

The heavy-quark hybrid meson spectrum in lattice QCD

Colin Morningstar

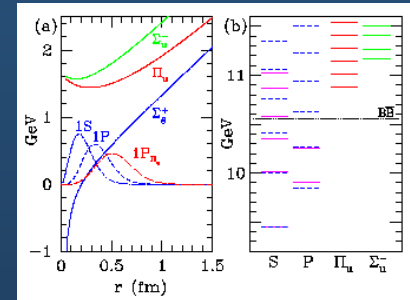
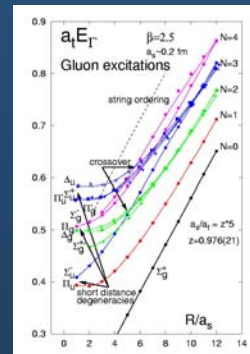
Carnegie Mellon University

Workshop on Gluonic Excitations, JLab

May 14, 2003

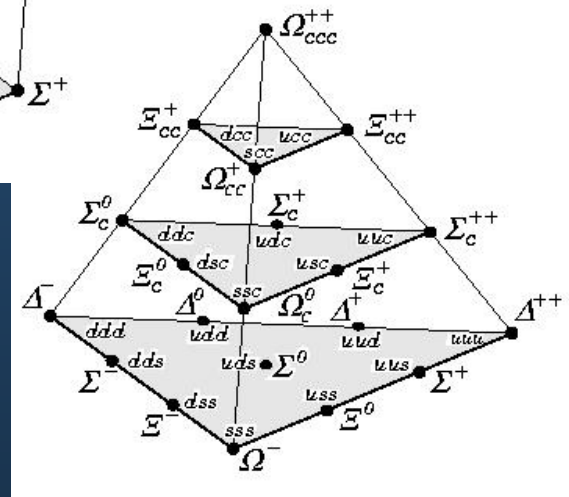
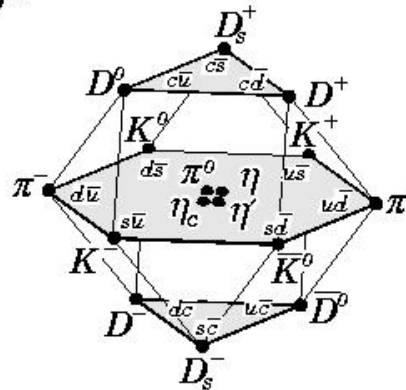
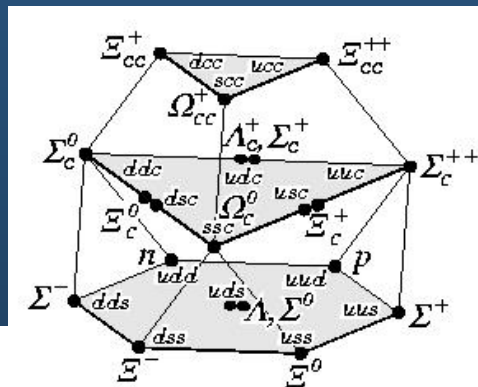
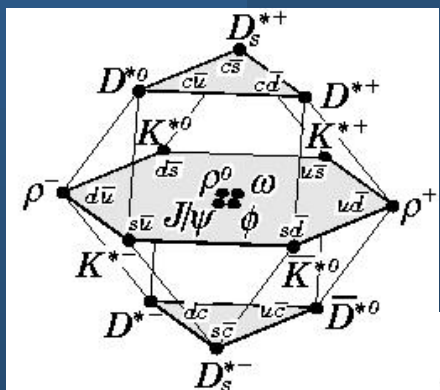
Outline

- Introduction
- Heavy-quark mesons: leading Born-Oppenheimer approximation
 - Stationary states of gluons in presence of static $q\bar{q}$ -pair
 - Leading order spectrum (no light quark pairs)
 - Testing the leading Born-Oppenheimer approximation
- Quark spin effects
- Incorporation of light quark loops
- Other tidbits
- Conclusion



Constituent quark model

- much of our understanding of hadron formation comes from the *constituent quark model*
 - motivated by QCD
 - valence quarks interacting via Coulomb + linear potential
 - gluons: source of the potential, *dynamics ignored*

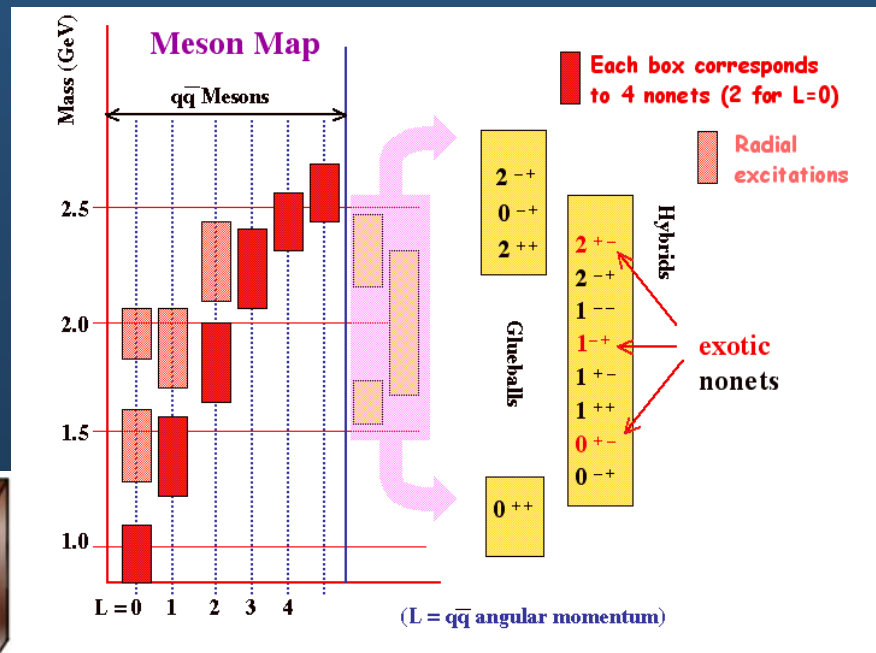


Quark model (continued)

- *most* of observed low-lying hadron spectrum described reasonably well by quark model
 - agreement is amazing given the crudeness of the model
- mesons: only certain J^{PC} allowed:
 - $0^{+-}, 0^{-+}, 1^{-+}, 2^{+-}, 3^{-+}, 4^{+-}, \dots$ forbidden
$$P = (-1)^{L+1} \quad L = 0, 1, 2, \dots$$
$$C = (-1)^{L+S} \quad S = 0, 1$$
- experimental results now need input beyond the quark model
 - over-abundance of states
 - forbidden 1^{-+} states

Gluonic excitations (new form of matter)

- QCD suggests existence of states in which *gluon* field is excited
 - glueballs (*excited glue*)
 - hybrid mesons ($q\bar{q}$ + *excited glue*)
 - hybrid baryons (qqq + *excited glue*)
- such states not well understood
 - quark model fails
 - perturbative methods fail
- lack of understanding makes identification difficult!
- confront gluon field behavior
 - bags, strings, ...
- clues to confinement

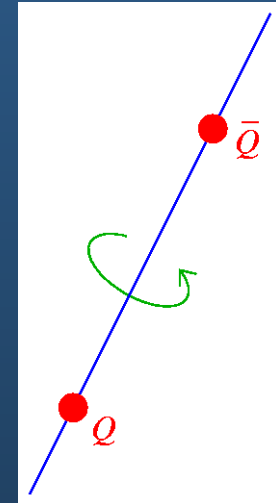


Heavy-quark hybrid mesons

- more amenable to theoretical treatment than light-quark hybrids
- early work: Hasenfratz, Horgan, Kuti, Richards (1983), Perantonis, Michael (1990)
- possible treatment like diatomic molecule (Born-Oppenheimer)
 - slow heavy quarks \leftrightarrow nuclei
 - fast gluon field \leftrightarrow electrons
(and light quarks)
- gluons provide adiabatic potentials $V_{Q\bar{Q}}(r)$
 - gluons fully relativistic, interacting
 - potentials computed in lattice simulations
- nonrelativistic quark motion described in *leading order* by solving Schrodinger equation for each $V_{Q\bar{Q}}(r)$

$$\left\{ \frac{p^2}{2\mu} + V_{Q\bar{Q}}(r) \right\} \psi_{Q\bar{Q}}(r) = E \psi_{Q\bar{Q}}(r)$$

- conventional mesons from Σ_g^+ ; hybrids from Π_u, Σ_u^-, \dots

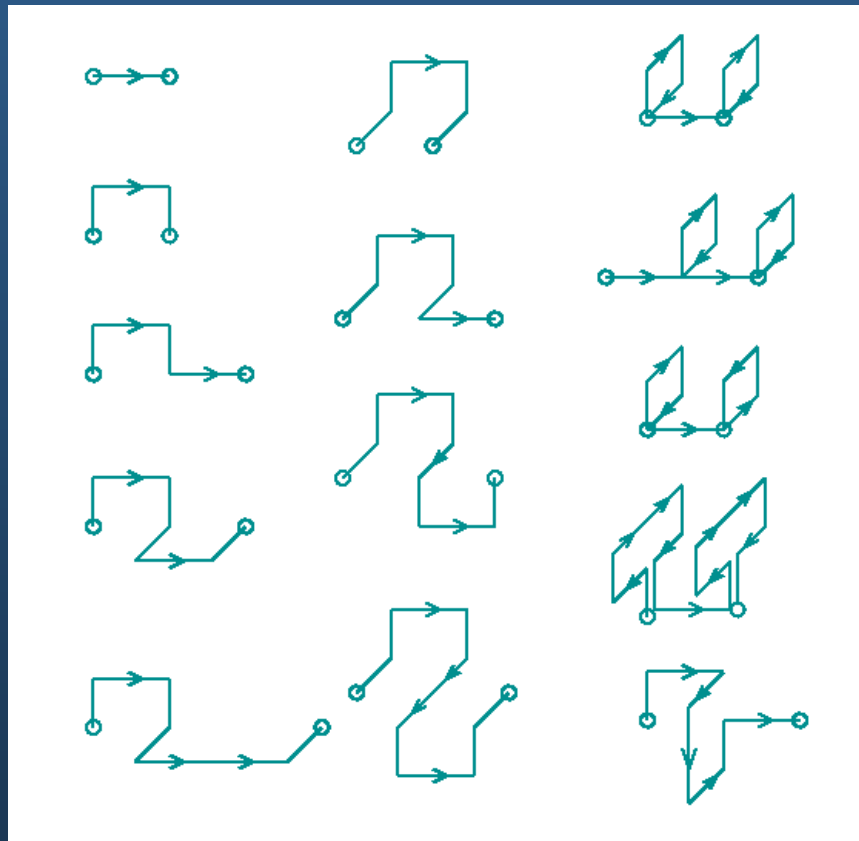


First step in Born-Oppenheimer

- first step in Born-Oppenheimer approximation
 - determine the “gluonic terms”
- calculational approach → resort to Monte Carlo methods
 - familiar perturbative Feynman diagram techniques fail
 - Schwinger-Dyson equations intractable
- estimate path integrals: very high dimensionality
 - Markov chain methods
- lattice regularization permits formulation of field theory suitable for computer simulations

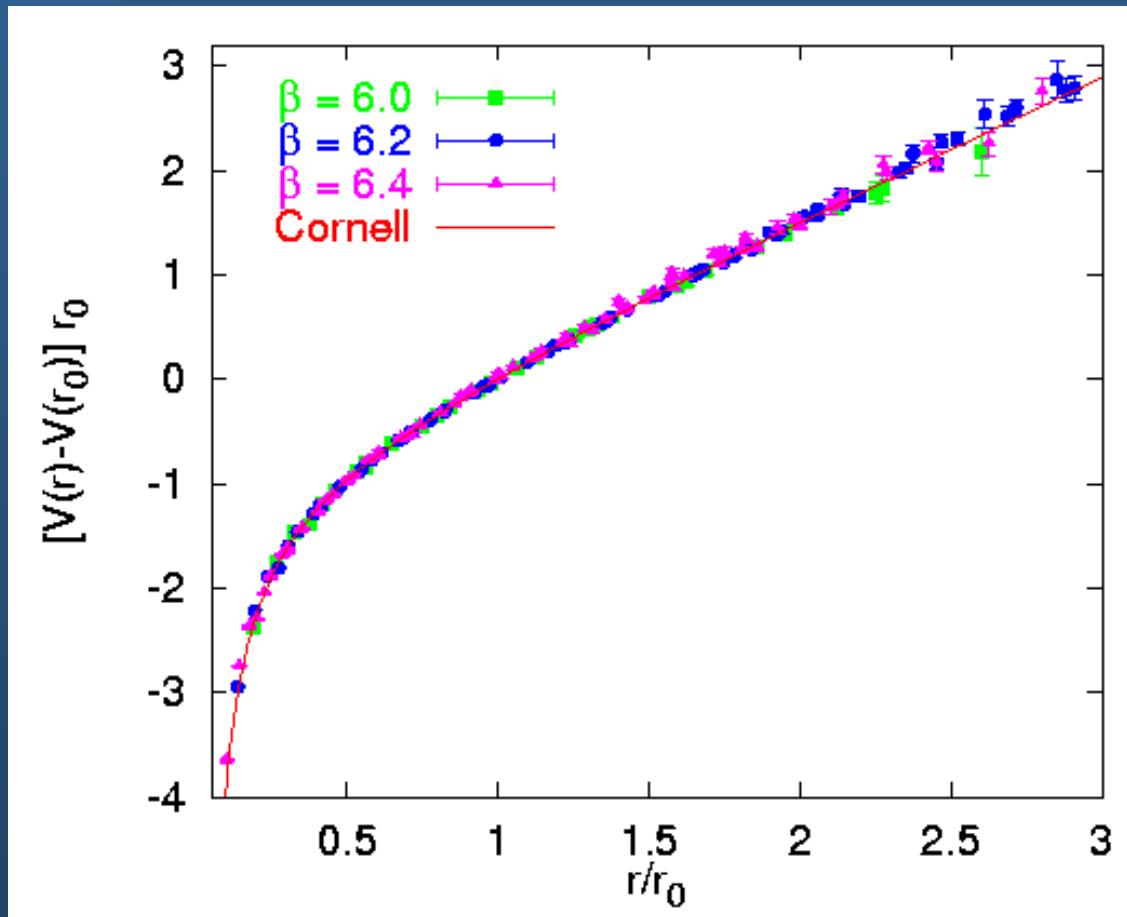
Generalized Wilson loops

- gluonic terms extracted from generalized Wilson loops
- large set of gluonic operators \rightarrow correlation matrix



Static quark-antiquark potential

- lattice simulations confirm linearly rising potential from gluon exchange

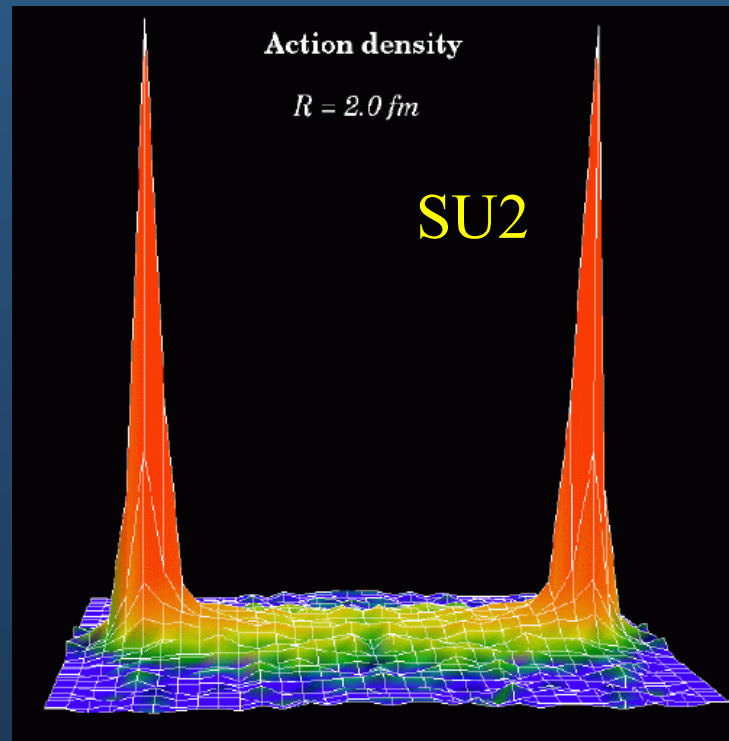


Bali *et al.*

$r_0 = 0.5$ fm

Gluonic flux profile

- computation of gluonic flux profile suggests that gluon field forms a *string*-like object between quark-antiquark



Bali *et al.*

Excitations of static quark potential

- gluon field in presence of static quark-antiquark pair can be *excited*
- classification of states: (notation from molecular physics)

- magnitude of glue spin
projected onto molecular axis

$$\Lambda = 0, 1, 2, \dots$$

$$= \Sigma, \Pi, \Delta, \dots$$

- charge conjugation + parity
about midpoint

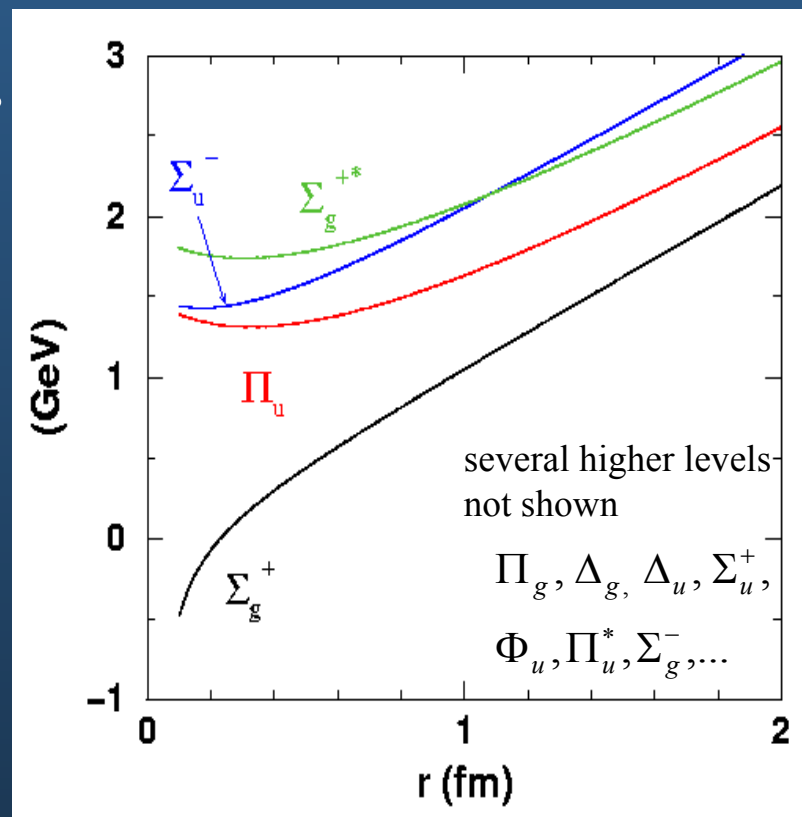
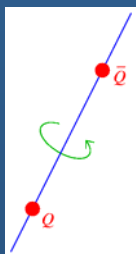
$$\eta = g \text{ (even)}$$

$$= u \text{ (odd)}$$

- chirality (reflections in plane
containing axis) Σ^+, Σ^-

Π, Δ, \dots doubly degenerate

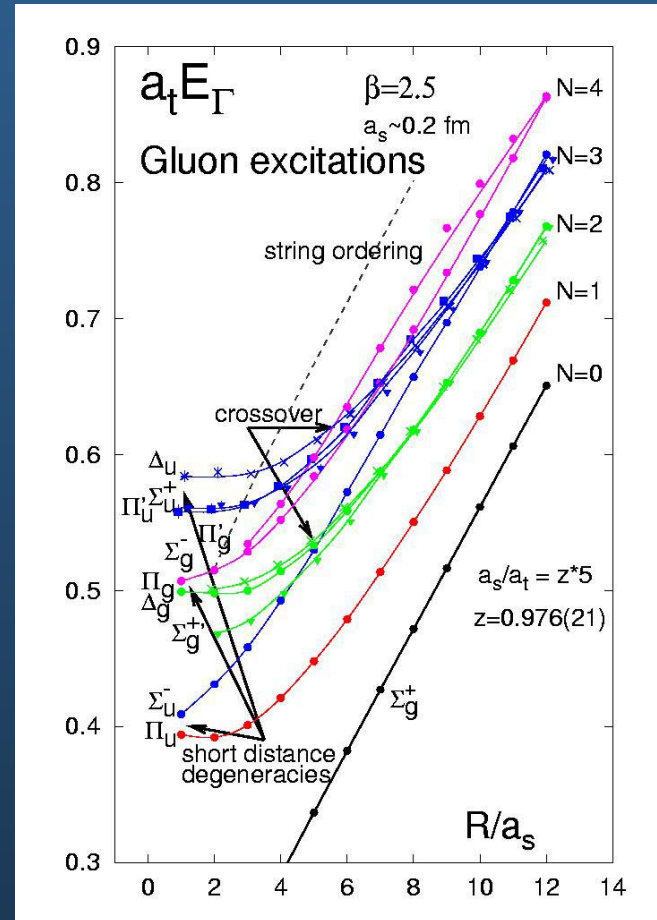
(Λ doubling)



Juge, Kuti, Morningstar, PRL 90, 161601 (2003)

Three scales

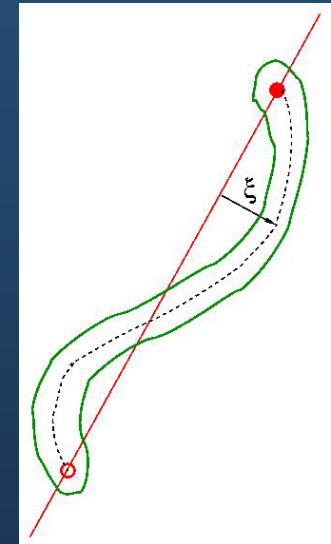
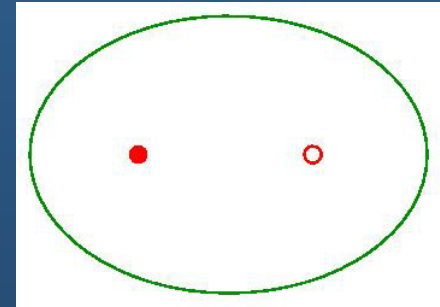
- small quark-antiquark separations r
 - excitations consistent with states from multipole OPE
- crossover region $0.5\text{fm} < r < 2\text{fm}$
 - dramatic level rearrangement
- large separations $r > 2\text{fm}$
 - excitations consistent with expectations from string models



Juge, Kuti, Morningstar, PRL 90, 161601 (2003)

Possible interpretation

- small r
 - strong E field of $q\bar{q}$ -pair repels physical vacuum (dual Meissner effect) creating a *bubble*
 - separation of degrees of freedom
 - gluonic modes inside bubble (low lying)
 - bubble surface modes (higher lying)
- large r
 - bubble stretches into thin tube of flux
 - separation of degrees of freedom
 - collective motion of tube (low lying)
 - internal gluonic modes (higher lying)
 - low-lying modes described by an effective string theory ($N\pi/r$ gaps – Goldstone modes)



Leading Born-Oppenheimer

- replace covariant derivative \vec{D}^2 by $\vec{\nabla}^2$ → neglects retardation
- neglect quark spin effects
- solve radial Schrodinger equation

$$\frac{-1}{2\mu} \frac{d^2 u(r)}{dr^2} + \left\{ \frac{\langle L_{q\bar{q}}^2 \rangle}{2\mu r^2} + V_{q\bar{q}}(r) \right\} u(r) = E u(r)$$

- angular momentum

$$\vec{J} = \vec{L} + \vec{S} \quad \vec{S} = \vec{s}_q + \vec{s}_{\bar{q}} \quad \vec{L} = \vec{L}_{q\bar{q}} + \vec{J}_g$$

- in LBO, L and S are good quantum numbers

- centrifugal term

$$\langle \vec{L}_{q\bar{q}}^2 \rangle = L(L+1) - 2\Lambda^2 + \langle \vec{J}_g^2 \rangle \quad \langle \vec{J}_g^2 \rangle = 0 \quad (\Sigma_g^+)$$

- J^{PC} eigenstates → Wigner rotations

$$= 2 \quad (\Pi_u, \Sigma_u^-)$$

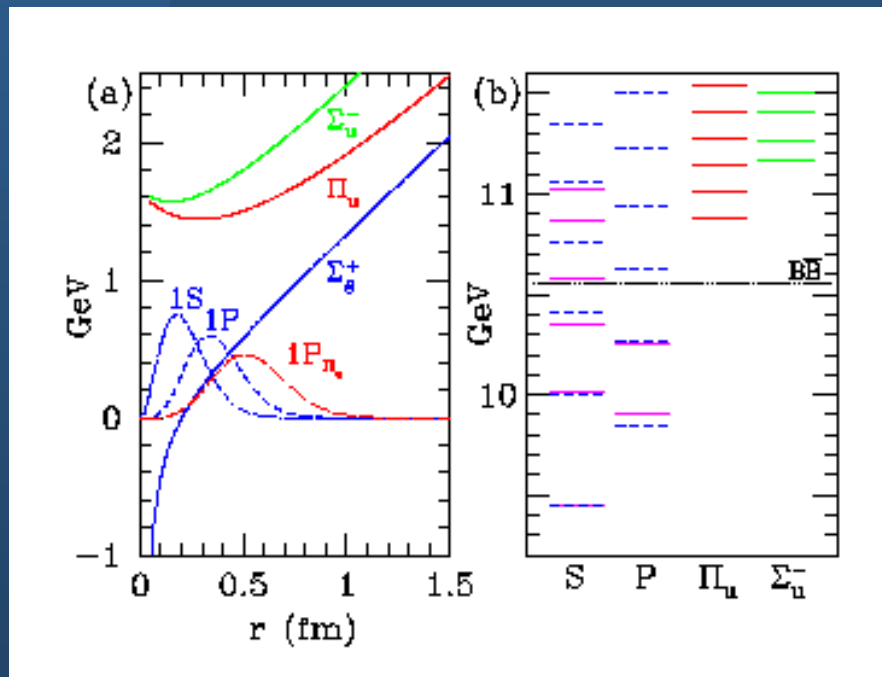
$$|LSJM; \Lambda \eta\rangle + \varepsilon |LSJM; -\Lambda \eta\rangle$$

□ η is CP, $\varepsilon = \pm 1$ for $\Lambda \geq 1$, $\varepsilon = \pm 1$ for Σ^\pm

- LBO allowed $J^{PC} \rightarrow P = \varepsilon(-1)^{L+\Lambda+1}$, $C = \eta\varepsilon(-1)^{L+S+\Lambda}$

Leading Born-Oppenheimer spectrum

- results obtained (in absence of light quark loops)
- good agreement with experiment below $B\bar{B}$ threshold
- plethora of hybrid states predicted (caution! quark loops)
- but is a Born-Oppenheimer treatment valid?



LBO degeneracies:

$$\Sigma_g^+(S): 0^{++}, 1^{--}$$

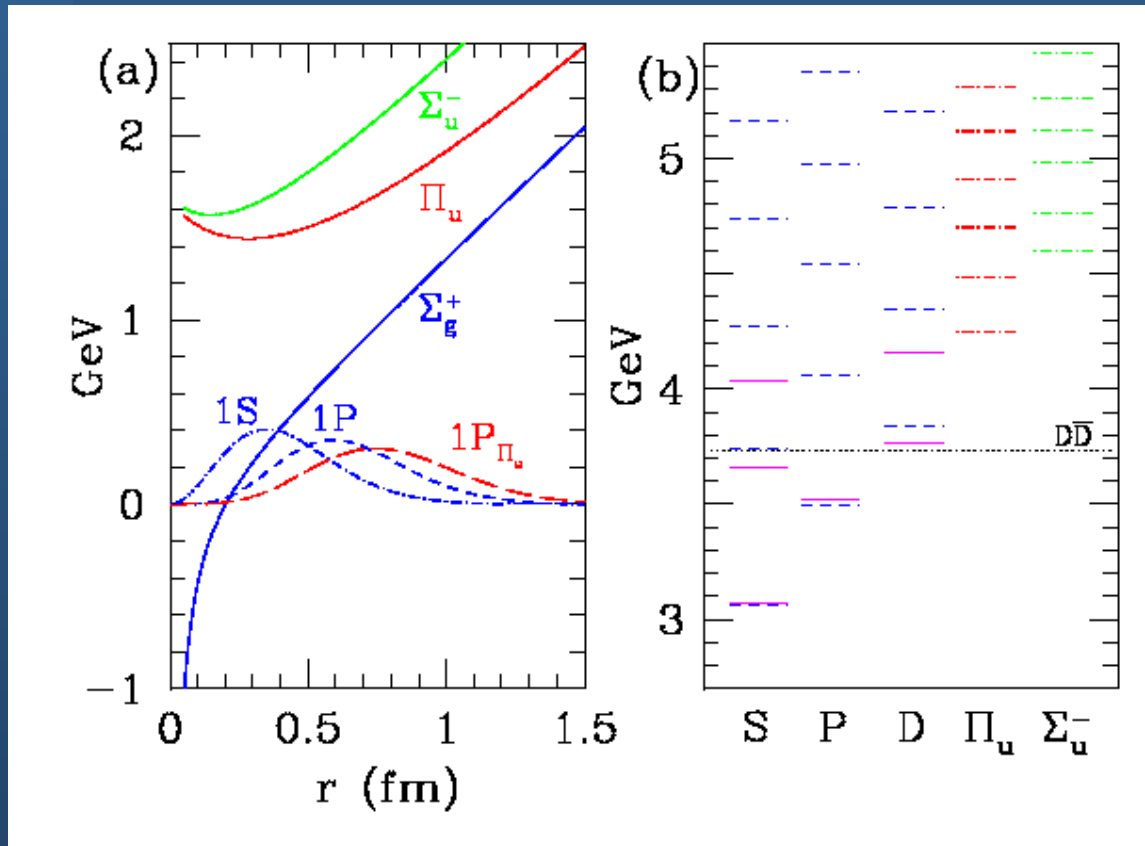
$$\Sigma_g^+(P): 0^{++}, 1^{++}, 2^{++}, 1^{+-}$$

$$\Pi_u(P): 0^{+-}, 0^{++}, 1^{++}, 1^{--}, \\ 1^{+-}, 1^{+-}, 2^{+-}, 2^{+-}$$

Juge, Kuti, Morningstar, Phys Rev Lett **82**, 4400 (1999)

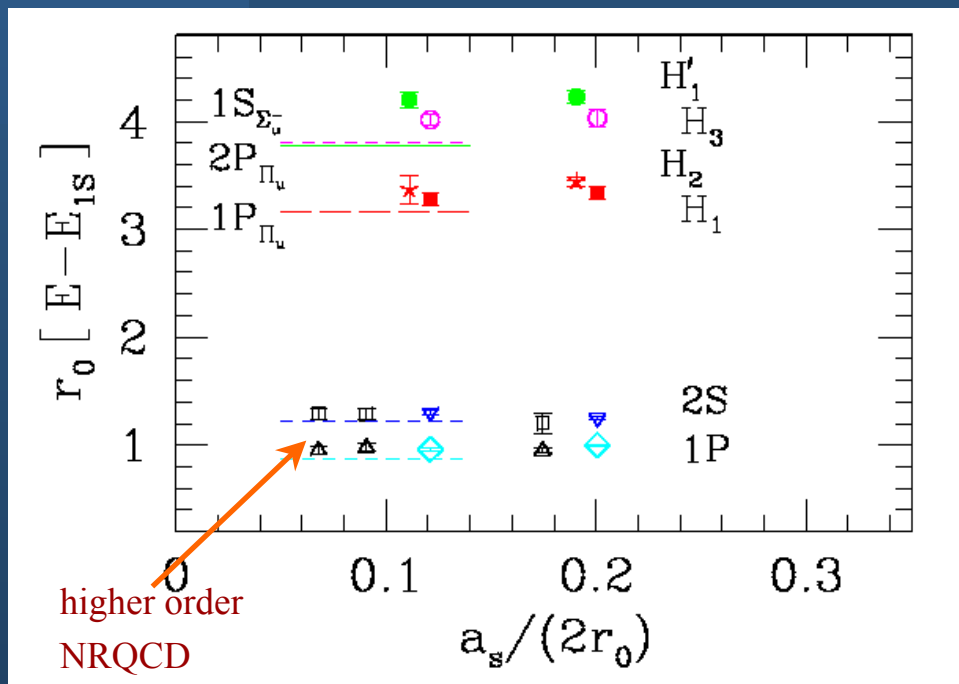
Charmonium LBO

- same calculation, but for charmonium



Testing LBO

- test LBO by comparison of spectrum with NRQCD simulations
 - include retardation effects, but no quark spin, no \vec{p}^4 , no light quarks
 - allow possible mixings between adiabatic potentials
- dramatic evidence of validity of LBO
 - level splittings agree to 10% for 2 conventional mesons, 4 hybrids



$$H_1, H'_1 = 1^{--}, 0^{++}, 1^{+-}, 2^{+-}$$

$$H_2 = 1^{++}, 0^{+-}, 1^{+-}, 2^{+-}$$

$$H_3 = 0^{++}, 1^{+-}$$

J^{PC}		Degeneracies	Operator
0^{-+}	S wave	1^{--}	$\chi^\dagger [\hat{\Delta}^{(2)}]^P \psi$
1^{+-}	P wave	$0^{++}, 1^{++}, 2^{++}$	$\chi^\dagger \hat{\Delta} \psi$
1^{--}	H_1 hybrid	$0^{-+}, 1^{-+}, 2^{-+}$	$\chi^\dagger \hat{B} [\hat{\Delta}^{(2)}]^P \psi$
1^{++}	H_2 hybrid	$0^{+-}, 1^{+-}, 2^{+-}$	$\chi^\dagger \hat{B} \times \hat{\Delta} \psi$
0^{++}	H_3 hybrid	1^{+-}	$\chi^\dagger \hat{B} \cdot \hat{\Delta} \psi$

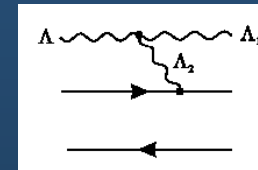
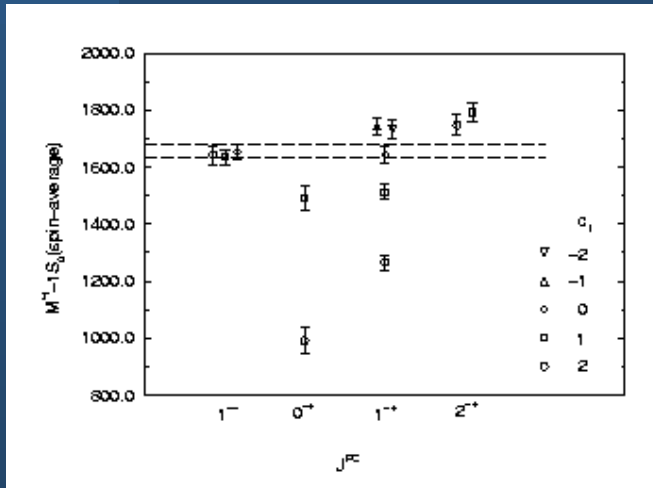
lowest hybrid 1.49(2)(5) GeV above 1S

Compelling physical picture

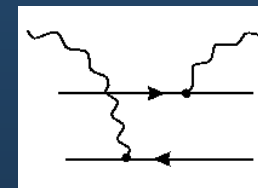
- Born-Oppenheimer provides simple physical picture for heavy-quark conventional and hybrid meson states
 - partial explanation of quark model success
 - insight into light quarks?
 - allows incorporation of gluon dynamics (beyond quark model)
- does this BO picture survive inclusion of
 - quark spin?
 - light-quark effects?

Quark spin effects

- quark spin: recent studies *suggest* BO picture survives
 - Drummond *et al.* Phys.Lett.B478, 151 (2000)
 - looked at 4 hybrids degenerate in LBO using NRQCD
 - found significant shifts from $c_1 \sigma \cdot B / M$ but used bag model to interpret results as not arising from surface mixing effects
 - suggestive, but not definitive



dominant (but does not spoil BO)



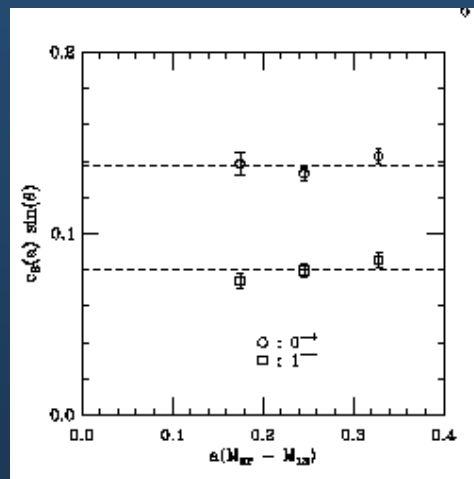
mixes adiabatic surfaces, but very small (10^{-4})

Quark spin effects (continued)

- Burch and Toussaint, hep-lat/0305008
 - NRQCD simulations, measured mixing via $c_1 \sigma \cdot B / M$
 - mixing in bottomonium seems not to spoil BO picture
 - larger effect in charmonium

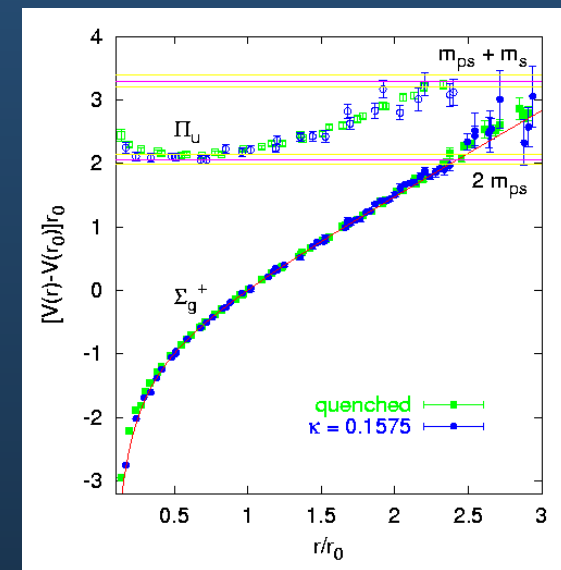
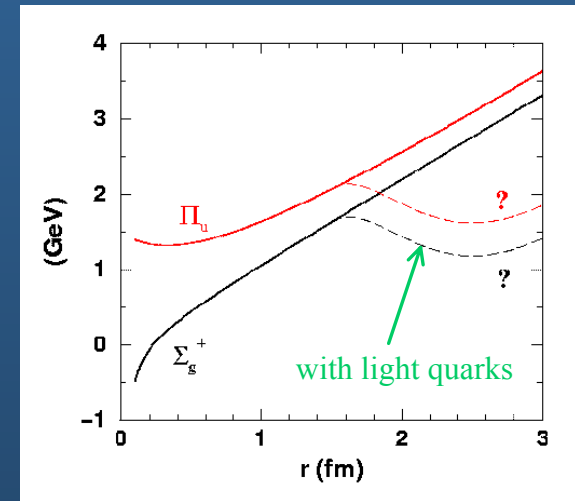
$$\langle 1H | Y \rangle \approx 0.076 - 0.11 \quad \langle 1H | J / \Psi \rangle \approx 0.18 - 0.25$$

$$\langle 1H | \eta_b \rangle \approx 0.13 - 0.19 \quad \langle 1H | \eta_c \rangle \approx 0.29 - 0.4$$



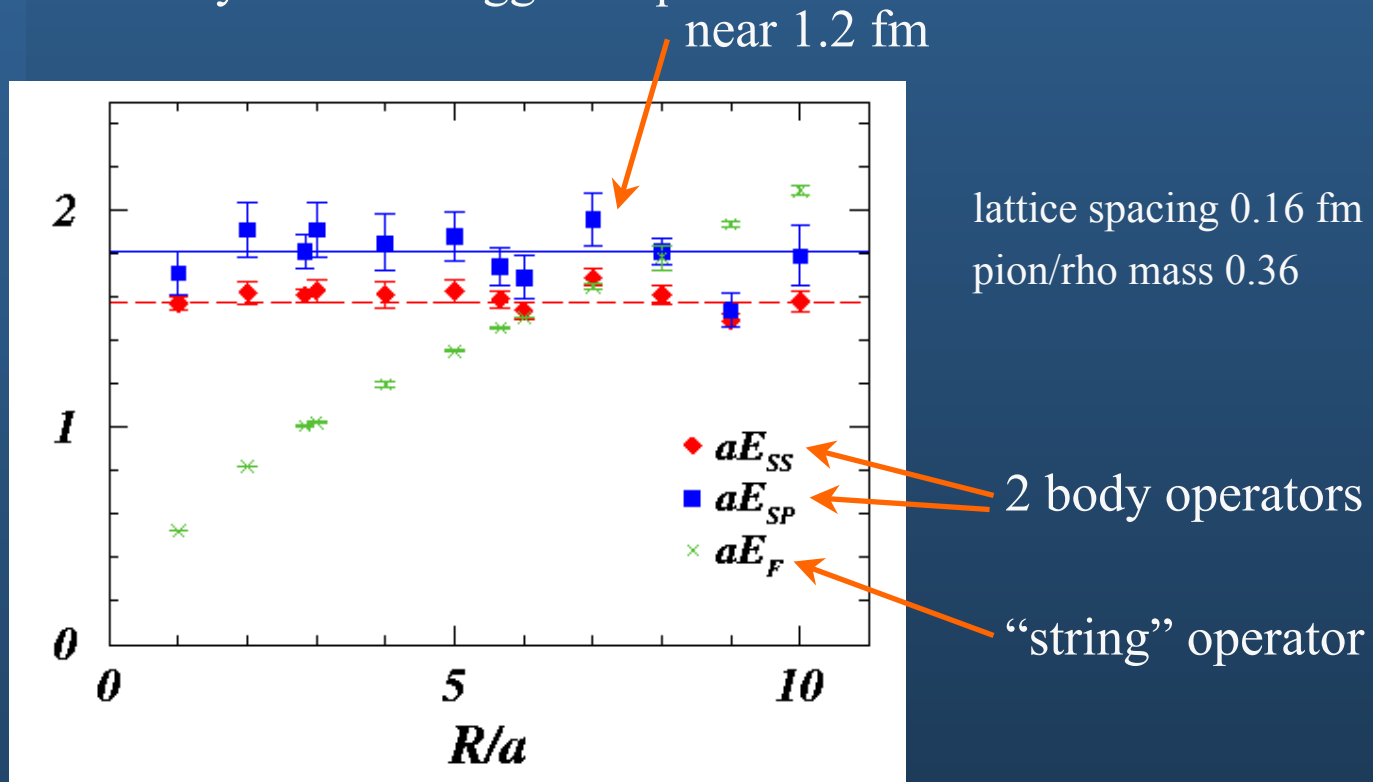
Light quark spoiler?

- spoil B.O.? → unknown
- light quarks change $V_{Q\bar{Q}}(r)$
 - small corrections at small r
 - fixes low-lying spectrum
 - large changes for $r > 1$ fm
 - fission into $(Qq)(\bar{Q}q)$
- states with diameters over 1 fm
 - most likely *cannot exist* as observable resonances
- dense spectrum of states from pure glue potentials will not be realized
 - survival of a few states conceivable given results from Bali *et al.*
- discrepancy with experiment above $B\bar{B}$
 - most likely due to light quark effects



String breaking

- string “breaking” using 2 body operators
- two flavors of dynamical staggered quarks

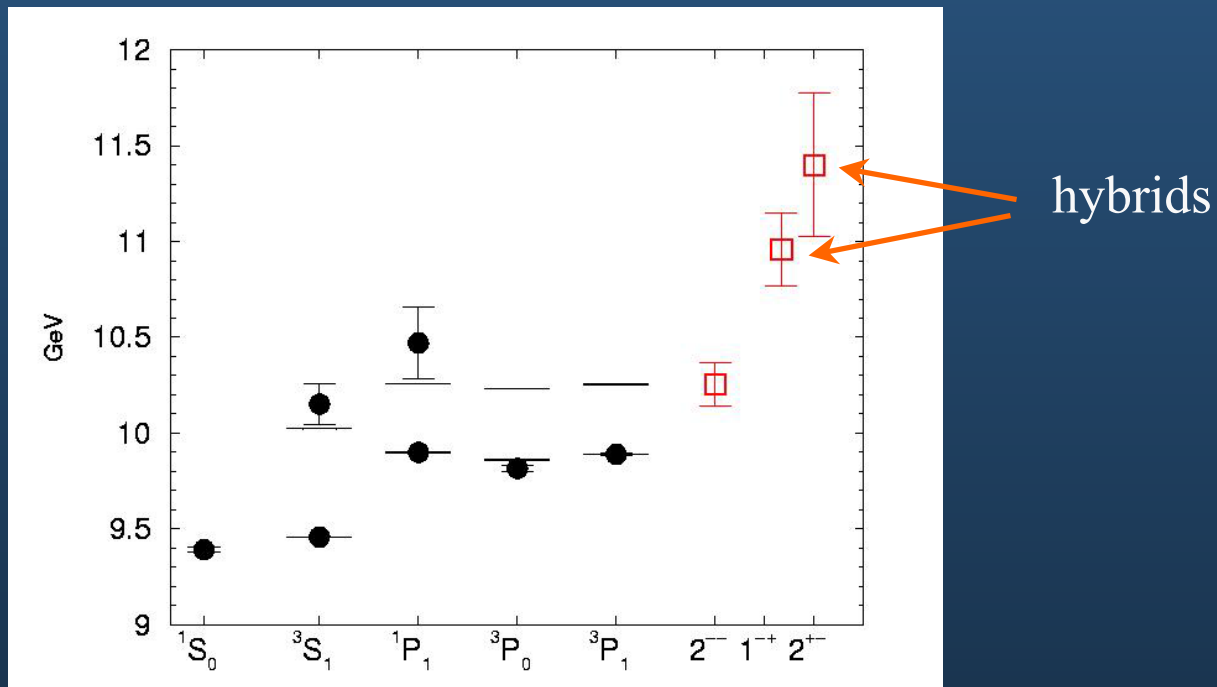


Bernard et al., PRD64, 074509 (2001)

Bottomonium hybrids

- recent calculation of bottomonium hybrids confirms earlier results
 - quenched, several lattice spacings so $a \rightarrow 0$ limit taken
 - improved anisotropic gluon and fermion (clover) actions
 - good agreement with Born-Oppenheimer (but errors large)

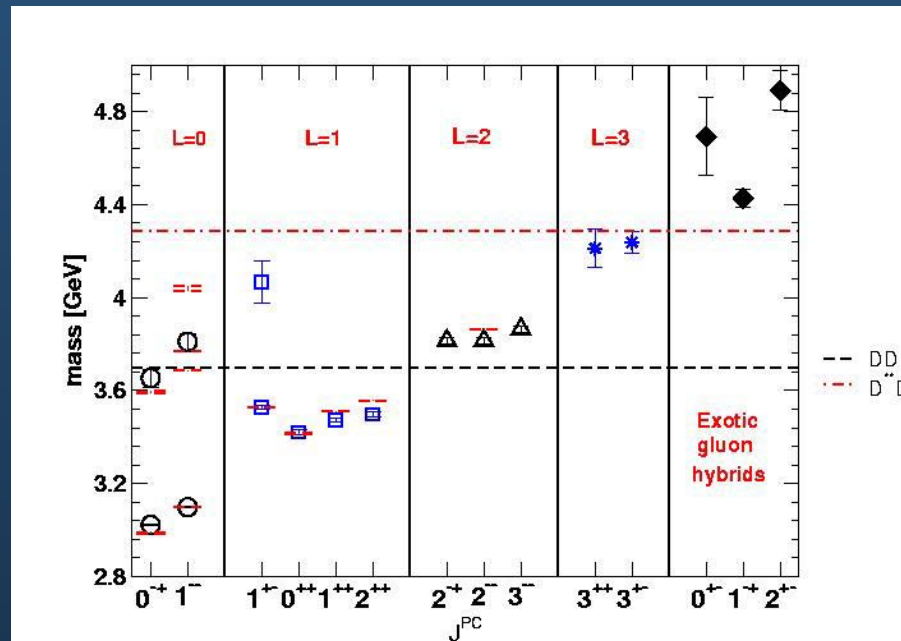
Liao, Manke, PRD65, 074508 (2002)



Charmonium hybrids

- recent determination of some charmonium hybrids
 - quenched, several lattice spacings for continuum limit
 - improved, anisotropic gluon and fermion (clover) actions
 - results suggest significant (but not large) corrections from LBO

Liao, Manke, hep-lat/0210030



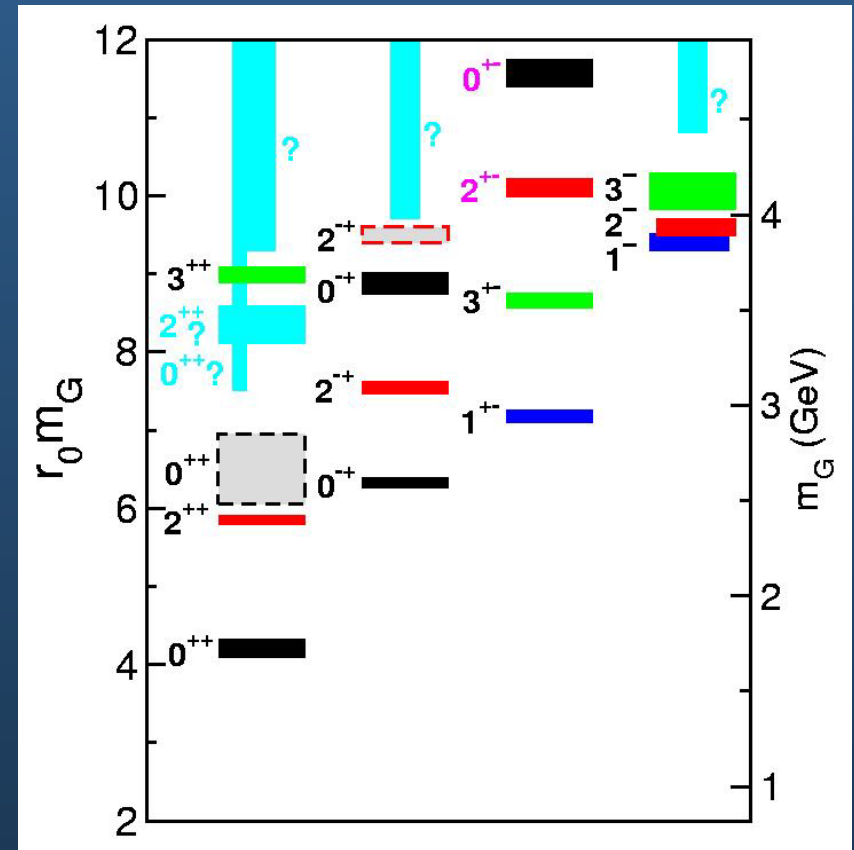
Tidbits

- glueballs
- light quark hybrids
- static three quark potential

Yang-Mills SU(3) Glueball Spectrum

- gluons can bind to form glueballs
- first glimpse of rich spectrum
- probe of confinement
- “experimental” results in simpler world (no quarks) to help build models of gluons
- add quarks for QCD glueballs
- future work: glueball structure
 - bag model, flux loops?

C. Morningstar and M. Peardon,
Phys. Rev. D 60, 034509 (1999)



$r_0^{-1} = 410(20)$ MeV, states labeled by J^{PC}

Glueballs (qualitative features)

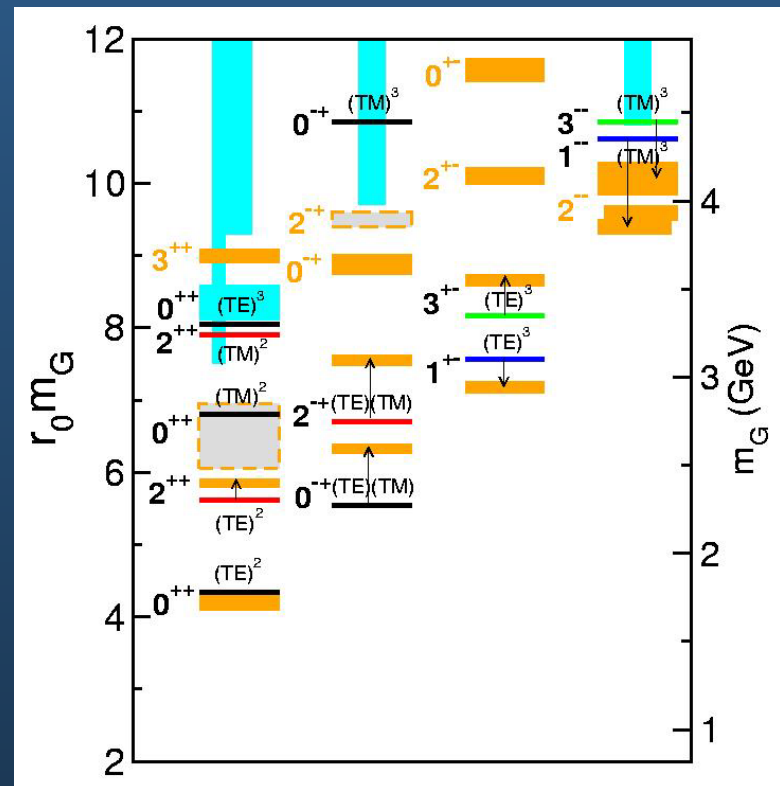
- spectrum can be qualitatively understood in terms of interpolating operators of minimal dimension (Jaffe,Johnson,Ryzak, Ann. Phys. **168**, 344 (1986))
 - dimension 4: $\text{Tr } F_{\mu\nu} F_{\alpha\beta} \Rightarrow 0^{++}, 0^{-+}, 2^{++}, 2^{-+}$
 - dimension 5: $\text{Tr } F_{\mu\nu} D_{\rho} F_{\alpha\beta} \Rightarrow 1^{++}, 3^{++}$
 - dimension 6: $\text{Tr } F_{\mu\nu} F_{\delta\sigma} F_{\alpha\beta} \Rightarrow 0^{\pm\pm}, 1^{\pm\pm}, 2^{\pm\pm}, 3^{\pm-}$
 $\text{Tr } F_{\mu\nu} \{D_{\rho}, D_{\sigma}\} F_{\alpha\beta} \Rightarrow 1^{-+}, 3^{-+}, 4^{\pm\pm}$
- of lightest 6 states, 4 have the J^{PC} of the dimension 4 operators
- absence of low-lying $0^{\pm-}, 1^{-+}$ glueballs explained

Glueballs (bag model)

- qualitative agreement with bag model
 - constituent gluons are TE or TM modes in spherical cavity
 - Hartree modes with residual perturbative interactions
 - center-of-mass correction

Carlson, Hansson, Peterson, PRD27, 1556 (1983);
J. Kuti (private communication)

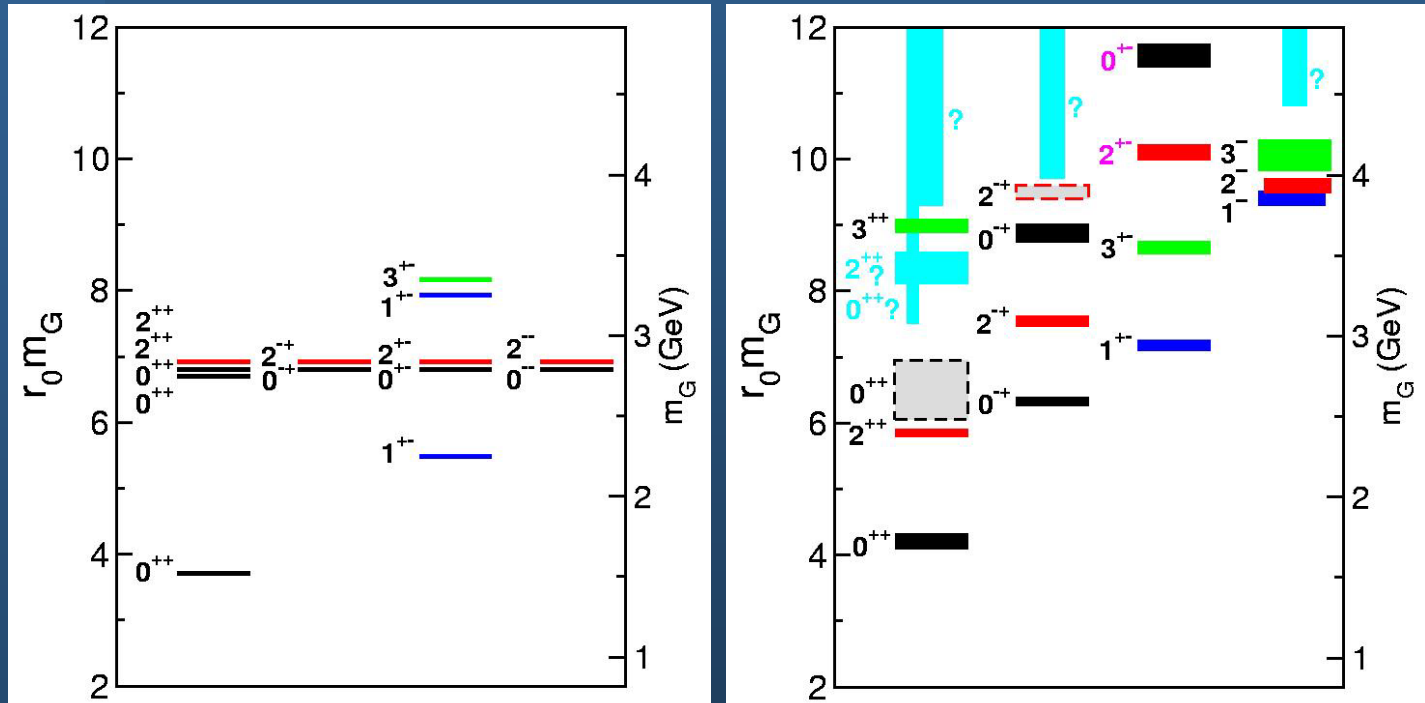
	1983	1993
	light baryon spectroscopy	static-quark potential
α_s	1.0	0.5
$B^{1/4}$	230 MeV	280 MeV



Glueballs (flux tube model)

- disagreement with one particular string model

Isgur, Paton, PRD31, 2910 (1985)

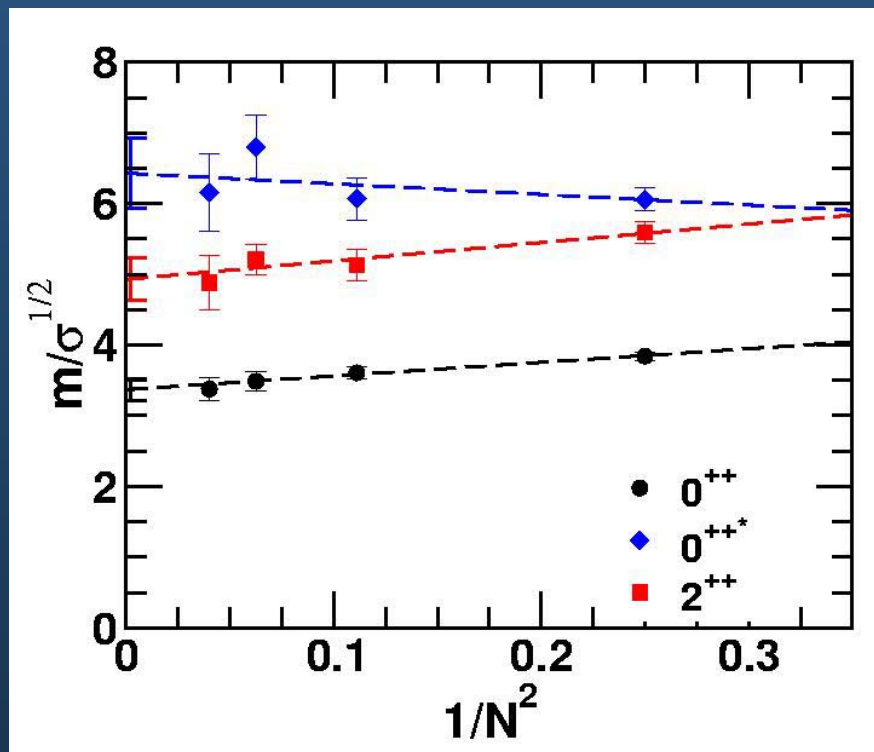


- future comparisons:
 - with more sophisticated string models (soliton knots)
 - AdS theories, duality

SU(N) Glueballs

- recent study of $0^{++}, 2^{++}, 0^{++*}$ glueballs in SU(N), $N=2,3,4,5$
- masses depend linearly on $1/N^2$
- large $N \rightarrow \infty$ limits differ little from $N=3$

Lucini, Teper, JHEP 06, 050 (2001).



Light-quark hybrids

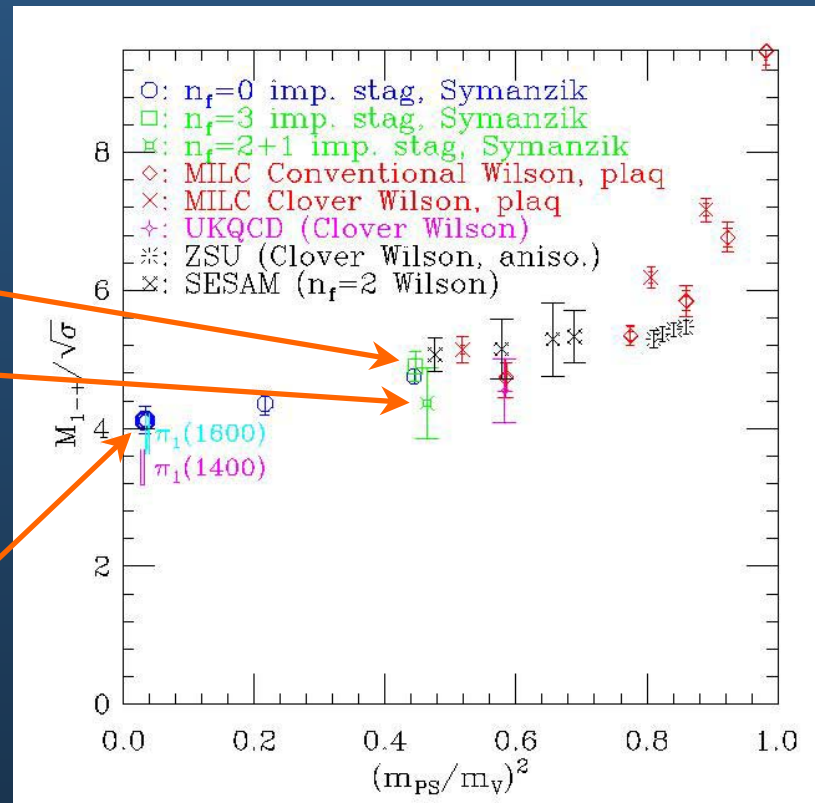
- recent new determination of exotic 1^{+-} hybrid meson
 - improved staggered fermions (lighter quark masses)
 - quenched and unquenched, Wilson gluon action
 - $a \approx 0.09$ fm
 - lightest mass still above experiment

MILC, hep-lat/0301024

$N_f = 3, \quad m_u = m_d = m_s$
(around strange quark mass)

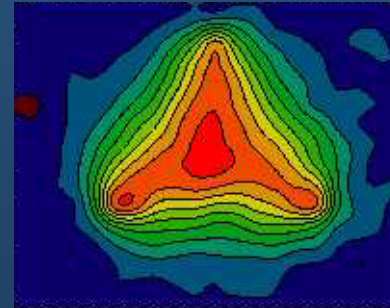
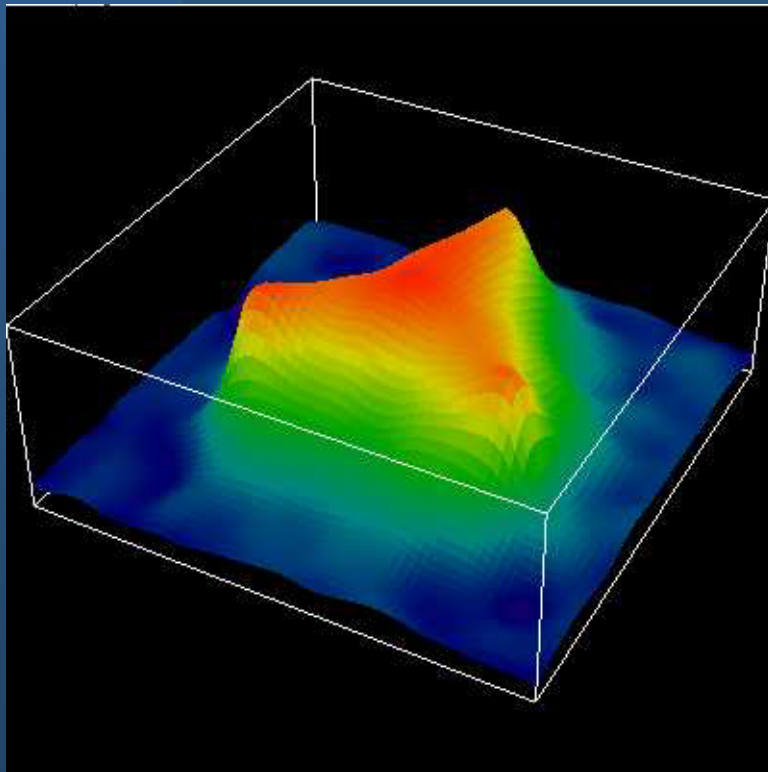
$m_u = m_d = 0.4m_s$

quenched continuum limit



Static three-quark system

- recent determination of the abelian action distribution of gluons and light quarks in the presence of three static quarks
 - supports a Y-type flux configuration

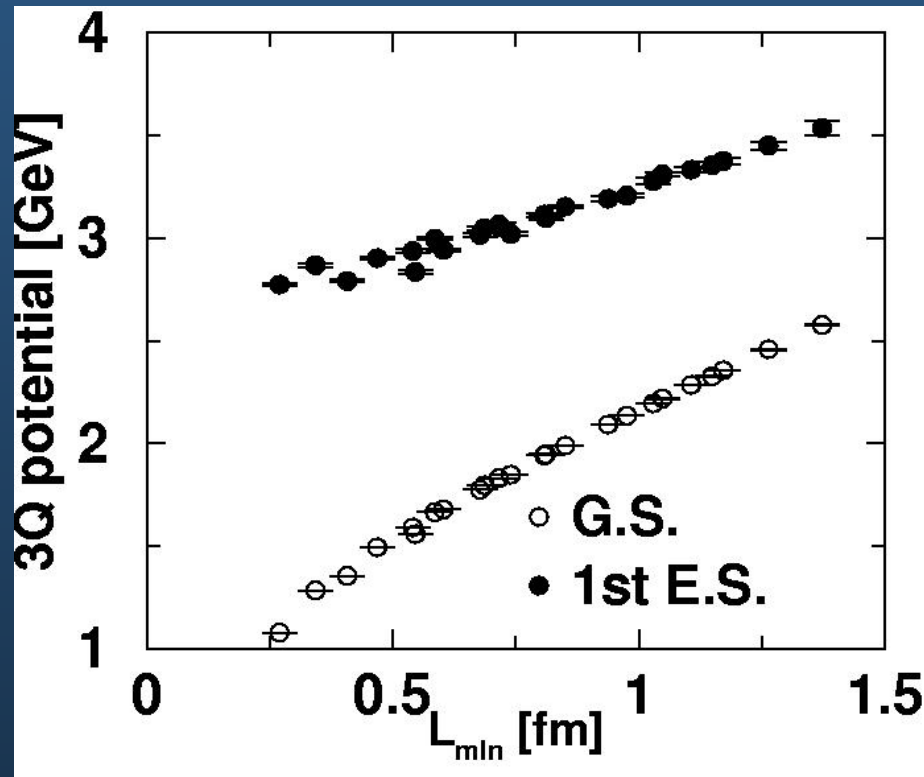


Ichie, Bornyakov, Struer, Schierholz,
hep-lat/0212024

Excitation of the static $3q$ system

- first excitation of the static $3q$ system recently determined
 - excitation energy about 1 GeV
 - finite spacing, finite volume errors still to be studied

Takahashi, Suganuma, hep-lat/0210024



Conclusion

- hadronic states bound by an *excited* gluon field
 - interesting new form of matter
 - shed new light on confinement in QCD
- heavy-quark hybrid mesons
 - validity of a Born-Oppenheimer treatment
 - relationship to excitations of the static quark potential
 - compelling physical picture
 - quark spin effects do not spoil BO
 - light quark loops → survival issue