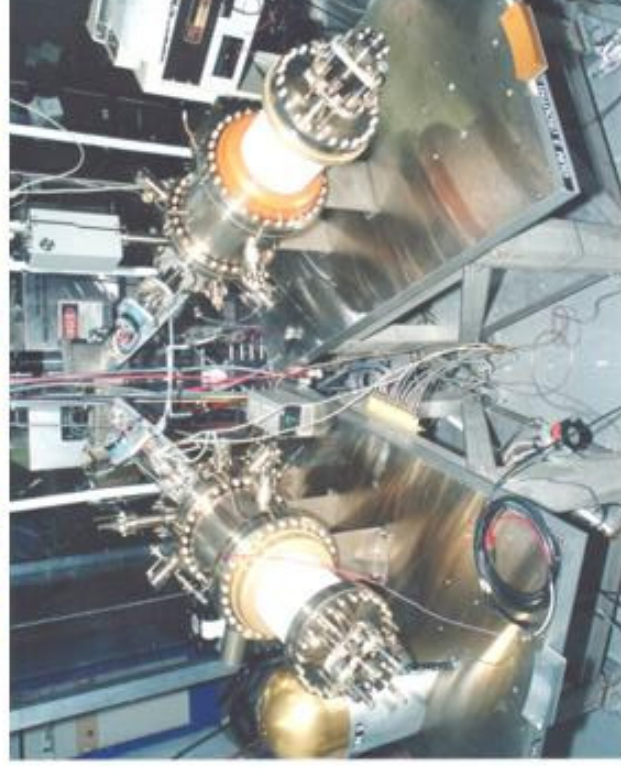
Polarized Injector/Accelerator

- Challenging specifications – all met!
 - 32 ns pulse spacing for t.o.f.
 - 40 μA beam current
 - higher bunch charge
 - run concurrently with small energy spread for Hall A

Beam Parameter	Achieved	“Specs”
Charge asymmetry	-0.14 ± 0.32 ppm	1 ppm
x position differences	3 ± 4 nm	20 nm
y position differences	4 ± 4 nm	20 nm
x angle differences	1 ± 1 nrad	2 nrad
y angle differences	1.5 ± 1 nrad	2 nrad
Energy differences	29 ± 4 eV	75 eV



New Tiger laser system for G0



JLab polarized injector

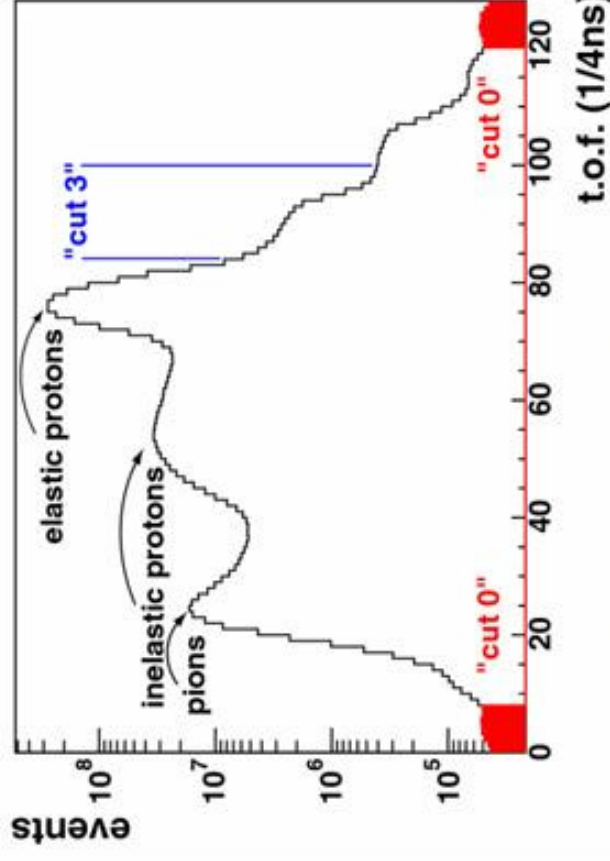
Leakage Beam Measurement

- Use "cut0" region in actual data to measure leakage yield, asymmetry throughout run
- Cut0 certified during test runs with only leakage beam
 - uncertainty determined in 3 ways

- compare lumi monitor (direct) measurements to cut0
- cut3 asymmetry independent of beam current (10, 20, 40 μA)
- variation of corrected cut3 asymmetry (should be constant over run)

- methods consistent at 20% level

- $\delta A_{\text{false,leak}} = -0.71 \pm 0.14 \text{ ppm}$

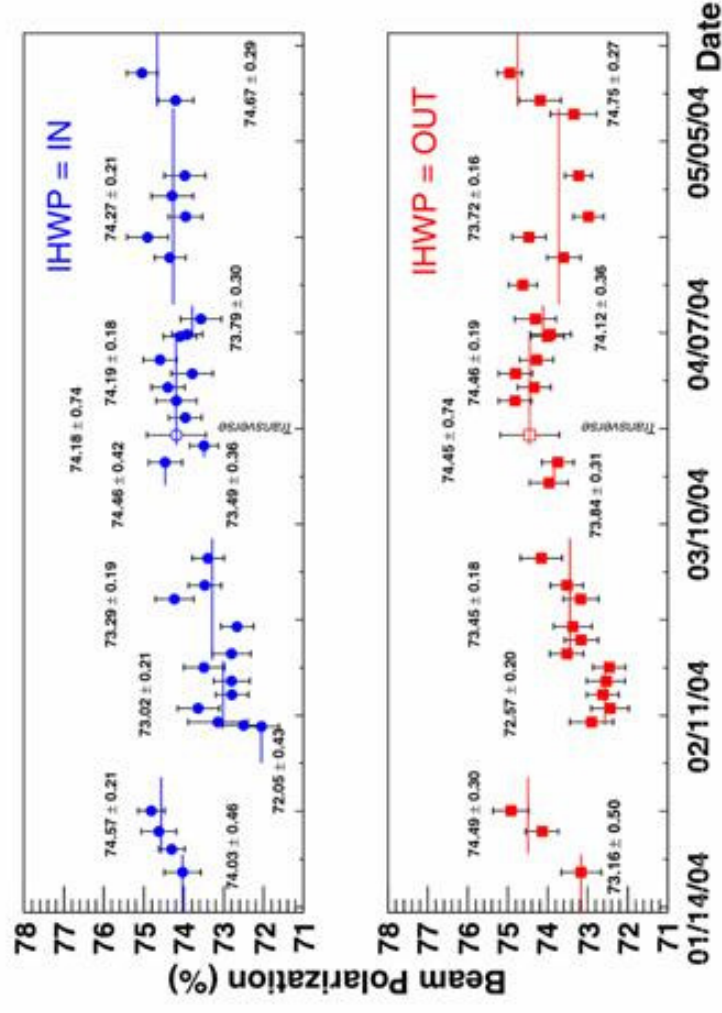


Leakage beam measurement regions

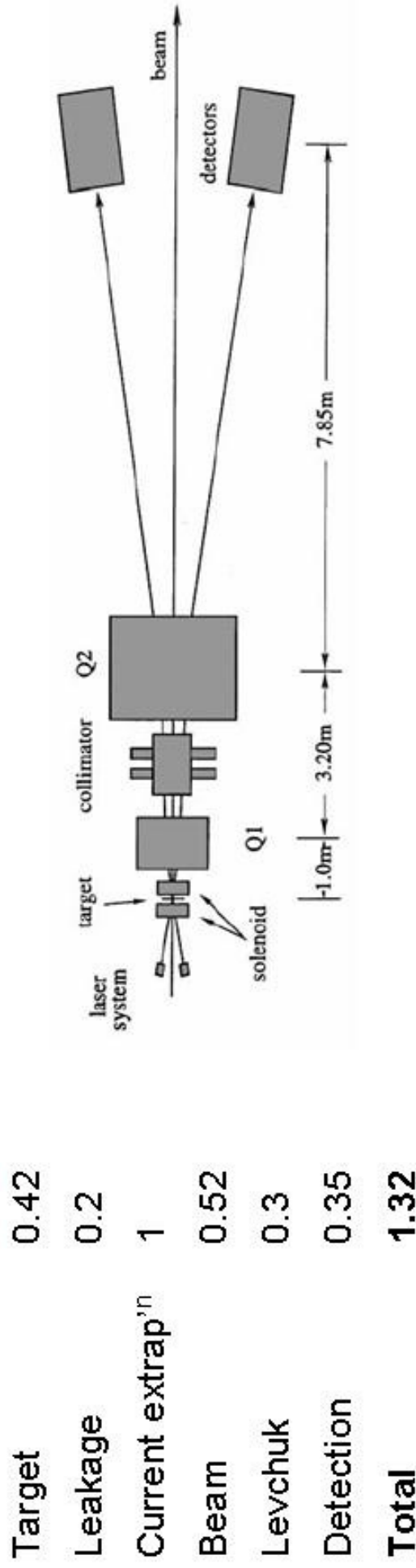
I (μA)	$A_{3,\text{meas}}$ (ppm)	$A_{3,\text{corr}}$ (ppm)
40	0.14 ± 0.43	-2.5 ± 0.43
20	-29.6 ± 2.1	-7.2 ± 2.1
10	-51.3 ± 3.9	-9.5 ± 3.9

Beam Polarization

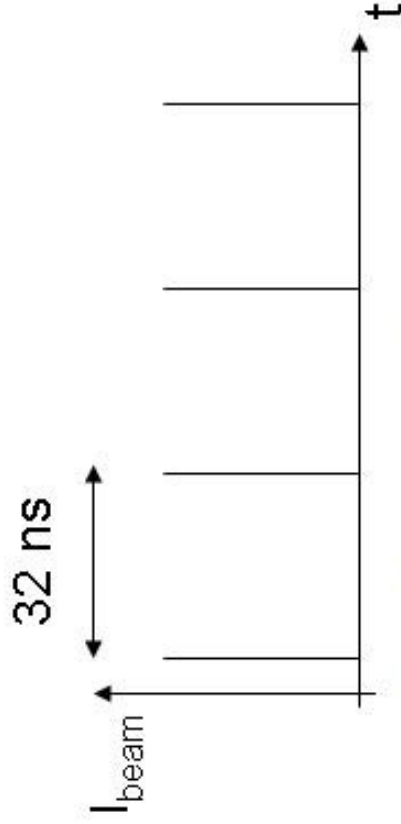
- **Beam polarization measured with interleaved Møller measurements**
 - std Hall C polarimeter (M. Hauger, et al. NIM **A462** (2001) 382.)
 - apply for groups of runs as shown
 - average: $P = 73.7\%$



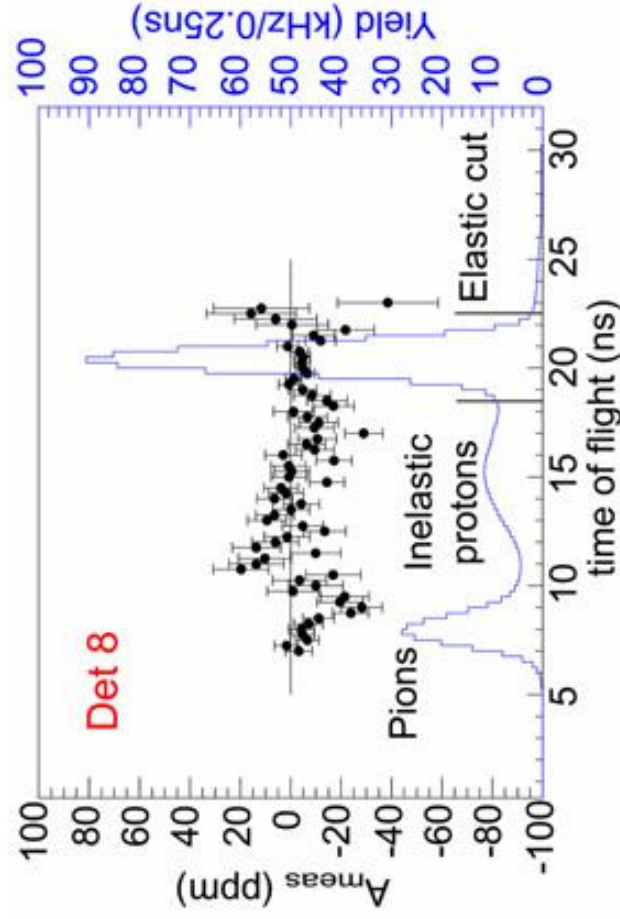
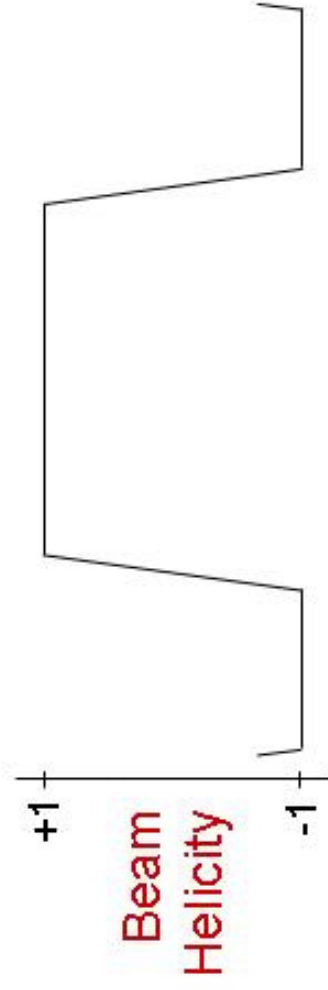
Source Rel. uncertainty (%)



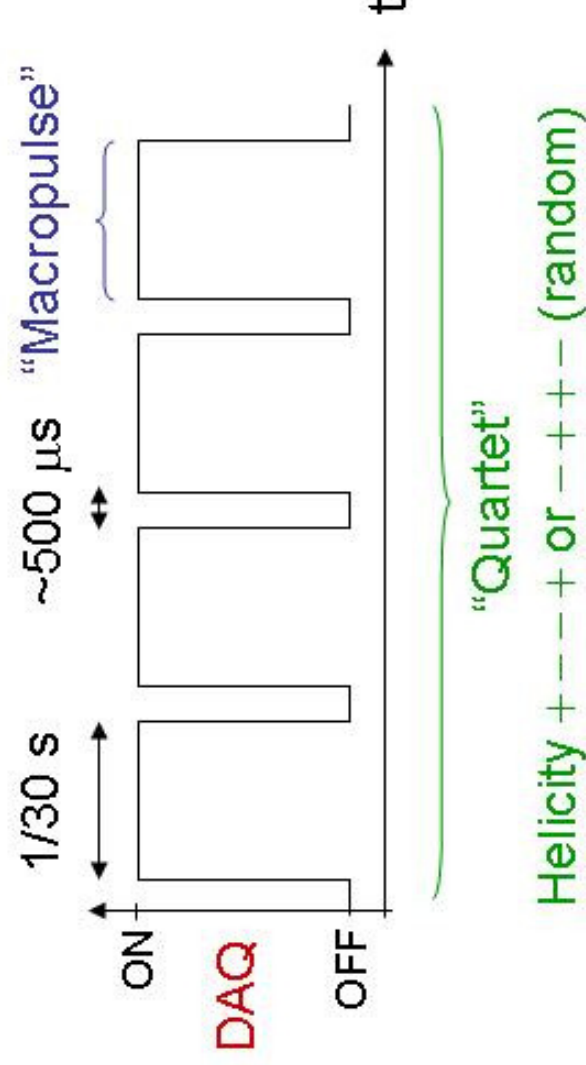
Timing in the Experiment



Accelerator pulse structure

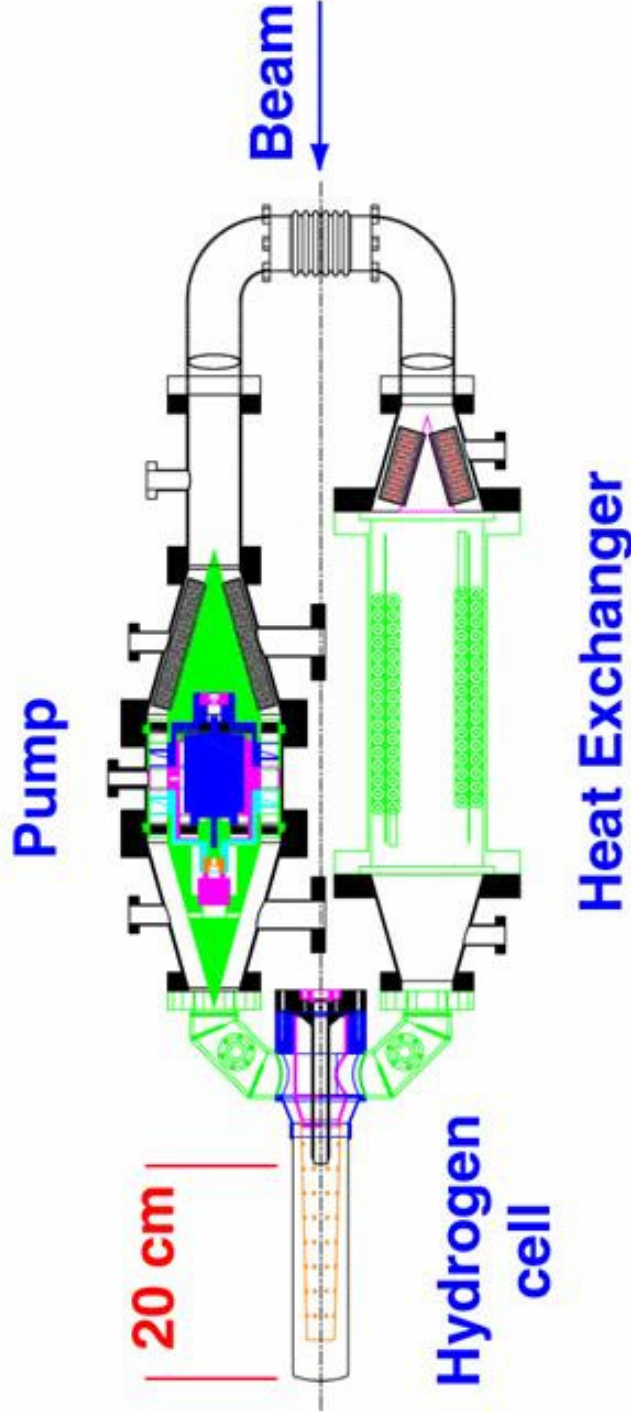


Typical t.o.f. spectrum

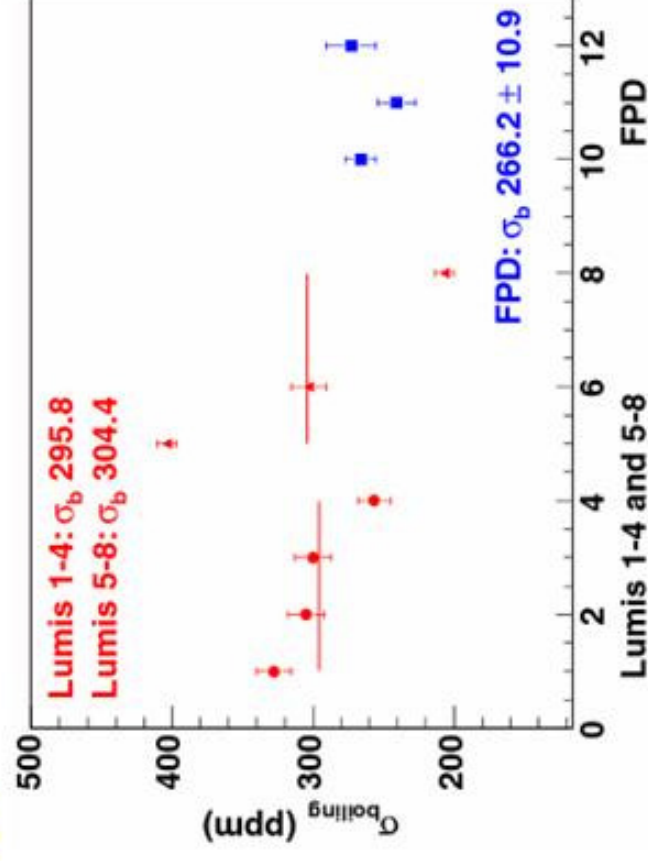


Measurement timing

Target

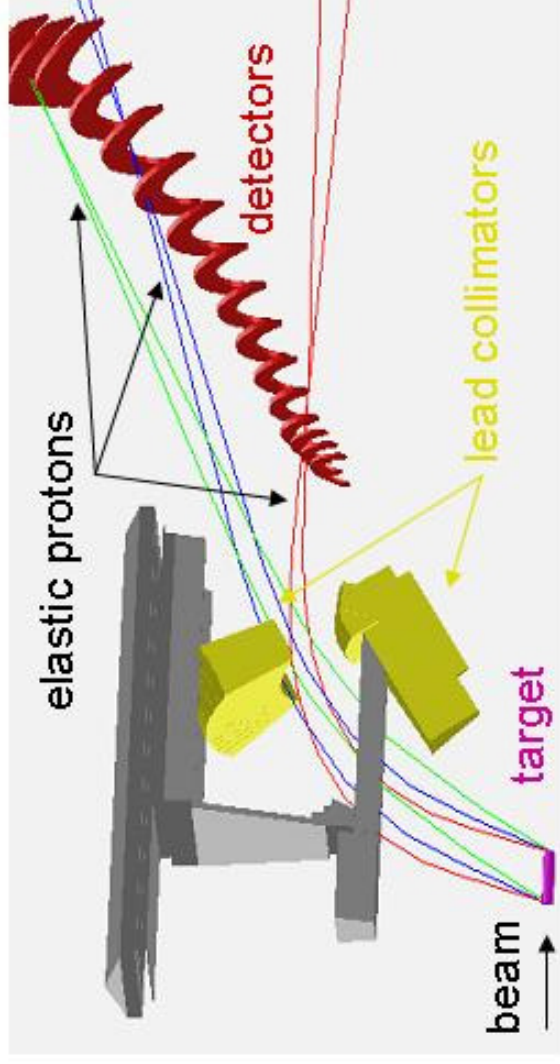


- 20 cm LH₂, aluminum target cell
- longitudinal flow, $v \sim 8$ m/s,
P > 1000 W!
- negligible density change < 1.5%
- measured small boiling contribution
 - 260 ppm/1200 ppm statistical width



Spectrometer Optics

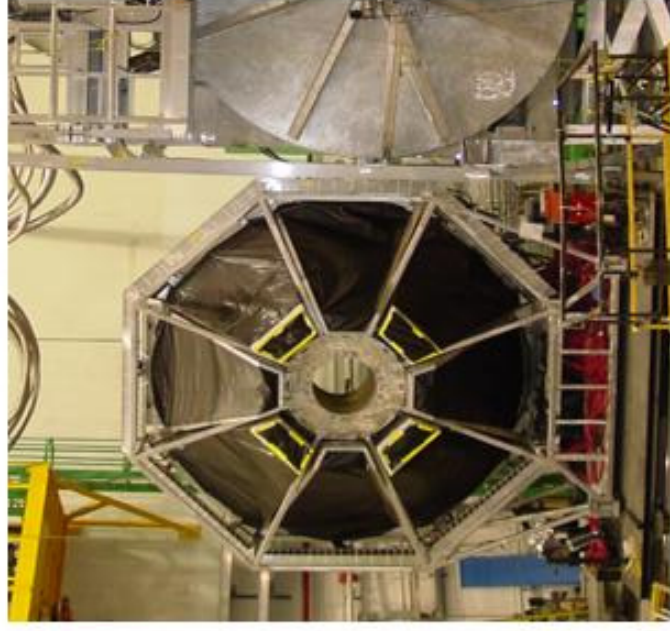
- zero magnification along beam axis
- elastic protons dispersed in Q^2 along focal surface



- acceptance $0.12 < Q^2 < 1.0 \text{ GeV}^2$ for 3 GeV incident beam
- detector 15 acceptance: $0.44 - 0.88 \text{ GeV}^2$
 - 3 Q^2 bins at 0.51, 0.63 and 0.78 GeV^2
- detector 14: $Q^2 = 0.41, 1.0 \text{ GeV}^2$
- det. 16: no elastic acceptance
 - important for measuring backgrounds

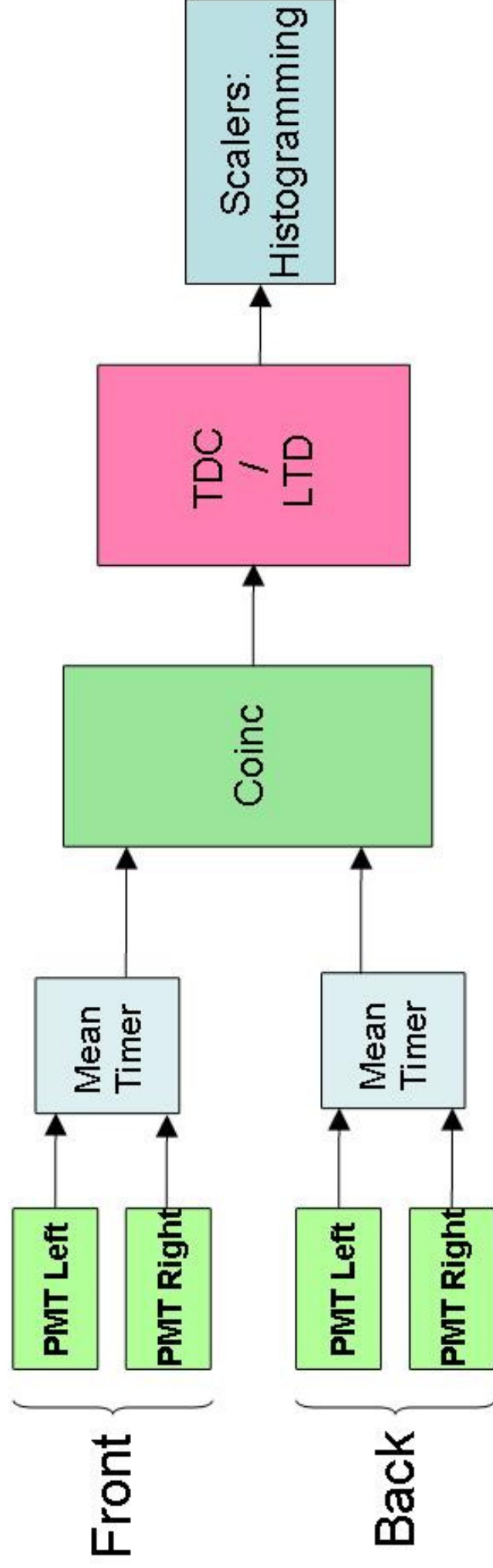
Detectors

- 16 detectors per octant
- Arc shape (const. Q^2), protons at normal incidence
- Each detector: scintillator pair
 - BC408: 0.5, 1.0 cm thick
 - 1/8 in. shielding in-between
- PMT at each end of each scintillator
 - XP2262B (NA), XP2282B (Fr)
- Signal: mean-time-front .AND. mean-time-back
- Assembled with ~ 2 mm accuracy
- Octants in light-tight enclosures



Electronics

- Measure time-of-flight target to detectors
- Counting rates ≤ 4 MHz per scintillator pair
- Fast time encoding
 - NA: dual 500 MHz shift registers \rightarrow scalars (1 ns resolution)
 - “latching time digitizer” (LTD)
 - Fr: flash TDC \rightarrow DSP \rightarrow scalars (1/4 ns resolution)



Electronics Deadtime Corrections

- Residual effect on asymmetry

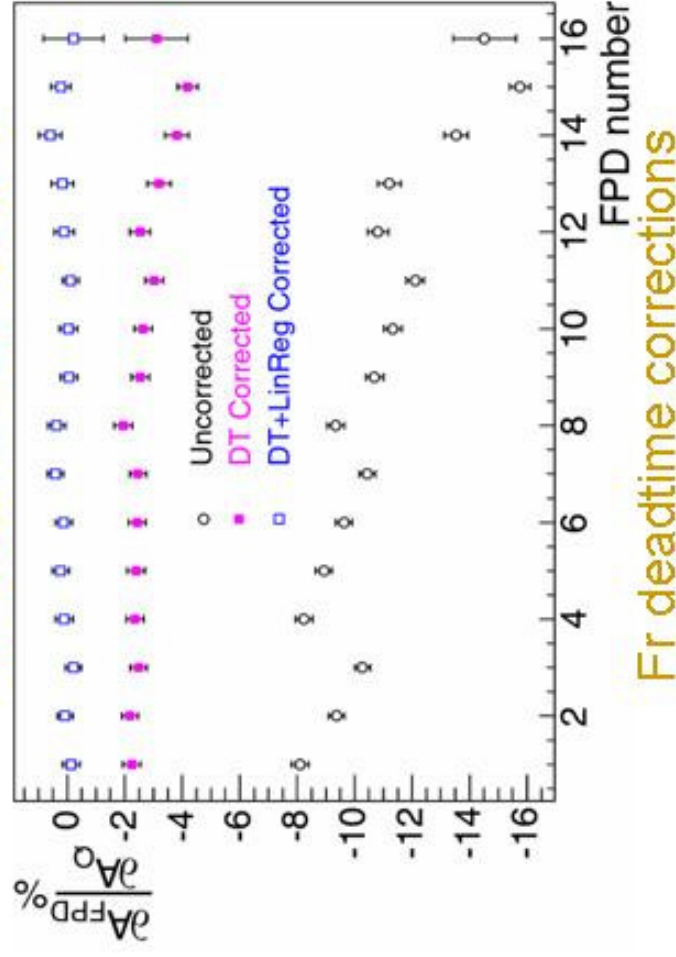
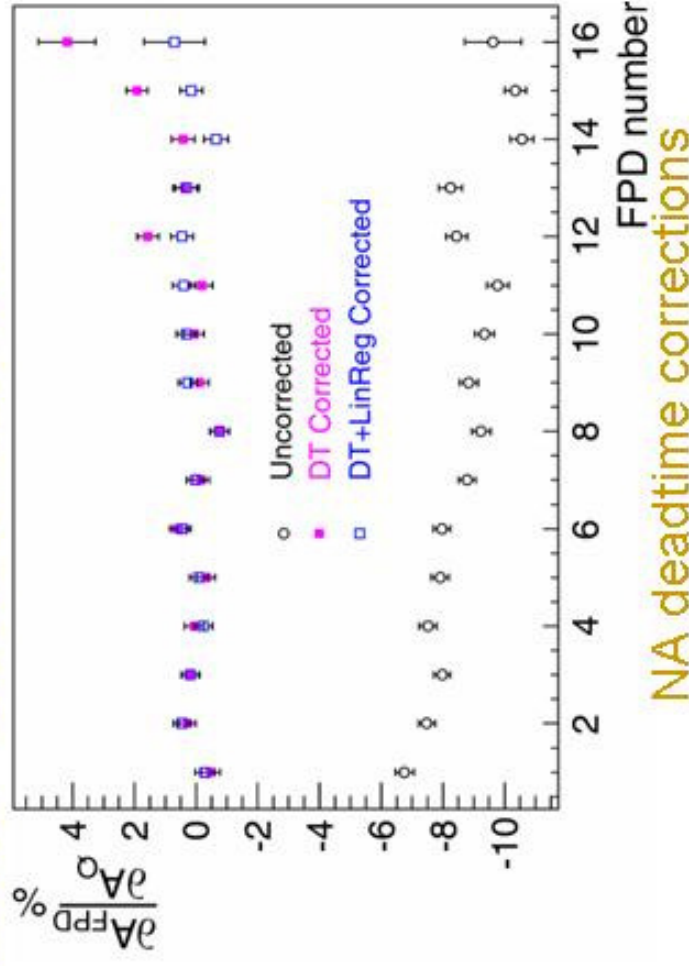
– scale factor

$$A_{meas} = \frac{R_+(1 - \tau R_+) - R_-(1 - \tau R_-)}{R_+(1 - \tau R_+) + R_-(1 - \tau R_-)}$$

$$\simeq A \left(1 - \tau \frac{R_+ + R_-}{2} \right)$$

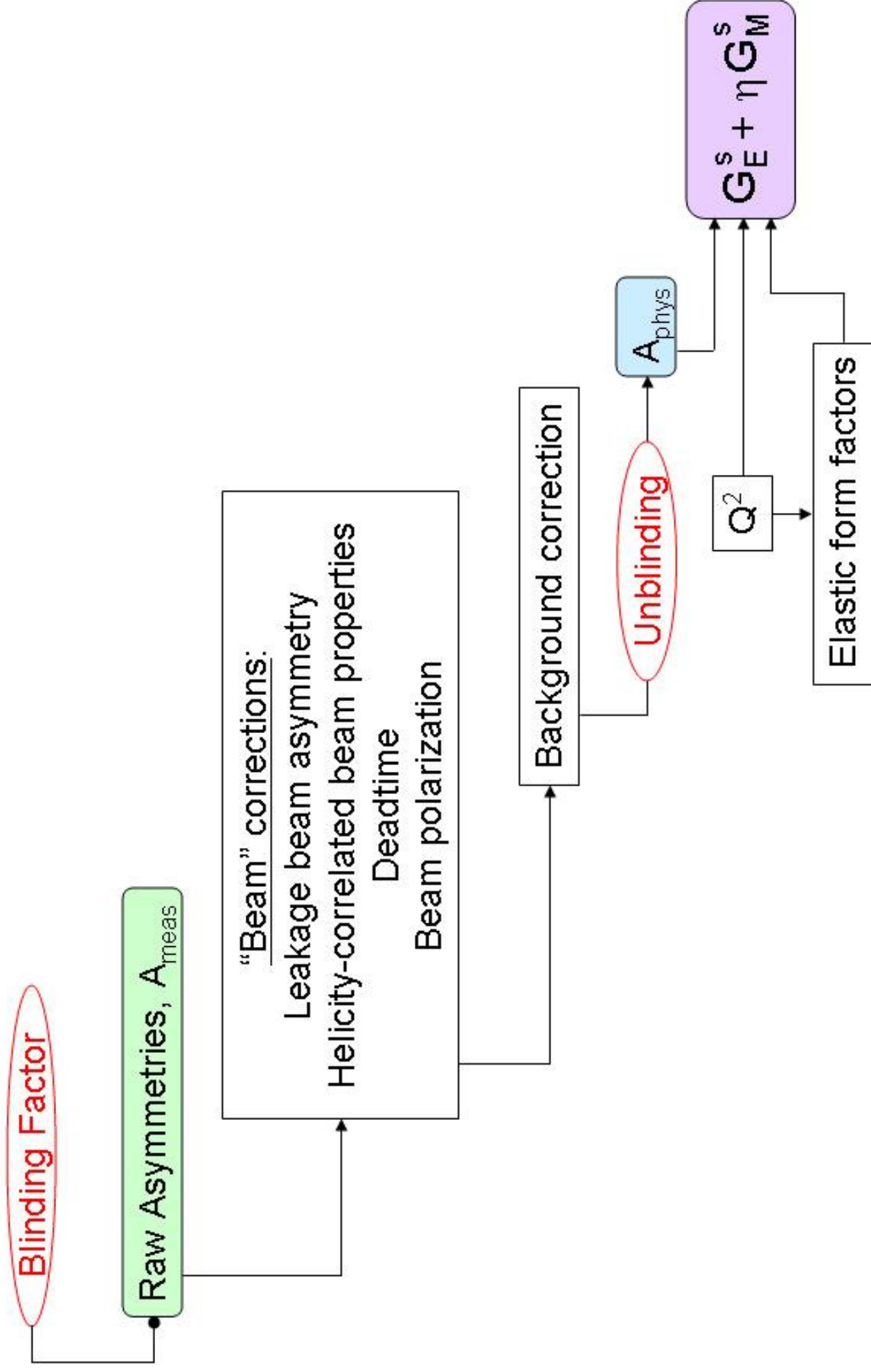
- A is sum of physics and charge asymmetries

- helicity-correlated beam current changes corrected in linear regression analysis
- correction for residual effect $\sim 0.05 \pm 0.05$ ppm (pt-pt systematic unc.)



Analysis

Analysis Overview



Forward Data Summary

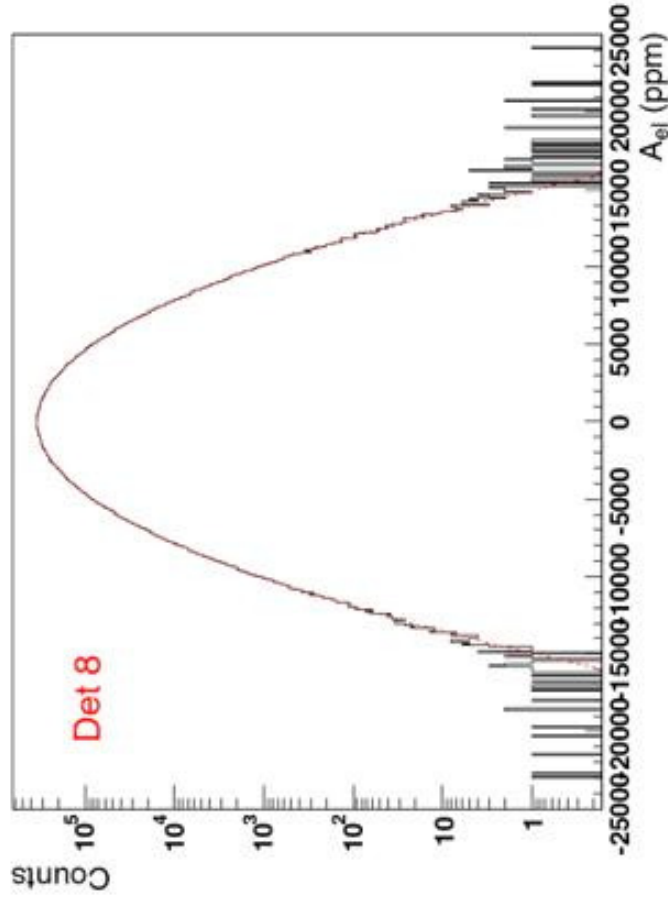
- **101 Coulombs of parity-quality beam**
 - cuts on helicity-correlated beam parameter are 4 x std. dev. for given run:

Quantity	Std. dev.
charge asymmetry	600 ppm
x, y position differences	8, 10 μm
x, y angle difference	0.6, 1.1 μrad
energy difference	7.5 keV

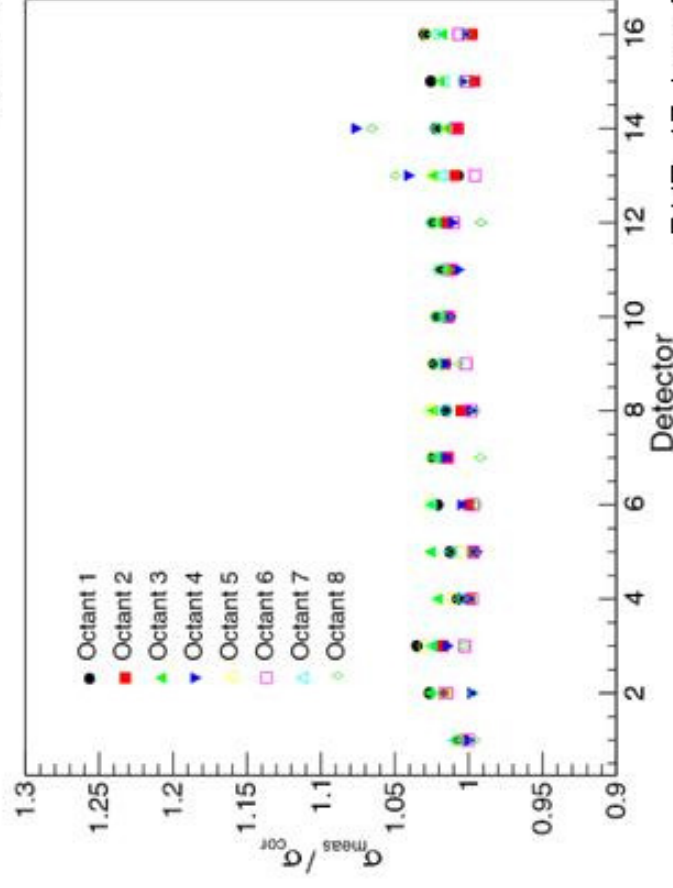
- **Includes running with both Hall A and Hall B (leakage beam asymmetry measured satisfactorily)**
- **Corresponds to: 701 h at 40 μA**
 - 19 x 10⁶ quartets**
 - 76 x 10⁶ MPS**

Statistical Properties of the Data

- Asymmetry distributions very clean over range of 10^5

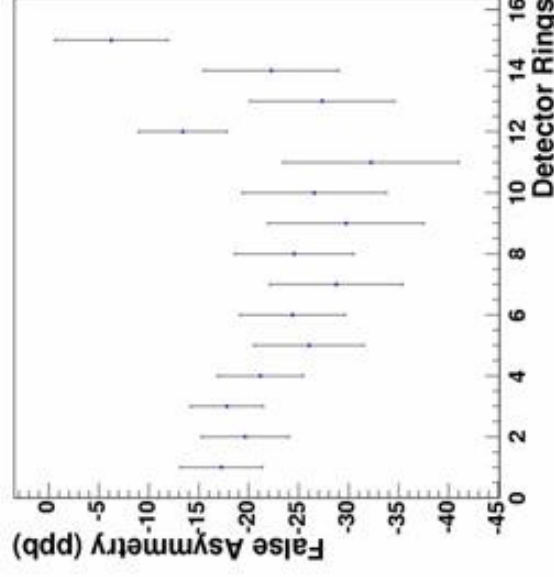
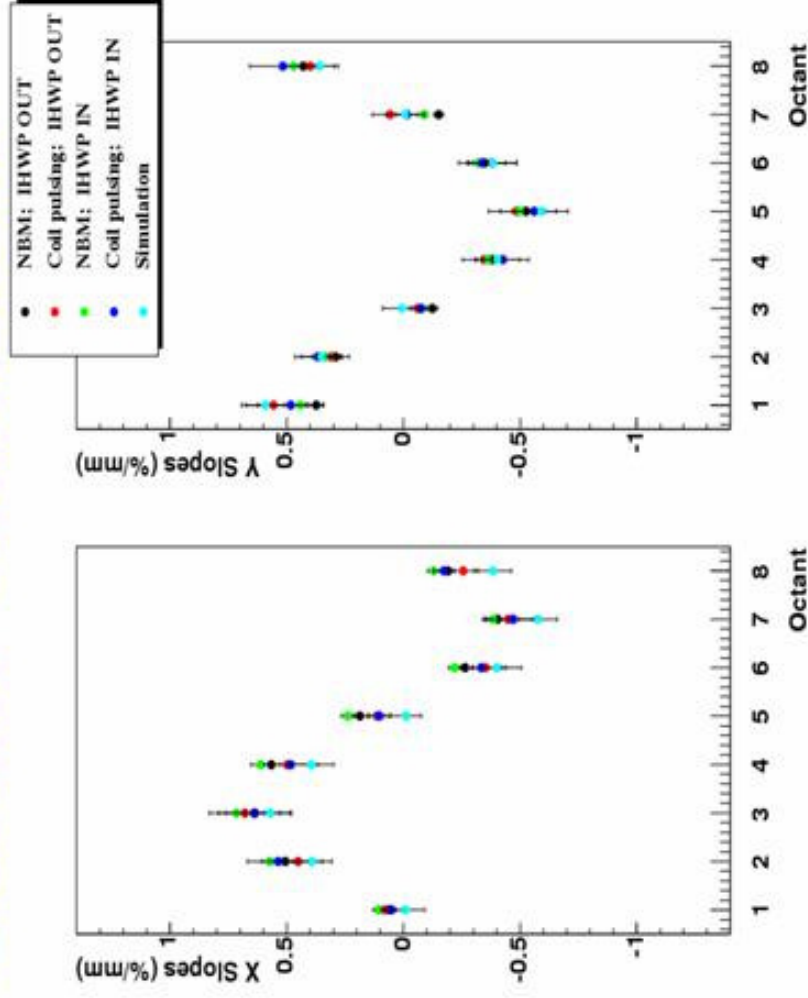


- Measured and expected widths agree at few % level



Helicity-Correlated Beam Parameters

- **Response of spectrometer to beam changes well understood**
- **Average helicity-correlated beam parameters very small**
- **False asymmetries due to helicity-correlated beam parameters very small**
 - overall about -0.02 ppm
 - largest is 0.01 ppm from residual charge asymmetry
 - uncertainties small as well: 0.01 ppm



Background Overview

- Measure yield and asymmetry of entire spectrum
- Correct asymmetry according to

$$A_{meas} = (1 - f)A_{el} + fA_{back}$$

where A_{el} is the raw elastic asymmetry,

$$f = \frac{Y_{back}}{Y_{meas}}$$

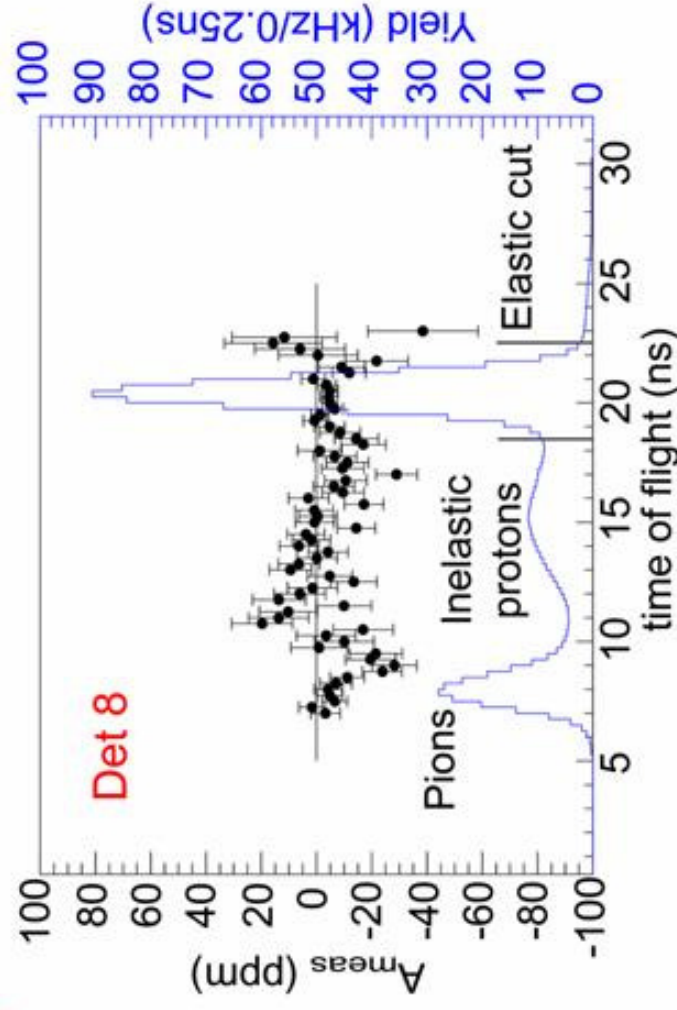
- Actual analysis: $f = f(t)$

– det. 1-14

- fit Y_{back} (polynomial of degree 4), Gaussian for elastic peak
- then fit A_{back} (polynomial of degree 2), constant A_{el}

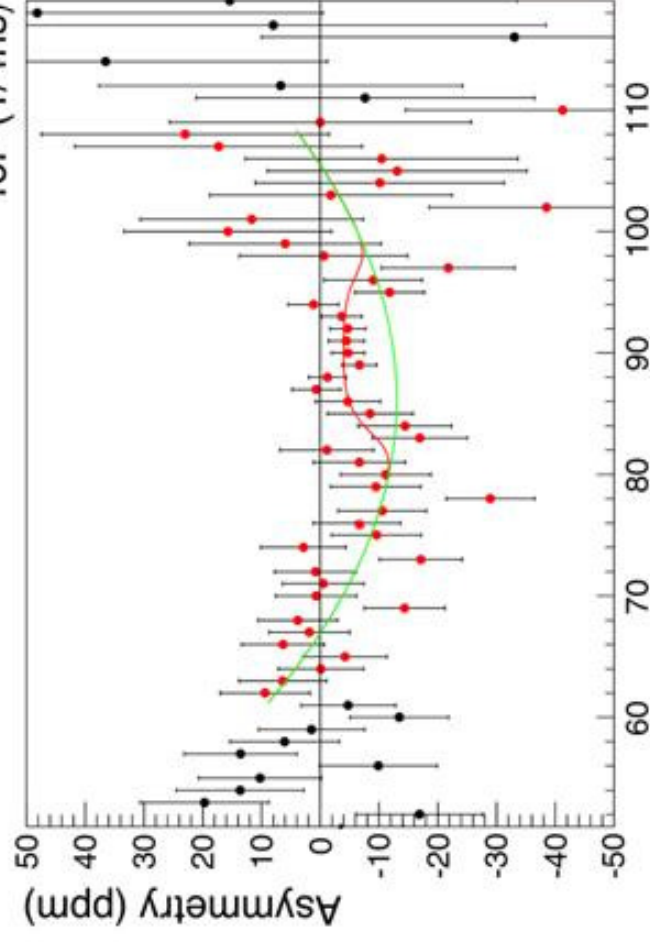
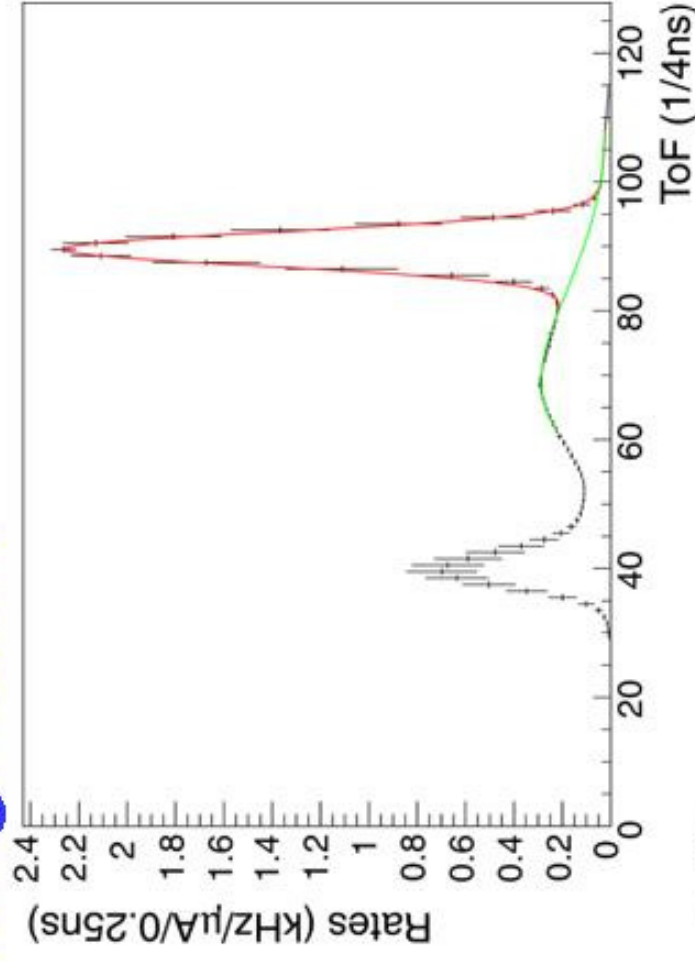
– det. 15

- interpolate over detectors for Y_{back} A_{back}
- fit 3 constants for A_{el}



Det 1-14 Background

- **Results of 2-step fitting procedure: det 8**
 - fit Y_{back} (polyⁿ of degree 4), Gaussian for elastic peak
 - then fit A_{back} (polyⁿ of degree 2), constant A_{el}
 - example fits
 - **yield: $\chi^2 = 31.1/40$**
 - **asym: $\chi^2 = 37.5/44$**
 - f determined from Y_{back} , Y_{meas} in subsequent analysis
 - **don't use detailed shape of elastic peak**
- **Det 14 similar except it has 2 elastic peaks**
 - $Q^2 = 0.41, 1.0 \text{ GeV}^2$



Det. 1-14 Background Uncertainty

- Statistical uncertainty includes that from A_{el} and from A_{back}

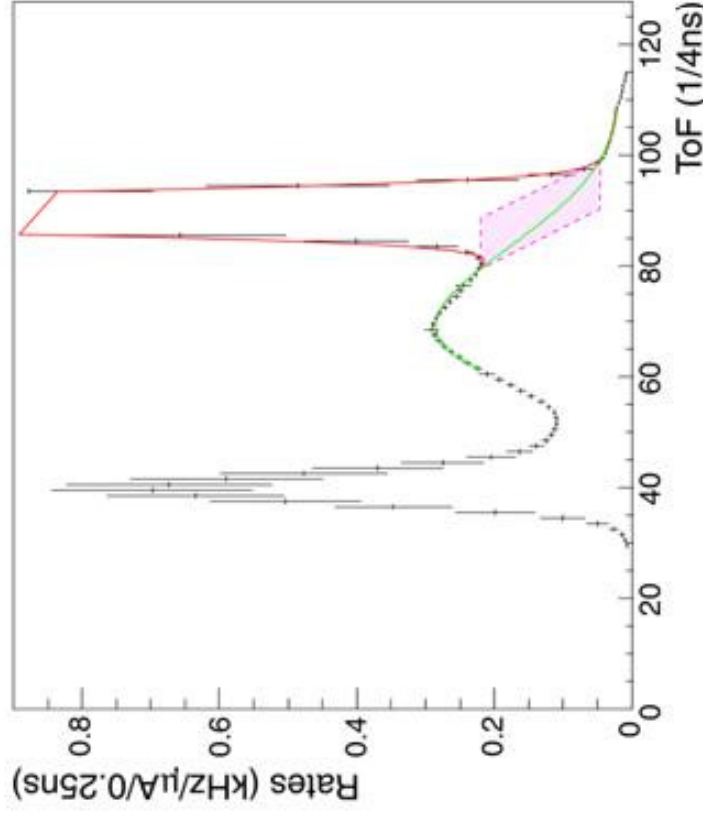
$$A_{meas} = (1-f)A_{el} + fA_{back}$$

- **Systematic uncertainty: general philosophy**
 - vary background yield and asymmetry over plausible ranges
 - consider distributions of results for A_{el}
 - unweighted
 - weighted by χ^2
 - systematic uncertainty is average of std. dev. of these two distributions

Det. 1-14 Background Uncertainty

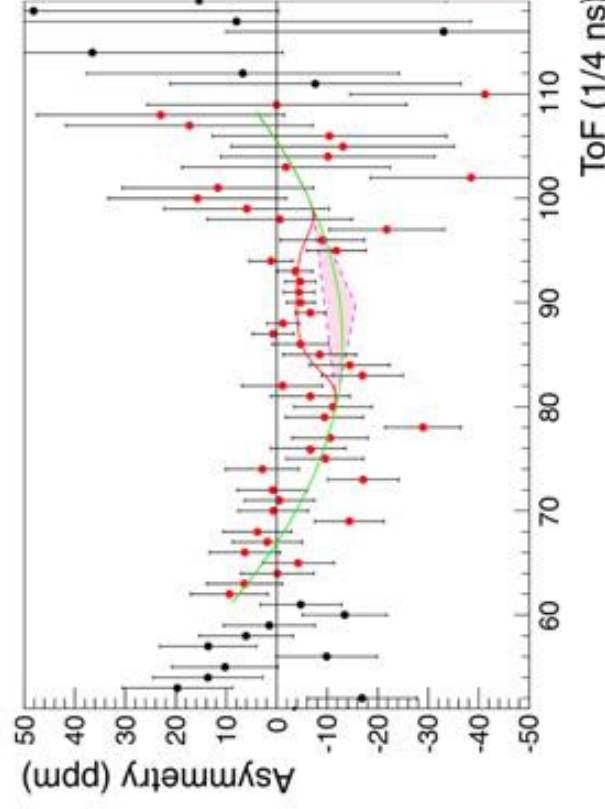
- Background yield varied within “lozenge”

- use a variety of shapes



- Similar approach for asymmetry

- vary throughout range



Correlations in Det 1-14 Backgrounds

- **Separate point-to-point (pt-pt) uncertainties in background correction from global uncertainties**
 - e.g. changing from linear to quadratic model for background asymmetry changes all det. 1 -14 asymmetries downward on average
- **Again using the distributions of results for A_{el}**
 - calculate ~ correlation coefficient
 - correlated uncertainty is change in centroid of distribution for given background model compared to width of overall distribution (L total systematic uncertainty)
- **For det. 1-14**

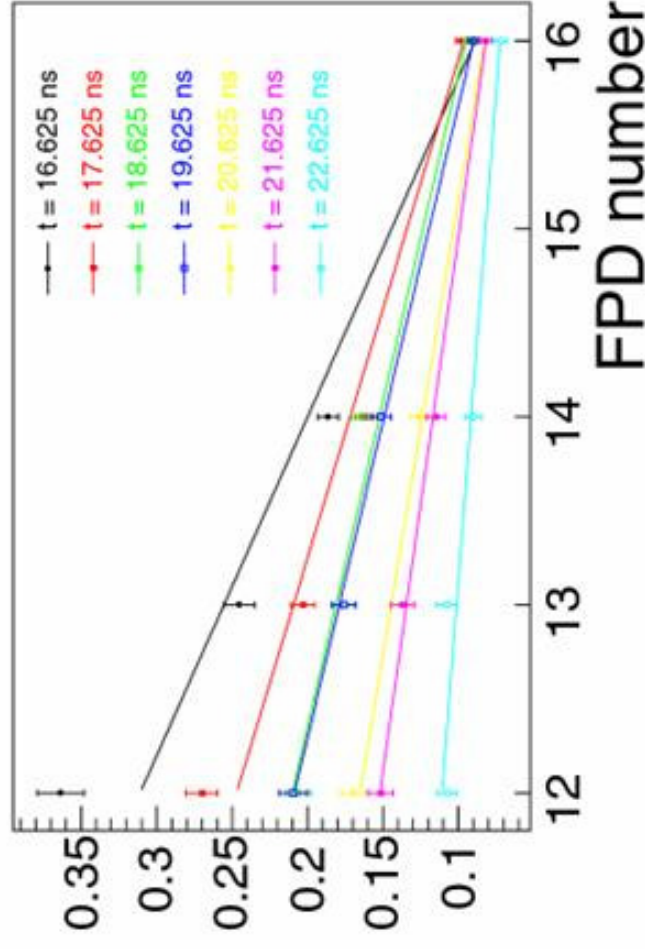
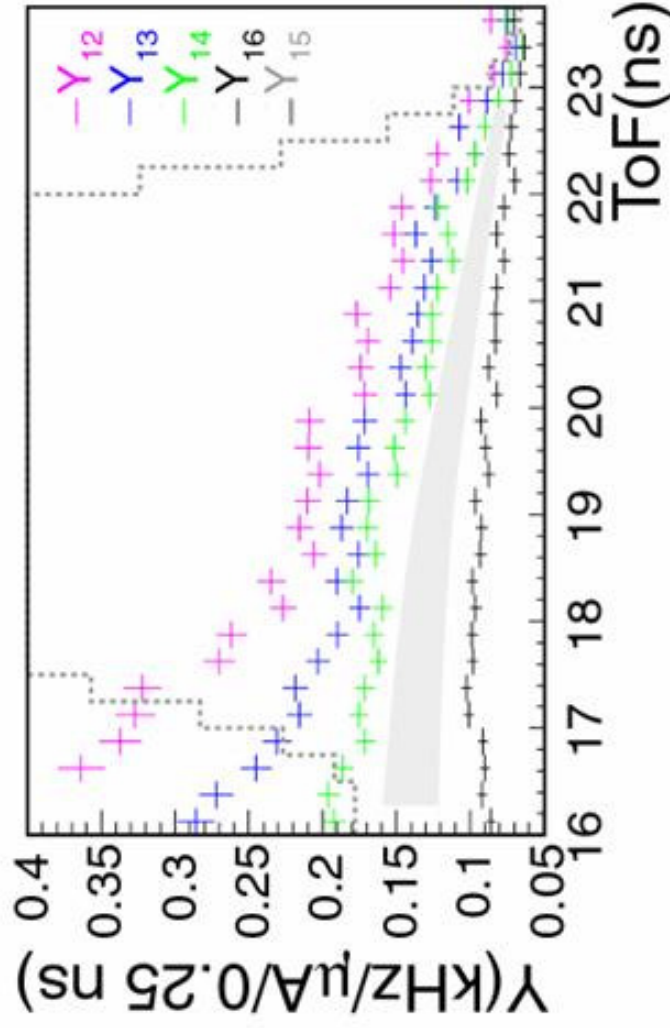
$$\Delta^2 A_{el,sys} = \Delta^2 A_{el,pt-pt} + \Delta^2 A_{el,glob}$$

$$\Delta^2 A_{el,pt-pt} = \frac{1}{4} \Delta^2 A_{el,sys}$$

$$\Delta^2 A_{el,glob} = \frac{3}{4} \Delta^2 A_{el,sys}$$

Det. 15 Background Yields

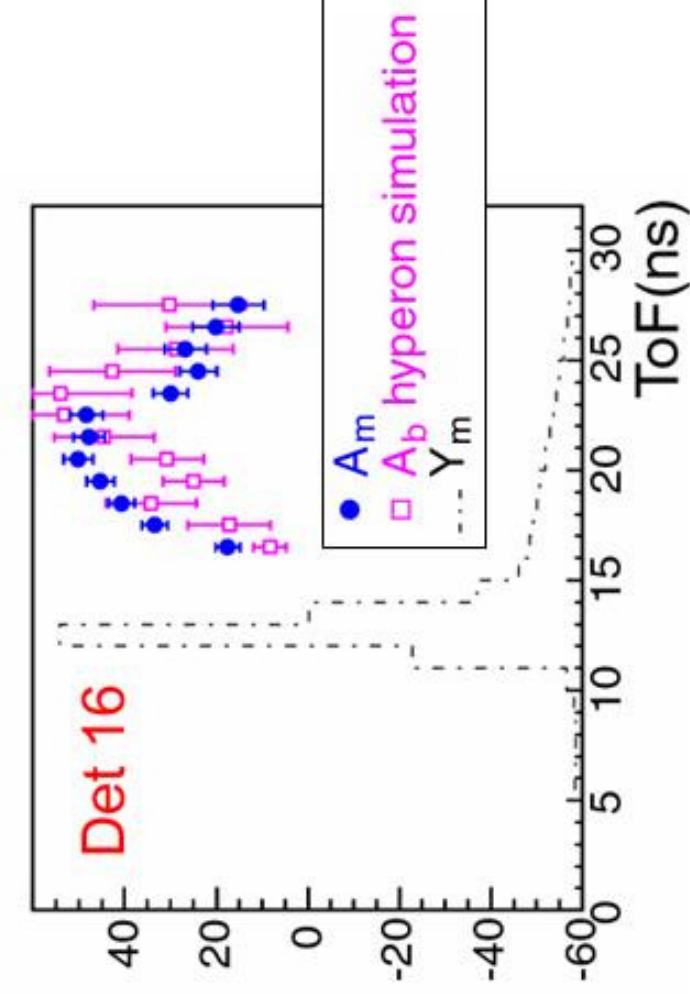
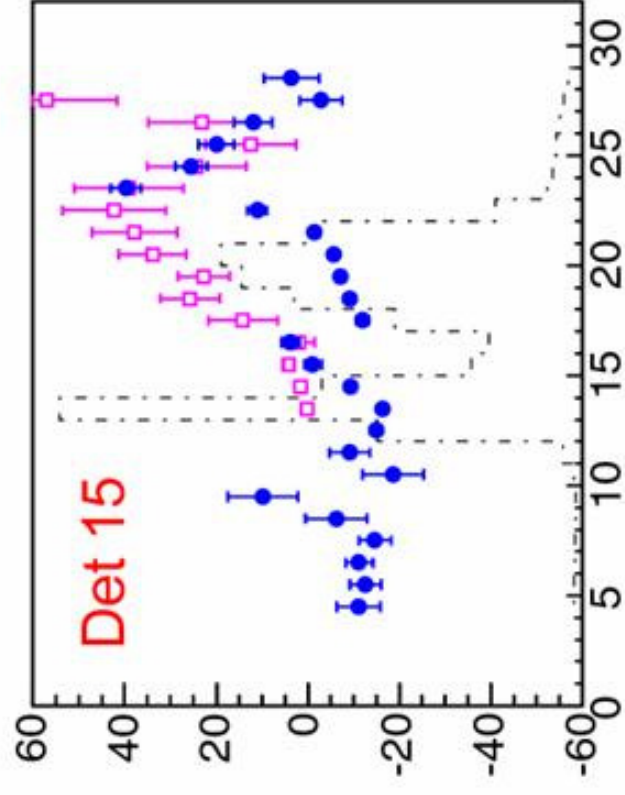
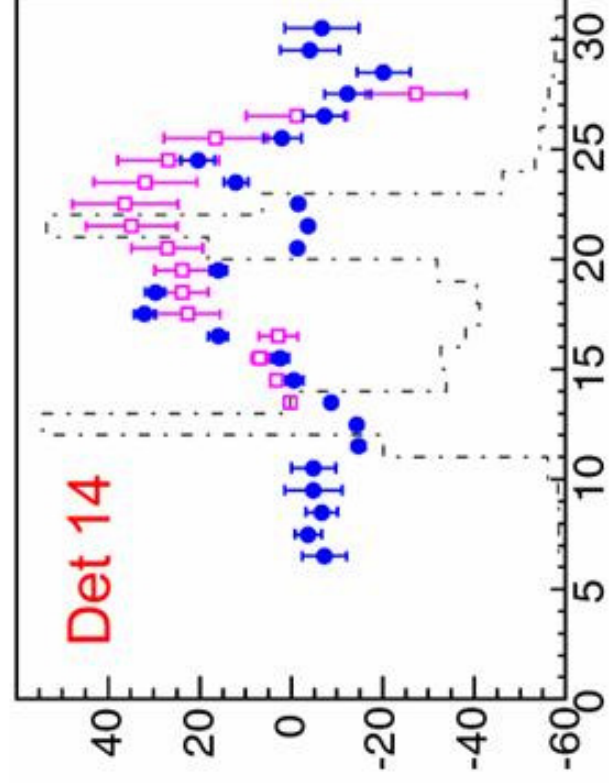
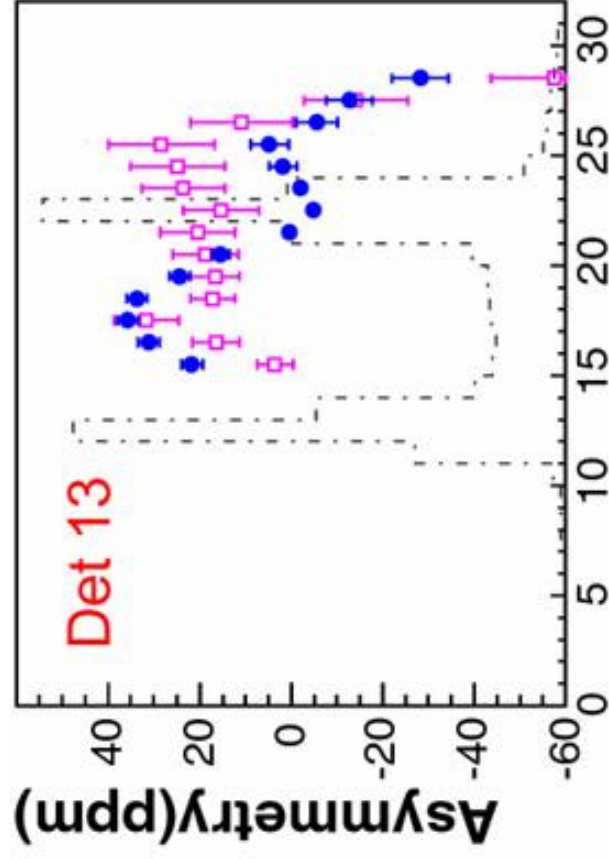
- Elastic protons shifted to lower t.o.f.
- Elastic peak broadened because of increased Q^2 acceptance
- Interpolate over detector range 12-14, 16
 - take out changing acceptance first



Positive Background Asymmetries

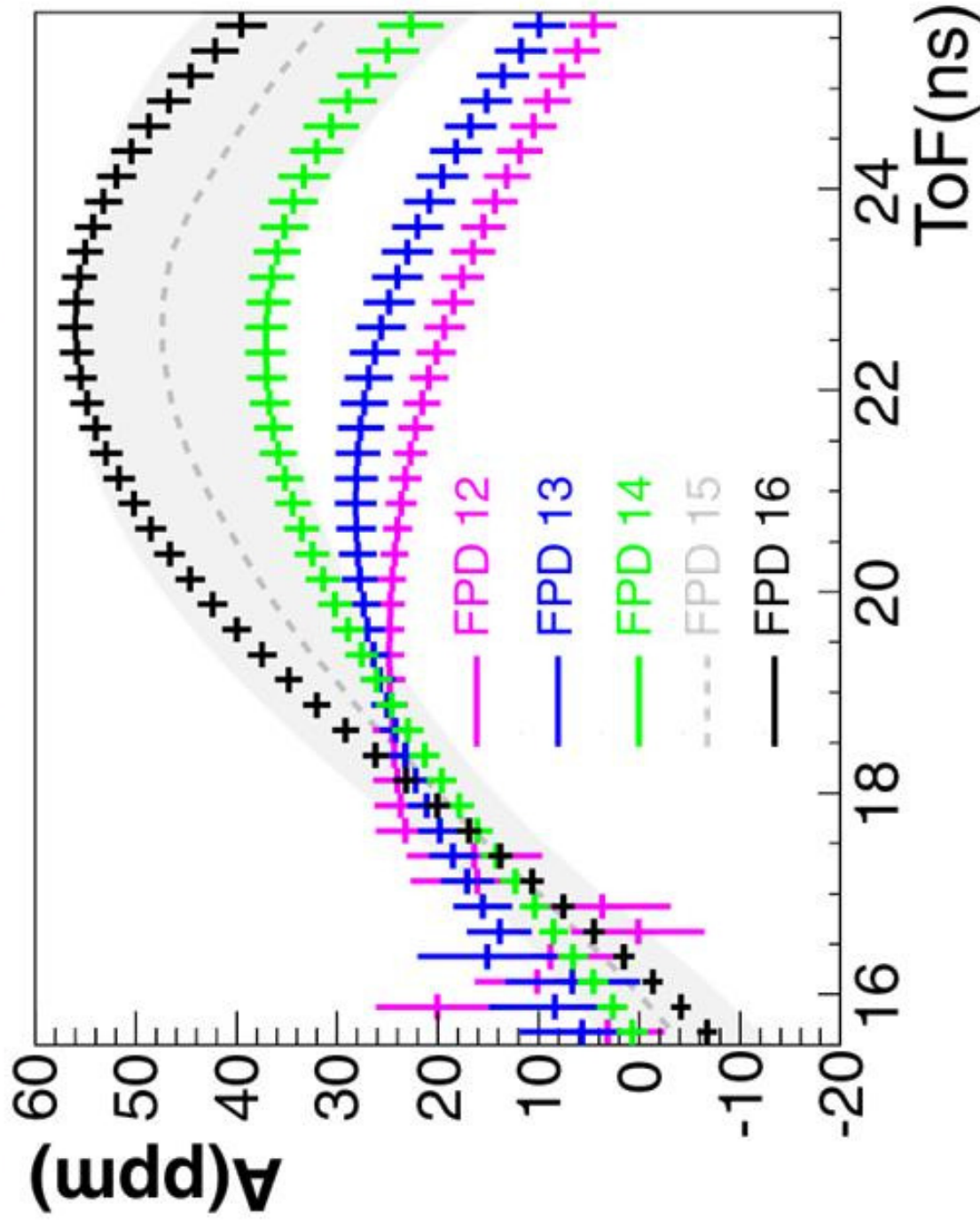
- Det. 12-16 see smoothly varying peak in background asymmetries
 - maximum magnitude $\sim +45$ ppm
- Source is protons from hyperon weak decay scattering inside spectrometer
 - GEANT simulation with generator for hyperon production based on CLAS data
 - simulate both Λ and $\Sigma^{+,0}$ decays
 - polarization transfer for Λ 100%
 - assume 70% for Σ^+
 - Σ^0 asymmetry scaled by further factor of $-1/3$ (CG coefficient)
 - simulation explains source; use measured data for actual analysis

Positive Background Asymmetries: GEANT



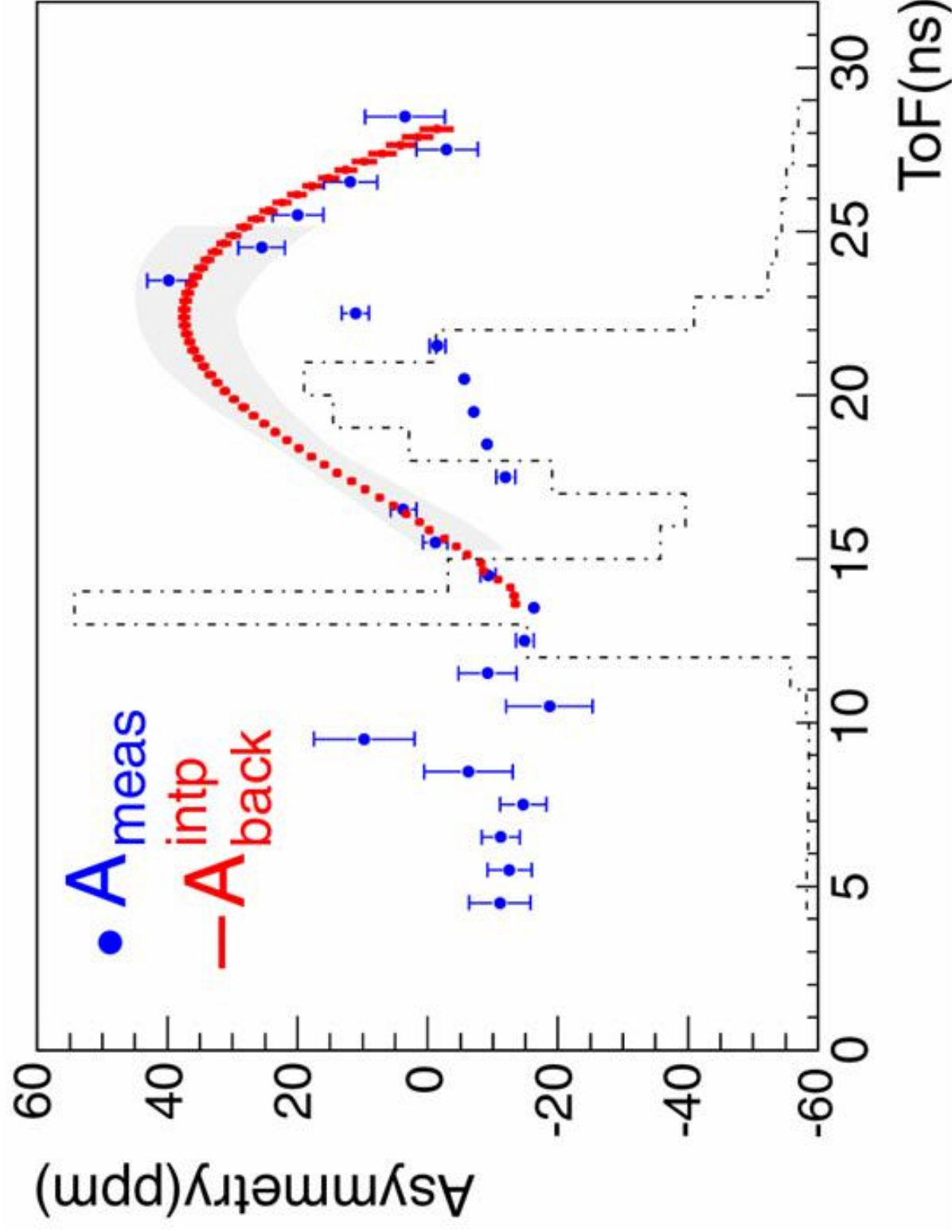
Det. 15 Background Asymmetry

- Use smoothed interpolation of A_{back} from det. 12-14, 16
- Uncertainties are ± 1 detector AND ± 0.5 ns time shift



Det. 15 Asymmetry

- Compare interpolated background asymmetry and data



Correlations in Det. 15 Backgrounds

- Separate point-to-point (pt-pt) uncertainties in background correction from global uncertainties
 - in det. 15, correlations larger because bins are contiguous
- Consider distributions of results for A_{el}
 - for variety of randomly generated models determine correlation coefficient
- For det. 15

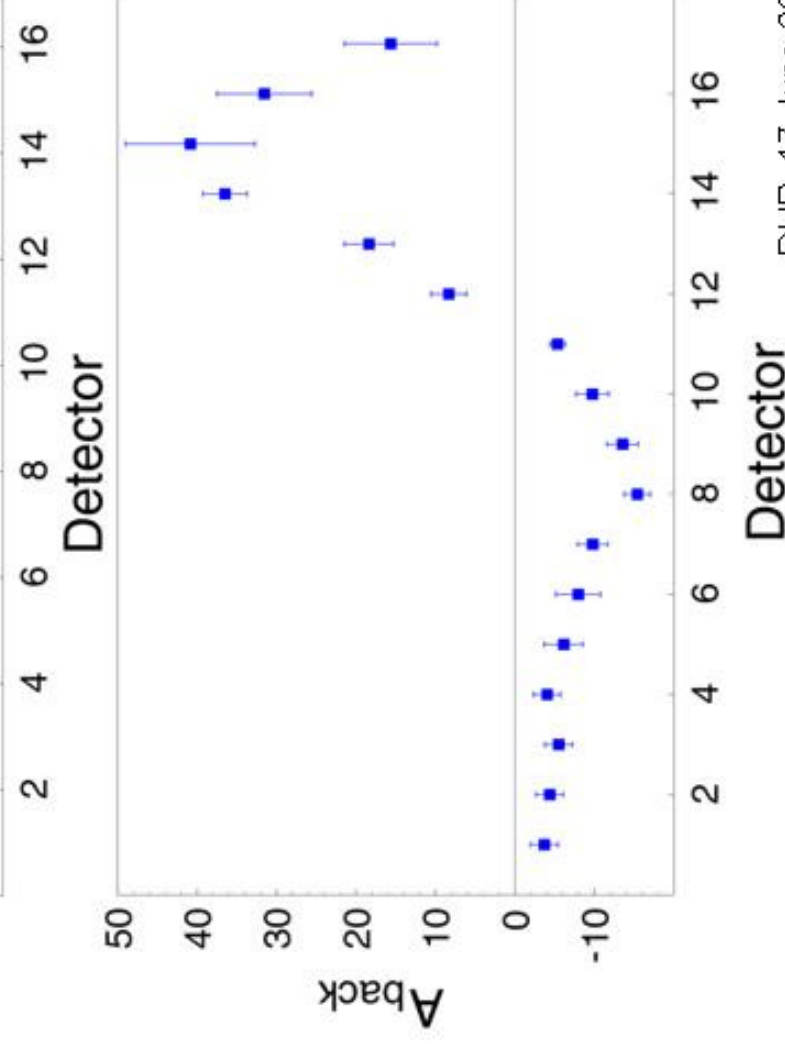
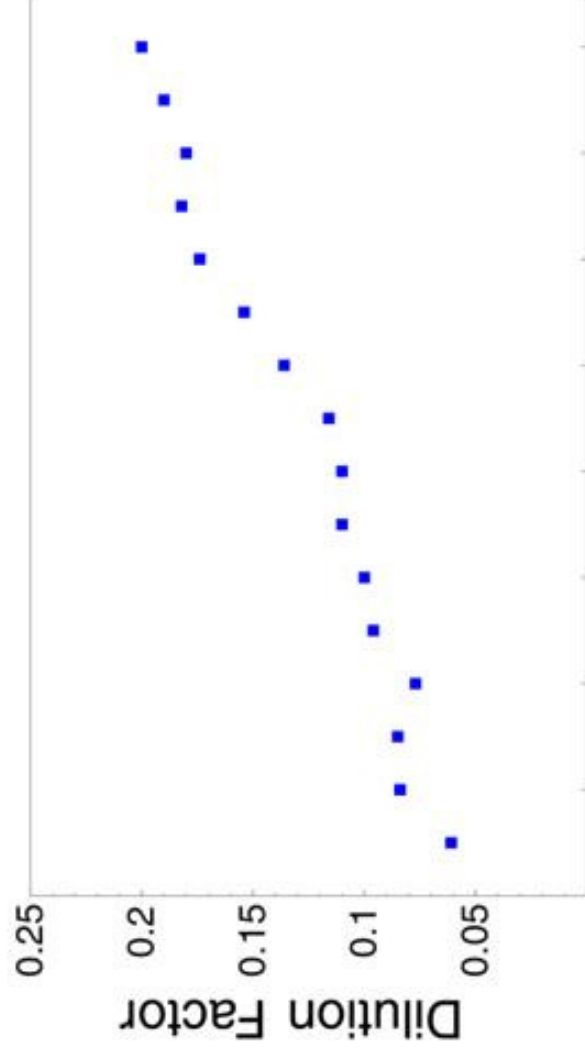
$$\Delta^2 A_{el,sys} = \Delta^2 A_{el,pt-pt} + \Delta^2 A_{el,glob}$$

$$\Delta^2 A_{el,pt-pt} = \frac{1}{2} \Delta^2 A_{el,sys}$$

$$\Delta^2 A_{el,glob} = \frac{1}{2} \Delta^2 A_{el,sys}$$

Dilution factor and Background Asymmetry

- **Smooth, systematic progression**
 - dilution factor
 - background asymmetry
 - both averaged over t.o.f. for demonstration



GO results

Experimental Results

- A_{phys} corrected for all beam, electronics, background factors

Det	Q^2 (GeV^2)	A_{phys} (ppm)	ΔA_{stat} (ppm)	$\Delta A_{\text{sys,pt}}$ (ppm)	$\Delta A_{\text{sys,glob}}$ (ppm)	f (ppm)	ΔA_{meas}
1	0.122	-1.513	0.436	0.224	0.176	0.061	-1.380
2	0.128	-0.972	0.409	0.198	0.173	0.084	-1.070
3	0.136	-1.298	0.424	0.174	0.170	0.085	-1.340
4	0.144	-2.707	0.433	0.183	0.176	0.077	-2.670
5	0.153	-2.223	0.431	0.284	0.214	0.096	-2.460
6	0.164	-2.880	0.434	0.324	0.234	0.100	-3.130
7	0.177	-3.949	0.426	0.251	0.205	0.110	-4.470
8	0.192	-3.850	0.485	0.218	0.192	0.110	-5.010
9	0.210	-4.683	0.475	0.258	0.212	0.116	-5.730
10	0.232	-5.267	0.505	0.301	0.232	0.136	-6.080
11	0.262	-5.260	0.520	0.108	0.166	0.154	-5.550
12	0.299	-7.715	0.602	0.531	0.349	0.174	-5.400
13	0.344	-8.400	0.676	0.850	0.521	0.182	-3.650
14 a	0.410	-10.25	0.674	0.895	0.551	0.180	-1.700
15 a	0.511	-16.81	0.889	1.478	1.498	0.190	-5.800
15 b	0.631	-19.96	1.112	1.277	1.306	0.200	-9.740
15 c	0.788	-30.83	1.857	2.556	2.589	0.400	-12.660
14 b	0.997	-37.93	7.237	9.000	0.519	0.780	4.210

Asymmetry with EW Radiative Corrections

- Full form of asymmetry used to extract $G_E^S + \eta G_M^S$

$$A = -\frac{G_F Q^2}{4\pi\alpha\sqrt{2}} \frac{1}{\varepsilon G_E^P + \tau G_M^P} \left\{ (1 - 4\sin^2\theta_W) (\varepsilon G_E^P + \tau G_M^P) (1 + R_V^P) - (\varepsilon G_E^P G_E^N + \tau G_M^P G_M^N) (1 + R_V^N) - (\varepsilon G_E^P G_E^S + \tau G_M^P G_M^S) (1 + R_V^{(0)}) - \varepsilon' (1 - 4\sin^2\theta_W) G_M^P G_A^e \right\}$$

where

$$G_A^e = -G_A^P (1 + R_A^{T=1}) + \left[\frac{1}{2} (3F - D) R_A^{T=0} + \Delta s (1 + R_A^{(0)}) \right] G_A^{dip}$$

and

$$G_A^P = g_A G_A^{dip} = (F + D) G_A^{dip} = \frac{g_A}{(1 + Q^2/\Lambda_A^2)^2}$$

Standard Parameters

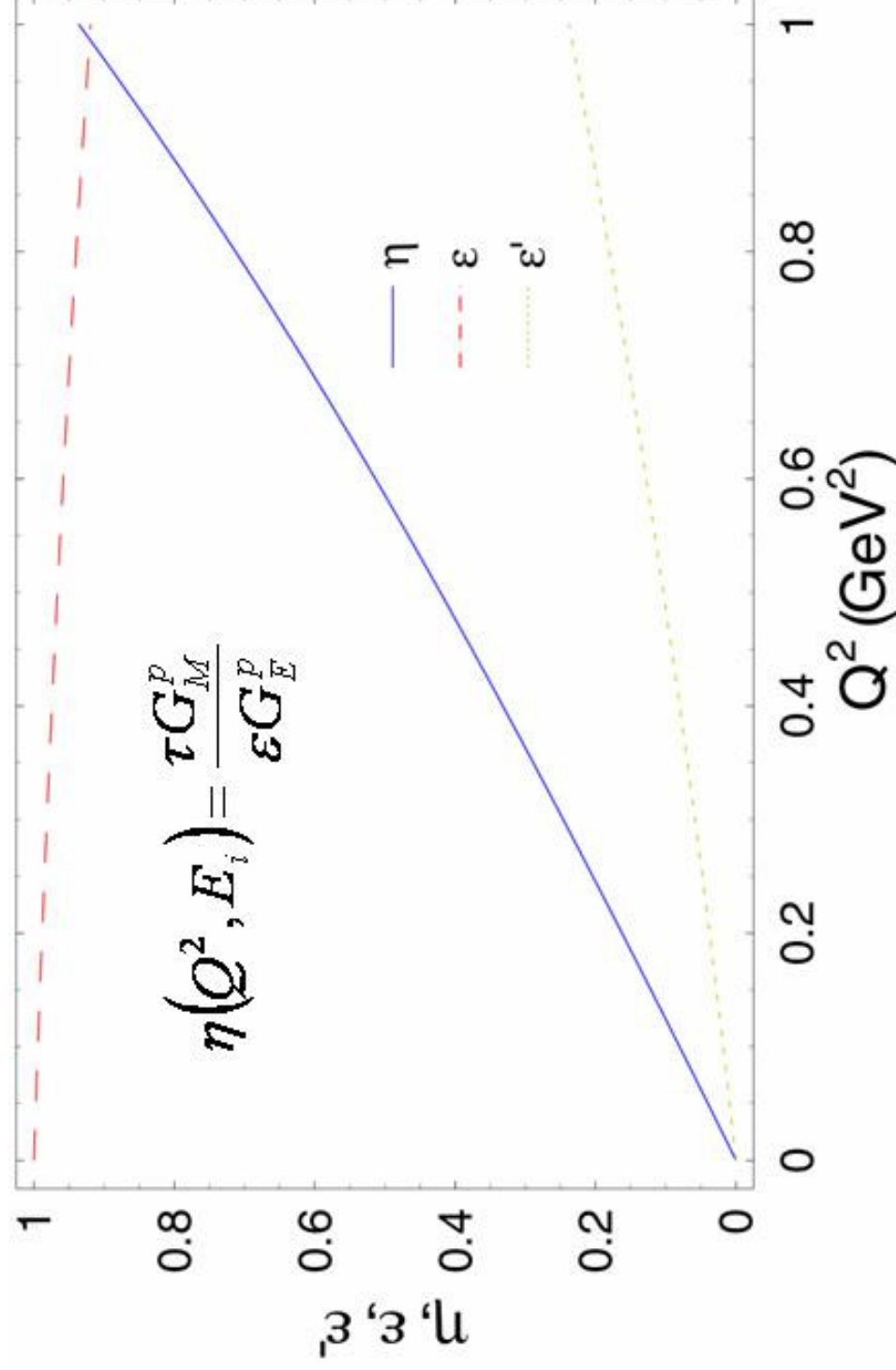
- em form factors: Kelly PRC 70 (2004) 068202
- use Zhu, et al. $R_A^{T=1}$, $R_A^{T=0}$ including anapole
- incident energy 3.028 GeV

α	1/137.03599976	R_V^p	-0.0447091
$\sin^2\theta_W$	0.2312	R_V^0	-0.011789
G_F	0.0000116639 GeV ²	$R_V^{(0)}$	-0.011789
m_p	0.938272 GeV	$R_A^{T=1}$	-0.259163
μ_p	2.79285	$R_A^{T=0}$	-0.23826
μ_n	-1.91304	$R_A^{(0)}$	-0.551753
Λ^2	0.711 GeV ²		
Λ_A^2	1.00 GeV ²		
g_A/g_V	1.2695		
3F-D	0.585		

Strange Quark Contribution

- Strange quark contribution to asymmetry

$$G_E^S + \eta G_M^S = \frac{4\pi\alpha\sqrt{2}}{G_F Q^2} \frac{\varepsilon G_E^{p^2} + \tau G_M^{p^2}}{\varepsilon G_E^p (1+R_V^{(0)})} (A_{phys} - A_{NVS})$$



A_{NVS} Sensitivities

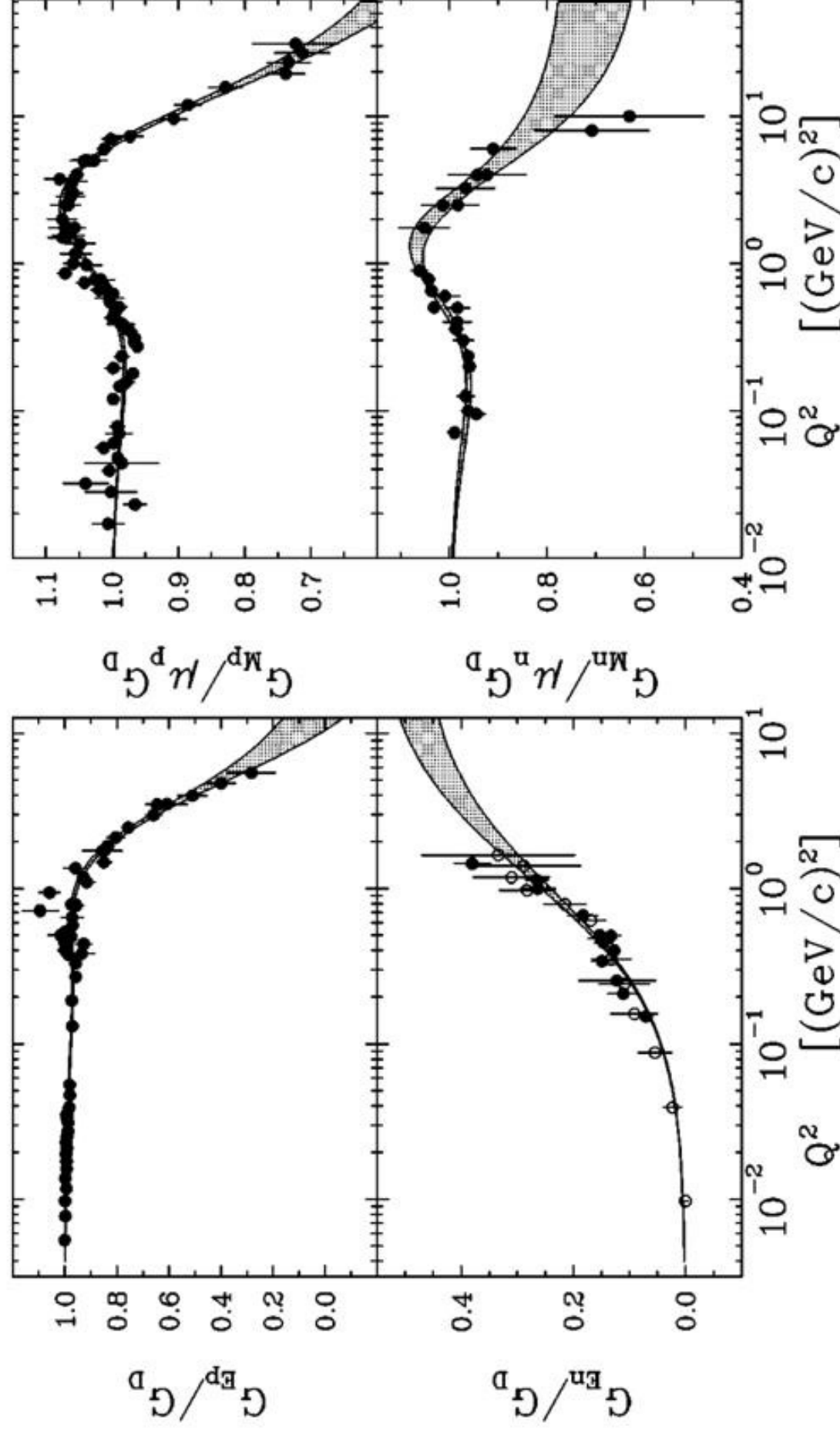
- Base calculation uses Kelly electromagnetic form factors and uncertainties
 - PRC **70** (2004) 068202
 - Friedrich & Walcher, Eur. Phys. J. A **17** (2003) 607.
 - Arrington, Phys. Rev. C **69** (2004) 022201.
 - “Rosenbluth” fit for G_E^p , G_M^p
 - Kelly for G_E^n , G_M^n

- Key factors, uncertainties

$\Delta G_A^p / G_A^p$	2.2%
$R_A^{T=0}$	-0.24 \pm 0.20
$R_A^{T=1}$	-0.26 \pm 0.35
$R_A^{(0)}$	-0.55 \pm 0.55
3F - D	0.585 \pm 0.025

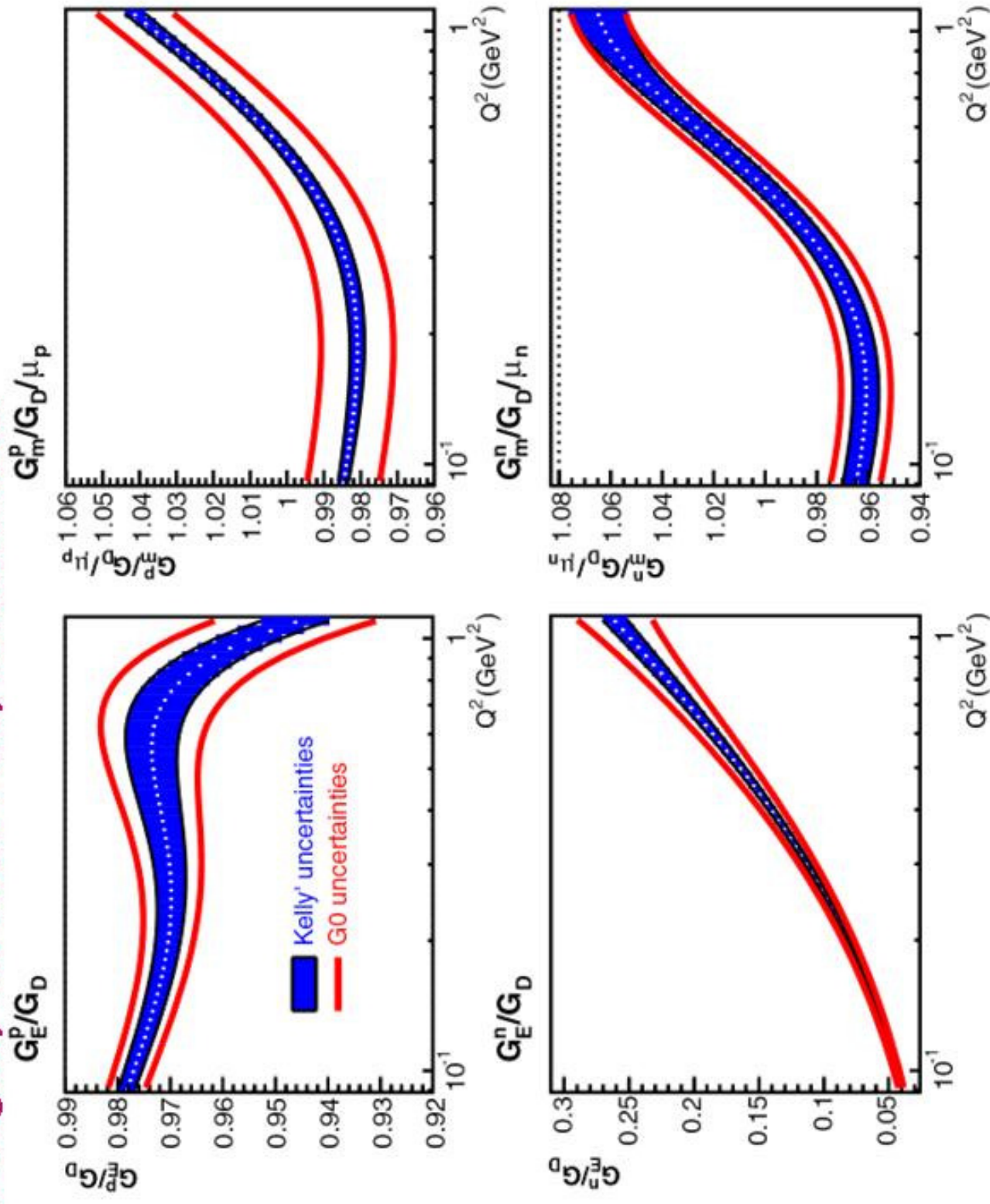
Electromagnetic Form Factors

- Use Kelly parameterization for baseline results
 - omits Rosenbluth data for G_E^p , $Q^2 > 1 \text{ GeV}^2$
 - omits some neutron data using associate particle calibration



Electromagnetic Form Factor Uncertainties

- Use slightly modified Kelly uncertainties



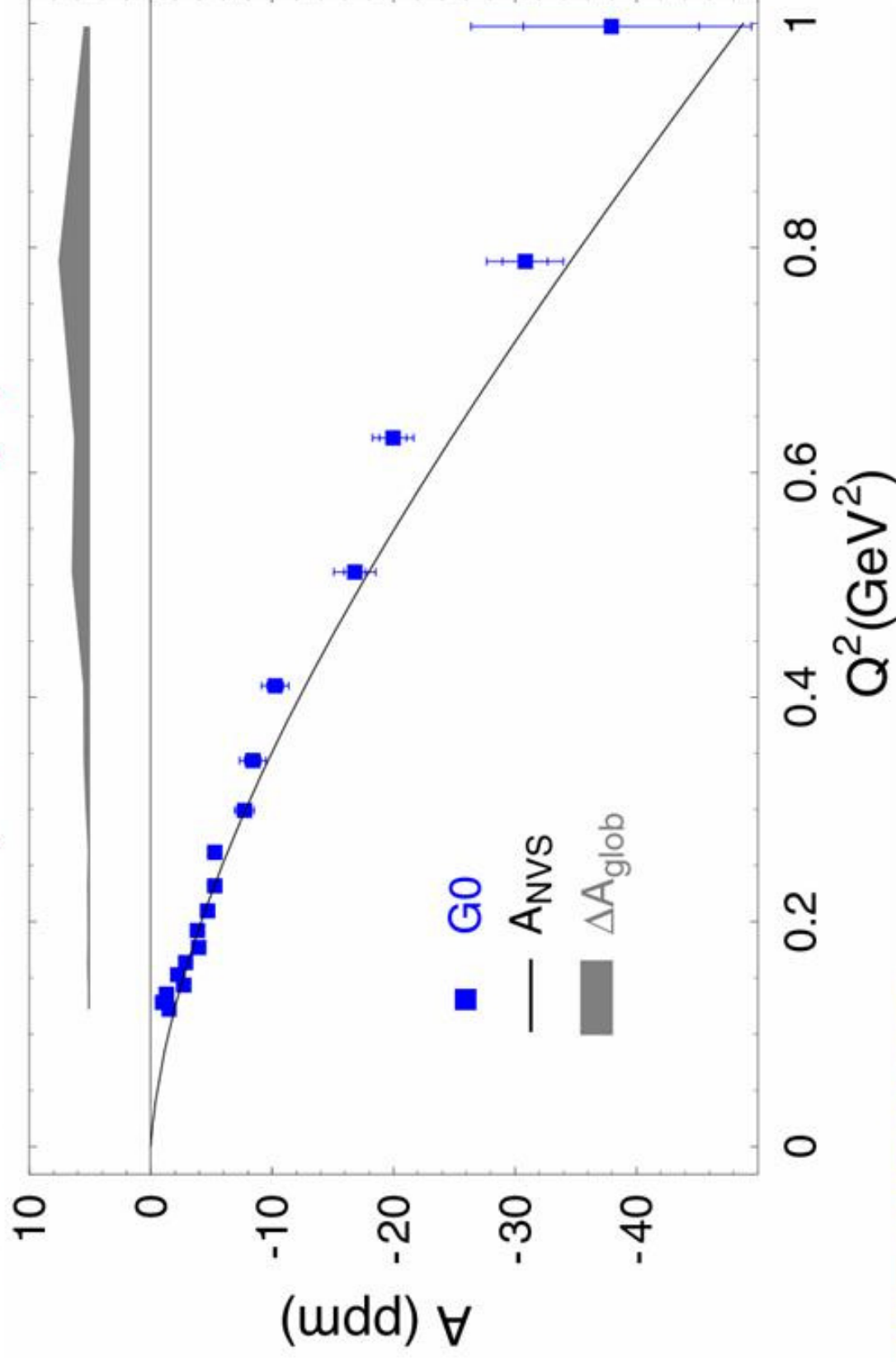
G0 Sensitivities for G_E^s

$$\frac{\partial G_E^s}{\partial p} :$$

	Det 1	Det 8	Det 14a
$A_{\text{meas}} (1/\text{ppm})$	8.344×10^{-2}	5.068×10^{-2}	2.113×10^{-2}
$E_i (1/\text{GeV})$	-1.911×10^{-3}	-3.021×10^{-3}	-6.755×10^{-3}
$Q^2 (1/\text{GeV}^2)$	1.680	1.575	8.718×10^{-1}
$G_E^{\gamma p}$	-2.396×10^{-1}	-3.909×10^{-1}	-6.987×10^{-1}
$G_M^{\gamma p}$	5.488×10^{-2}	6.063×10^{-2}	1.010×10^{-1}
$G_E^{\gamma n}$	-1.000	-1.000	-1.000
$G_M^{\gamma n}$	-9.839×10^{-2}	-1.559×10^{-1}	-3.407×10^{-1}
$G_A^{\gamma p}$	3.992×10^{-3}	6.471×10^{-3}	1.516×10^{-2}
G_A^s	-1.721×10^{-3}	-2.472×10^{-3}	-4.140×10^{-3}
R_V^p	6.896×10^{-2}	6.589×10^{-2}	5.861×10^{-2}
R_V^n	9.440×10^{-2}	1.307×10^{-1}	2.060×10^{-1}
$R_V^{(0)}$	-3.739×10^{-2}	-2.846×10^{-3}	-5.320×10^{-2}
R_A^0	-1.123×10^{-3}	-1.613×10^{-3}	-2.702×10^{-3}
R_A^1	4.875×10^{-3}	7.002×10^{-3}	1.173×10^{-2}
$R_A^{(0)}$	3.226×10^{-4}	4.633×10^{-4}	7.759×10^{-4}
3F-D	4.575×10^{-4}	6.571×10^{-4}	1.100×10^{-3}

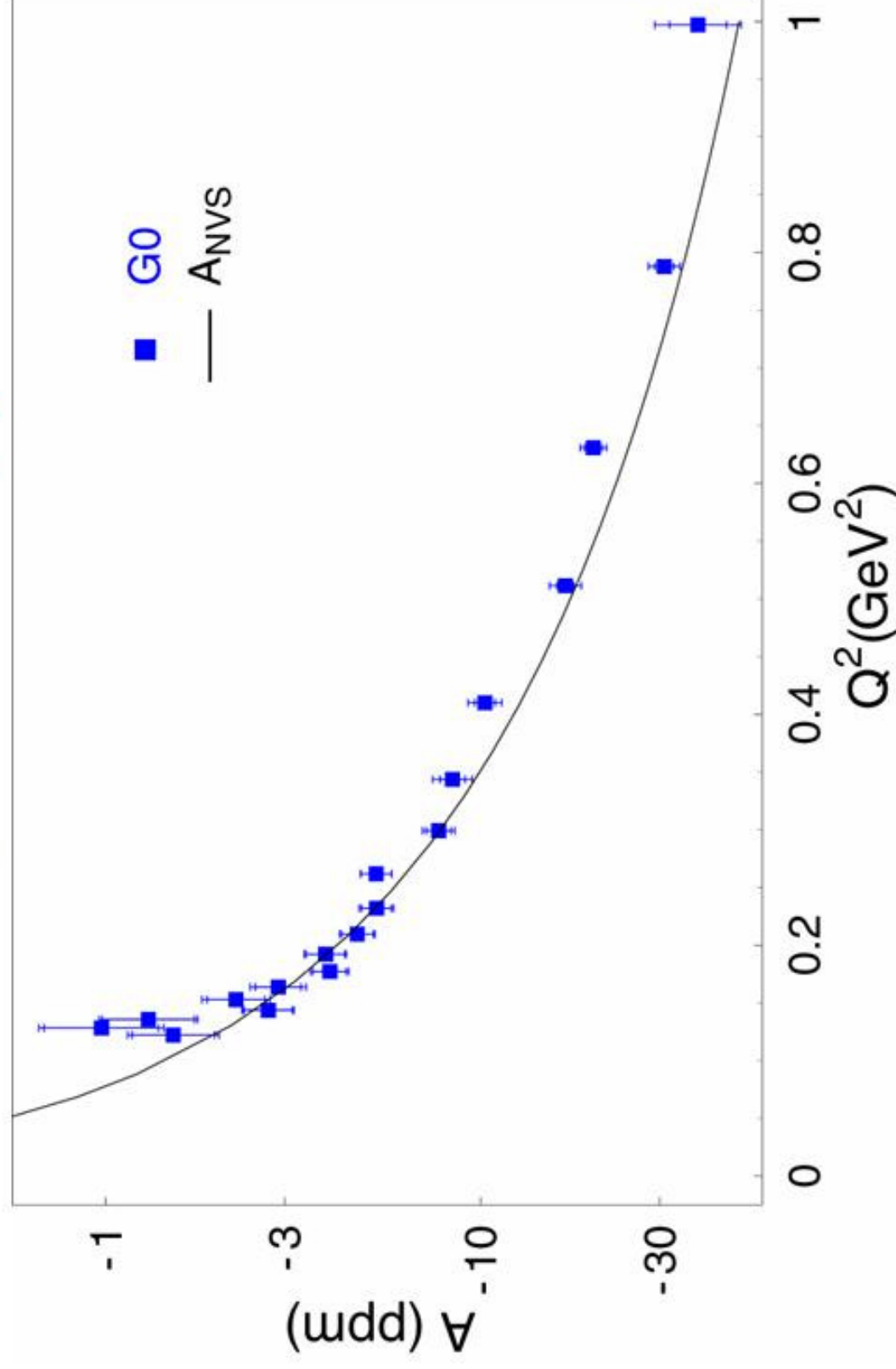
Experimental Asymmetries

- em form factors: Kelly PRC 70 (2004) 068202
- “no vector strange” asymmetry, A_{NVS} , is $A(G_E^s, G_M^s = 0)$
- inside error bars: stat, outside: stat & pt-pt

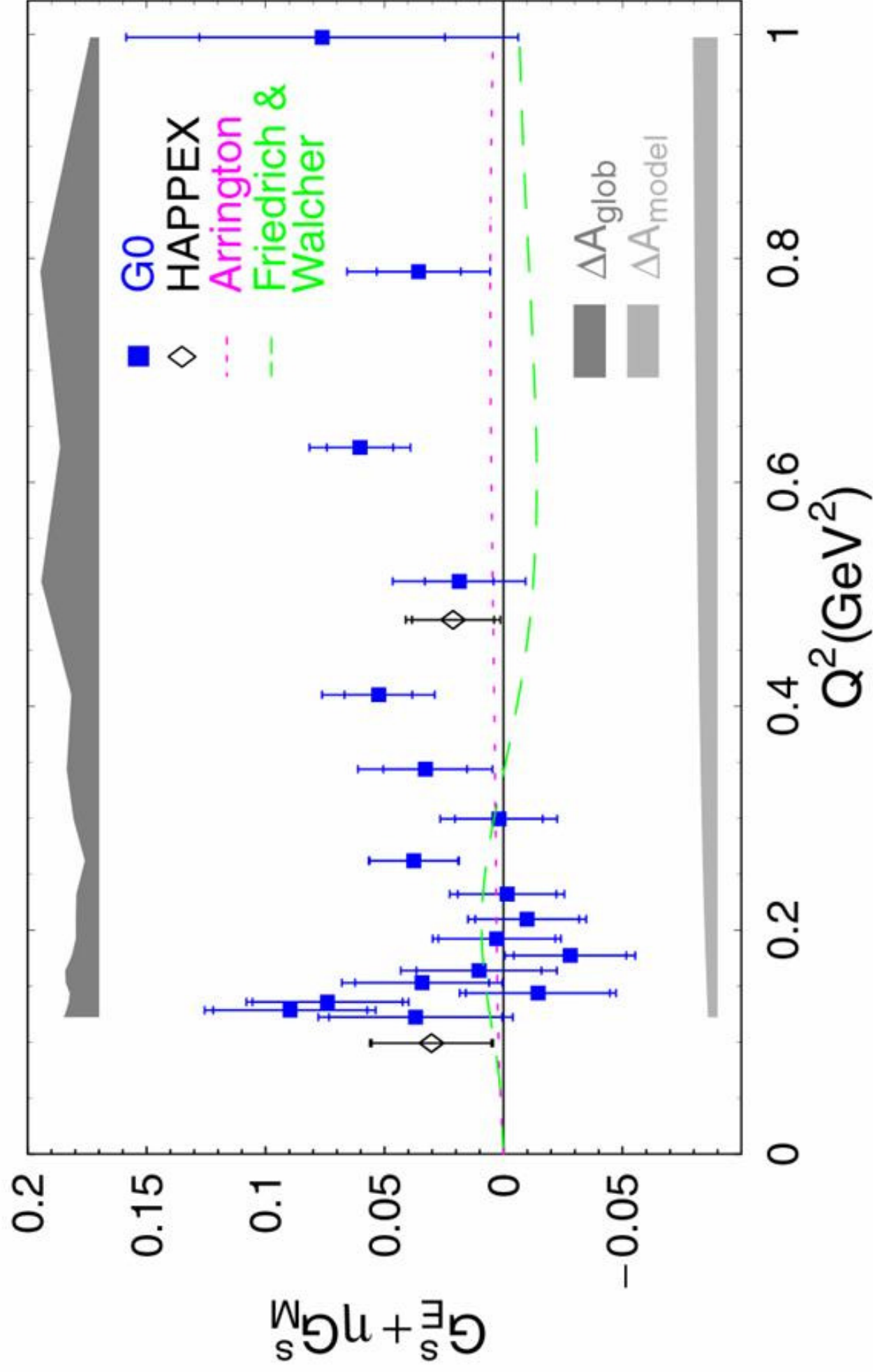


Experimental Asymmetries

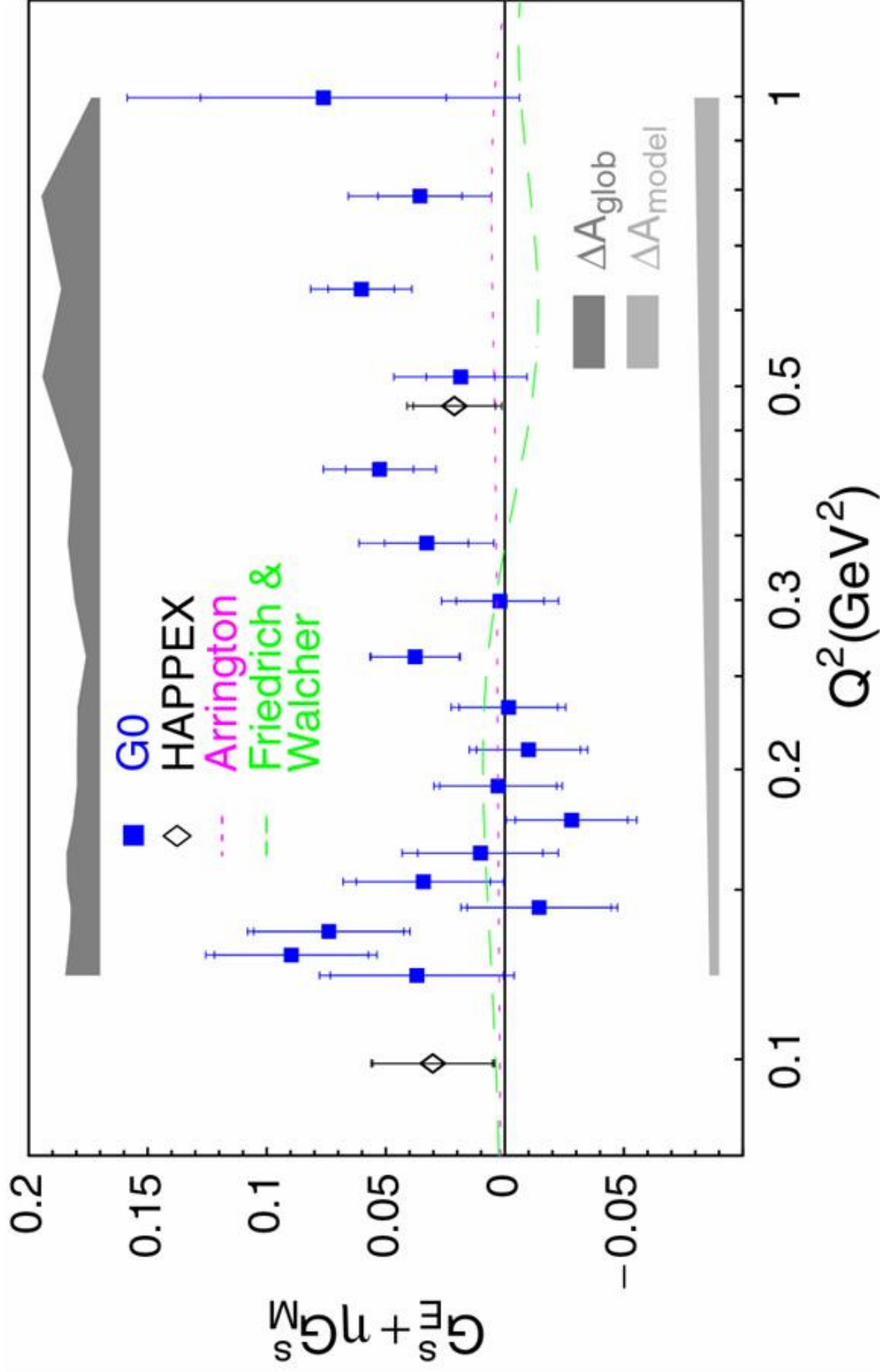
- em form factors: Kelly PRC 70 (2004) 068202
- “no vector strange” asymmetry, A_{NVs} , is $A(G_E^S, G_M^S = 0)$
- inside error bars: stat, outside: stat & pt-pt



Strange Quark Contribution to Proton



Strange Quark Contribution to Proton



Speculation

Simple Fits to World Hydrogen Data

- Fit

$$G_E^s(Q^2) + \eta(Q^2, E_i) G_M^s(Q^2) = \frac{4\pi\alpha\sqrt{2} \varepsilon G_E^p{}^2 + \tau G_M^p{}^2}{G_F Q^2 \varepsilon G_E^p (1 + R_V^{(0)})} (A_{phys} - A_{MVS}(Q^2, E_i))$$

with simple forms for G_E^s , G_M^s

$$G_E^s(Q^2) = \frac{c_2 Q^4}{1 + d_1 Q^2 + d_2 Q^4 + d_3 Q^6}$$

$$G_M^s(Q^2) = \frac{G_M^s(Q^2 = 0)}{(1 + Q^2 / \Lambda_M^s)^2}$$

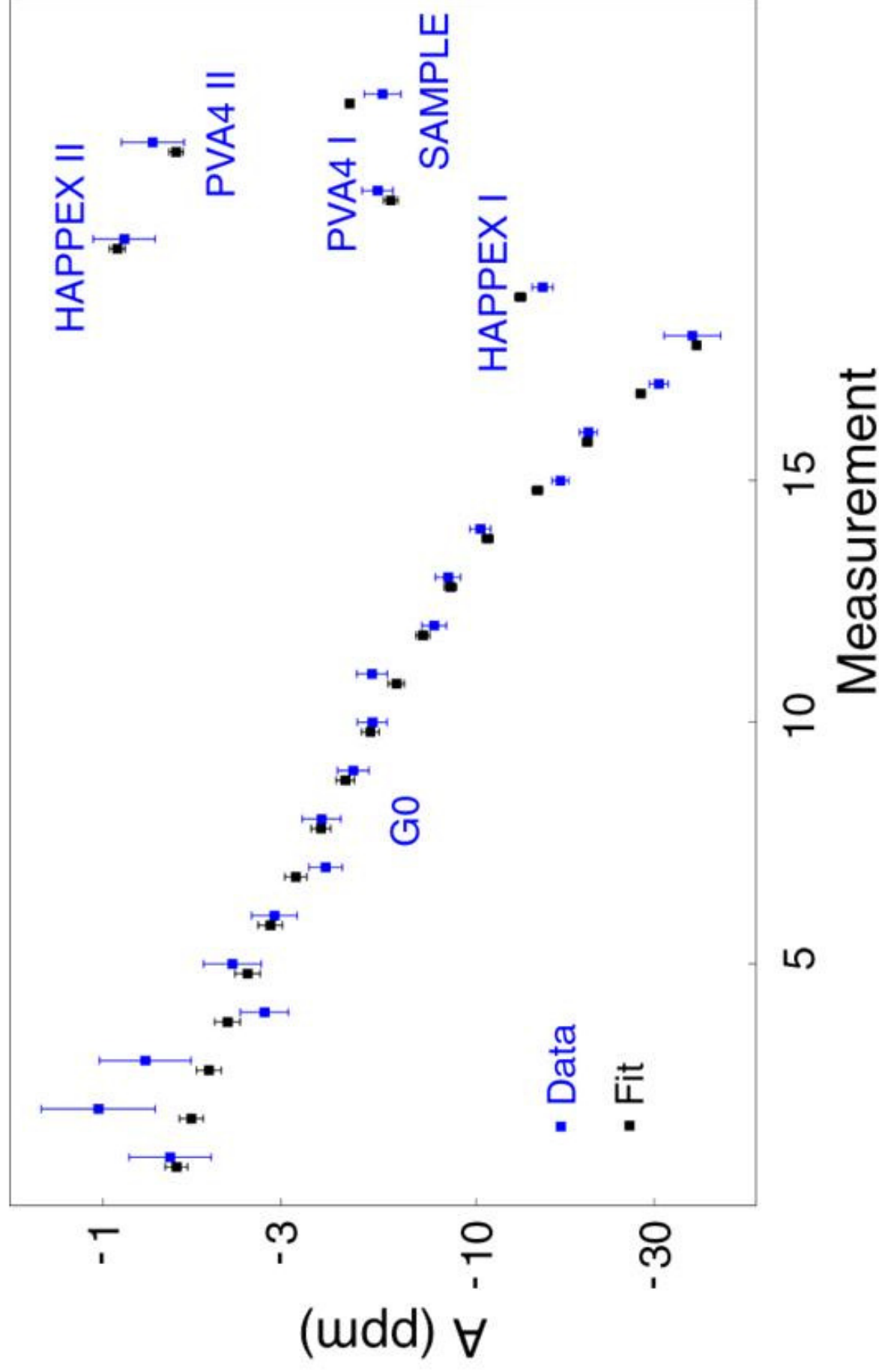
à la Kelly

with

$$G_M^s(Q^2 = 0) = 0.81 \quad \text{from } Q^2 = 0.1 \text{ GeV}^2 \text{ plot, dipole ff}$$

“Fit” to World Hydrogen Data

- $\chi^2 = 31/20$



“Fit” to World Hydrogen Data

$$c_2 = -0.51 \pm 0.25$$

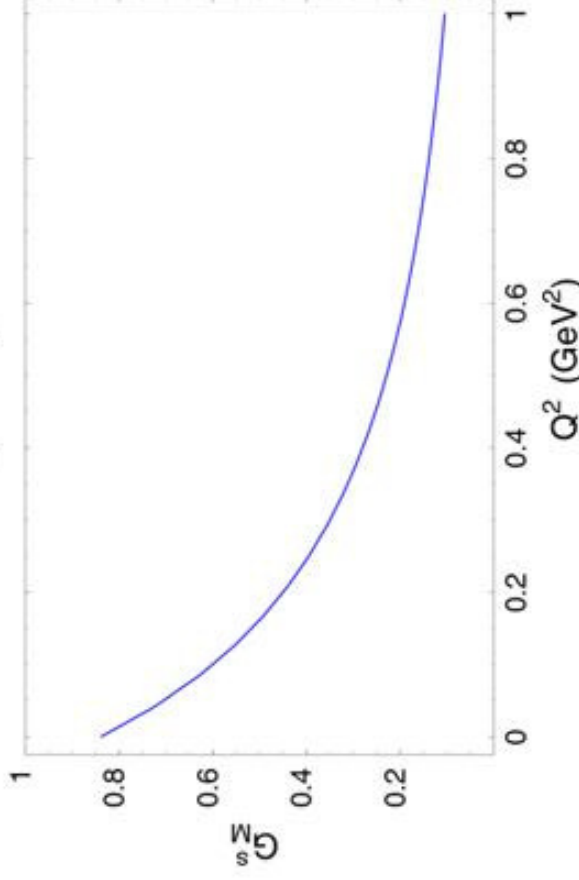
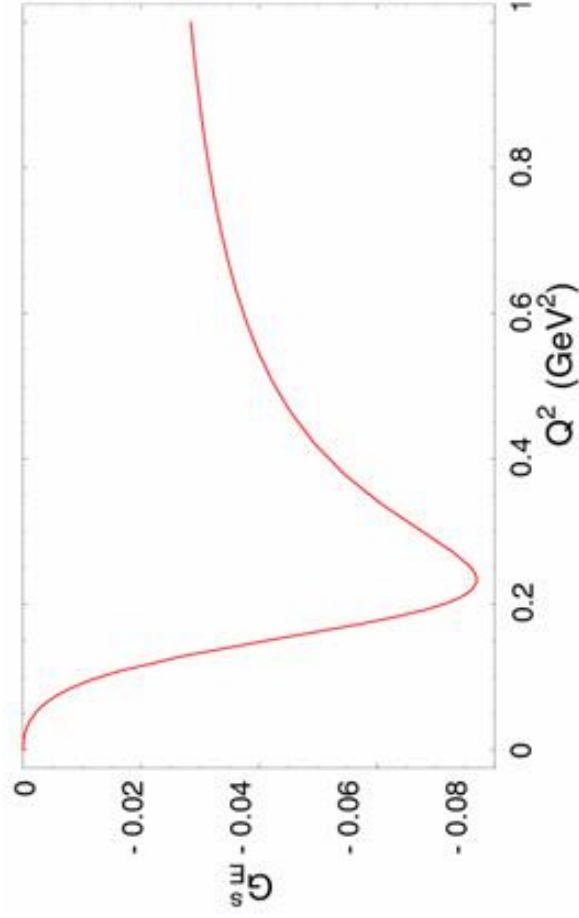
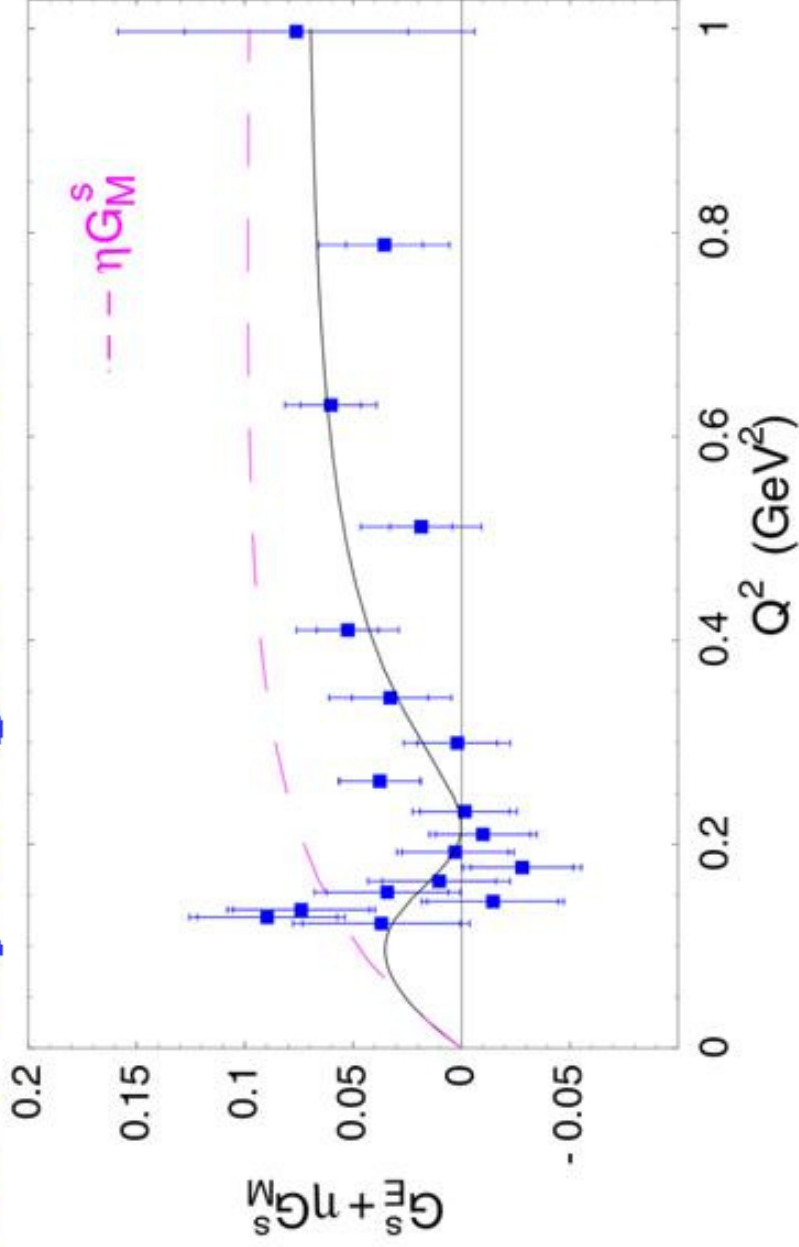
$$d_1 = -8.5 \pm 0.9$$

$$d_2 = 24 \pm 6$$

$$d_3 = 1$$

$$\Lambda_M^{s^2} = \Lambda^2 / 1.3$$

Remember the factor of $-1/3$



G0 Summary

- First measurement of parity-violating asymmetries over broad Q^2 range
- Excellent performance of accelerator, experimental equipment
- Conservative estimates of uncertainties
 - careful assessment of backgrounds
- Results consistent with previous measurements
- Emerging picture
 - $G_M^S > 0$ at low Q^2
 - $G_E^S < 0$ at medium Q^2 a possibility
 - $G_E^S + \eta G_M^S$ positive at higher Q^2