

# Neutrino Mass

The neutrino mass search technique sensitive to the smallest  $\nu$  masses is the search for  $\nu$  oscillation.

Under a number of circumstances,

$$P(\nu_{\ell} \rightarrow \nu_{\ell' \neq \ell}) = \sin^2 2\theta \sin^2 \left[ 1.27 \delta M^2 (\text{eV}^2) \frac{L(\text{km})}{E_{\nu}(\text{GeV})} \right]$$

Mixing angle  $\theta$  points to  $\sin^2 2\theta$ .  
 Distance  $\nu$  travels points to  $L(\text{km})$ .  
 $\nu$  energy points to  $E_{\nu}(\text{GeV})$ .  
 $M_2^2 - M_1^2$  points to  $\delta M^2$ .  
 $\ell, \mu, \tau$  points to  $\nu_{\ell}$ .

## 5.1] Neutrino Oscillation

There are 3 pieces of evidence that neutrinos oscillate, which implies that they have mass.

Neutrinos

Evidence of Oscillation

Atmospheric

Compelling

Solar

Strongish

LSND

Unconfirmed

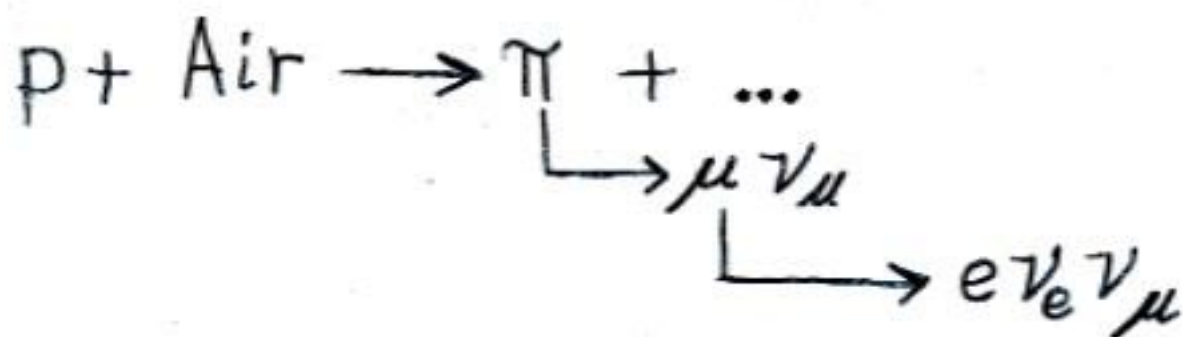
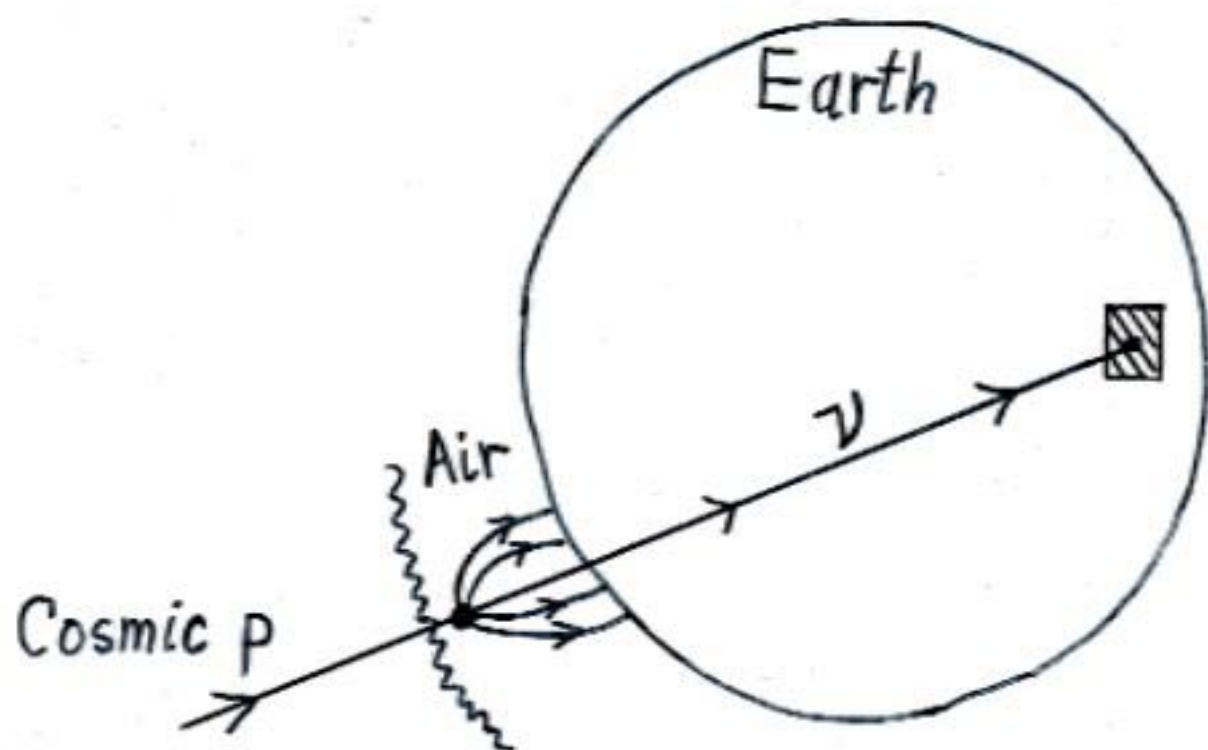
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Recent review of the current situation:

[hep-ph/9906244](https://arxiv.org/abs/hep-ph/9906244)

(Fisher, B.K., McFarland)

# Evidence for Neutrino Mass From the Atmospheric Neutrinos



$$\therefore \nu_{\mu} : \nu_e = 2 : 1 \quad [\text{with corrections}]$$

Try to measure something that is as independent of theoretical flux calculations as possible:



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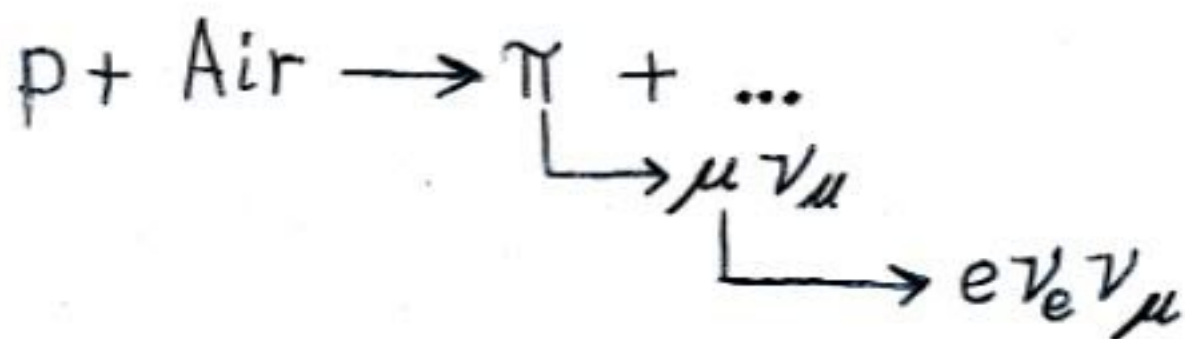
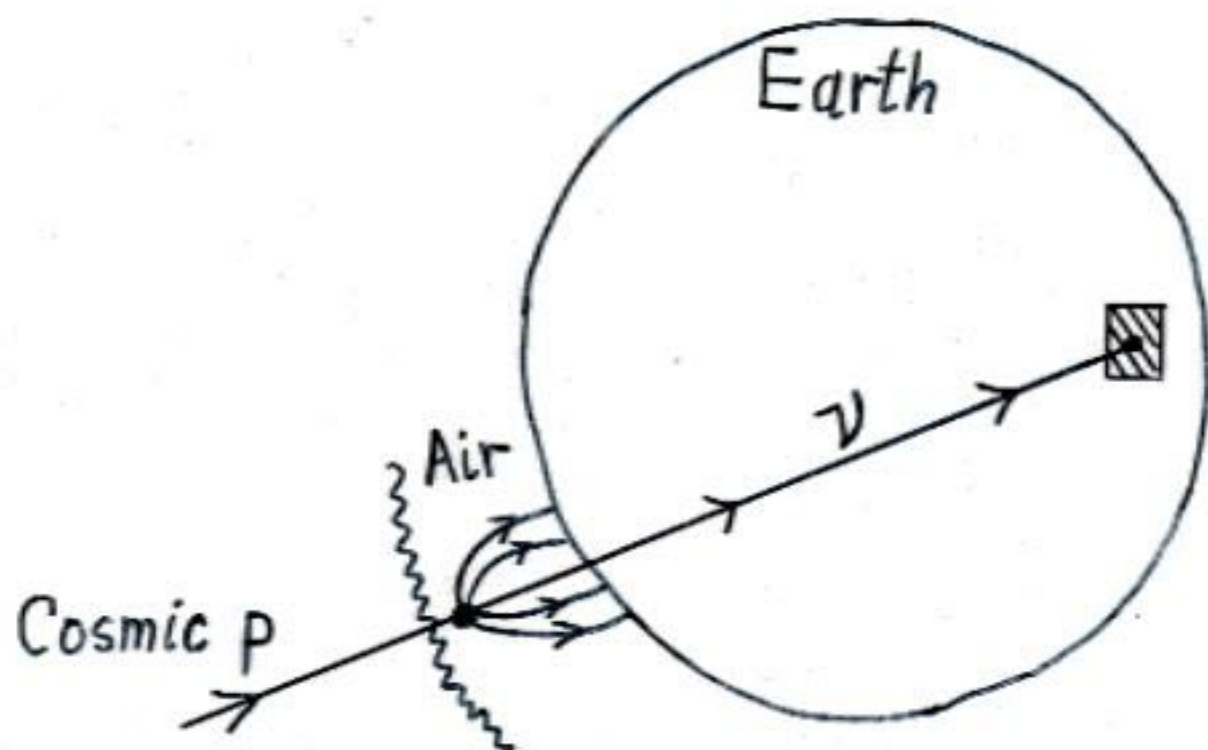
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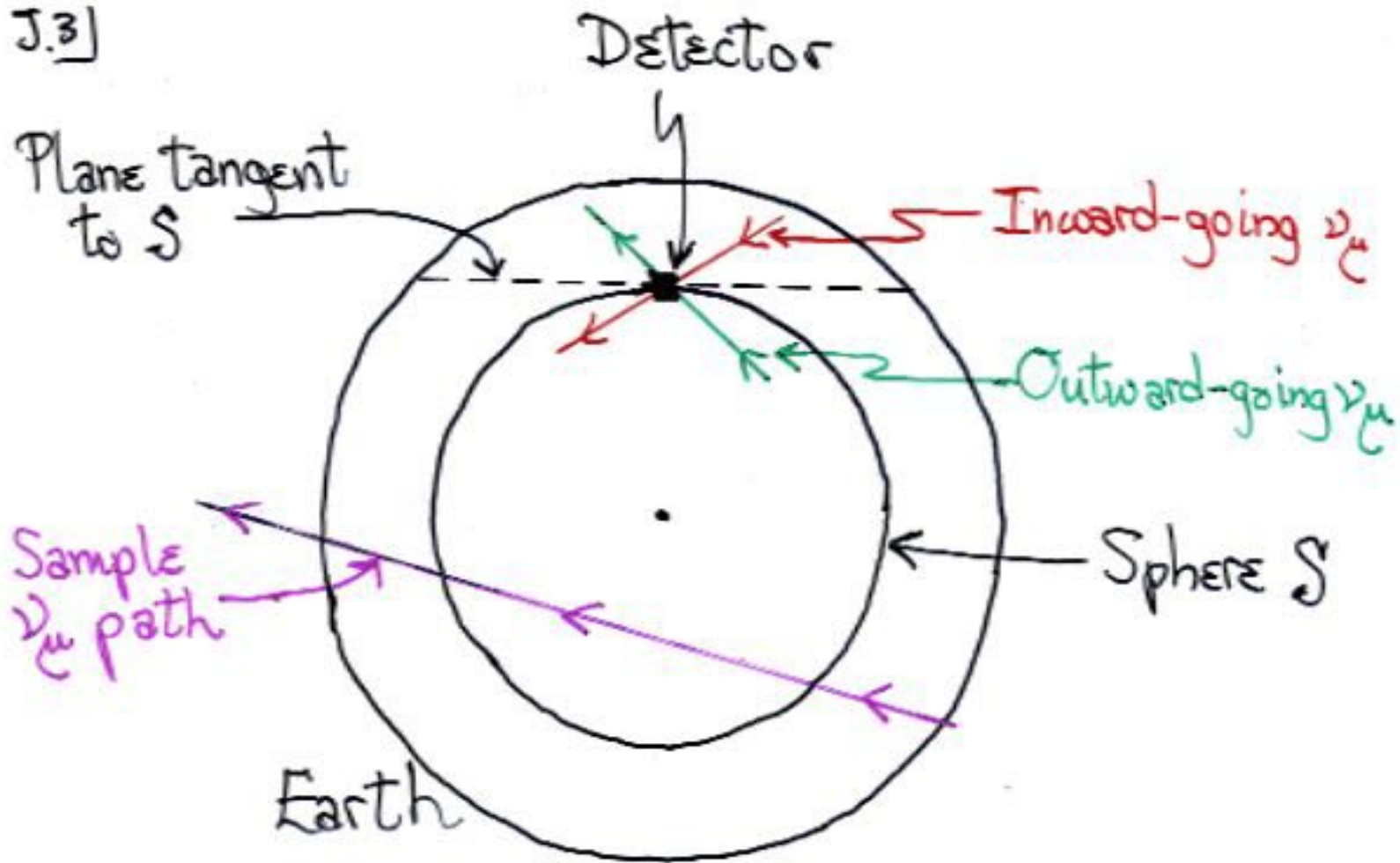
## Flux Down vs. Flux Up

Suppose nothing increases or decreases the atmospheric  $\nu_\mu$  flux during earth traversal.

Then, for  $E_\nu >$  a few GeV, where the cosmic ray flux is  $\approx$  isotropic,

$$\nu_\mu \text{ Flux Down} \approx \nu_\mu \text{ Flux Up}$$

J.3]



Any  $\nu_\mu$  that enters  $S$  later exits  $S$ .

$\therefore$  In a steady state—

Total  $\nu_\mu$  flux into  $S$  = Total  $\nu_\mu$  flux out of  $S$ .

Cosmic-ray isotropy  $\Rightarrow$  Spherical symmetry.  
 $\Rightarrow$  At each point of  $S$ —

$\nu_\mu$  flux into  $S$  =  $\nu_\mu$  flux out of  $S$

$\therefore$   $\nu_\mu$  Flux Down =  $\nu_\mu$  Flux Up



J.4]

With just a bit more work —

$$\nu_{\mu} \text{ Flux}(\theta_z) = \nu_{\mu} \text{ Flux}(\pi - \theta_z).$$

## SuperK Down vs. Up

Down :  $+0.2 < \cos \theta_z < +1.0$

Up :  $-1.0 < \cos \theta_z < -0.2$

In the Multi-GeV data,

$$\frac{\nu_{\mu}(\text{Up})}{\nu_{\mu}(\text{Down})} = 0.53 \pm 0.05.$$

(Learned)

Something is adding or removing muon neutrinos within the earth.

$\nu_{\mu} \rightarrow \nu_{\tau}$  or  $\nu_{\mu} \rightarrow \nu_{\text{sterile}}$  oscillation?

Not  $\nu_{\mu} \rightarrow \nu_e$  because  $\nu_e(\text{Up})/\nu_e(\text{Down}) = 0.92 \pm 0.11$ .

C.I.2

The zenith-angle distributions are well described by

$$\nu_{\mu} \rightarrow \nu_{\tau} \text{ with } \delta M^2 = 3.5 \times 10^{-3} \text{ eV}^2, \sin^2 2\theta = 1.0$$

Combined analysis of all SuperK  $\nu_{\text{Atmos}}$  data  $\Rightarrow$

$$\nu_{\mu} \rightarrow \nu_{\tau} \text{ with}$$

$$2 \times 10^{-3} \text{ eV}^2 \lesssim \delta M^2 \lesssim 6 \times 10^{-3} \text{ eV}^2$$

and

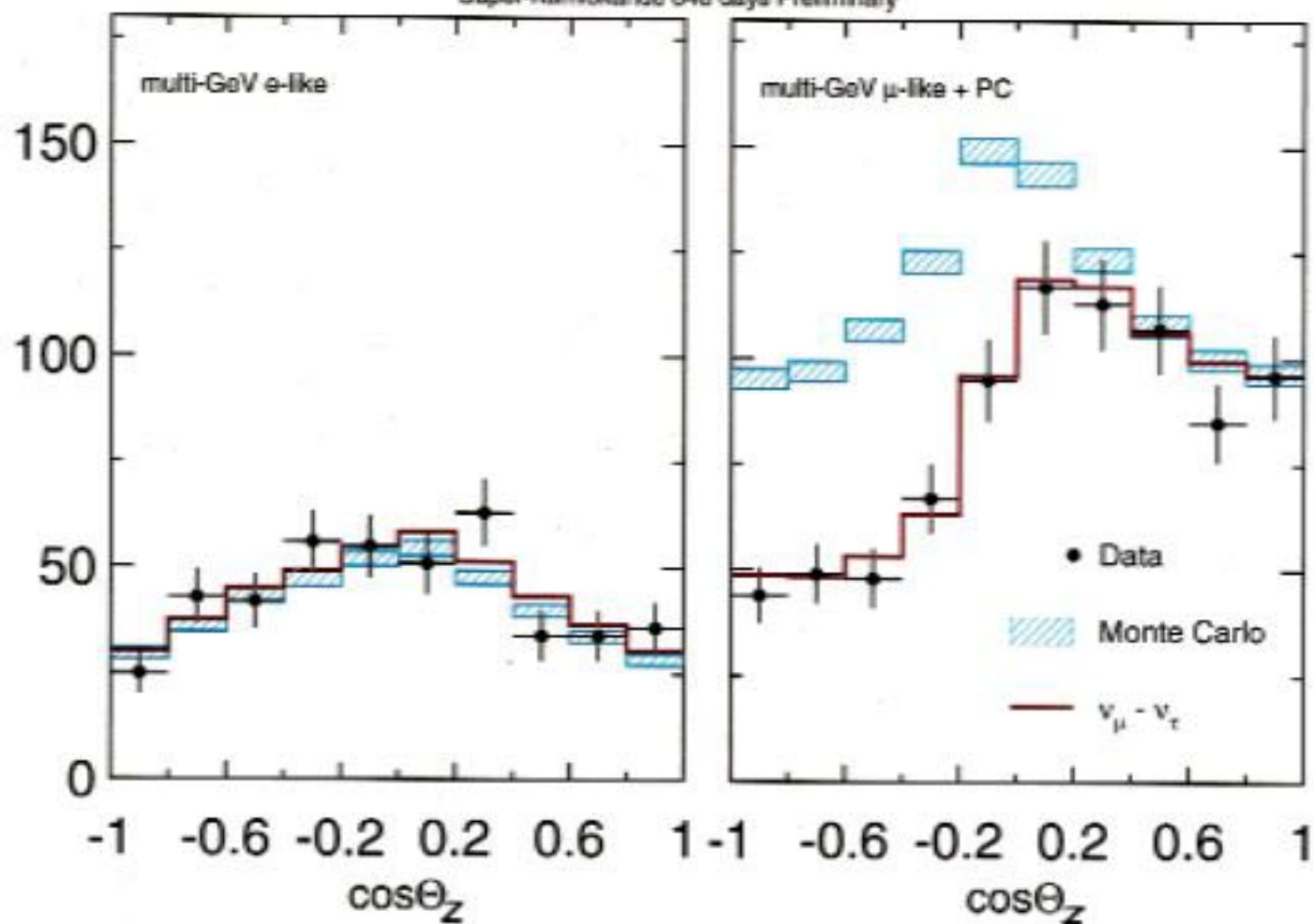
$$\sin^2 2\theta \approx 1.$$



Similar favored regions from  
MACRO, Soudan 2, and KAM.

When all data are considered,  $\nu$  decay and  $\nu$ -matter interaction do not work.

(Barger, Learned, Pakvasa, Weiler; Gonzalez-Garcia et al.; Lipari, Lusignoli; Fogli et al.)

Super-Kamiokande 848 days Preliminary



 , hence  , does depend on theory.



S.1)

Oscillation

Matter Effect in the Earth

$$\nu_{\mu} \rightarrow \nu_{\tau}$$

Yes

$$\nu_{\mu} \rightarrow \nu_{\tau}$$

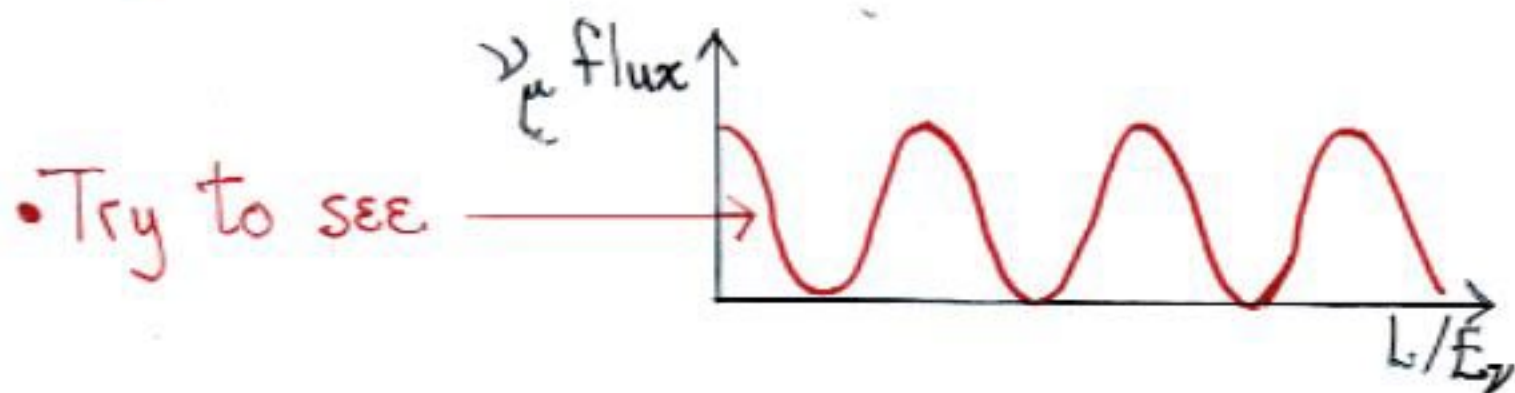
No

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Higher energy  $\nu_{\text{Atmos}}$  data (upward-going muons, and partially-contained events) disfavor  $\nu_{\mu} \rightarrow \nu_{\tau}$  at  $\sim 2\sigma$ . (Learned)

## Some key experiments

Test the  $\nu_\mu \rightarrow ?$  oscillation interpretation of the  $\nu_{\text{Atmos}}$  data by seeking  $\nu_\mu \rightarrow ?$  in long-baseline accelerator-generated  $\nu_\mu$  beams.



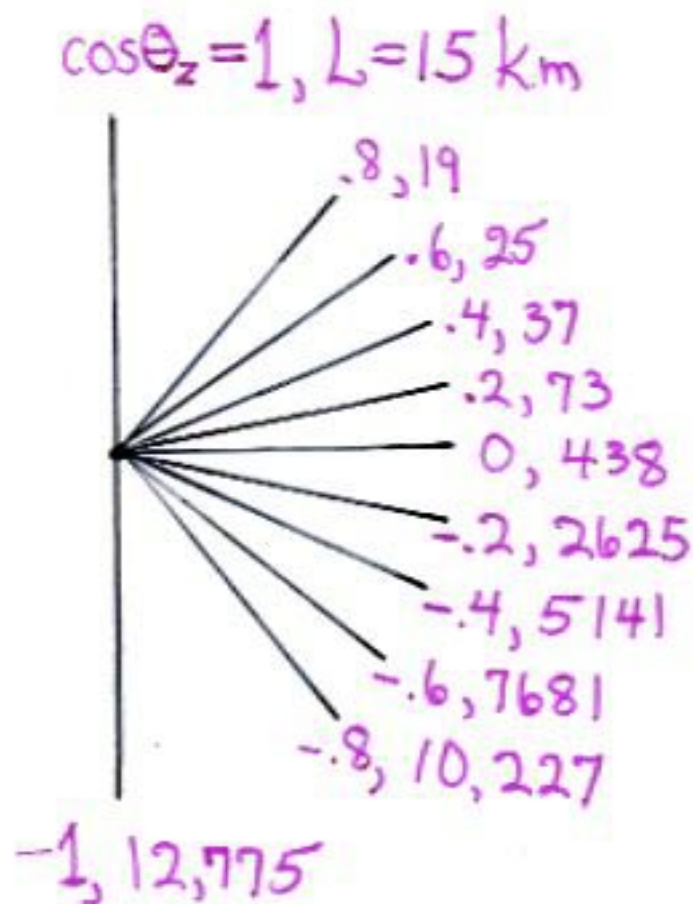
$$\left\{ P(\nu_\mu \rightarrow ?) \propto \sin^2 \left[ 1.27 \delta M^2 L / E_\nu \right] \right\}$$

• Try to see if  $? = \nu_\tau$  by seeking  $\nu_\tau$  appearance.

Near term: Conventionally produced  $\nu$  beams  
 Long term:  $\nu$  beams from muon facilities

W For  $\delta M^2 = 3.5 \times 10^{-3} \text{ eV}^2$ , and  $E_\nu = 1 \text{ GeV}$ ,  
 $\lambda_{\text{osc}} = 707 \text{ km}$ .

Then -



(Di Lella)



# I.2] The hints of neutrino mass

## Solar Neutrinos

Experimental reports:

Target	Exp.	$E_{\nu}$ Threshold (MeV)	Event Rate Std. Solar Model*
$H_2O(e)$	KAM	$\sim 7$	$.54 \pm .12$
$\gg$	SuperK	$\sim 6$	$.47 \pm .07$
$^{37}Cl$	HMSTK	0.8	$.33 \pm .06$
$^{71}Ga$	SAGE	0.2	$.52 \pm .07$
$\gg$	GALLEX	0.2	$.60 \pm .07$

\*Bahcall & Pinsonneault, 1998

It has proved very difficult to explain these departures from theory via solar and nuclear physics, without invoking neutrino mass.

(Updated analysis:  
Hata & Langacker.)

$\nu_e \rightarrow \nu_{\mu, \tau, \text{ or } s}$  [which implies  $\nu$  mass]  
explains all the  $\nu_o$  event rates.

[MSW]

Flavor Conversion  
Mechanism

Requirements

<u>Flavor Conversion Mechanism</u>	<u><math>\Delta M^2</math> (eV<sup>2</sup>)</u>	<u><math>\sin^2 2\theta</math></u>
MSW Effect	$(0.4-1.0) \times 10^{-5}$	$10^{-3} - 10^{-2}$
» »	$(2-8) \times 10^{-5}$	0.7-0.9
Vacuum Oscillation	$10^{-11} - 10^{-9}$	0.7-1.0

---

(Hata & Langacker)  
(Krauss & Petcov)

If  $P(\nu_e \rightarrow \nu_{\mu, \tau}) = \frac{1}{2}$

At ALL  $E_\nu$ ,

We Expect

.57

.57

.50

.50

.50



$P(\nu_e \rightarrow \nu_{\mu, \tau}) = \frac{1}{2}$ , independent of  $E_\nu$ ,  
could be explained by vacuum  
oscillation with  $\delta M^2$  as large  
as  $\sim 10^{-3} \text{ eV}^2$ .

E.2]

## Some key experiments

Seek a  $\nu_\mu$  or  $\nu_\tau$  component in the solar neutrino flux, produced by

$$\nu_e \xrightarrow[\text{Vac. Osc.}]{\text{MSW or}} \nu_\mu \text{ or } \nu_\tau$$

Caution: If  $\nu_e \rightarrow \nu_{\text{sterile}}$ , there is no  $\nu_{\mu, \tau}$  component in the solar flux.

Continue to probe the dependence of the solar  $\nu_e$  flux on —

- Neutrino Energy
- Whether it is Day or Night
- The Season of the Year

# LSND

A so-far unconfirmed further hint of  $\nu$  mass.

From  $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$   $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$   $\uparrow$  Appearance

$L/E_\nu \sim 30 \text{ m} / 30 \text{ MeV}$ , so  $\delta M^2 \gtrsim 1 \text{ eV}^2$ .

Taking into account constraints from other experiments (KARMEN, Bugey, NOMAD, ...):

$$0.2 \text{ eV}^2 \lesssim \delta M^2 \lesssim 10 \text{ eV}^2$$

$$0.002 \lesssim \sin^2 2\theta \lesssim 0.03$$



### E.3] Some key experiments

Confirm or disprove the  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  interpretation of the LSND data in experiments at other facilities.

## A Cosmological Constraint on $\nu$ Mass

There are 115 neutrinos/species/cm<sup>3</sup> in the universe.

Even with tiny masses, neutrinos can contribute significantly to

$$\frac{\text{Mass Density}}{\text{Critical Density}} \equiv \Omega.$$

Evidence from several sources implies

$$\Omega_{\text{all matter}} \lesssim 0.4.$$

To allow the observed galactic structure to have formed,

$$\Omega_{\nu} \lesssim 0.2.$$

(N. Bahcall, Primack, Steigman)

C.2 If  $\nu$  mass eigenstate  $\nu_m$  has mass  $M_{\nu_m}$ ,

$$\Omega_\nu = (M_{\nu_1} + M_{\nu_2} + M_{\nu_3}) / 40 \text{ eV}.$$

$$\therefore M_{\nu_1} + M_{\nu_2} + M_{\nu_3} \lesssim 8 \text{ eV}.$$

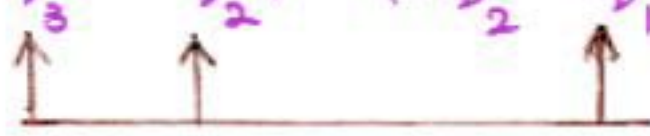


# Where Does the $\nu$ Oscillation Evidence Point?

$\nu_e$ ,  $\nu_{\text{Atmos}}$ , and  $\nu_{\text{LSND}}$  oscillations cannot all be explained in terms of just 3 neutrinos:

With only 3 neutrinos,

$$\sum \delta M^2 = (M_{\nu_3}^2 - M_{\nu_2}^2) + (M_{\nu_2}^2 - M_{\nu_1}^2) + (M_{\nu_1}^2 - M_{\nu_3}^2) = 0.$$

 Mass eigenstates

But—

Oscillating Neutrinos

Solar  
Atmospheric  
LSND

Required  $|\delta M^2|$  (eV<sup>2</sup>)

$10^{-10}$  or  $10^{-5}$   
 $10^{-3}$   
1

---

$$\sum \delta M^2 \neq 0$$

$\therefore$  Must add a sterile neutrino  $\nu_s$ .

Must not contribute to  $\Gamma_{z^0}$

There is a loophole in this argument.

No attempt to exploit the loophole has succeeded in fitting all the data.

(Cardall & Fuller; Teshima, Sakai, Inagaki;  
Thun & McKee; Barenboim & Scheck;  
Ohlsson & Snellman)

Even if we disregard HMSTK, and suppose  $\delta M_{\text{solar}}^2$  can be large, it is still very hard to make 3 neutrinos work.

(Giunti)

Confirmation of a  $\nu_s$  would be groundbreaking.

$\therefore$  We need to check each of the 3 reported signals of neutrino oscillation.



## CP in Neutrino Oscillation

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_{e'}; U) \stackrel{\text{CPT}}{=} P(\nu_e \rightarrow \nu_{e'}; U^*)$$

↑  
Neutrino mixing matrix

If  $U$  is not real, we can have CP differences

$$\Delta_{CP}(ll') \equiv P(\nu_e \rightarrow \nu_{e'}) - P(\bar{\nu}_e \rightarrow \bar{\nu}_{e'}) \neq 0.$$

Such differences would cleanly establish that  $U$  contains a complex phase.

$$\text{CPT} \Rightarrow P(\nu_e \rightarrow \nu_e) = P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$$

$\therefore$  No CP in a disappearance neutrino oscillation experiment.

It takes at least 3 non-degenerate neutrino mass eigenstates  $\nu_m$  to produce CP.



F.4] To see  $\mathcal{CP}$ , an experiment must have  $L/E_\nu$  big enough so that

$$\delta M_{mm'}^2 (\text{eV}^2) \frac{L (\text{km})}{E_\nu (\text{GeV})} \gtrsim 1$$

for at least 3  $\delta M_{mm'}^2 \equiv M_m^2 - M_{m'}^2$ .  
Mass of  $\nu_m$

If there are only 3 neutrinos,

$$\delta M^2 (\text{MSW for } \nu_\odot) \sim 10^{-5} \text{ eV}^2$$

$$\Rightarrow \text{We need } \frac{L (\text{km})}{E_\nu (\text{GeV})} \gtrsim 10^5.$$

Physics for a  $\nu$  Factory?

(37)

$\mathcal{CP}$  analyses:

Arafune & Sato; Fisher, B.K., McFarland;  
Bernabéu; Schubert; Dick, Freund,  
Lindner, Romanino; Gago, Pleitez, Funchal

CP  
If there are only 3 neutrinos, then  
for  $ll' = e\mu, \mu\tau, \text{ or } \tau e,$

$$\Delta_{CP}(ll') = -16 J s_{12} s_{23} s_{31} .$$

Here  $J \equiv \text{Im}(U_{e1} U_{e2}^* U_{\mu 1}^* U_{\mu 2}),$

and  $s_{mm'} \equiv \sin \left[ 1.27 \delta M_{mm'}^2 (\text{eV}^2) \frac{L (\text{km})}{E_\nu (\text{GeV})} \right],$

where  $\delta M_{mm'}^2 \equiv M_m^2 - M_{m'}^2.$

(Arafune & Sato  
Fisher, B.K., McFarland)

$\Delta_{CP}(ll')$  is the same for  $ll' = e\mu, \mu\tau, \text{ and } \tau e.$

C.7] Questions

Q: How do we test/falsify the various  $\nu$  mass scenarios?

A: LSND confirmation  $\Rightarrow$  ~~3  $\nu$  (No LSND)~~

$\nu_{\mu, \tau}$  from the sun  $\Rightarrow$   ~~$\nu_{\mu}$~~   $\rightarrow$   ~~$\nu_{\tau}$~~   $\rightarrow$   $\nu_{\text{sterile}}$

Q: Are sterile neutrinos unnatural?

A: Heavy ones are common in a see-saw picture.

Even light ones can be natural in some schemes. But if one, maybe three.

(Foot & Volkas  
Bereziani & Mohapatra)

But: Ockham's razor.



Q: Is  $\sim$  maximal mixing of  $\nu_{\mu}^{\text{Atmos}}$  with some other neutrino unnatural?

A: No,  $\nu_{\mu}$  may be maximally mixed with a  $\nu_s$ . This has no quark analogue.

Or, large mixing may arise from a family symmetry, or from the details of the see-saw mechanism.

(Foot & Volkas; Ramond; Smirnov)

Q: How do we measure  $\nu$  masses, and not just their splittings?

A: Good question.

## 5 Conclusion

This is an exciting time in neutrino physics.

We have lots of interesting work to do.